

Higgs and beyond at the LHC with a little help from QCD

Gavin Salam

CERN, Princeton University & LPTHE/CNRS (Paris)

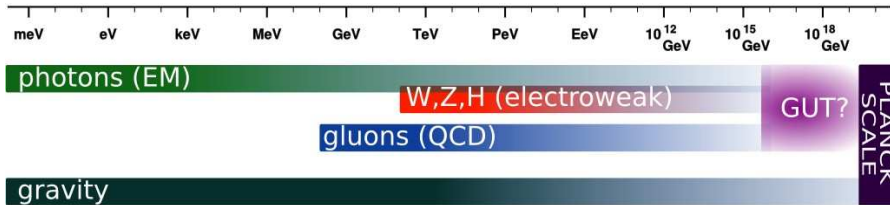
MIT Physics Colloquium
6 September 2012

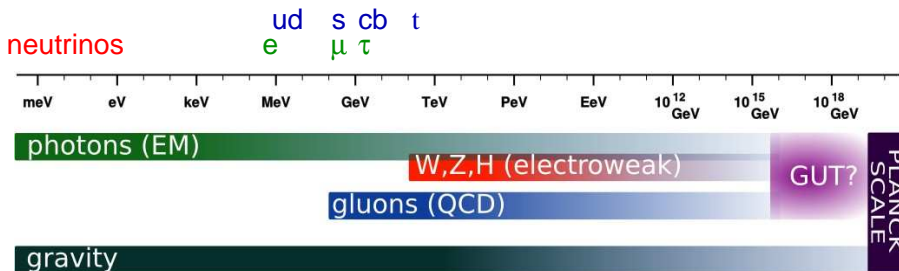
The LHC has been colliding protons since late 2009

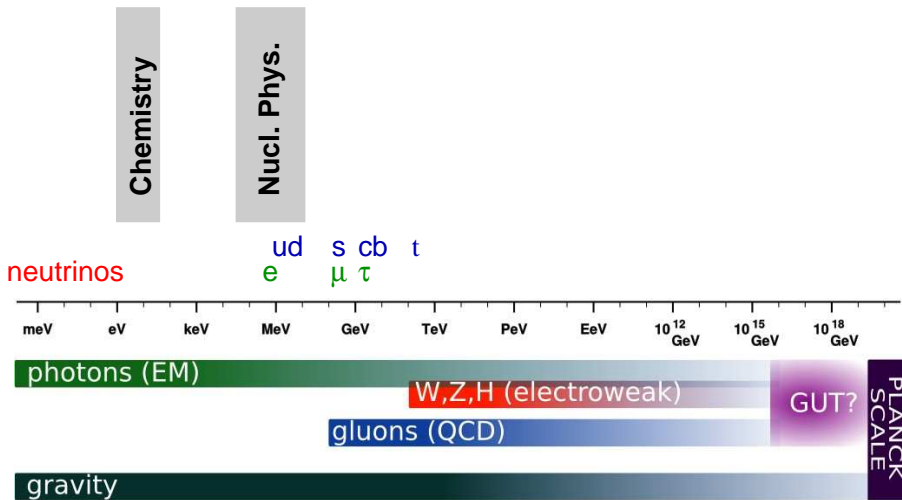
The world's largest fundamental physics endeavour

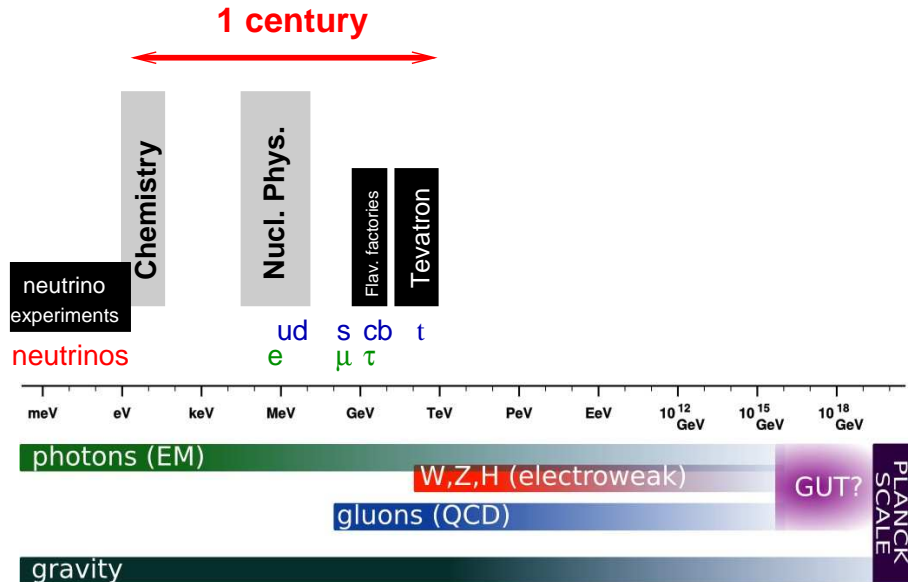
Involving $\mathcal{O}(10\,000)$ scientists and engineers
From about 60 countries across the world
At a cost of several billion US dollars



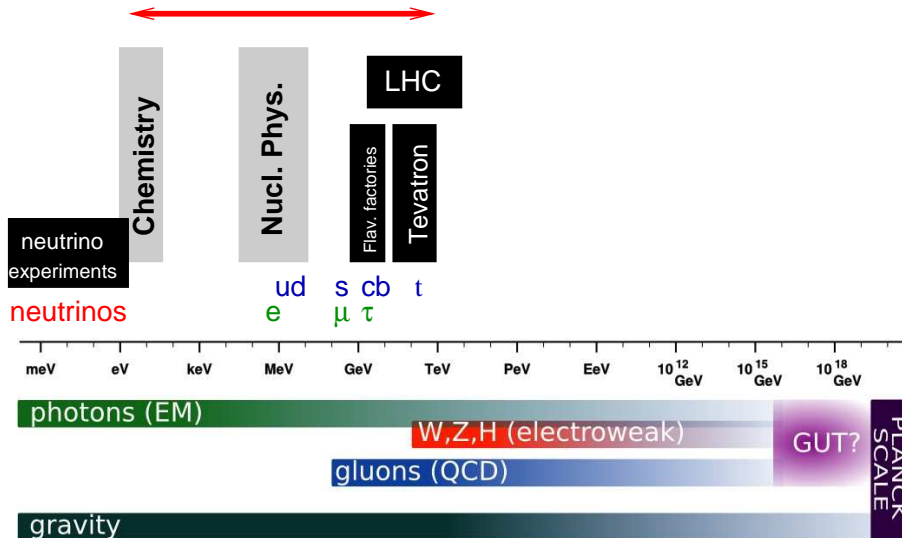


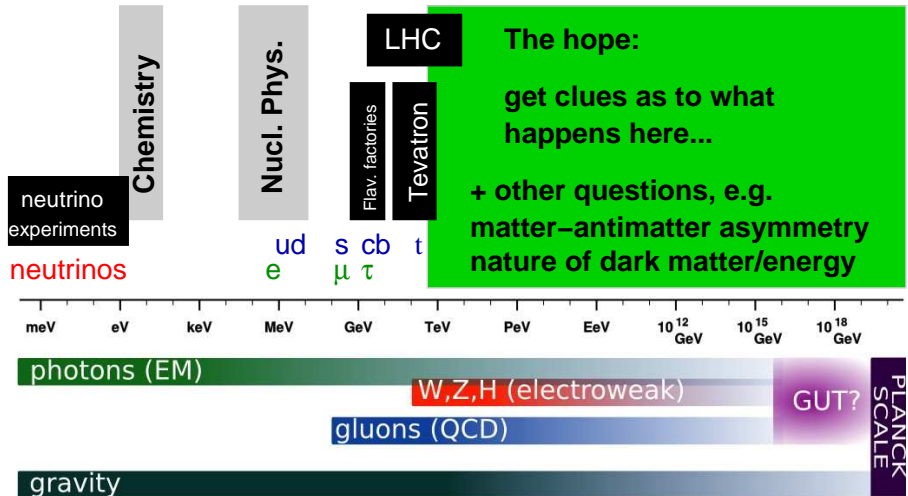


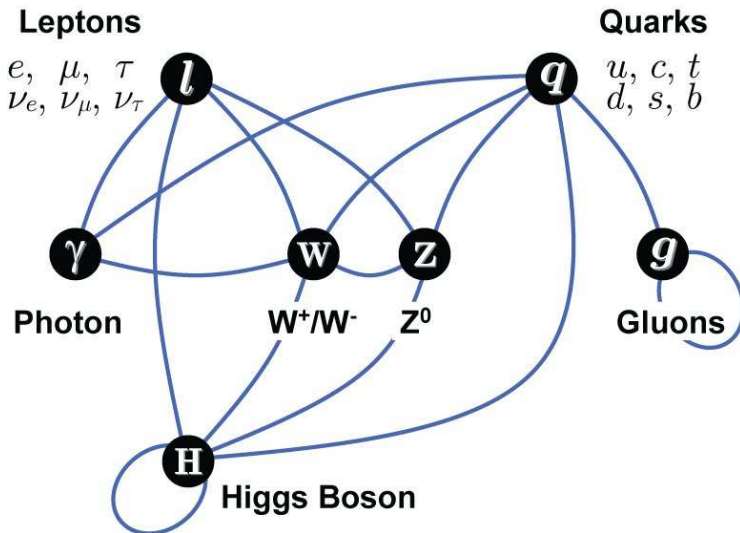


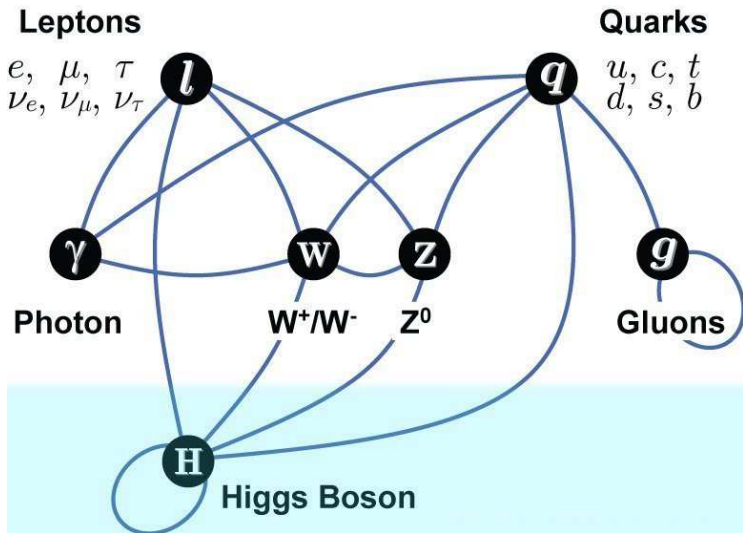


1 century









Among the terms in the
Standard Model Lagrangian:

$$g_W^2 \phi^2 Z_\mu Z^\mu$$

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↖
Z-boson fields

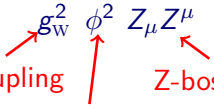
Among the terms in the
Standard Model Lagrangian:

$$g_w^2 \phi^2 Z_\mu Z^\mu$$

gauge coupling

Z-boson fields

Among the terms in the
Standard Model Lagrangian:



The diagram shows three mathematical terms in blue: g_w^2 , ϕ^2 , and $Z_\mu Z^\mu$. Red arrows point from these terms to their physical interpretations written in red text below them. An arrow points from g_w^2 to "gauge coupling". An arrow points from ϕ^2 to "scalar field". An arrow points from $Z_\mu Z^\mu$ to "Z-boson fields".

g_w^2 ϕ^2 $Z_\mu Z^\mu$

gauge coupling scalar field Z-boson fields

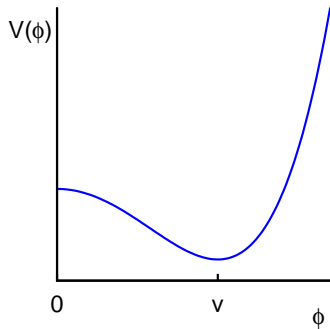
Among the terms in the Standard Model Lagrangian:

g_w^2 ϕ^2 $Z_\mu Z^\mu$
 gauge coupling scalar field Z-boson fields

Potential for scalar field is

$$V(\phi) = -\mu^2 \phi^2 + \lambda \phi^4$$

Universe lives at minimum of potential, $\phi \simeq v$. Rewrite ϕ in terms of perturbations H around minimum



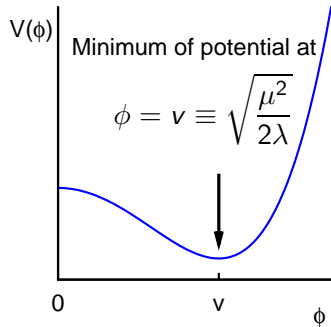
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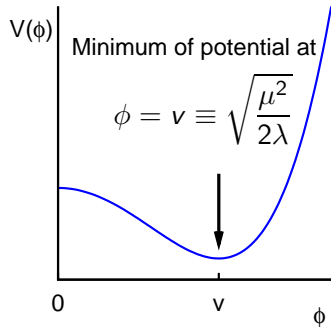
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$$\phi \equiv v + H \rightarrow \phi^2 = v^2 + 2vH + H^2$$

H is the **Higgs-boson** field

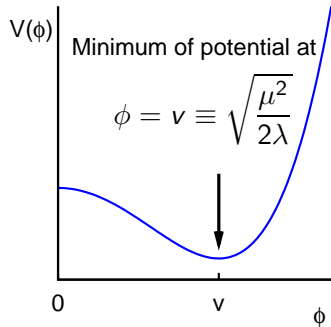
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H is the **Higgs-boson** field

$$g_w^2 \phi^2 Z_\mu Z^\mu \rightarrow \underbrace{g_w^2 v^2}_{\text{Z mass}^2} Z_\mu Z^\mu + \underbrace{2g_w^2 v}_{\text{HZZ coupling}} H Z_\mu Z^\mu$$

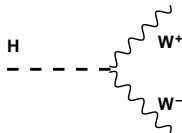
Mechanism generates particle masses
And a "Higgs" boson

A similar mechanism holds for fermions, with a Yukawa coupling y_f ,

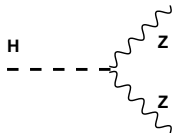
$$y_f \phi f \bar{f} \rightarrow \underbrace{y_f v f \bar{f}}_{\text{fermion mass}} + \underbrace{y_f H f \bar{f}}_{\text{interaction}}$$

Higgs mechanism gives mass to all fundamental particles (except maybe neutrinos).

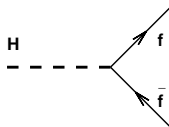
It predicts characteristic relationship between their masses and their interactions with the Higgs boson.



$$2i \frac{M_W^2}{v} g_{\mu\nu}$$



$$2i \frac{M_Z^2}{v} g_{\mu\nu}$$



$$i \frac{m_f}{v}$$

$v \simeq 246 \text{ GeV}$ is known as vacuum expectation value of Higgs field

$$V(\phi) = -\mu^2\phi^2 + \lambda\phi^4 \quad \longrightarrow \quad M_H = \sqrt{2\lambda}v$$

quartic coupling λ unknown, so **no prediction about Higgs mass**

Still, strong arguments say

- ▶ it cannot be below 70 GeV, because λ too small — renormalisation group evolution drives it negative and our universe is unstable
- ▶ if $M_H \gtrsim 800$ GeV, λ is large and we see new non-perturbative physics at ~ 1 TeV

So ideally build a collider that can discover Higgs-boson up to 800 GeV and perform WW scattering up to $\simeq 1$ TeV

F.A.Q.s

ABOUT THE

HADRON COLLIDER



Q: How does the Hadron Collider work?

A: You didn't even understand eleventh-grade math, so why are you asking?



Q: What would happen if I went inside it?

A: Just. Don't.



Q: How many miles of pipes and whatnot are in it?

A: A bajillion.

Q: How much did it cost?

A: Forty squillion.



Q: What does this thing do?

A: Don't touch that.



Q: What would happen if you, like, put a cat inside it?

A: I don't know.

Q: If I concentrate ultra-hard, will I ever be able to understand it?

A: No.

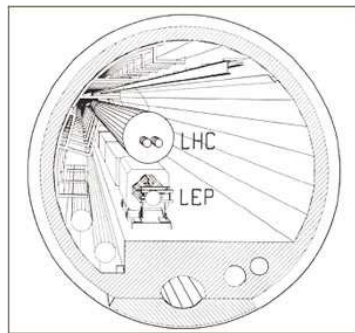


R. Chan

From the CERN Courier in **1984**:

The installation of a hadron collider in the [27km] LEP tunnel, using superconducting magnets, has always been foreseen by ECFA and CERN as the natural long term extension of the CERN facilities beyond LEP. [...]

Although the installation of such a hadron collider in the LEP tunnel might appear still a long way off [...], it [is] an opportune moment for ECFA, in collaboration with CERN, to organize a 'Workshop on the Feasibility of a Hadron Collider in the LEP Tunnel' [...]

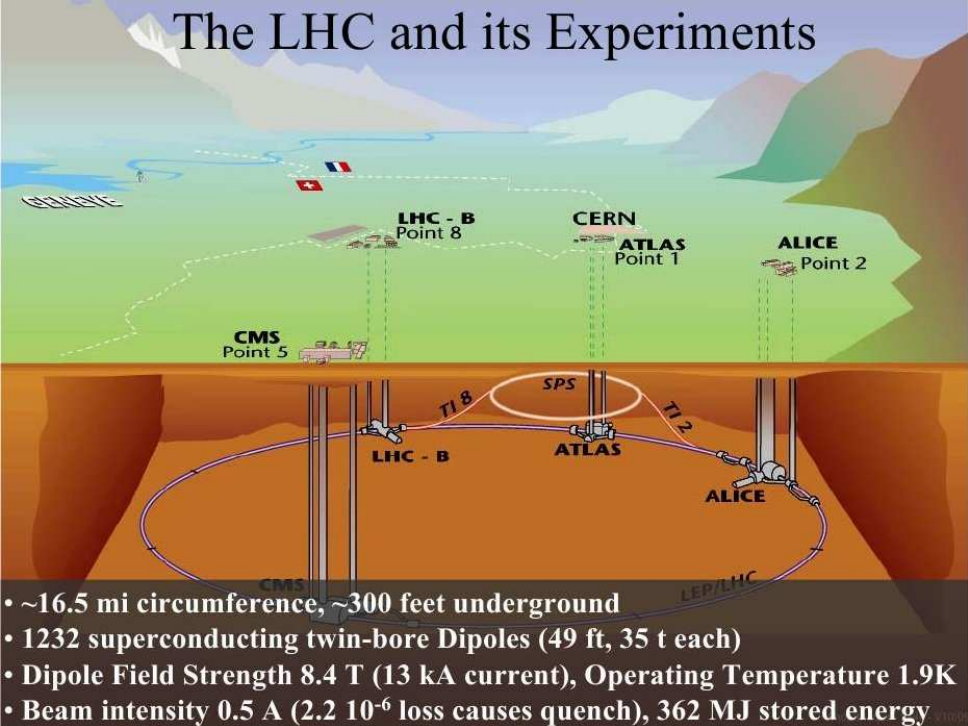


$$E \propto BR$$

ring radius $R \sim 4 \times \text{Tevatron}$
 superconduction magnets: $B = 8 \text{ T}$
 ($2 \times \text{Tevatron}$)

Tevatron $\sim 2 \text{ TeV} \rightarrow \text{LHC} \sim 14 \text{ TeV}$

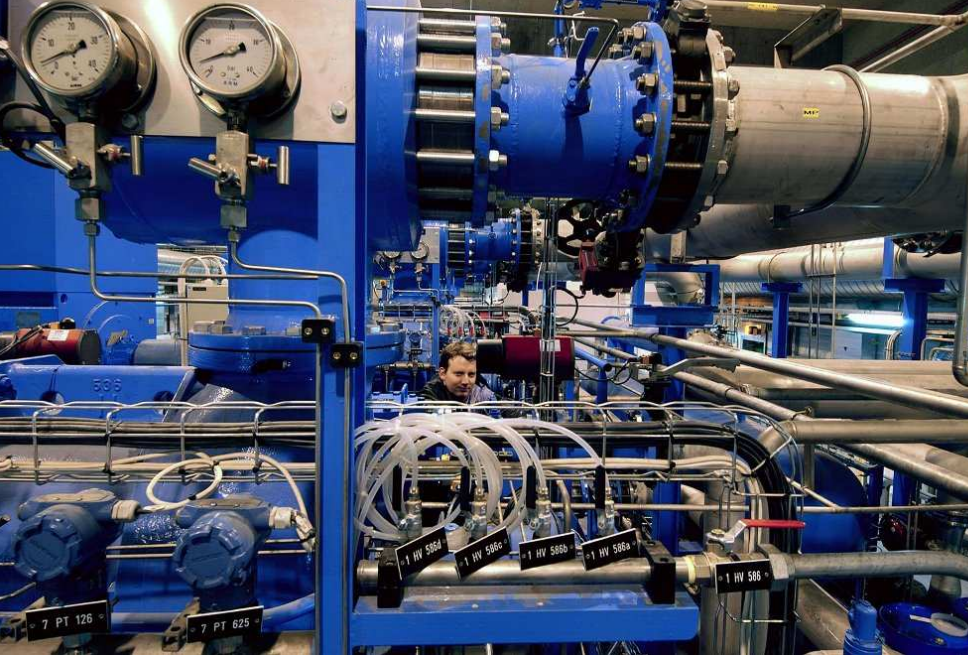
The LHC and its Experiments



- ~16.5 mi circumference, ~300 feet underground
- 1232 superconducting twin-bore Dipoles (49 ft, 35 t each)
- Dipole Field Strength 8.4 T (13 kA current), Operating Temperature 1.9K
- Beam intensity 0.5 A ($2.2 \cdot 10^{-6}$ loss causes quench), 362 MJ stored energy



Interconnection between two “dipoles” (bending magnets) in the LHC tunnel.



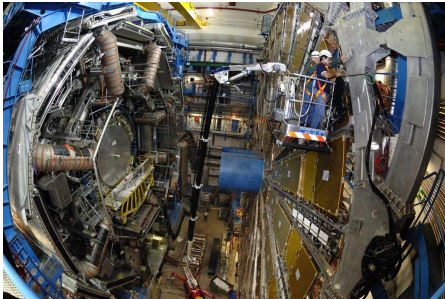
Cryogenics plant: 96 000 kg of Helium circulate through the machine at 1.9K

The detectors:

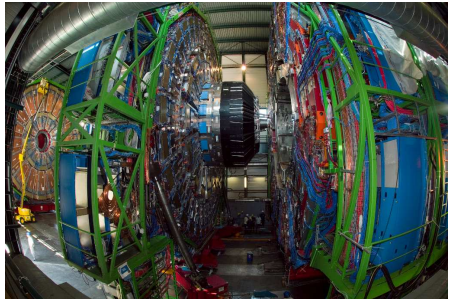
To accumulate 5×10^{16} collisions over a few years, they have to be able to handle a pp collision rate of 10^9 Hz
[25 collisions every 25 ns]

Typically, about 100 000 000 channels to read out.
[must be examined 40 000 000 times/s,
interesting events written to long-term storage ~ 400 – 1000 times/s]

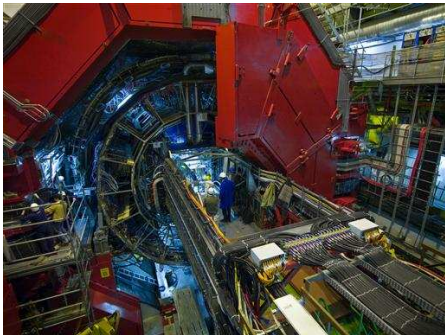
ATLAS: general purpose



CMS: general purpose



ALICE: heavy-ion physics



LHCb: B-physics



+ TOTEM, LHCf

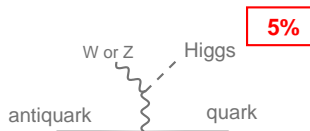
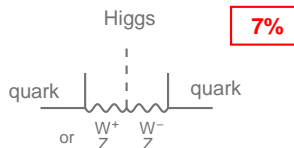
HIGGS PRODUCTION CHANNELS

(always indirect, because H couples
only weakly to light quarks)

cross section $\simeq 20$ pb

~ 1 Higgs every 5×10^9 pp collisions

[for $m_H = 125.5$ GeV]

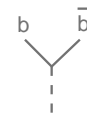


DECAY CHANNELS

[for $m_H = 125.5 \text{ GeV}$]

WW and ZZ suppressed relative to simple coupling proportionality, because they cannot be produced on-shell

Best channels for detection are $\gamma\gamma$ and $ZZ^*(\rightarrow 4e, \mu)$, because of excellent experimental mass resolution and manageable backgrounds



57%



22%

1% for $2 e, \mu + 2\nu$

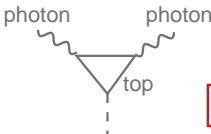


2.8%

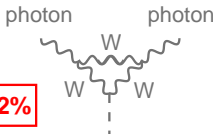
0.01% for $4 e, \mu$

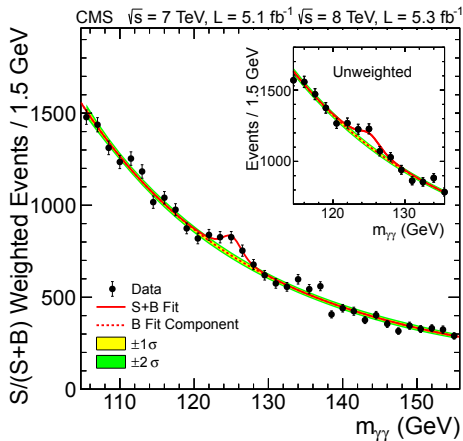
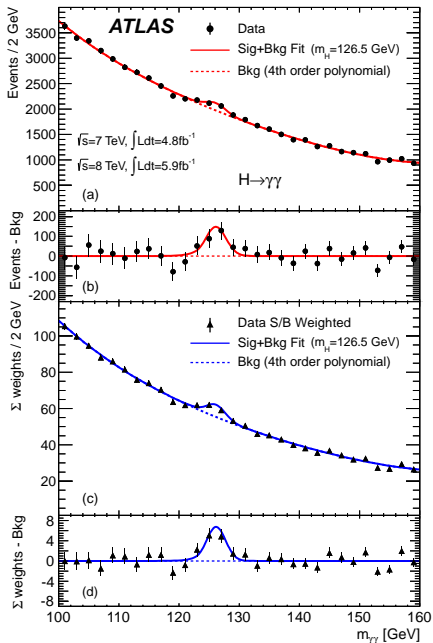


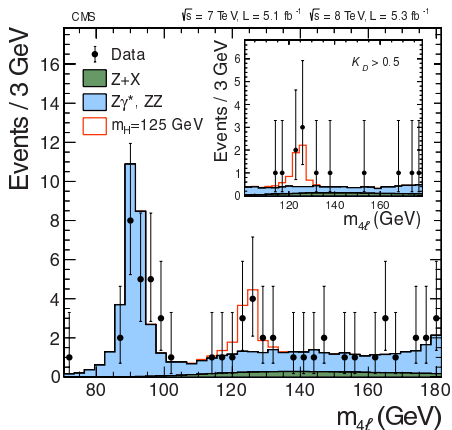
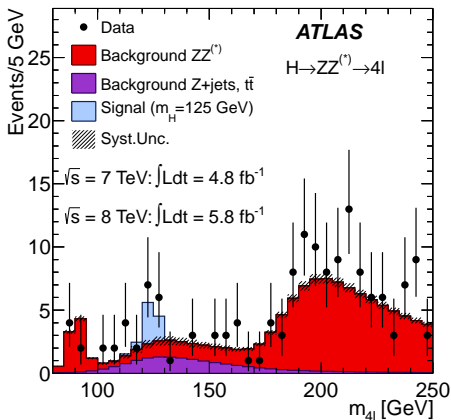
6%



0.2%

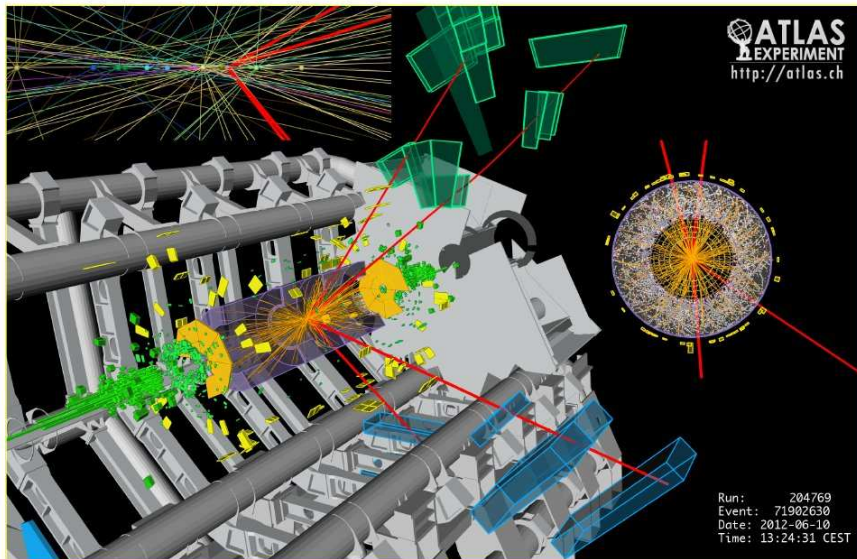




ZZ^* (4-lepton) invariant mass distribution

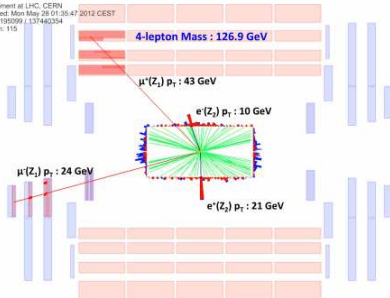
$$m_{4\mu} = 125.1 \text{ GeV}$$

p_T (muons) = 36.1, 47.5, 26.4, 71.7 [GeV]
 $m_{12} = 86.3 \text{ GeV}$, $m_{34} = 31.6 \text{ GeV}$
 15 reconstructed vertices

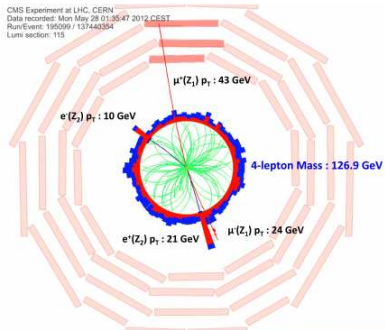


Event Display 2e2 μ event

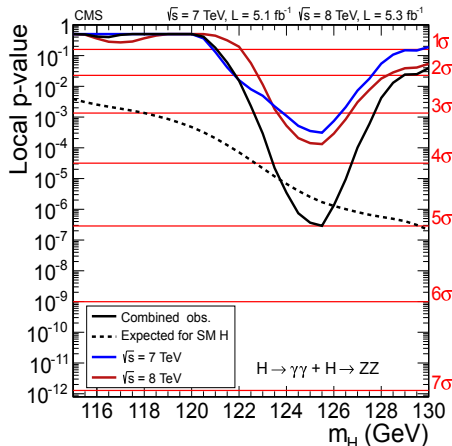
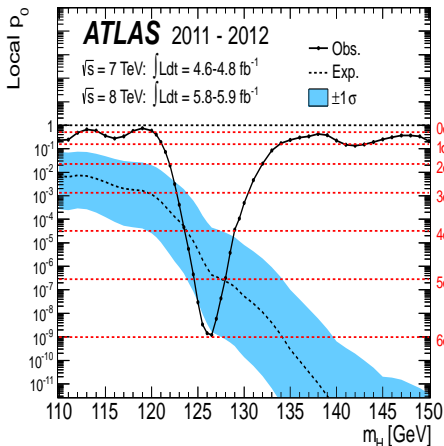
CMS Experiment at LHC, CERN
Data recorded: Mon May 28 01:35:47 2012 CEST
Run/Event: 195099 / 137440354
Lumi section: 115



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Data recorded: Mon May 28 01:35:47 2012 CEST
Run/Event: 195099 / 137440354
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p_0 value = probability background fluctuated to produce observed “bumps”.



**$\geq 5\sigma$: OBSERVATION OF A NEW PARTICLE
BY BOTH ATLAS AND CMS!**

Properties of this new “Higgs-like” particle

- ▶ decay to $\gamma\gamma$ indicates either spin 0 or 2
 - data by end of year should pin down spin and parity
- ▶ Mass:
 - ATLAS: $126.0 \pm 0.4 \pm 0.4$ GeV
 - CMS: $125.3 \pm 0.4 \pm 0.5$ GeV
- ▶ Couplings to other SM particles (key prediction of Higgs mechanism)

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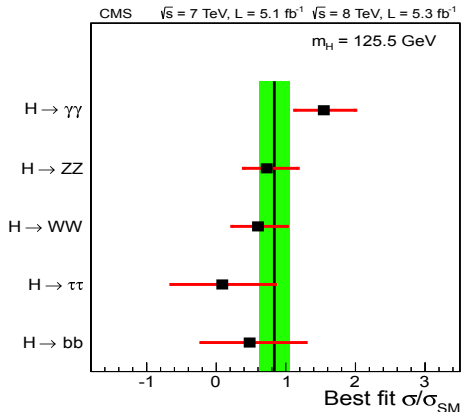
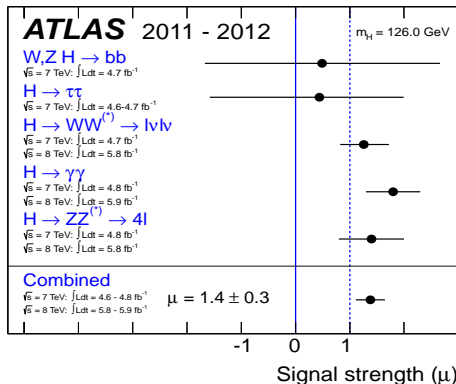
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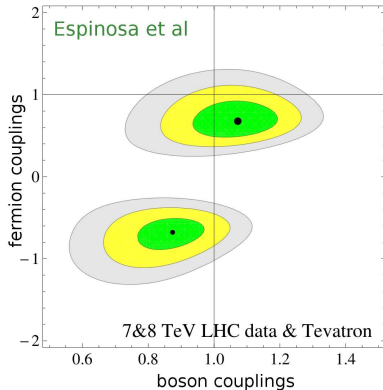
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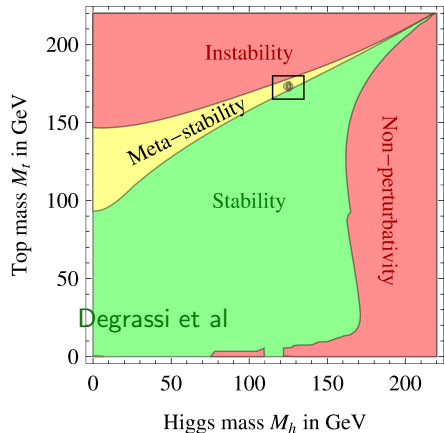
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Fits to underlying couplings



Stability of universe at M_{planck}



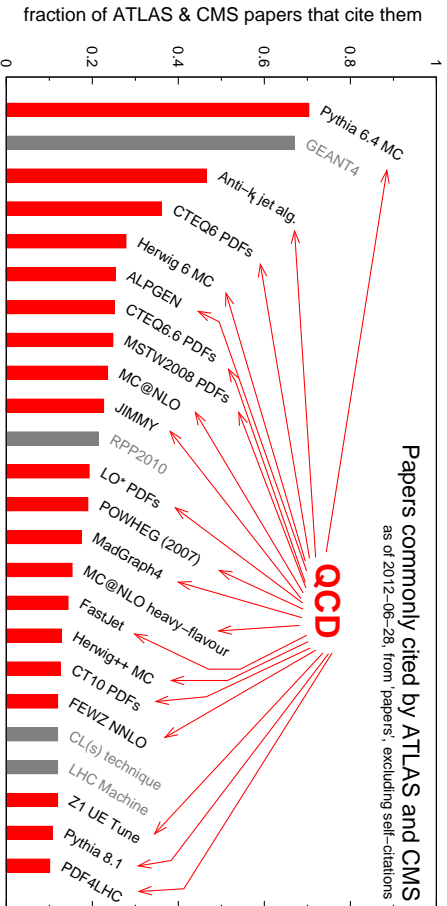
Since the discovery, a slew of papers have discussed couplings, proposed new physics explanations of deviations, etc.

The forthcoming installments of data will tell us more.

Behind the scenes...

One of the enablers for the LHC's success is our understanding of the strong interaction

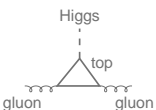
Quantum Chromodynamics — QCD



Rough combination of ATLAS and CMS $H \rightarrow \gamma\gamma$ signal strengths: they see 1.6 ± 0.3 times the standard model expectation — a little high, but still consistent

The standard-model cross section for gluon-fusion Higgs production is written as a perturbative expansion in powers of the strong coupling constant, with a leading-order (LO) term:

Higgs production cross section



strong coupling constant ~ 0.11

number of colliding gluon pairs per pp collision ~ 100

$$\sigma_{\text{LO}} = \frac{\alpha_s^2}{576\pi v^2} \mathcal{L}_{gg}$$

Higgs field vacuum expectation value

If ATLAS and CMS had used the previous page's formula they would have found

$$\frac{\sigma_{\text{observed}}}{\sigma_{\text{LO}}} = 5.6 \pm 1.1$$

This would have been a strong sign (4σ) of physics beyond the standard model!

Where's the catch? Higher orders of QCD perturbation theory:

$$\begin{aligned}\sigma_{gg \rightarrow H} &= \sigma_{\text{LO}} (1 + 11.4 \alpha_s + 63 \alpha_s^2 + \dots) \\ &= \sigma_{\text{LO}} (1 + 1.27 + 0.79 + \dots) \\ &\simeq \sigma_{\text{LO}} \times 3.4\end{aligned}$$

NLO: Dawson '91; Djouadi, Spira & Zerwas '91

NNLO: Harlander & Kilgore '02; Anastasiou & Melnikov '02; Ravindran, Smith & van Neerven '03

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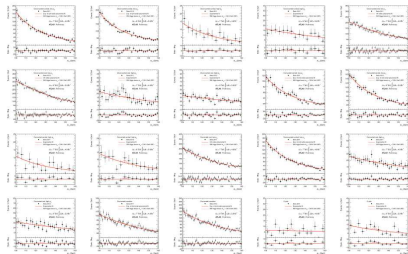
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Corrections to total $\sigma_{gg \rightarrow H}$ are tip of QCD iceberg

Mass spectrum

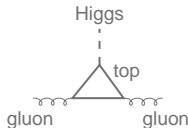
Mass spectra of the individual categories consisting the final result



Krisztian Peters

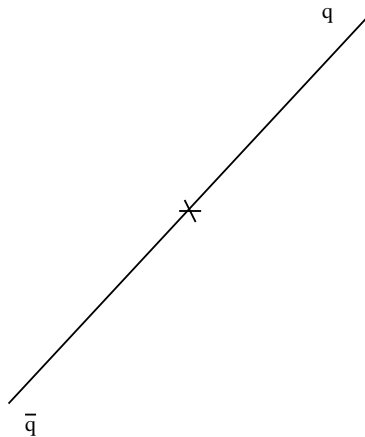
ATLAS $H \rightarrow \gamma\gamma$ 16

Experimental searches break analysis into sub-channels, which can differ in terms of signal process, backgrounds, resolutions, etc. Need QCD predictions in each sub-channel.



One of the richest aspects of this kind of event characterization involves **jets**, which help count the number of energetic quarks and gluons in an event.

Quarks & gluons? We only ever see “jets”



Start off with quark and anti-quark, $q\bar{q}$

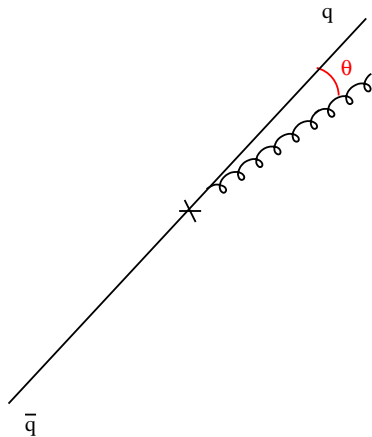
Quarks & gluons? We only ever see “jets”

In perturbative quantum chromodynamics (QCD), probability that a quark or gluon emits a gluon:

$$\sim \alpha_s \frac{dE}{E} \frac{d\theta}{\theta}$$

Diverges for small gluon energies E

Diverges for small angles θ



**A quark never survives unchanged
it always emits a gluon (usually low-energy, at small angles)**

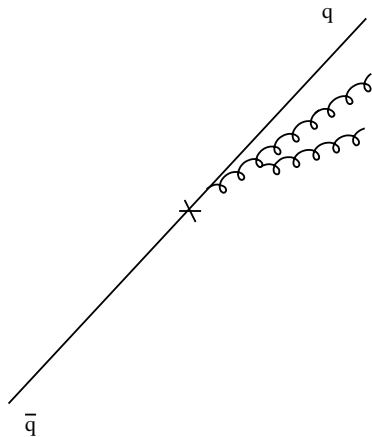
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Each gluon radiates a further gluon

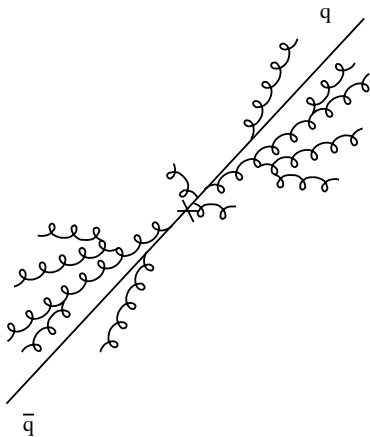
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And so forth

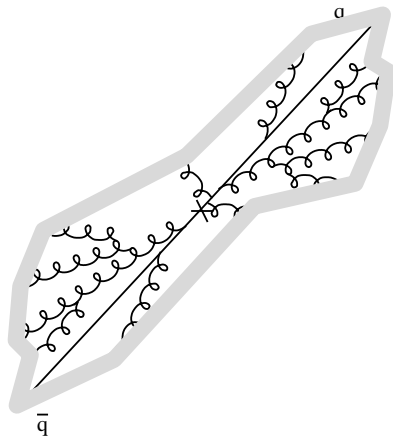
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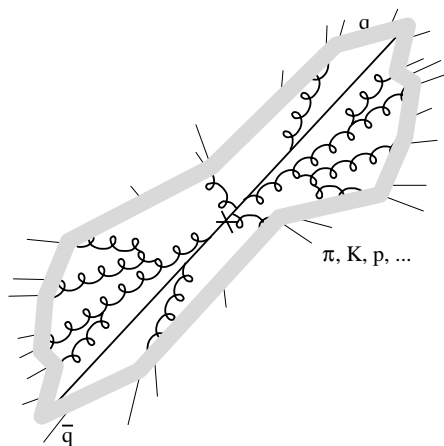
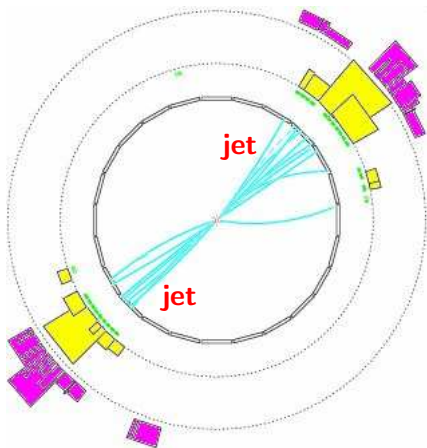
Diverges for small gluon energies E

Diverges for small angles θ



And then a non-perturbative transition occurs

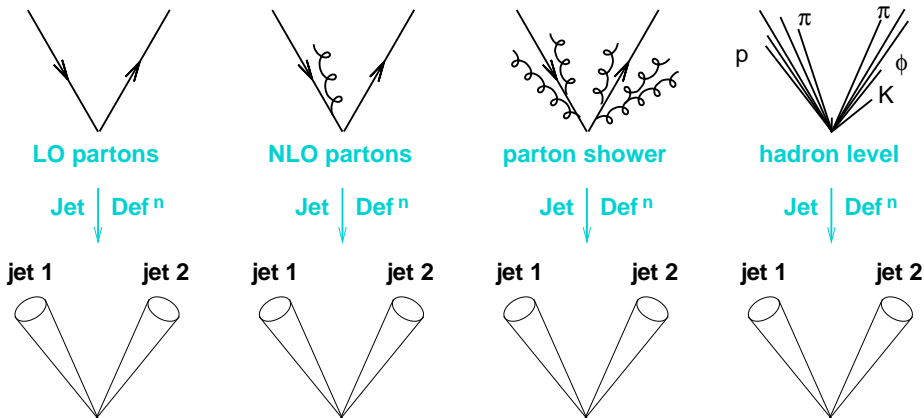
Quarks & gluons? We only ever see “jets”



Giving a pattern of hadrons that “remembers” the gluon branching

Hadrons mostly produced at small angle wrt $q\bar{q}$ directions or with low energy

Jets made systematic: jet definitions



LHC events may be discussed in terms of quarks, quarks+gluon, or hadrons

A **jet definition** provides common representation of different “levels” of event complexity.

QCD theorists have spent the past 10–15 years making accurate calculations of signals and backgrounds at the LHC, many of them with jets (with remarkable advances in field theory on the way)

$\mathcal{O}(100)$ people \times 10 years \simeq \$100 000 000

Problem 1: the jet definitions originally foreseen by LHC experiments were not compatible with these calculations — they “leaked” infinities:

$$\sigma = \sigma_{\text{LO}} (1 + c_1 \alpha_s + c_2 \alpha_s^2 + \infty \alpha_s^3 + \dots)$$

Problem 2: the jet definitions advocated by theorists since 1990's had been mostly shunned by proton-collider experiments

a) bad response to experimental noise

b) severe computational issues (1 minute/event $\times 10^{10}$ recorded events)

Discovered a link between QCD jet-finding and problems of 2D computational geometry

Cacciari & GPS '05

Jet clustering reduces to 2D dynamic nearest neighbour problem
time to cluster N particles reduced from $N^3 \rightarrow N \ln N$ (or $N^{\frac{3}{2}}$)

Developed a theory of the interplay between jet-finding, QCD radiation and experimental noise

Cacciari, GPS & Soyez '08

A crucial element was linearity of response

Spin-off applications for γ and lepton ID in Higgs searches

Proposed a new jet-definition based on what we'd learnt

anti- k_t

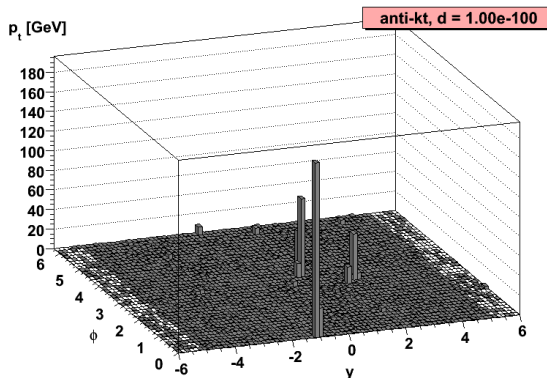
Cacciari, GPS & Soyez '08

simple agglomerative clustering algorithm

repeatedly cluster objects with smallest $d_{ij} = \frac{\Delta R_{ij}^2}{\max(k_{ti}^2, k_{tj}^2)}$

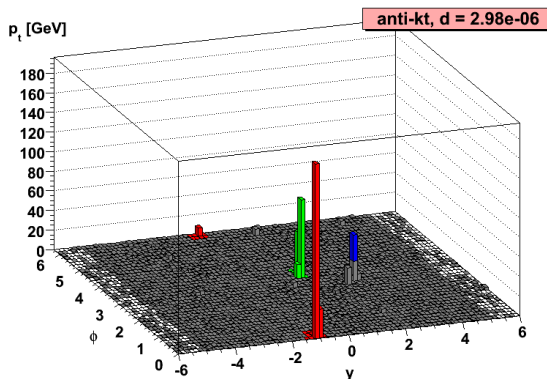
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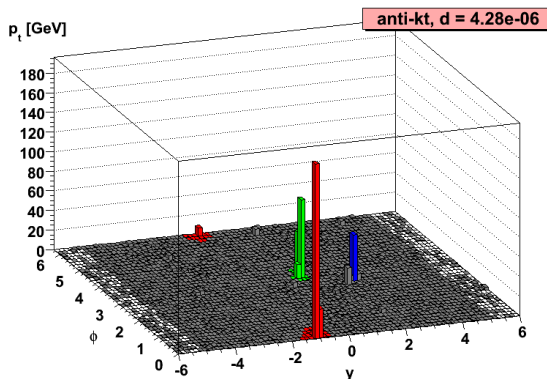
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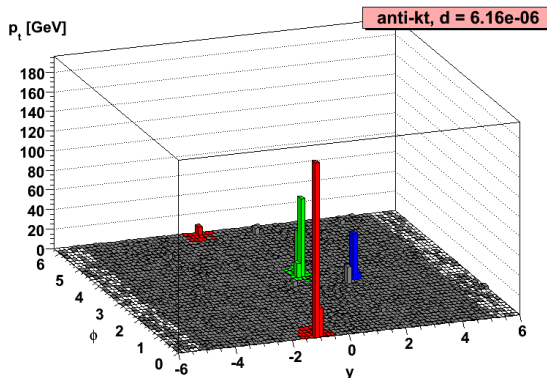
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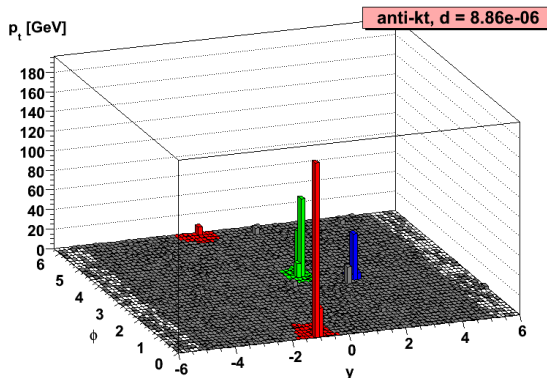
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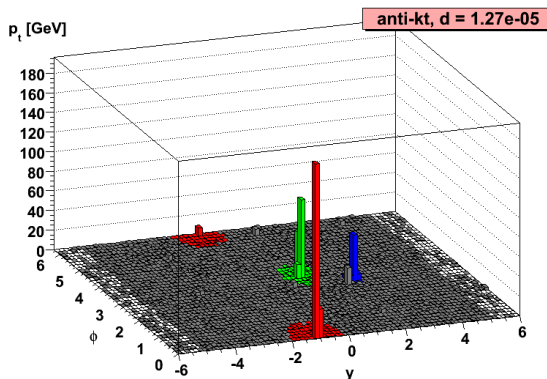
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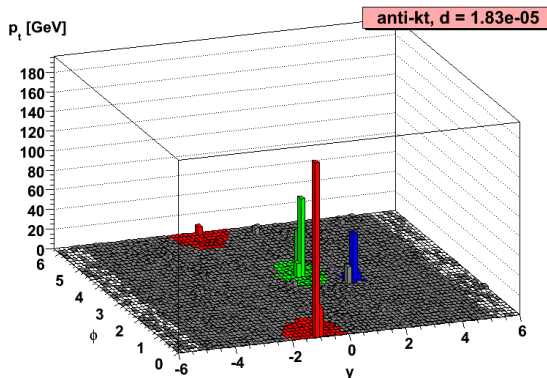
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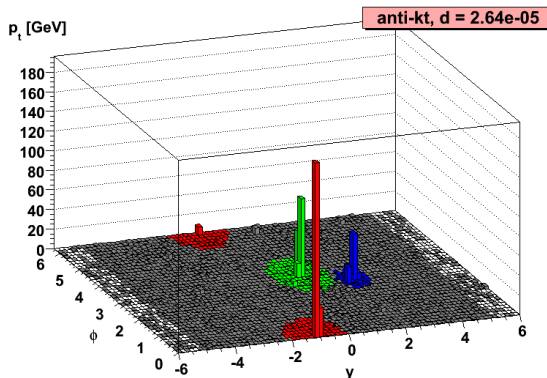
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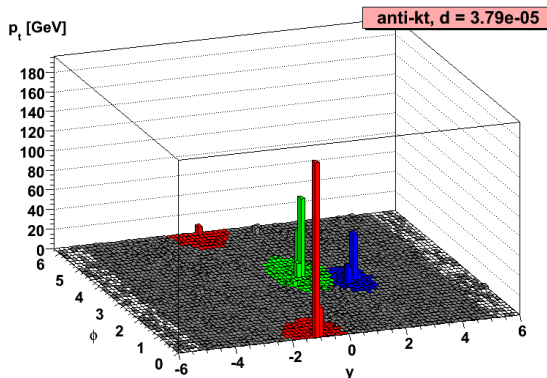
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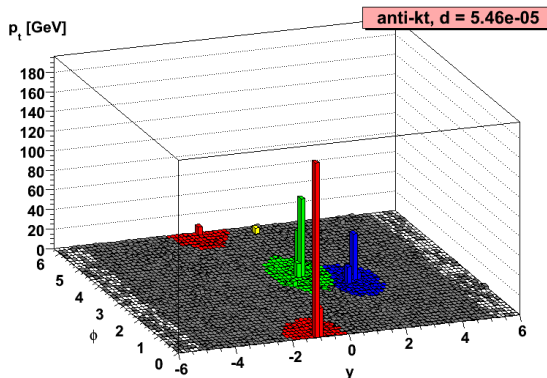
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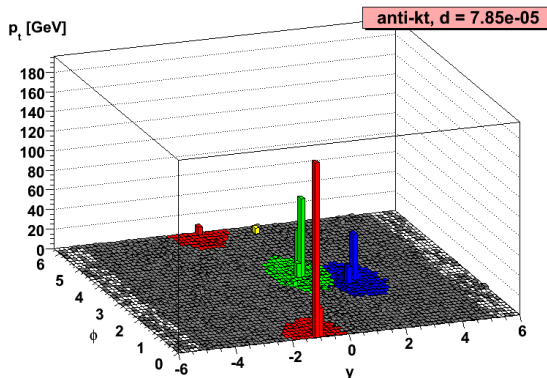
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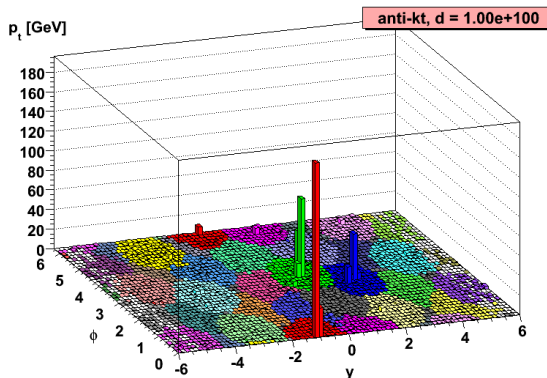
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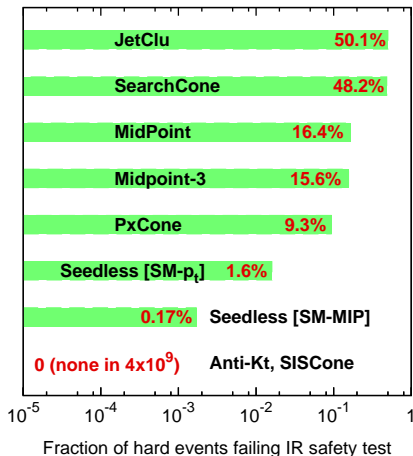


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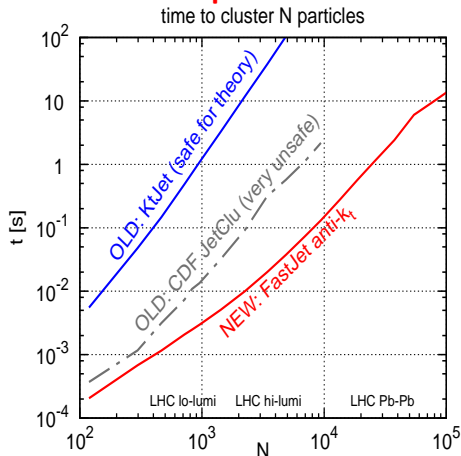


Coefficient of “infinity”



Safe for perturbative QCD:
No infinities

Speed



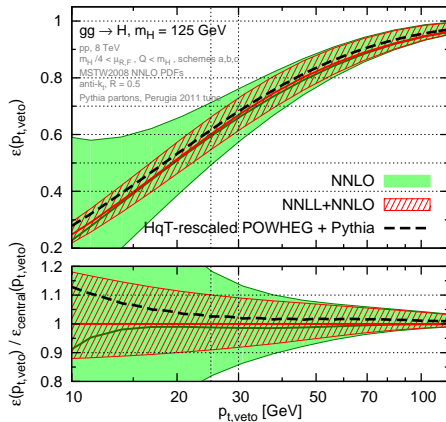
~ 1000 times faster than previous codes

e.g.: fraction of Higgs events that pass a “jet veto,” which matters for measuring $H \rightarrow WW$

anti- k_t used in nearly all jet measurements at the LHC

→ possible to make accurate predictions for range of measurements, including Higgs

allowing the experiments to pin down Higgs couplings more accurately in years to come



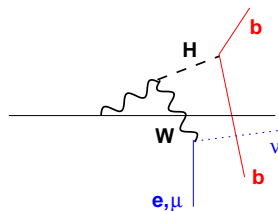
Banfi, Monni, GPS & Zanderighi '12
cf. talk tomorrow by Pier Monni

QCD is, in part, about predicting properties of collisions

But also about devising techniques to carry out more
effective searches

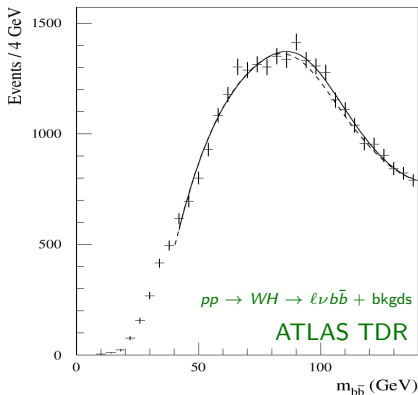
$H \rightarrow b\bar{b}$ (57% of decays) v. hard to see

Best hope is $pp \rightarrow W^\pm H$ (and ZH), $W^\pm \rightarrow \ell^\pm \nu$, $H \rightarrow b\bar{b}$.



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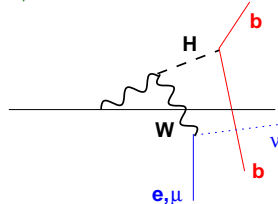


Conclusion (ATLAS TDR):

"The extraction of a signal from $H \rightarrow b\bar{b}$ decays in the WH channel will be very difficult at the LHC, even under the most optimistic assumptions [...]"

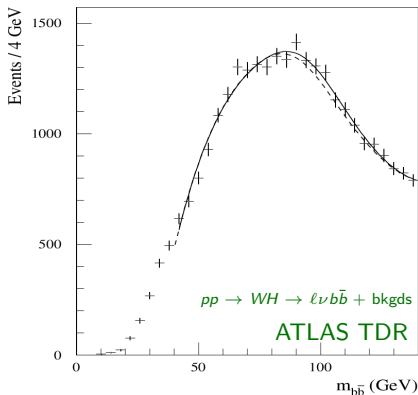
Low efficiency, huge backgrounds, e.g. $t\bar{t}$

NB: Evidence of this channel seen recently at Tevatron, but similar difficulties



$H \rightarrow b\bar{b}$ (57% of decays) v. hard to see

Best hope is $pp \rightarrow W^\pm H$ (and ZH), $W^\pm \rightarrow \ell^\pm \nu$, $H \rightarrow b\bar{b}$.



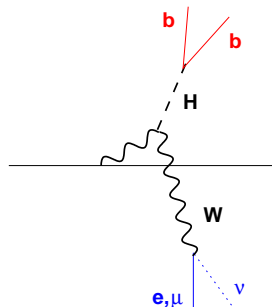
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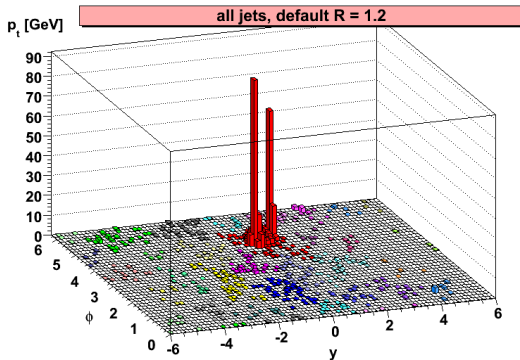
Low efficiency, huge backgrounds, e.g. $t\bar{t}$

Analysis of signal/bkgd suggests:

- ▶ Go to high p_t ($p_{tH}, p_{tW} > 200$ GeV)
- ▶ Lose 95% of signal, but more efficient?
- ▶ Maybe kill $t\bar{t}$ & gain clarity?



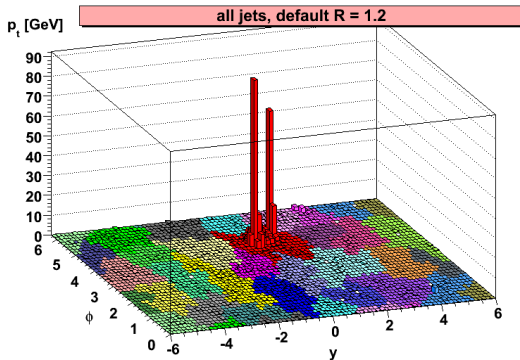
Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



Cluster event, C/A, R=1.2

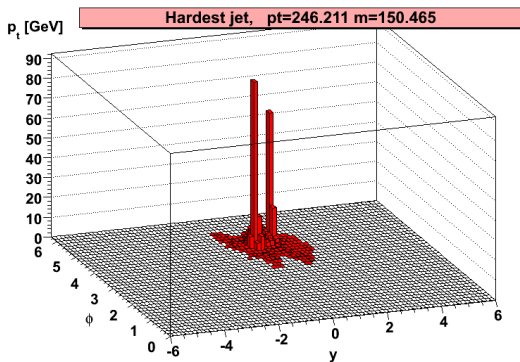
Butterworth, Davison, Rubin & GPS '08
also earlier work by Seymour; Butterworth et al

Herwig 6.510 + Jimmy 4.31 + FastJet 2.3

Fill it in, \rightarrow show jets more clearly

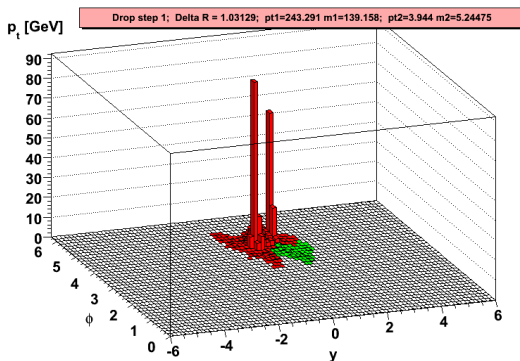
Butterworth, Davison, Rubin & GPS '08
 also earlier work by Seymour; Butterworth et al

Herwig 6.510 + Jimmy 4.31 + FastJet 2.3

Consider hardest jet, $m = 150$ GeV

Butterworth, Davison, Rubin & GPS '08
 also earlier work by Seymour; Butterworth et al

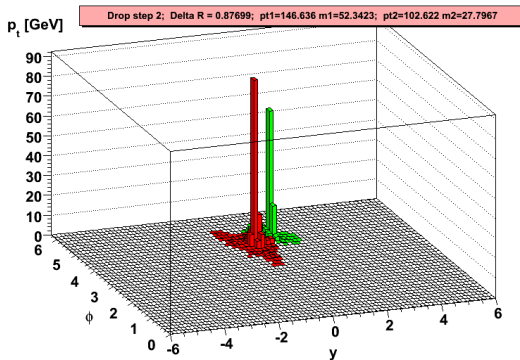
Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



split: $m = 150$ GeV, $\frac{\max(m_1, m_2)}{m} = 0.92 \rightarrow$ repeat

Butterworth, Davison, Rubin & GPS '08
also earlier work by Seymour; Butterworth et al

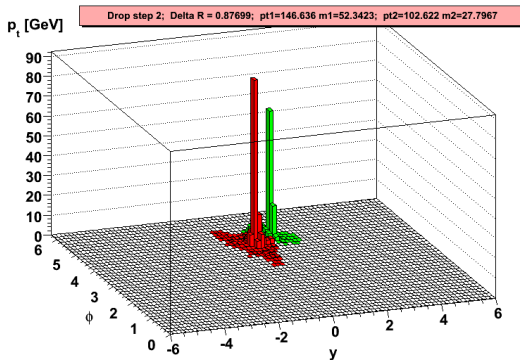
Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



split: $m = 139$ GeV, $\frac{\max(m_1, m_2)}{m} = 0.37 \rightarrow$ mass drop

Butterworth, Davison, Rubin & GPS '08
also earlier work by Seymour; Butterworth et al

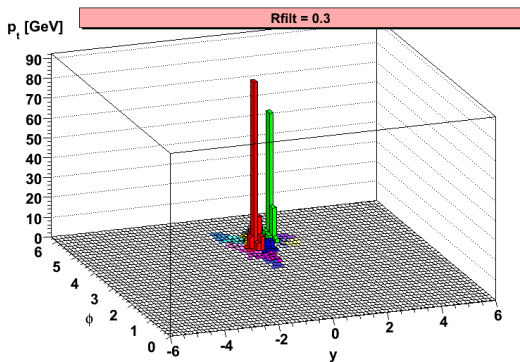
Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



check: $y_{12} \simeq \frac{p_{T2}}{p_{T1}} \simeq 0.7 \rightarrow \text{OK} + 2 \text{ } b\text{-tags (anti-QCD)}$

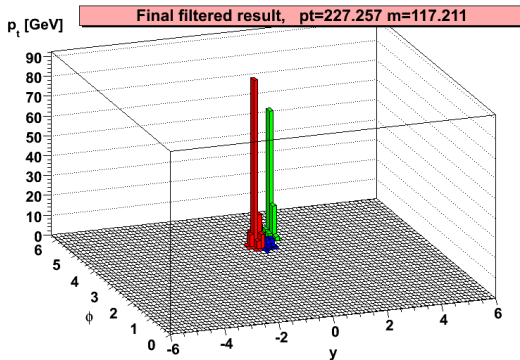
Butterworth, Davison, Rubin & GPS '08
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 $R_{filt} = 0.3$

Butterworth, Davison, Rubin & GPS '08
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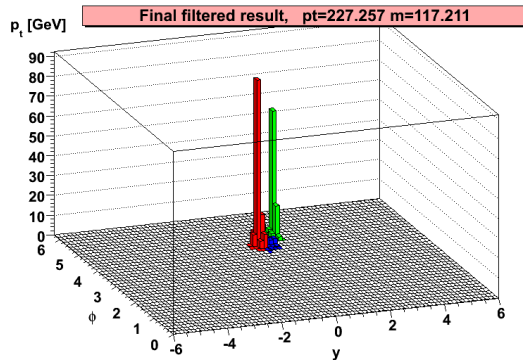
Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



$R_{filt} = 0.3$: take 3 hardest, $m = 117$ GeV

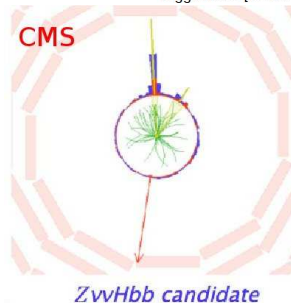
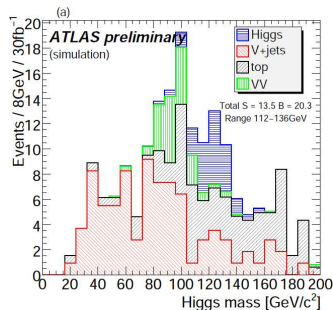
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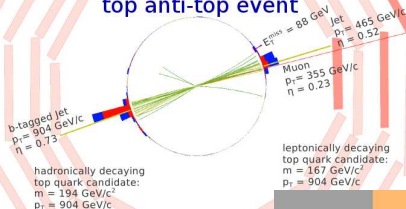
Butterworth, Davison, Rubin & GPS '08
also earlier work by Seymour; Butterworth et al





CMS Experiment at LHC, CERN
Data recorded: Tue, Aug 9 13:57:08 2011 CEST
Run/Event: 172952 /1031053741
Lumi section: 887

~2 TeV (boosted)
top anti-top event

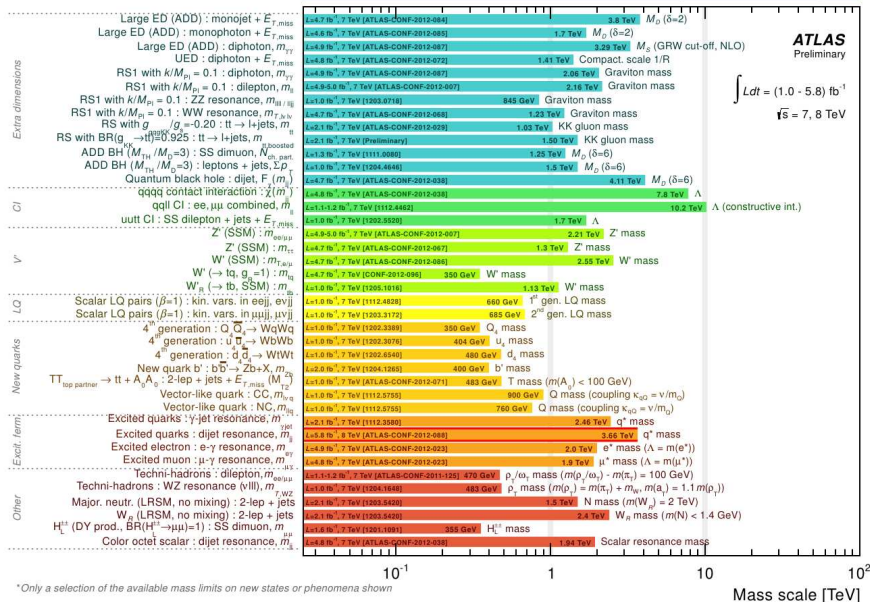


“substructure” techniques have applications in many searches, e.g. $t\bar{t}$ resonances

very active field, (including MIT group!)

	Traditional		Jet substructure	
ATLAS	$e+e, \mu+\mu$ and $e+\mu$ 6.5%			
Article/ Note	arXiv: 1205.5371	arXiv: 1205.5371	arXiv: 1207.2409	ATLAS-CONF-2012-102
Integrated Luminosity	2 fb ⁻¹	2 fb ⁻¹	2 fb ⁻¹	4.7 fb ⁻¹
Z' limits $\Gamma/m = 1.2\%$	-	0.5-0.88 TeV	0.6-1.15 TeV	0.7-1.3 TeV
KKG limits $\Gamma/m = 15.3\%$	0.5-1.08 TeV	0.5-1.13 TeV	0.6-1.5 TeV	0.7-1.5 TeV

An overview of some new-physics searches

ATLAS Exotics Searches* - 95% CL Lower Limits (Status: ICHEP 2012)


rest of 2012

continue running at 8 TeV to reach $\sim 30 \text{ fb}^{-1}$ per experiment (a factor of 3 more data than used for discovery). Higgs features should emerge more clearly

2013

short proton-lead run

followed by ~ 18 month shutdown to complete LHC repairs

late 2014-

Running at 13 – 14 TeV

Accumulating several 100 fb^{-1} over a few years

Cover a large chunk of LHC's potential to search for new physics

Beyond

LHC luminosity upgrades: factor 5-10?

Linear collider (to study Higgs in detail)?

Higher-energy LHC?

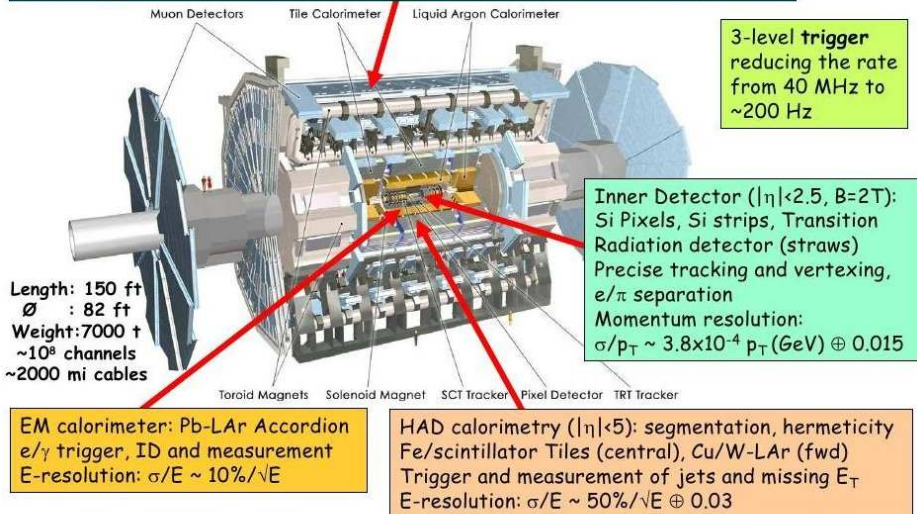
EXTRAS

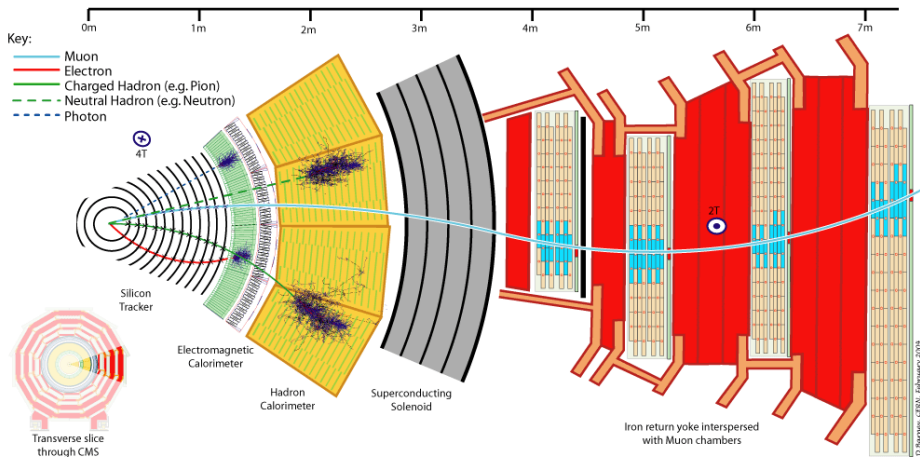
Energy–frontier colliders of the past 25 years

Collider	Lab	Date	Collided		C.o.M. Energy
Tevatron	Fermilab/USA	1987 –	$p\bar{p}$	@	1960 GeV
SLC	SLAC/USA	1989 – 1998	e^+e^-	@	100 GeV
LEP	CERN/Europe	1989 – 2000	e^+e^-	@	209 GeV
HERA	DESY/Germany	1992 – 2007	$e^\pm p$	@	330 GeV

Protons are made of quarks, anti-quarks and gluons. It's the individual quarks and gluons that collide. Only a fraction of the proton's energy is actually available in a single *quark/gluon* collision.

Muon Spectrometer ($|\eta| < 2.7$): air-core toroids with gas-based muon chambers
Muon trigger and measurement with momentum resolution $< 10\%$ up to $E_\mu \sim 1$ TeV





1995: LHC approved

2000: LEP closed

2008/09: LHC beams circulated

2008/09: severe “incident”

poor electrical connection, arc, catastrophic Helium release, much damage

Followed by reviews, the fixes that could be made in ~ 1 year

2009/11: LHC starts up again, 900 GeV pp collisions

2009/12: 2360 GeV pp collisions

2010/03: 7000 GeV pp collisions ($\sim 50 \text{ pb}^{-1}$)

“reduced-energy” target for safe operation

2010/11: 2760 GeV PbPb collisions

2011/03: 7 TeV pp collisions ($\sim 5 \text{ fb}^{-1}$)

2011/11: 2760 GeV PbPb collisions

2012/03: 8 TeV pp collisions ($\sim 6 \text{ fb}^{-1}$ published, $\sim 14 \text{ fb}^{-1}$ delivered)

Circular e^+e^- collider. Basic issue is synchrotron radiation

$$\text{Energy loss per orbit} \sim \frac{E^4}{m^4 R}$$

At LEP the numbers are $\mathcal{O}(10\%)$ of the electron's energy per orbit.

Circular pp collider

Proton mass 2000 times larger, so synchrotron radiation not a problem.
The limitation is magnetic field needed to bend the protons round

$$B \sim \frac{E}{R}$$

Tevatron: $R \sim 1 \text{ km}$, $E_{c.o.m} \sim 2 \text{ TeV} \implies B = 4 \text{ TeV}$.

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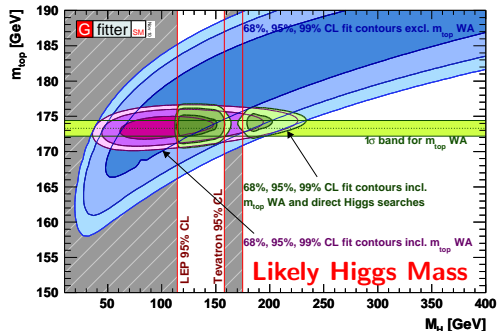
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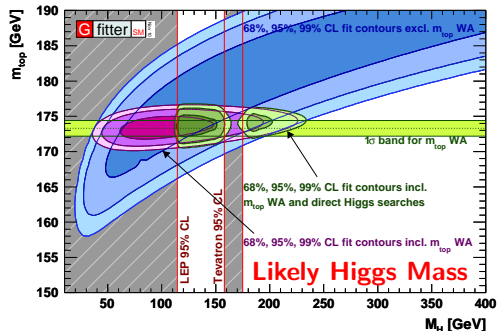
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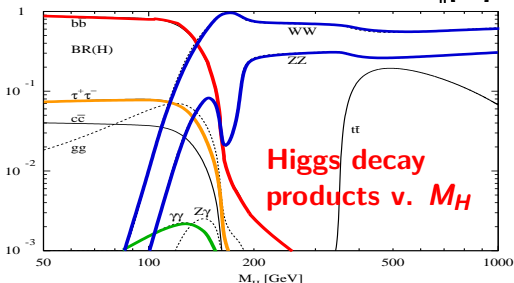
So what energy do you need?



There's some likelihood that the Higgs boson will be "light", $M_H \sim 120$ GeV



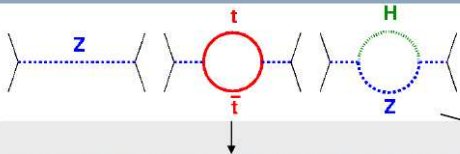
There's some likelihood that the Higgs boson will be "light", $M_H \sim 120$ GeV



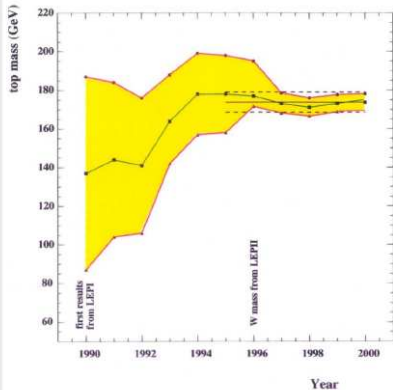
If it is, crucial test of whether it **is** the Higgs, will come from measuring several different decays

Remember: Higgs couplings intimately related to origin of particle masses

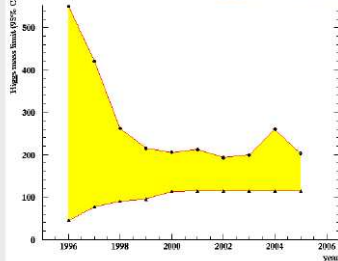
Test of the SM at the Level of Quantum Fluctuations



indirect determination of the top mass



prediction of the range for the Higgs mass

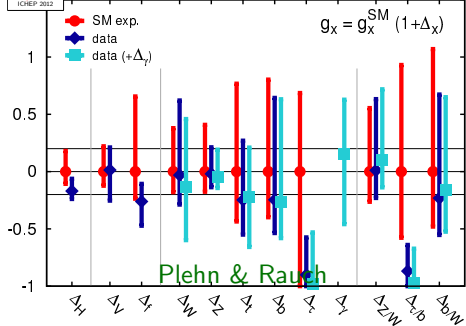


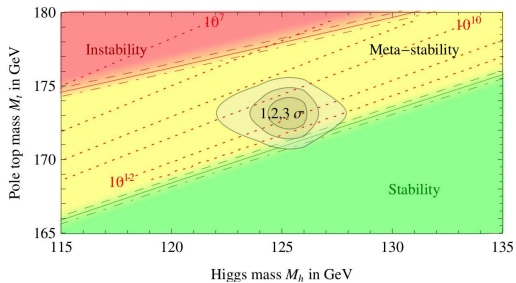
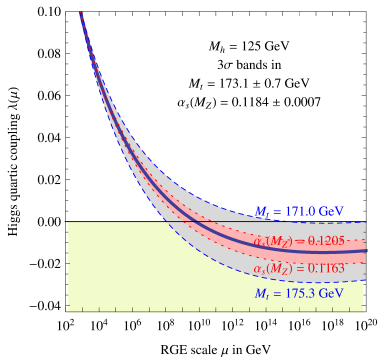
possible due to

- precision measurements
- **known higher order electroweak corrections**

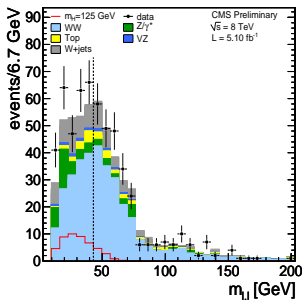
$$\propto \left(\frac{M_t}{M_W} \right)^2, \ln \left(\frac{M_h}{M_W} \right)$$

$L=4.6-5.1(7\text{ TeV})+5.1-5.9(8\text{ TeV})\text{ fb}^{-1}$, 68% CL: ATLAS + CMS



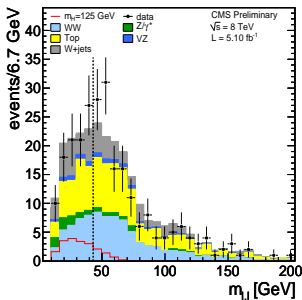


0-jet events



mainly WW^*
background

1-jet events

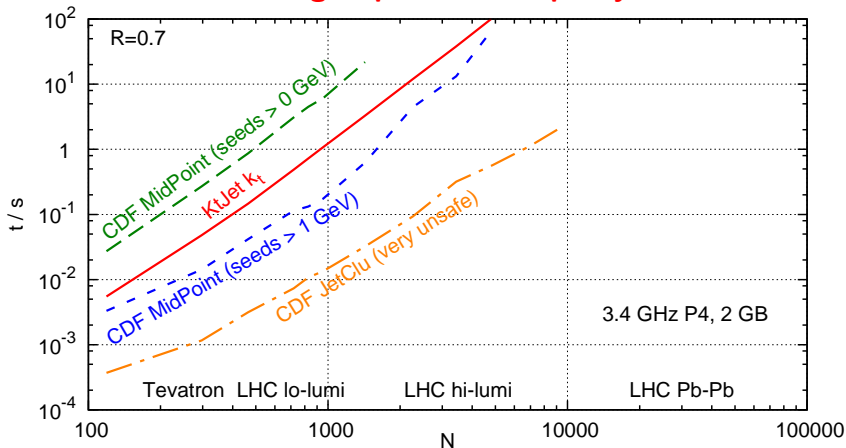


also some $t\bar{t}$
background

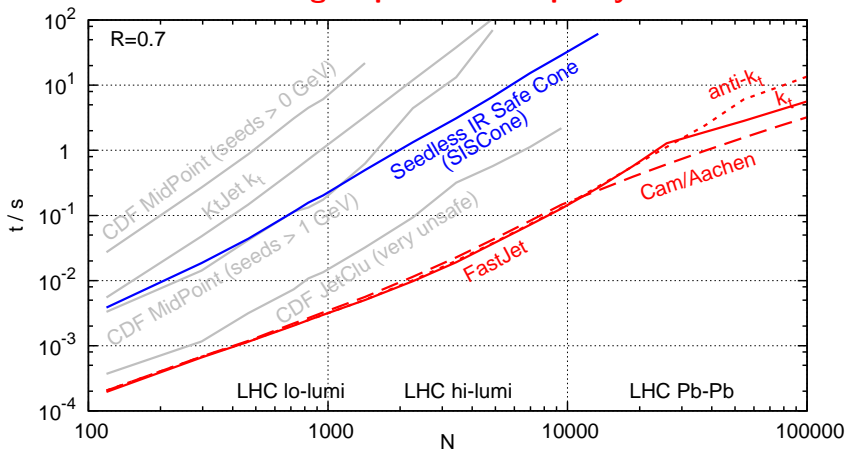
≥ 2 -jet events

not used
too much background

Timing v. particle multiplicity 2005

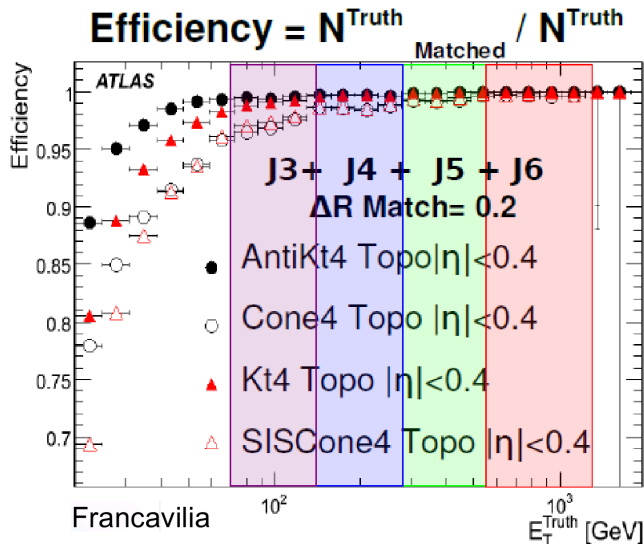


Timing v. particle multiplicity 2008

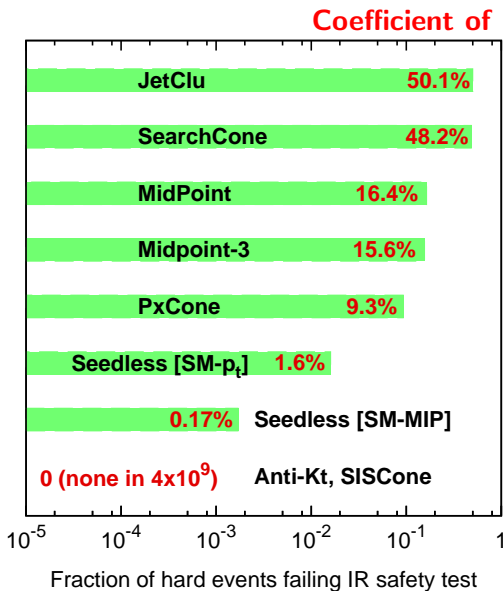


in critical region of $N \sim 2000 - 4000$
1000 times faster than previous attempts with similar jet algorithms

Experimental sensitivity to noise



As good as, or
 better than all
 previous
 experimentally-
 favoured
 algorithms



Safe for perturbative QCD predictions:

No “leakage” of infinities to higher orders