13th ICFA Seminar on **Future Perspectives in High-Energy Physics** 

28 November – 1 December 2023 DESY, Hamburg

# The future of HEP: our motivation

Gavin Salam University of Oxford & All Souls College



THE ROYAL SOCIETY





Science and Technology **Facilities Council** 



# Medium/large projects: community knows how to motivate and get them funded





DUNE, HK, JUNO, and neutrino observatories will enable a bona fide precision physics program in the neutrino sector

The Electron Ion Collider



## Status of WIMP Searches: from the sky and underground

Jianglai Liu





# <u>top-down</u>

# figure out the best collider you can realistically build

establish what physics it will probe

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# <u>bottom up</u>

establish what you want to learn

figure out how to build a collider that will best achieve it





Dear Santa Claus,

We have been good these past decades. Please could you now bring us

- a dark matter candidate
- an explanation for the fermion masses
- an explanation of matter-antimatter asymmetry
- an axion, to solve the strong CP problem
- a solution to fine tuning the EW scale
- a solution to fine tuning the cosmological constant

Thank you, Particle Physicists

ps: please, no anthropics

# we have so far been unlucky in getting answers to these many questions





# Snowmass Dark Matter report, <u>2209.07426</u>



# **30 orders** of magnitude in interaction strength

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# magnitude in mass





# the standard-model particle set is complete







#### PORTUGIESENVIERTEL

### Baumwall (Elbphilharmonie)

# ap San Diego 🔟



# Hamburger

# Am Sanatorna

# Elbphilharmonie Hamburg

**a** 

Nordera

# Grasbrookhafen

Strandhafen







# the standard-model particle set is complete

but we have been lucky with the Higgs boson's 125 GeV mass it opens a door to the most mysterious part of the Standard Model





# desirable features of the next major HEP project(s)?

# an important target to be reached $\sim$ guaranteed discovery

exploration into the unknown by a significant factor in energy

likelihood of success, robustness (e.g. multiple experiments)

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major progress on a broad array of particle physics topics

cost-effective construction & operation, low carbon footprint, novel technologies

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# Higgs physics Higgs is the last particle of the SM. So the SM is complete, right?

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# The Lagrangian and Higgs interactions: two out of three qualitatively new!

 $\mathscr{L}_{SM} = \cdots + |D_{\mu}\phi|^2 + \psi_i y_{ij}\psi_j\phi -$ 

Gauge interactions, structurally like those in QED, QCD, EW, studied for many decades (but now with a scalar)

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Yukawa interactions. Responsible for fermion masses, and induces "fifth force" between fermions. Direct study started only in 2018!

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Higgs potential  $\rightarrow$ self-interaction ("sixth?" force between scalars). Holds the SM together. Unobserved





# Almost every problem of the Standard Model originates from Higgs interactions

# $\mathscr{L} = y H \psi \bar{\psi} + \mu^2 |H|^2 - \lambda |H|^4 - V_0$ stability naturalness flavour cosmological constant

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typeset from Gian Giudice original









#### **Thermal History of** Universe

#### **Naturalness**

#### **Fundamental** or Composite?

#### Is it unique?

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# Yukawa interaction hypothesis

Yukawa couplings ~ fermion mass

first fundamental interaction that we probe at the quantum level where interaction strength (y<sub>ii</sub>) not quantised (i.e. no underlying unit of conserved charge across particles)





# Protons are lighter than neutrons $\rightarrow$ protons are stable. Giving us the hydrogen atom, & chemistry and biology as we know it

### $\simeq 938.3$ MeV

## $\simeq 939.6$ MeV







Protons are **lighter** than neutrons  $\rightarrow$  protons are stable. Giving us the hydrogen atom, & chemistry and biology as we know it Supposedly because up quarks interact more weakly with the Higgs field than down quarks

~ 938.3 MeV

 $\simeq 939.6$  MeV





# Why do Yukawa couplings matter? (2) Because, within SM conjecture, they're what give masses to all leptons



# electron mass determines size of all atoms

it sets energy levels of all chemical reactions

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# currently we have no evidence that up and down quarks and electron get their masses from Yukawa interactions — it's in textbooks, but is it nature?

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![](_page_20_Figure_0.jpeg)

to be conclusively established at the LHC within 3 – 10 years

![](_page_20_Picture_2.jpeg)

![](_page_21_Figure_0.jpeg)

to be conclusively established at the LHC within 3 – 10 years

![](_page_21_Picture_2.jpeg)

![](_page_22_Figure_0.jpeg)

![](_page_23_Figure_0.jpeg)

# A side comment on the near future at LHC

- with the world as we experience it
- > LHC will reach 5 $\sigma$  sensitivity for  $H \rightarrow \mu\mu$  in the coming years (if it is SM-like), Yukawa mechanism
- that will be a crucial step on the way from 3rd generation Yukawas to 1st
- it deserves a big event with the world's press to announce it
- fundamental particles that we are made of

> particle physics normally deals with esoteric particles that have [almost] no relation

offering first proof that particles other than 3rd generation also get their mass from

> an opportunity to explain the quest for understanding the origin of the mass of the

![](_page_25_Picture_0.jpeg)

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# the Higgs potential

![](_page_25_Picture_5.jpeg)

# **Higgs potential**

![](_page_26_Figure_1.jpeg)

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$$|\phi|^4 + V_0$$

**Standard Model** 

the Higgs mechanism gives mass to particles because the Higgs field  $\phi$  is non-zero That happens because the minimum of the SM potential is at non-zero φ

![](_page_26_Picture_7.jpeg)

# Higgs potential

![](_page_27_Figure_1.jpeg)

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φ

# **Standard Model**

depth is 
$$\frac{m_H^2 v^2}{8} (m_H \simeq 125 \text{ GeV}, v \simeq 246 \text{ GeV})^4$$
  
a fairly innocuous sounding  $(104 \text{ GeV})^4$ 

![](_page_27_Picture_6.jpeg)

![](_page_27_Picture_7.jpeg)

# Higgs potential – remember: it's an energy density

# $V(\phi)$ , SM

![](_page_28_Figure_2.jpeg)

# **Standard Model**

Corresponds to an energy density of  $1.5 \times 10^{10} \, \text{GeV/fm}^3$ i.e. >10 billion times nuclear density Mass density of  $2.6 \times 10^{28} \text{ kg/m}^3$ 

![](_page_28_Picture_8.jpeg)

![](_page_28_Picture_9.jpeg)

![](_page_29_Picture_0.jpeg)

# Earth at neutron star density

https://en.wikipedia.org/wiki/Globe#/media/File:World\_Glob https://en.wikipedia.org/wiki/Old fashioned glass#/media/File:Old Fashioned https://de.wikipedia.org/wiki/Volksparkstadion#/media/Datei:RK\_1009\_9831\_Volkspark

<u>e Map.j</u>	ipg
<u>Glass.</u>	ipg
<u>tadion.j</u>	ipg

![](_page_30_Picture_0.jpeg)

# Earth at neutron star density

https://en.wikipedia.org/wiki/Globe#/media/File:World\_Glob https://en.wikipedia.org/wiki/Old fashioned glass#/media/File:Old Fashioned https://de.wikipedia.org/wiki/Volksparkstadion#/media/Datei:RK\_1009\_9831\_Volkspark

![](_page_30_Picture_3.jpeg)

![](_page_30_Picture_4.jpeg)

![](_page_30_Picture_5.jpeg)

<u>e Map.j</u>	pg
Glass.	<u>pg</u>
<u>tadion.j</u>	ipg

# cosmological constant & fine-tuning [classically]

$$V_{min} = \begin{bmatrix} -\mu^{2} |\phi|^{2} + \lambda |\phi|^{4} \end{bmatrix}_{\phi_{0}} + V_{0}$$
  
=  $-2.6 \times 10^{28} kg/m^{3} + V_{0} = \begin{bmatrix} 5.96 \times 10^{-27} kg/m^{3} \end{bmatrix}$ 

- $\succ$  V<sub>0</sub> needs to be fine tuned for cosmological constant to have today's size (also with respect to various sources of quantum correction)
- not the only fine-tuning problem in fundamental physics, -- arguably special in that it appears already classically
- $\blacktriangleright$  collider physics cannot tell us anything about  $V_0$ — but it would seem negligent not to try and establish the rest of the potential

# The potential expanded around the minimum

 $\blacktriangleright$  take *h* as the Higgs field excitation in units of the field at minimum

$$V = \frac{m_H^2 v^2}{8} \left( -1 + 4h \right)$$

the Higgs boson mass term

# prediction of the strength of HHH interaction [modifier may be called $\kappa_{\lambda}$ or $\kappa_3$ ]

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![](_page_32_Figure_7.jpeg)

![](_page_32_Picture_8.jpeg)

- $\blacktriangleright$  equivalent for an interaction is a bit ambiguous but better than  $\pm 20\%$ determination is probably a reasonable target
- ➤ for something of this importance, we may be wary of relying on 20% only from a combination of N experiments — a result's robustness comes from confirmation by independent experiments
- ► indirect v. direct:
  - > all measurements are indirect (we measure hadrons and leptons...)
  - single H is good to have
  - > but HH & kinematic structure brings assurance that what we are seeing is indeed HHH coupling

► NB there exist different points of view on this

when would we claim discovery? [5o in each of two independent experiments is our gold standard]

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![](_page_33_Picture_13.jpeg)

- > equivalent for an interaction is a bit ambiguous but better than  $\pm 20\%$ determination is probably a reasonable target
- ► for something of this importance, we may be combination of N experiments independent experim
- ► indirec
  - ► all me
  - ► single
  - HHH coupling

► NB there exist different points of view on this

![](_page_34_Figure_10.jpeg)

> but HH & kinematic structure brings assurance that what we are seeing is indeed

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![](_page_34_Picture_14.jpeg)

# **Higgs potential**

![](_page_35_Figure_1.jpeg)

# **Standard Model**

### $\succ$ this is a cartoon

- ► caution needed: e.g. realistic BSM models do not just modify the potential, but may bring extra scalars (often modify other couplings, but not always, e.g. <u>2209.00666</u>)
- even if we take the picture seriously we may want to consider impact of limited constraints on  $\lambda_4$

![](_page_35_Picture_9.jpeg)

![](_page_35_Picture_21.jpeg)

# **Higgs potential**

![](_page_36_Figure_1.jpeg)

# **Standard Model**

# know today $-0.4 < \lambda_3 / SM < 6.3$

### $\succ$ this is a cartoon

- ► caution needed: e.g. realistic BSM models do not just modify the potential, but may bring extra scalars (often modify other couplings, but not always, e.g. <u>2209.00666</u>)
- even if we take the picture seriously we may want to consider impact of limited constraints on  $\lambda_4$

![](_page_36_Picture_10.jpeg)

![](_page_36_Picture_22.jpeg)

![](_page_37_Figure_0.jpeg)

# **Standard Model**

### what we may know in 2040 $0.5 < \lambda_3 / SM < 1.6$

### $\succ$ this is a cartoon

- ► caution needed: e.g. realistic BSM models do not just modify the potential, but may bring extra scalars (often modify other couplings, but not always, e.g. <u>2209.00666</u>)
- even if we take the picture seriously we may want to consider impact of limited constraints on  $\lambda_4$

![](_page_37_Picture_9.jpeg)

![](_page_37_Picture_21.jpeg)

# **Higgs potential**

# V(φ), 2060 (FCC-ee, 4IP)

![](_page_38_Figure_2.jpeg)

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# **Standard Model**

### what we may know in 2060 $0.76 < \lambda_3 / SM < 1.24$

### $\succ$ this is a cartoon

- ► caution needed: e.g. realistic BSM models do not just modify the potential, but may bring extra scalars (often modify other couplings, but not always, e.g. <u>2209.00666</u>)
- even if we take the picture seriously we may want to consider impact of limited constraints on  $\lambda_4$

![](_page_38_Picture_12.jpeg)

![](_page_38_Picture_24.jpeg)

# **Higgs potential**

# V(φ), 2080 (FCC-hh)

![](_page_39_Figure_2.jpeg)

# **Standard Model**

### what we may know in 2080 $0.97 < \lambda_3 / SM < 1.03$

### $\succ$ this is a cartoon

- ► caution needed: e.g. realistic BSM models do not just modify the potential, but may bring extra scalars (often modify other couplings, but not always, e.g. <u>2209.00666</u>)
- even if we take the picture seriously we may want to consider impact of limited constraints on  $\lambda_4$

![](_page_39_Picture_11.jpeg)

![](_page_39_Picture_23.jpeg)

![](_page_40_Figure_0.jpeg)

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# **Standard Model**

### what we may know in 2080

 $0.97 < \lambda_3 / SM < 1.03$  $-1 < \lambda_4 / SM < 6.5$ 

### $\succ$ this is a cartoon

- ► caution needed: e.g. realistic BSM models do not just modify the potential, but may bring extra scalars (often modify other couplings, but not always, e.g. <u>2209.00666</u>)
- even if we take the picture seriously we may want to consider impact of limited constraints on  $\lambda_4$

![](_page_40_Picture_12.jpeg)

![](_page_40_Picture_24.jpeg)

# desirable features of the next major HEP project(s)?

an important target to be reached  $\sim$  guaranteed discovery

exploration into the unknown by a significant factor in energy

major progress on a broad array of particle physics topics

likelihood of success, robustness (e.g. multiple experiments)

cost-effective construction & operation, low carbon footprint, novel technologies

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![](_page_41_Picture_8.jpeg)

![](_page_41_Picture_9.jpeg)

![](_page_41_Picture_10.jpeg)

# **mw measurements**

![](_page_42_Figure_1.jpeg)

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![](_page_42_Picture_5.jpeg)

# we don't know the precision limit of hadron colliders — but we may be close to reaching it

![](_page_43_Figure_1.jpeg)

### gg-lumi, ratio to PDF4LHC15 @ m<sub>H</sub>

PDF4LHC15	1.0000	$\pm$	0.0184	
PDF4LHC21	0.9930	±	0.0155	
CT18	0.9914	$\pm$	0.0180	_ ×3
MSHT20	0.9930	$\pm$	0.0108	
NNPDF40	0.9986	$\pm$	0.0058	

Parton Distribution Functions are one of several elements that may limit LHC/FCC-hh precision:

- essential for hadron-collider interpretation
- > PDF fits are complex, e.g. involve (sometimes inconsistent) data, some of it close to nonperturbative scale
- only partial understanding of their limits

![](_page_43_Picture_11.jpeg)

![](_page_43_Figure_12.jpeg)

![](_page_43_Picture_13.jpeg)

![](_page_43_Picture_14.jpeg)

# we don't know the precision limit of hadron colliders — but we may be close to reaching it

![](_page_44_Figure_1.jpeg)

- perturbative scale
- only partial understanding of their limits

![](_page_44_Picture_10.jpeg)

![](_page_44_Figure_11.jpeg)

![](_page_44_Picture_12.jpeg)

![](_page_44_Picture_13.jpeg)

# desirable features of a worldwide HEP project?

- an important target that is guaranteed to be reached (no-lose theorem)
- exploration into the unknown by a significant factor in energy
  - major progress on a broad array of particle physics topics
- likelihood of success, robustness (incl. multiple experiments)
- cost-effective construction & operation, low carbon footprint

![](_page_45_Picture_14.jpeg)

![](_page_45_Picture_16.jpeg)

![](_page_45_Picture_17.jpeg)

![](_page_45_Picture_18.jpeg)

![](_page_45_Picture_19.jpeg)

![](_page_45_Picture_20.jpeg)

![](_page_45_Picture_21.jpeg)

![](_page_45_Picture_22.jpeg)

# what should we expect as a step up in energy?

# **Tevatron** *pp*, 1.96 TeV, 10 fb<sup>-1</sup>

# Exclusion limit ~ 1.2 TeV

(if they had analysed all their data in electron and muon channels: actual CDF limit 1.071 TeV, 4.7fb<sup>-1</sup>, µµ only)

# I like the $Z'_{SSM}$ as a simple measure of progress (perhaps not very "exciting", but simple and most experiments look for it)

# × 5.6

# LHC *pp*, **14 TeV**, **3000 fb**<sup>-1</sup>

# Exclusion limit ~ 6.7 TeV

(electron and muon channels. single experiment)

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# increase in precision

#### precision reach on effective couplings from SMEFT global fit

![](_page_47_Figure_2.jpeg)

https://arxiv.org/abs/2206.08326

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![](_page_47_Picture_7.jpeg)

![](_page_47_Picture_8.jpeg)

# increase in precision is like $\times 4 - 5$ increase in energy reach

### 95% CL scale limits on 4–fermion contact interactions from $O_{2B}$

![](_page_48_Figure_2.jpeg)

https://arxiv.org/abs/2206.08326

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![](_page_48_Picture_7.jpeg)

# step up in energy for direct searches?

# LHC *pp*, **13 TeV**, **3000 fb**<sup>-1</sup>

# Exclusion limit ~ 6.7 TeV

(electron and muon channels, single experiment)

# I like the $Z'_{SSM}$ as a simple measure of progress (perhaps not very "exciting", but simple and most experiments look for it)

# × 6.1

# FCC-hh *pp*, **100 TeV**, **20 ab**<sup>-1</sup>

# Exclusion limit ~ 41 TeV

(based on PDF luminosity scaling, assuming detectors can handle muons and electrons at these energies)

![](_page_49_Picture_13.jpeg)

![](_page_49_Picture_14.jpeg)

# step up in energy for direct searches?

# LHC *pp*, **13 TeV**, **3000 fb**<sup>-1</sup>

# Exclusion limit ~ 6.7 TeV

(electron and muon channels. single experiment)

![](_page_50_Picture_7.jpeg)

# I like the $Z'_{SSM}$ as a simple measure of progress (perhaps not very "exciting", but simple and most experiments look for it)

![](_page_50_Picture_9.jpeg)

# SppC 125 TeV. 5 ab<sup>-1</sup>

# Exclusion limit ~ 43 TeV

(based on PDF luminosity scaling, assuming detectors can handle muons and electrons at these energies)

![](_page_50_Picture_14.jpeg)

![](_page_50_Picture_15.jpeg)

# desirable features of the next major HEP project(s)?

an important target to be reached  $\sim$  guaranteed discovery

exploration into the unknown by a significant factor in energy

likelihood of success, robustness (e.g. multiple experiments)

cost-effective construction & operation, low carbon footprint, novel technologies

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major progress on a broad array of particle physics topics

![](_page_51_Picture_10.jpeg)

![](_page_51_Picture_11.jpeg)

![](_page_51_Picture_12.jpeg)

# illustration is for FCC — but message is comparable for other colliders

![](_page_52_Figure_1.jpeg)

Gavir Slide from C. Grojean @ FCC Week'22

![](_page_52_Figure_3.jpeg)

![](_page_52_Picture_11.jpeg)

# conclusions

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. . . . . . . . . . . . . . . .

![](_page_53_Picture_3.jpeg)

![](_page_53_Picture_4.jpeg)

# Conclusions

- at higher-energy colliders
  - origin of electron mass at circular e<sup>+</sup>e<sup>-</sup> colliders?
- The step up in energy reach that we expect is  $\sim \times 4 5$ 
  - increase  $\sim \times 18$
  - some scenarios)

> There is a guaranteed discovery: directly establishing Higgs self-interaction, which holds the SM together, via robust precision of Higgs factory and direct measurement

➤ is there a chance of a second no-lose theorem in establishing (or disproving) SM

> e+e- colliders deliver that mostly in "indirect" sensitivity, through precision

► FCC-hh/SppS deliver that in direct search sensitivity (muon collider does for

> Diversity and robustness of the programme = essential part of their strength

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![](_page_54_Figure_16.jpeg)

![](_page_55_Picture_0.jpeg)

Looking at the wide variety of alternatives which have been proposed, it might appear that theorists are in disarray but it seems to me that the present situation is an inevitable consequence of the successes of the 1970's. The problems of the 1960's - the nature of hadrons, the nature of the strong force, the nature of the weak force - have We now confront deeper problems - the origin of mass, the choice of been solved. fundamental building blocks (the problem of flavour), the question of further unification of forces including gravity, the origin of charge and of gauge symmetry. It is only to be expected that many of the first attempts to grapple with these problems will be misguided. As ever, we must reply on experiment to reveal the truth.

#### via Nathaniel Craig @ CERN-TH naturalness workshop

#### PHYSICS WITH A MULTI-TeV HADRON COLLIDER

#### C.H. Llewellyn Smith,

![](_page_55_Picture_7.jpeg)

# backup

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![](_page_56_Picture_3.jpeg)

# What does $2.6 \times 10^{28}$ kg/m<sup>3</sup> mean?

![](_page_57_Picture_1.jpeg)

Gavin Salam

![](_page_57_Picture_7.jpeg)

# What does $2.6 \times 10^{28}$ kg/m<sup>3</sup> mean?

![](_page_58_Picture_1.jpeg)

Gavin Salam

![](_page_58_Picture_4.jpeg)

![](_page_58_Picture_8.jpeg)

# What does $2.6 \times 10^{28}$ kg/m<sup>3</sup> mean?

![](_page_59_Picture_1.jpeg)

# fit the mass of the sun into a standard 40ft shipping container

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![](_page_59_Picture_5.jpeg)

![](_page_59_Picture_9.jpeg)

https://upload.wikimedia.org/wikipedia/commons/4/44/Mecanismo de Higgs PH.png https://en.wikipedia.org/wiki/Neutron\_star#/media/File:Neutron\_Star\_X-ray\_beaming\_with\_accretion\_disk.jpg

![](_page_60_Figure_1.jpeg)

![](_page_60_Picture_2.jpeg)

![](_page_60_Picture_4.jpeg)

Table 3.3: Values for  $1\sigma$  sensitivity on the S and T parameters. In all cases the value shown is after combination with HL-LHC. For ILC and CLIC the projections are shown with and without dedicated running at the Z-pole. All other oblique parameters are set to zero. The intrinsic theory uncertainty is also set to zero.

	Current	HL-LHC	II	$LC_{250}$	CEPC	FCC-ee	C	LIC <sub>380</sub>
				(& ILC <sub>91</sub> )				(& CLIC <sub>91</sub> )
S	0.13	0.053	0.012	0.009	0.0068	0.0038	0.032	0.011
T	0.08	0.041	0.014	0.013	0.0072	0.0022	0.023	0.012

improvements of up to  $\times 14 - 18$ 

![](_page_61_Picture_7.jpeg)

![](_page_61_Figure_8.jpeg)

![](_page_61_Picture_9.jpeg)

![](_page_62_Figure_0.jpeg)

It's not inconceivable that the top mass could be sufficiently mis-measured at hadron colliders that the SM-universe is stable all the way to the Planck scale

condition in terms of the pole top mass. We can express the stability condition of eq. (64) as  $M_t < (171.53 \pm 0.15 \pm 0.23_{\alpha_3} \pm 0.15_{M_h}) \,\text{GeV} = (171.53 \pm 0.42) \,\text{GeV}.$ (66)

arXiv:1307.3536

![](_page_62_Figure_6.jpeg)

# Searches at muon collider

Plots being shown suggest: 4 TeV muon collider beats a 100 TeV pp collider in searches for new physics.

Useful to nuance the statement:

- ► 100 TeV pp, 20 ab<sup>-1</sup> can discover Z' up to  $m_{Z'} \sim 38 \text{ TeV}$
- For  $\mu\mu$  collider to discover Z' at  $m_{Z'} \sim 38$  TeV, it needs  $\sqrt{s} \sim 38$  TeV (with lower  $\sqrt{s}$  you would see deviation from SM, but not know what it is)
- TeV pp machine

 $s_p$ 

#### Fig. 3 of Snowmass Muon Collider Forum Report

![](_page_63_Figure_9.jpeg)

 $\sqrt{s_{\mu}}$  [TeV]

 $\blacktriangleright$  However a 38 TeV muon collider would be much better at studying the Z' than the 100

![](_page_63_Picture_12.jpeg)

![](_page_63_Picture_13.jpeg)