LONDON United Kingdom



https://cern.ch/fccweek2023

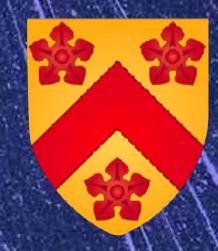


PHYSICS PERSPECTIVES

Gavin Salam







A preamble

- ➤ this type of talk is often given by a theorist who builds models of new physics
- > such a theorist can tell you with authority about the landscape of models that any given collider might probe

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- ➤ this type of talk is often given by a theorist who builds models of new physics
- ➤ such a theorist can tell you with authority about the landscape of models that any given collider might probe

- > there are many kinds of theorist
- > while I'm a theorist, I am not a BSM model-builder
- ➤ my "day job" is to calculate phenomena in QCD (jets, parton showers, etc.), in order to help augment colliders' capabilities
- ➤ this talk will not involve specifics of models, but rather attempt to explore the case for new colliders more generically

desirable features of a worldwide HEP project?

an important target that is guaranteed to be reached (no-lose theorem)

exploration into the unknown by a significant factor in energy

major progress on a broad array of particle physics topics

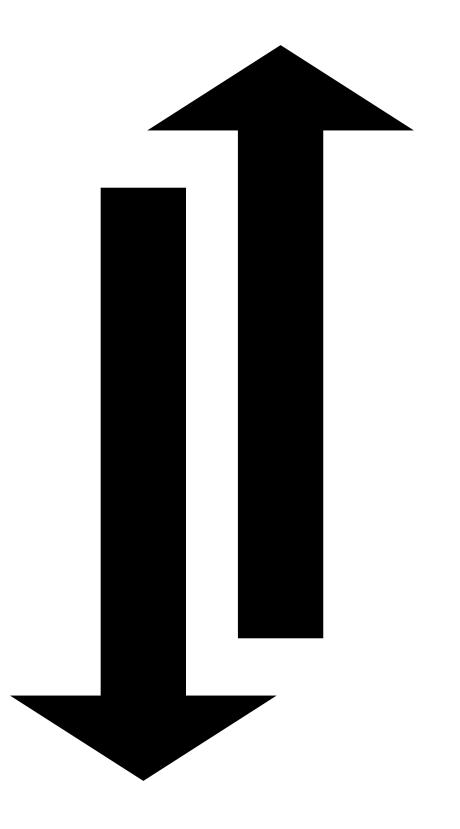
likelihood of success, robustness (multiple experiments)

cost-effective construction & operation, low carbon footprint

top-down

figure out the best collider you can realistically build

establish what physics it will probe



bottom up

establish what you want to learn

figure out how to build a collider that will best achieve it



Dear Santa Claus,

We have been good these past decades.
Please could you now bring us

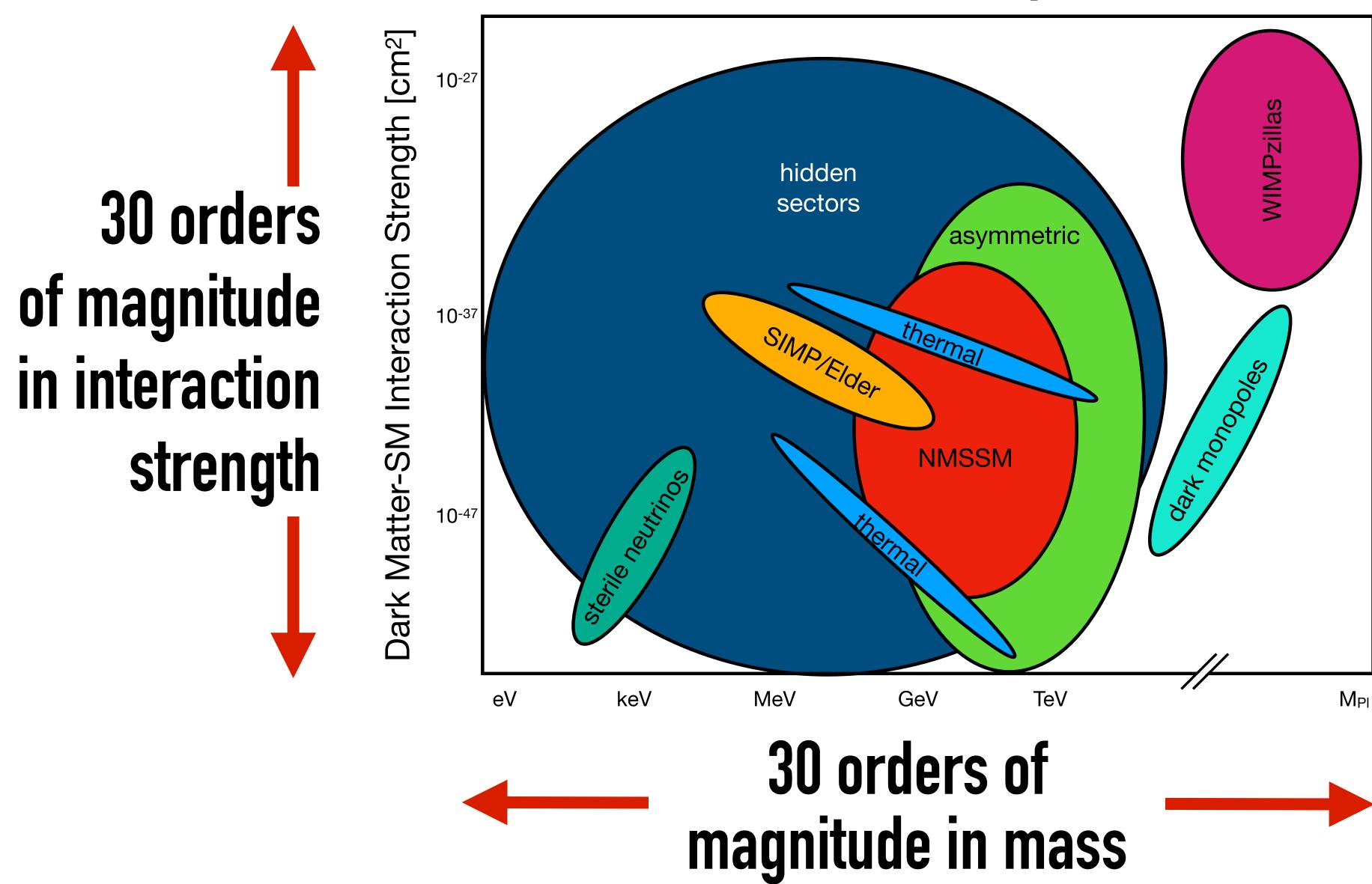
- a dark matter candidate
- an explanation for the fermion masses
- an explanation of matter-antimatter asymmetry
- an axion, to solve the strong CP problem
- a solution to fine tuning the EW scale
- a solution to fine tuning the cosmological constant

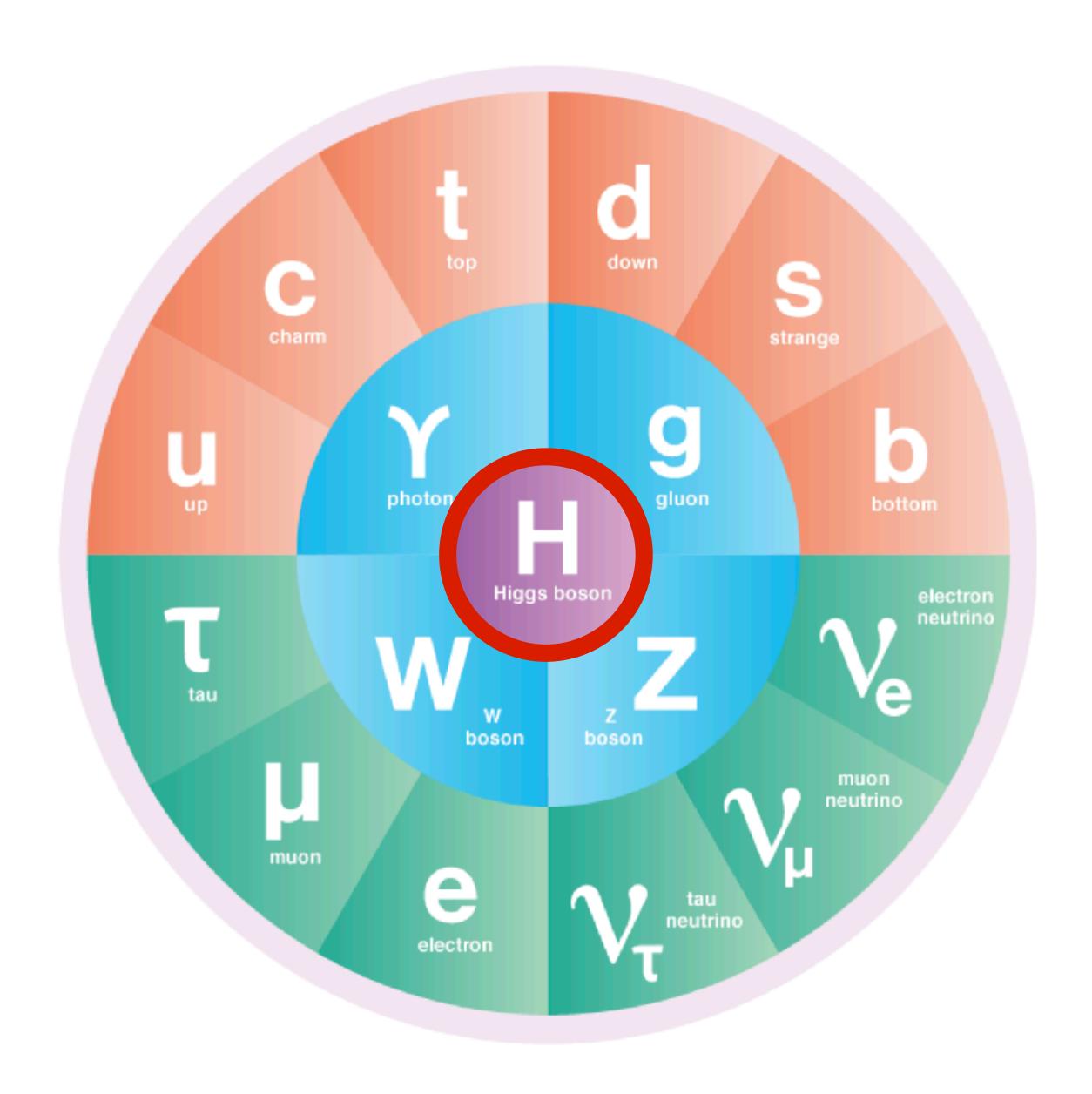
Thank you, Particle Physicists

ps: please, no anthropics

we have so far been unlucky in getting answers to these many questions

Snowmass Dark Matter report, <u>2209.07426</u>





but we have been lucky in discovering a 125 GeV Higgs boson

it opens a door to the most mysterious part of the Standard Model

desirable features of a worldwide HEP project?

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Higgs physics

Higgs is the last particle of the SM. So the SM is complete, right?

The Lagrangian and Higgs interactions: two out of three qualitatively new!

$$\mathcal{L}_{SM} = \cdots + |D_{\mu}\phi|^2 + \psi_i y_{ij} \psi_j \phi - V(\phi)$$

Gauge interactions, structurally like those in QED, QCD, EW, studied for many decades (but now with a scalar)

Yukawa interactions. Responsible for fermion masses, and induces "fifth force" between fermions. Direct study started only in 2018!

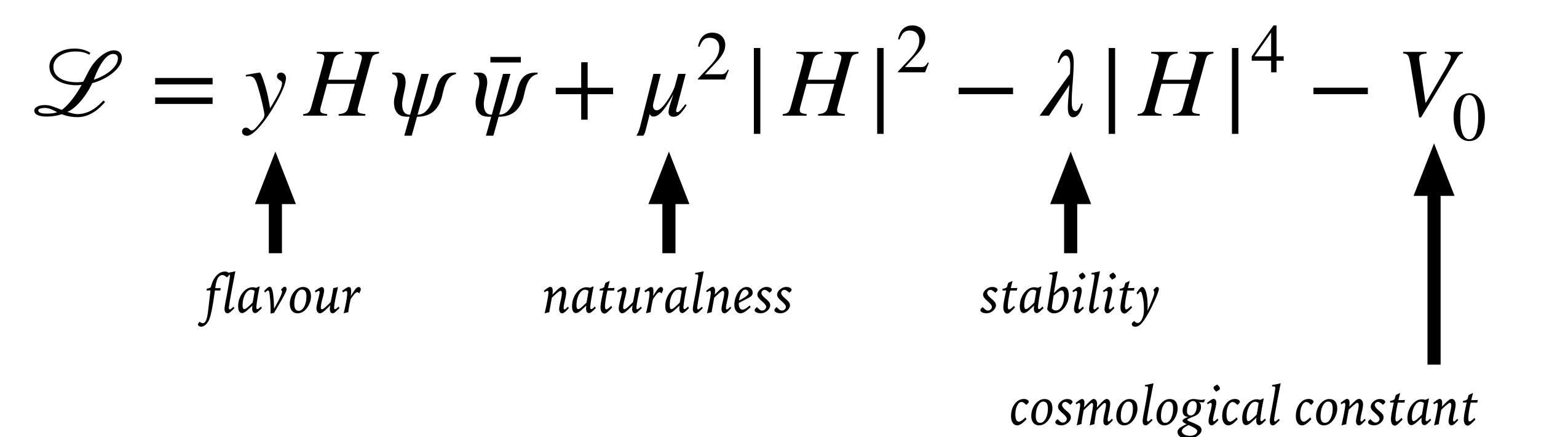


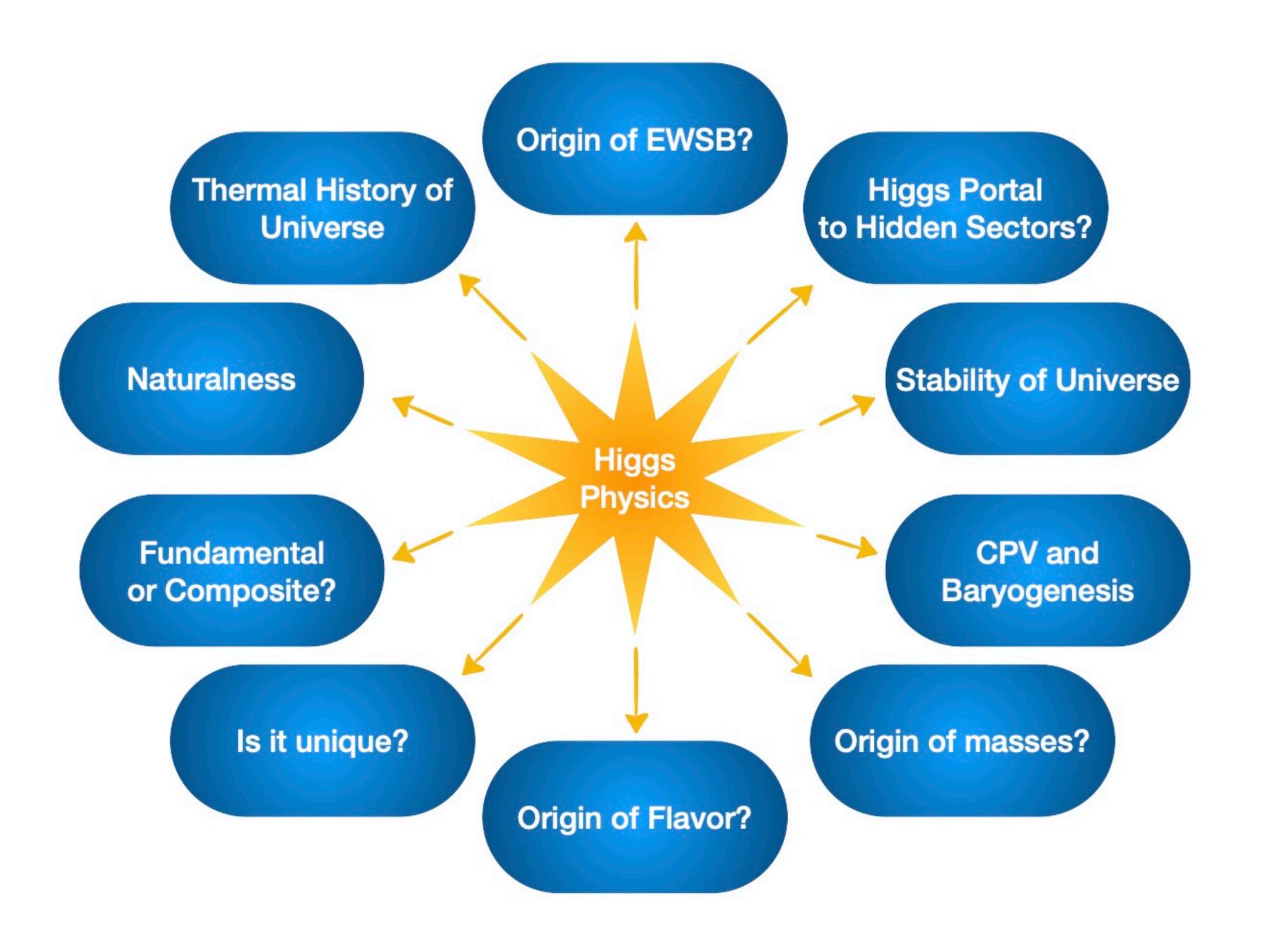
Higgs potential → self-interaction ("sixth?" force between scalars).

Holds the SM together.

Unobserved

Almost every problem of the Standard Model originates from Higgs interactions





Yukawa interaction hypothesis

Yukawa couplings ~ fermion mass

first fundamental interaction that we probe at the quantum level where interaction strength (y_{ij}) not quantised

(i.e. no underlying unit of conserved charge across particles)

2.2 MeV 2.2 MeV 4.7 MeV

proton:



+electromagnetic & strong forces

~ 938.3 MeV

neutron:



≃ 939.6 MeV

2.2 MeV 4.7 MeV 4.7 MeV

Protons are **lighter** than neutrons→ protons are stable. Giving us the hydrogen atom, & chemistry and biology as we know it

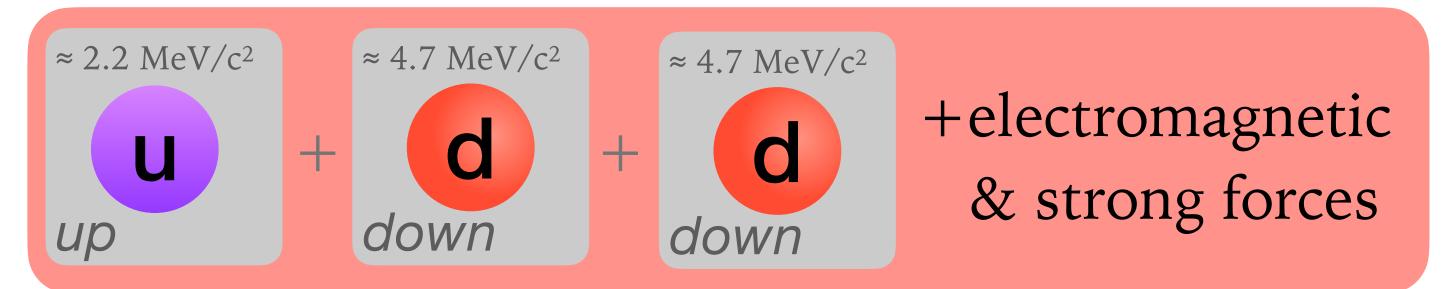
2.2 MeV 2.2 MeV 4.7 MeV

proton:



 $\simeq 938.3 \text{ MeV}$

neutron:



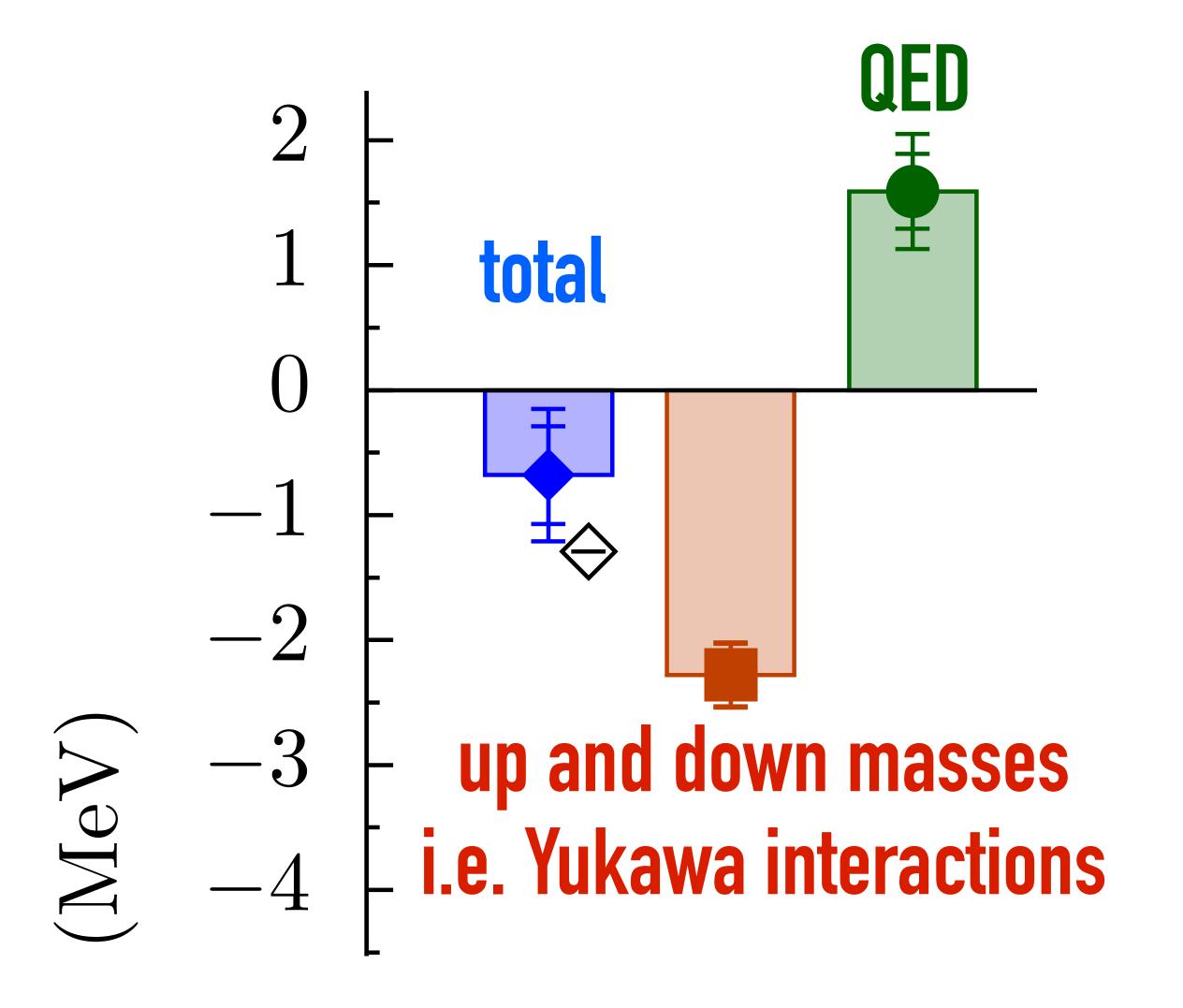
~ 939.6 MeV

2.2 MeV 4.7 MeV 4.7 MeV

Protons are **lighter** than neutrons→ protons are stable. Giving us the hydrogen atom, & chemistry and biology as we know it

Supposedly because up quarks interact more weakly with the Higgs field than down quarks

proton – neutron mass difference



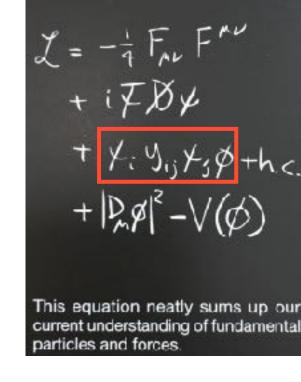
Lattice calculation
(BMW collab.)

1306.2287

1406.4088

Why do Yukawa couplings matter?

(2) Because, within SM conjecture, they're what give masses to all leptons



Bohr radius

$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{m_e e^2} = \frac{\hbar}{m_e c\alpha} \propto \frac{1}{y_e}$$

electron mass determines size of all atoms

it sets energy levels of all chemical reactions

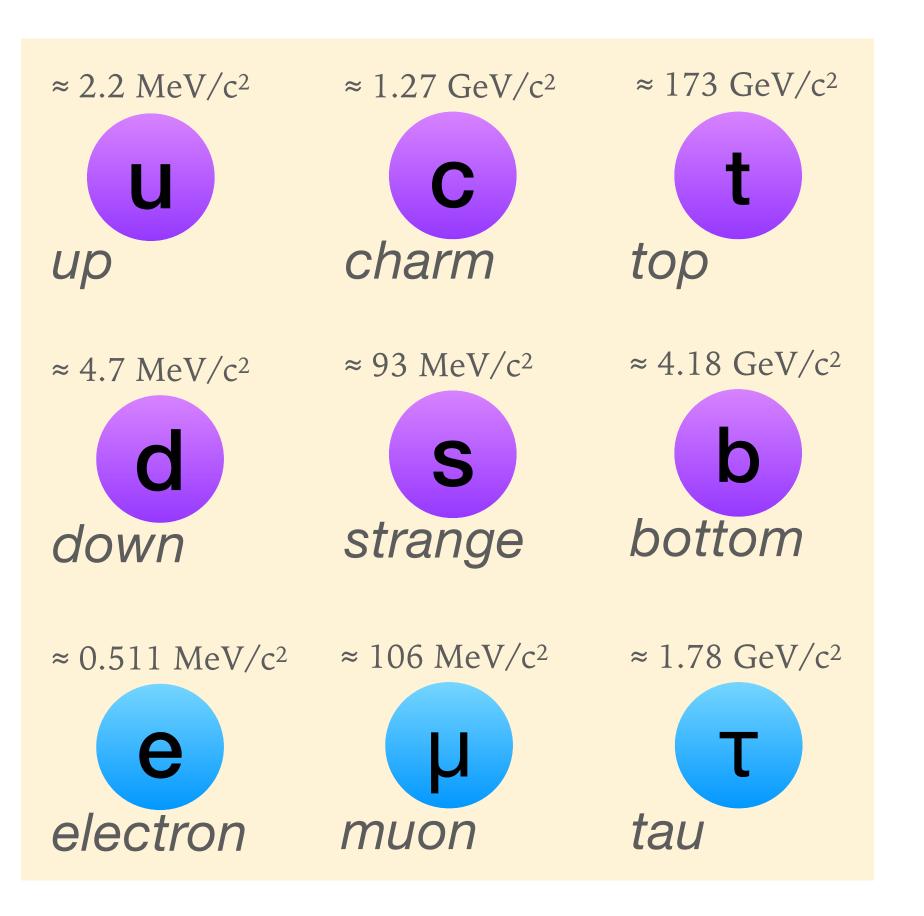
currently we have no evidence that up and down quarks and electron get their masses from Yukawa interactions—it's in textbooks, but is it nature?

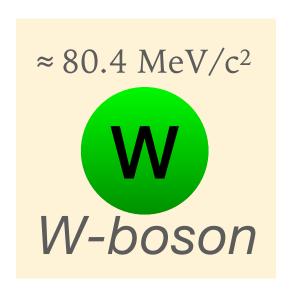
First

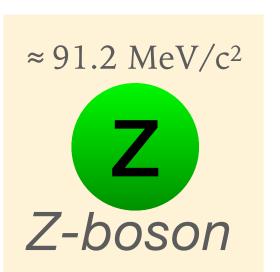
Second generation generation

Third

H interactions





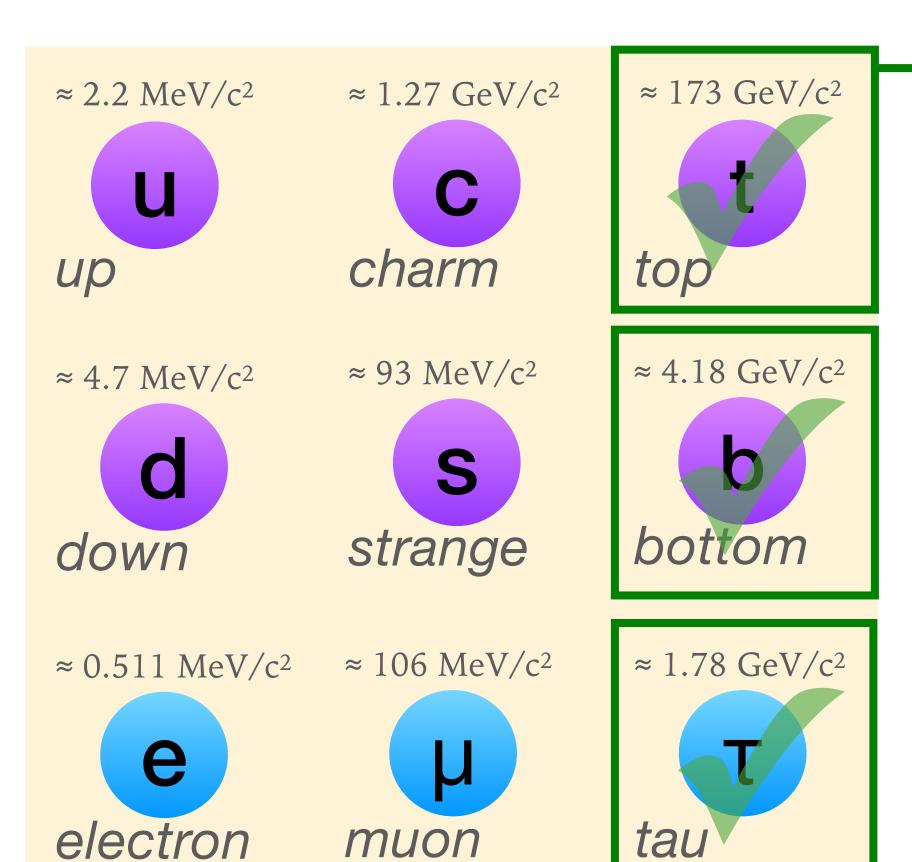


First generation

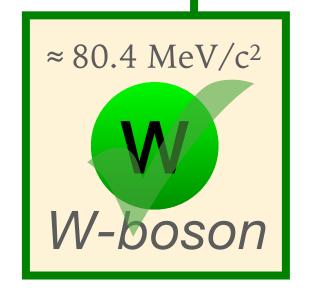
Second generation generation

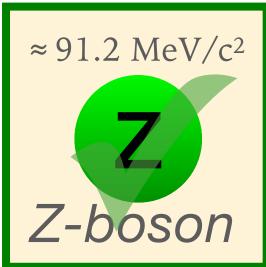
Third

Hinteractions



established (5σ) at LHC by observation of direct interaction with H

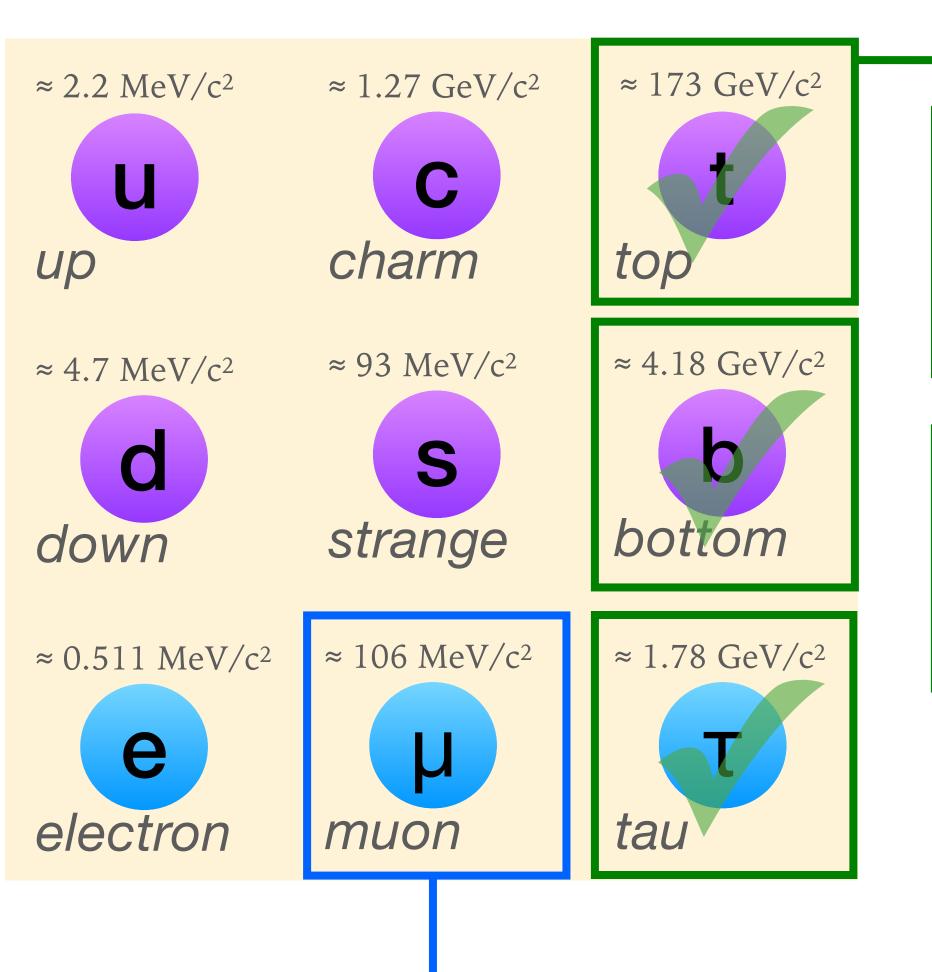




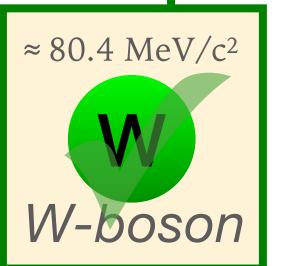
First generation Second

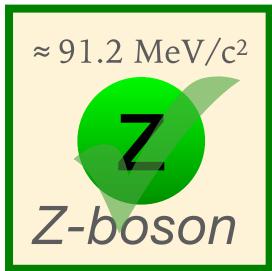
Third generation generation

H interactions



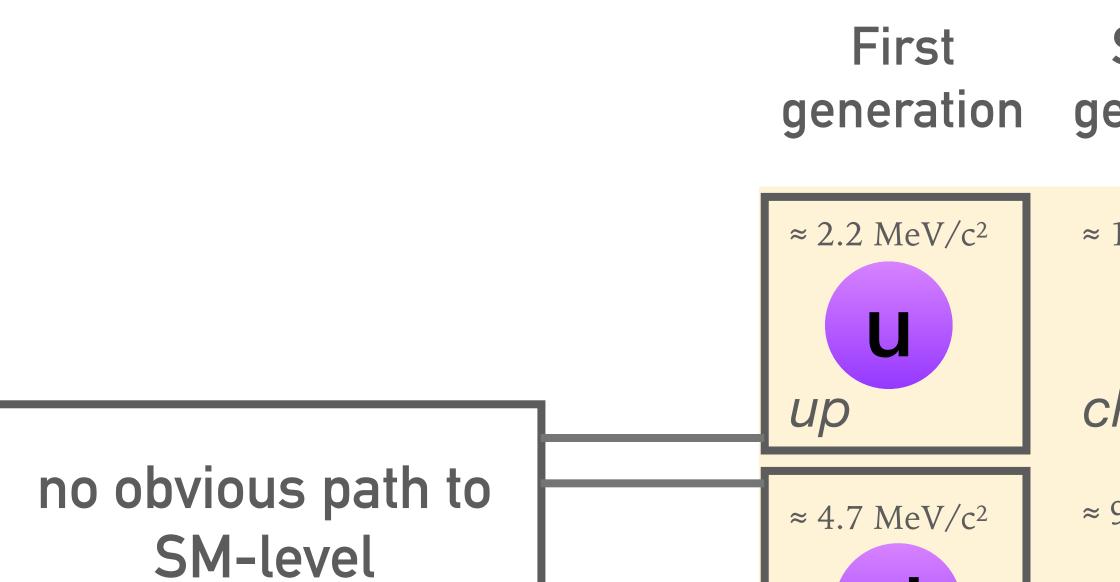
established (5σ) at LHC by observation of direct interaction with H





first evidence (3σ)

to be conclusively established at the LHC within 3 - 10 years



measurement

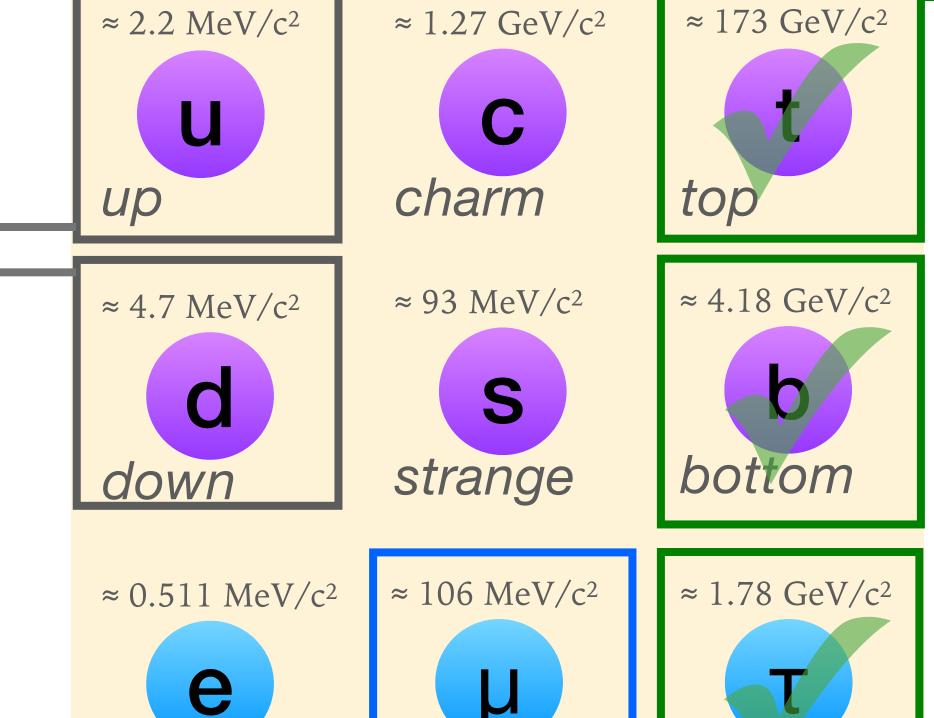
bright ideas

needed!



tau

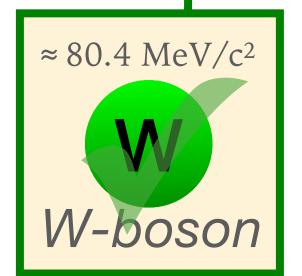
H interactions

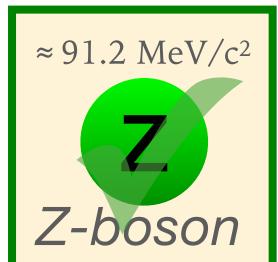


muon

electron

established (5σ) at LHC by observation of direct interaction with H





first evidence (3σ)

to be conclusively established at the LHC within 3 – 10 years

no evidence yet

guaranteed at FCC-ee

no obvious path to SM-level measurement bright ideas needed!

First generation

 $\approx 2.2 \text{ MeV/c}^2$

 $\approx 4.7 \text{ MeV/c}^2$

d

 $\approx 0.511 \text{ MeV/c}^2$

e

electron

down

up

Second generation generation

 $\approx 1.27 \text{ GeV/c}^2$

C

charm

 $\approx 93 \text{ MeV/c}^2$

S

strange

≈ 106 MeV/c²

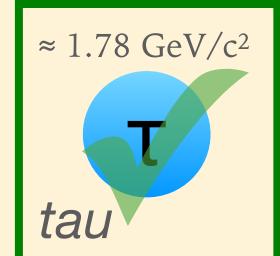
muon

Third

≈ 173 GeV/c² top

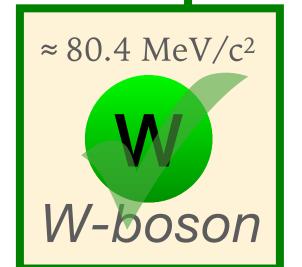


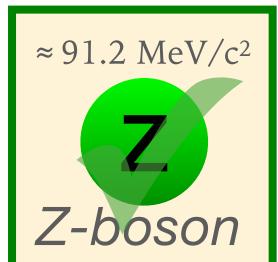




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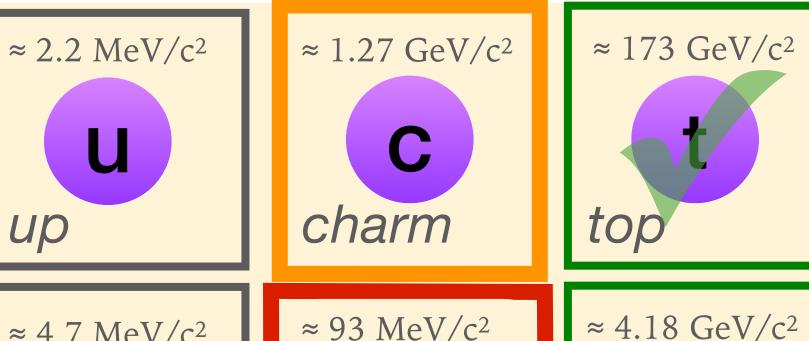
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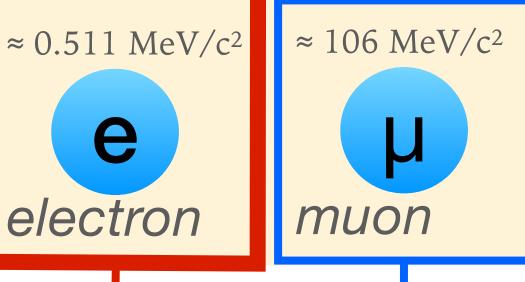
tantalisingly close to reach of FCC-ee

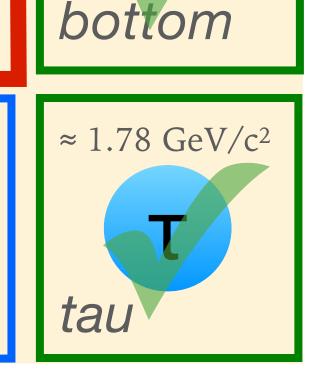
Second Third First generation generation generation





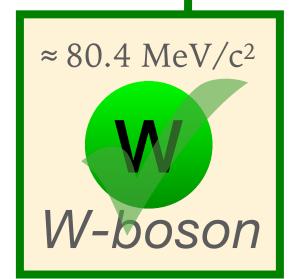
up

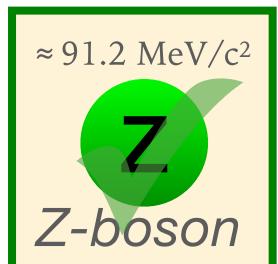




H interactions

established (5σ) at LHC by observation of direct interaction with H





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Teaser from the analysis front [FCC-ee, H → hadrons]

Tools fully incorporated in FCCSW [details]

• Example: $Z(\rightarrow vv)H(\rightarrow qq)$

Signal extraction: 2D fit

Categorize events: bb, cc, ss, gg Sub-categories w/ different S/B

ParticleNet-ee

Z - 0.32 0.17 0.25 0.2 0.06 0 0 0 0 0 0

HWW - 0.19 0.049 0.43 0.11 0.23 0 0 0 0 0 0 0

HZZ - 0.2 0.25 0.2 0.19 0.16 0 0 0 0 0 0

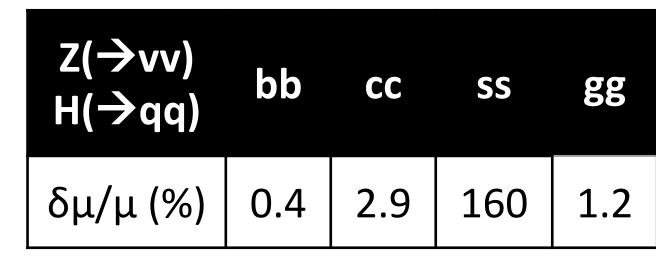
033 0.65 0.15 0.032 0.14 0 0 0 0 0 0 0

q b c s g tau ZZ WW Z HWWHZZqqH

Predicted

Rest (syst: S





m(rec)

FCCAnalyses: FCC-ee Simulation (Delphes)

0 20 40 60 80 100 120 140 160 180 200 220 240

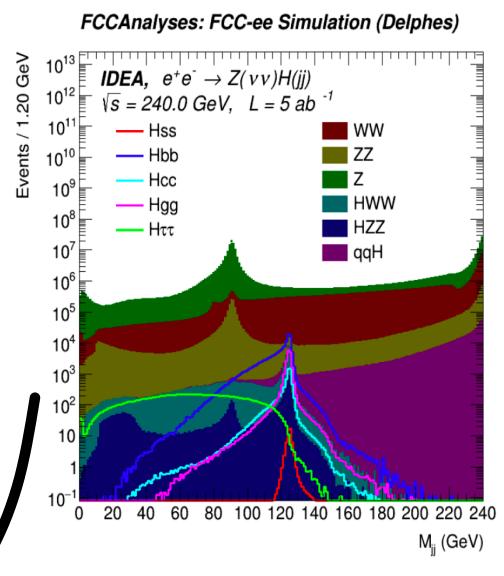
Hpp

M_{rec} (GeV)

 $\sqrt{s} = 240.0 \text{ GeV}, L = 5 \text{ ab}$

* $|\kappa_{\rm S}|$ < 1.9





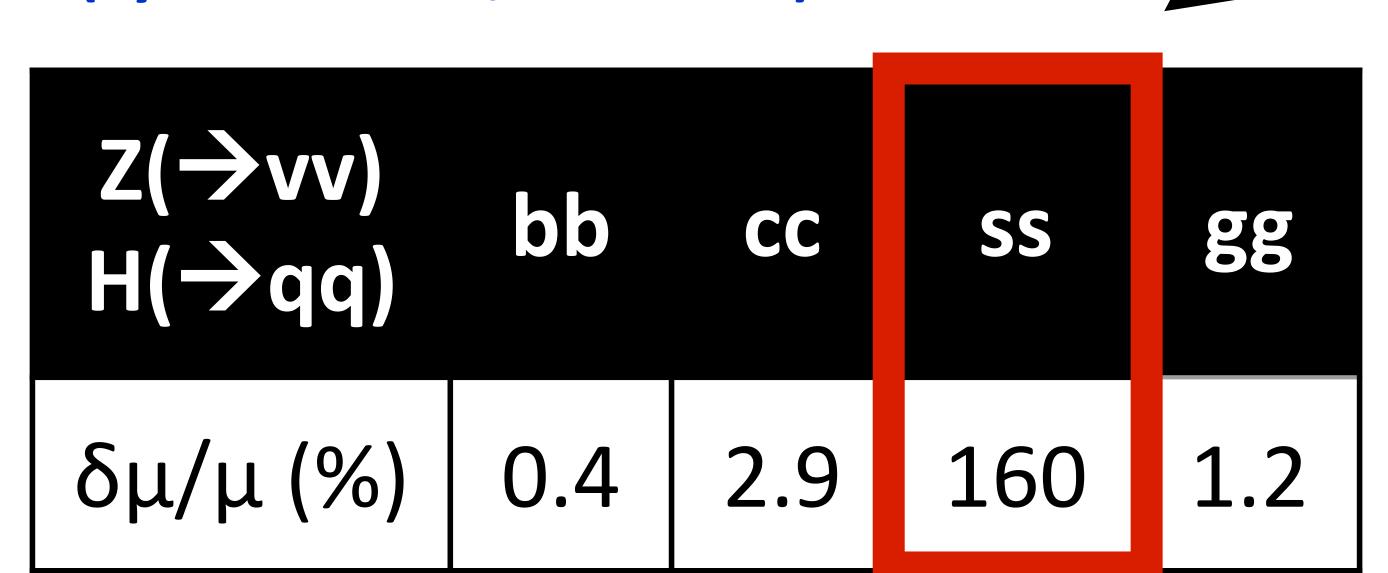
More on Friday:

G. Marchiori

Gavin Salam

Results @ 5ab⁻¹

(syst: 5% BKG, 0.1% SIG)

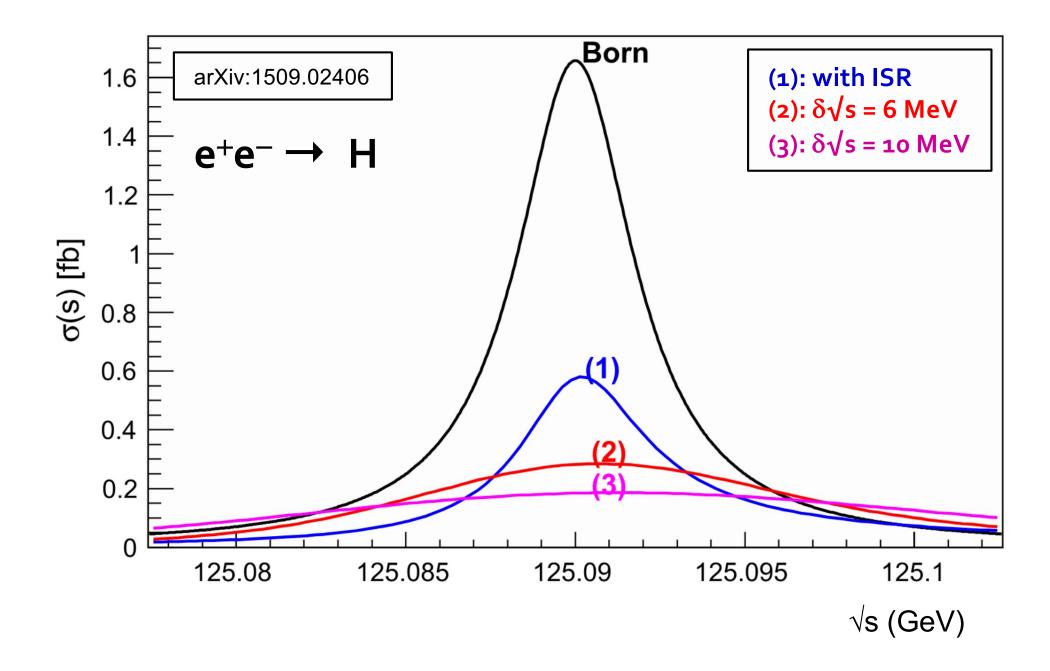


 $*|\kappa_{S}|<1.9$

strange Yukawa tantalisingly close to being within reach would complete 2nd generation Yukawas

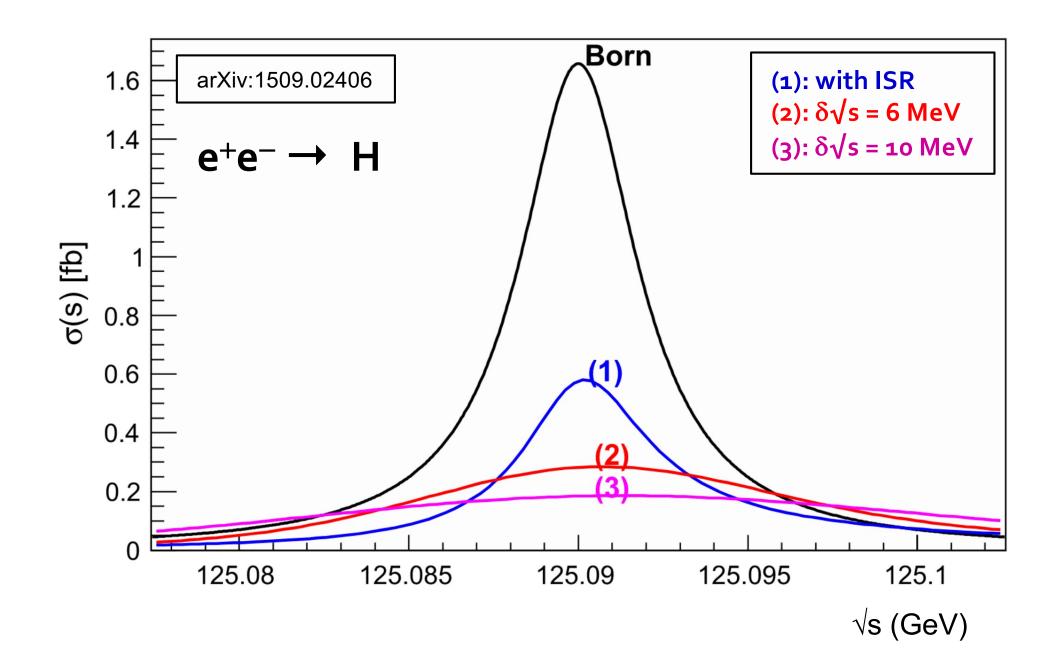
(not yet in the baseline)

- One of the toughest challenges, which requires in particular, at $\sqrt{s} = 125$ GeV
 - \bullet Higgs boson mass prior knowledge to a couple MeV, requires at least the design lumi at \sqrt{s} = 240 GeV
 - Huge luminosity, achievable with with several years of running and possibly 4 IPs
 - ♦ $\sqrt{\text{s}}$ monochromatisation : Γ_{H} (4.2 MeV) \ll natural beam energy spread (~100 MeV)



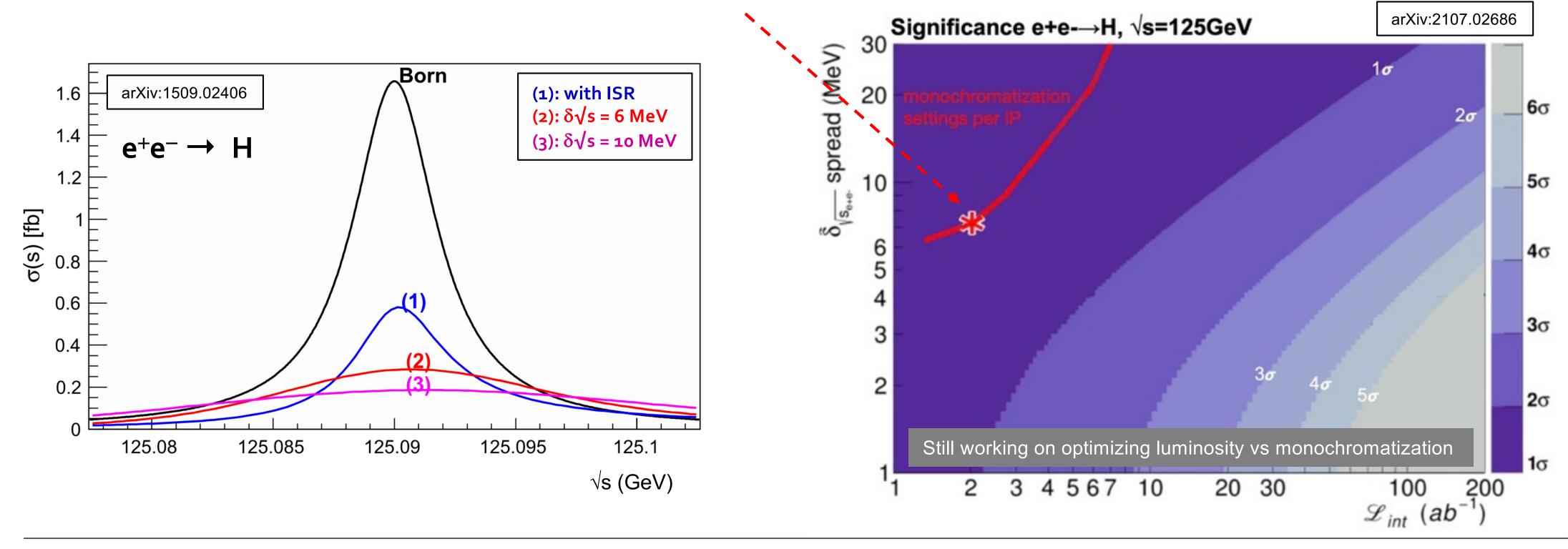
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- First studies indicate a significance of 0.4 σ with one detector in one year



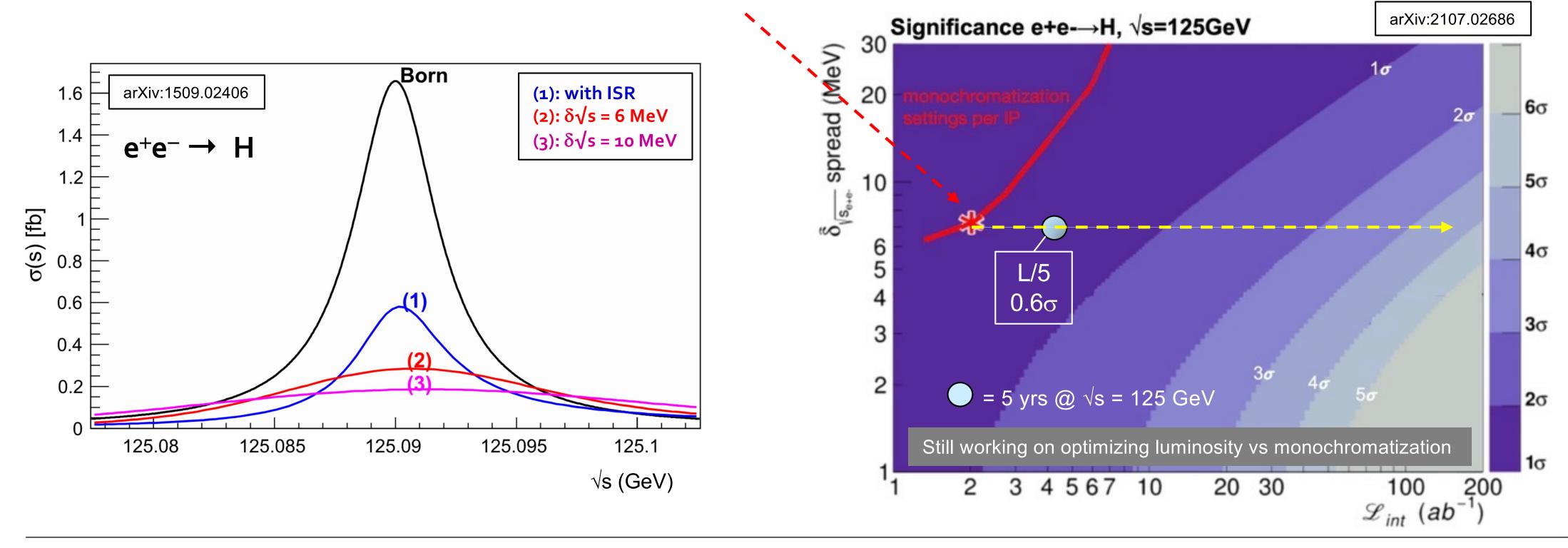
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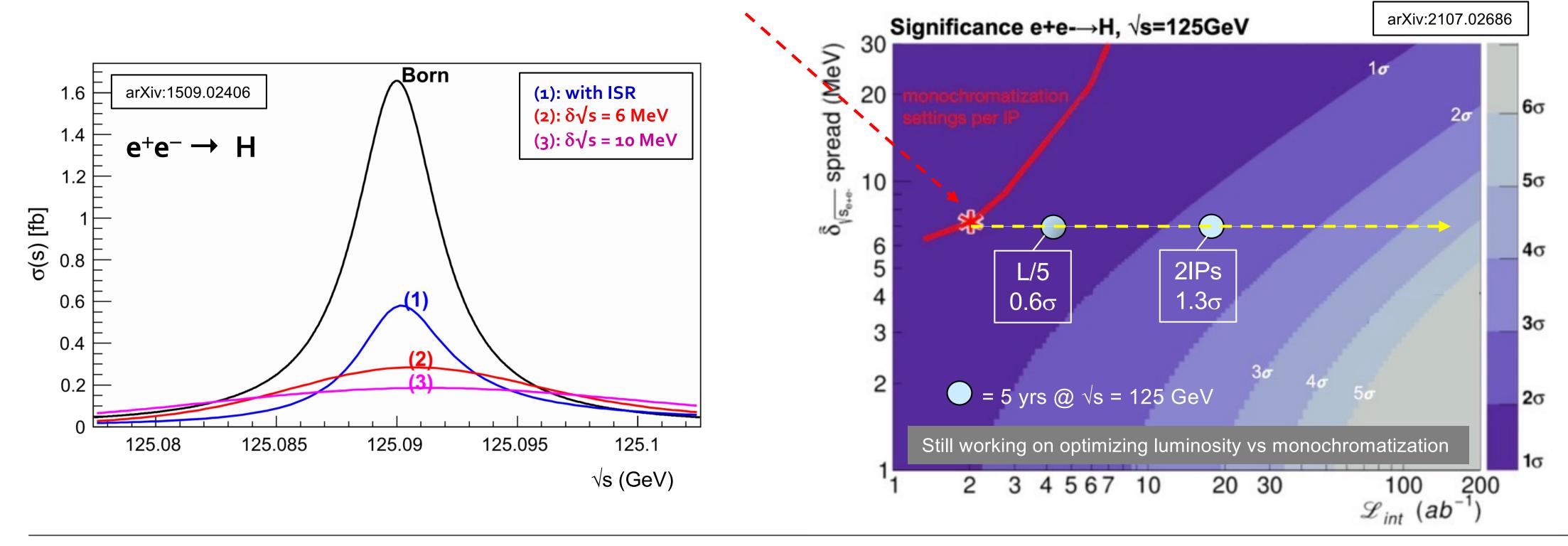
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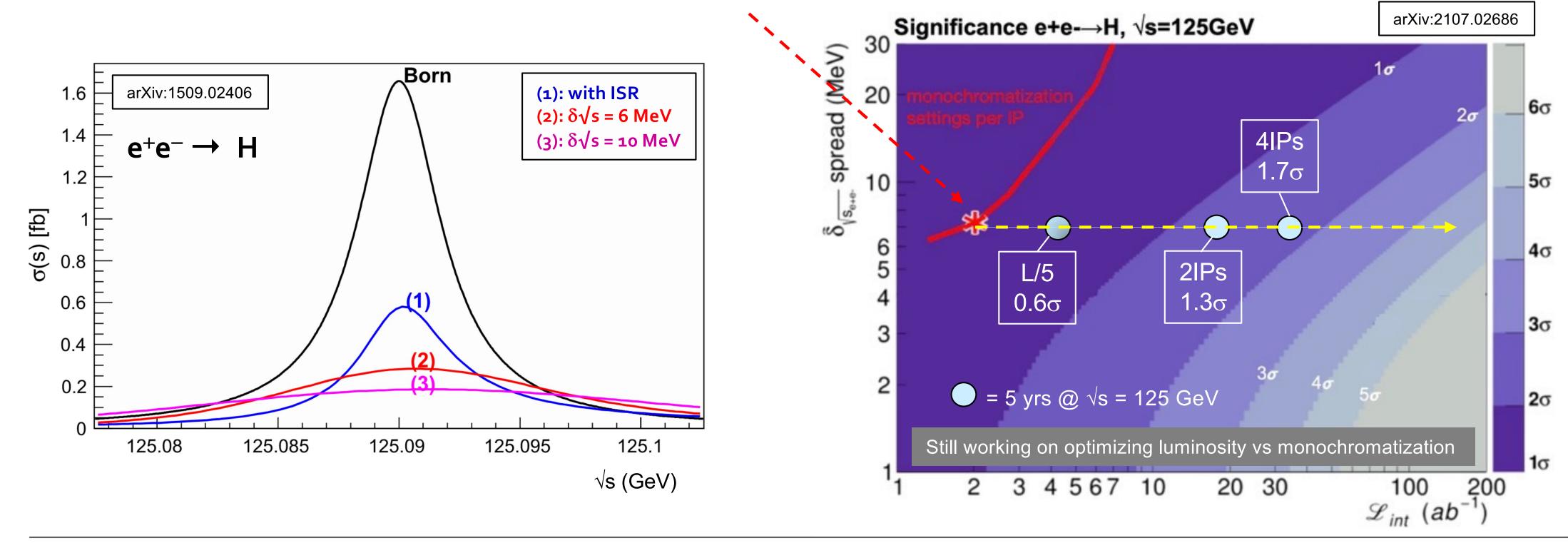
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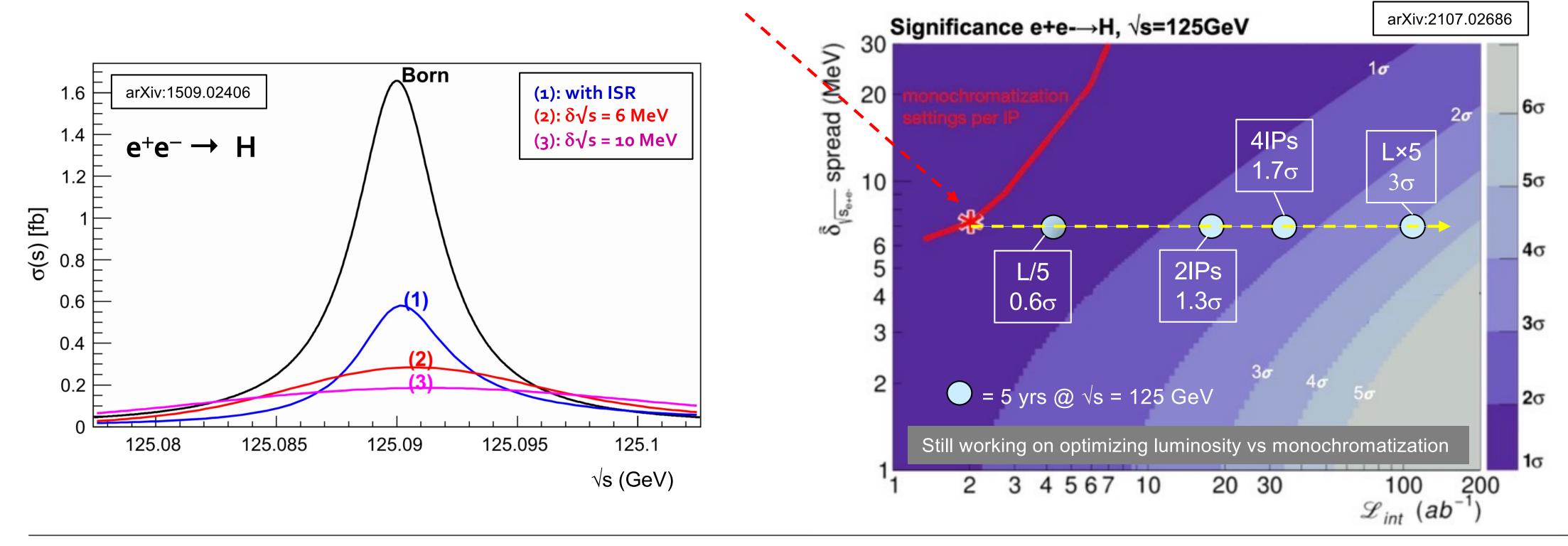
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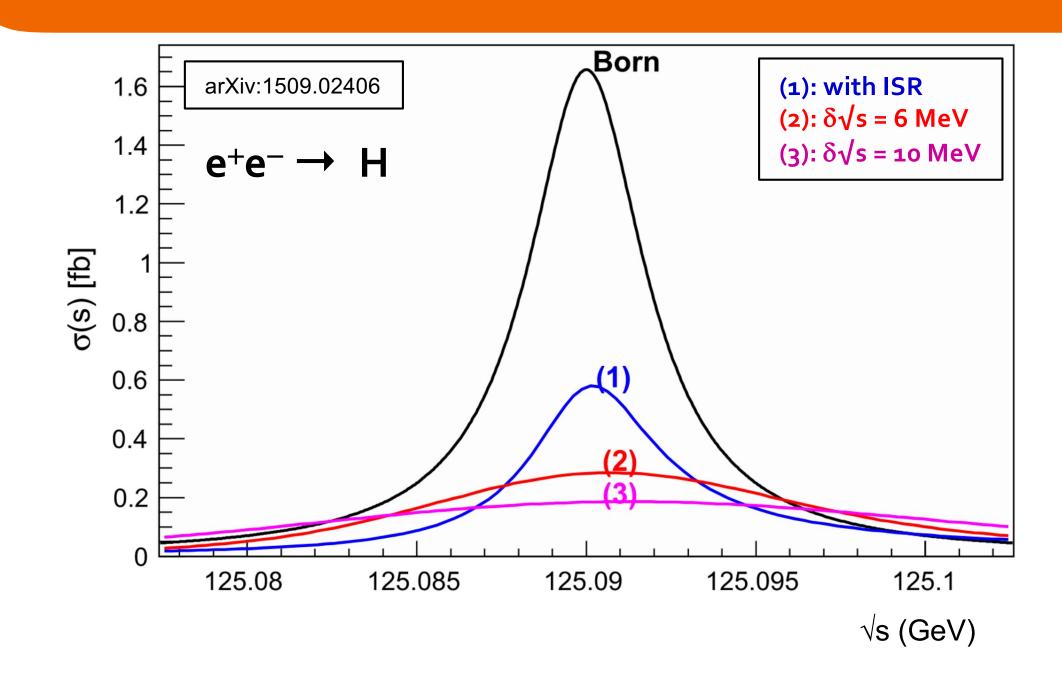
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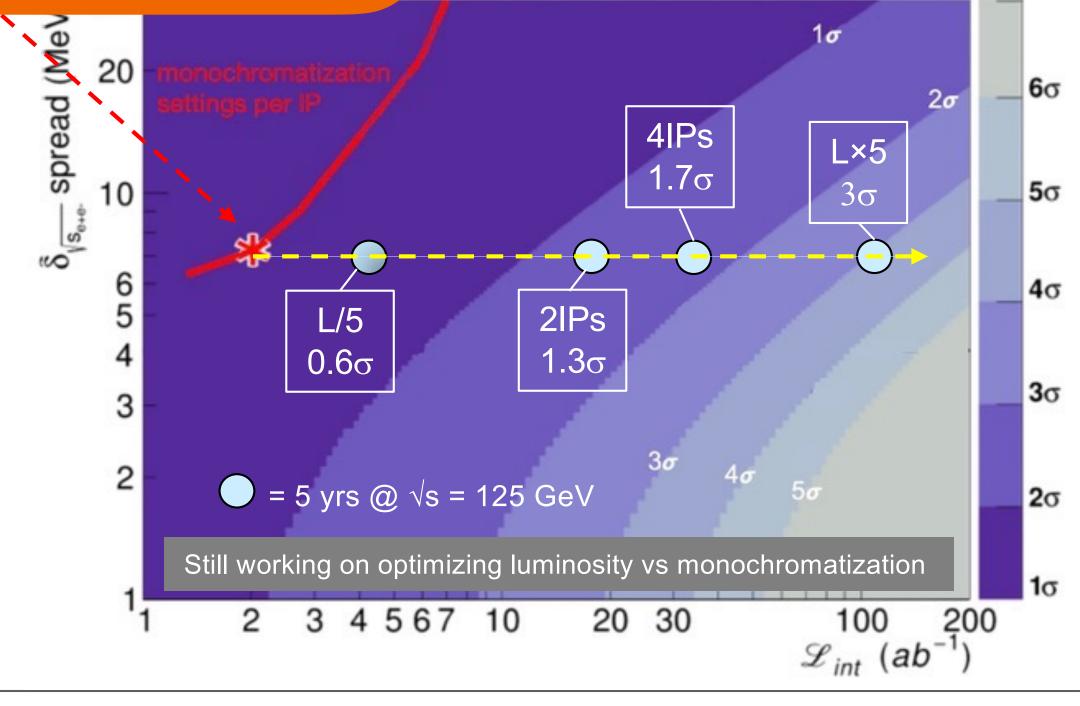
some caution needed with the numbers (cf. Soyez @ 2022 FCC Physics Week on state-of-the art tagging of H→gg)

possibly 4 IPs d (~100 MeV) ne year

→H, √s=125GeV

arXiv:2107.02686





(not yet in the baseline)

6σ

5σ

4σ

- One of the toughest challenges, which requires in particular, at $\sqrt{s} = 125$ GeV
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d (~100 MeV)

ne year

→H, √s=125GeV

4IPs

 1.7σ

arXiv:2107.02686

L×5

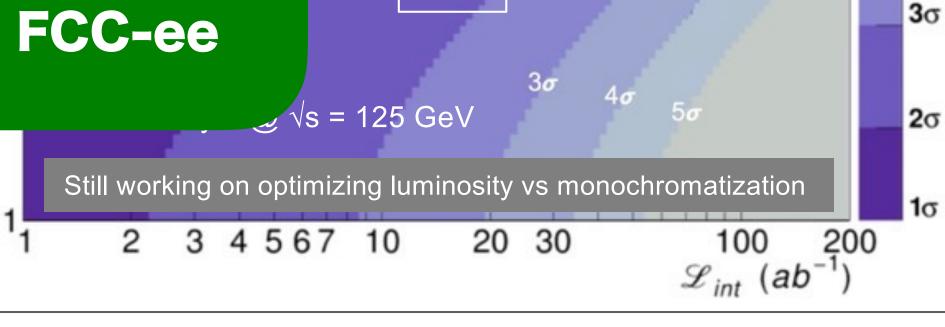
3σ

still a couple of bright ideas away from concrete path to 5σ discovery of the origin of the electron mass; may simply not be feasible

— but would be a clear no-lose theorem for FCC-ee

0 125.08 125.085 125.09 125.095 125.1 √s (GeV)

Born

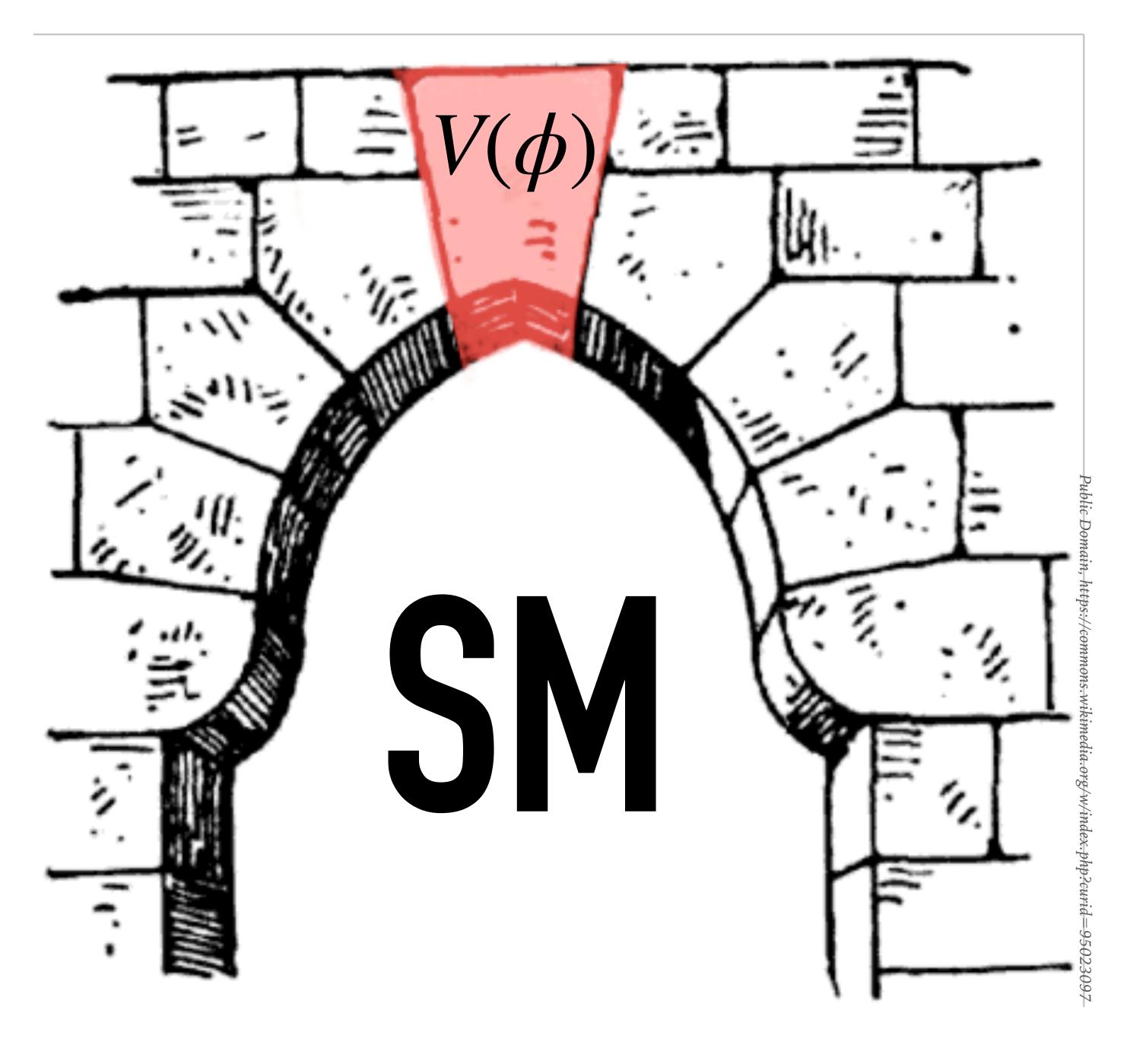


2IPs

 1.3σ

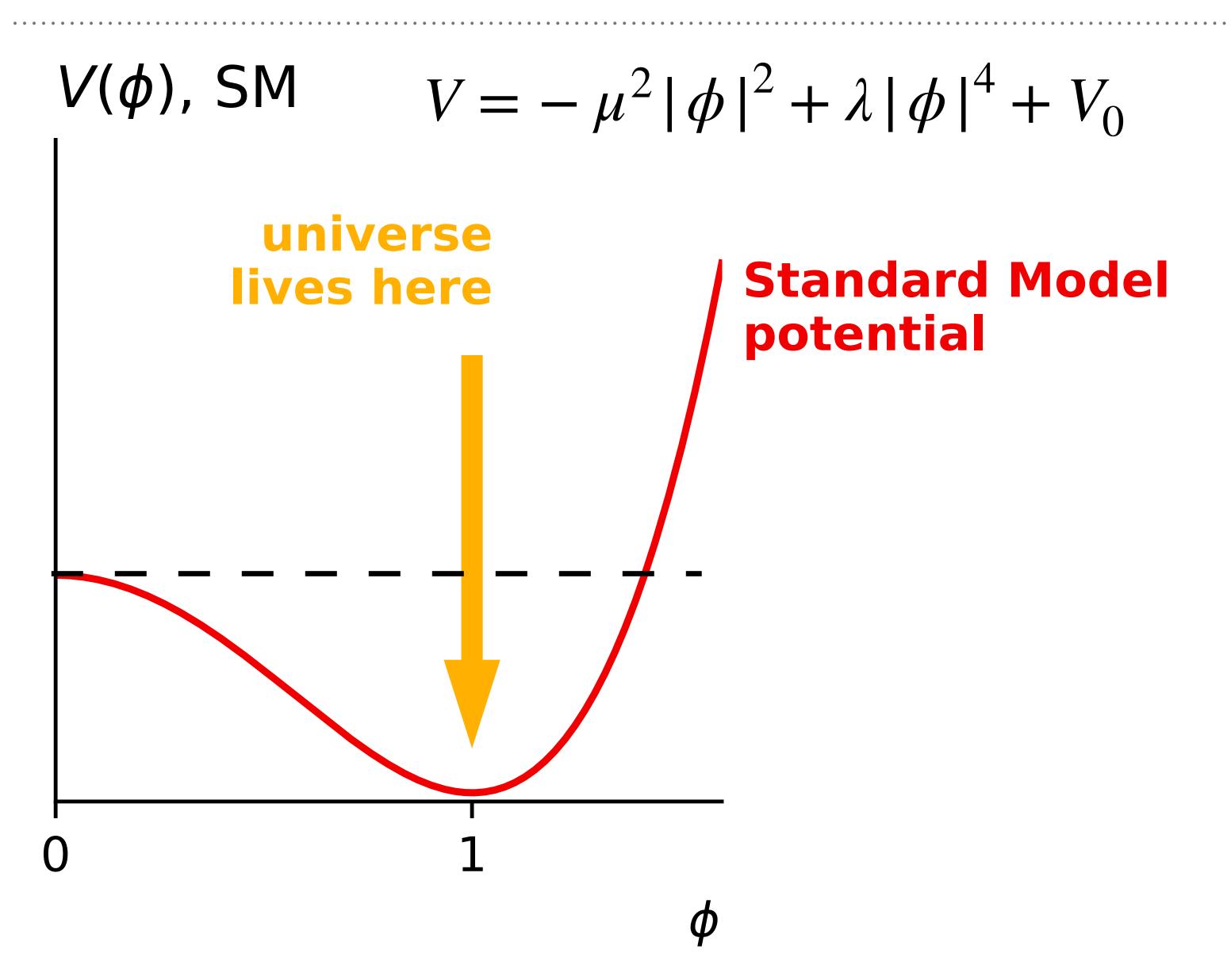
A side comment on the near future at LHC

- ➤ particle physics normally deals with esoteric particles that have [almost] no relation with the world as we experience it
- ► LHC will reach 5σ sensitivity for H → μμ in the coming years (if it is SM-like), offering first proof that particles other than 3rd generation also get their mass from Yukawa mechanism
- ➤ that will be a crucial step on the way from 3rd generation Yukawas to 1st
- > it deserves a big event with the world's press to announce it
- ➤ an opportunity to explain the quest for understanding the origin of the mass of the fundamental particles that <u>we are made of</u>



the Higgs potential

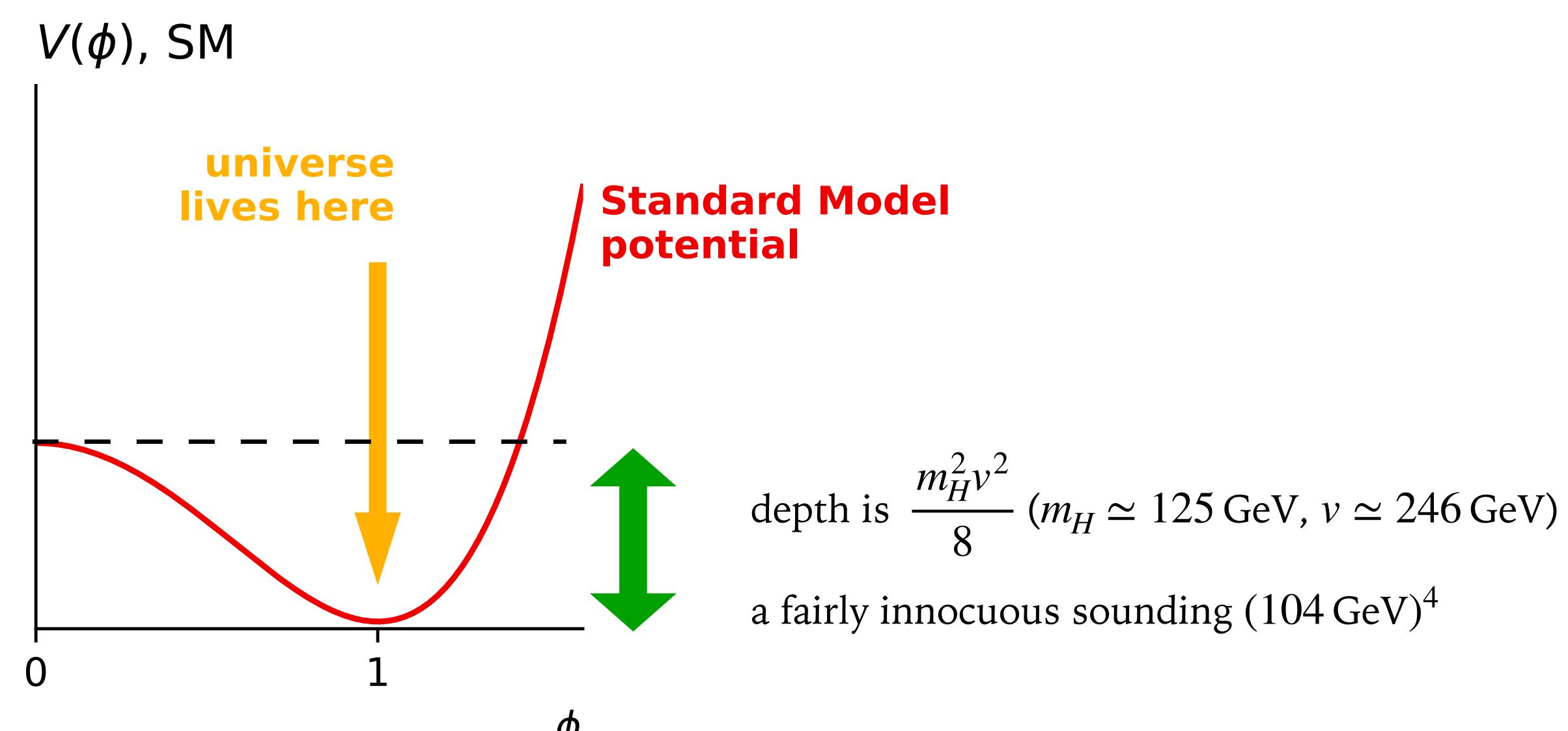
Higgs potential



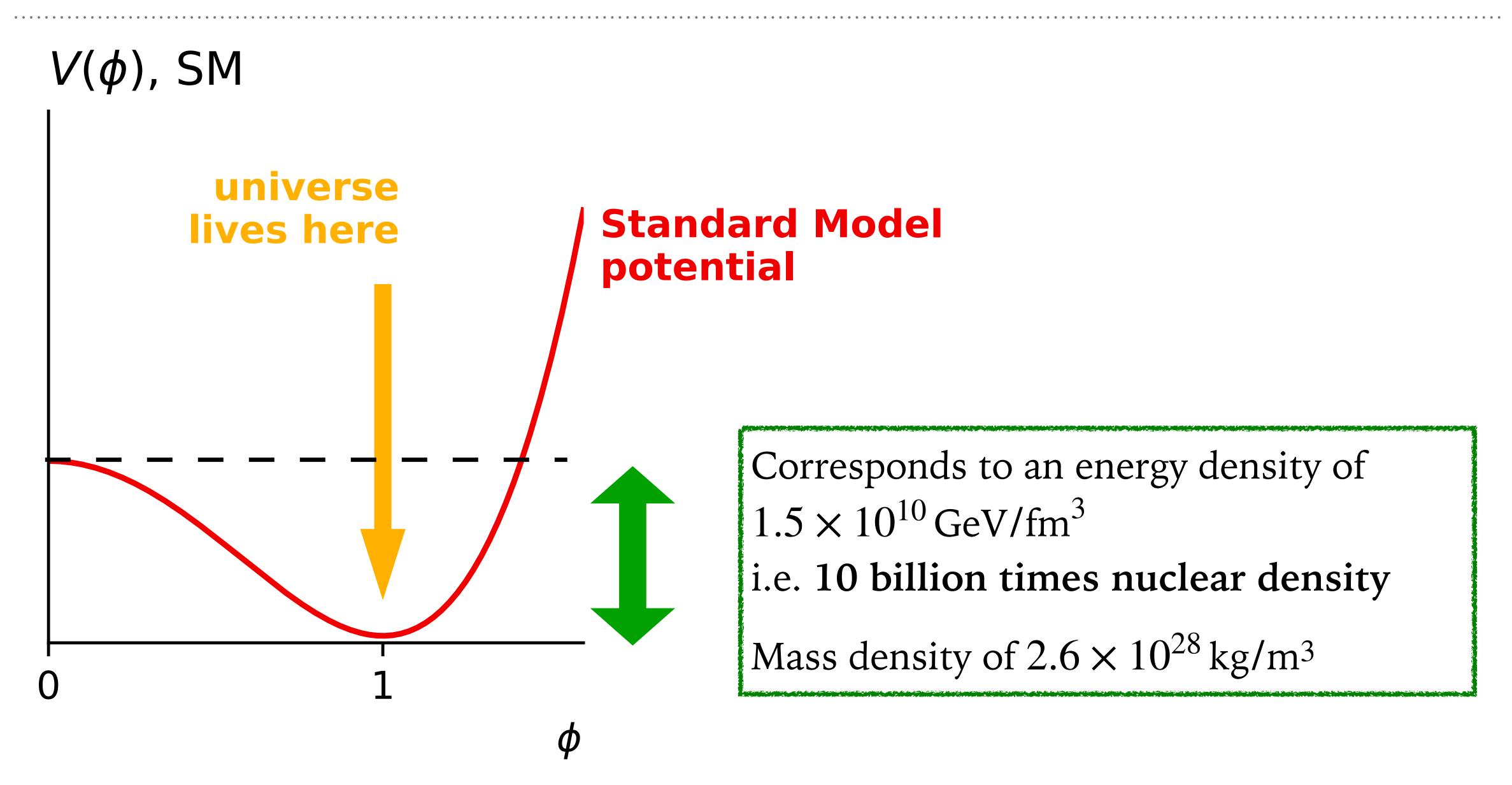
the Higgs
mechanism gives
mass to particles
because the Higgs
field φ is non-zero

That happens because the minimum of the SM potential is at non-zero φ

Higgs potential



Higgs potential — remember: it's an energy density



What does 2.6×10²⁸ kg/m³ mean?



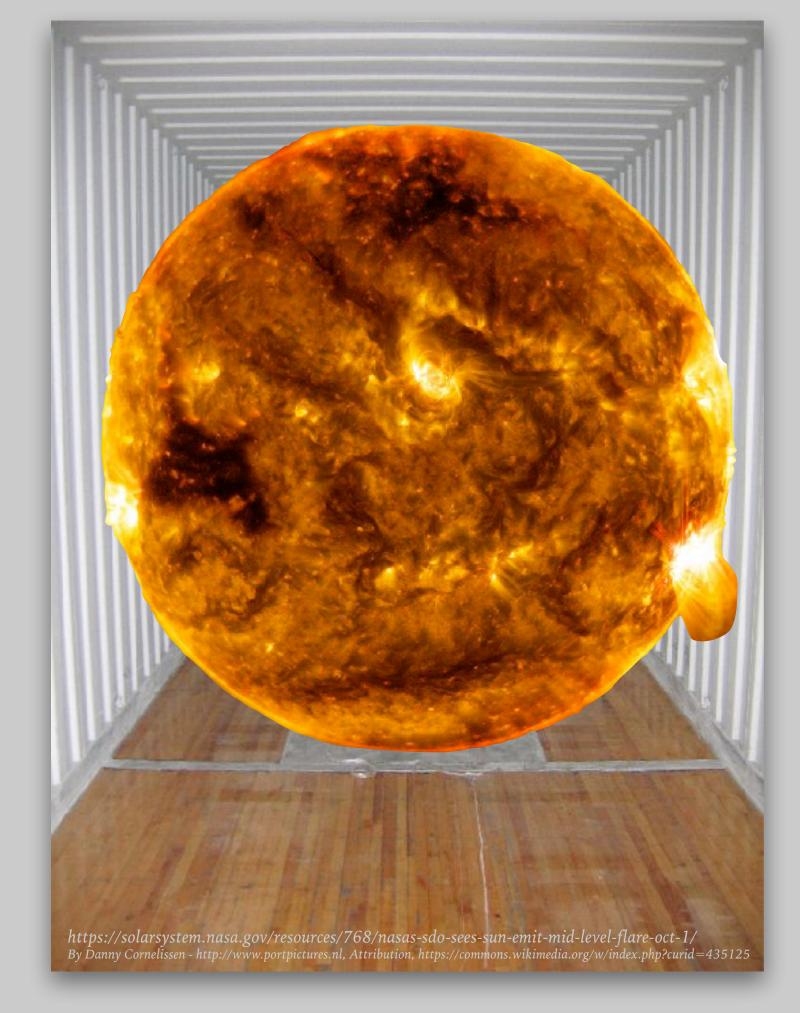
What does 2.6×10²⁸ kg/m³ mean?





What does 2.6×10²⁸ kg/m³ mean?





fit the mass of the sun into a standard 40ft shipping container

cosmological constant & fine-tuning [classically]

$$V_{min} = \left[-\mu^2 |\phi|^2 + \lambda |\phi|^4 \right]_{\phi_0} + V_0$$

$$= -2.6 \times 10^{28} \, kg/m^3 + V_0 = 5.96 \times 10^{-27} \, kg/m^3$$

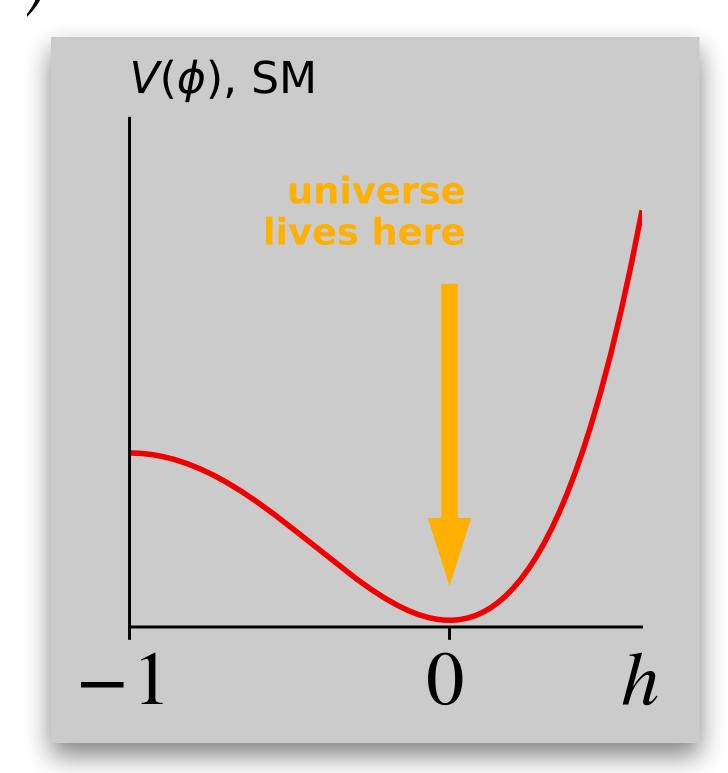
- $ightharpoonup V_0$ needs to be fine tuned for cosmological constant to have today's size (also with respect to various sources of quantum correction)
- > not the only fine-tuning problem in fundamental physics,
 - arguably special in that it appears already classically
- \succ collider physics cannot tell us anything about V_0
 - but it would seem negligent not to try and establish the rest of the potential

The potential expanded around the minimum

➤ take *h* as the Higgs field excitation in units of the field at minimum

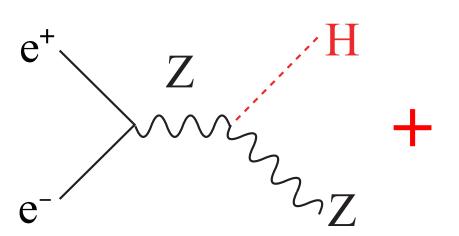
$$V = \frac{m_H^2 v^2}{8} \left(-1 + 4h^2 + 4h^3 + h^4\right)$$
the Higgs boson mass term

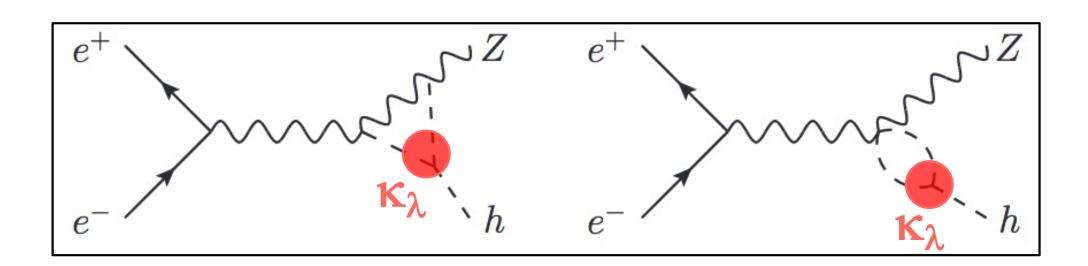
prediction of the strength of HHH interaction [modifier may be called κ_{λ} or κ_{3}]

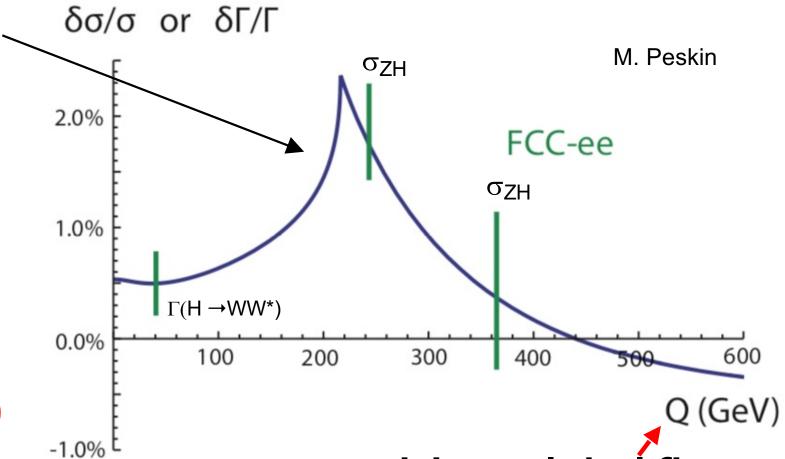


Higgs self-coupling at FCC-ee

- Statistics-limited sensitivity comes from $\sigma_{ee \to ZH}$ measurements at 240 and 365 GeV
 - Thanks to the relative change with centre-of-mass energy





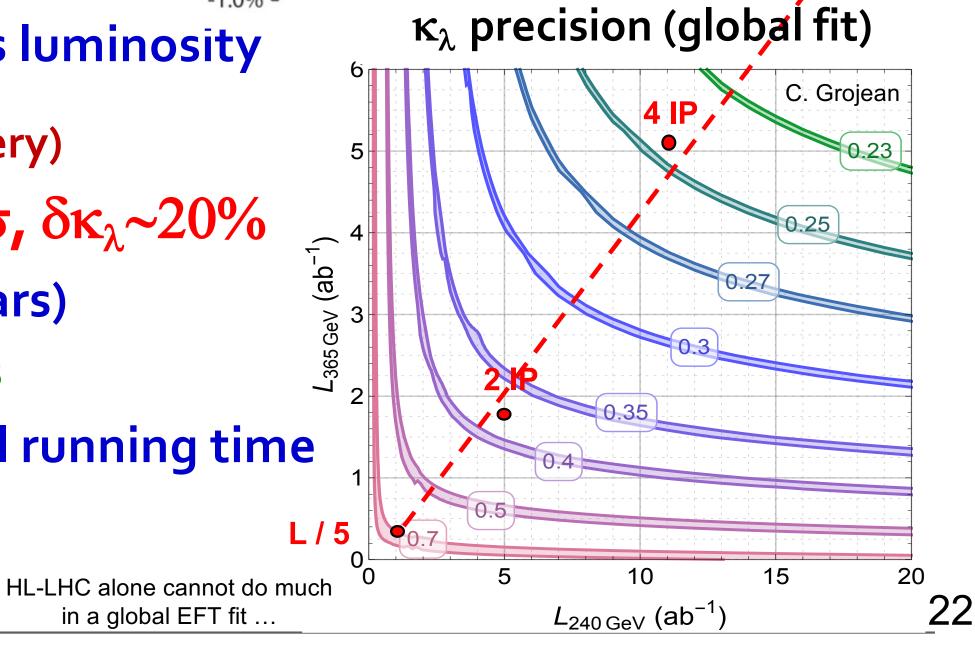


- Estimate with present run plan and 2 IPs: $\geq 2\sigma$ from $\kappa_{\lambda} = 0$
 - Analyses will improve, but no hope with 5 times less luminosity

(Discovery)

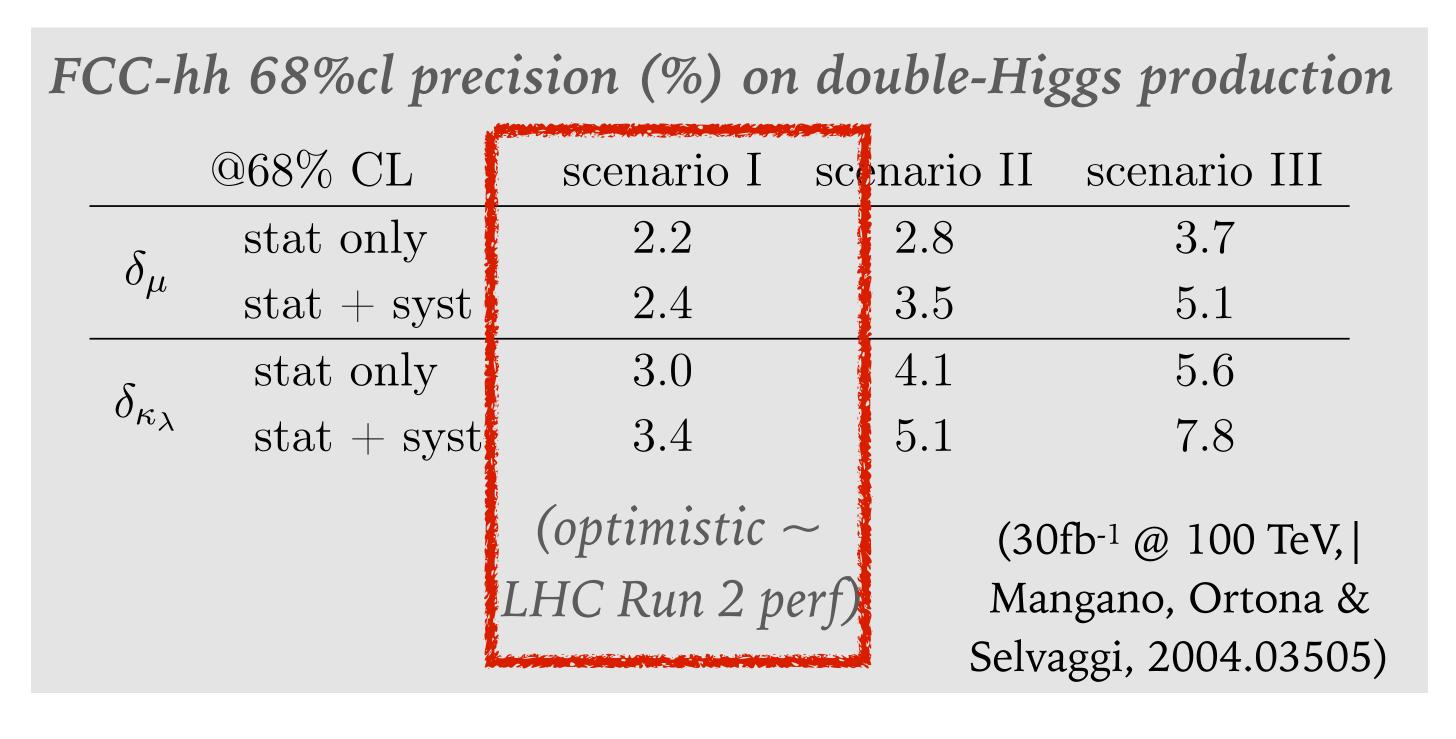
in a global EFT fit ...

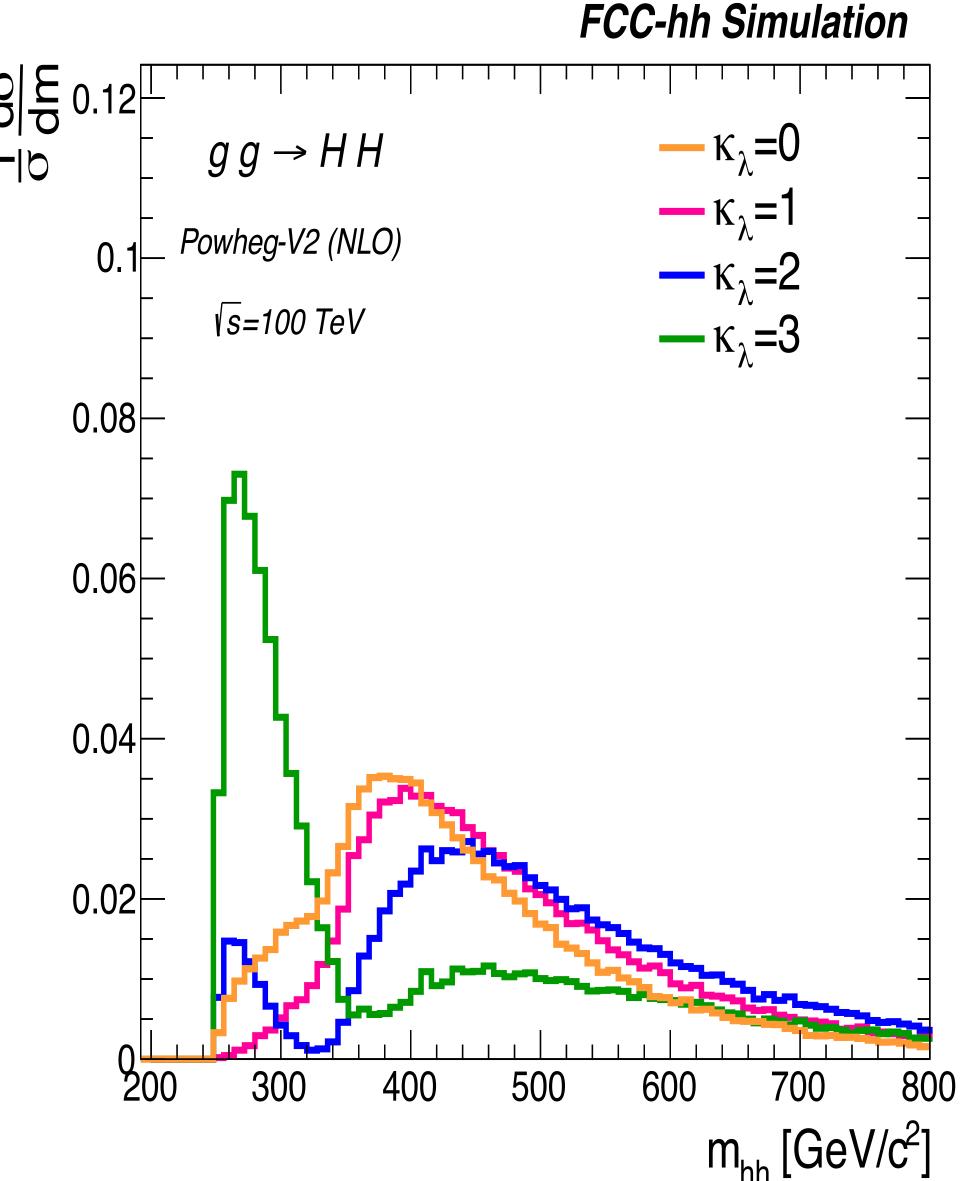
- With 4 IPs and optimization of run plan: target $\geq 5\sigma$, $\delta\kappa_{\lambda} \sim 20\%$
 - Increase duration at 240 and 365 GeV (to 4 and 7 years)
 - Reduce Z and WW run duration @ constant statistics
 - Or better: increase specific luminosity and/or overall running time
 - If it is worth doing, it is worth doing well



Testing SM V(φ) by measuring HH production at FCC:~3-5% accuracy

- ➤ kinematic shape of HH pair clearly distinguishes independent HH production from correlated HH
- ➤ FCC-hh → few % determination (needs accurate $t\bar{t}Z$ and Higgs couplings from FCC-ee)



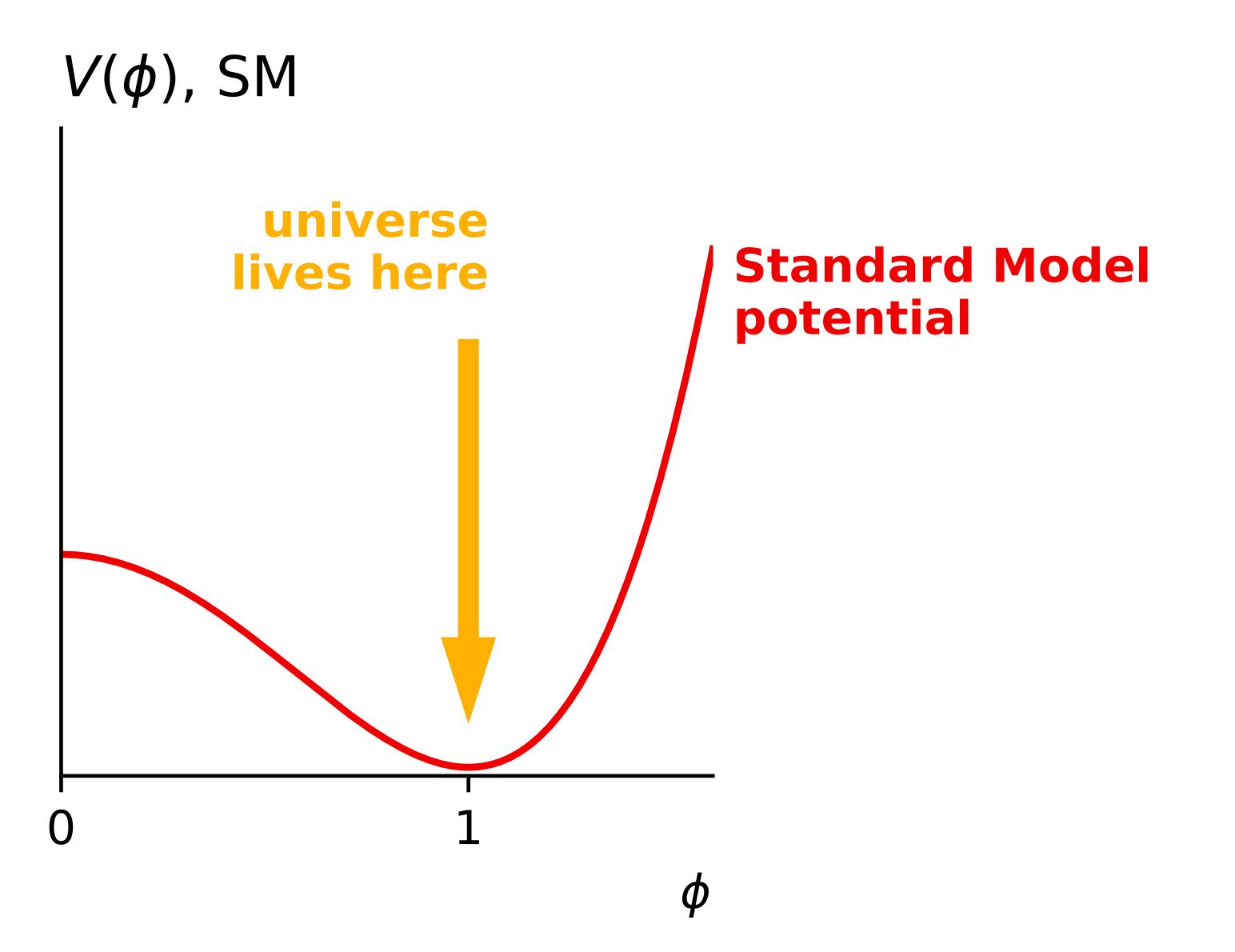


when would we claim discovery? [5σ in each of two independent experiments is our gold standard]

- ➤ equivalent for an interaction is a bit ambiguous but better than ±20% determination is probably a reasonable target
- ➤ for something of this importance, I am wary of relying on 20% only from a combination of N experiments a result's robustness comes from confirmation by independent experiments
- ➤ indirect v. direct:
 - ➤ all measurements are indirect (we measure hadrons and leptons...)
 - > single H is good to have
 - ➤ but HH & kinematic structure brings assurance that what we are seeing is indeed HHH coupling
- > NB there exist different points of view on this

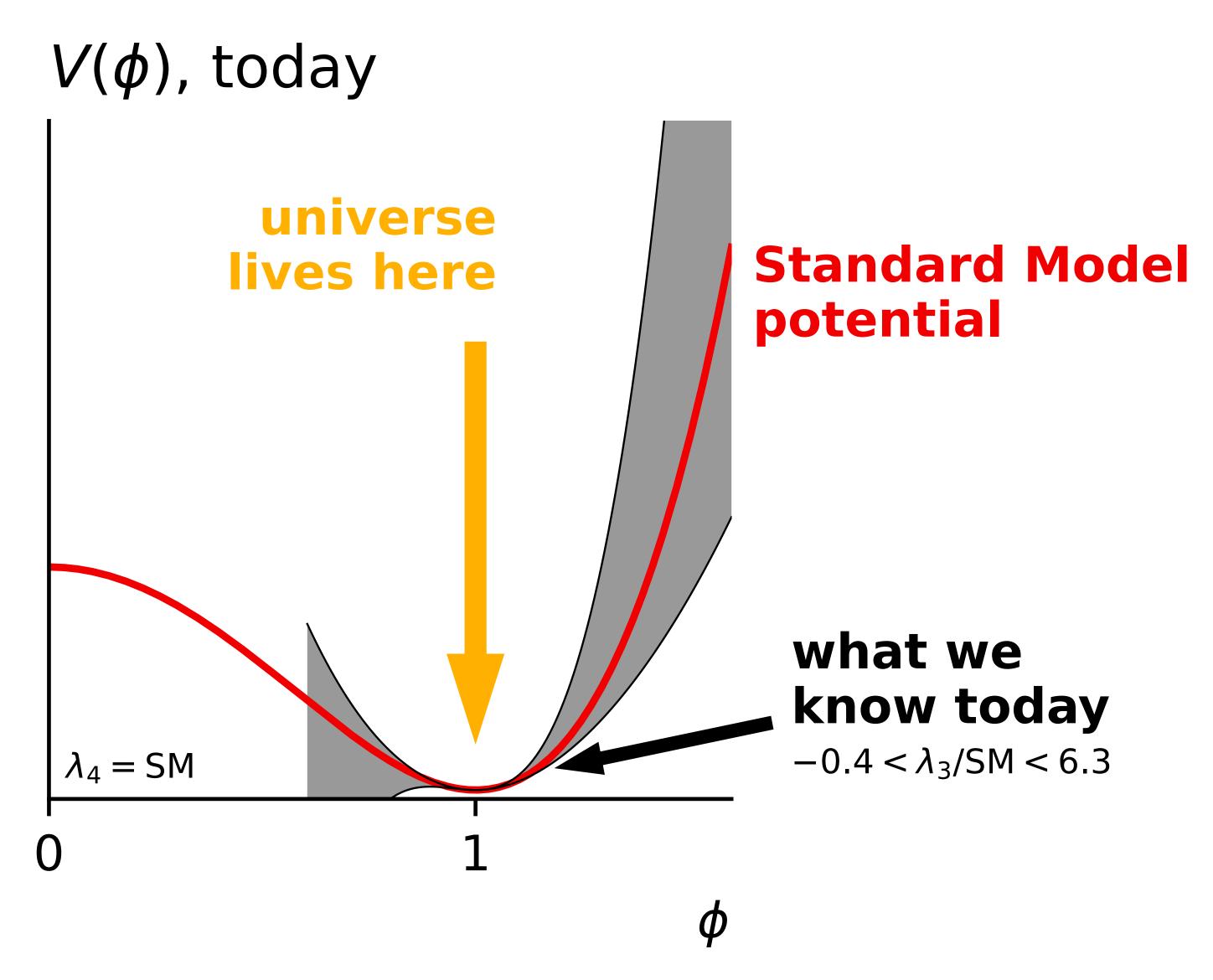
when would we claim discovery? [5σ in each of two independent experiments is our gold standard]

- > equivalent for an interaction is a bit ambiguous but better than +2
- tion by
 - - > single H w have
 - ➤ but HH & kinematic structure brings assurance that what we are seeing is indeed HHH coupling
 - ➤ NB there exist different points of view on this



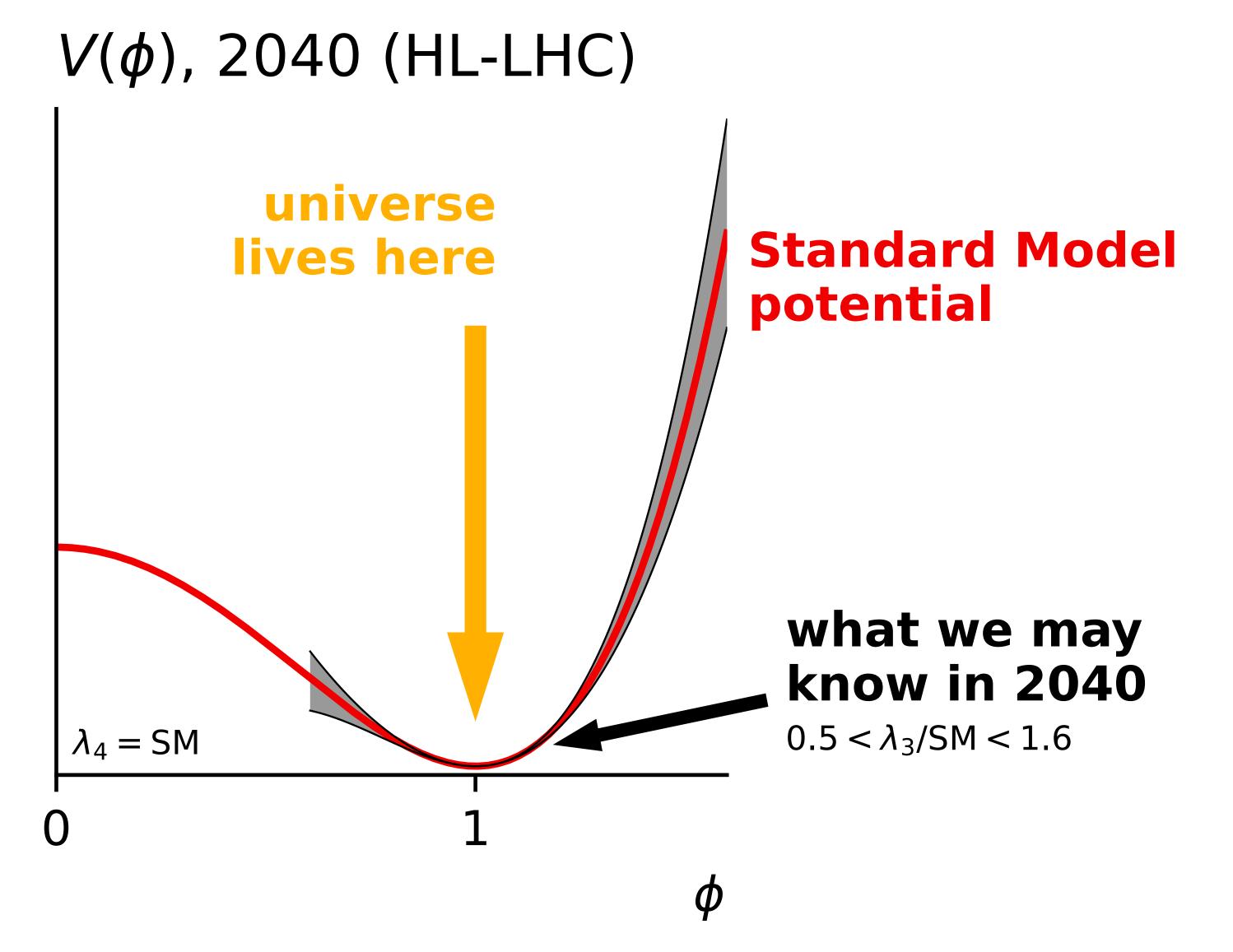
- > this is a cartoon
- ➤ caution needed: e.g.

 realistic BSM models do not
 just modify the potential,
 but may bring extra scalars
 (often modify other couplings, but not
 always, e.g. 2209.00666)
- rightharpoonup even if we take the picture seriously we may want to consider impact of limited constraints on λ_4



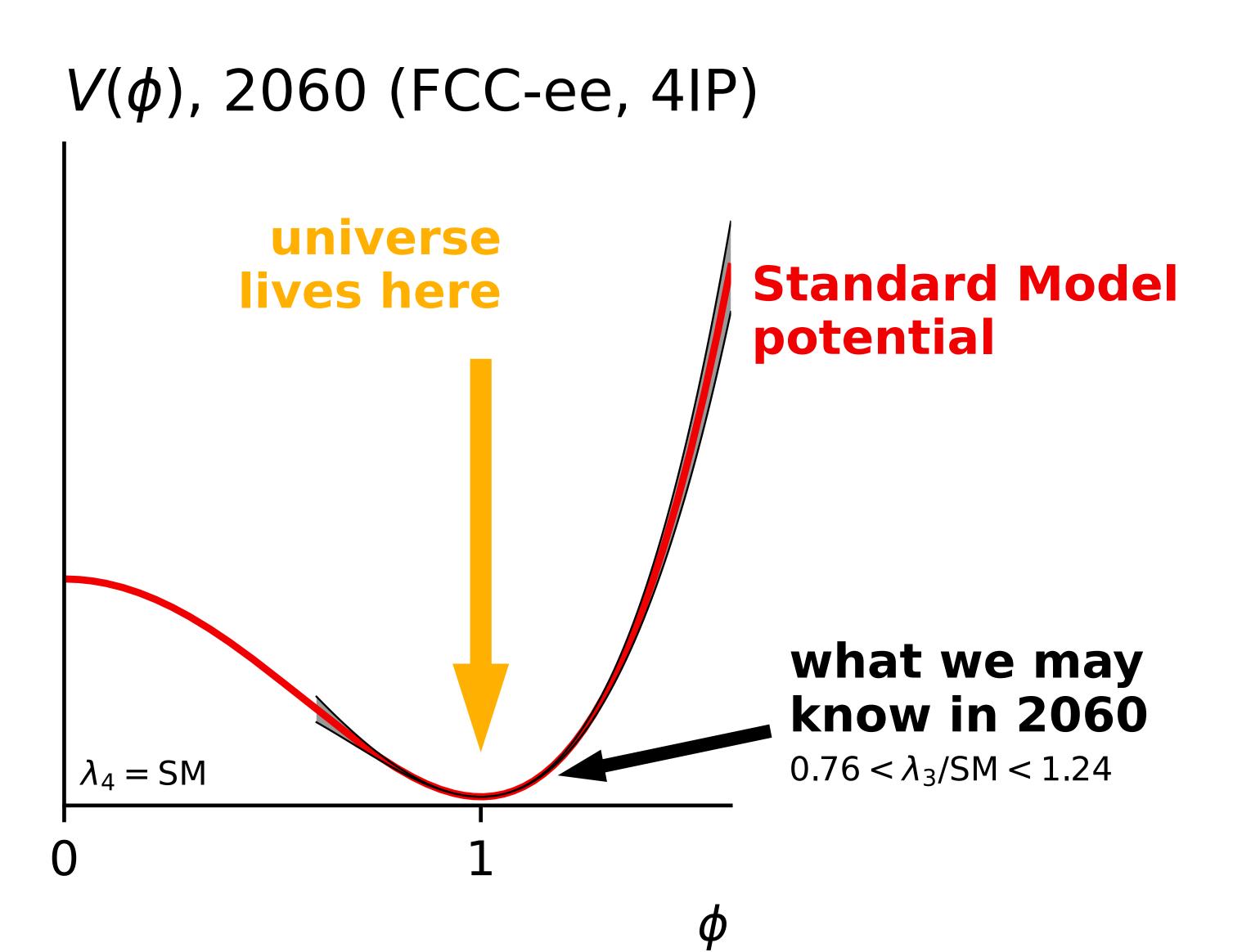
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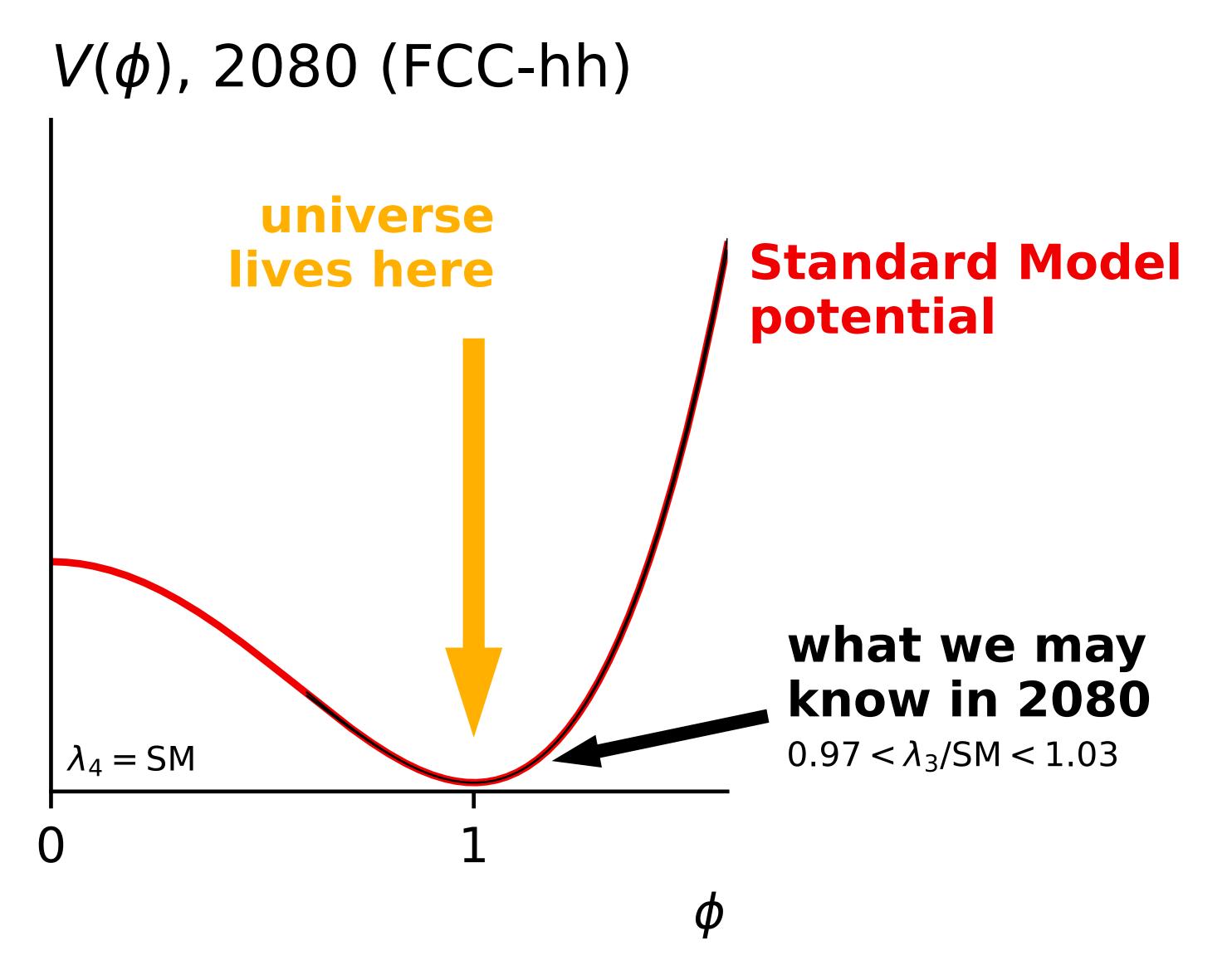
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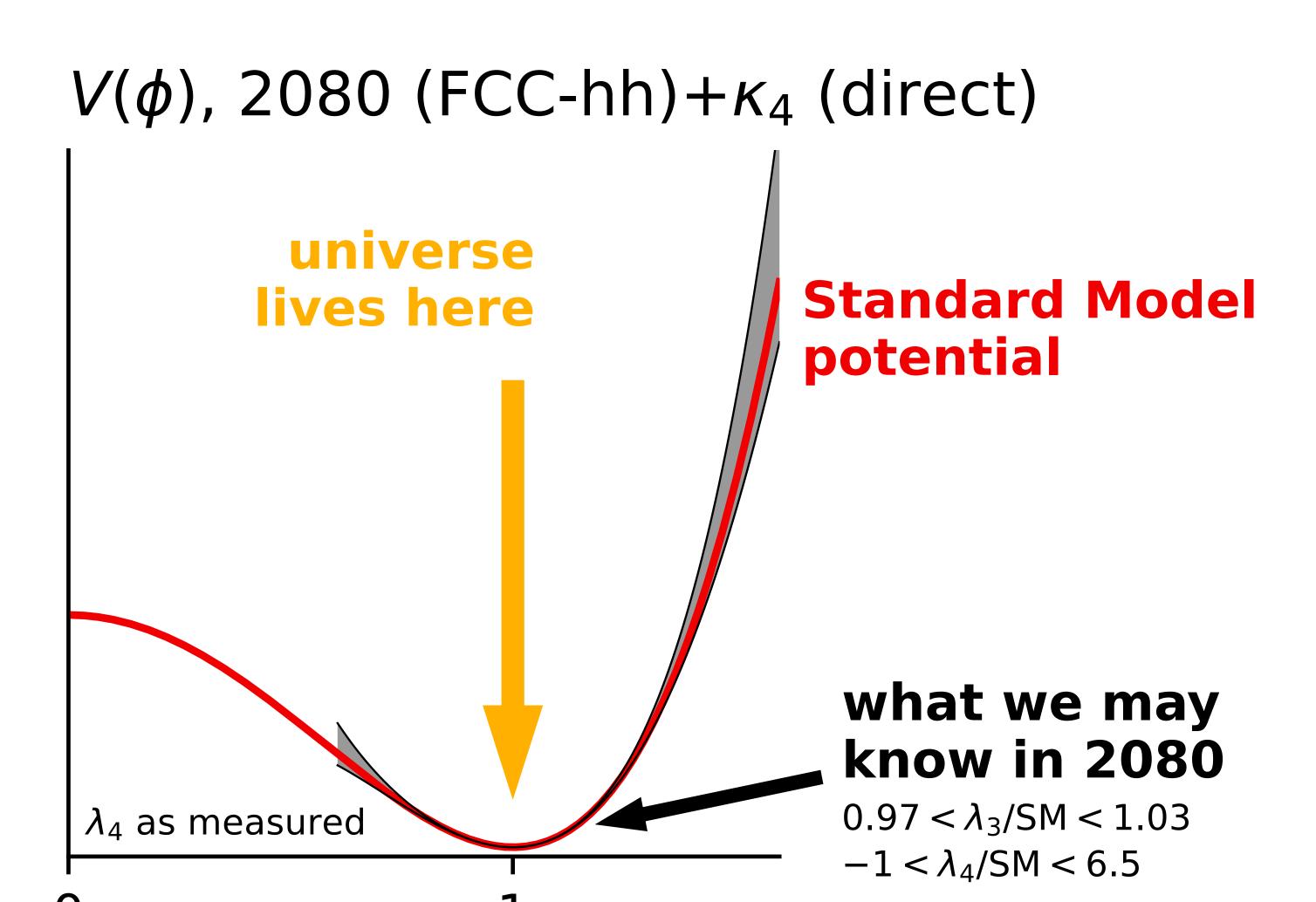
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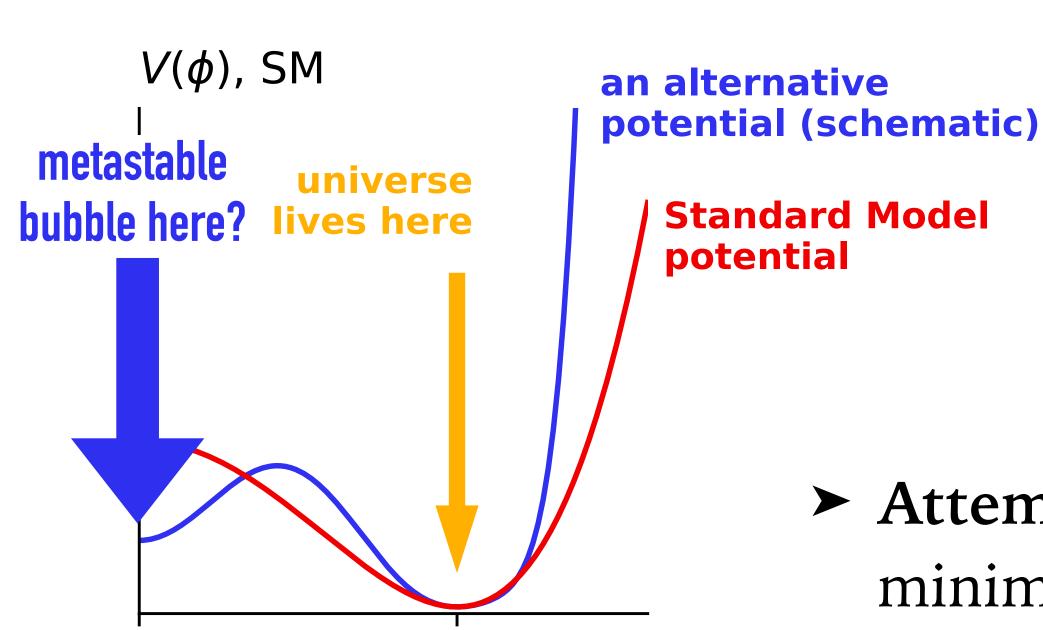
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A wildly speculative aside [science fiction!]



- common argument for fundamental research: it may pay off in terms of technological advances in a century or two
- ➤ in particle physics, it's hard to conceive of a way in which this could be true
- ➤ Attempt at counterexample: if there were 2nd minimum in Higgs potential, could we create metastable bubbles of alternative vacuum? (cf. EW phase transition)
- ➤ likely very short lifetime, unless some kind of protection
- ➤ what might we do with it? E.g. very different nuclear physics, if light quarks get all mass from Yukawa interactions, long-range strong force (pion ~ massless), etc.
- ➤ this scenario is very far fetched: do not take it seriously! (But we can't even tell how far fetched it is if we haven't measured the potential)

FCC week, London, June 2023

desirable features of a worldwide HEP project?

an important target that is guaranteed to be reached (no-lose theorem)

exploration into the unknown by a significant factor in energy

major progress on a broad array of particle physics topics

likelihood of success, robustness (incl. multiple experiments)

cost-effective construction & operation, low carbon footprint

various arguments favour a circular e+e- collider [you all know them well]

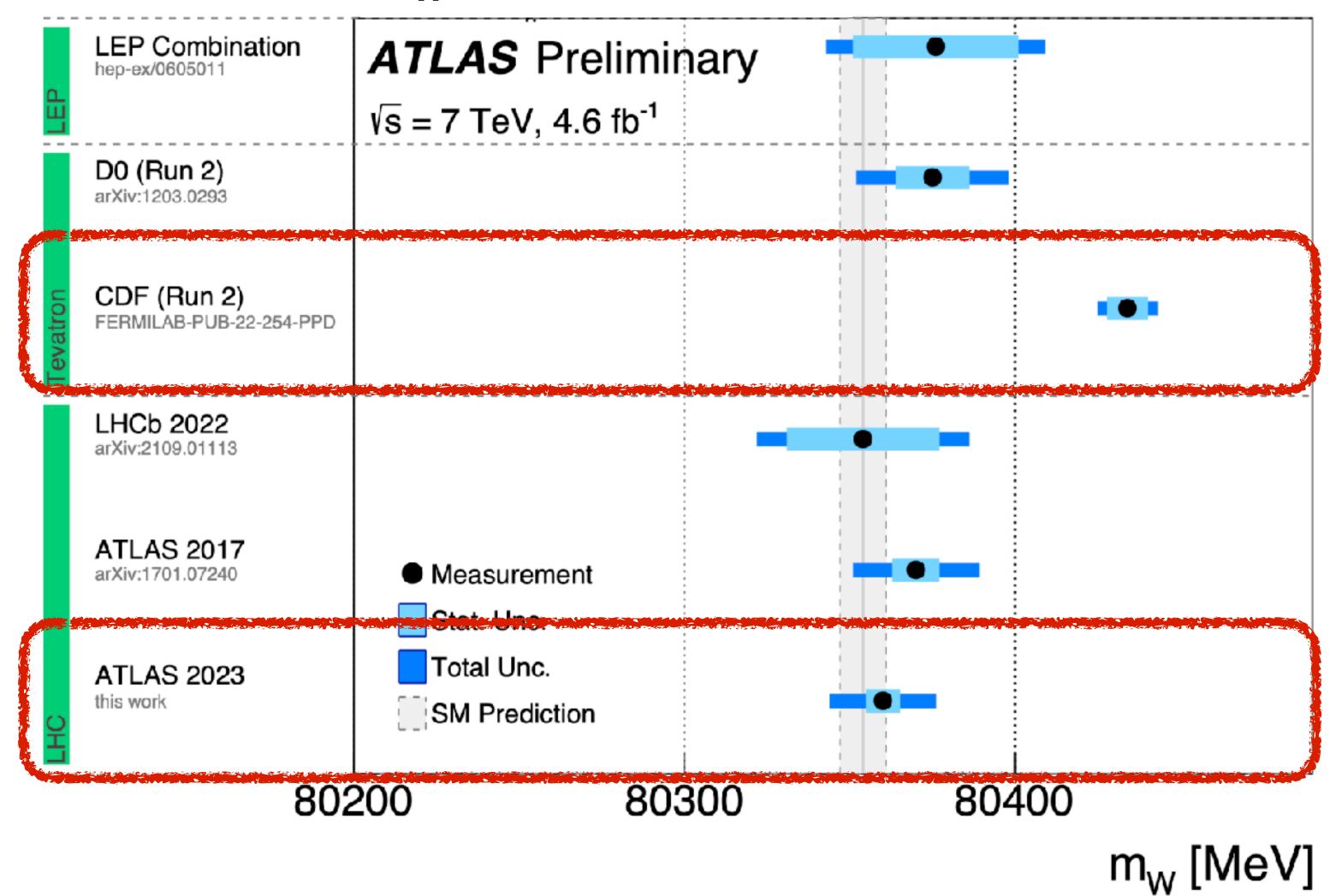
- ➤ historical track record of delivering luminosity [LEP]
- > unlike linear colliders, they naturally accommodate multiple experiments
- > energy efficiency/unit luminosity from Z-pole to ZH
- > electrons are a lot easier than muons

But some people ask if we need a lepton collider at all; should we not just go for the next hadron collider?

[practical arguments against: we don't really know how to build the magnets for a 100 TeV collider; cost of 91km collider is high even with LHC-type magnets]

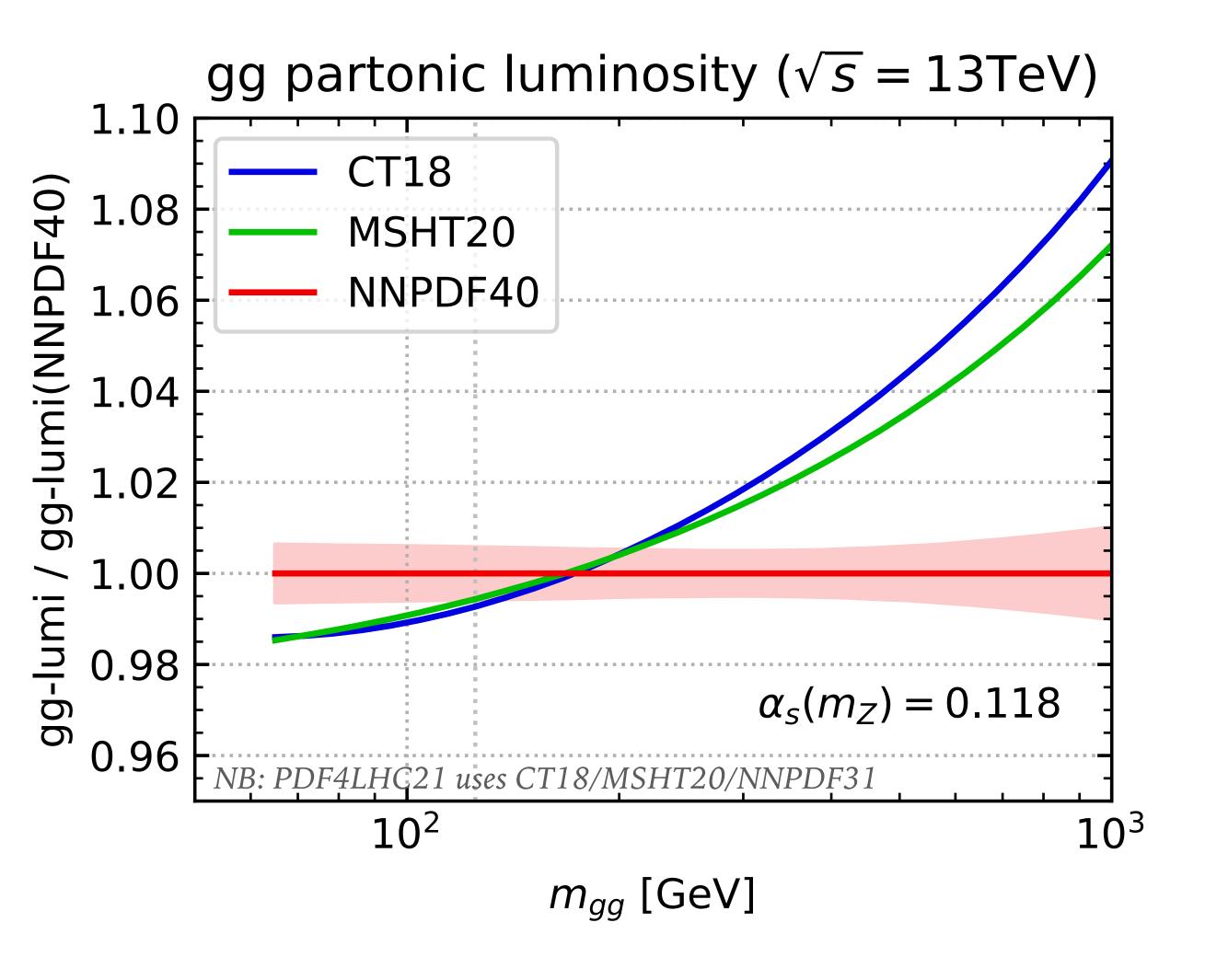
36

mw measurements



do you believe the measurement when it **disagrees** with your expectations?

we don't know the precision limit of hadron colliders — but we may be close to reaching it



gg-lumi, ratio to PDF4LHC15 @ m_H

PDF4LHC15	1.0000	\pm	0.0184	k
PDF4LHC21	0.9930	\pm	0.0155	
CT18	0.9914	\pm	0.0180	× 3
MSHT20	0.9930	土	0.0108	
NNPDF40	0.9986	<u>±</u>	0.0058	

Parton Distribution Functions are one of several elements that may limit LHC/FCC-hh precision:

- > essential for hadron-collider interpretation
- ➤ PDF fits are complex, e.g. involve (sometimes inconsistent) data, some of it close to non-perturbative scale
- > only partial understanding of their limits

first approx N3L0 PDFs

$gg \rightarrow H \; (\mu = m_H/2)$ [adapted], all with N3LO σ 50 NNLO PDFs 45 aN 3 LO $(H_{ij} + K_{ij})^{-1}$ PDFs aN 3 LO H'_{ij} PDFs 40 Light: PDF + Scale uncertainty Dark: PDF uncertainty

σ order	PDF order	$\sigma \text{ (pb)} + \Delta \sigma_{+} - \Delta \sigma_{-} \text{ (\%)}$
		PDF uncertainties
	aN ³ LO (no theory unc.)	44.164 + 3.03% - 3.13%
	aN ³ LO $(H_{ij} + K_{ij})$	44.164 + 3.34% - 3.15%
	$aN^3LO(H'_{ij})$	44.164 + 3.43% - 3.07%
	NNLO	47.817 + 1.17% - 1.22%

Approximate N³LO Parton Distribution Functions with Theoretical Uncertainties:

MSHT20aN³LO PDFs

arXiv:2207.04739v1

J. McGowan^a, T. Cridge^a, L. A. Harland-Lang^b, and R.S. Thorne^a

- ➤ includes approximations & datadriven fits to parts of N3LO currently unknown
- > 7.6% decrease in Higgs cross section (w. N3LO σ)
- ➤ PDF part of uncertainty goes up by ×2.5–3
- ➤ fairly surprising; starting point for many future investigations

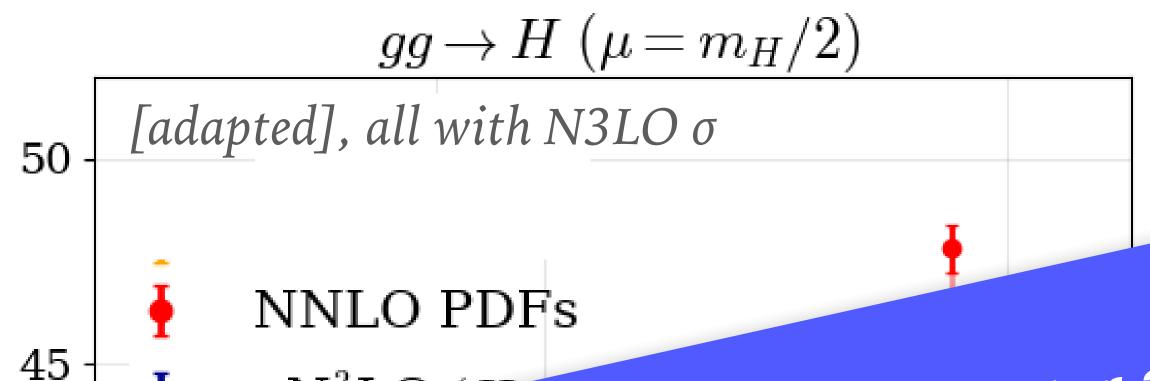
first approx N3L0 PDFs

Approximate N³LO Parton Distribution Functions with Theoretical Uncertainties:



arXiv:2207.04739v1

J. McGowan^a, T. Cridge^a, L. A



a lepton collider as a next step ensures solid foundations for the field e.g. measurement of H→gg at 1% at FCC-ee underpins precision of FCC-hh

 $44.164 + 3.03\% - \overline{3.13\%}$ neory unc.) $aN^3LO(H_{ij} + K_{ij})$ 44.164 + 3.34% - 3.15% N^3LO $aN^3LO(H'_{ij})$ 44.164 + 3.43% - 3.07%47.817 + 1.17% - 1.22%NNLO

- PDF part of uncertainty goes up by $\times 2.5 - 3$
- ➤ fairly surprising; starting point for many future investigations

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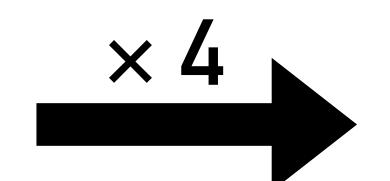
what should we expect as a step up in energy?

I like the $Z'_{\rm SSM}$ as a simple measure of progress (perhaps not very "exciting", but simple and most experiments look for it)

Tevatron $p\bar{p}$, 1.96 TeV, 10 fb⁻¹

Exclusion limit ~ 1.2 TeV

(if they had analysed all their data in electron and muon channels; actual CDF limit 1.071 TeV, 4.7fb⁻¹, μμ only)



LHC pp, 13.6 TeV, 139 fb⁻¹

Exclusion limit ~ 5.1 TeV

(electron and muon channels, single experiment)

what should we expect as a step up in energy?

I like the $Z'_{\rm SSM}$ as a simple measure of progress (perhaps not very "exciting", but simple and most experiments look for it)

LHC *pp*, 13 TeV, 139 fb⁻¹

Exclusion limit ~ 5.1 TeV

(electron and muon channels, single experiment)

× 7.8

FCC-hh pp, 100 TeV, 20 ab⁻¹

Exclusion limit ~ 41 TeV

(based on PDF luminosity scaling, assuming detectors can handle muons and electrons at these energies)

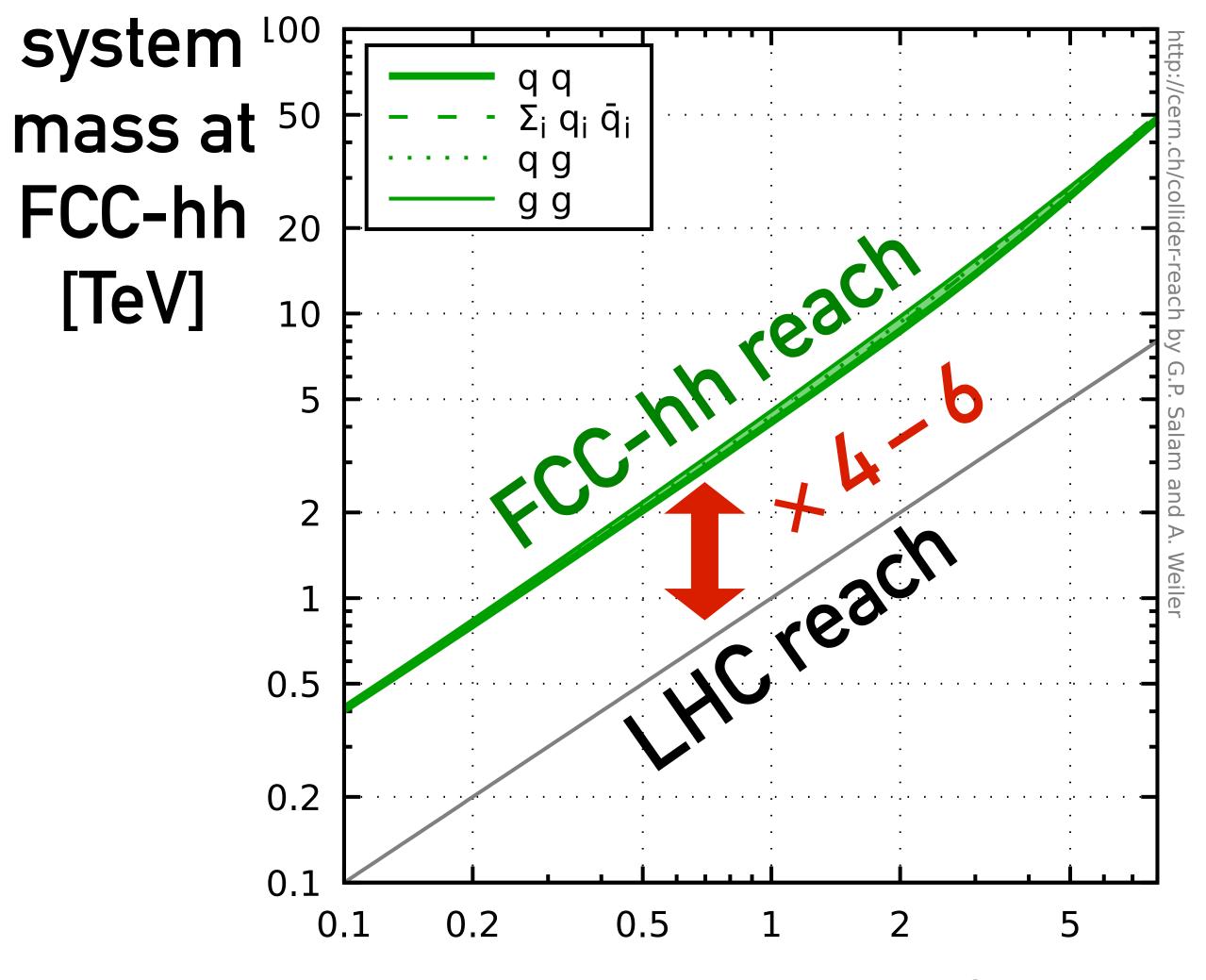
LHC $3 \text{ ab}^{-1} \rightarrow \text{FCC-hh } 20 \text{ ab}^{-1}$

Collider 1: CoM energy 14 TeV, integrated luminosity Collider 2: CoM energy 100 TeV, integrated luminosity

3000 fb⁻¹

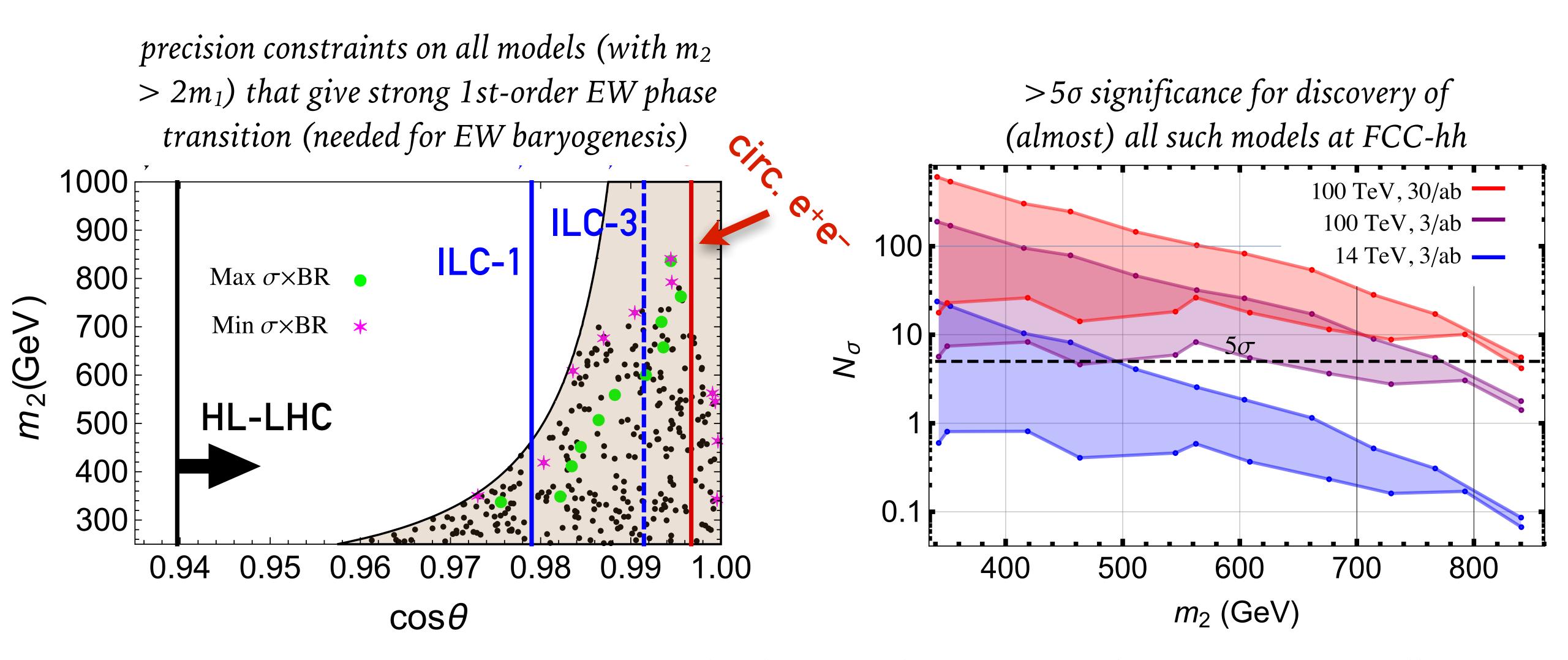
20000 fb⁻¹

FCC-hh delivers the kind of step up in direct-search sensitivity (\times 4 – 6) that we would hope for



system mass at LHC [TeV]

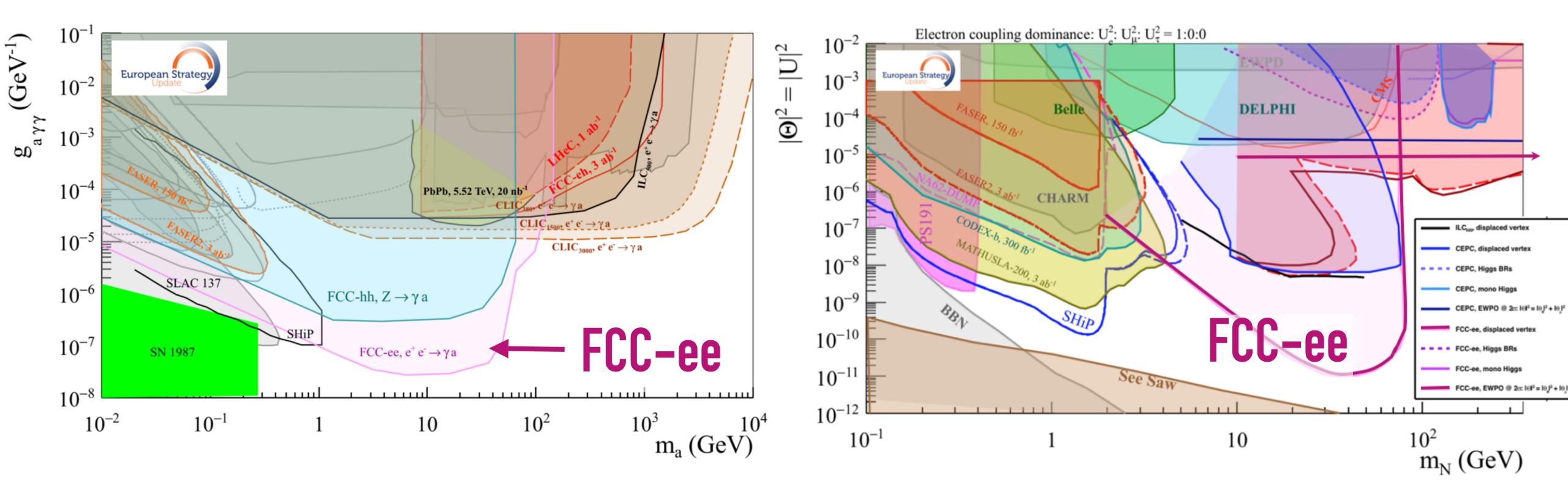
Extension of SM with one extra scalar ("h2", gauge singlet)



1605.06123

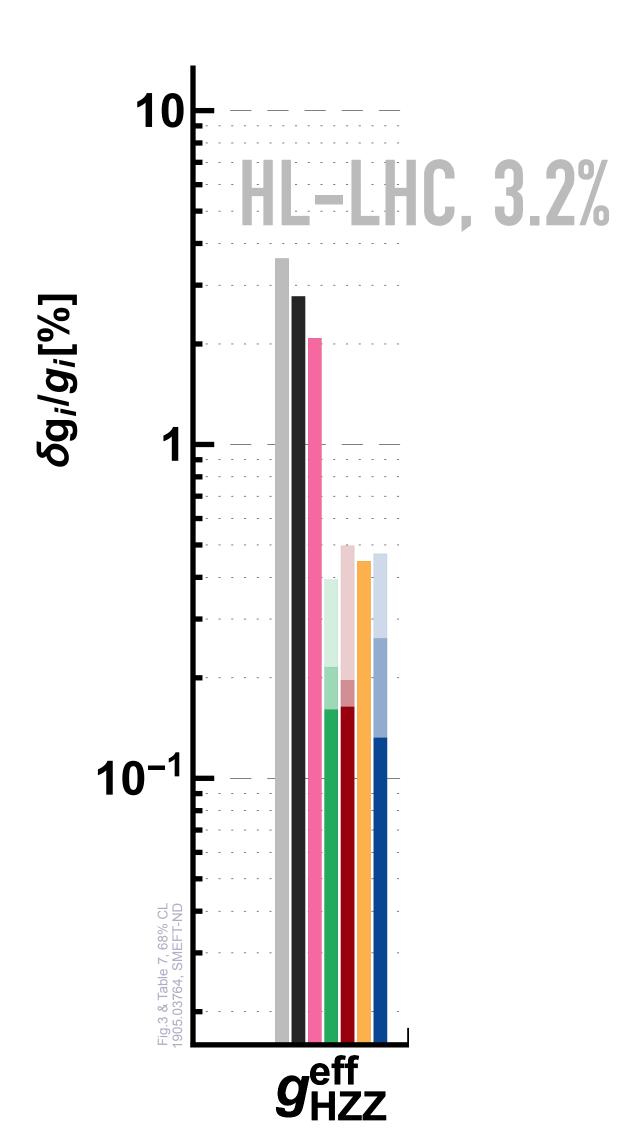
It is important to take these conclusions somewhat impressionistically, as we have made a number of simplifying assumptions in order to paint the broad picture.

FCC-ee, e.g. axion and heavy-neutral lepton searches

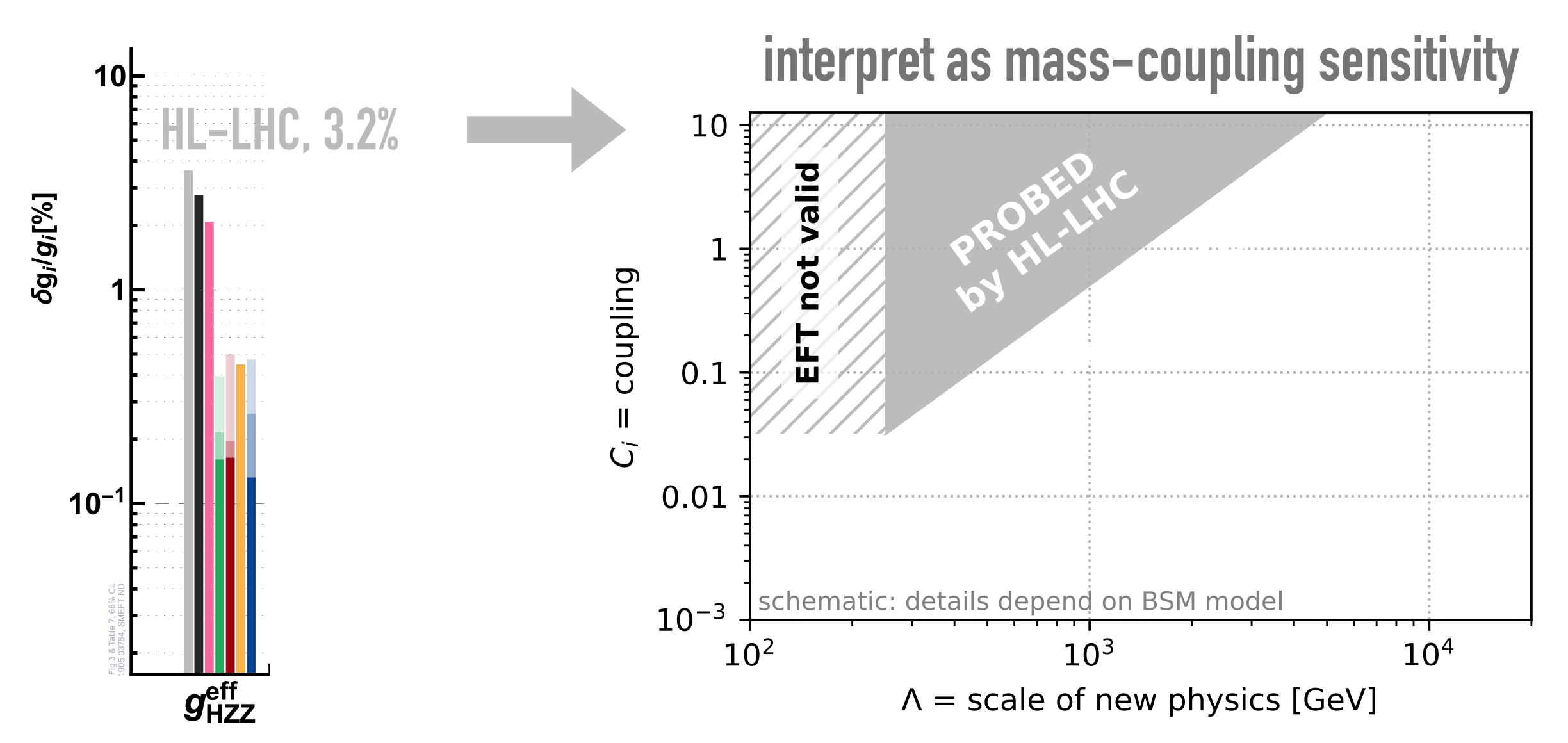


benefits from huge Z-pole luminosity (some models in these regions have potential to connect with dark matter, baryon asymmetry, neutrino masses, etc.)

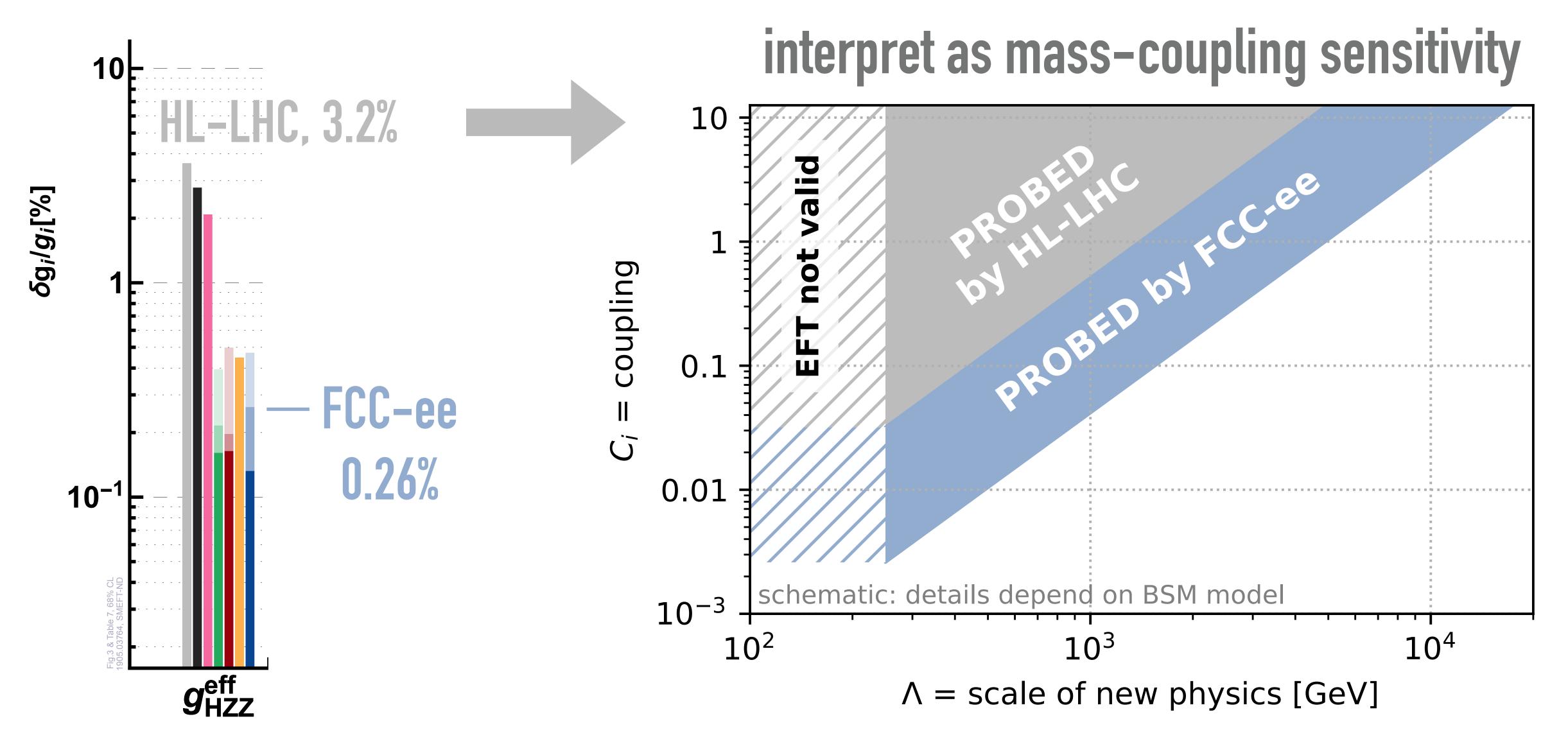
Interpret higher precision as increase in indirect reach



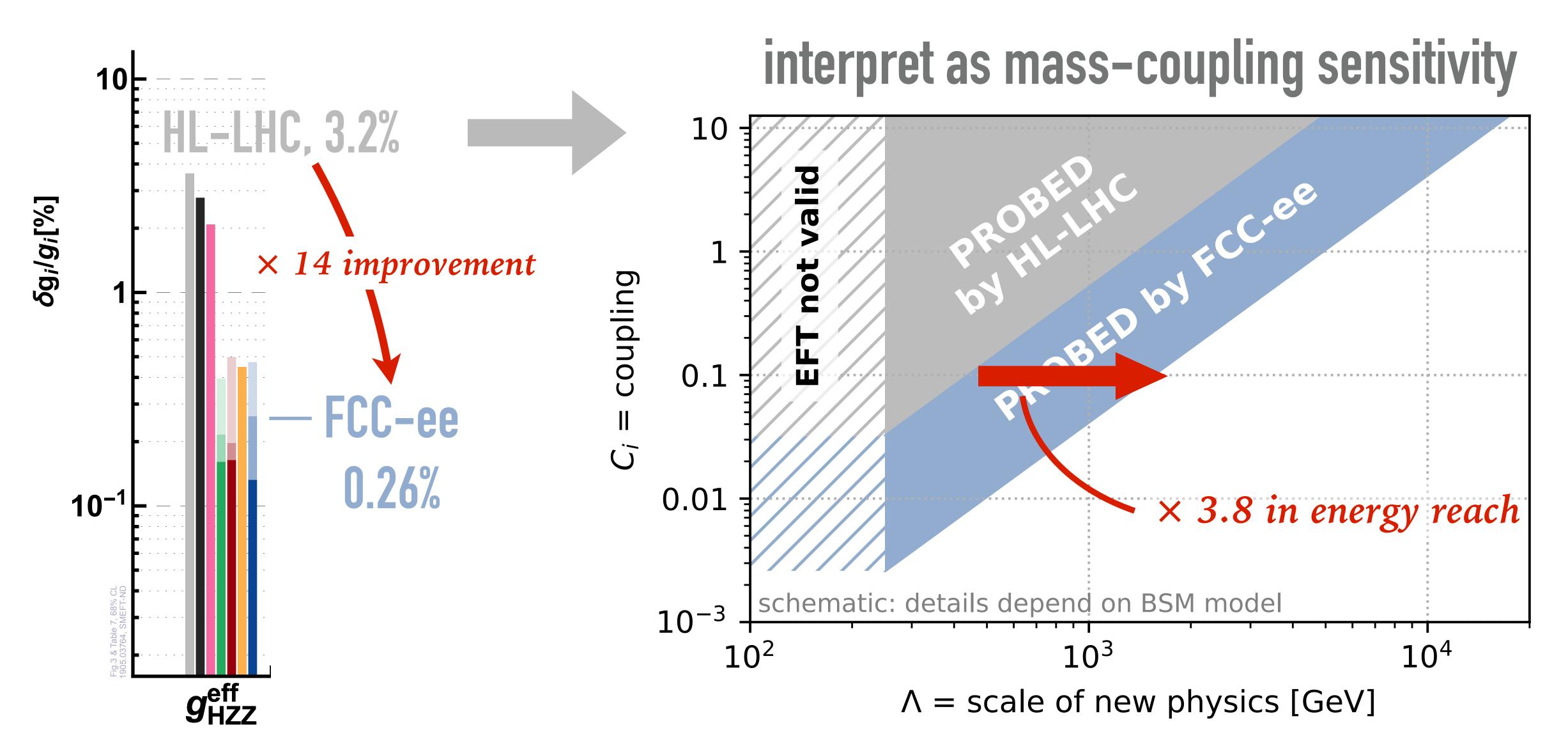
Interpret higher precision as increase in indirect reach



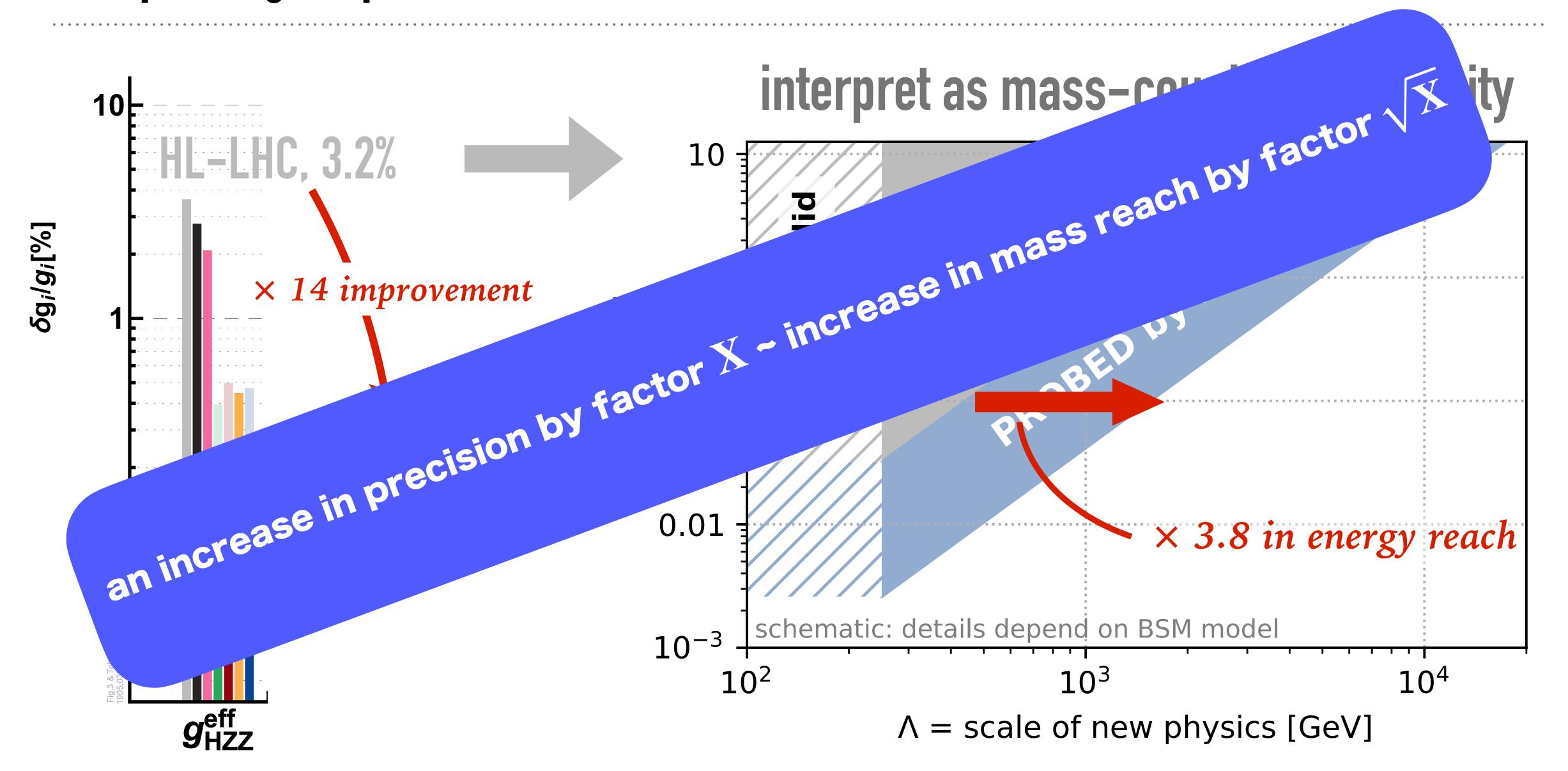
Interpret higher precision as increase in indirect reach



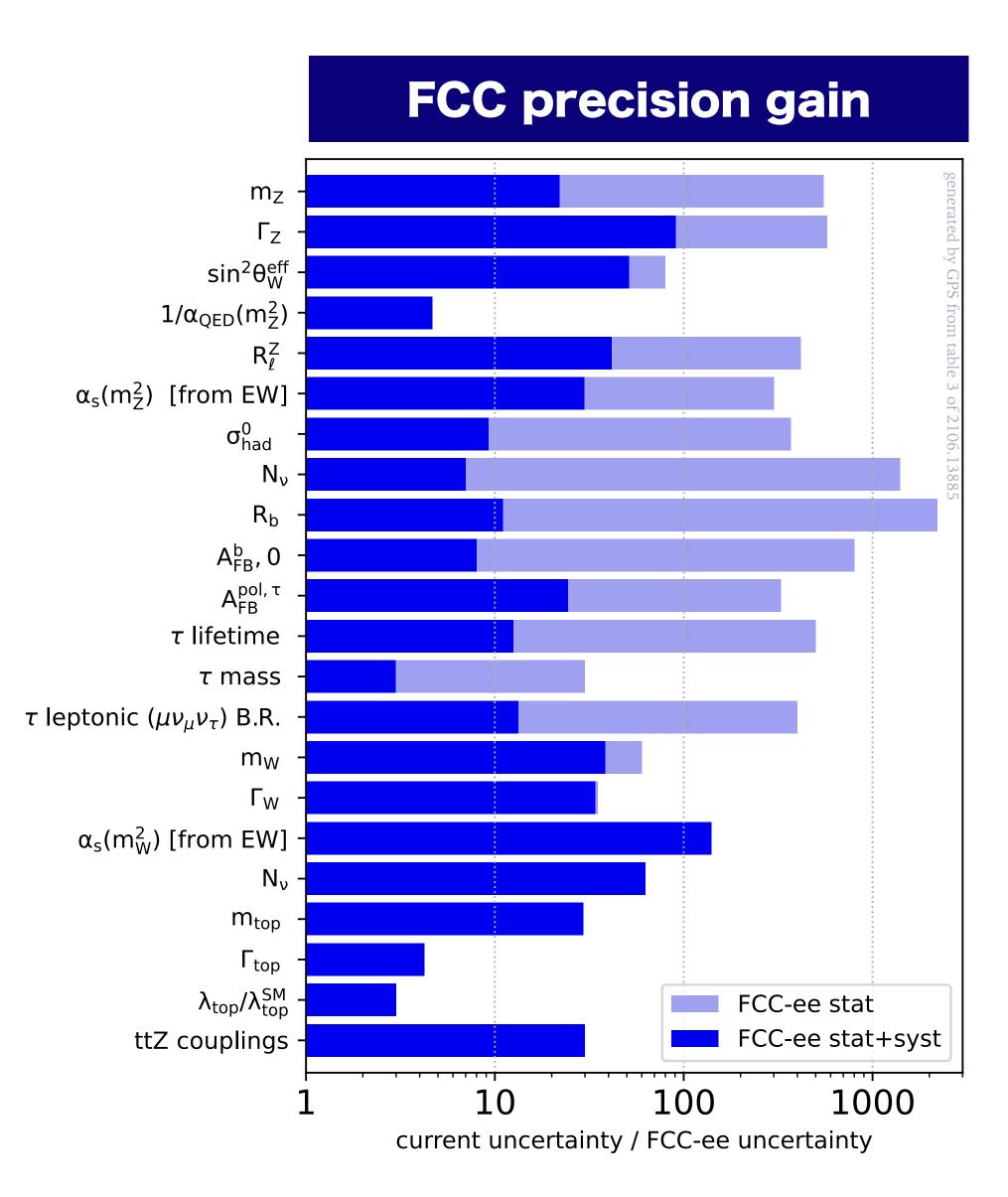
Interpret higher precision as increase in indirect reach

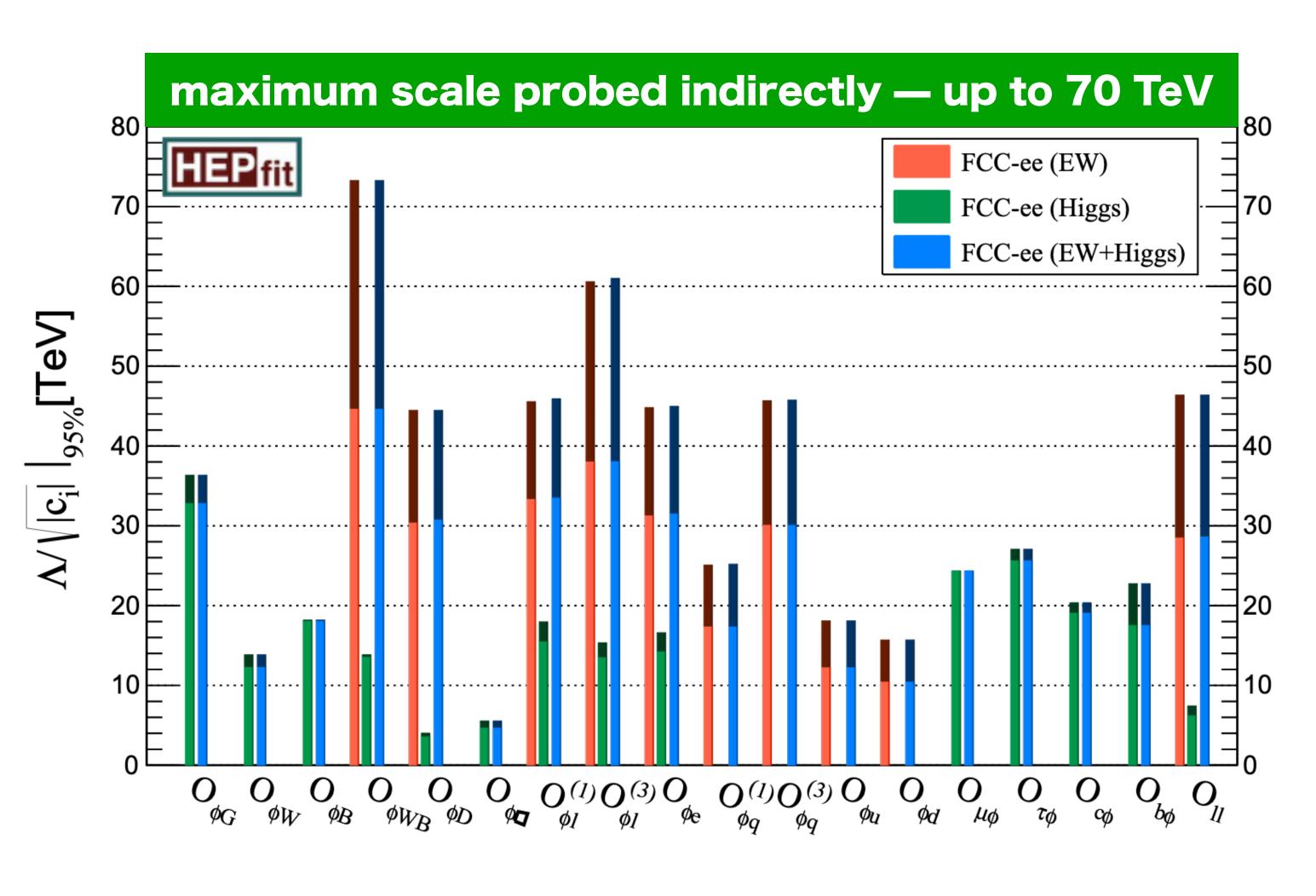


Interpret higher precision as increase in indirect reach

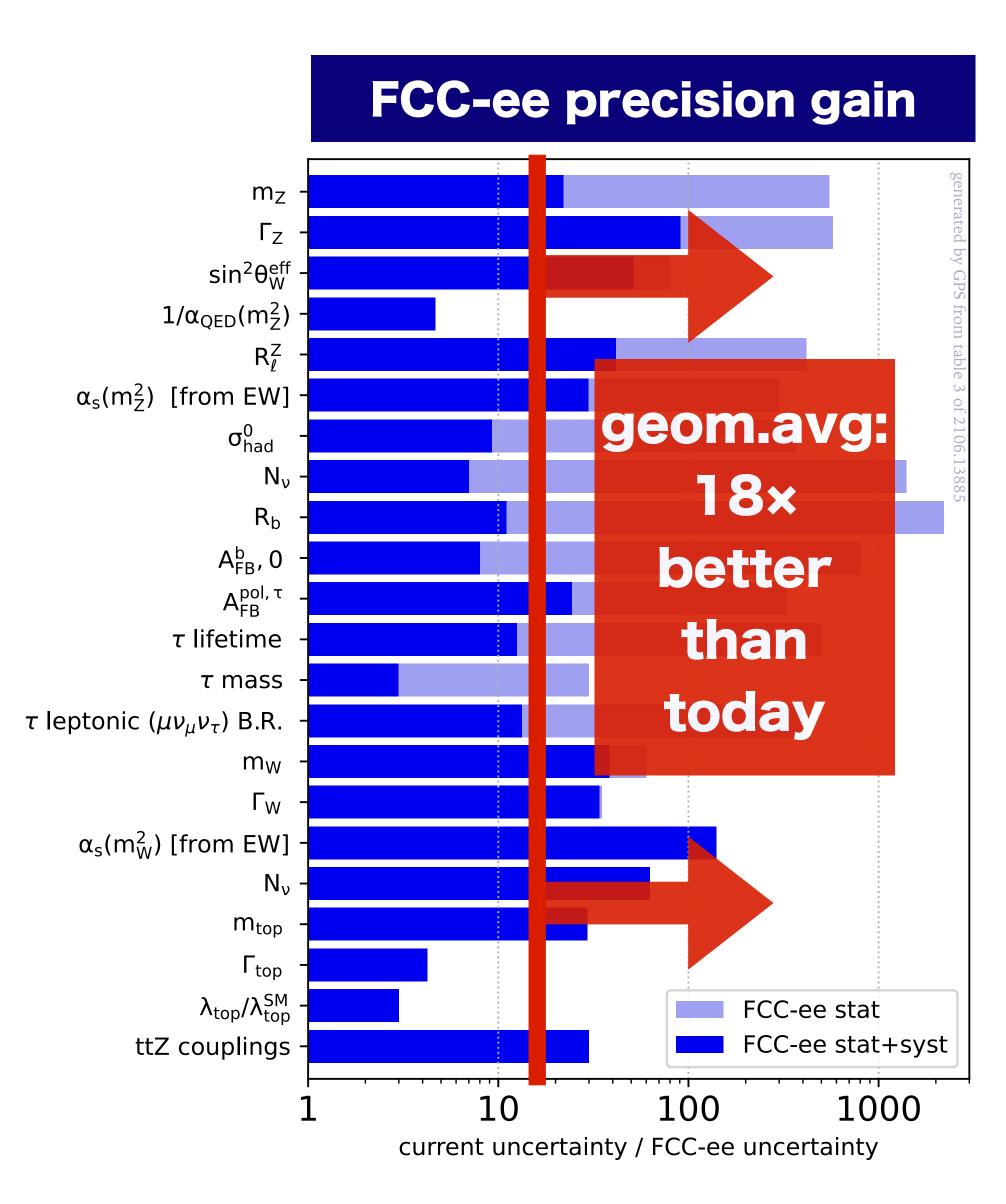


increase in precision at FCC-ee is equivalent to \times 4 – 5 increase in energy reach





increase in precision at FCC-ee is equivalent to \times 4 – 5 increase in energy reach



Two messages

- ➤ with a rough estimate for systematics, FCC brings a big step forward (geom.avg.
 - $= \times 18$, across $\gtrsim 20$ observables)
- still huge scope for thinking about how to improve systematics (gain of up to further × 100 in some cases)

This is the fun part for us as physicists!

and will call for joint efforts by

experiment/theory/accelerator

physicists

desirable features of a worldwide HEP project?

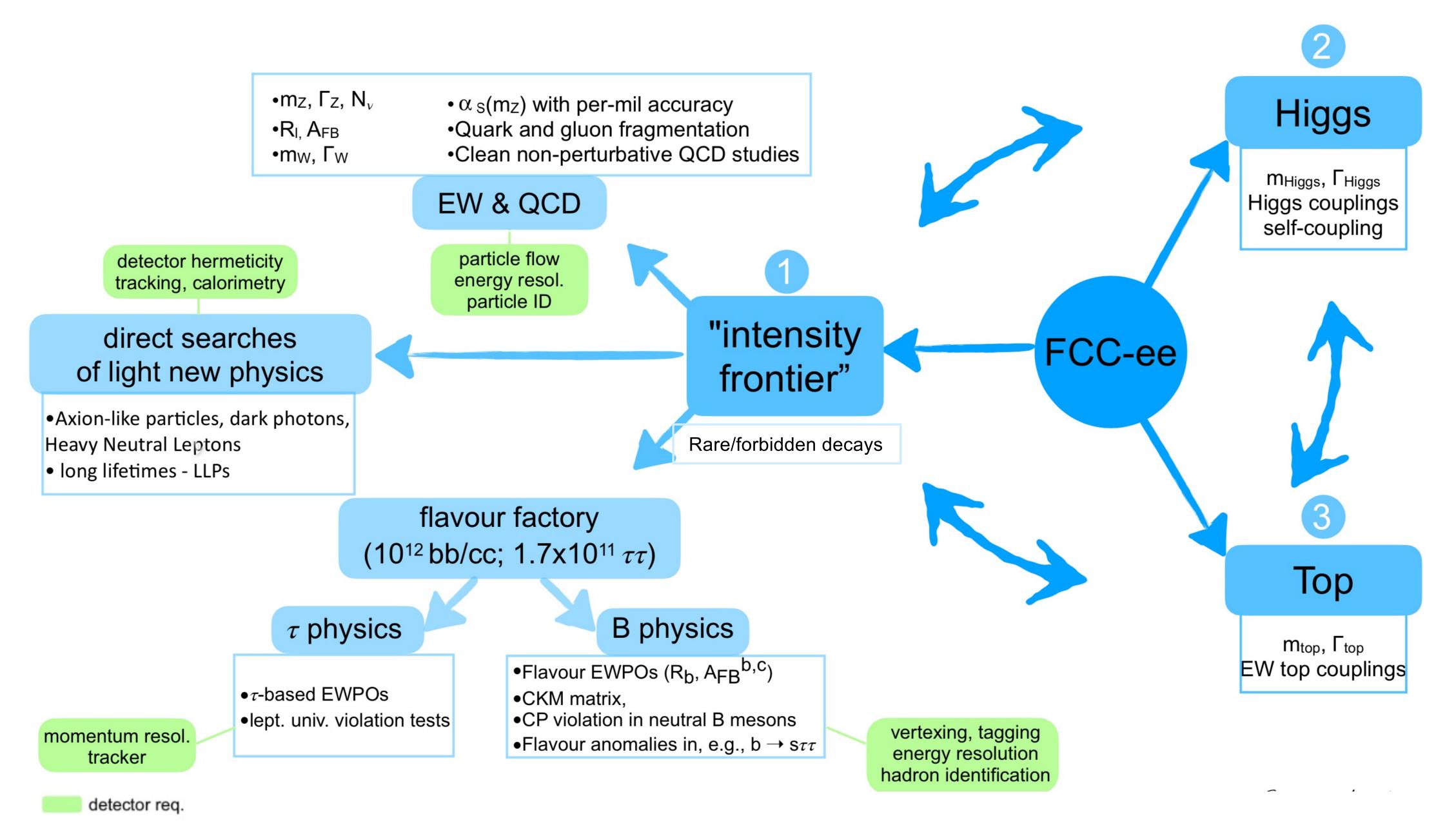
an important target that is guaranteed to be reached (no-lose theorem)

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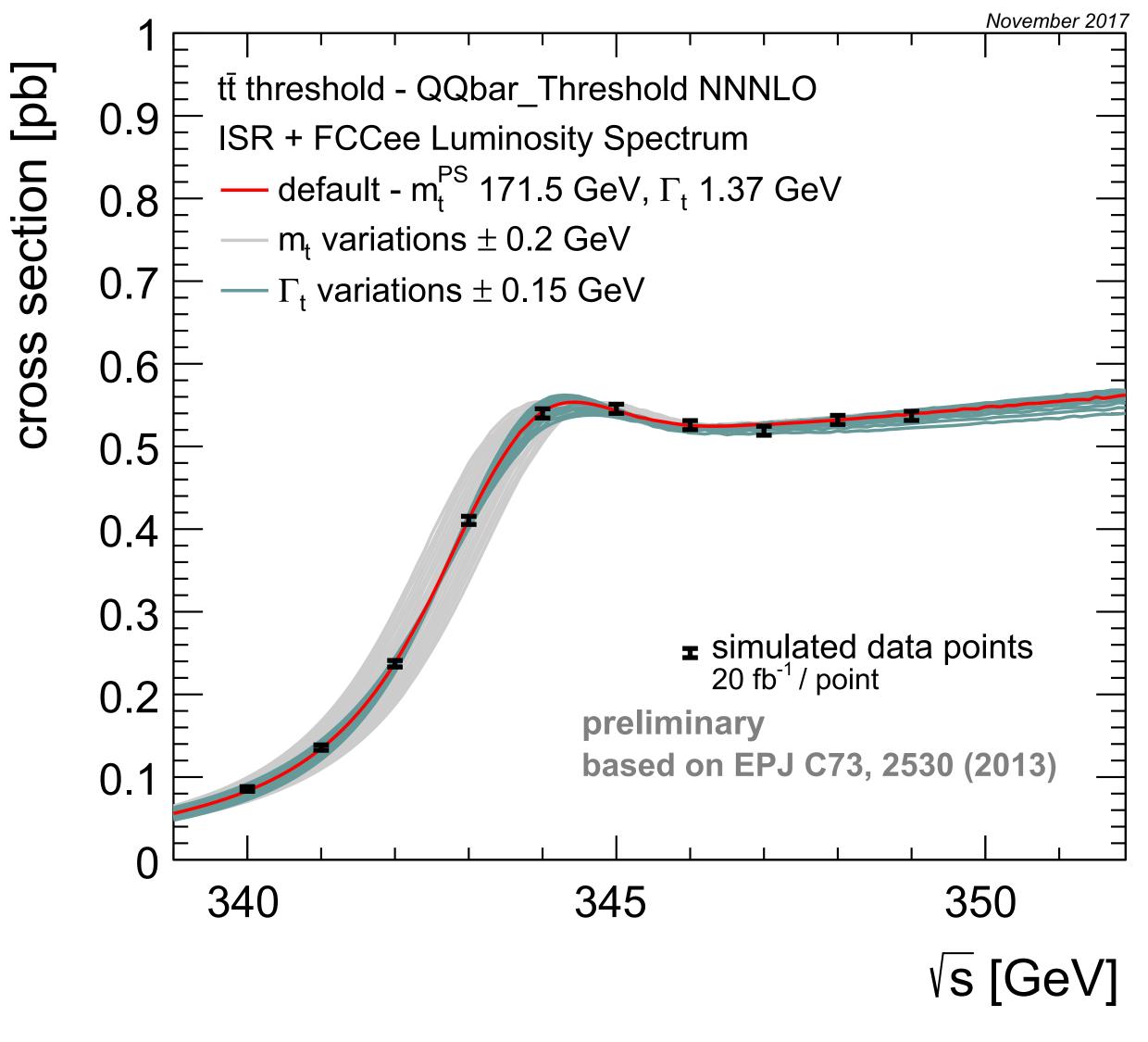
likelihood of success, robustness (incl. multiple experiments)

cost-effective construction & operation, low carbon footprint

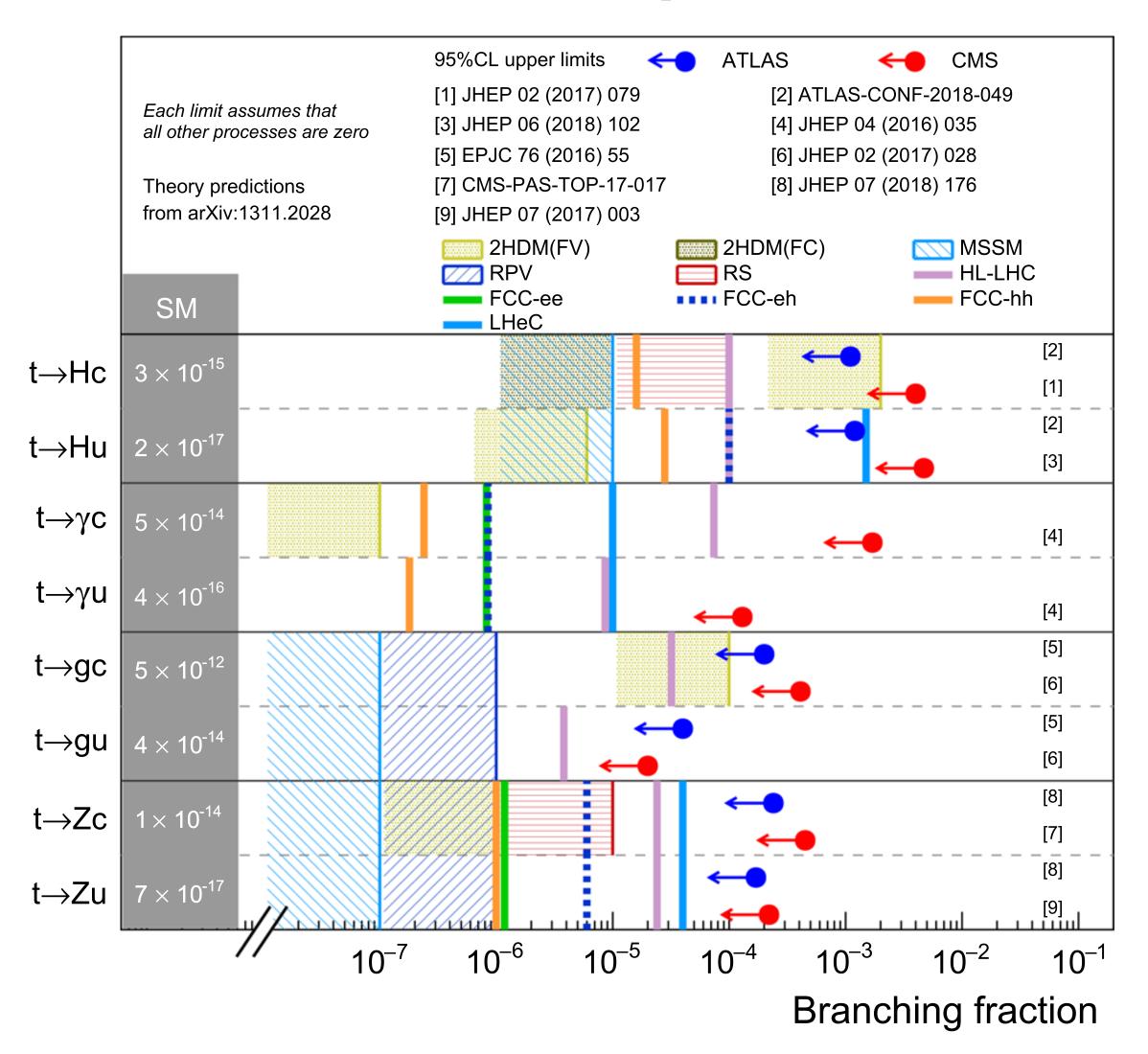


Slide from C. Grojean @ FCC Week'22

threshold scan for top mass



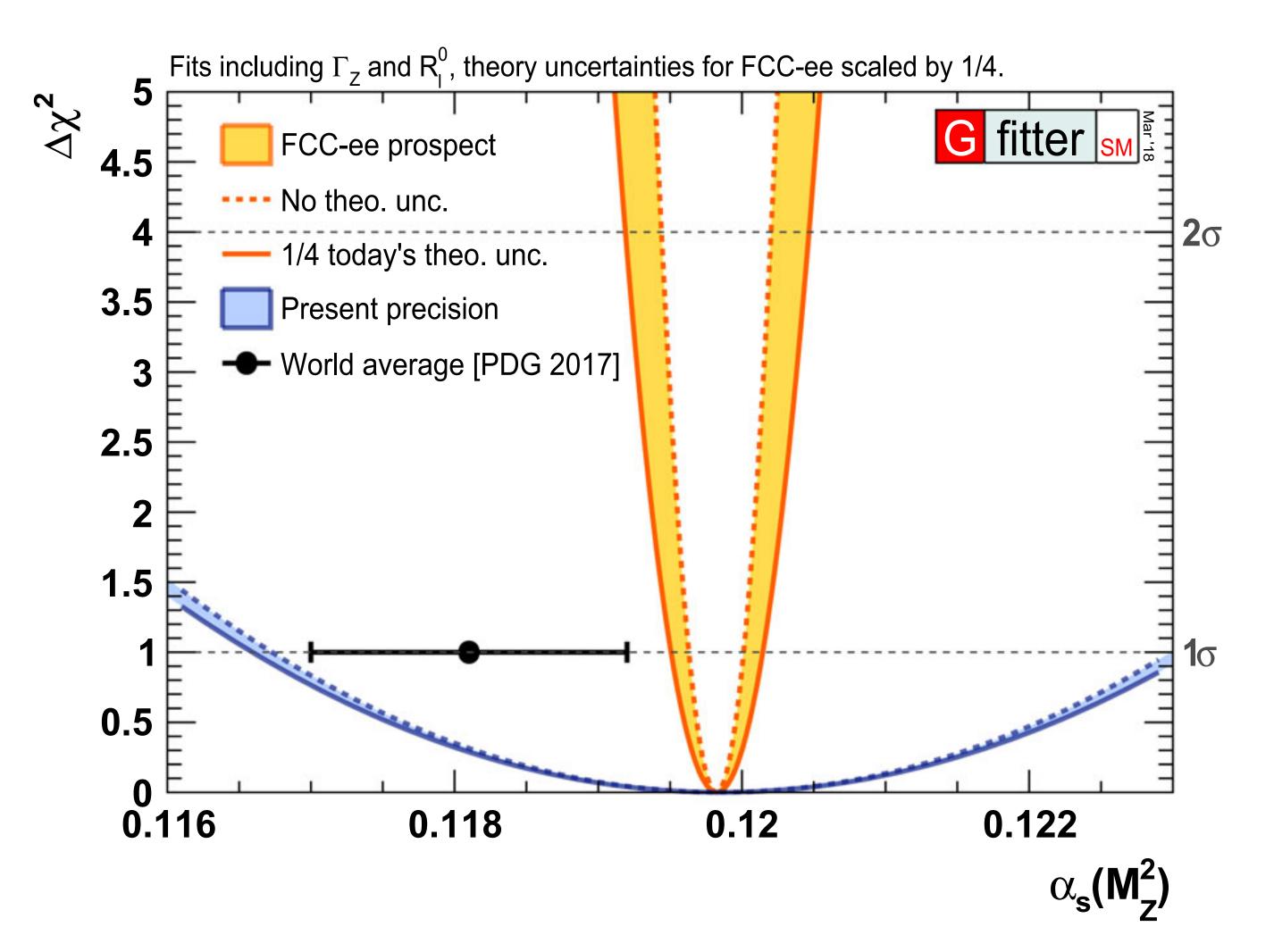
limits on top FCNF



Flavour physics: 15× more b-pairs at FCC-ee than at Belle II

2106.01259			FCC-ee
Attribute	$\Upsilon(4S)$	pp	Z^0
All hadron species			
High boost			
Enormous production cross-section			
Negligible trigger losses			
Low backgrounds			
Initial energy constraint			

FCC-ee & QCD: strong coupling, etc.

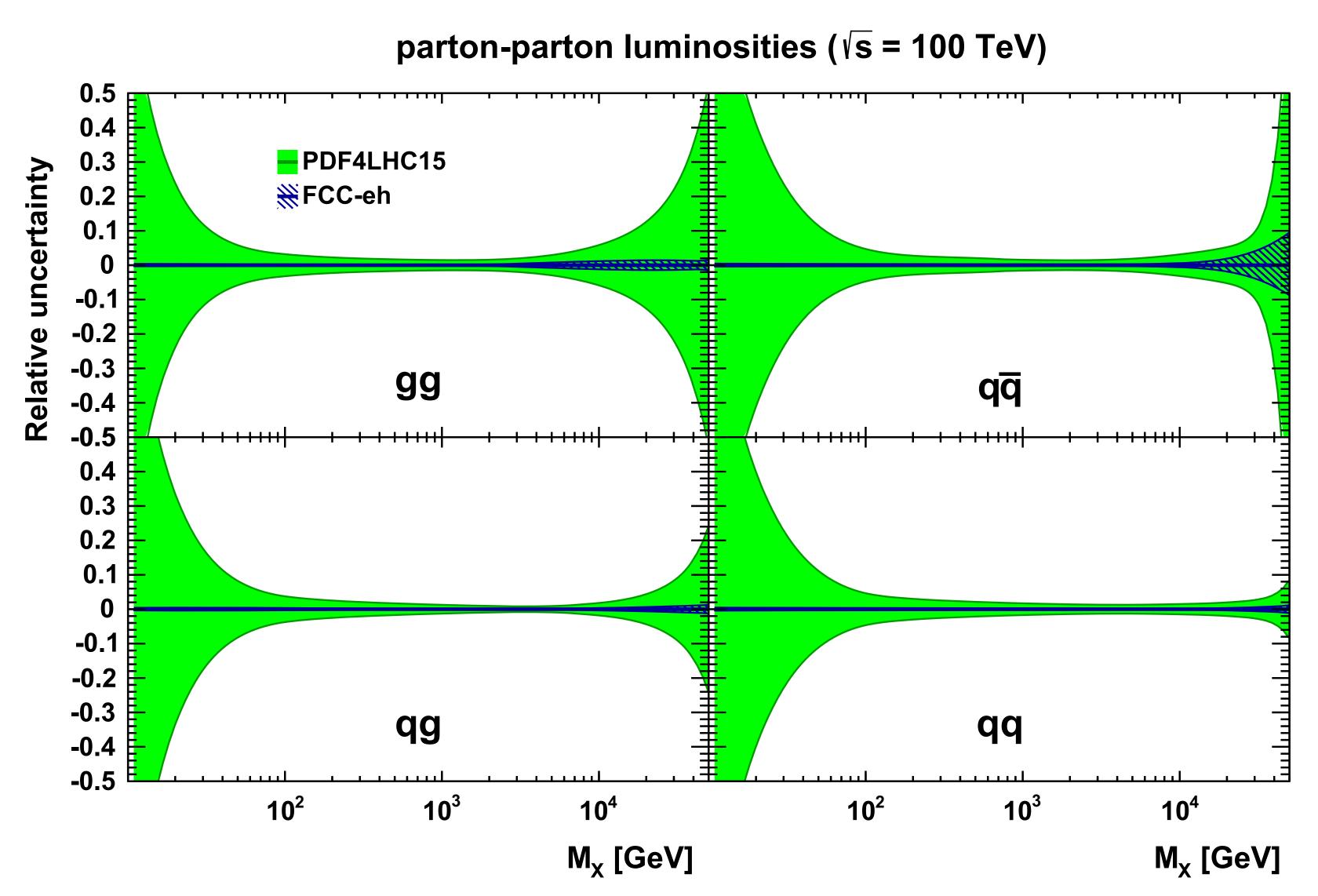


- strong coupling from EW precision to per-mil accuracy
- studies of colour reconnection in W-pair events
- ➤ jet rates, substructure, flavour, fragmentation

54

> etc.

FCC-eh: huge improvement partonic luminosities



PDFs from FCC-eh are potentially crucial for full exploitation of FCC-hh physics programme.

NB: potential worries about non-perturbative contributions in PDF fits to moderate- Q^2 DIS data

& reliance on data from single experiment

FCC-hh PbPb collisions: top & W decays probe q/g-plasma across yoctosecond time-scales

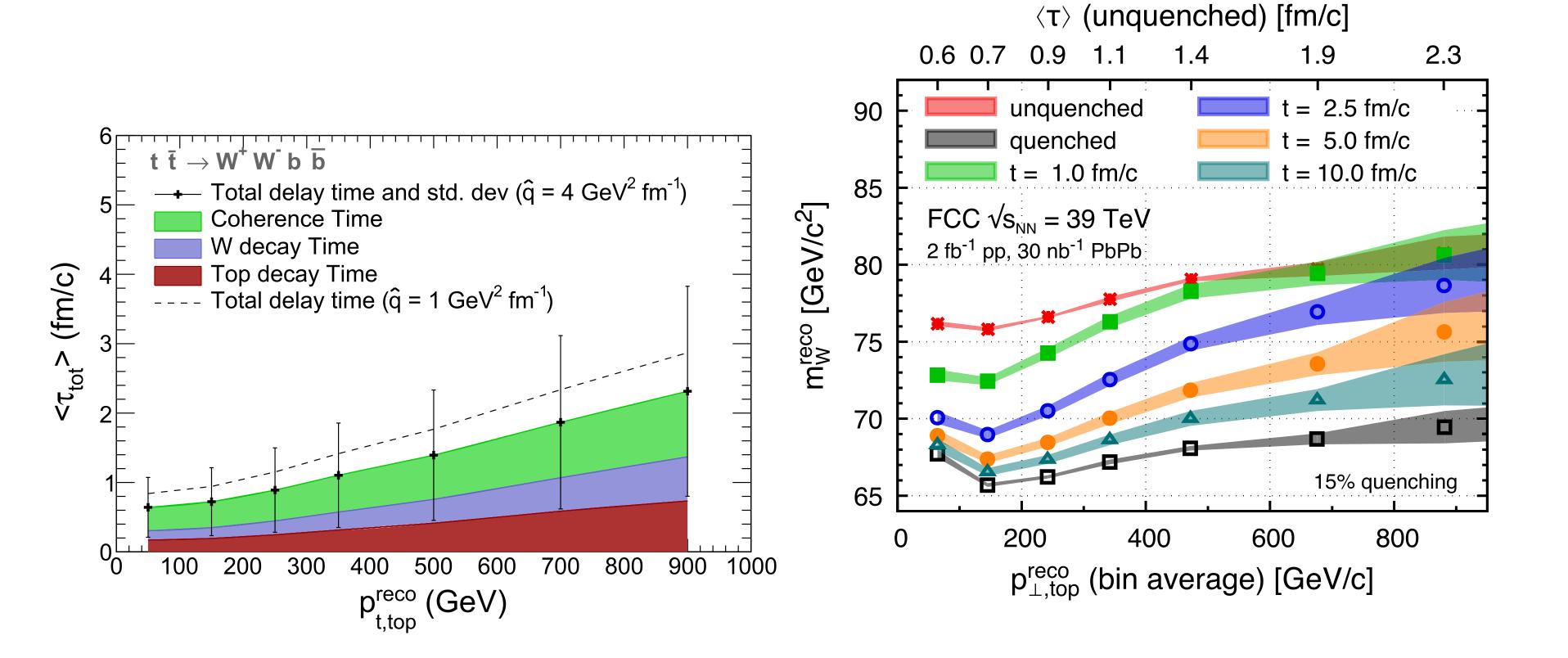


Fig. S.6 Left: total delay time for the QGP energy-loss parameter $\hat{q} = 4 \text{ GeV}^2/\text{fm}$ as a function of the top transverse momentum (black dots) and its standard deviation (error bars). The average contribution of each component is shown as a coloured stack band. The dashed line

corresponds to a $\hat{q} = 1 \text{ GeV}^2/\text{fm}$. Right: reconstructed W boson mass, as a function of the top p_T . The upper axis refers to the average total time delay of the corresponding top p_T bin

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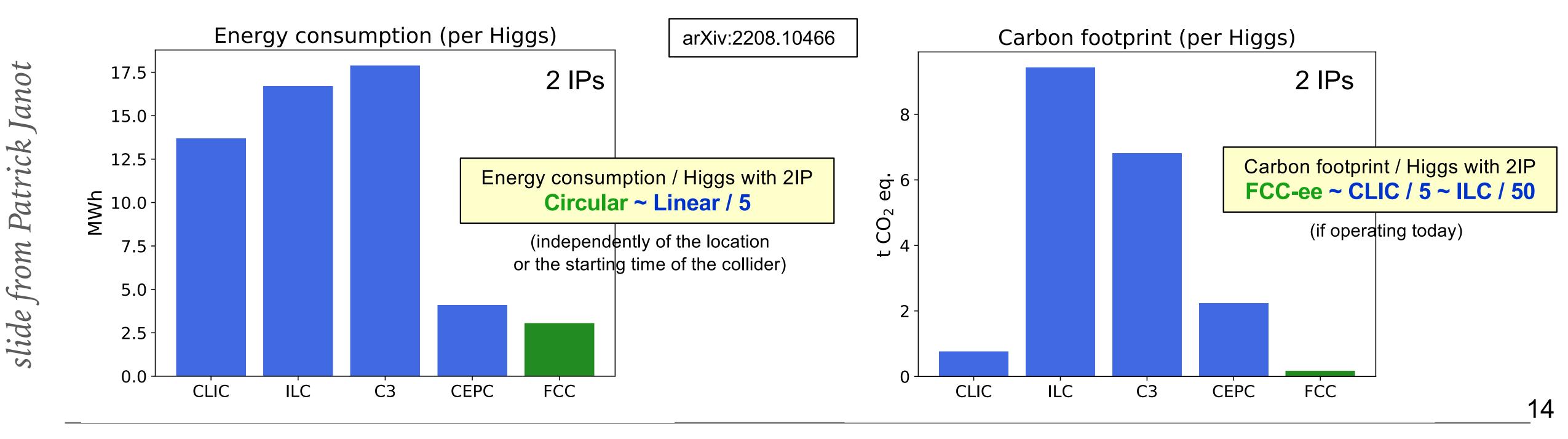
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Energy consumption and carbon footprint @ 240 GeV

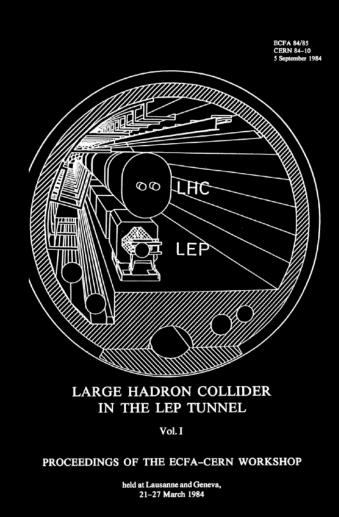
- Our first responsibility (as particle physicists) is to do the maximum of science
 - ◆ With the minimum energy consumption and the minimum environmental impact for our planet
 - Should become one of our top-level decision criteria for design, choice and optimization of a collider
- All Higgs factories have a "similar" physics outcome (ESU'20 and Snowmass'21)
 - Natural question: what is their energy consumption or carbon footprint for the same physics outcome?
 - Circular colliders have a much larger instantaneous luminosity and operate several detectors
 - FCC-ee is at CERN, where electricity is already almost carbon-free (and will be even more so in 2048)



conclusions

Conclusions

- ➤ There is a **no-lose theorem:** directly establishing Higgs self-interaction (it holds the SM together), which is made solid by precision of FCC-ee and direct measurement at FCC-hh
 - ➤ is there a chance of a second no-lose theorem in establishing (or disproving) SM origin of electron mass?
- The step up in energy reach that we expect is $\sim \times 4 5$
 - ➤ FCC-ee delivers that in "indirect" sensitivity, through precision increase ~ × 18
 - > FCC-hh delivers that in direct search sensitivity
- ➤ The programme is diverse and robust
- ➤ One issue: timeline.
 - ➤ Probably no realistic faster route to a new collider of any kind, but the field as a whole is at risk if we don't soon consolidate the path to a new collider that starts in next c. 20 years.



PHYSICS WITH A MULTI-TEV HADRON COLLIDER

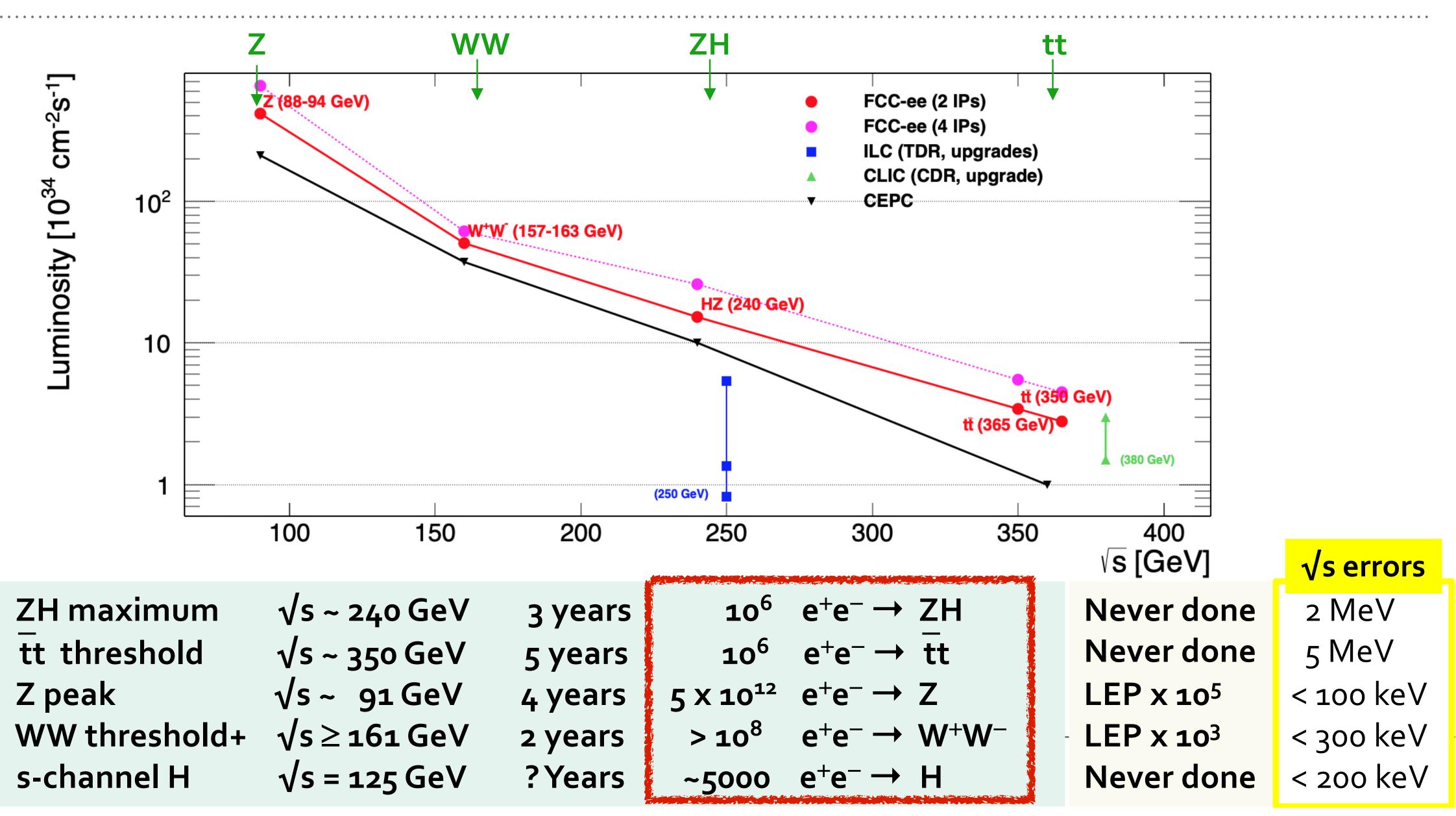
C.H. Llewellyn Smith,

Looking at the wide variety of alternatives which have been proposed, it might appear that theorists are in disarray but it seems to me that the present situation is an inevitable consequence of the successes of the 1970's. The problems of the 1960's - the nature of hadrons, the nature of the strong force, the nature of the weak force - have been solved. We now confront deeper problems - the origin of mass, the choice of fundamental building blocks (the problem of flavour), the question of further unification of forces including gravity, the origin of charge and of gauge symmetry. It is only to be expected that many of the first attempts to grapple with these problems will be misguided. As ever, we must reply on experiment to reveal the truth.

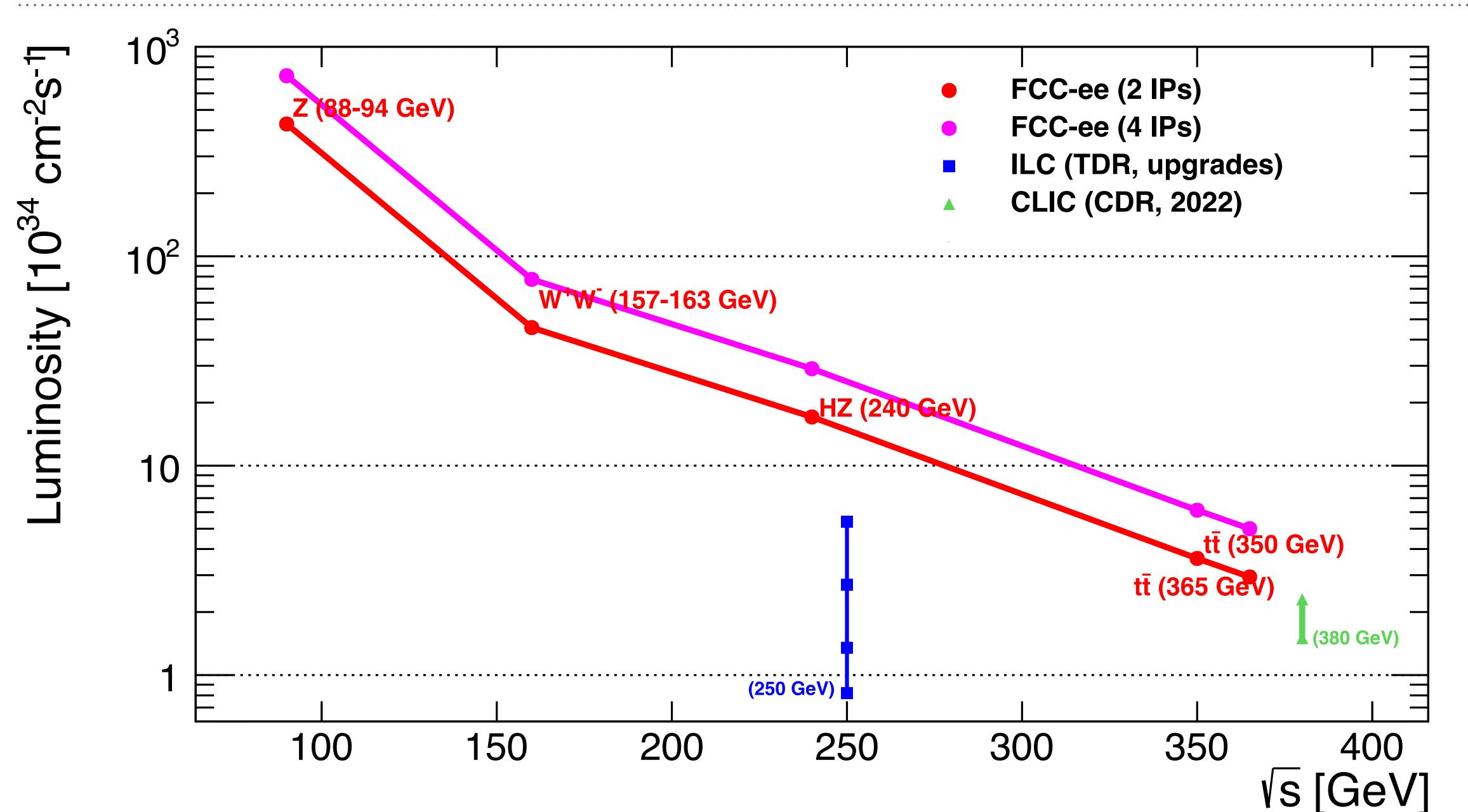
backup

Recalling the basic numbers

FCC-ee (numbers of events are for 2 detectors — baseline is now 4)



FCC-ee (updated plot for 4 detectors)



65

Gavin

FCC-hh: what do 20/30ab⁻¹ @ 100 TeV buy you?

- ➤ ~ ×5 in mass reach of new-physics searches relative to HL-LHC (fairly independently of the new physics scenario)
- ➤ 100 → 500 × higher numbers of Higgs bosons, $t\bar{t}$ pairs, etc. than HL-LHC (much more at high-p_T & for high-mass pairs)

Table 1.1. Higgs production event rates for selected processes at $100 \,\text{TeV}$ (N_{100}) and statistical increase with respect to the statistics of the HL-LHC $(N_{100} = \sigma_{100 \,\text{TeV}} \times 30 \,\text{ab}^{-1}, N_{14} = \sigma_{14 \,\text{TeV}} \times 3 \,\text{ab}^{-1})$.

	$gg \rightarrow H$	VBF	WH	ZH	$t \overline{t} H$	HH
N_{100}	24×10^9	2.1×10^9	4.6×10^{8}	3.3×10^{8}	9.6×10^{8}	3.6×10^7
N_{100}/N_{14}	180	170	100	110	530	390

together with PbPb [and maybe ep and ePb options]

NB ee numbers
are outdated
(2IP, should be 4)

	√s	L /IP (cm ⁻² s ⁻¹)	Int. L /IP(ab ⁻¹)	Comments
e ⁺ e ⁻ FCC-ee	~90 GeV Z 160 WW 240 H ~365 top	230 x10 ³⁴ 28 8.5 1.5	75 5 2.5 0.8	2-4 experiments Total ~ 15 years of operation
pp FCC-hh	100 TeV	5 x 10 ³⁴ 30	20-30	2+2 experiments Total ~ 25 years of operation
PbPb FCC-hh	√ <u>s_{NN}</u> = 39TeV	3 x 10 ²⁹	100 nb ⁻¹ /run	1 run = 1 month operation
ep Fcc-eh	3.5 TeV	1.5 10 ³⁴	2 ab ⁻¹	60 GeV e- from ERL Concurrent operation with pp for ~ 20 years
e-Pb Fcc-eh	√s _{eN} = 2.2 TeV	0.5 10 ³⁴	1 fb ⁻¹	60 GeV e- from ERL Concurrent operation with PbPb

FCC as a Higgs factory [NB numbers are for 2 IP — new baseline is 4 IP]

Higgs provides a very good reason why we need both e⁺e⁻ <u>AND</u> pp colliders

- FCC-ee measures g_{HZZ} to 0.2% (absolute, model-independent, standard candle) from σ_{7H}
 - Γ_{H} , g_{Hbb} , g_{Hcc} , $g_{H\tau\tau}$, g_{Hww} follow
 - Standard candle fixes all HL-LHC / FCC-hh couplings
- ♦ FCC-hh produces over 10¹º Higgs bosons
 - (1st standard candle \rightarrow) $g_{H\mu\mu}$, $g_{H\gamma\gamma}$, $g_{HZ\gamma}$, Br_{inv}
- ♦ FCC-ee measures top EW couplings ($e^+e^- \rightarrow t\bar{t}$)
 - Another standard candle
- ♦ FCC-hh produces 10⁸ ttH and 2. 10⁷ HH pairs
 - (2nd standard candle \rightarrow) g_{Htt} and g_{HHH}
- FCC-ee / FCC-hh complementarity is outstanding
 - Unreachable by high-energy lepton colliders

		·		
Collider	HL-LHC	$FCC-ee_{240\rightarrow365}$	FCC-INT	
Lumi (ab^{-1})	3	5+0.2+1.5	30	
Years	10	3 + 1 + 4	25	
g_{HZZ} (%)	1.5	0.18 / 0.17	0.17/0.16	
g_{HWW} (%)	1.7	0.44 / 0.41	0.20/0.19*	
g_{Hbb} (%)	5.1	0.69 / 0.64	0.48/0.48	∫ ee
$g_{ m Hcc}~(\%)$	SM	1.3 / 1.3	0.96/0.96	
$g_{ m Hgg}~(\%)$	2.5	1.0 / 0.89	0.52/0.5	
$g_{\mathrm{H} au au}$ (%)	1.9	0.74 / 0.66	0.49/0.46	
$g_{\mathrm{H}\mu\mu}$ (%)	4.4	8.9 / 3.9	0.43/0.43	
$g_{\mathrm{H}\gamma\gamma}$ (%)	1.8	3.9 / 1.2	0.32/0.32	
$g_{\mathrm{HZ}\gamma}$ (%)	11.	- / 10.	0.71/0.7	
$g_{ m Htt}$ (%)	3.4	10. / 3.1	1.0/0.95	pp
	50.	44./33.	3	
g_{HHH} (%)	50.	27./24.	3	
$\Gamma_{ m H}$ (%)	SM	1.1	0.91	ee
BR _{inv} (%)	1.9	0.19	0.024	pp
BR_{EXO} (%)	SM(0.0)	1.1	1	ee

* g_{HWW} includes also ep

FCC-ee is also the most pragmatic, safest, and most effective way toward FCC-hh

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Patrick Janot

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Observable	m _Z	2m _W	HZ max. 240-250 GeV	2M top 340-380 GeV	500 GeV	1.5 TeV	3 TeV	28 TeV 37 TeV 48 TeV	100 TeV	Leading Physics Questions
Precision EW (Z, W, top)	Transverse polarization	Transverse polarization		m_{top} ($m_{W_{m{\prime}}}$ $lpha_{S}$)						Existence of more SM- Interacting particles
QCD (α_{S}) QED (α_{QED})	5×10 ¹² Z	3×10 ⁸ W	1o⁵ H→gg							Fundamental constants and tests of QED/QCD
Model-independent Higgs couplings		→ H = m _H		d 75k WW→H energies					<1% precision (*)	Test Higgs nature
Higgs rare decays									<1% precision (*)	Portal to new physics
Higgs invisible decays									10 ⁻⁴ BR sensitivity	Portal to dark matter
Higgs self-coupling				oop corrections oss sections					3% (HH prod) (*)	Key to EWSB
Flavours (b, τ)	5×10 ¹² Z									Portal to new physics Test of symmetries
RH v's, Feebly interacting particles	5×10 ¹² Z								10 ¹¹ W	Direct NP discovery At low couplings
Direct search at high scales					M _χ <250GeV Small ΔM	M _x <750GeV Small ΔM	M _X <1.5TeV Small ∆M		Up to 40 TeV	Direct NP discovery At high mass
Precision EW at high energy							Y		W, Z	Indirect Sensitivity to Nearby new physics
Quark-gluon plasma Physics w/ injectors										QCD at origins

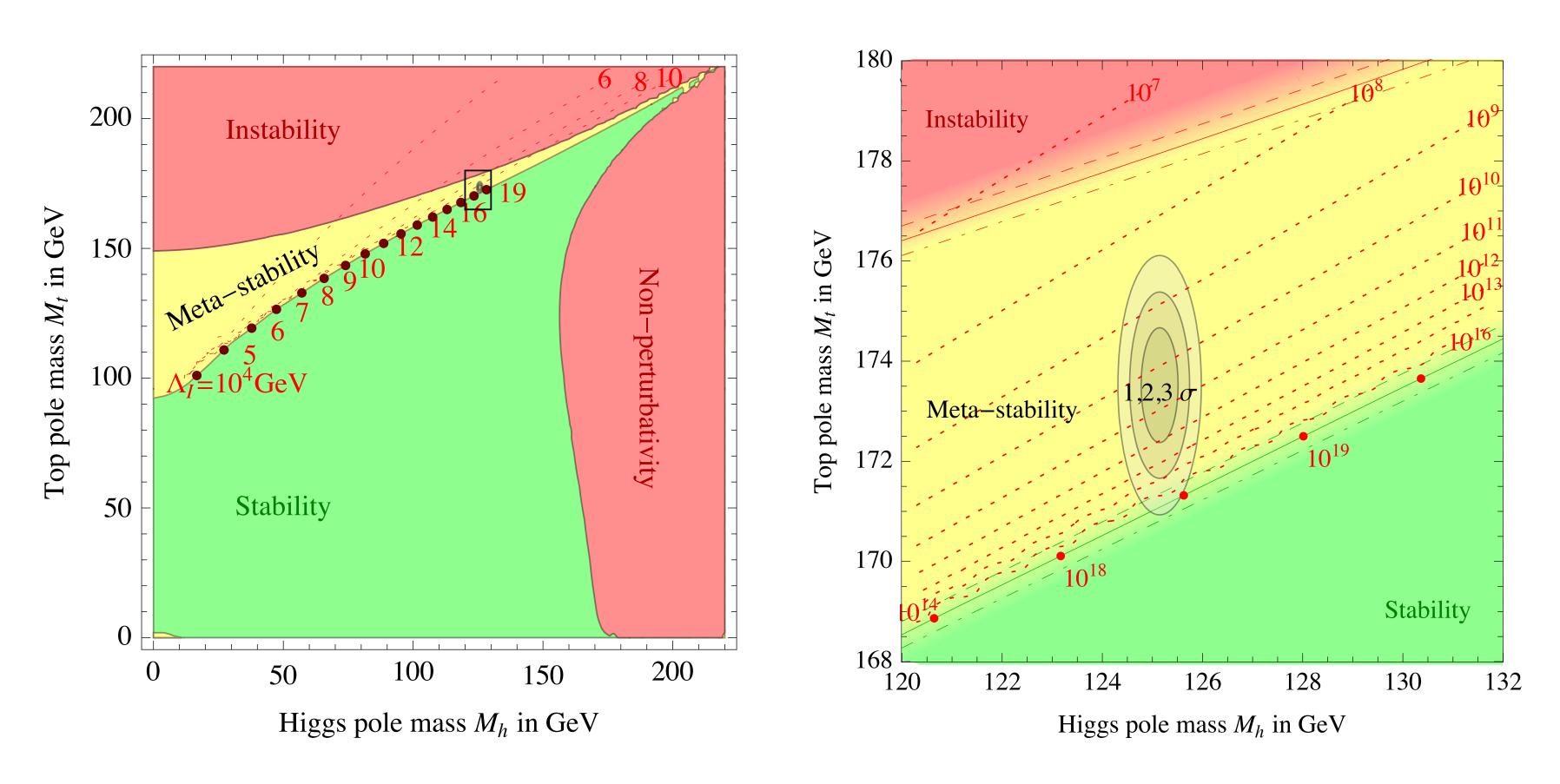
Green = Unique to FCC; Blue = Best with FCC; (*) = if FCC-hh is combined with FCC-ee; Pink = Best with other colliders;

Electroweak fits (1910.11775), e.g. S & T parameters (i.e. specific EFT operator combinations)

Table 3.3: Values for 1σ sensitivity on the S and T parameters. In all cases the value shown is after combination with HL-LHC. For ILC and CLIC the projections are shown with and without dedicated running at the Z-pole. All other oblique parameters are set to zero. The intrinsic theory uncertainty is also set to zero.

	Current	HL-LHC	ILC ₂₅₀		CEPC	FCC-ee	\mathbf{C}	LIC ₃₈₀
			(& ILC ₉₁)					(& CLIC ₉₁)
S	0.13	0.053	0.012	0.009	0.0068	0.0038	0.032	0.011
T	0.08	0.041	0.014	0.013	0.0072	0.0022	0.023	0.012

FCC-ee brings × 14-18 increase in precision



It's not inconceivable that the top mass could be sufficiently mis-measured at hadron colliders that the SM-universe is stable all the way to the Planck scale

condition in terms of the pole top mass. We can express the stability condition of eq. (64) as

$$M_t < (171.53 \pm 0.15 \pm 0.23_{\alpha_3} \pm 0.15_{M_h}) \,\text{GeV} = (171.53 \pm 0.42) \,\text{GeV}.$$
 (66)

arXiv:1307.3536

muon colliders

Higgs at muon collider

Table 6: 68% probability sensitivity to the Higgs couplings, assuming no BSM Higgs decay chann

	HL-LHC	HL-LHC	HL-LHC	HL-LHC	HL-LHC
		$+~125~{ m GeV~MuC}$	+ 3 TeV MuC	+ 10 TeV MuC	$ \hspace{.1cm} + \hspace{.1cm} 10 \hspace{.1cm} \mathrm{TeV} \hspace{.1cm} \mathrm{MuC} \hspace{.1cm} $
Coupling		$5 / 20 \text{ fb}^{-1}$	$1/2 \text{ ab}^{-1}$	10 ab^{-1}	+ FCC-ee
$\kappa_W [\%]$	1.7	1.3~/~0.9	$\boxed{0.4\ /\ 0.3}$	0.1	0.1
$\kappa_Z [\%]$	1.5	$1.3 \ / \ 1.0$	$0.9 \; / \; 0.7$	0.4	0.1
$\kappa_g [\%]$	2.3	$1.7\ /\ 1.4$	$1.2 \; / \; 1.0$	0.7	0.6
$\kappa_{\gamma} [\%]$	1.9	$1.6\ /\ 1.5$	$1.3 \; / \; 1.2$	0.8	0.8
$\kappa_{Z\gamma} [\%]$	10	10 / 10	$9.3 \; / \; 8.6$	7.2	7.1
$\kappa_c ~[\%]$	-	$12\ /\ 5.9$	$igg 6.2 \ / \ 4.4$	2.3	1.1
$\kappa_b [\%]$	3.6	$1.6 \; / \; 1.0$	$0.8 \; / \; 0.7$	0.4	0.4
$\kappa_{\mu} [\%]$	4.6	$0.6 \ / \ 0.3$	$igg 4.2 \ / \ 4.0$	3.4	3.2
$\kappa_{ au} [\%]$	1.9	$1.4\ /\ 1.2$	$1.2 \; / \; 1.0$	0.6	0.4
$\kappa_t^\dagger~[\%]$	3.3	$3.2 \; / \; 3.1$	$igg 3.1 \; / \; 3.1$	3.1	3.1
Γ_H^{\ddagger} [%]	5.3	$2.7 \; / \; 1.7$	$1.3 \ / \ 1.0$	0.5	0.4

[†] No input used for μ collider.

Table 7: 68% probability intervals for the Higgs trilinear coupling.

	HL-LHC	3 TeV MuC $\text{L} \approx 1 \text{ ab}^{-1} / 2 \text{ ab}^{-1}$	10 TeV MuC $L=10 \text{ ab}^{-1}$	14 TeV MuC $L \approx 20 \text{ ab}^{-1}$	30 TeV MuC $L=90 \text{ ab}^{-1}$
$\delta \kappa_{\lambda}$	[-0.5, 0.5]	$[-0.27, 0.35] \cup [0.85, 0.94] / [-0.15, 0.16]$	[-0.035, 0.037]	[-0.024, 0.025]	[-0.011, 0.012]
comb. w HL-LHC		[-0.2,0.22] / [-0.13,0.14]	[-0.035,0.036]	[-0.024,0.025]	[-0.011,0.012]

[‡] Prediction assuming only SM Higgs decay channels. Not a free parameter in the fits.

triple Higgs at muon collider from 2003.13628

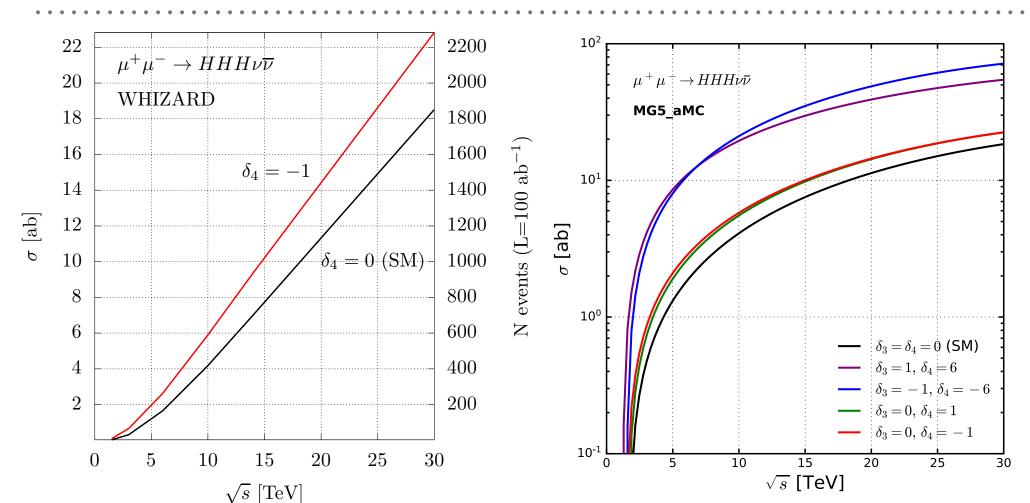


Figure 2: Expected cross sections (left) and signal event numbers for a reference integrated luminosity of 100 ab⁻¹ (right) for $\mu^+\mu^- \to HHH\nu\bar{\nu}$ versus the c.m. collision energy, for $M_{\bar{\nu}\nu} \gtrsim 150 \text{GeV}$. Cross sections for different assumptions of the trilinear and quartic couplings are presented, as well as for the SM case, obtained by WHIZARD (left-hand side) and MAD-GRAPH5_AMC@NLO (right-hand side). Details on the scenarios are given in the text.

			Constraints on	$\delta_4 \text{ (with } \delta_3 = 0)$
\sqrt{s} (TeV)	Lumi (ab^{-1})	x-sec only	x-sec only	threshold + $M_{HHH} > 1$ TeV
		1σ	2σ	1σ
6	12	[-0.60, 0.75]	[-0.90, 1.00]	[-0.55, 0.85]
10	20	[-0.50, 0.55]	[-0.70, 0.80]	[-0.45, 0.70]
14	33	[-0.45, 0.50]	[-0.60, 0.65]	[-0.35, 0.55]
30	100	[-0.30, 0.35]	[-0.45, 0.45]	[-0.20, 0.40]
3	100	[-0.35, 0.60]	[-0.50, 0.80]	[-0.45, 0.65]

Table 5: Summary of the constraints on the quartic deviations δ_4 , assuming $\delta_3 = 0$, for various muon collider energy/luminosity options, as obtained from the total expected cross sections (1 σ and 2 σ CL). The third column shows the bounds obtained from the combination of the constraints corresponding to the setups $M_{HHH} < 1$ TeV and $M_{HHH} > 1$ TeV.

Searches at muon collider

Plots being shown suggest: 4 TeV muon collider beats a 100 TeV pp collider in searches for new physics.

Useful to nuance the statement:

- ➤ 100 TeV pp, 20 ab⁻¹ can discover Z' up to $m_{Z'} \sim 38$ TeV
- For $\mu\mu$ collider to discover Z' at $m_{Z'} \sim 38$ TeV, it needs $\sqrt{s} \sim 38$ TeV (with lower \sqrt{s} you $\sqrt{s_{\mu}}$ [TeV] would see deviation from SM, but not know what it is)
- \blacktriangleright However a 38 TeV muon collider would be much better at studying the Z' than the 100 TeV pp machine

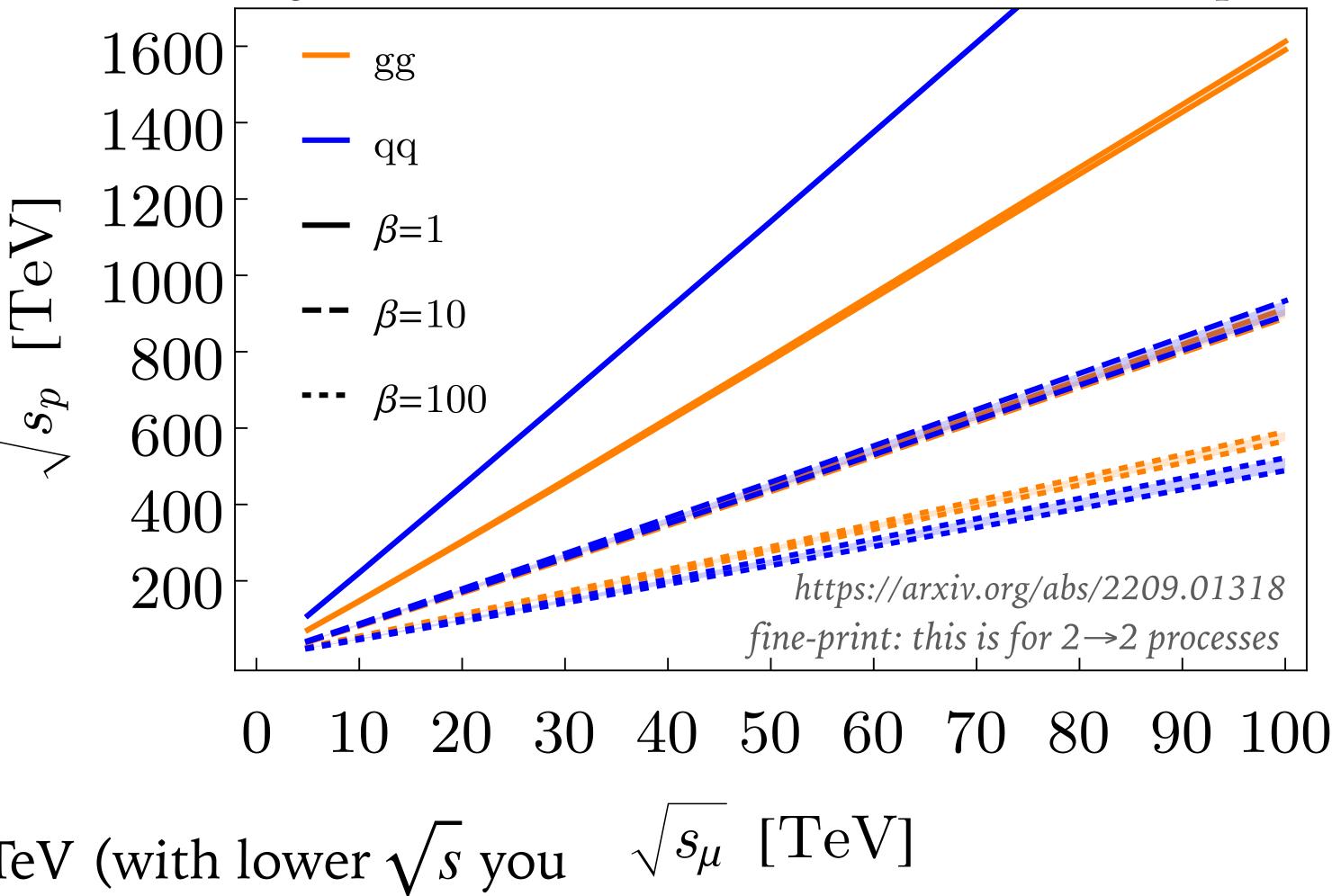
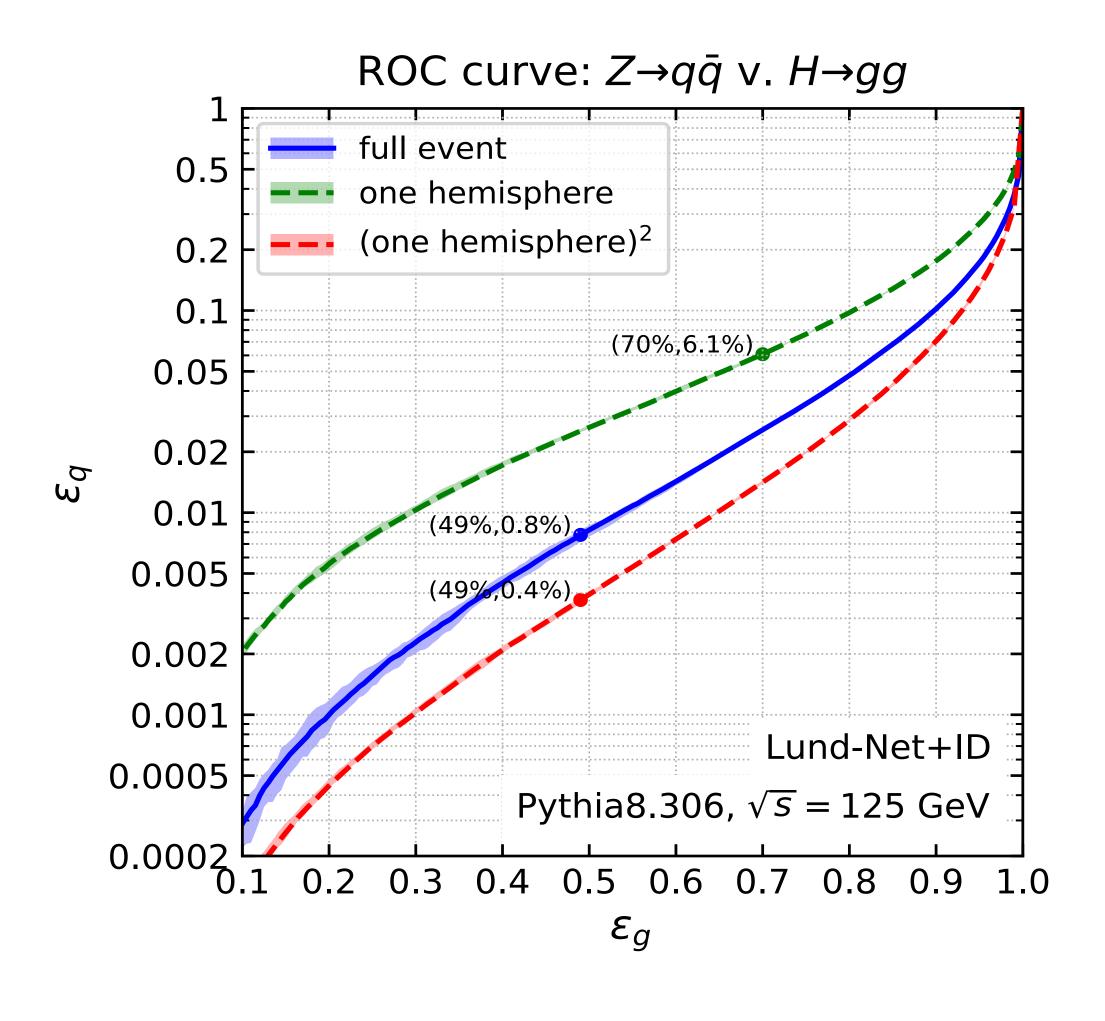


Fig. 3 of Snowmass Muon Collider Forum Report

$H \rightarrow gg$ at FCC-ee

$$e^+e^-
ightarrow Z
ightarrow q \overline{q}$$
 v. $e^+e^-
ightarrow H
ightarrow gg$ $(\sqrt{s}=125$ GeV, no ISR)



Observed performance:

- per jet: 6% quark mistag
 for 70% gluon efficiency
 Not quite the 1% quark mistag in 2107.02686
- full event: 0.8% quark mistag for 49% gluon efficiency full event worse that (jet)²