

Event shapes for hadron colliders

Gavin P. Salam

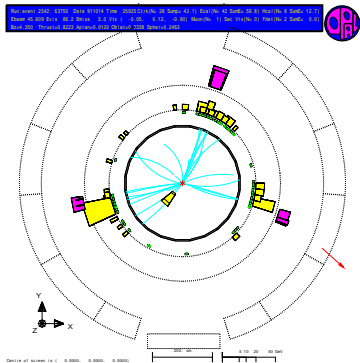
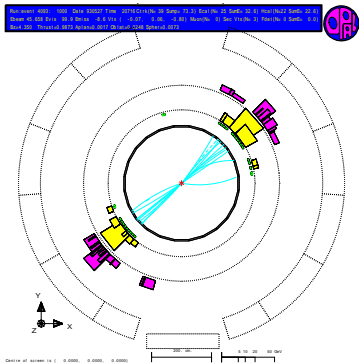
in collaboration with Andrea Banfi & Giulia Zanderighi

LPTHE, Universities of Paris VI and VII and CNRS

Max Planck Institut für Physik

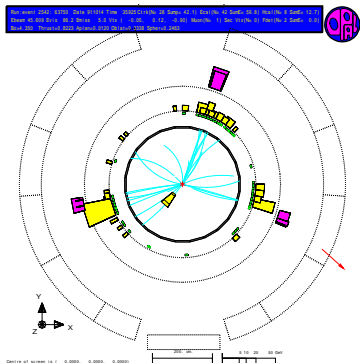
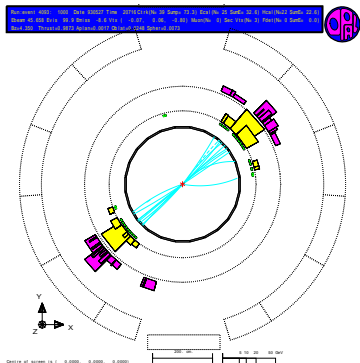
Munich, July 2007

A wealth of information about QCD lies in its final states. Problem is how to extract it.



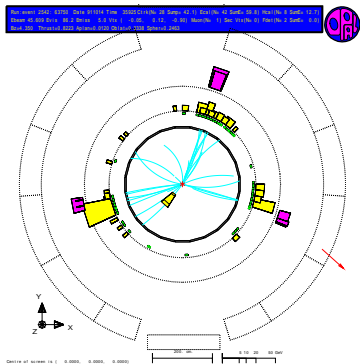
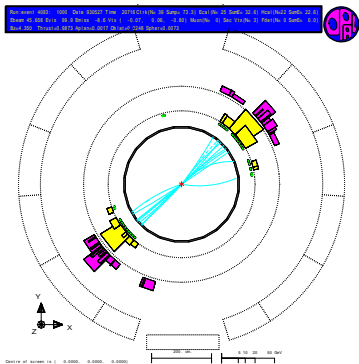
One option is to use a jet-algorithm and *classify* events – 2 jets, 3 jets, . . .
 But this does not capture *continuous nature* of variability of events.

A wealth of information about QCD lies in its final states. Problem is how to extract it.



One option is to use a jet-algorithm and *classify* events – 2 jets, 3 jets, ...
 But this does not capture *continuous nature* of variability of events.

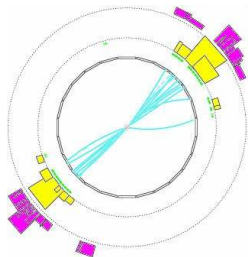
A wealth of information about QCD lies in its final states. Problem is how to extract it.



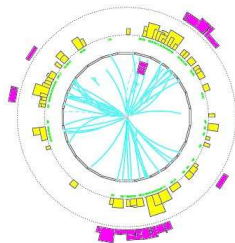
One option is to use a jet-algorithm and *classify* events – 2 jets, 3 jets, ...
 But this does not capture *continuous nature* of variability of events.

First discussion goes back to 1964. Serious work got going in late '70s. Various proposals to measure *shape* of events. Most famous example is **Thrust**:

$$T = \max_{\vec{n}_T} \frac{\sum_i |\vec{p}_i \cdot \vec{n}_T|}{\sum_i |\vec{p}_i|},$$



2-jet event: $T \simeq 1$

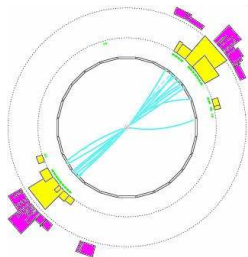


3-jet event: $T \simeq 2/3$

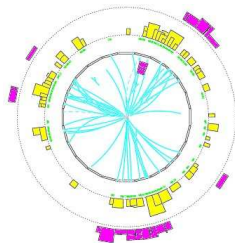
There exist many other measures of aspects of the shape: Thrust-Major, C-parameter, broadening, heavy-jet mass, jet-resolution parameters, ...

First discussion goes back to 1964. Serious work got going in late '70s. Various proposals to measure *shape* of events. Most famous example is **Thrust**:

$$T = \max_{\vec{n}_T} \frac{\sum_i |\vec{p}_i \cdot \vec{n}_T|}{\sum_i |\vec{p}_i|},$$



2-jet event: $T \simeq 1$

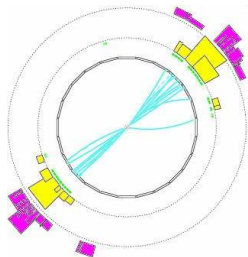


3-jet event: $T \simeq 2/3$

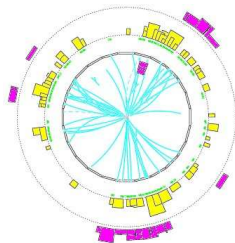
There exist many other measures of aspects of the shape: Thrust-Major, C-parameter, broadening, heavy-jet mass, jet-resolution parameters,...

First discussion goes back to 1964. Serious work got going in late '70s. Various proposals to measure *shape* of events. Most famous example is **Thrust**:

$$T = \max_{\vec{n}_T} \frac{\sum_i |\vec{p}_i \cdot \vec{n}_T|}{\sum_i |\vec{p}_i|},$$



2-jet event: $T \simeq 1$

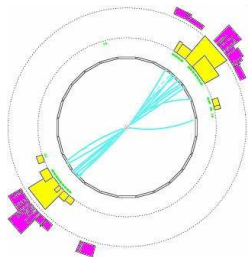


3-jet event: $T \simeq 2/3$

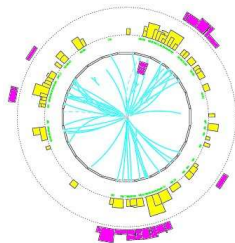
There exist many other measures of aspects of the shape: Thrust-Major, C-parameter, broadening, heavy-jet mass, jet-resolution parameters, ...

First discussion goes back to 1964. Serious work got going in late '70s. Various proposals to measure *shape* of events. Most famous example is **Thrust**:

$$T = \max_{\vec{n}_T} \frac{\sum_i |\vec{p}_i \cdot \vec{n}_T|}{\sum_i |\vec{p}_i|},$$

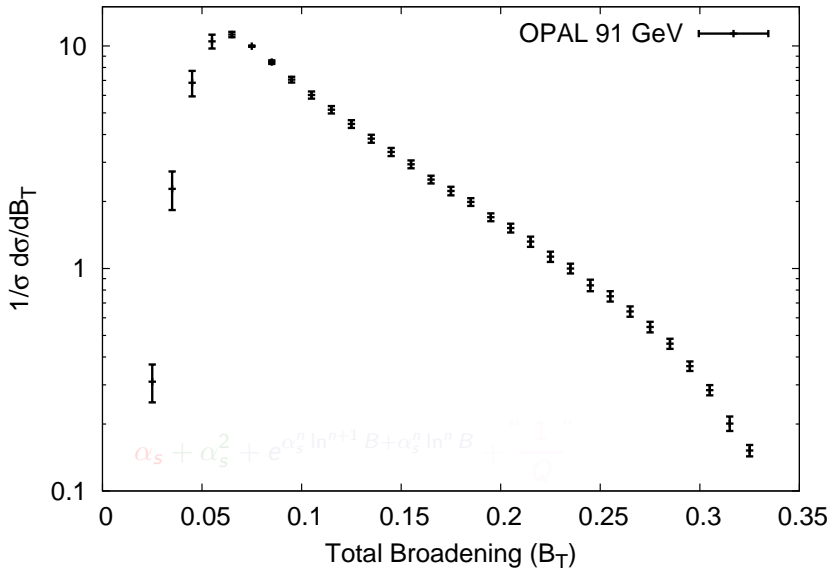


2-jet event: $T \simeq 1$

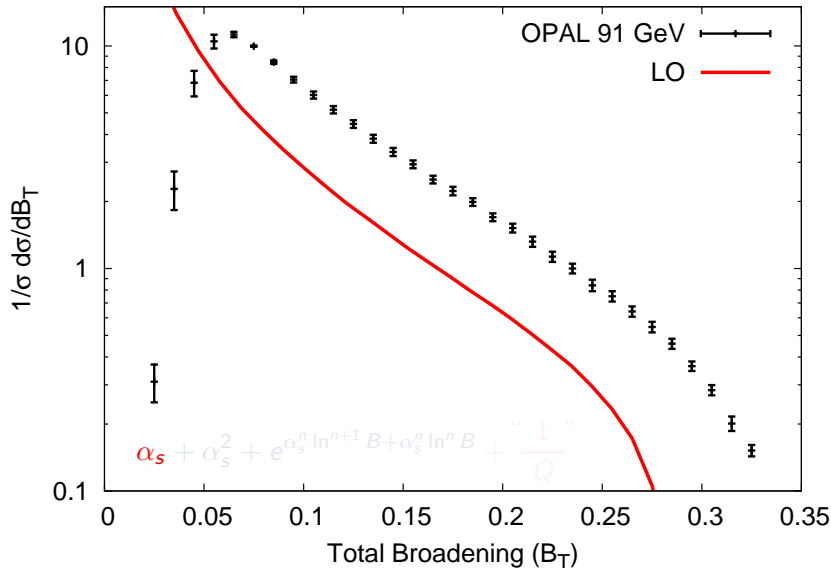


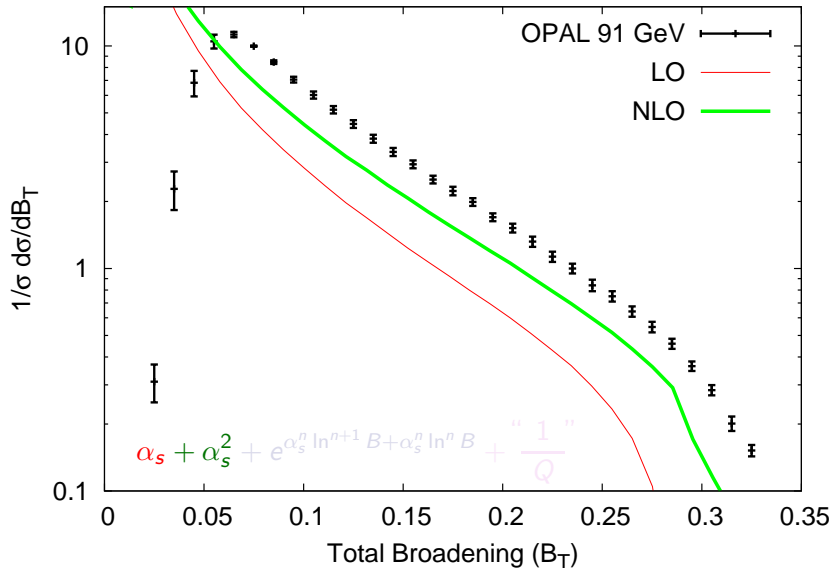
3-jet event: $T \simeq 2/3$

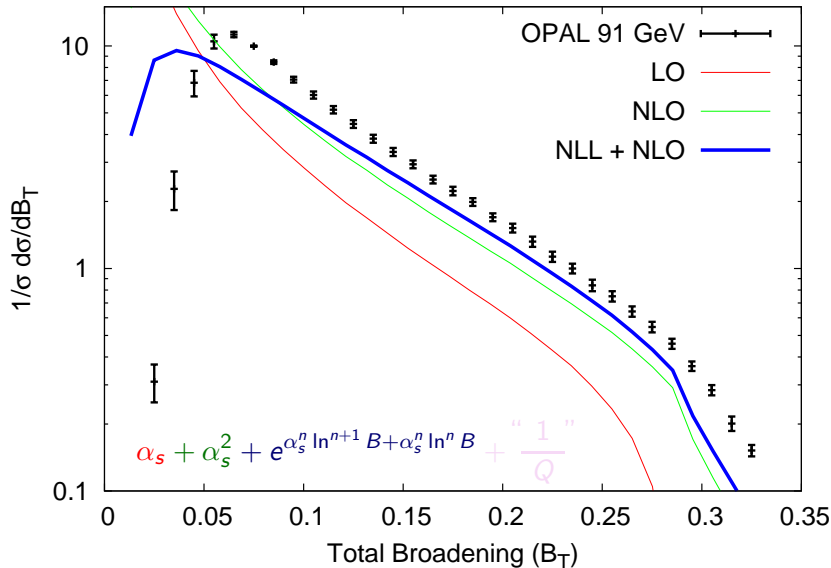
There exist many other measures of aspects of the shape: **Thrust-Major**, **C-parameter**, **broadening**, **heavy-jet mass**, **jet-resolution parameters**,...

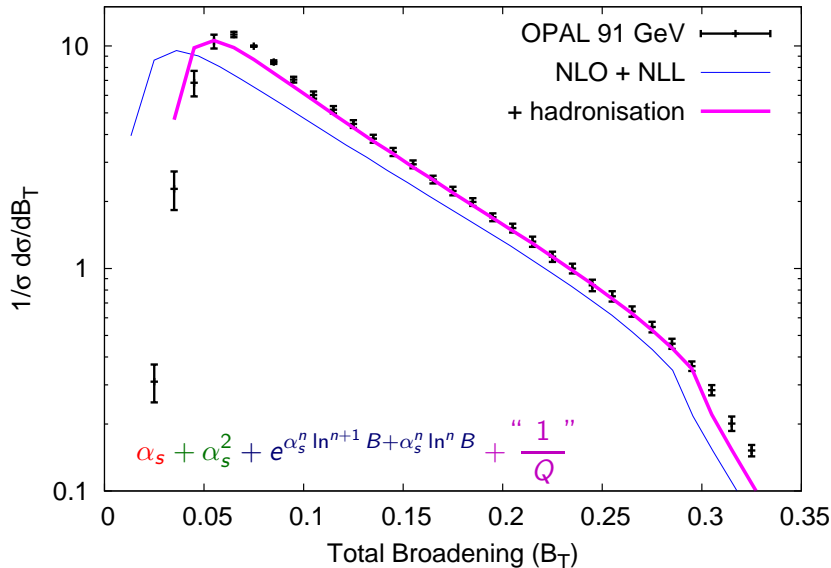


Event-shapes: probe range of physics

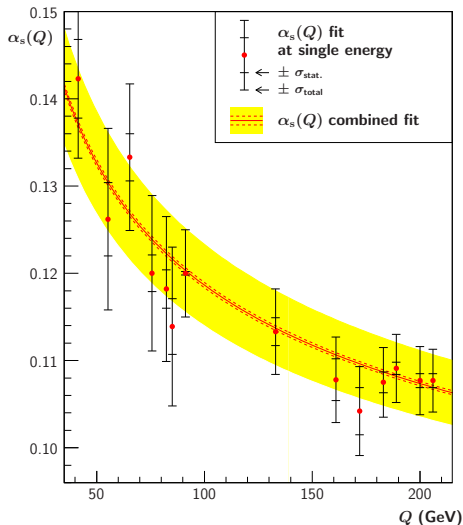








Event shapes: high information content



Much knowledge has been extracted from event-shapes in e^+e^- and DIS:

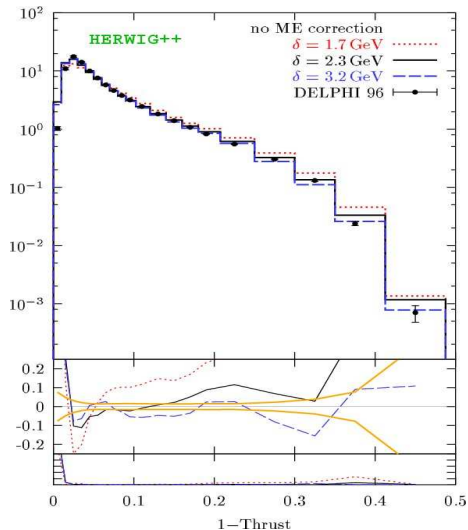
- α_s fits
- Tuning of Monte Carlos
- Colour factor fits (C_A, C_F, \dots)
- Studies of analytical hadronisation models ($1/Q$, shape functions, ...)

But mostly neglected so far at hadron colliders

except: CDF broadening ('91)

D0 Thrust ('02)

Event shapes: high information content



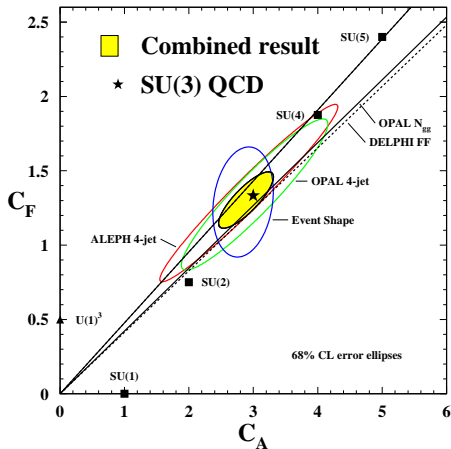
Much knowledge has been extracted from event-shapes in e^+e^- and DIS:

- α_s fits
- Tuning of Monte Carlos
- Colour factor fits (C_A, C_F, \dots)
- Studies of analytical hadronisation models ($1/Q$, shape functions, ...)

But mostly neglected so far at hadron colliders

except: CDF broadening ('91)
D0 Thrust ('02)

Event shapes: high information content



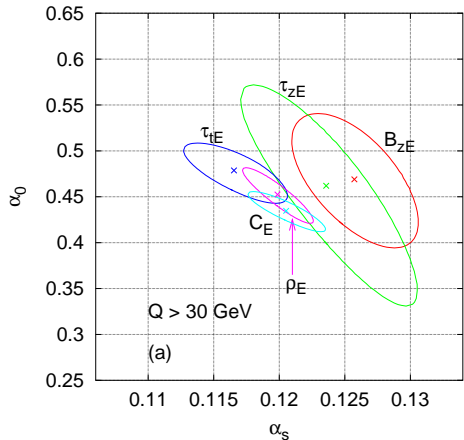
Much knowledge has been extracted from event-shapes in e^+e^- and DIS:

- α_s fits
- Tuning of Monte Carlos
- Colour factor fits (C_A, C_F, \dots)
- Studies of analytical hadronisation models ($1/Q$, shape functions, ...)

But mostly neglected so far at hadron colliders

except: CDF broadening ('91)
 D0 Thrust ('02)

Event shapes: high information content



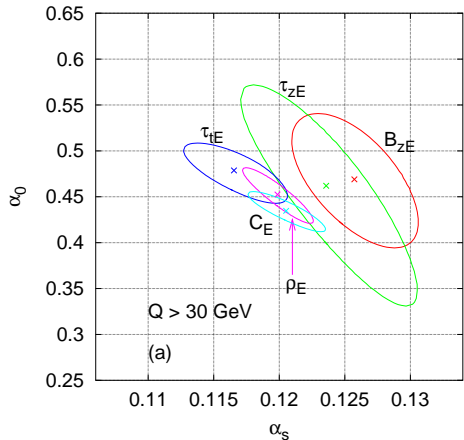
Much knowledge has been extracted from event-shapes in e^+e^- and DIS:

- α_s fits
- Tuning of Monte Carlos
- Colour factor fits (C_A, C_F, \dots)
- Studies of analytical hadronisation models ($1/Q$, shape functions, ...)

But mostly neglected so far at hadron colliders

except: CDF broadening ('91)
 D0 Thrust ('02)

Event shapes: high information content



Much knowledge has been extracted from event-shapes in e^+e^- and DIS:

- α_s fits
- Tuning of Monte Carlos
- Colour factor fits (C_A, C_F, \dots)
- Studies of analytical hadronisation models ($1/Q$, shape functions, ...)

But mostly neglected so far at hadron colliders

except: CDF broadening ('91)
 D0 Thrust ('02)

Various processes:

- $pp \rightarrow W/Z/H \text{ boson} + \text{jet}$
- $pp \rightarrow 2 \text{ jets}$

Banfi Marchesini Smye Zanderighi '01

Main subject of this talk

Standard applications (e.g.)

- Measure α_s
- As for 3-jet/2-jet ratio in $p\bar{p}$,
reduce dependence on PDFs
- But for event-shapes \rightarrow
distribution
- Far more information than
3-jet/2-jet ratio

New territory

- 4-jet (2 + 2) topology \rightarrow novel
perturbative structures
soft colour evln matrices
- 3 & 4-jet topologies (& g-jets)
 \rightarrow rich environment for
analytical non-pert. studies
- Underlying event — test models
(analytical & MC).

Variety of event-shape observables \rightarrow complementary information \rightarrow
disentangle the different physics issues.

Various processes:

- $pp \rightarrow W/Z/H \text{ boson} + \text{jet}$
- $pp \rightarrow 2 \text{ jets}$

Standard applications (e.g.)

- Measure α_s
- As for 3-jet/2-jet ratio in $p\bar{p}$, reduce dependence on PDFs
- But for event-shapes \rightarrow *distribution*
- Far more information than 3-jet/2-jet ratio

Banfi Marchesini Smye Zanderighi '01

Main subject of this talk

New territory

- 4-jet (2 + 2) topology \rightarrow novel perturbative structures
soft colour evln matrices
- 3 & 4-jet topologies (& g-jets) \rightarrow rich environment for analytical non-pert. studies
- Underlying event — test models (analytical & MC).

Variety of event-shape observables \rightarrow complementary information \rightarrow disentangle the different physics issues.

Various processes:

- $pp \rightarrow W/Z/H$ boson + jet
- $pp \rightarrow 2$ jets

Standard applications (e.g.)

- Measure α_s
- As for 3-jet/2-jet ratio in $p\bar{p}$, reduce dependence on PDFs
- But for event-shapes \rightarrow *distribution*
- Far more information than 3-jet/2-jet ratio

Banfi Marchesini Smye Zanderighi '01

Main subject of this talk

New territory

- 4-jet (2 + 2) topology \rightarrow novel perturbative structures
soft colour evln matrices
- 3 & 4-jet topologies (& g-jets) \rightarrow rich environment for analytical non-pert. studies
- Underlying event — test models (analytical & MC).

Variety of event-shape observables \rightarrow complementary information \rightarrow disentangle the different physics issues.

Various processes:

- $pp \rightarrow W/Z/H \text{ boson} + \text{jet}$
- $pp \rightarrow 2 \text{ jets}$

Standard applications (e.g.)

- Measure α_s
- As for 3-jet/2-jet ratio in $p\bar{p}$, reduce dependence on PDFs
- But for event-shapes \rightarrow *distribution*
- Far more information than 3-jet/2-jet ratio

Banfi Marchesini Smye Zanderighi '01

Main subject of this talk

New territory

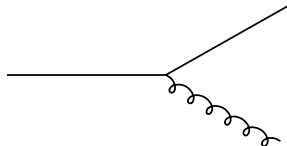
- 4-jet (2 + 2) topology \rightarrow novel perturbative structures
soft colour evln matrices
- 3 & 4-jet topologies (& g-jets) \rightarrow rich environment for analytical non-pert. studies
- Underlying event — test models (analytical & MC).

Variety of event-shape observables \rightarrow complementary information \rightarrow disentangle the different physics issues.

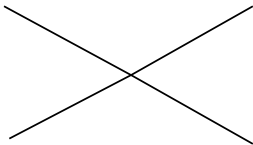
Multi-jet final states: relative colour of pairs of hard partons determines soft large-angle radiation.



2 jets: always in a *colour singlet*



3 jets: colour state of any pair *fixed by third parton* (colour conservation).



4 jets: a given pair can be in various colour states. Soft virtual corrections mix colour states.

Resummation leads to *matrix evolution equation for colour state of amplitudes* ('soft anomalous dimensions')

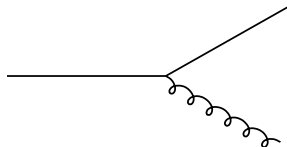
Developed at Stony Brook: Botts, Kidonakis, Oderda & Sterman '89-99

Interesting to test it (NB: used also for top threshold corrections).

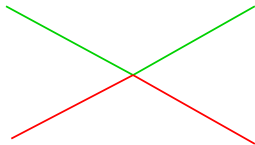
Multi-jet final states: relative colour of pairs of hard partons determines soft large-angle radiation.



2 jets: always in a *colour singlet*



3 jets: colour state of any pair *fixed by third parton* (colour conservation).



4 jets: a given pair can be in various colour states. Soft virtual corrections mix colour states.

Resummation leads to *matrix evolution equation for colour state of amplitudes* ('soft anomalous dimensions')

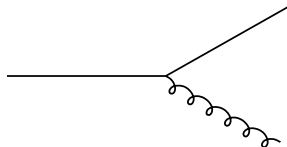
Developed at Stony Brook: Botts, Kidonakis, Oderda & Sterman '89–99

Interesting to test it (NB: used also for top threshold corrections).

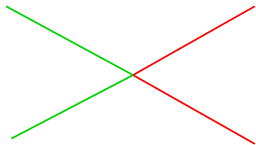
Multi-jet final states: relative colour of pairs of hard partons determines soft large-angle radiation.



2 jets: always in a *colour singlet*



3 jets: colour state of any pair *fixed by third parton* (colour conservation).



4 jets: a given pair can be in various colour states. Soft virtual corrections mix colour states.

Resummation leads to *matrix evolution equation for colour state of amplitudes* ('soft anomalous dimensions')

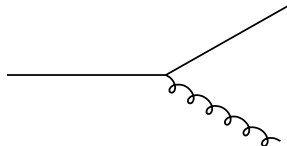
Developed at Stony Brook: Botts, Kidonakis, Oderda & Sterman '89–99

Interesting to test it (NB: used also for top threshold corrections).

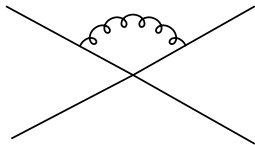
Multi-jet final states: relative colour of pairs of hard partons determines soft large-angle radiation.



2 jets: always in a *colour singlet*



3 jets: colour state of any pair *fixed by third parton* (colour conservation).



4 jets: a given pair can be in various colour states. *Soft virtual corrections mix colour states.*

Resummation leads to *matrix evolution equation for colour state of amplitudes* ('soft anomalous dimensions')

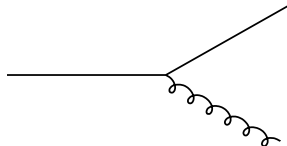
Developed at Stony Brook: Botts, Kidonakis, Oderda & Sterman '89–99

Interesting to test it (NB: used also for top threshold corrections).

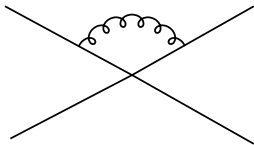
Multi-jet final states: relative colour of pairs of hard partons determines soft large-angle radiation.



2 jets: always in a *colour singlet*



3 jets: colour state of any pair *fixed by third parton* (colour conservation).



4 jets: a given pair can be in various colour states. *Soft virtual corrections mix colour states.*

Resummation leads to *matrix evolution equation for colour state of amplitudes* ('soft anomalous dimensions')

Developed at Stony Brook: Botts, Kidonakis, Oderda & Sterman '89–99

Interesting to test it (NB: used also for top threshold corrections).

Fixed order

- Event shapes trivial for Born events (e.g. $p\bar{p} \rightarrow 2$ jets, thrust=1)
- First non-trivial order (LO) is Born + 1 parton, i.e. $p\bar{p} \rightarrow 3$ jets
- For NLO, need a program like NLOJET++ ($p\bar{p} \rightarrow 3$ jets @ NLO)

Nagy, '01 & '03

- Also:

- Kilgore-Giele code ($p\bar{p} \rightarrow 3$ jets @ NLO),
- MCFM ($p\bar{p} \rightarrow W/Z/H + 2$ jets @ NLO)

Campbell & Ellis '02

Resummation

- In e^+e^- it was always done by hand, one observable at a time.
- *Next-to-leading logs* (NLL) are tedious, complicated, error-prone.
- Recently automated: Computer-Automated Expert Semi-Analytical Resummer (CAESAR). Banfi, GPS & Zanderighi '01-'04
- For $p\bar{p} \rightarrow 2$ jets, uses 'Stony Brook' soft-colour evolution matrices.
- Currently restricted to *continuously-global* observables

Fixed order

- Event shapes trivial for Born events (e.g. $p\bar{p} \rightarrow 2$ jets, thrust=1)
- First non-trivial order (LO) is Born + 1 parton, i.e. $p\bar{p} \rightarrow 3$ jets
- For NLO, need a program like NLOJET++ ($p\bar{p} \rightarrow 3$ jets @ NLO)

Nagy, '01 & '03

- Also:

- Kilgore-Giele code ($p\bar{p} \rightarrow 3$ jets @ NLO),
- MCFM ($p\bar{p} \rightarrow W/Z/H + 2$ jets @ NLO)

Campbell & Ellis '02

Resummation

- In e^+e^- it was always done by hand, one observable at a time.
- *Next-to-leading logs* (NLL) are tedious, complicated, error-prone.
- Recently automated: Computer-Automated Expert Semi-Analytical Resummer (CAESAR).
- For $p\bar{p} \rightarrow 2$ jets, uses 'Stony Brook' soft-colour evolution matrices.
- Currently restricted to *continuously-global* observables

Banfi, GPS & Zanderighi '01-'04

$e^+e^- \rightarrow 2 \text{ jets}$

S. Catani, G. Turnock, B. R. Webber and L. Trentadue, *Thrust distribution in e^+e^- annihilation*, Phys. Lett. B **263** (1991) 491.

S. Catani, G. Turnock and B. R. Webber, *Heavy jet mass distribution in e^+e^- annihilation*, Phys. Lett. B **272** (1991) 368.

S. Catani, Yu. L. Dokshitzer, M. Olsson, G. Turnock and B. R. Webber, *New clustering algorithm for multi-jet cross-sections in e^+e^- annihilation*, Phys. Lett. B **269** (1991) 432.

S. Catani, L. Trentadue, G. Turnock and B. R. Webber, *Resummation of large logarithms in e^+e^- event shape distributions*, Nucl. Phys. B **407** (1993) 3.

S. Catani, G. Turnock and B. R. Webber, *Jet broadening measures in e^+e^- annihilation*, Phys. Lett. B **295** (1992) 269.

G. Dissertori and M. Schmelling, *An Improved theoretical prediction for the two jet rate in e^+e^- annihilation*, Phys. Lett. B **361** (1995) 167.

Y. L. Dokshitzer, A. Lucenti, G. Marchesini and GPS, *On the QCD analysis of jet broadening*, JHEP **9801** (1998) 011

S. Catani and B. R. Webber, *Resummed C-parameter distribution in e^+e^- annihilation*, Phys. Lett. B **427** (1998) 377

S. J. Burby and E. W. Glover, *Resumming the light hemisphere mass and narrow jet broadening distributions in e^+e^- annihilation*, JHEP **0104** (2001) 029

M. Dasgupta and GPS, *Resummation of non-global QCD observables*, Phys. Lett. B **512** (2001) 323

C. F. Berger, T. Kucs and G. Sterman, *Event shape / energy flow correlations*, Phys. Rev. D **68** (2003) 014012

DIS 1+1 jet

V. Antonelli, M. Dasgupta and GPS, *Resummation of thrust distributions in DIS*, JHEP **0002** (2000) 001

M. Dasgupta and GPS, *Resummation of the jet broadening in DIS*, Eur. Phys. J. C **24** (2002) 213

M. Dasgupta and GPS, *Resummed event-shape variables in DIS*, JHEP **0208** (2002) 032

 e^+e^- , DY, DIS 3 jets

A. Banfi, G. Marchesini, Y. L. Dokshitzer and G. Zanderighi, *QCD analysis of near-to-planar 3-jet events*, JHEP **0007** (2000) 002

A. Banfi, Y. L. Dokshitzer, G. Marchesini and G. Zanderighi, *Near-to-planar 3-jet events in and beyond QCD perturbation theory*, Phys. Lett. B **508** (2001) 269

A. Banfi, Y. L. Dokshitzer, G. Marchesini and G. Zanderighi, *QCD analysis of D-parameter in near-to-planar three-jet events*, JHEP **0105** (2001) 040

A. Banfi, G. Marchesini, G. Smye and G. Zanderighi, *Out-of-plane QCD radiation in hadronic Z0 production*, JHEP **0108** (2001) 047

A. Banfi, G. Marchesini, G. Smye and G. Zanderighi, *Out-of-plane QCD radiation in DIS with high $p(t)$ jets*, JHEP **0111** (2001) 066

A. Banfi, G. Marchesini and G. Smye, *Azimuthal correlation in DIS*, JHEP **0204** (2002) 024

Average: 1 observable per paper

Analytical work (done once and for all)

- A1. derive a master formula for a generic observable in terms of simple properties of the observable
- A2. formulate the exact applicability conditions for the master formula

Numerical work (to be repeated for each observable)

- N1. let an "expert system" investigate the applicability conditions
- N2. it also determines the inputs for the master formula
- N3. straightforward evaluation of the master formula, including phase space integration etc.

Note: N1 and N2 are core of automation

- a) they require high precision arithmetic to take asymptotic (soft & collinear) limits
- b) validation of hypotheses uses methods inspired by "Experimental Mathematics"

Analytical work (done once and for all)

- A1. derive a **master formula** for a **generic observable** in terms of simple properties of the observable
- A2. formulate the exact **applicability conditions** for the master formula

Numerical work (to be repeated for each observable)

- N1. let an "**expert system**" investigate the applicability conditions
- N2. it also determines the **inputs for the master formula**
- N3. straightforward **evaluation of the master formula**, including phase space integration etc.

Note: N1 and N2 are core of automation

- a) they require **high precision arithmetic** to take asymptotic (soft & collinear) limits
- b) validation of hypotheses uses methods inspired by "**Experimental Mathematics**"

Analytical work (done once and for all)

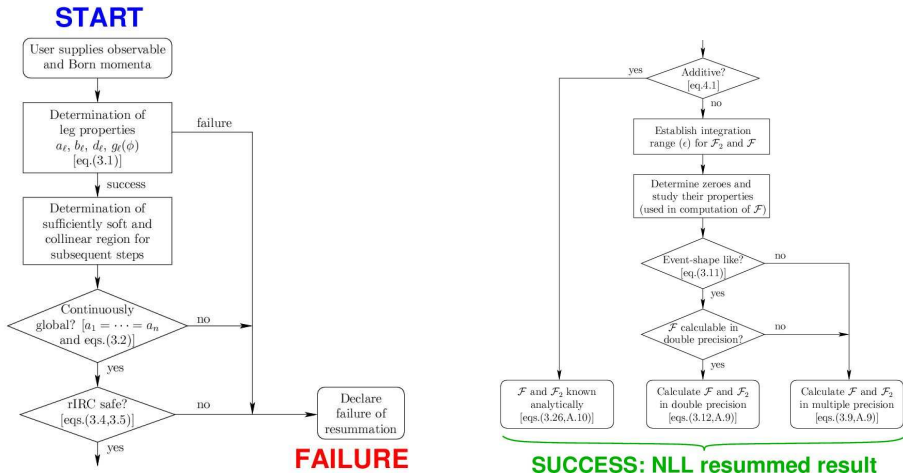
- A1. derive a **master formula** for a **generic observable** in terms of simple properties of the observable
- A2. formulate the exact **applicability conditions** for the master formula

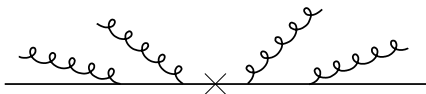
Numerical work (to be repeated for each observable)

- N1. let an "**expert system**" investigate the applicability conditions
- N2. it also determines the **inputs for the master formula**
- N3. straightforward **evaluation of the master formula**, including phase space integration etc.

Note: N1 and N2 are core of automation

- a) they require **high precision arithmetic** to take asymptotic (soft & collinear) limits
- b) validation of hypotheses uses methods inspired by "**Experimental Mathematics**"



Global observable:e.g. total e^+e^- Broadening, B 

making $B \ll 1$ restricts emissions everywhere.

Coherence + globalness:

→ emissions can be resummed as if independent (*proved*)

Answers guaranteed to NLL accuracy

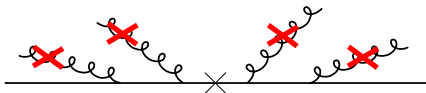
Non-Global observable:Right-hemisphere Broadening, B_R

making $B_R \ll 1$ restricts emissions in right-hand hemisphere (\mathcal{H}_R).

Tempting to *assume* one can:

- ignore left hemisphere (\mathcal{H}_L)
- use independent emission approximation in \mathcal{H}_R .

WRONG AT NLL ACCURACY

Global observable:e.g. total e^+e^- Broadening, B making $B \ll 1$ restricts emissions everywhere.

Coherence + globalness:

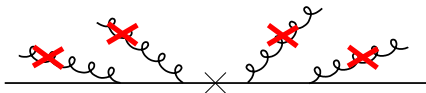
→ emissions can be resummed as if independent (*proved*)

Answers guaranteed to NLL accuracy

Non-Global observable:Right-hemisphere Broadening, B_R making $B_R \ll 1$ restricts emissions in right-hand hemisphere (\mathcal{H}_R).Tempting to *assume* one can:

- ignore left hemisphere (\mathcal{H}_L)
- use independent emission approximation in \mathcal{H}_R .

WRONG AT NLL ACCURACY

Global observable:e.g. total e^+e^- Broadening, B making $B \ll 1$ restricts emissions everywhere.

Coherence + globalness:

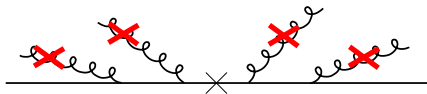
→ emissions can be resummed as if independent (*proved*)

Answers guaranteed to NLL accuracy

Non-Global observable:Right-hemisphere Broadening, B_R making $B_R \ll 1$ restricts emissions in right-hand hemisphere (\mathcal{H}_R).Tempting to *assume* one can:

- ignore left hemisphere (\mathcal{H}_L)
- use independent emission approximation in \mathcal{H}_R .

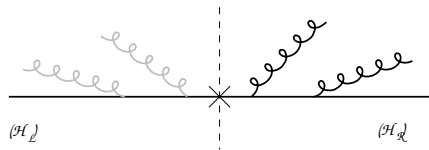
WRONG AT NLL ACCURACY

Global observable:e.g. total e^+e^- Broadening, B making $B \ll 1$ restricts emissions everywhere.

Coherence + globalness:

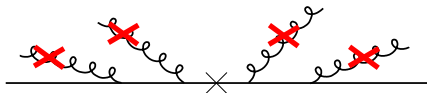
→ emissions can be resummed as if independent (*proved*)

Answers guaranteed to NLL accuracy

Non-Global observable:Right-hemisphere Broadening, B_R making $B_R \ll 1$ restricts emissions in right-hand hemisphere (\mathcal{H}_R).Tempting to *assume* one can:

- ignore left hemisphere (\mathcal{H}_L)
- use independent emission approximation in \mathcal{H}_R .

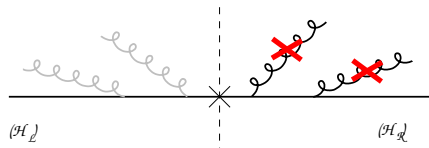
WRONG AT NLL ACCURACY

Global observable:e.g. total e^+e^- Broadening, B making $B \ll 1$ restricts emissions everywhere.

Coherence + globalness:

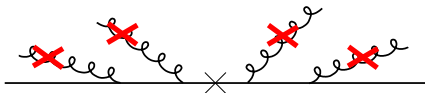
→ emissions can be resummed as if independent (*proved*)

Answers guaranteed to NLL accuracy

Non-Global observable:Right-hemisphere Broadening, B_R making $B_R \ll 1$ restricts emissions in right-hand hemisphere (\mathcal{H}_R).Tempting to *assume* one can:

- ignore left hemisphere (\mathcal{H}_L)
- use independent emission approximation in \mathcal{H}_R .

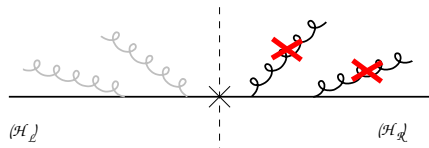
WRONG AT NLL ACCURACY

Global observable:e.g. total e^+e^- Broadening, B making $B \ll 1$ restricts emissions everywhere.

Coherence + globalness:

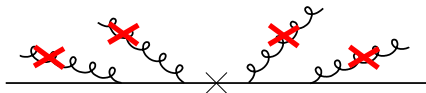
→ emissions can be resummed as if independent (*proved*)

Answers guaranteed to NLL accuracy

Non-Global observable:Right-hemisphere Broadening, B_R making $B_R \ll 1$ restricts emissions in right-hand hemisphere (\mathcal{H}_R).Tempting to *assume* one can:

- ignore left hemisphere (\mathcal{H}_L)
- use independent emission approximation in \mathcal{H}_R .

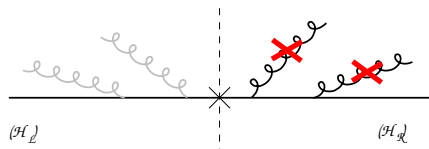
WRONG AT NLL ACCURACY

Global observable:e.g. total e^+e^- Broadening, B making $B \ll 1$ restricts emissions everywhere.

Coherence + globalness:

→ emissions can be resummed as if independent (*proved*)

Answers guaranteed to NLL accuracy

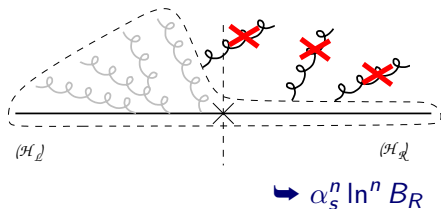
Non-Global observable:Right-hemisphere Broadening, B_R making $B_R \ll 1$ restricts emissions in right-hand hemisphere (\mathcal{H}_R).Tempting to *assume* one can:

- ignore left hemisphere (\mathcal{H}_L)
- use independent emission approximation in \mathcal{H}_R .

WRONG AT NLL ACCURACY

All-orders:

Forbid coherent radiation from energy-ordered ensembles of large-angle gluons

Difficulties:

- Logarithms resummed so far only in large- N_c limit
- In general, boundary between the two regions may have arbitrary shape.
- It may depend on the pattern of emissions (e.g. with jet algorithm).

Appleby & Seymour '02, '03

Banfi & Dasgupta '05

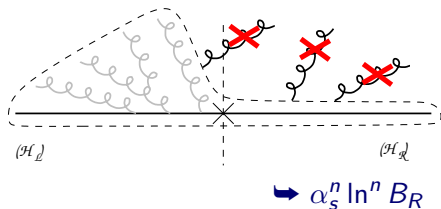
Delenda, A, B & D '06

Resummation of a general non-global observable is tricky.
For time-being CAESAR deals only with global observables.

NB: (most) Monte Carlo's are also best suited to global observables

All-orders:

Forbid coherent radiation from energy-ordered ensembles of large-angle gluons

Difficulties:

- Logarithms resummed so far only in large- N_c limit
- In general, boundary between the two regions may have arbitrary shape.
- It may depend on the pattern of emissions (e.g. with jet algorithm).

Appleby & Seymour '02, '03

Banfi & Dasgupta '05

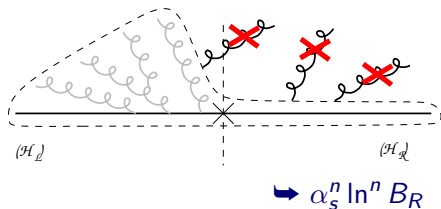
Delenda, A, B & D '06

Resummation of a general non-global observable is tricky.
For time-being CAESAR deals only with global observables.

NB: (most) Monte Carlo's are also best suited to global observables

All-orders:

Forbid coherent radiation from energy-ordered ensembles of large-angle gluons

Difficulties:

- Logarithms resummed so far only in large- N_c limit
- In general, boundary between the two regions may have arbitrary shape.
- It may depend on the pattern of emissions (e.g. with jet algorithm).

Appleby & Seymour '02, '03

Banfi & Dasgupta '05

Delenda, A, B & D '06

Resummation of a general non-global observable is tricky.
For time-being CAESAR deals only with global observables.

NB: (most) Monte Carlo's are also best suited to global observables

Contradiction?

Theoretical calculations are for global observables.

But experiments only have detectors in limited rapidity range.

(Strictly: series of sub-detectors, of worsening quality as rapidity increases)

Model by cut around beam $|\eta| < \eta_{\max}$

➡ Problems with globalness

Take cut as being edge of most forward detector with momentum or energy resolution:

	Tevatron	LHC
η_{\max}	3.5	5.0

Contradiction?

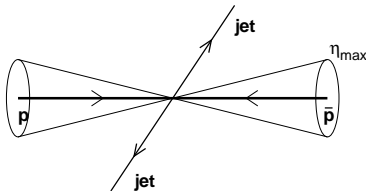
Theoretical calculations are for global observables.

But experiments only have detectors in limited rapidity range.

(Strictly: series of sub-detectors, of worsening quality as rapidity increases)

Model by cut around beam $|\eta| < \eta_{\max}$

➡ Problems with **globalness**



Take cut as being edge of most forward detector with momentum or energy resolution:

	Tevatron	LHC
η_{\max}	3.5	5.0

Contradiction?

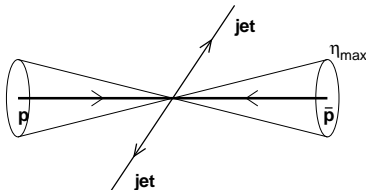
Theoretical calculations are for global observables.

But experiments only have detectors in limited rapidity range.

(Strictly: series of sub-detectors, of worsening quality as rapidity increases)

Model by cut around beam $|\eta| < \eta_{\max}$

➡ Problems with **globalness**



Take cut as being edge of most forward detector with momentum or energy resolution:

	Tevatron	LHC
η_{\max}	3.5	5.0

Select events with central, hard jets (x_1, x_2 not too small), with transverse momentum P_\perp .

From kinematics, emissions (k) near forward detector edges typically have small transverse momentum:

$$k_\perp \sim P_\perp e^{-\eta_0} \ll P_\perp$$

If event-shape value is always sufficiently large that such an emission contributes negligibly, then:

we can ignore rapidity cut & pretend measurement is global

Proceed as follows:

- Calculate distribution without any rapidity cutoff
- Determine smallest 'typical' value of observable
- Check self-consistency: *i.e.* that in comparison, emissions beyond cutoff contribute negligibly.

Select events with central, hard jets (x_1, x_2 not too small), with transverse momentum P_\perp .

From kinematics, emissions (k) near forward detector edges typically have small transverse momentum:

$$k_\perp \sim P_\perp e^{-\eta_0} \ll P_\perp$$

If event-shape value is always sufficiently large that such an emission contributes negligibly, then:

we can ignore rapidity cut & pretend measurement is global

Proceed as follows:

- Calculate distribution without any rapidity cutoff
- Determine smallest 'typical' value of observable
- Check self-consistency: *i.e.* that in comparison, emissions beyond cutoff contribute negligibly.

Select events with central, hard jets (x_1, x_2 not too small), with transverse momentum P_\perp .

From kinematics, emissions (k) near forward detector edges typically have small transverse momentum:

$$k_\perp \sim P_\perp e^{-\eta_0} \ll P_\perp$$

If event-shape value is always sufficiently large that such an emission contributes negligibly, then:

we can ignore rapidity cut & pretend measurement is global

Proceed as follows:

- Calculate distribution without any rapidity cutoff
- Determine smallest 'typical' value of observable
- Check self-consistency: *i.e.* that in comparison, emissions beyond cutoff contribute negligibly.

Select events with central, hard jets (x_1, x_2 not too small), with transverse momentum P_\perp .

From kinematics, emissions (k) near forward detector edges typically have small transverse momentum:

$$k_\perp \sim P_\perp e^{-\eta_0} \ll P_\perp$$

If event-shape value is always sufficiently large that such an emission contributes negligibly, then:

we can ignore rapidity cut & pretend measurement is global

Proceed as follows:

- Calculate distribution without any rapidity cutoff
- Determine smallest 'typical' value of observable
- Check self-consistency: *i.e.* that in comparison, emissions beyond cutoff contribute negligibly.

Results that follow based on this (illustrative) event selection:

- Run longitudinally invariant inclusive k_t jet algorithm (could also use Cambridge/Aachen or SISCone)
- Require hardest jet to have $P_{\perp,1} > P_{\perp,\min} = 50 \text{ GeV}$
- Require two hardest jets to be central $|\eta_1|, |\eta_2| < \eta_c = 0.7$

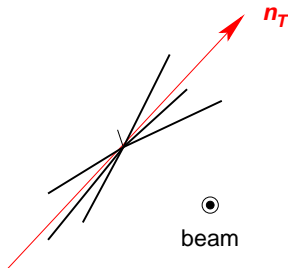
Pure resummed results
no matching to NLO (or even LO)
Shown for Tevatron run II

Some observables are naturally defined in terms of all particles in the event, e.g. *Global Transverse Thrust*

$$T_{\perp,g} \equiv \max_{\vec{n}_T} \frac{\sum_i |\vec{q}_{\perp i} \cdot \vec{n}_T|}{\sum_i q_{\perp i}}, \quad \tau_{\perp,g} = 1 - T_{\perp,g},$$

and *Global Thrust Minor*

$$T_{m,g} \equiv \frac{\sum_i |\vec{q}_i \cdot \vec{n}_m|}{\sum_i q_{\perp i}}, \quad \vec{n}_m \cdot \vec{n}_T = 0$$

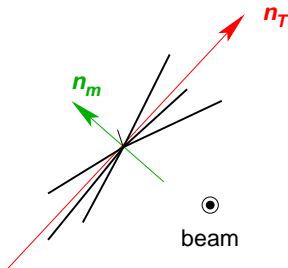


Some observables are naturally defined in terms of all particles in the event, e.g. *Global Transverse Thrust*

$$T_{\perp,g} \equiv \max_{\vec{n}_T} \frac{\sum_i |\vec{q}_{\perp i} \cdot \vec{n}_T|}{\sum_i q_{\perp i}}, \quad \tau_{\perp,g} = 1 - T_{\perp,g},$$

and *Global Thrust Minor*

$$T_{m,g} \equiv \frac{\sum_i |\vec{q}_i \cdot \vec{n}_m|}{\sum_i q_{\perp i}}, \quad \vec{n}_m \cdot \vec{n}_T = 0$$

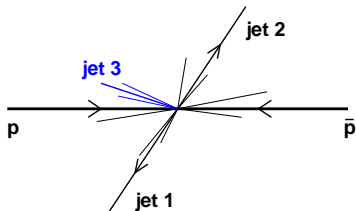
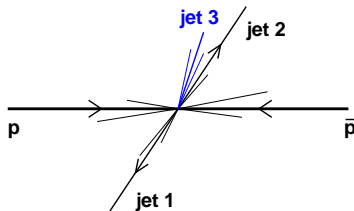


Use *exclusive* long. inv. k_t algorithm: successive recombination of pair with smallest closeness measure d_{kl} , d_{kB} :

$$d_{kB} = q_{\perp k}^2, \quad d_{kl} = \min\{q_{\perp k}^2, q_{\perp l}^2\} ((\eta_k - \eta_l)^2 + (\phi_k - \phi_l)^2).$$

Define $d^{(n)}$ as smallest d_{kl} , d_{kB} when only n pseudo-jets left. Examine (normalised) 3-jet resolution threshold

$$y_{23} = \frac{1}{(E_{\perp,1} + E_{\perp,2})^2} d^{(3)}$$



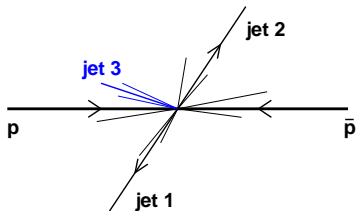
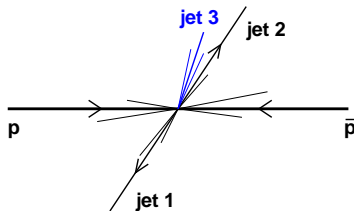
Generalisation of 3-jet cross section

Use *exclusive* long. inv. k_t algorithm: successive recombination of pair with smallest closeness measure d_{kl} , d_{kB} :

$$d_{kB} = q_{\perp k}^2, \quad d_{kl} = \min\{q_{\perp k}^2, q_{\perp l}^2\} ((\eta_k - \eta_l)^2 + (\phi_k - \phi_l)^2).$$

Define $d^{(n)}$ as smallest d_{kl} , d_{kB} when only n pseudo-jets left. Examine (normalised) 3-jet resolution threshold

$$y_{23} = \frac{1}{(E_{\perp,1} + E_{\perp,2})^2} \max_{n \geq 3} \{d^{(n)}\},$$



Generalisation of 3-jet cross section

Probability $P(v)$ that event shape is smaller than some value v :

$$P(v) = \exp \left[-G_{12} \frac{\alpha_s L^2}{2\pi} + \dots \right], \quad L = \ln \frac{1}{v}$$

Ev. Shp.	G_{12}
$\tau_{\perp,g}$	$2C_B + C_J$
$T_{m,g}$	$2C_B + 2C_J$
y_{23}	$\frac{1}{2}C_B + \frac{1}{2}C_J$

C_B = total colour of Beam partons

C_J = total colour of Jet partons

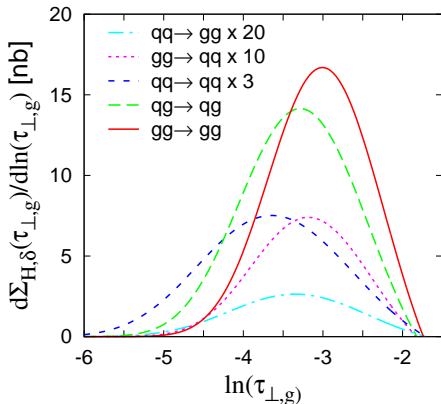
Probability $P(v)$ that event shape is smaller than some value v :

$$P(v) = \exp \left[-G_{12} \frac{\alpha_s L^2}{2\pi} + \dots \right], \quad L = \ln \frac{1}{v}$$

Ev.Shp.	G_{12}
$\tau_{\perp,g}$	$2C_B + C_J$
$T_{m,g}$	$2C_B + 2C_J$
y_{23}	$\frac{1}{2}C_B + \frac{1}{2}C_J$

C_B = total colour of Beam partons

C_J = total colour of Jet partons



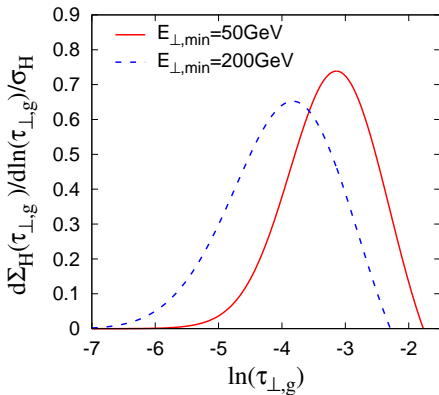
Probability $P(v)$ that event shape is smaller than some value v :

$$P(v) = \exp \left[-G_{12} \frac{\alpha_s L^2}{2\pi} + \dots \right], \quad L = \ln \frac{1}{v}$$

Ev. Shp.	G_{12}
$\tau_{\perp,g}$	$2C_B + C_J$
$T_{m,g}$	$2C_B + 2C_J$
y_{23}	$\frac{1}{2}C_B + \frac{1}{2}C_J$

C_B = total colour of Beam partons

C_J = total colour of Jet partons



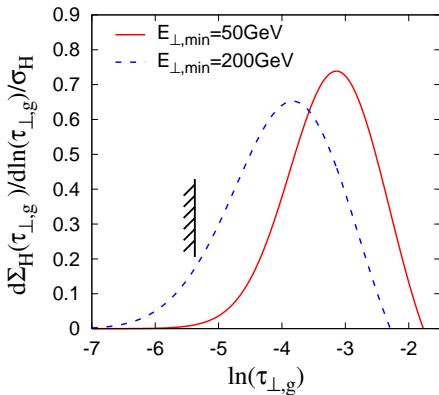
Probability $P(v)$ that event shape is smaller than some value v :

$$P(v) = \exp \left[-G_{12} \frac{\alpha_s L^2}{2\pi} + \dots \right], \quad L = \ln \frac{1}{v}$$

Ev. Shp.	G_{12}
$\tau_{\perp,g}$	$2C_B + C_J$
$T_{m,g}$	$2C_B + 2C_J$
y_{23}	$\frac{1}{2}C_B + \frac{1}{2}C_J$

C_B = total colour of Beam partons

C_J = total colour of Jet partons



Beam cut: $\tau_{\perp,g} \gtrsim 0.15e^{-\eta_{\max}}$

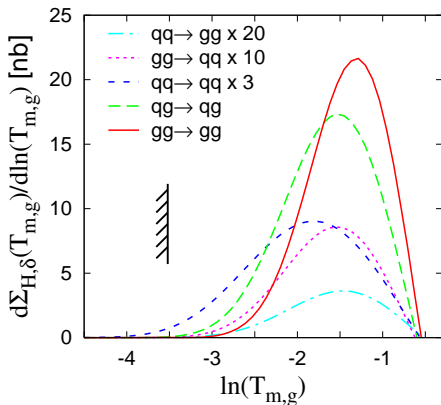
Probability $P(v)$ that event shape is smaller than some value v :

$$P(v) = \exp \left[-G_{12} \frac{\alpha_s L^2}{2\pi} + \dots \right], \quad L = \ln \frac{1}{v}$$

Ev. Shp.	G_{12}
$\tau_{\perp,g}$	$2C_B + C_J$
$T_{m,g}$	$2C_B + 2C_J$
y_{23}	$\frac{1}{2}C_B + \frac{1}{2}C_J$

C_B = total colour of Beam partons

C_J = total colour of Jet partons



Beam cut: $T_{m,g} \gtrsim e^{-\eta_{\max}}$

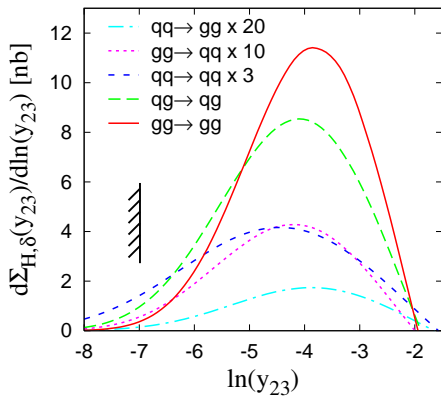
Probability $P(v)$ that event shape is smaller than some value v :

$$P(v) = \exp \left[-G_{12} \frac{\alpha_s L^2}{2\pi} + \dots \right], \quad L = \ln \frac{1}{v}$$

Ev. Shp.	G_{12}
$\tau_{\perp, g}$	$2C_B + C_J$
$T_{m, g}$	$2C_B + 2C_J$
y_{23}	$\frac{1}{2}C_B + \frac{1}{2}C_J$

C_B = total colour of Beam partons

C_J = total colour of Jet partons



Beam cut: $y_{23} \gtrsim e^{-2\eta_{\max}}$ [because $y_{23} \sim k_t^2$]

Forward-suppressed observables

Divide event into central region (\mathcal{C} , say $|\eta| < 1.1$) and rest of event ($\bar{\mathcal{C}}$).

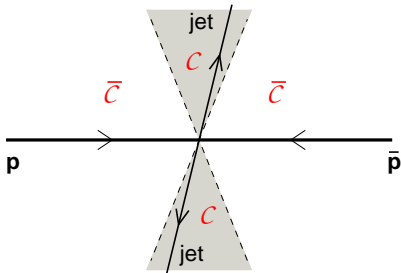
[NB: \exists considerable freedom in definition of \mathcal{C} : e.g. can also be two hardest jets]

Define central \perp mom., and rapidity:

$$Q_{\perp, \mathcal{C}} = \sum_{i \in \mathcal{C}} q_{\perp i}, \quad \eta_{\mathcal{C}} = \frac{1}{Q_{\perp, \mathcal{C}}} \sum_{i \in \mathcal{C}} \eta_i q_{\perp i}$$

and an *exponentially suppressed forward term*,

$$\mathcal{E}_{\bar{\mathcal{C}}} = \frac{1}{Q_{\perp, \mathcal{C}}} \sum_{i \notin \mathcal{C}} q_{\perp i} e^{-|\eta_i - \eta_{\mathcal{C}}|}.$$



Define a non-global event-shape in \mathcal{C} . Then add on $\mathcal{E}_{\bar{\mathcal{C}}}$.

Result is a global event shape, with suppressed sensitivity to forward region.

Forward-suppressed observables

Divide event into central region (\mathcal{C} , say $|\eta| < 1.1$) and rest of event ($\bar{\mathcal{C}}$).

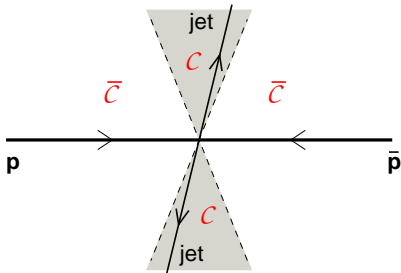
[NB: \exists considerable freedom in definition of \mathcal{C} : e.g. can also be two hardest jets]

Define central \perp mom., and rapidity:

$$Q_{\perp, \mathcal{C}} = \sum_{i \in \mathcal{C}} q_{\perp i}, \quad \eta_{\mathcal{C}} = \frac{1}{Q_{\perp, \mathcal{C}}} \sum_{i \in \mathcal{C}} \eta_i q_{\perp i}$$

and an *exponentially suppressed forward term*,

$$\mathcal{E}_{\bar{\mathcal{C}}} = \frac{1}{Q_{\perp, \mathcal{C}}} \sum_{i \notin \mathcal{C}} q_{\perp i} e^{-|\eta_i - \eta_{\mathcal{C}}|}.$$



Define a non-global event-shape in \mathcal{C} . Then add on $\mathcal{E}_{\bar{\mathcal{C}}}$.

Result is a global event shape, with suppressed sensitivity to forward region.

- Split \mathcal{C} into two pieces: *Up, Down*
- Define *jet masses* for each

$$\rho_{X,\mathcal{C}} \equiv \frac{1}{Q_{\perp,\mathcal{C}}^2} \left(\sum_{i \in \mathcal{C}_X} q_i \right)^2, \quad X = U, D,$$

Define sum and heavy-jet masses

$$\rho_{S,\mathcal{C}} \equiv \rho_{U,\mathcal{C}} + \rho_{D,\mathcal{C}}, \quad \rho_{H,\mathcal{C}} \equiv \max\{\rho_{U,\mathcal{C}}, \rho_{D,\mathcal{C}}\},$$

Define global extension, with extra forward-suppressed term

$$\rho_{S,\mathcal{E}} \equiv \rho_{S,\mathcal{C}} + \mathcal{E}_{\bar{\mathcal{C}}}, \quad \rho_{H,\mathcal{E}} \equiv \rho_{H,\mathcal{C}} + \mathcal{E}_{\bar{\mathcal{C}}}.$$

- Similarly: *total and wide jet-broadenings*

$$B_{T,\mathcal{E}} \equiv B_{T,\mathcal{C}} + \mathcal{E}_{\bar{\mathcal{C}}}, \quad B_{W,\mathcal{E}} \equiv B_{W,\mathcal{C}} + \mathcal{E}_{\bar{\mathcal{C}}}.$$

$$P(v) = \exp \left[-G_{12} \frac{\alpha_s L^2}{2\pi} + \dots \right], \quad L = \ln \frac{1}{v}$$

Ev.Shp.	G_{12}
$\rho_{S,\mathcal{E}}$	$C_B + C_J$
$\rho_{H,\mathcal{E}}$	$C_B + C_J$
$B_{T,\mathcal{E}}$	$C_B + 2C_J$
$B_{W,\mathcal{E}}$	$C_B + 2C_J$
\vdots	\vdots

C_B = total colour of Beam partons

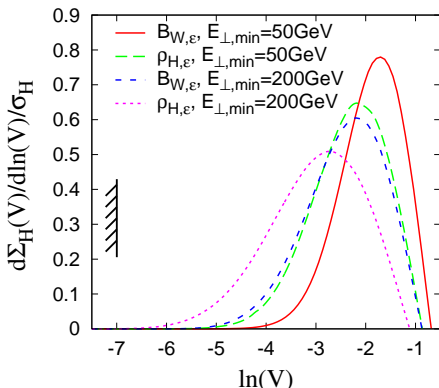
C_J = total colour of Jet partons

$$P(v) = \exp \left[-G_{12} \frac{\alpha_s L^2}{2\pi} + \dots \right], \quad L = \ln \frac{1}{v}$$

Ev. Shp.	G_{12}
$\rho_{S,\mathcal{E}}$	$C_B + C_J$
$\rho_{H,\mathcal{E}}$	$C_B + C_J$
$B_{T,\mathcal{E}}$	$C_B + 2C_J$
$B_{W,\mathcal{E}}$	$C_B + 2C_J$
\vdots	\vdots

C_B = total colour of Beam partons

C_J = total colour of Jet partons



Beam cuts: $B_{X,\mathcal{E}}, \rho_{X,\mathcal{E}} \gtrsim e^{-2\eta_{\max}}$ [because $\mathcal{E}_{\bar{c}} \sim k_t e^{-|\eta|}$]

By momentum conservation

$$\sum_{i \in \mathcal{C}} \vec{q}_{\perp i} = - \sum_{i \notin \mathcal{C}} \vec{q}_{\perp i}$$

Use central particles to define *recoil term*, which is *indirectly sensitive* to non-central emissions

$$\mathcal{R}_{\perp, \mathcal{C}} \equiv \frac{1}{Q_{\perp, \mathcal{C}}} \left| \sum_{i \in \mathcal{C}} \vec{q}_{\perp i} \right|,$$

Define event shapes exclusively in terms of *central particles*:

$$\rho_{X, \mathcal{R}} \equiv \rho_{X, \mathcal{C}} + \mathcal{R}_{\perp, \mathcal{C}}, \quad B_{X, \mathcal{R}} \equiv B_{X, \mathcal{C}} + \mathcal{R}_{\perp, \mathcal{C}}, \dots$$

These observables are *indirectly global*

First studied at HERA (B_{zE} broadening)

CAESAR resummation works for observables having *direct exponentiation*:

$$P(v) = e^{Lg_1(\alpha_s L) + g_2(\alpha_s L) + \dots}$$

For recoil observables, exponentiation holds fully only after Fourier & other integral transforms (**generalised b -space resummation**).

Manifestation: NLLs ($g_2(\alpha_s L)$) diverge at some $\alpha_s L \sim 1$.

Consequently, cannot extend distribution to $v = 0$ — must cut before divergence.

CAESAR resummation works for observables having *direct exponentiation*:

$$P(v) = e^{Lg_1(\alpha_s L) + g_2(\alpha_s L) + \dots}$$

For recoil observables, exponentiation holds fully only after Fourier & other integral transforms (**generalised b -space resummation**).

Manifestation: NLLs ($g_2(\alpha_s L)$) diverge at some $\alpha_s L \sim 1$.

Consequently, cannot extend distribution to $v = 0$ — must cut before divergence.

CAESAR resummation works for observables having *direct exponentiation*:

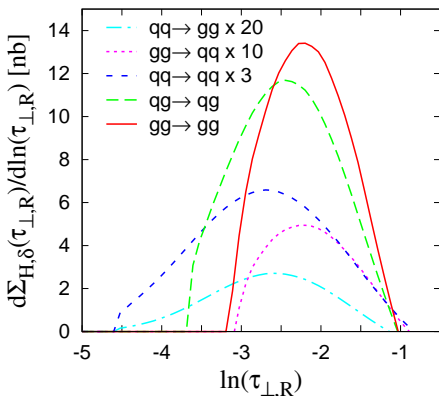
$$P(v) = e^{Lg_1(\alpha_s L) + g_2(\alpha_s L) + \dots}$$

For recoil observables, exponentiation holds fully only after Fourier & other integral transforms (*generalised b -space resummation*).

Manifestation: NLLs ($g_2(\alpha_s L)$) diverge at some $\alpha_s L \sim 1$.

Consequently, cannot extend distribution to $v = 0$ — must cut before divergence.

recoil transverse thrust



Quite large effect: $\sim 15\%$ of X-sct is beyond cutoff

CAESAR resummation works for observables having *direct exponentiation*:

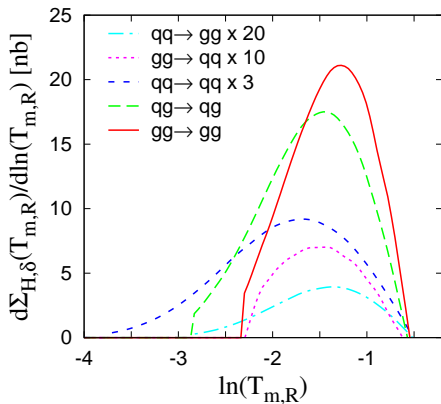
$$P(v) = e^{Lg_1(\alpha_s L) + g_2(\alpha_s L) + \dots}$$

For recoil observables, exponentiation holds fully only after Fourier & other integral transforms (*generalised b -space resummation*).

Manifestation: NLLs ($g_2(\alpha_s L)$) diverge at some $\alpha_s L \sim 1$.

Consequently, cannot extend distribution to $v = 0$ — must cut before divergence.

recoil thrust minor



Moderate effect: few % of X-sct is beyond cutoff

Event-shape	Impact of η_{\max}	Resummation breakdown	Underlying Event	Jet hadronisation
$\tau_{\perp,g}$	tolerable	none	$\sim \eta_{\max}/Q$	$\sim 1/Q$
$T_{m,g}$	tolerable	none	$\sim \eta_{\max}/Q$	$\sim 1/(\sqrt{\alpha_s}Q)$
y_{23}	tolerable	none	$\sim \sqrt{y_{23}}/Q$	$\sim \sqrt{y_{23}}/Q$
$\tau_{\perp,\mathcal{E}}, \rho_{X,\mathcal{E}}$	negligible	none	$\sim 1/Q$	$\sim 1/Q$
$B_{X,\mathcal{E}}$	negligible	none	$\sim 1/Q$	$\sim 1/(\sqrt{\alpha_s}Q)$
$T_{m,\mathcal{E}}$	negligible	serious	$\sim 1/Q$	$\sim 1/(\sqrt{\alpha_s}Q)$
$y_{23,\mathcal{E}}$	negligible	none	$\sim 1/Q$	$\sim \sqrt{y_{23}}/Q$
$\tau_{\perp,\mathcal{R}}, \rho_{X,\mathcal{R}}$	none	serious	$\sim 1/Q$	$\sim 1/Q$
$T_{m,\mathcal{R}}, B_{X,\mathcal{R}}$	none	tolerable	$\sim 1/Q$	$\sim 1/(\sqrt{\alpha_s}Q)$
$y_{23,\mathcal{R}}$	none	intermediate	$\sim \sqrt{y_{23}}/Q$	$\sim \sqrt{y_{23}}/Q$

NB: there may be surprises after more detailed study, e.g. matching to NLO...

Grey entries are definitely subject to uncertainty

Note complementarity between observables

Summary of observables

Event-shape	Impact of η_{\max}	Resummation breakdown	Underlying Event	Jet hadronisation
$\tau_{\perp,g}$	tolerable	none	$\sim \eta_{\max}/Q$	$\sim 1/Q$
$T_{m,g}$	tolerable	none	$\sim \eta_{\max}/Q$	$\sim 1/(\sqrt{\alpha_s}Q)$
y_{23}	tolerable	none	$\sim \sqrt{y_{23}}/Q$	$\sim \sqrt{y_{23}}/Q$
$\tau_{\perp,\mathcal{E}}, \rho_{X,\mathcal{E}}$	negligible	none	$\sim 1/Q$	$\sim 1/Q$
$B_{X,\mathcal{E}}$	negligible	none	$\sim 1/Q$	$\sim 1/(\sqrt{\alpha_s}Q)$
$T_{m,\mathcal{E}}$	negligible	serious	$\sim 1/Q$	$\sim 1/(\sqrt{\alpha_s}Q)$
$y_{23,\mathcal{E}}$	negligible	none	$\sim 1/Q$	$\sim \sqrt{y_{23}}/Q$
$\tau_{\perp,\mathcal{R}}, \rho_{X,\mathcal{R}}$	none	serious	$\sim 1/Q$	$\sim 1/Q$
$T_{m,\mathcal{R}}, B_{X,\mathcal{R}}$	none	tolerable	$\sim 1/Q$	$\sim 1/(\sqrt{\alpha_s}Q)$
$y_{23,\mathcal{R}}$	none	intermediate	$\sim \sqrt{y_{23}}/Q$	$\sim \sqrt{y_{23}}/Q$

NB: there may be surprises after more detailed study, e.g. matching to NLO...

Grey entries are definitely subject to uncertainty

Note complementarity between observables

First NLO+NLL+1/Q matching for multi-jet ev. shapes

Banfi & Zanderighi *prelim.*

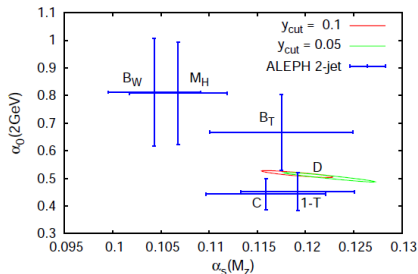
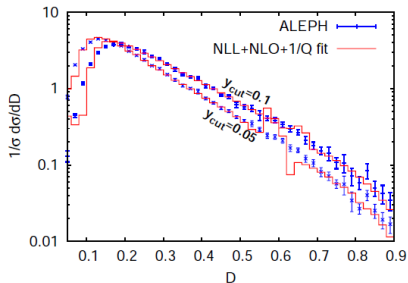
e^+e^- D -parameter and thrust minor

- confirms that framework can work in multi-jet context
- progress on road to full matching in pp

Tests of power corrections for D and T_m

- Select 3-jet events with $y_3 > y_{\text{cut}}$
- Differential distributions obtained with CAESAR at $Q = 91.2 \text{ GeV}$

[AB, G. Salam, G. Zanderighi hep-ph/0407286]

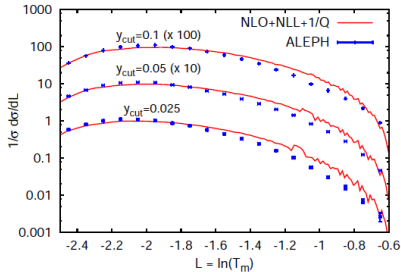


- D -parameter: first ever α_s - α_0 fits in a three-jet event shapes!
- Good fits only for $D > 0.2$: $\chi^2/\text{d.o.f.}(y_{\text{cut}} = 0.1) = 12/20$
⇒ Small- D region: shape function or large subleading logs?

Tests of power corrections for D and T_m

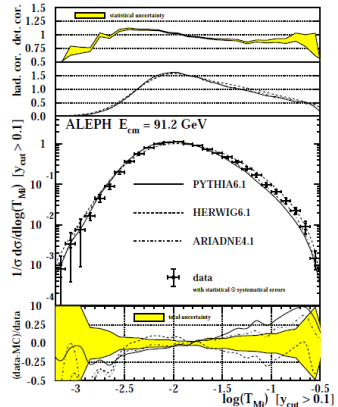
- Select 3-jet events with $y_3 > y_{\text{cut}}$
- Differential distributions obtained with CAESAR at $Q = 91.2 \text{ GeV}$

[AB, G. Salam, G. Zanderighi hep-ph/0407286]



Thrust minor PC should be positive

MC say PC are negative at large T_m
 \Rightarrow PC from 4-jet configurations?



Key ingredient in all resummations is **coherence**

Large-angle reals/virtuals not affected by small-angle emissions

Implies: interjet energy-flow type resummations involve single-logs, $\alpha_s^n L^n$

Calculation by Forshaw, Kyrieleis & Seymour '06 finds $\alpha_s^4 L^5 \times 1/N_c^2$

-
- If these terms exist they could affect resummations for $\tau_{\perp,\mathcal{E}}$, $\rho_{\mathcal{X},\mathcal{E}}$, $B_{\mathcal{X},\mathcal{E}}$, $y_{23,\mathcal{E}}$.
 \equiv Observables with η dependence in forward regions
 - FKS paper alone is not sufficient to prove existence — coefficient of result depends on (arbitrary) choice of ordering variable.
 - FKS find they are numerically small (N_c suppressed phase interference) — perhaps not serious in practice even if conceptually important

One should keep an eye on this issue

Key ingredient in all resummations is **coherence**

Large-angle reals/virtuals not affected by small-angle emissions

Implies: interjet energy-flow type resummations involve single-logs, $\alpha_s^n L^n$

Calculation by Forshaw, Kyrieleis & Seymour '06 finds $\alpha_s^4 L^5 \times 1/N_c^2$

-
- If these terms exist they could affect resummations for $\tau_{\perp,\mathcal{E}}$, $\rho_{\mathcal{X},\mathcal{E}}$, $B_{\mathcal{X},\mathcal{E}}$, $y_{23,\mathcal{E}}$.
 \equiv Observables with η dependence in forward regions
 - FKS paper alone is not sufficient to prove existence — coefficient of result depends on (arbitrary) choice of ordering variable.
 - FKS find they are numerically small (N_c suppressed phase interference) — perhaps not serious in practice even if conceptually important

One should keep an eye on this issue

Key ingredient in all resummations is **coherence**

Large-angle reals/virtuals not affected by small-angle emissions

Implies: interjet energy-flow type resummations involve single-logs, $\alpha_s^n L^n$

Calculation by Forshaw, Kyrieleis & Seymour '06 finds $\alpha_s^4 L^5 \times 1/N_c^2$

-
- If these terms exist they could affect resummations for $\tau_{\perp,\mathcal{E}}$, $\rho_{\mathcal{X},\mathcal{E}}$, $B_{\mathcal{X},\mathcal{E}}$, $y_{23,\mathcal{E}}$.
 \equiv Observables with η dependence in forward regions
 - FKS paper alone is not sufficient to prove existence — coefficient of result depends on (arbitrary) choice of ordering variable.
 - FKS find they are numerically small (N_c suppressed phase interference) — perhaps not serious in practice even if conceptually important

One should keep an eye on this issue

Main difficulty:

$$\sigma_{qq \rightarrow qq}^{NLO} e^{-2C_F \alpha_s L^2 / 2\pi} + \sigma_{gg \rightarrow gg}^{NLO} e^{-2C_A \alpha_s L^2 / 2\pi} + \dots$$

$$\neq \frac{\sigma^{NLO}}{\sigma^{LO}} \left[\sigma_{qq \rightarrow qq}^{LO} e^{-2C_F \alpha_s L^2 / 2\pi} + \sigma_{gg \rightarrow gg}^{LO} e^{-2C_A \alpha_s L^2 / 2\pi} + \dots \right]$$

In order to guarantee $\alpha_s^n L^{2n-2}$ (NNLL in expanded result), part at least of matching must be done channel by channel. Never an issue before

Problems:

- Flavour channel definition not IR safe with normal jet algs
 Use special *flavour- k_t* algorithm, Banfi, GPS & Zanderighi '06
- NLO Monte Carlos for pp do not provide information on flavour of partons. Can be disentangled in NLOJET++ (1 month of hard work)
- We got distracted: used flavour- k_t alg. to *reduce uncertainty on b -jet spectrum from 40 – 60% to 10 – 20%* Banfi, GPS & Zanderighi '07

Main difficulty:

$$\begin{aligned} & \sigma_{qq \rightarrow qq}^{NLO} e^{-2C_F \alpha_s L^2 / 2\pi} + \sigma_{gg \rightarrow gg}^{NLO} e^{-2C_A \alpha_s L^2 / 2\pi} + \dots \\ & \neq \frac{\sigma^{NLO}}{\sigma^{LO}} \left[\sigma_{qq \rightarrow qq}^{LO} e^{-2C_F \alpha_s L^2 / 2\pi} + \sigma_{gg \rightarrow gg}^{LO} e^{-2C_A \alpha_s L^2 / 2\pi} + \dots \right] \end{aligned}$$

In order to guarantee $\alpha_s^n L^{2n-2}$ (NNLL in expanded result), part at least of matching must be done channel by channel. Never an issue before

Problems:

- Flavour channel definition not IR safe with normal jet algs
Use special *flavour- k_t* algorithm, Banfi, GPS & Zanderighi '06
- NLO Monte Carlos for pp do not provide information on flavour of partons. Can be disentangled in NLOJET++ (1 month of hard work)
- We got distracted: used flavour- k_t alg. to *reduce uncertainty on b -jet spectrum from 40 – 60% to 10 – 20%* Banfi, GPS & Zanderighi '07

Main difficulty:

$$\sigma_{qq \rightarrow qq}^{NLO} e^{-2C_F \alpha_s L^2 / 2\pi} + \sigma_{gg \rightarrow gg}^{NLO} e^{-2C_A \alpha_s L^2 / 2\pi} + \dots$$

$$\neq \frac{\sigma^{NLO}}{\sigma^{LO}} \left[\sigma_{qq \rightarrow qq}^{LO} e^{-2C_F \alpha_s L^2 / 2\pi} + \sigma_{gg \rightarrow gg}^{LO} e^{-2C_A \alpha_s L^2 / 2\pi} + \dots \right]$$

In order to guarantee $\alpha_s^n L^{2n-2}$ (NNLL in expanded result), part at least of matching must be done channel by channel. Never an issue before

Problems:

- Flavour channel definition not IR safe with normal jet algs
Use special *flavour- k_t* algorithm, Banfi, GPS & Zanderighi '06
- NLO Monte Carlos for pp do not provide information on flavour of partons.
Can be disentangled in NLOJET++ (1 month of hard work)
- We got distracted: used flavour- k_t alg. to *reduce uncertainty on b -jet spectrum from 40 – 60% to 10 – 20%* Banfi, GPS & Zanderighi '07

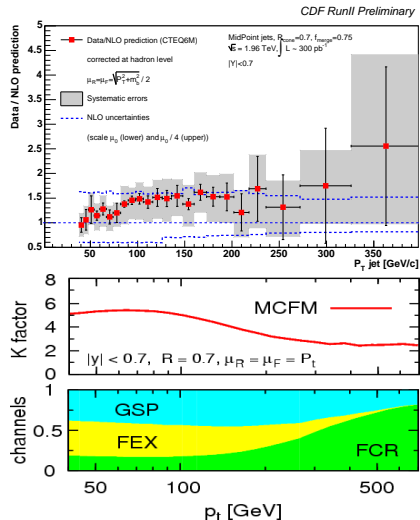
Main difficulty:

$$\begin{aligned} & \sigma_{qq \rightarrow qq}^{NLO} e^{-2C_F \alpha_s L^2 / 2\pi} + \sigma_{gg \rightarrow gg}^{NLO} e^{-2C_A \alpha_s L^2 / 2\pi} + \dots \\ & \neq \frac{\sigma^{NLO}}{\sigma^{LO}} \left[\sigma_{qq \rightarrow qq}^{LO} e^{-2C_F \alpha_s L^2 / 2\pi} + \sigma_{gg \rightarrow gg}^{LO} e^{-2C_A \alpha_s L^2 / 2\pi} + \dots \right] \end{aligned}$$

In order to guarantee $\alpha_s^n L^{2n-2}$ (NNLL in expanded result), part at least of matching must be done channel by channel. Never an issue before

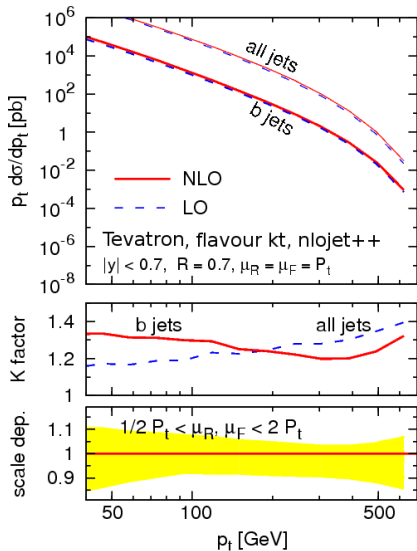
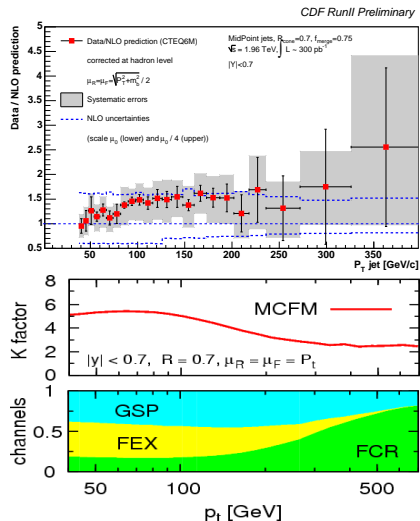
Problems:

- Flavour channel definition not IR safe with normal jet algs
Use special *flavour- k_t* algorithm, Banfi, GPS & Zanderighi '06
- NLO Monte Carlos for pp do not provide information on flavour of partons.
Can be disentangled in NLOJET++ (1 month of hard work)
- We got distracted: used flavour- k_t alg. to *reduce uncertainty on b -jet spectrum from 40 – 60% to 10 – 20%*
Banfi, GPS & Zanderighi '07



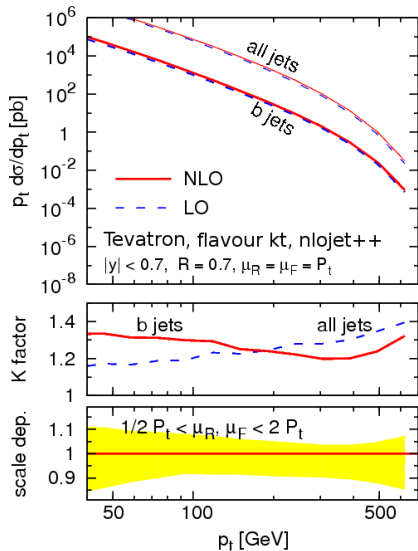
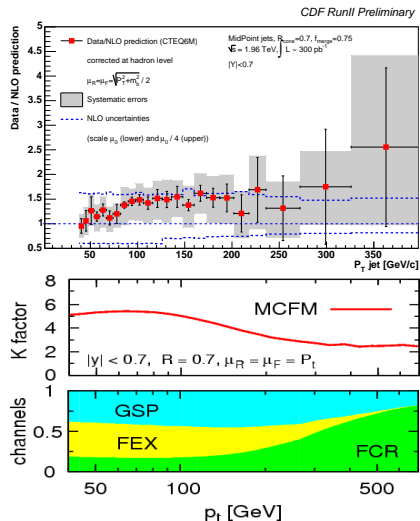
Now that this is done, we will be moving back towards the matching. . .

Flavour separation only to α_s^3 , but enough for $\alpha_s^2 L^{2n-2}$



Now that this is done, we will be moving back towards the matching. . .

Flavour separation only to α_s^3 , but enough for $\alpha_s^2 L^{2n-2}$



Now that this is done, we will be moving back towards the matching. . .

Flavour separation only to α_s^3 , but enough for $\alpha_s^2 L^{2n-2}$

Groundwork

- Important that multijet event shapes also be studied in DIS and e^+e^- .
 - Measurements available from LEP and HERA.
 - Theoretical comparisons now appearing — automation facilitates this.

Hadron-collider specificities

- New domain for “rigorous” QCD studies:
 - non-perturbative: **underlying event**
 - perturbative: Stony Brook **soft colour resummation**
 - surprises: **super-leading logs?**
- Tension between theoretical simplicity (**globalness**) and experimental measurability (**limited rapidity**) — can be resolved

Next step: matching to NLO

- Technology now exists for decent matching.
 - flavour-separated NLOJET++, flavour jet algs
- Concrete matching still to be done.

Further info: hep-ph/0407287 and <http://qcd-caesar.org>

Groundwork

- Important that multijet event shapes also be studied in DIS and e^+e^- .
 - Measurements available from LEP and HERA.
 - Theoretical comparisons now appearing — automation facilitates this.

Hadron-collider specificities

- New domain for “rigorous” QCD studies:
 - non-perturbative: **underlying event**
 - perturbative: Stony Brook **soft colour resummation**
 - surprises: **super-leading logs?**
- Tension between theoretical simplicity (**globalness**) and experimental measurability (**limited rapidity**) — **can be resolved**

Next step: matching to NLO

- Technology now exists for decent matching.
 - flavour-separated NLOJET++, flavour jet algs
- Concrete matching still to be done.

Further info: hep-ph/0407287 and <http://qcd-caesar.org>

Groundwork

- Important that multijet event shapes also be studied in DIS and e^+e^- .
 - Measurements available from LEP and HERA.
 - Theoretical comparisons now appearing — automation facilitates this.

Hadron-collider specificities

- New domain for “rigorous” QCD studies:
 - non-perturbative: **underlying event**
 - perturbative: Stony Brook **soft colour resummation**
 - surprises: **super-leading logs?**
- Tension between theoretical simplicity (**globalness**) and experimental measurability (**limited rapidity**) — **can be resolved**

Next step: matching to NLO

- Technology now exists for decent matching.
flavour-separated NLOJET++, flavour jet algs
- Concrete matching still to be done.

Further info: hep-ph/0407287 and <http://qcd-caesar.org>

Groundwork

- Important that multijet event shapes also be studied in DIS and e^+e^- .
 - Measurements available from LEP and HERA.
 - Theoretical comparisons now appearing — automation facilitates this.

Hadron-collider specificities

- New domain for “rigorous” QCD studies:
 - non-perturbative: **underlying event**
 - perturbative: Stony Brook **soft colour resummation**
 - surprises: **super-leading logs?**
- Tension between theoretical simplicity (**globalness**) and experimental measurability (**limited rapidity**) — **can be resolved**

Next step: matching to NLO

- Technology now exists for decent matching.
flavour-separated NLOJET++, flavour jet algs
- Concrete matching still to be done.

Further info: hep-ph/0407287 and <http://qcd-caesar.org>