

# Characterising non-perturbative effects in jets

Gavin Salam

LPTHE, Universities of Paris VI and VII and CNRS

work in progress with M. Cacciari  
+ close links with with M. Dasgupta & L. Magnea

Ringberg workshop on non-perturbative QCD of jets  
8–10 January 2007

Much work done on non-perturbative effects in  $e^+e^-$  and DIS event shapes. But little understood about jets.

**Webber hep-ph/9510283**: 3-jet resolution,  $y_3$ , gets  $\Lambda^2/Q^2$  corrections  
'Higher' orders give  $\sqrt{y_3}\Lambda/Q$  or  $\sqrt{y_3}\ln y_3\Lambda/Q$

**Seymour, NPB513(1998)269**: differential jet shape at angular distance  $r$  from jet axis gets correction  $\frac{\Lambda}{r^2 p_T}$

**Mangano, hep-ph/9911256**: hadron-collider inclusive jet-spectrum gets a roughly  $p_T$ -independent shift of order  $\Lambda$ .

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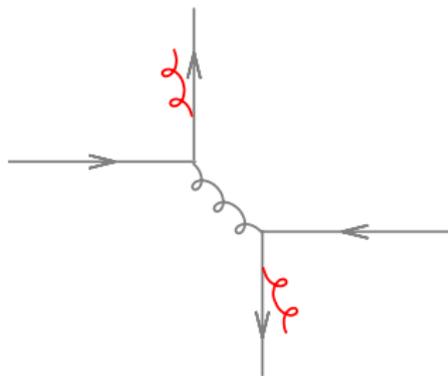
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▶ **Underlying event:**

emissions from proton remnants, (multiple) interaction between two proton remnants.

▶ **Pileup:**

at high luminosity, contribution from simultaneous  $pp$  collisions in the same bunch crossing.



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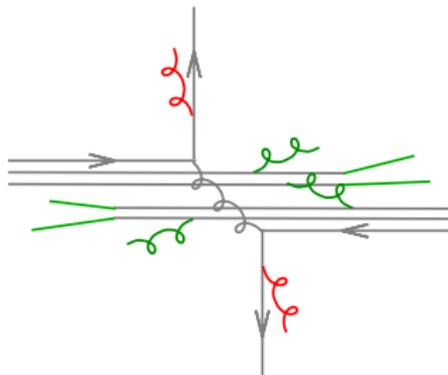
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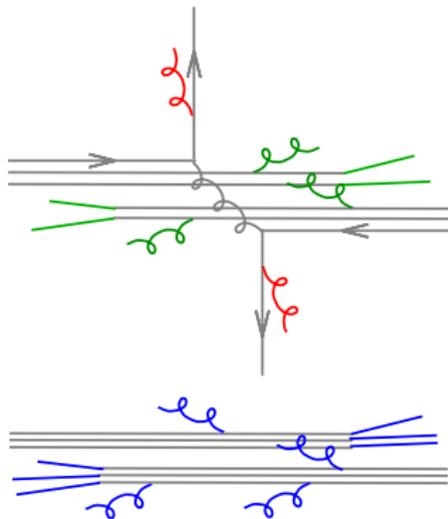
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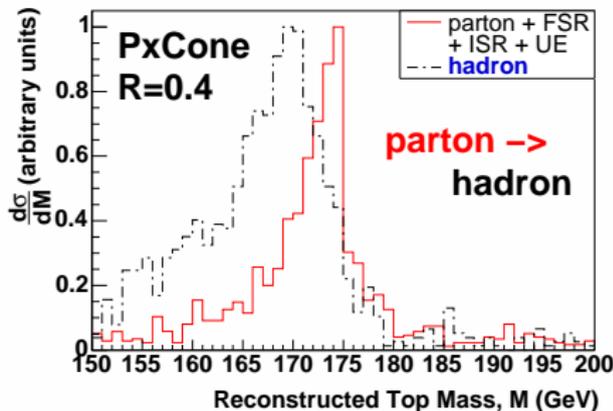
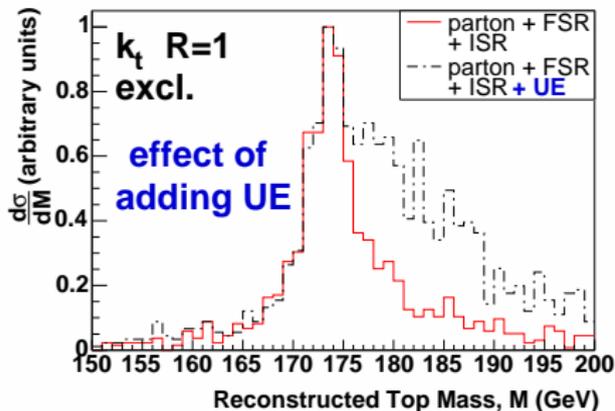
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- ▶  **$k_t$ :**  
Combine pair of particles closest in  $k_t$ -distance; repeat until all particles separated by angular ( $\Delta R^2 = \Delta y^2 + \Delta \phi^2$ ) distance  $> R$  [inclusive] or  $k_t$  distance  $> d_{cut}$  [exclusive].  
Catani et al '93; Ellis & Soper '93
- ▶ **Cambridge/Aachen:**  
Combine pair of particles closest in angular-distance; repeat until all particles separated by ang. dist.  $> R$  [inclusive], or make a particle into jet if about to cluster with harder particle and  $k_t$  dist.  $> d_{cut}$  [exclusive].  
Dokshitzer et al '97; Wobisch & Wengler '99
- ▶ **Cone:**  
Find 'stable cones' of half-angle  $R$ ; run a split-merge procedure on stable cones that overlap so as to get final jets. Serman & Weinberg '77  
Many variants since then...  
Seedless IR Safe cone (SISCone): GPS & Soyez '07



adapted from Seymour & Tevlin '06

Common statements:

- ▶  $k_t$  has larger UE & pileup corrections
- ▶ cone has larger hadronization corrections

But  $k_t$  and cone often used with different parameters ( $R = 1$  v.  $R = 0.4$ ).

**Can we get analytic understanding of parameter & algorithm dependence for hadronisation and UE/pileup effects?**

We will try to calculate N.P. corrections to **jet transverse momentum**.

Easily related, e.g., to mass reconstructions

Starting point, as for many NP-calculations, is 1 hard parton (jet) + 1 soft gluon:

- ▶ This is a valid approximation only if the observable is **linear** in effects of multiple soft momenta. cf. Milan factor, Dokshitzer et al. '97-'98  
crucial input in Lee & Sterman '06
- ▶ Many  $e^+e^-$  & DIS event shapes had some form of linearity.
- ▶ Jet algorithms **are not linear**.  
But 1-gluon approx. may still be useful for getting first picture
- ▶  $k_t$ , Cam/Aachen & cone are **identical** @ 1 soft-gluon level

Assume soft gluon produced uniformly in  $y$  (rapidity) and  $\phi$  with transv. mom. density (averaged over many events):

$$\left\langle \frac{dp_{t,NP}}{dy d\phi} \right\rangle = \rho_{U.E.} \sim \Lambda, \quad \text{or} \quad \rho_{P.U.} \sim n_{P.U.} \Lambda,$$

independently of hard event (marginal for U.E.? Fine for P.U.).

The soft gluon (g) will be clustered into jet (j) if  $\Delta R_{gj} < R$ . This defines a *jet area*  $A$  in  $y, \phi$  space,  $A = \pi R^2$ , and the jet  $p_t$  is increased proportionally to its area:

$$\Delta p_{t,jet,UE} = \pi R^2 \rho_{UE} \quad (\text{P-scheme})$$

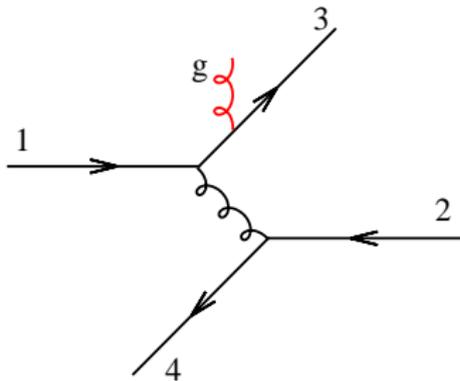
$$\Delta p_{t,jet,UE} = 2\pi R J_1(R) \rho_{UE} = \left( \pi R^2 - \frac{\pi}{8} R^4 + \dots \right) \rho_{UE} \quad (\text{E-scheme})$$

Note:  $\mathcal{O}(R^4)$  depends on recombination scheme

Universality & Milan factor: calculate hadronisation by calculating effect of a *trigger gluon* (gluer)  $k$  on the observable. [keeping it simple!]

$$\delta V = C \sum_{\text{dipoles}} \int d\eta_{k,dip} d\phi_{k,dip} dk_{t,dip} \delta(k_{t,dip} - \Lambda) (V(\{\tilde{p}_i\}, k) - V(\{p_i\}))$$

with  $C$  known from many event shapes in  $e^+e^-$ :  $C\Lambda \simeq 0.5 \text{ GeV}$ .



NB: recoiled hard momenta  $\{\tilde{p}_i\}$  v. orig.  $\{p_i\}$ .

Event shapes:  $V(\{p_i\}) = 0$ , recoil irrelevant;

For jets:  $V(\{p_i\}) = p_{t,3} \neq 0$   
 $V(\{\tilde{p}_i\}, k) = \tilde{p}_{t,3}[+k_t]$

∃ **ambiguity** in decision about how to assign  $k$ 's recoil between  $\tilde{p}_3$  and  $\tilde{p}_4$

Recoil ambiguity foils any 'traditional' calculation of hadronization corrections to jet  $p_t$ 's.

Similar issue e.g. for thrust in 3-jet region

Two approximate solutions:

- ▶ Go to threshold limit (recoil uniquely defined) ↳ talk by Magnea
- ▶ Consider only small  $R$ : hadronisation dominated by gluon emission close to hard parton; *assume* recoil dominantly taken by that hard parton.

gluon in jet:  $p_{t,jet} = k_t + \tilde{p}_{t,3} = p_{t,3}$

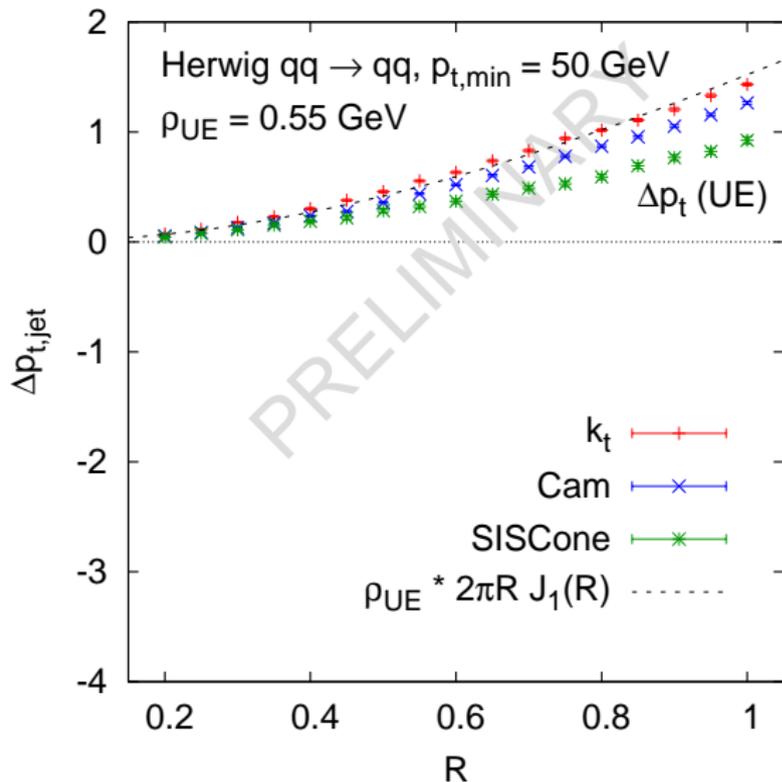
gluon out of jet:  $p_{t,jet} = \tilde{p}_{t,3} = p_{t,3} - k_t$

$$\delta p_{t,jet} = C \int^{-\ln \tan R/2} d\eta_{dip} (-\Lambda \sinh \eta_{dip}) = C\Lambda \left( -\frac{1}{R} + \mathcal{O}(1) \right)$$

Gluonic jet has extra factor  $C_A/C_F$

$1/R$  structure coincides with threshold result by Dasgupta & Magnea

Less accurate than D&M, but holds regardless of event structure



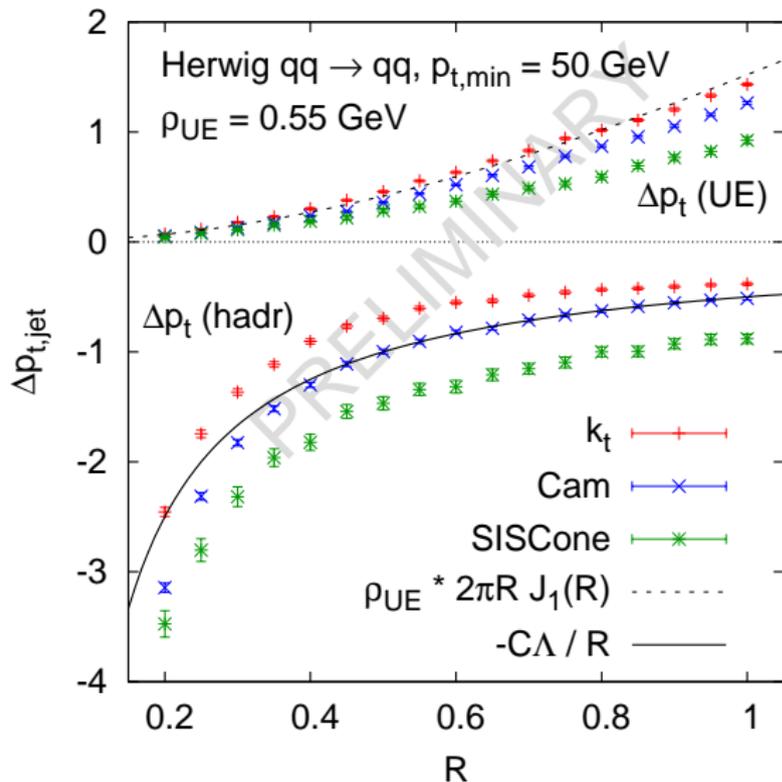
Analytical results have strong  $R$  dependence, but *do not depend on jet algorithm*.

Compare to MC:

- ▶ **Broad features agree** with MC
- ▶  $R$ -dependence deviates a little
- ▶ moderate jet. alg. dependence is present:

$$k_t > \text{Cam} > \text{Cone}$$

NB: normalisations depend on how one selects jets



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Jet algorithms are identical at level of 1 soft gluon. Can we understand nature of differences beyond 1 gluon?

Study just UE and pileup:

- ▶ They are easier, since no recoil to worry about
- ▶ UE larger than appears from previous page, often dominant  
default Herwig underestimates it
- ▶ Pileup will be huge at LHC, and will dominate over other effects.  
20  $pp$  interactions per bunch crossing

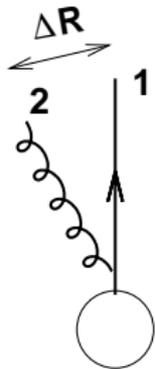
*BUT: don't study U.E., pileup effect directly.* Instead assume PT content of jet is independent of U.E. & pileup, so that effect of U.E. & pileup is proportional to **jet area, A**:

$$\Delta p_{t,jet} = \rho A$$

Consider jet composed of two  $p_t$ -ordered perturbative partons,

$$p_{t1} \gg p_{t2} \gg \Lambda$$

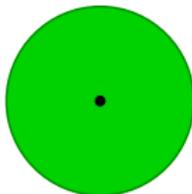
separated by  $\Delta R$ . Scan a NP gluon, '*ghost*', over the  $y$ - $\phi$  plane, and see when it goes into the jet containing  $p_1$ . From this deduce the jet area.



$k_t$



Cam/Aachen



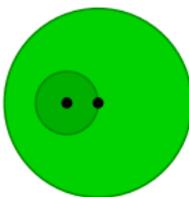
cone



$k_t$



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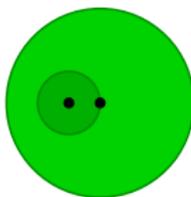


$$\Delta R = R/3$$

$k_t$



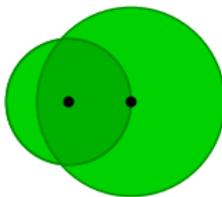
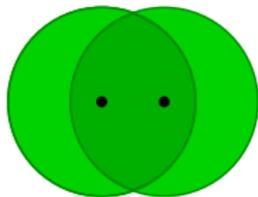
Cam/Aachen



cone



$$\Delta R = R/3$$



$$\Delta R = 2R/3$$

$k_t$



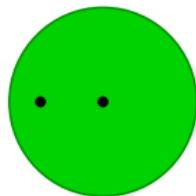
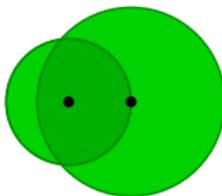
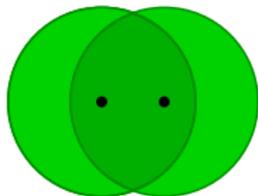
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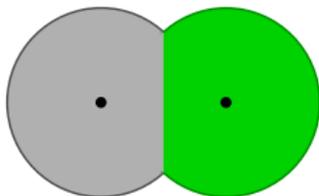
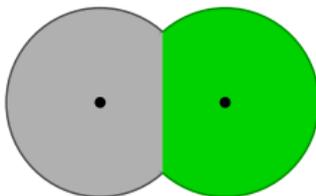
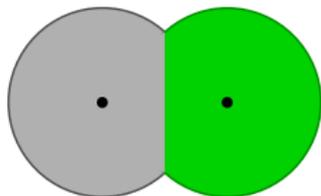
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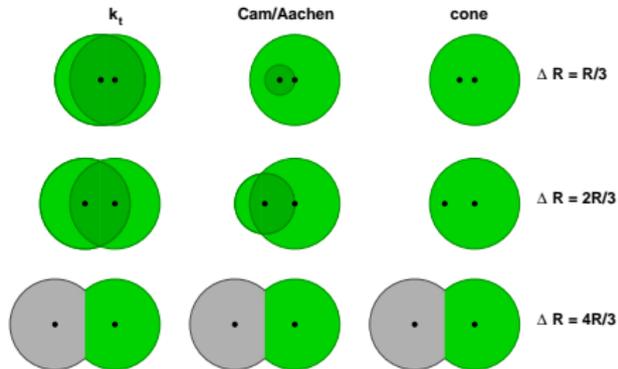
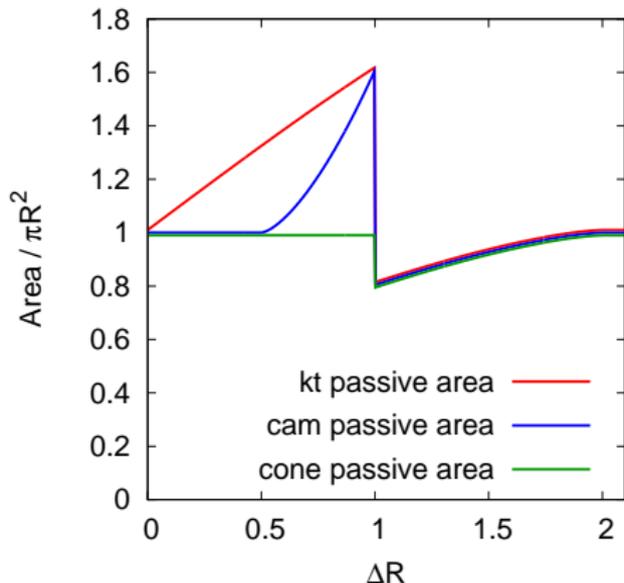
$$\Delta R = R/3$$



$$\Delta R = 2R/3$$



$$\Delta R = 4R/3$$



NB: difference in areas is independent of softness of  $p_{t2}$ .

$$\langle \Delta A \rangle = \frac{2\alpha_s C_F}{\pi} \int_{\Lambda}^{p_{t1}} \frac{dp_{t2}}{p_{t2}} \int_0^{2R} \frac{d\Delta R}{\Delta R} \Delta A(\Delta R)$$

Suppose incoming partons (colour charge  $C_i$ ) and outgoing jets (col. charge =  $C_o$ ) are not colour connected.

Mean outgoing jet area  $\langle A \rangle$  depends on jet  $P_t$  as follows:

$$\langle A \rangle = R^2 \left( \pi + (a_0 C_o + a_2 C_i R^2) \frac{\alpha_s}{\pi} \ln \frac{p_{t1}^2}{\Lambda^2} + \mathcal{O}(\alpha_s, \alpha_s^2 L^2) \right)$$

Have neglected  $\mathcal{O}(C_o R^2)$  term  
 $\alpha_s^n \ln^n p_t / \Lambda$  terms build up anomalous dimension

	$a_0$	$a_2$	comment
$k_t$	+1.771	+0.325	significant, positive
Cam / Aachen	+0.249	0	small, positive
Cone	-0.200	-0.325	small, negative

For  $\Lambda \sim 10$  GeV (pileup),  $P_t \sim 100 - 1000$  GeV,  $\frac{\alpha_s}{\pi} \ln P_t^2 / Q_0^2 \sim 0.2 - 0.4$

**NB: ordering of algorithms is that seen in MC**

## Passive area

- ▶ Having just 1 NP gluon in event is convenient analytically
- ▶ But not very realistic
- ▶ In presence of many NP gluons, approx. is equivalent to pretending NP gluons don't cluster between each other: **passive area**

## Active area

- ▶ Throw in  $\mathcal{O}(10^4)$  NP *'ghost'* gluons ( $10^{-100}$  GeV)
- ▶ Run clustering on event including ghosts
- ▶ Count how many ghosts end up in each jet — this is a more realistic measure/definition of area: **active area**

To run on  $10^4$  particles requires fast clustering  
 $k_t$  & Cam: FastJet [Cacciari & GPS '05]  $\sim N \ln N$   
cone more difficult: SIScone  $\sim N^2 \ln N$

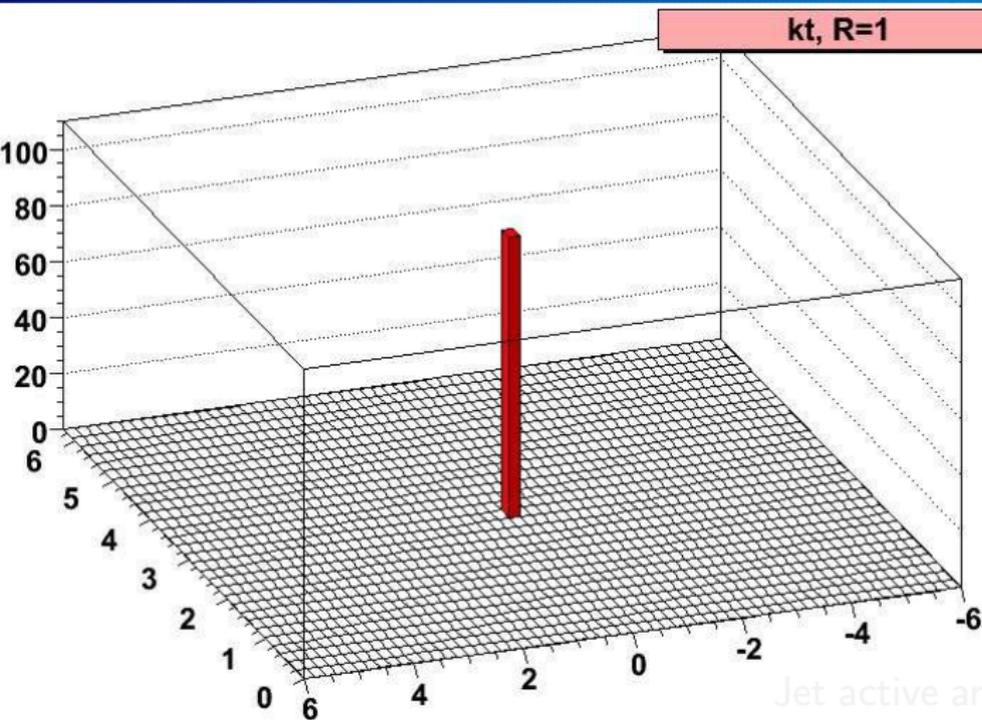
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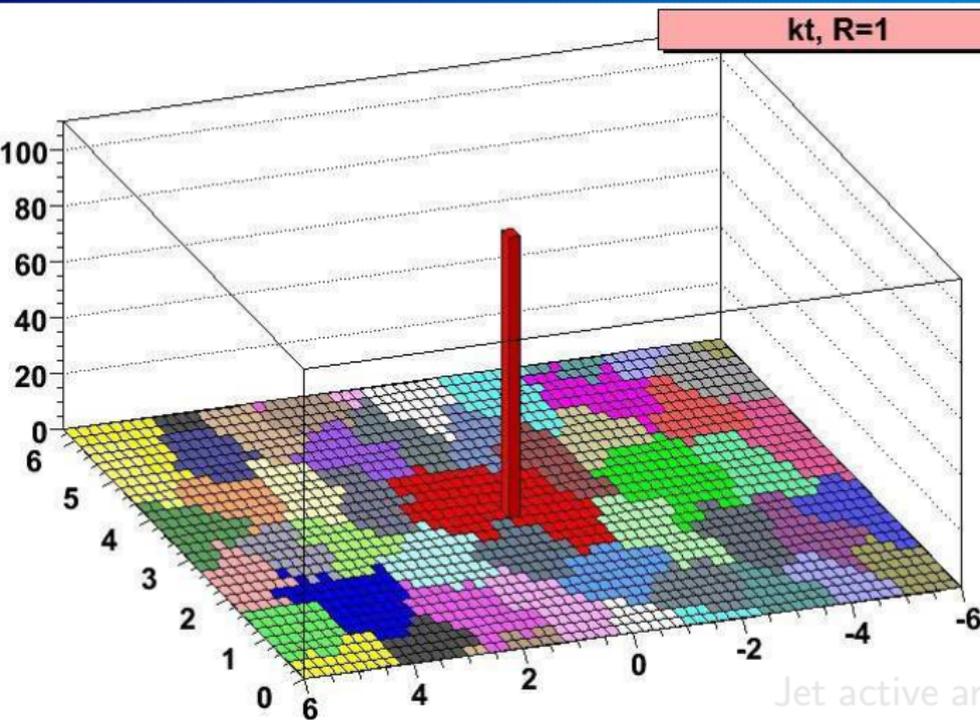


1 hard +  $10^4$  ghosts

- ▶ pure ghost jets  
model pileup jets
- ▶ hard jet has  
irregular  
boundary
- ▶ perturb ghosts  
→ change jet  
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Jet active area only defined after  
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Fluctuations in soft/hard jet active areas provide a model for fluctuations in pure pileup jets and of pileup in narrow hard jets.

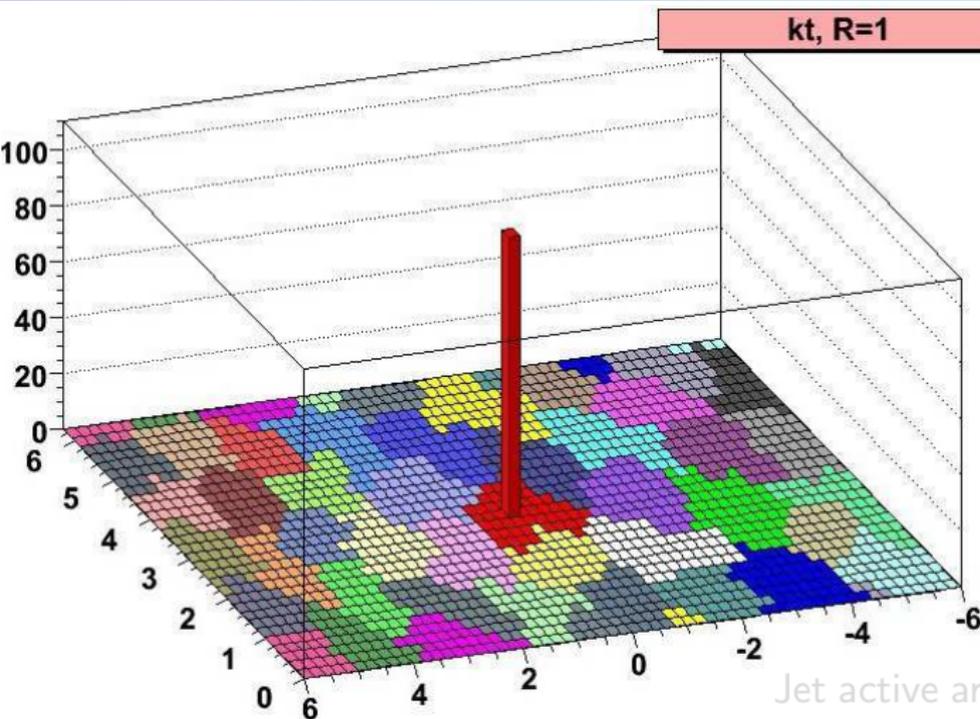


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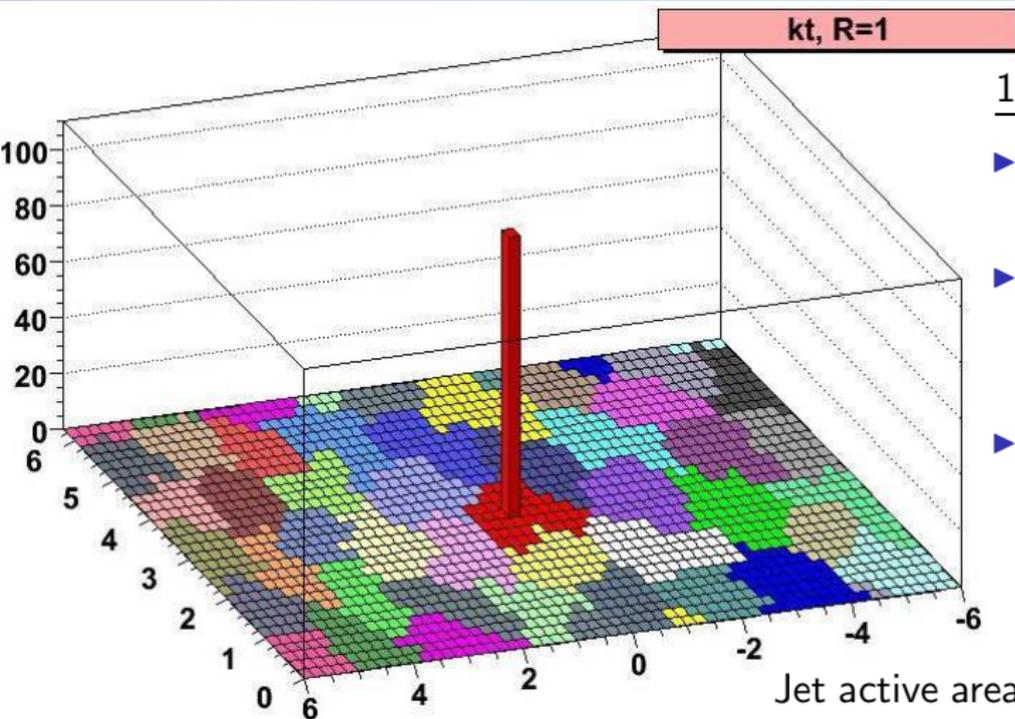


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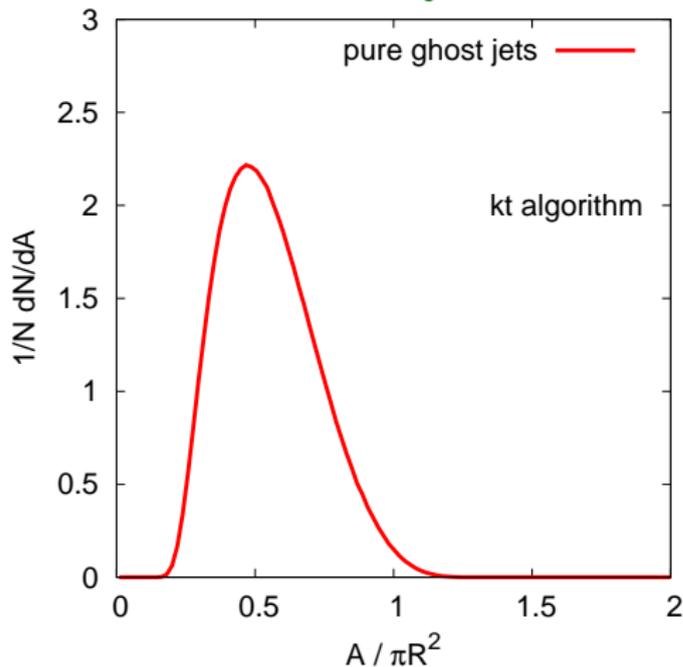
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## Distribution of jet areas



### Ghost v. hard jets:

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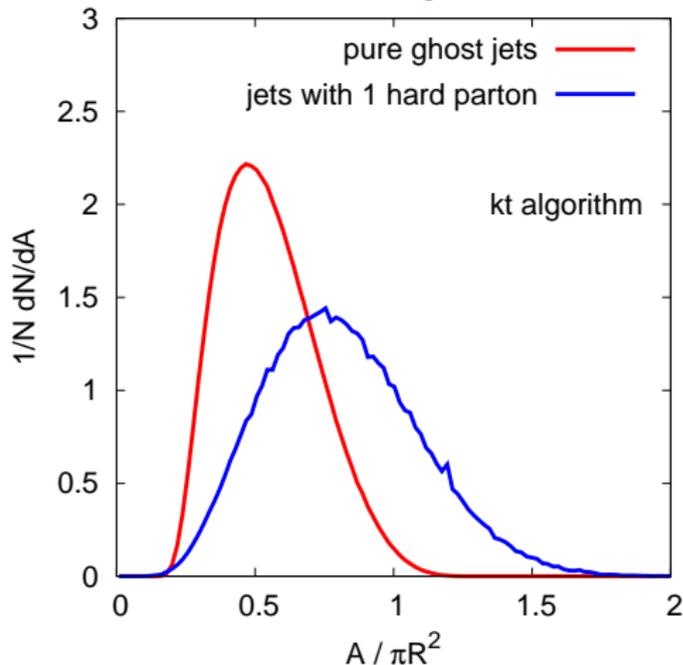
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Cone still being studied

### Conclusions:

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- ▶ Jet area expands when it is anchored by a hard parton.

- ▶ Can one obtain analytical insight into this? To some extent in 1D
- ▶ 'Hierarchical clustering' is used in many fields (bio, computing, ...) — are similar features of relevance there?

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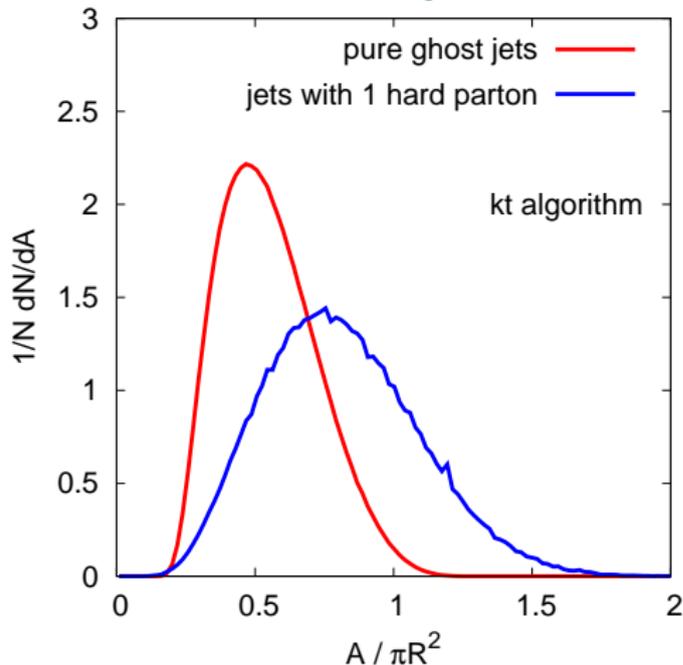
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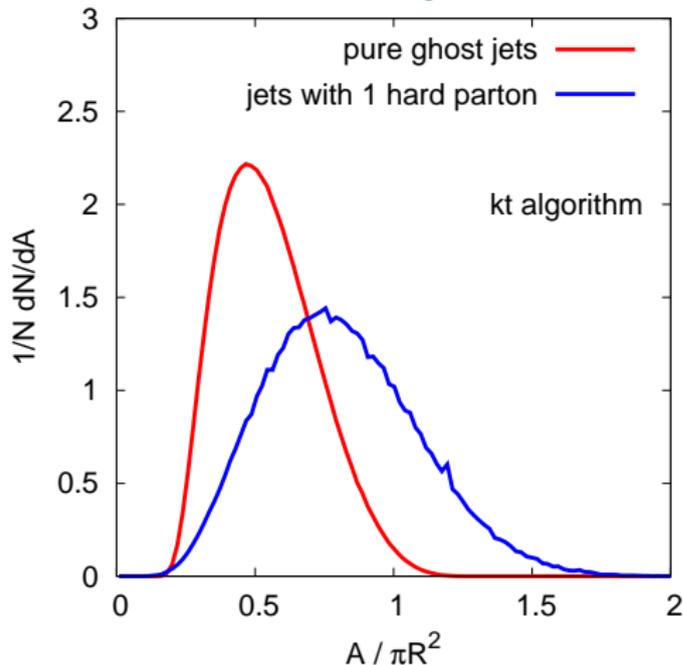
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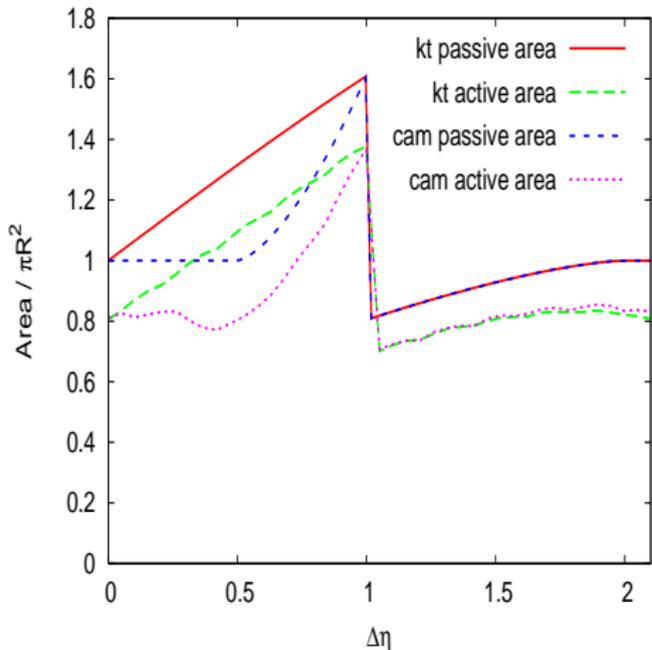
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Put 1 hard PT gluon, 1 soft PT gluon (separated by  $\Delta R$ ), as before.

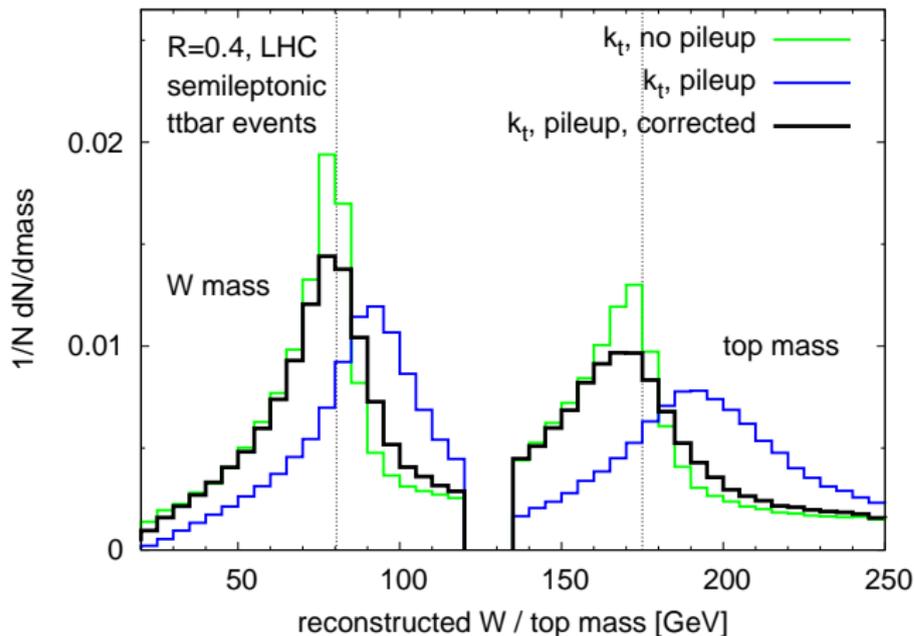
Calculate passive and active areas.

Picture is same for both, but  $\sim$  rescaled...

**2-parton anomalous dimension should hold also for active area**

Will cone also just be rescaled?

## Areas: also a tool for correcting jets



Areas not just theoretical tool.

Can be **measured jet-by-jet** in real events and used for pileup corrections.

Each jet corrected by  $\text{area} \times \text{median} (P_t/\text{area})$

E.g.: semileptonic  
 $t\bar{t}$  @ LHC with  
 $\langle n_{P.U.} \rangle \simeq 20$ .

Naive analysis: no cuts; assume both b's tagged  
 Take two hardest non-b jets — call them a W  
 Take correct sign b, combine with W → top

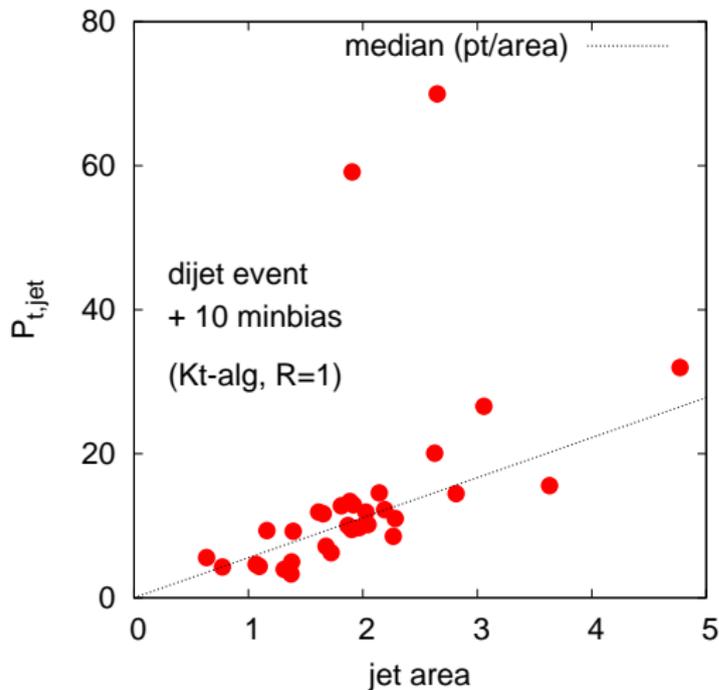
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  - ▶ hadronisation:  $-\Lambda/R$
  - ▶ UE & pileup:  $+\Lambda R^2$
- ▶ Differences that are often noted are mainly due to different  $R$ 's.
- ▶ *Jet areas* are a useful playground for understanding effects beyond 1-NP-gluon level.
  - ▶ Perturbative sub-structure  $\rightarrow$  algorithm-specific anomalous dimensions
  - ▶ Accounting for self-clustering  $\rightarrow$  rescaling of jet area
  - ▶ Full understanding of two together needs further work
- ▶ Jet areas are also a useful concept in jet-by-jet *corrections* of pileup contamination.

Les Houches '07 will have jets subgroup — input welcome!

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  - ▶ UE & pileup:  $+\Lambda R^2$
- ▶ Differences that are often noted are mainly due to different  $R$ 's.
- ▶ *Jet areas* are a useful playground for understanding effects beyond 1-NP-gluon level.
  - ▶ Perturbative sub-structure  $\rightarrow$  algorithm-specific anomalous dimensions
  - ▶ Accounting for self-clustering  $\rightarrow$  rescaling of jet area
  - ▶ Full understanding of two together needs further work
- ▶ Jet areas are also a useful concept in jet-by-jet *corrections* of pileup contamination.

**Les Houches '07 will have jets subgroup — input welcome!**

# EXTRA MATERIAL



Jet areas in  $k_t$  algorithm are quite varied

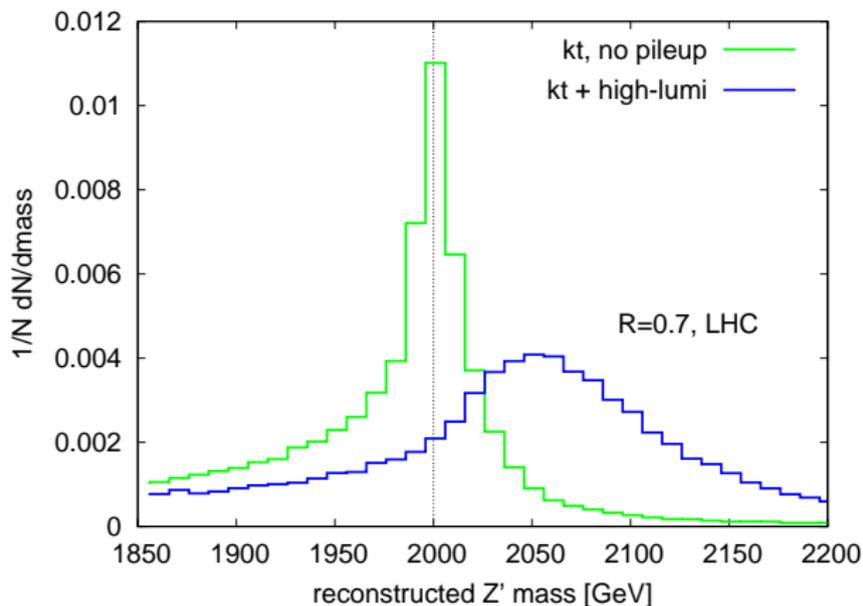
Because  $k_t$ -alg adapts to the jet structure

- ▶ Contamination from min-bias  $\sim$  area

Complicates corrections: min-bias subtraction is different for each jet.

Cone supposedly simpler  
Area =  $\pi R^2$ ? (Not quite...)

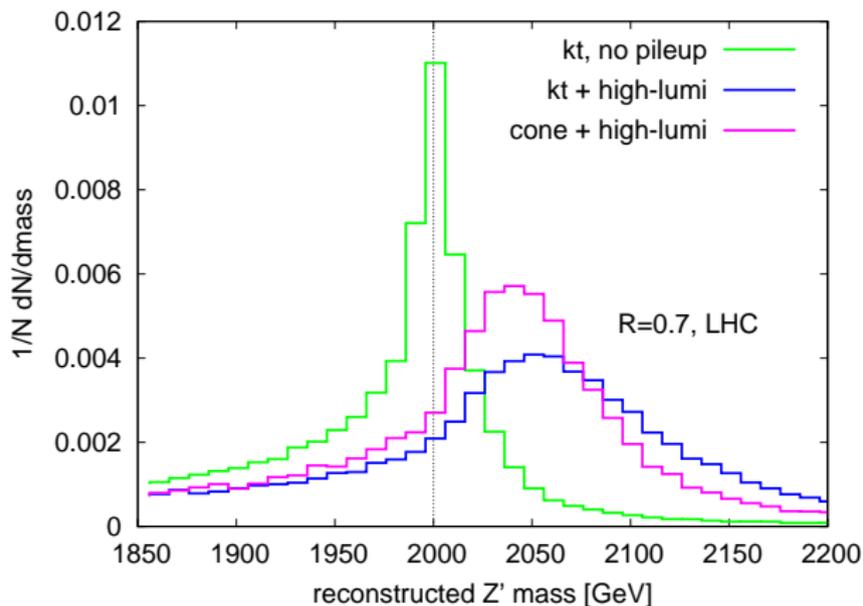
**But:** area can be measured for each jet, as can typical median  $p_t/area$ .



Uncorrected cone better than  $k_t$ .

Cam is intermediate ( $\langle A_{cam} \rangle \simeq \langle A_{cone} \rangle$ , but fluctuations larger)

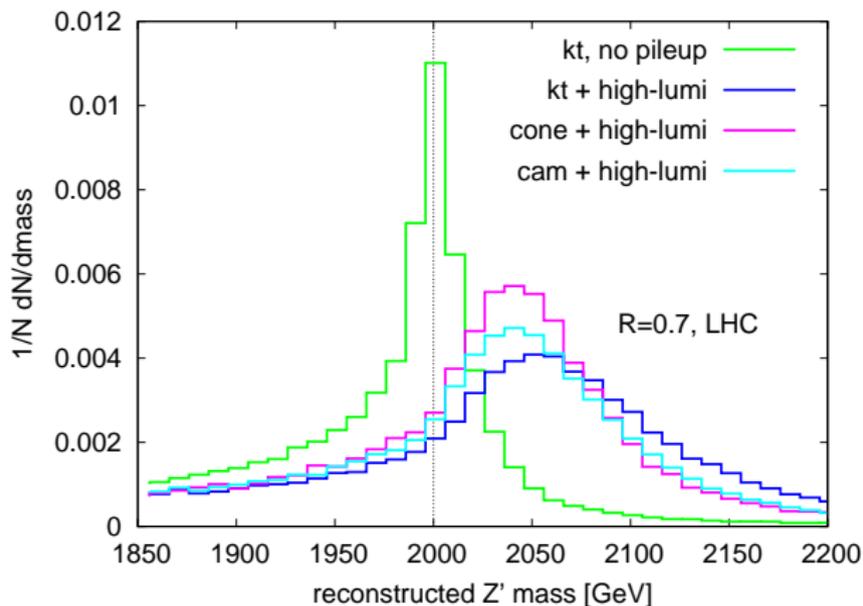
Corrected Cam (and  $k_t$ ) is best.



Uncorrected cone better than  $k_t$ .

Cam is intermediate ( $\langle A_{cam} \rangle \simeq \langle A_{cone} \rangle$ , but fluctuations larger)

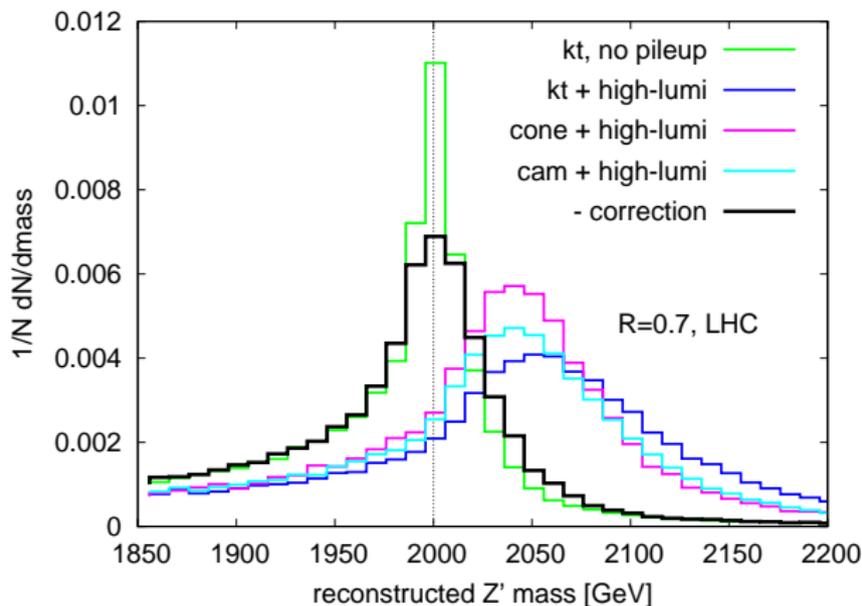
Corrected Cam (and  $k_t$ ) is best.



Uncorrected cone better than  $k_t$ .

Cam is intermediate ( $\langle A_{cam} \rangle \simeq \langle A_{cone} \rangle$ , but fluctuations larger)

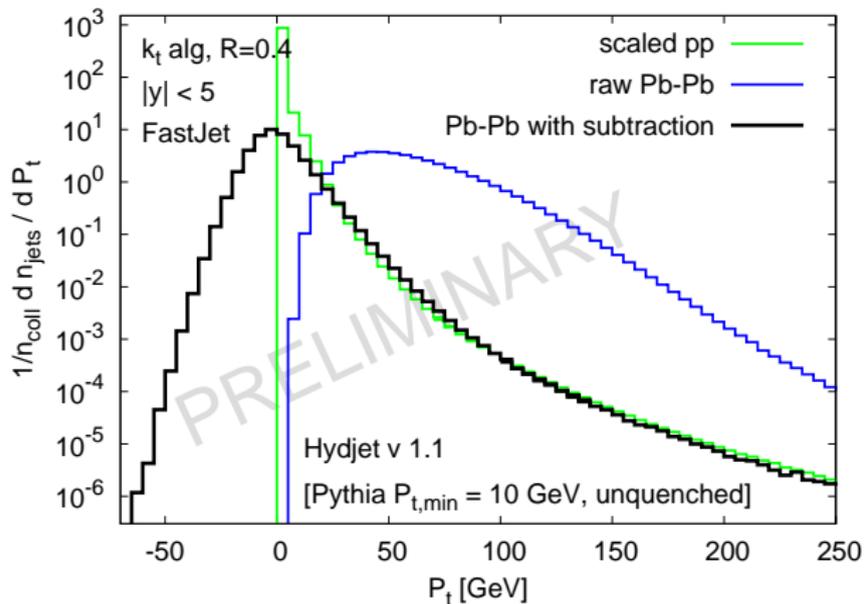
Corrected Cam (and  $k_t$ ) is best.



Uncorrected cone better than  $k_t$ .

Cam is intermediate ( $\langle A_{cam} \rangle \simeq \langle A_{cone} \rangle$ , but fluctuations larger)

Corrected Cam (and  $k_t$ ) is best.



Most HI studies use just particles with  $p_t >$  a few GeV  
IR unsafe  
affected by quenching

We use *all* particles and area-based subtraction.

Good results despite the huge subtraction being performed.