

# Theoretical aspects of jet-finding

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*Based on work with Cacciari (LPTHE) & Soyez (BNL), and also with Butterworth (ATLAS UCL), Dasgupta (Manchester), Davison (ATLAS UCL), Magnea (Turin), Rojo (LPTHE), Rubin (LPTHE)*

**Aim:** to provide a reminder/intro to the “basics” of jet-finding, with a couple of bleeding-edge subjects at the end.

▶ General considerations

*Some points that deserve to be at the back of your mind as you carry out the challenging technical work to prepare ATLAS to measure jets.*

▶ Modern jet definitions

*Cone algorithms (and SISCone)*

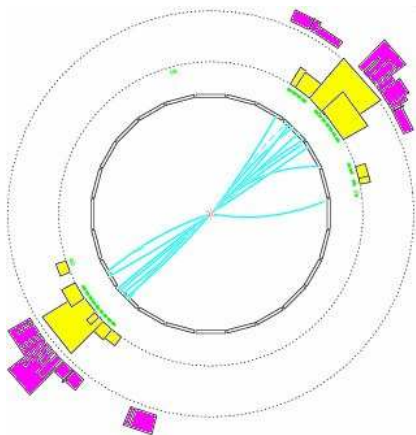
*Sequential recombination algorithms ( $k_t$ , etc.)*

*Learning to compare them*

▶ Jets at work

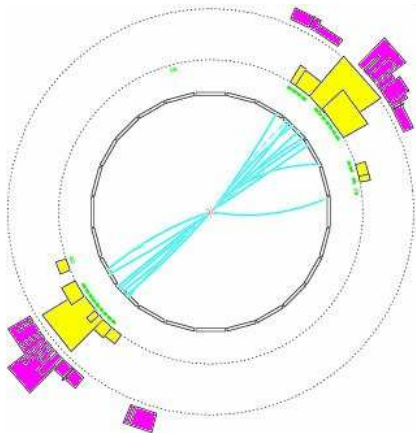
*Pileup characterisation and subtraction*

*A new jet-based Higgs search-channel*

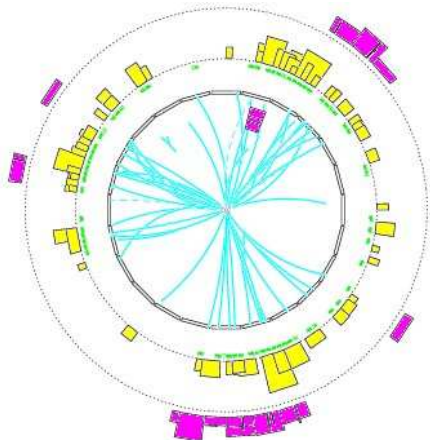


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Clearly(?) 2 jets here

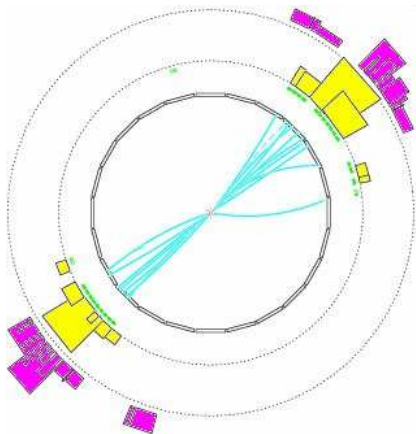
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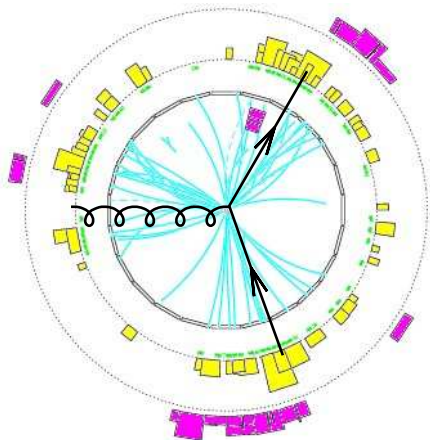
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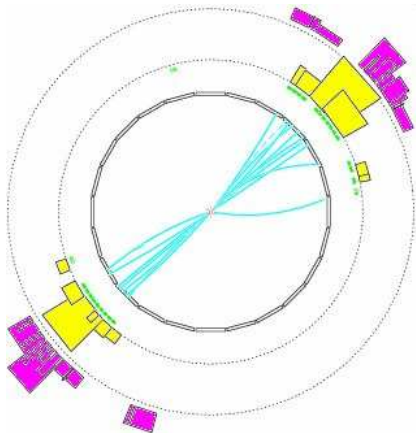
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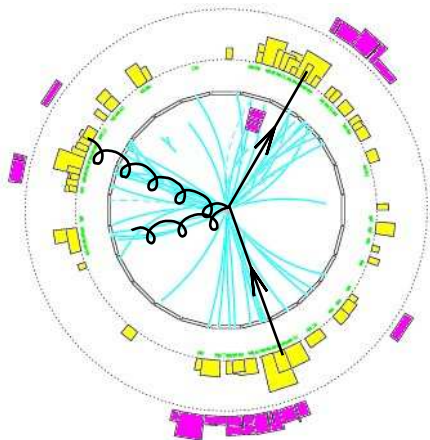
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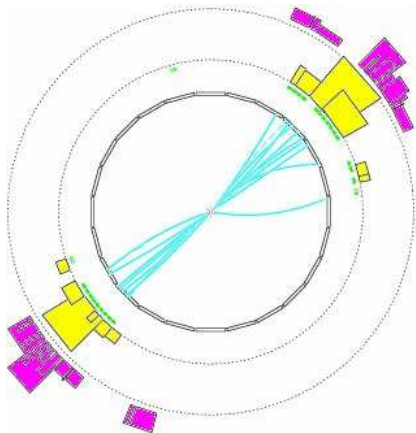
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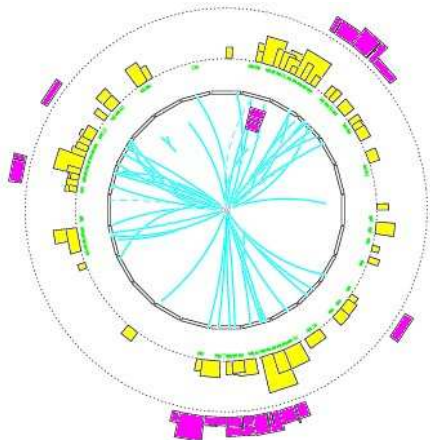
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▶ *A jet definition is a fully specified set of rules for projecting information from 100's of hadrons, onto a handful of parton-like objects:*

- ▶ or project 1000's of calorimeter towers
- ▶ or project dozens of (showered) partons
- ▶ or project a handful of (unshowered) partons

Idea of such a projection: Sterman & Weinberg '77

▶ Resulting objects (jets) used for many things, e.g. :

- ▶ reconstructing decaying massive particles
- ▶ constraining proton structure
- ▶ as a theoretical tool to attribute structure to an event

e.g. top  $\rightarrow$  3 jets

E.g. in CKKW matching

▶ You *lose much information* in projecting event onto jet-like structure:

- ▶ Sometimes information you had no idea how to use
- ▶ Sometimes information you may not trust, or of no relevance



The construction of a jet is unavoidably ambiguous. On at least two fronts:

1. which particles get put together into a common jet?      Jet algorithm  
+ parameters
2. how do you combine their momenta?      Recombination scheme  
Most commonly used: direct 4-vector sums ( $E$ -scheme)

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cf. Les Houches '07 nomenclature accord

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- ▶ Physical results (particle discovery, masses, PDFs, coupling) should be independent of your choice of jet definition
  - a bit like renormalisation scale/scheme invariance
  - Tests independence on modelling of radiation, hadronisation, etc.
- ▶ Sometimes there may be a good reason why one jet definition is more optimal than another
  - In such cases you should understand why
- ▶ But in general there is no single universal 'best' jet definition
  - How best to look at an event depends on what you're trying to see

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Much fine detail  
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photograph it  
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focus=40cm.

keep focus at 40cm

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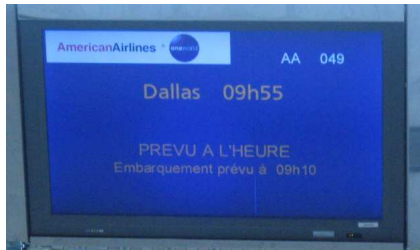
# “Jetography” like photography



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Jets should be **invariant** with respect to certain modifications of the event:

- ▶ infrared (IR) emission e.g. 1 GeV particles wrt 1 TeV jets
- ▶ collinear (C) splitting

This was one of the key ideas introduced by Sterman & Weinberg in '77.

## Why?

- ▶ They happen randomly, quantum-mechanically, all the time in QCD.
- ▶ IR/C sensitivity  $\leftrightarrow$  lose real-virtual cancellation in NLO/NNLO QCD calculations  $\rightarrow$  meaningless, divergent results
- ▶ Hadron-level IR/C modifications are fundamentally non-perturbative
- ▶ Detectors resolve neither full collinear nor full infrared event structure

Requirement known as: **infrared and collinear (IRC) safety**



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The majority of algorithms map each particle into (at most) one jet.

## Cone

- ▶ top-down
- ▶ centred around idea of an ‘invariant’, directed energy flow

## Sequential pairwise recombination ( $k_t$ , etc.)

- ▶ bottom-up
- ▶ successively undoes QCD branching

## Others

- ▶ “Optimal” Jet Finder (OJF): weight  $w_{iJ}$  for each particle  $i$  to be in jet  $J$
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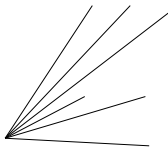
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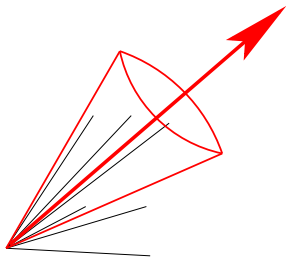
Many cone algs have two main steps:

- ▶ Find some/all stable cones
  - ≡ cone pointing in same direction as the momentum of its contents
- ▶ Resolve cases of overlapping stable cones
  - By running a 'split-merge' procedure [Blazey et al. '00 (Run II jet physics)]



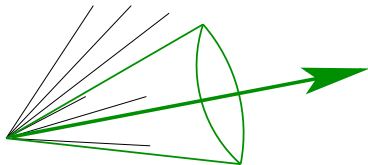
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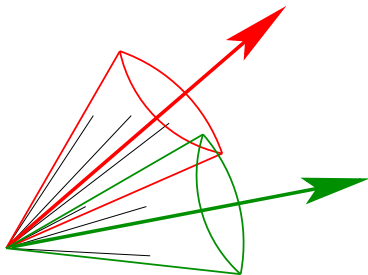
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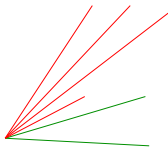
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**Qu: How do you find the stable cones?**

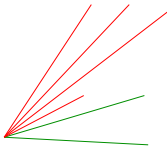
Until recently used iterative methods:

- ▶ use each particle as a starting direction for cone; use sum of contents as new starting direction; repeat.

**Iterative Cone with Split Merge (IC-SM)**

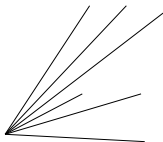
e.g. Tevatron cones (JetClu, midpoint)

ATLAS cone



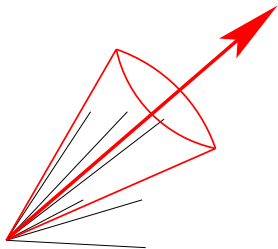
## Other cones avoid split-merge:

- ▶ Find one stable cone
  - ▶ Call it a jet; remove its particles from the event; repeat
- E.g. by iterating from hardest seed particle



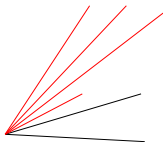
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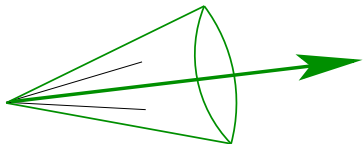
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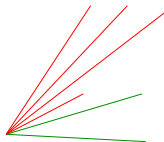
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- ▶ This is not the same algorithm
- ▶ Many physics aspects differ

## **Iterative Cone with Progressive Removal (IC-PR)**

e.g. CMS cone, Pythia Cone, [GetJet], ...



- ▶ What are the starting points for iteration?
- ▶ Start with hardest particle as seed (IC-PR): collinear unsafe
- ▶ Use all particles (IC-SM): extra soft one → new solution

Iterative cone finding **plagued by IR and collinear unsafety problems**

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Iterative cone finding **plagued by IR and collinear unsafety problems**

Among consequences of IR unsafety:

	<i>Last meaningful order</i>		
	ATLAS cone [IC-SM]	MidPoint [IC <sub>mp</sub> -SM]	CMS it. cone [IC-PR]
Inclusive jets	LO	NLO	NLO
W/Z + 1 jet	LO	NLO	NLO
3 jets	<b>none</b>	LO	LO
W/Z + 2 jets	<b>none</b>	LO	LO
$m_{\text{jet}}$ in $2j + X$	<b>none</b>	<b>none</b>	<b>none</b>

NB: \$30 – 50M investment in NLO

Algorithm	Type	IRC status	Notes
CDF JetClu	$IC_r$ -SM	$IR_{2+1}$	
CDF MidPoint cone	$IC_{mp}$ -SM	$IR_{3+1}$	
CDF searchcone	$IC_{se,mp}$ -SM	$IR_{2+1}$	
D0 Run II cone	$IC_{mp}$ -SM	$IR_{3+1}$	no seed threshold, but c
ATLAS Cone	IC-SM	$IR_{2+1}$	
PxCone	$IC_{mp}$ -SD	$IR_{3+1}$	no seed threshold, but c
CMS Iterative Cone	IC-PR	$Coll_{3+1}$	
PyCell/CellJet (from Pythia)	FC-PR	$Coll_{3+1}$	
GetJet (from ISAJET)	FC-PR	$Coll_{3+1}$	

**Too many cones, too many problems, needs consolidation.**

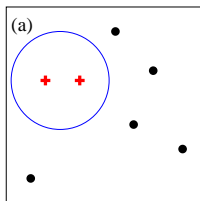
IC-SM → SIScone (finds all stable cones)  
 IC-PR → anti- $k_t$  [NEW!! More in a few minutes]

## Solve IR issue: find **all** stable cones

Cones are just *circles* in the  $y - \phi$  plane. To find all stable cones:

1. Find all distinct ways of enclosing a subset of particles in a  $y - \phi$  circle
2. Check, for each enclosure, if it corresponds to a stable cone

Finding all distinct circular enclosures of a set of points is **geometry**:



*Any enclosure can be moved until a pair of points lies on its edge.*

Result: Seedless Infrared Safe Cone algorithm (SISCone)

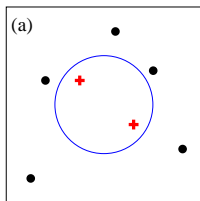
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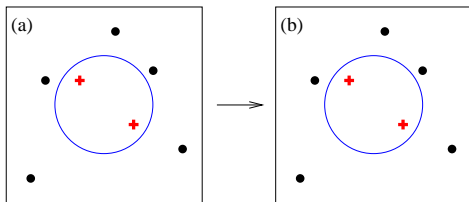
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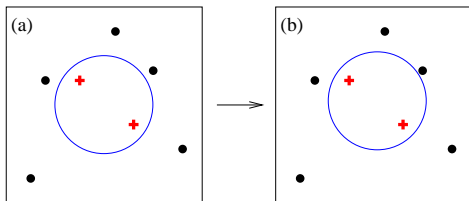
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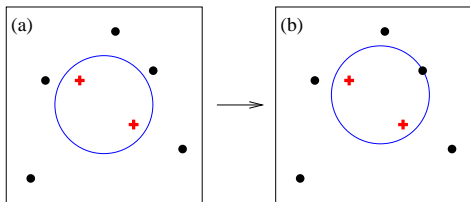
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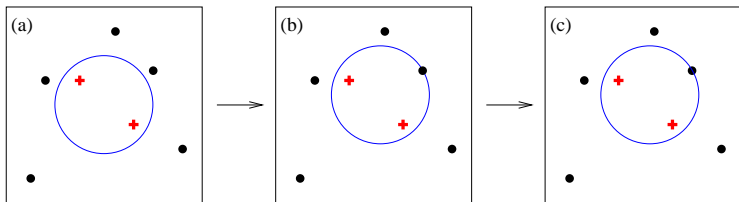
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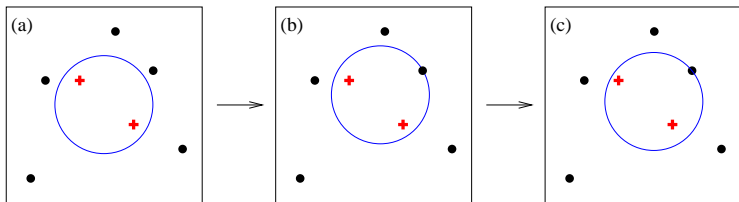
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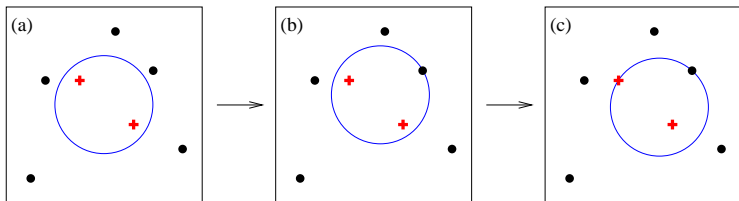
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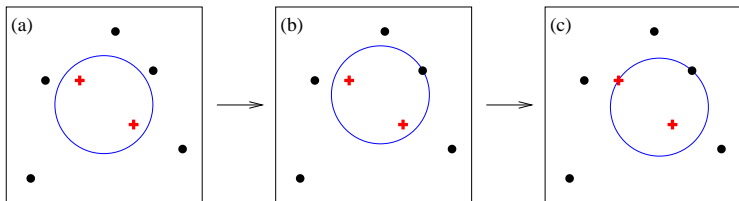
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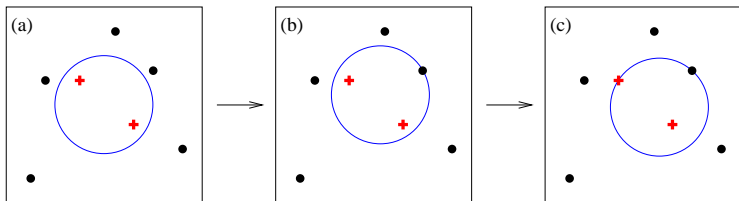
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- 1: Put the set of current particles equal to the set of all particles in the event.
- 2: **repeat**
- 3: Find *all* stable cones of radius  $R$  for the current set of particles, e.g. using algorithm 2.
- 4: For each stable cone, create a protojet from the current particles contained in the cone, and add it to the list of protojets.
- 5: Remove all particles that are in stable cones from the list of current particles.
- 6: **until** No new stable cones are found, or one has gone around the loop  $N_{\text{pass}}$  times.
- 7: Run a Tevatron Run-II type split–merge procedure, algorithm 3, on the full list of protojets, with overlap parameter  $f$  and transverse momentum threshold  $p_{t,\text{min}}$ .

- 1: For any group of collinear particles, merge them into a single particle.
- 2: **for** particle  $i = 1 \dots N$  **do**
- 3: Find all particles  $j$  within a distance  $2R$  of  $i$ . If there are no such particles,  $i$  forms a stable cone of its own.
- 4: Otherwise for each  $j$  identify the two circles for which  $i$  and  $j$  lie on the circumference. For each circle, compute the angle of its centre  $C$  relative to  $i$ ,  $\zeta = \arctan \frac{\Delta\phi_{iC}}{\Delta y_{iC}}$ .
- 5: Sort the circles into increasing angle  $\zeta$ .
- 6: Take the first circle in this order, and call it the current circle. Calculate the total momentum and checkxor for the cones that it defines. Consider all 4 permutations of edge points being included or excluded. Call these the "current cones".
- 7: **repeat**
- 8:     **for** each of the 4 current cones **do**
- 9:         If this cone has not yet been found, add it to the list of distinct cones.
- 10:         If this cone has not yet been labelled as unstable, establish if the in/out status of the edge particles (with respect to the cone momentum axis) is the same as when defining the cone; if it is not, label the cone as unstable.
- 11:     **end for**
- 12:     Move to the next circle in order. It differs from the previous one either by a particle entering the circle, or one leaving the circle. Calculate the momentum for the new circle and corresponding new current cones by adding (or removing) the momentum of the particle that has entered (left); the checkxor can be updated by XORing with the label of that particle.
- 13:     **until** all circles considered.
- 14: **end for**
- 15: **for** each of the cones not labelled as unstable **do**
- 16:     Explicitly check its stability, and if it is stable, add it to the list of stable cones (protojets).
- 17: **end for**



## 1: repeat

Remove all protojets with  $p_t < p_{t,\min}$ .

Identify the protojet ( $i$ ) with the highest  $\tilde{p}_t$  ( $\tilde{p}_{t,\text{jet}} = \sum_{i \in \text{jet}} |p_{t,i}|$ ).

Among the remaining protojets identify the one ( $j$ ) with highest  $\tilde{p}_t$  that shares particles (overlaps) with  $i$ .

## 5: if there is such an overlapping jet then

6: Determine the total  $\tilde{p}_{t,\text{shared}} = \sum_{k \in i \& j} |p_{t,k}|$  of the particles shared between  $i$  and  $j$ .

## 7: if $\tilde{p}_{t,\text{shared}} < f \tilde{p}_{t,j}$ then

Each particle that is shared between the two protojets is assigned to the one to whose axis it is closest. The protojet momenta are then recalculated.

## 9: else

Merge the two protojets into a single new protojet (added to the list of protojets, while the two original ones are removed).

## 11: end if

12: If steps 7–11 produced a protojet that coincides with an existing one, maintain the new protojet as distinct from the existing copy(ies).

## 13: else

Add  $i$  to the list of final jets, and remove it from the list of protojets.

## 15: end if

16: until no protojets are left.

- ▶ Generate event with  $2 < N < 10$  hard particles, find jets
- ▶ Add  $1 < N_{soft} < 5$  soft particles, find jets again [repeatedly]
- ▶ If the jets are different, algorithm is IR unsafe.

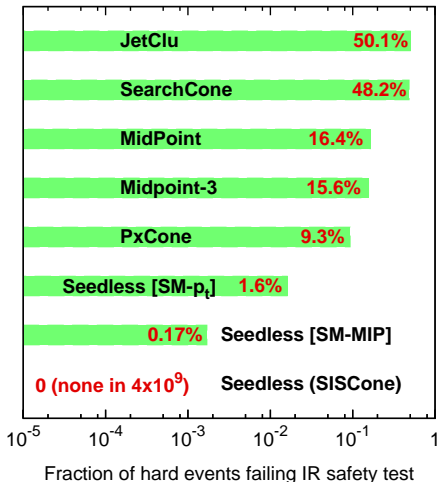
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2 hard + 1 soft	~ 50%
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Majority of QCD branching is soft & collinear, with following divergences:

$$|dk_j| |M_{g \rightarrow g_i g_j}^2(k_j)| \simeq \frac{2\alpha_s C_A}{\pi} \frac{dE_j}{\min(E_i, E_j)} \frac{d\theta_{ij}}{\theta_{ij}}, \quad (E_j \ll E_i, \theta_{ij} \ll 1).$$

**Invert branching process:** take pair with strongest divergence between them — they're the most *likely* to belong together.

→  $k_t$ /Durham algorithm ( $e^+e^-$ )

1. Calculate (or update) distances between all particles  $i$  and  $j$ :

$$y_{ij} = \frac{2 \min(E_i^2, E_j^2)(1 - \cos \theta_{ij})}{Q^2}$$

2. Find smallest of  $y_{ij}$

NB: relative  $k_t$  between particles

- ▶ If  $> y_{cut}$ , stop clustering
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Catani, Dokshitzer, Olsson, Turnock & Webber '91

## Exclusive long. inv. $k_t$ algorithm:

- ▶ Drop normalisation to  $Q^2$  (not fixed in pp):  $y_{ij} \rightarrow d_{ij}$ ,  $y_{cut} \rightarrow d_{cut}$
- ▶ Make it longitudinally boost invariant

$$d_{ij} = \min(p_{ti}^2, p_{tj}^2) \Delta R_{ij}^2, \quad \Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$$

- ▶ Introduce clustering with “beam jets”, distance:  $d_{iB} = p_{ti}^2$   
Catani, Dokshitzer, Seymour & Webber '93

## Inclusive long. inv. $k_t$ : introduce $R$ , drop $d_{cut}$

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3. repeat from step 1 until no particles left.

S.D. Ellis & Soper, '93; most “cone-like”  
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## Fast Hierarchical Clustering and Other Applications of Dynamic Closest Pairs

David Eppstein  
UC Irvine

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We develop data structures for dynamic closest pair problems with arbitrary distance functions, that do not necessarily come from any geometric structure on the objects. Based on a technique previously used by the author for Euclidean closest pairs, we show how to insert and delete objects from an  $n$ -object set, maintaining the closest pair, in  $O(n \log^2 n)$  time per update and  $O(n)$  space. With quadratic space, we can instead use a quadtree-like structure to achieve an optimal time bound,  $O(n)$  per update. We apply these data structures to hierarchical clustering, greedy matching, and TSP heuristics, and discuss other potential applications in machine learning, Gröbner bases, and local improvement algorithms for partition and placement problems. Experiments show our new methods to be faster in practice than previously used heuristics.

Categories and Subject Descriptors: F.2.2 [Analysis of Algorithms]: Nonnumeric Algorithms

General Terms: Closest Pair, Agglomerative Clustering

Additional Key Words and Phrases: TSP, matching, conga line data structure, quadtree, nearest neighbor heuristic

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### 1. INTRODUCTION

Hierarchical clustering has long been a mainstay of statistical analysis, and clustering based methods have attracted attention in other fields: computational biology (reconstruction of evolutionary trees; tree-based multiple sequence alignment), scientific simulation ( $n$ -body problems), theoretical computer science (network design and nearest neighbor searching) and of course the web (hierarchical indices such as Yahoo). Many clustering methods have been devised and used in these applications, but less effort has gone into algorithmic speedups of these methods.

In this paper we identify and demonstrate speedups for a key subroutine used in several clustering algorithms, that of maintaining closest pairs in a dynamic set of objects. We also describe several other applications or potential applications of the

Idea behind  $k_t$  alg. is to be found over and over in many areas of (computer) science.

# Sequential recombination of cacti

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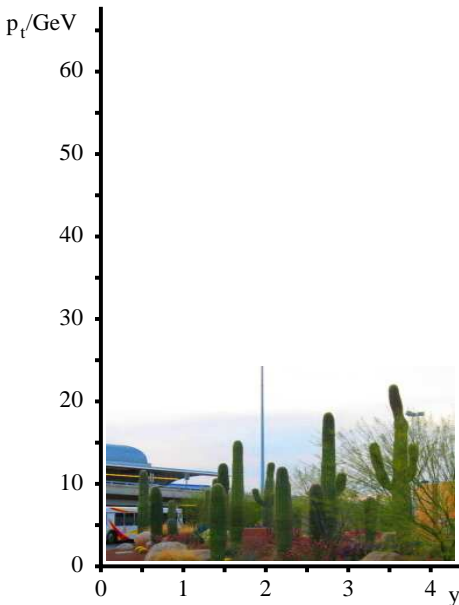
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$\phi$  assumed 0 for all towers



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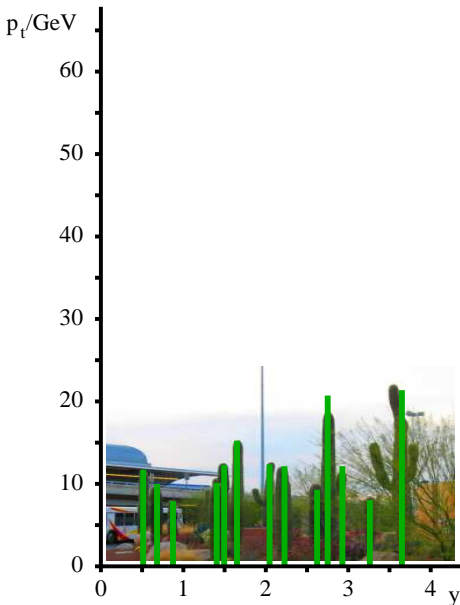
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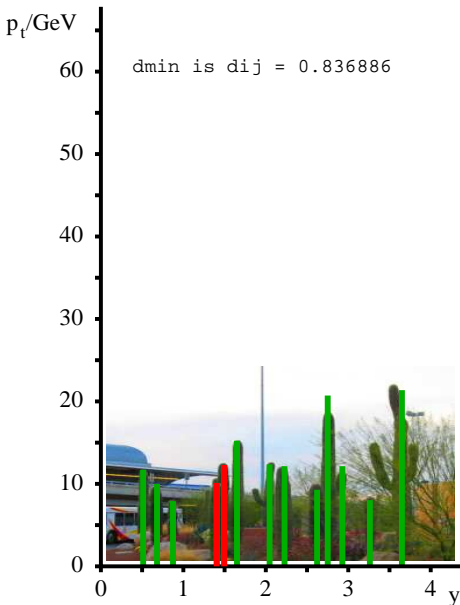


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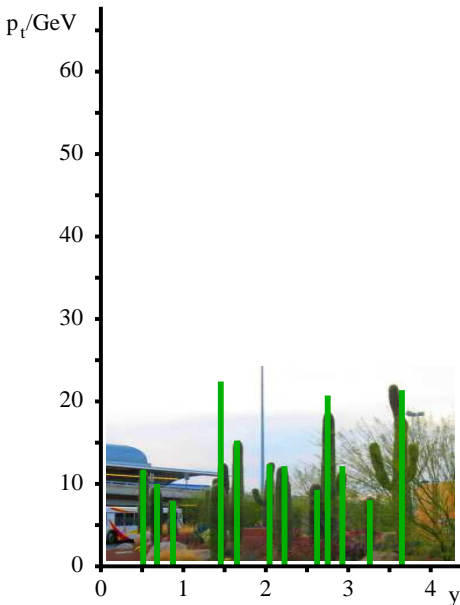
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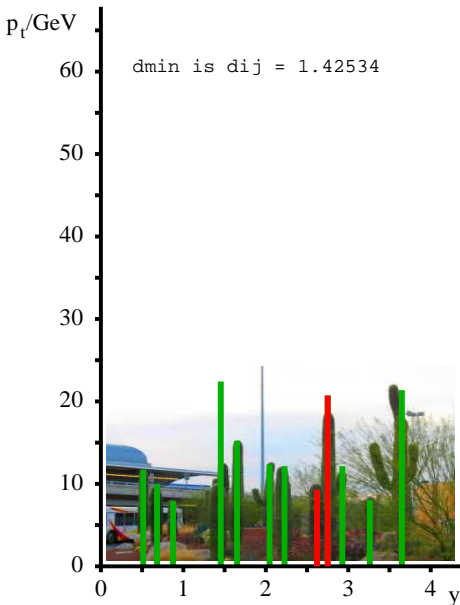


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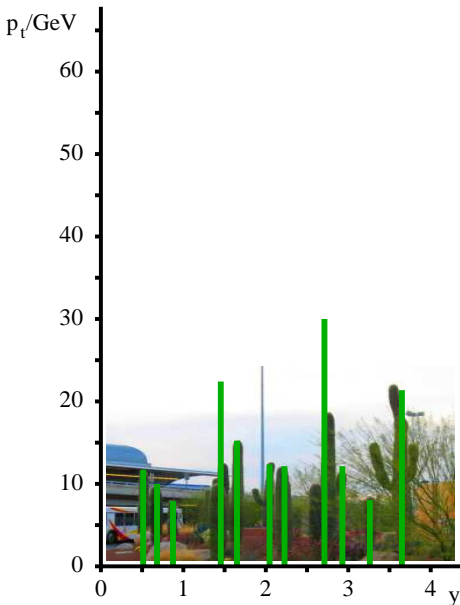
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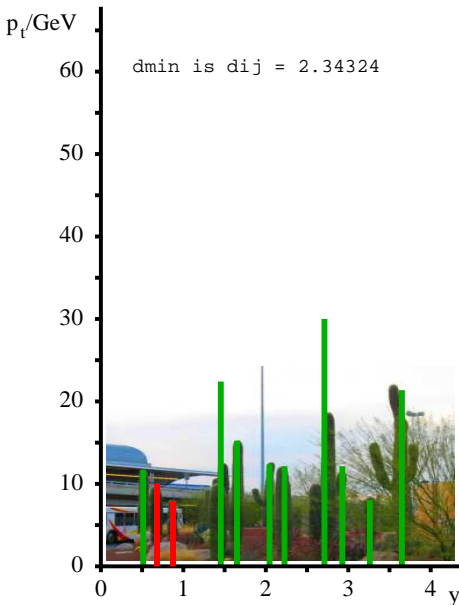
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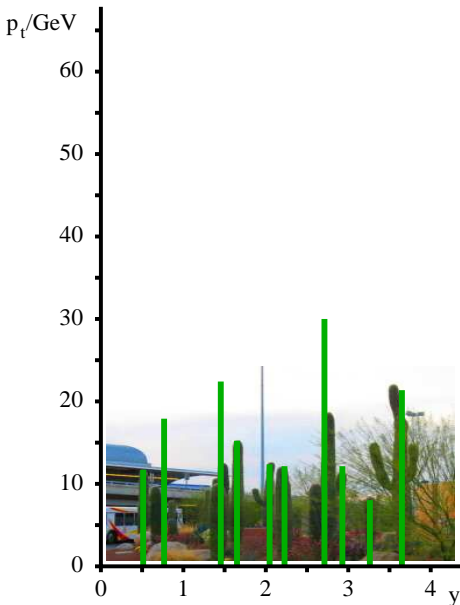
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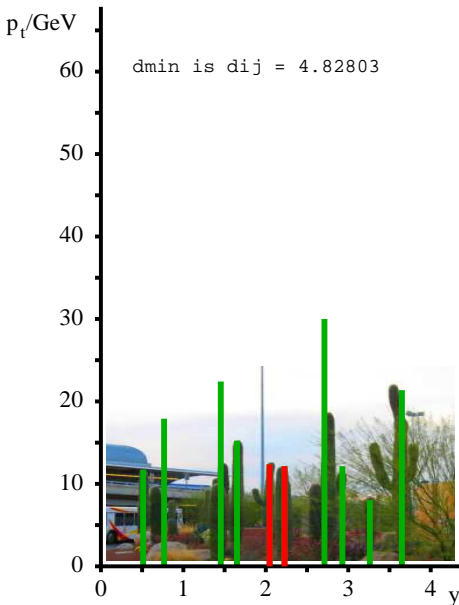


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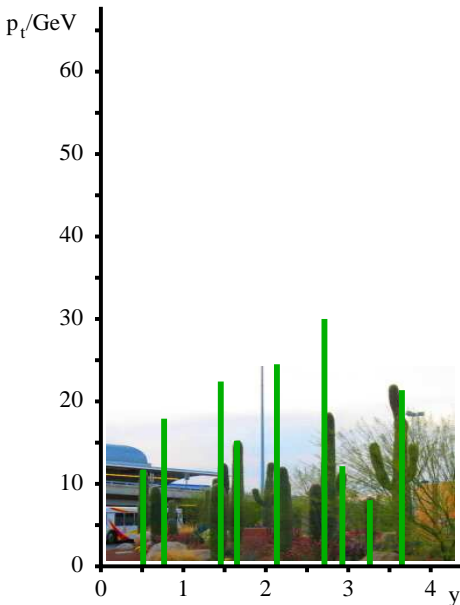
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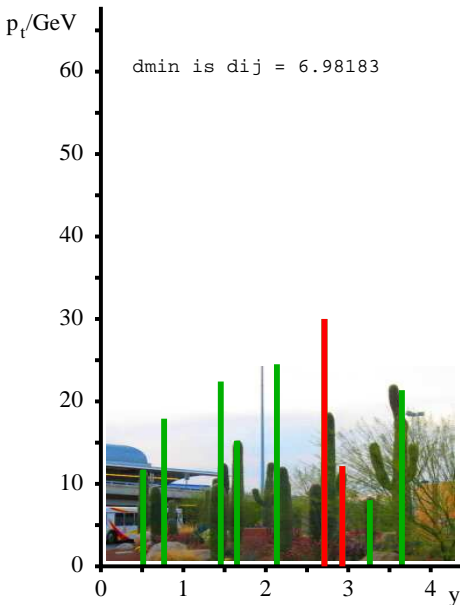


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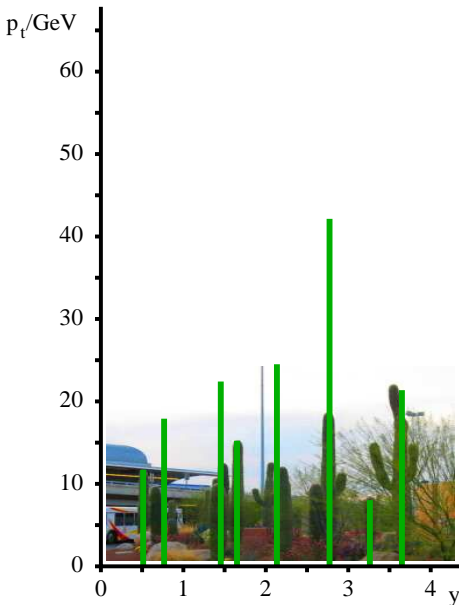
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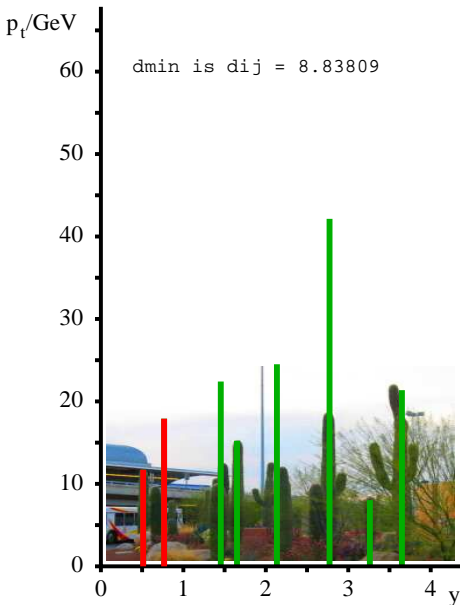


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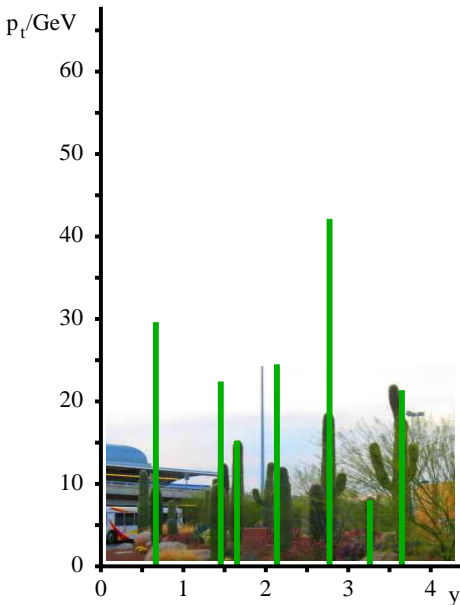
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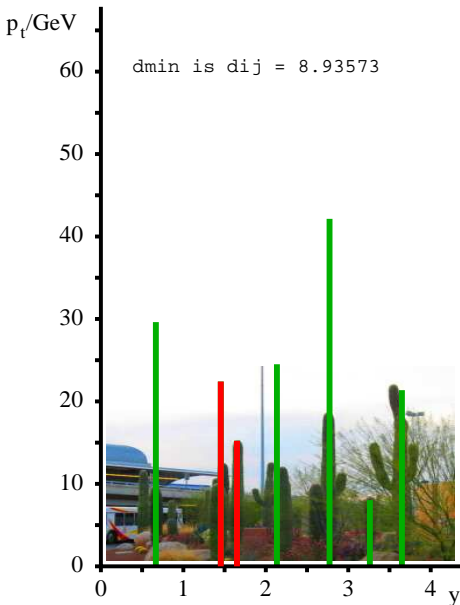
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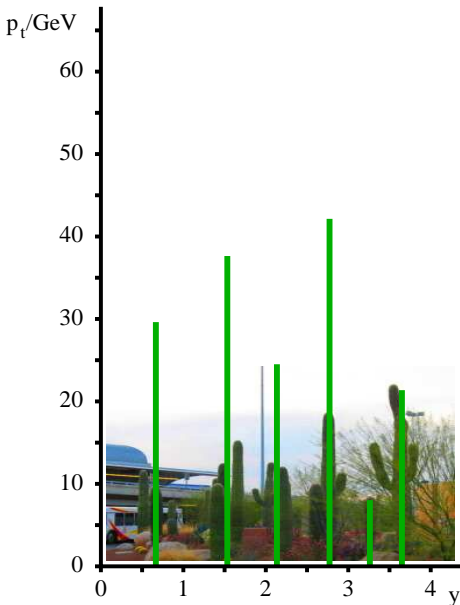
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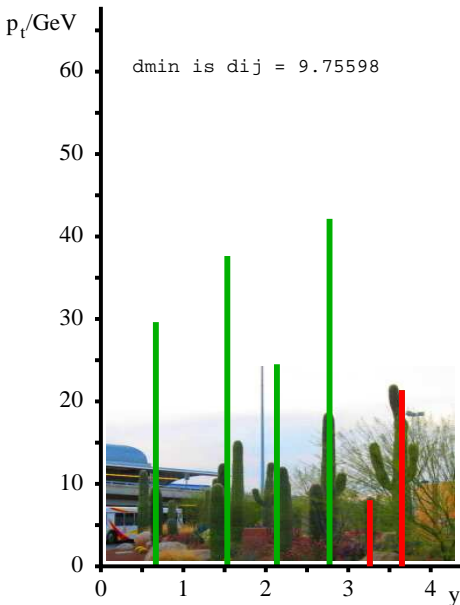
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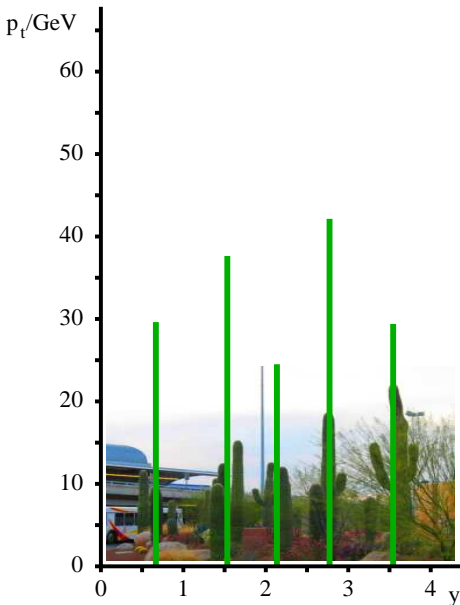
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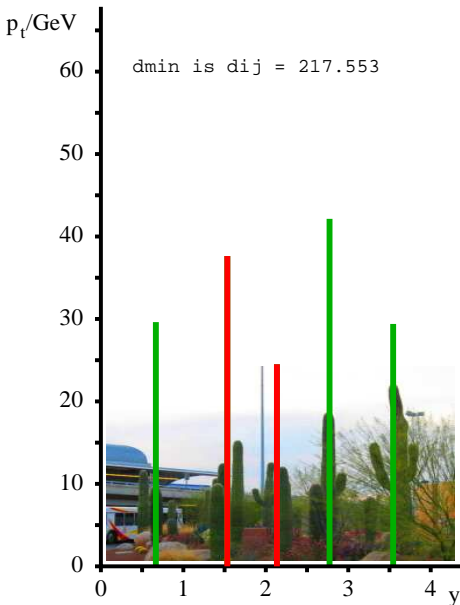


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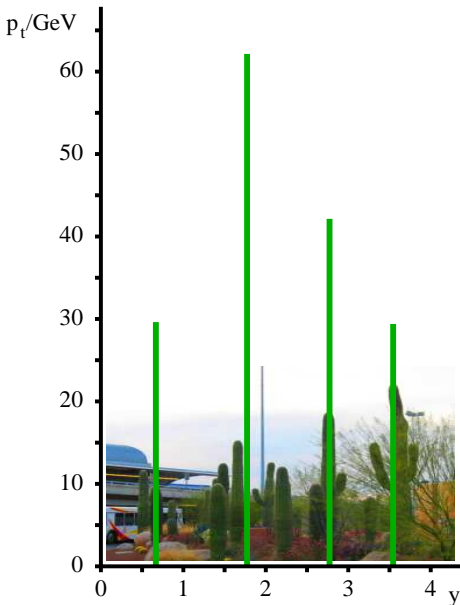
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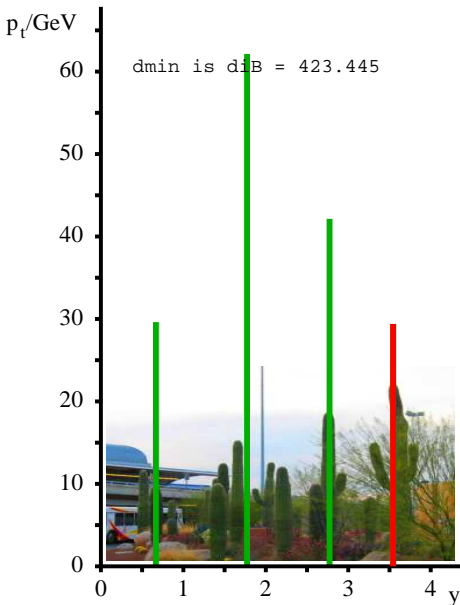
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# Sequential recombination of cacti



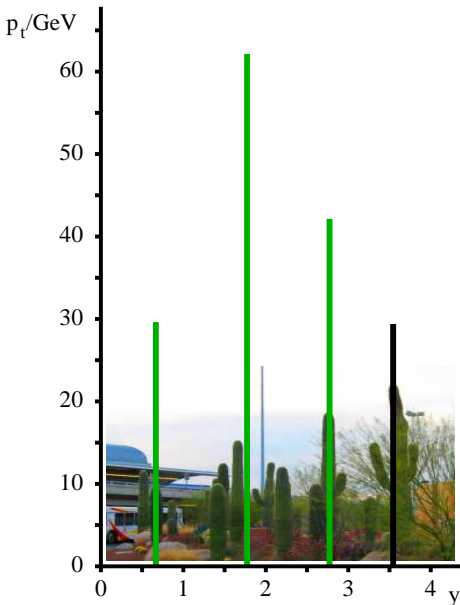
**$k_t$  alg.:** Find smallest of

$$d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2, \quad d_{iB} = k_{ti}^2$$

If  $d_{ij}$  recombine; if  $d_{iB}$ ,  $i$  is a jet  
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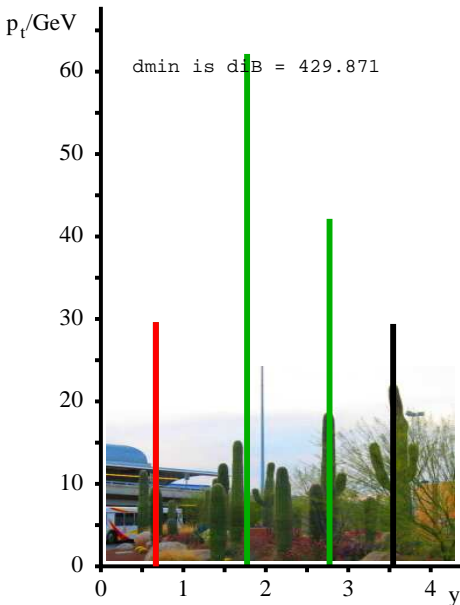
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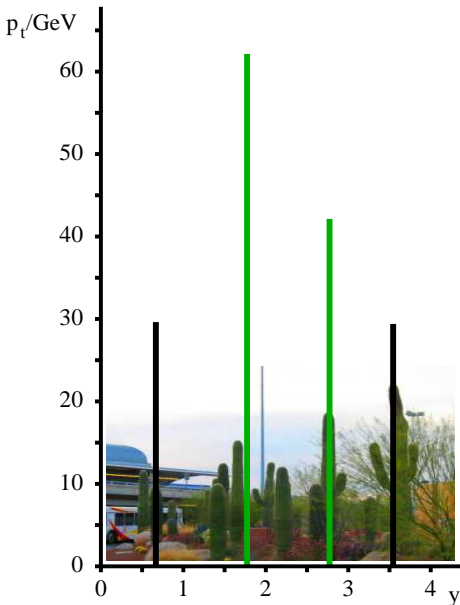
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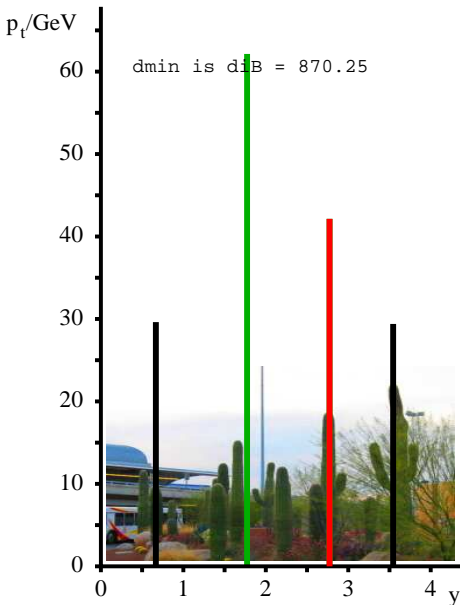
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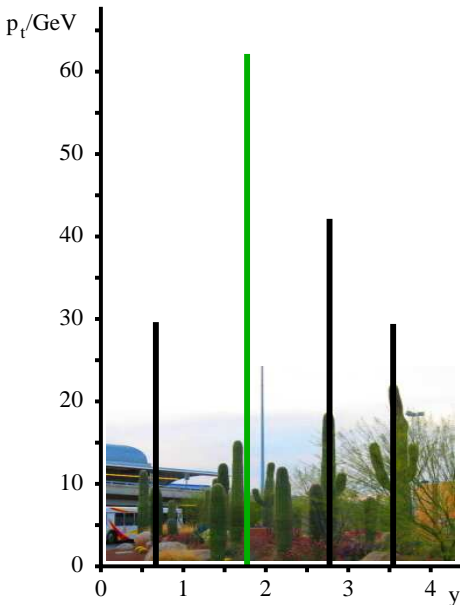
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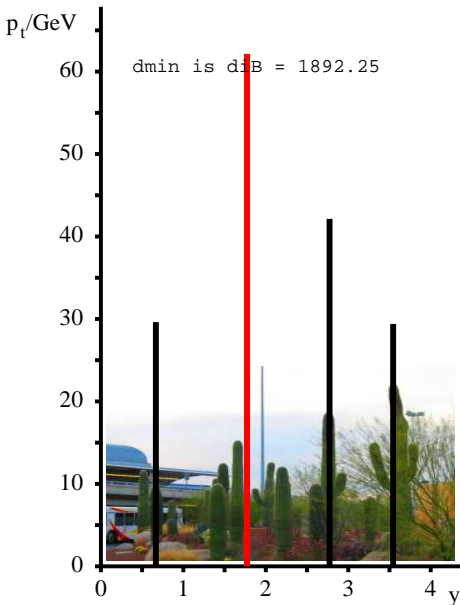
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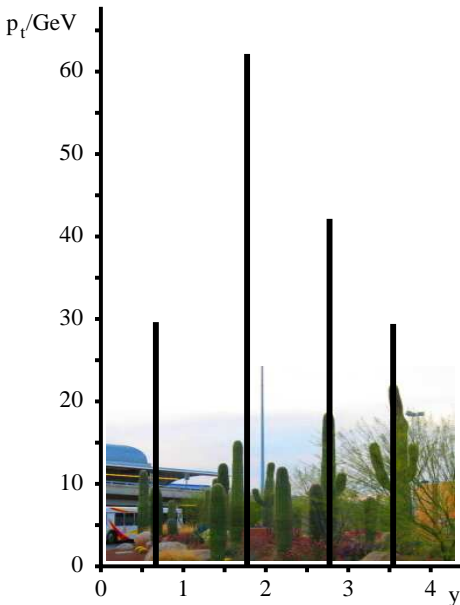
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 Example clustering with  $k_t$  algorithm,  $R = 0.7$

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Generalise inclusive-type sequential recombination with

$$d_{ij} = \min(k_{ti}^{2p}, k_{tj}^{2p}) \Delta R_{ij}^2 / R^2 \quad d_{iB} = k_{ti}^{2p}$$

	Alg. name	Comment	time
$p = 1$	$k_t$ CDOSTW '91-93; ES '93	Hierarchical in rel. $k_t$	$N \ln N$ exp.
$p = 0$	Cambridge/Aachen Dok, Leder, Moretti, Webber '97 Wengler, Wobisch '98	Hierarchical in angle Scan multiple $R$ at once $\leftrightarrow$ QCD angular orderin	$N \ln N$
$p = -1$	anti- $k_t$ Cacciari, GPS, Soyez '08 $\sim$ reverse- $k_t$ Delsart, Loch et al.	Hierarchy meaningless. Behaves like IC-PR	$N^{3/2}$
SC-SM	SISCone GPS Soyez '07 + Tevatron run II '00	Replacement for IC-SM notably "MidPoint" cones	$N^2 \ln N$ exp.

One could invent/try others (e.g. OJF, etc.). Our [Paris+BNL] philosophy: 4 algs is enough of a basis to develop first physics understanding.

We already have far more than can be shown here

non-COMMERCIAL BREAK



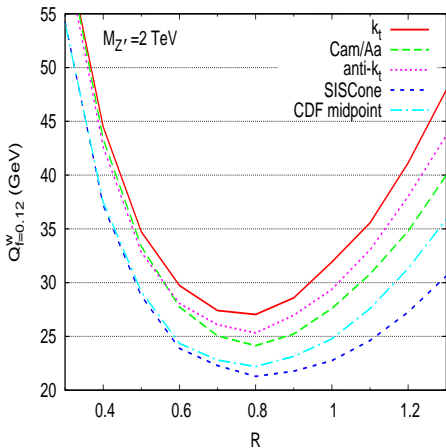
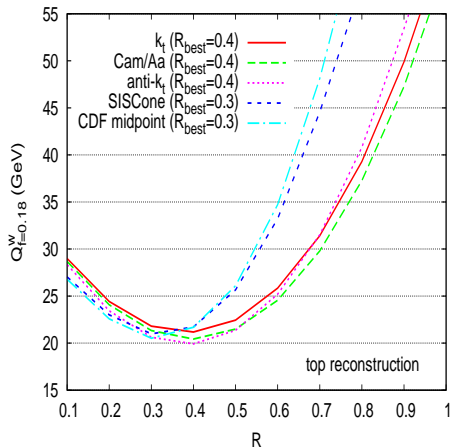
One place to stop for your jet-finding needs:

## FastJet

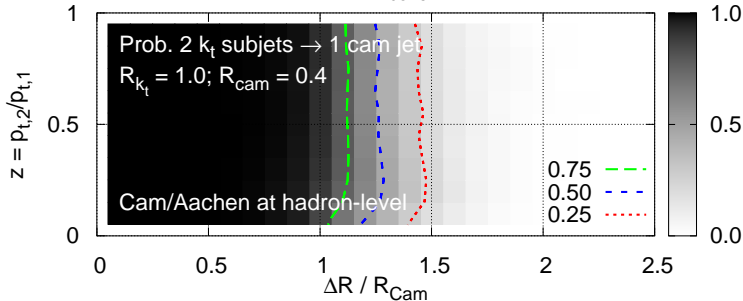
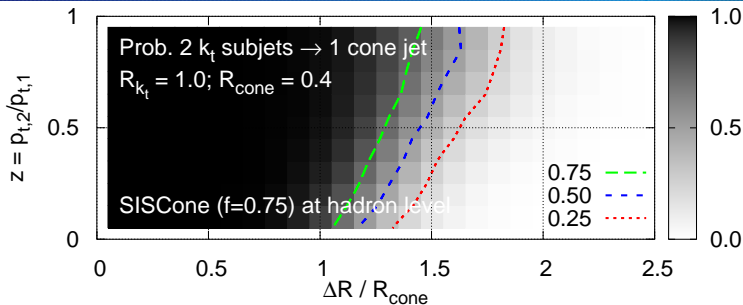
<http://www.lpthe.jussieu.fr/~salam/fastjet>  
Cacciari, GPS & Soyez '05–08

- ▶ Fast, native, computational-geometry methods for  $k_t$ , Cam/Aachen, anti- $k_t$  Cacciari & GPS '05-06
- ▶ Plugins for SISCone (plus some other, deprecated, legacy cones)
- ▶ Documented user interface for adding extra algorithms of your own
- ▶ Tools for jet areas, pileup characterisation & subtraction

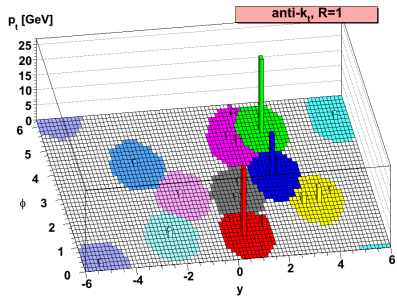
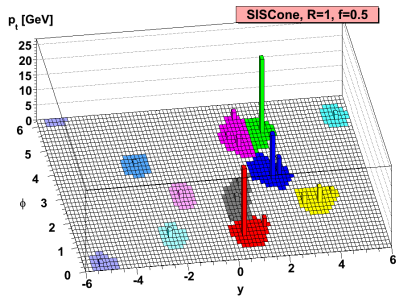
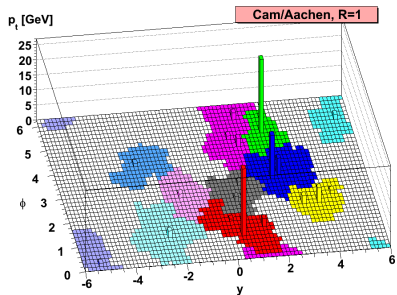
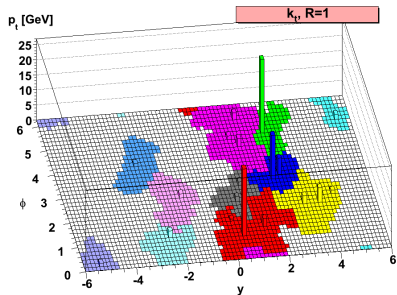
E.g. width of narrowest window (around mass-peak) containing  $\sim 25\%$  of events that pass basic selection cuts, in hadronic  $t\bar{t}$  and  $Z'$  events.



# Reach of jet algorithms



# Jet contours – visualised



## LEP

- ▶  $M_{BSM} \sim 1 \text{ TeV}$
- ▶  $M_{EW} \sim 100 \text{ GeV}$   $\sim \alpha_s \times M_{BSM}$
- ▶  $p_{t,\text{pileup}} \sim 25 - 100 \text{ GeV/unit rap.}$   $\sim M_{EW}$
- ▶  $p_{t,\text{UE}} \sim 2.5 - 5 \text{ GeV/unit rap.}$   $\sim \alpha_s \times M_{EW}$
- ▶  $p_{t,\text{hadr.}} \sim 0.5 \text{ GeV/unit rap.}$

Multitude of scales forces us to go beyond  
"1 parton = 1 jet" equivalence

## Tevatron

- ▶  $M_{BSM} \sim 1 \text{ TeV}$
- ▶  $M_{EW} \sim 100 \text{ GeV}$   $\sim \alpha_s \times M_{BSM}$
- ▶  $p_{t,\text{pileup}} \sim 25 - 100 \text{ GeV/unit rap.}$   $\sim M_{EW}$
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## LHC

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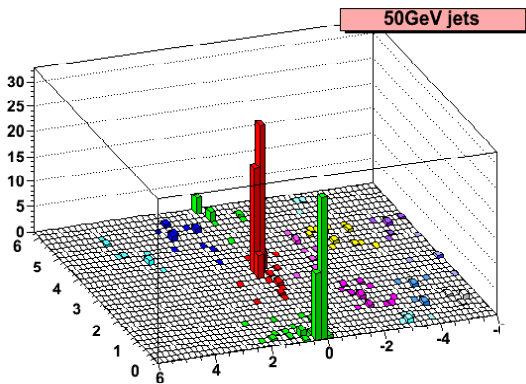
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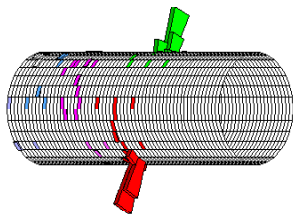


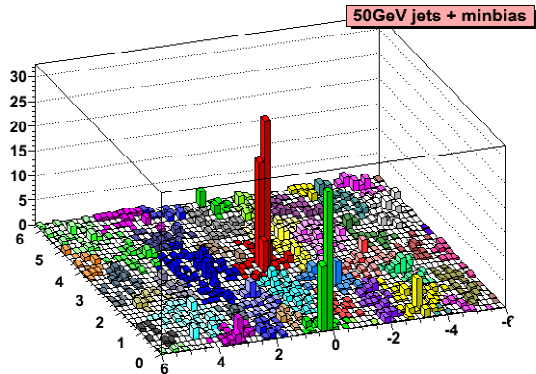
'Standard hard' event  
Two well isolated jets

50GeV jets

~ 200 particles

Clustering takes  $\lesssim 1ms$

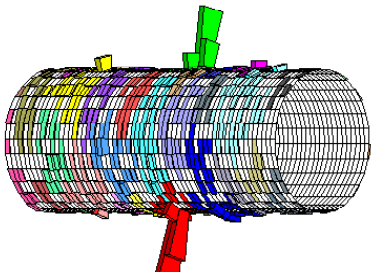


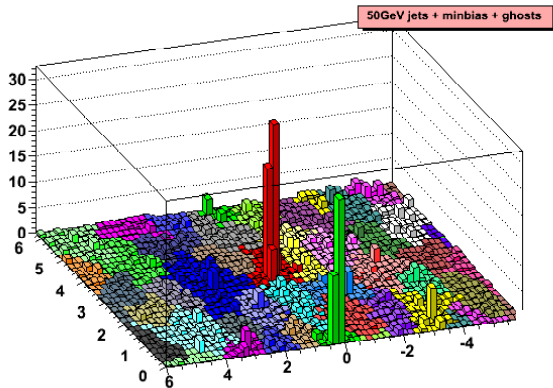


Add 10 min-bias events  
(moderately high lumi)

~ 2000 particles

Clustering takes ~ 10ms



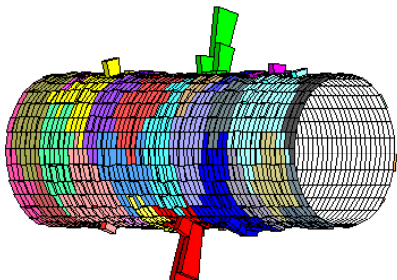


~ 10000 particles

Clustering takes ~ 0.2s

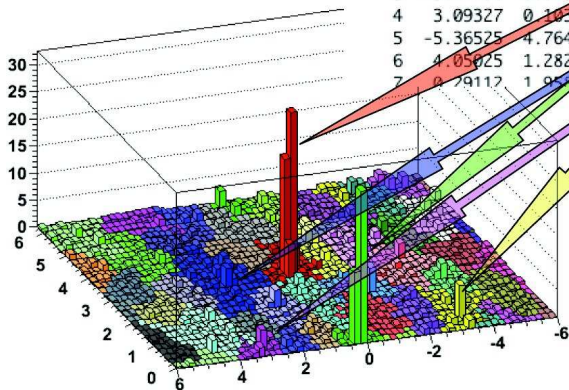
Add dense coverage of infinitely soft "ghosts"

See how many end up in jet to measure jet area



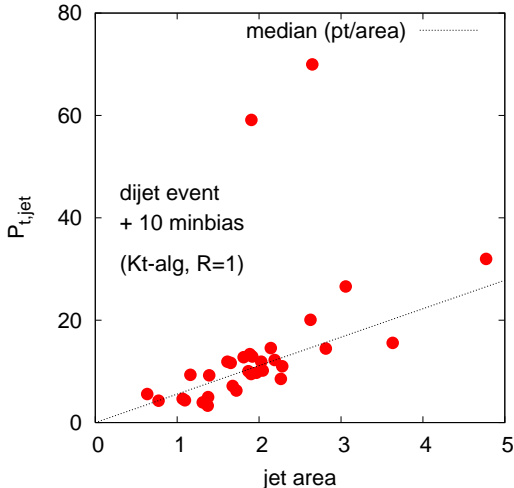
iev 0 (irepeat 24): number of particles = 1428  
 strategy used = NlnN  
 number of particles = 9051  
 Total area: 76.0265  
 Expected area: 76.0265

ijet	eta	phi	Pt	area +- err
0	0.15050	3.24498	69.970	2.625 +- 0.020
1	0.18579	0.13150	59.133	1.896 +- 0.020
2	2.33840	3.23960	31.976	4.749 +- 0.028
3	-3.41796	0.52394	26.595	3.084 +- 0.021
4	3.09327	0.10350	20.072	2.688 +- 0.023
5	-5.36525	4.76491	19.593	2.780 +- 0.012
6	4.05025	1.28279	15.861	3.592 +- 0.028
7	0.29112	1.95745	12.566	2.114 +- 0.018



Approximate linear relation  
 between Pt and area for  
 minimum bias jets.

Can be used on an event-by-  
 event basis to correct the hard  
 jets



Jet areas in  $k_t$  algorithm are quite varied

Because  $k_t$ -alg adapts to the jet structure

- ▶ Hard jets' contamination from min-bias  $\sim$  area  
Area varies even for SISConc
- ▶ Soft jets'  $p_t/area$  tells you about min-bias normalisation and fluctuations

Median  $p_t/area$  across the set of jets in an event is a good estimator of pileup+UE in *that event*

## Basic Procedure:

- ▶ Use  $p_t/A$  from majority of jets (pileup jets) to get level,  $\rho$ , of pileup and UE in event
- ▶ Subtract pileup from hard jets:

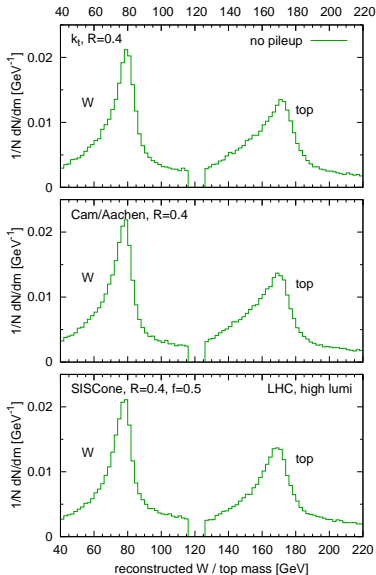
$$p_t \rightarrow p_{t,sub} = p_t - A\rho$$

Cacciari & GPS '07

## Illustration:

- ▶ semi-leptonic  $t\bar{t}$  production at LHC
- ▶ high-lumi pileup ( $\sim 20$  ev/bunch-X)

*Same simple procedure works for a range of algorithms*



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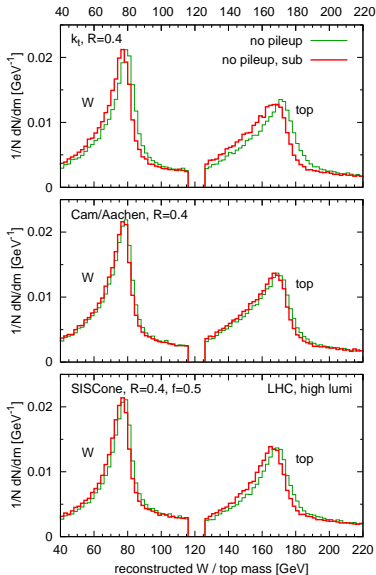
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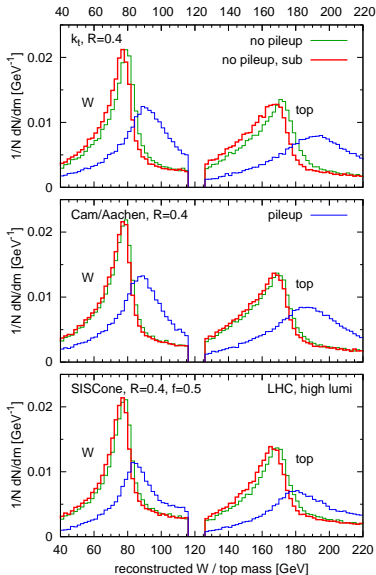
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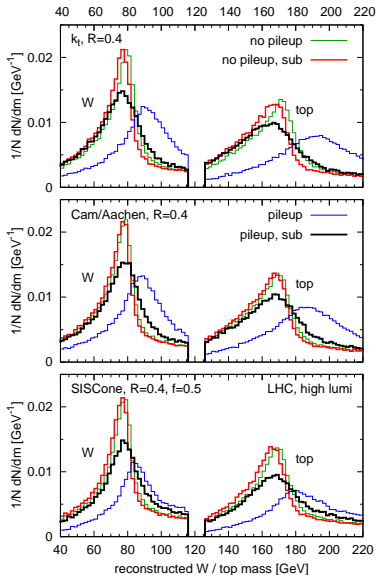
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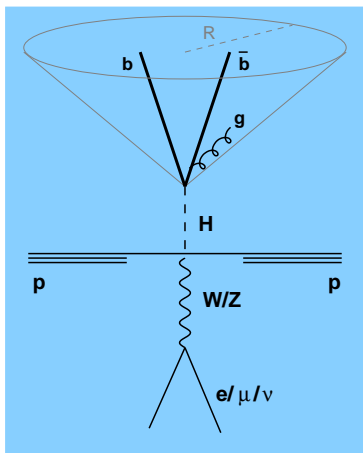
*Same simple procedure works for a range of algorithms*



# Searching for high- $p_t$ HW/WZ?

High- $p_t$  light Higgs decays to  $b\bar{b}$  inside a single jet. Can this be seen?

Butterworth, Davison, Rubin & GPS '08



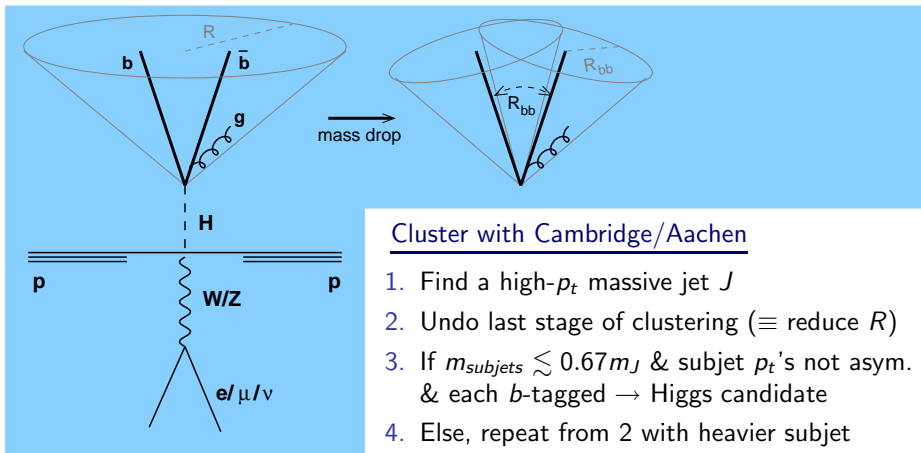
## Cluster with Cambridge/Aachen

1. Find a high- $p_t$  massive jet  $J$
2. Undo last stage of clustering ( $\equiv$  reduce  $R$ )
3. If  $m_{\text{subjects}} \lesssim 0.67m_J$  & subjet  $p_t$ 's not asym. & each  $b$ -tagged  $\rightarrow$  Higgs candidate
4. Else, repeat from 2 with heavier subjet

Then on the Higgs-candidate: *filter* away UE/pileup by reducing  $R \rightarrow R_{\text{filt}}$ , take *three hardest subjects* (keep LO gluon rad<sup>n</sup>) + require  $b$ -tags on two hardest.

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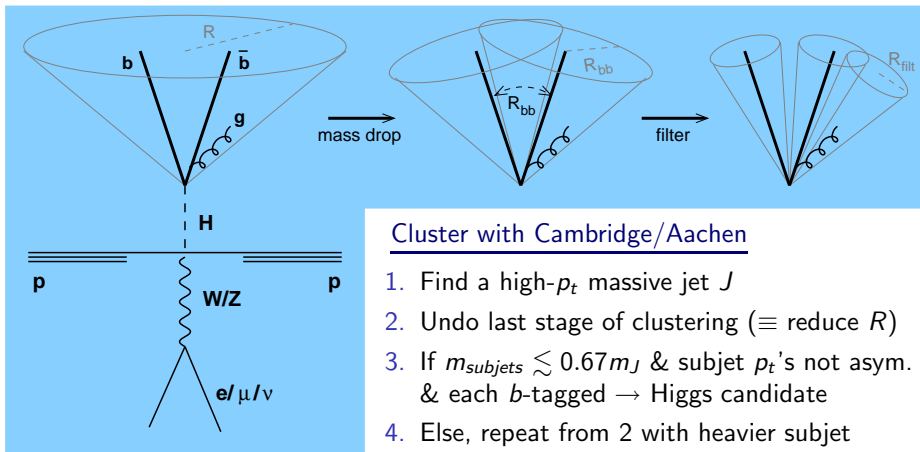
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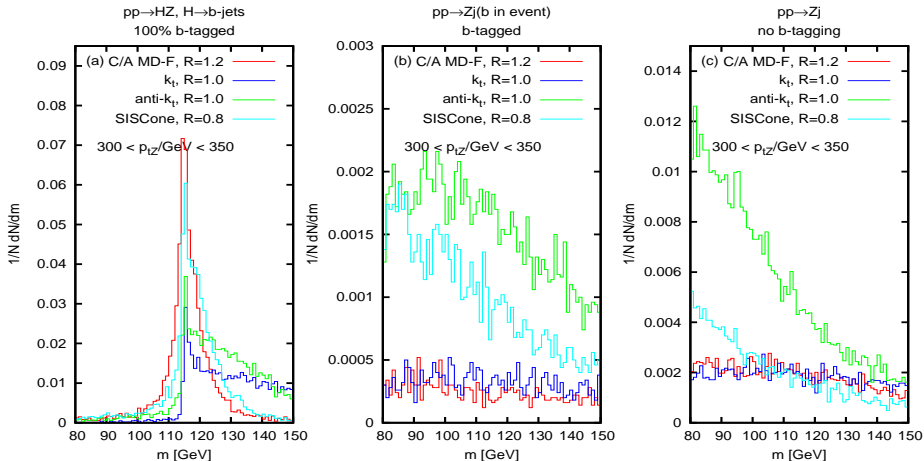
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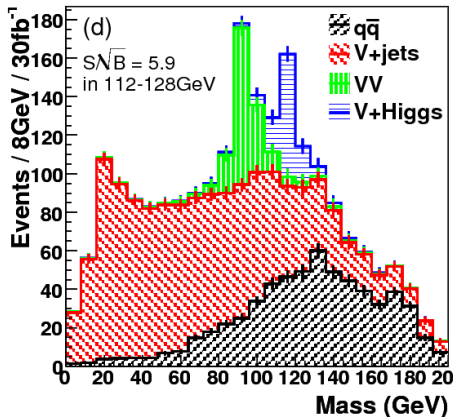


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Check mass spectra in HZ channel,  $H \rightarrow b\bar{b}$ ,  $Z \rightarrow \ell^+\ell^-$



Cambridge/Aachen (C/A) with mass-drop and filtering (MD/F) works best

Combined HZ + HW,  $p_t > 200$  GeV


- ▶ Take  $Z \rightarrow \ell^+ \ell^-$ ,  $Z \rightarrow \nu \bar{\nu}$ ,  
 $W \rightarrow \ell \nu$   $\ell = e, \mu$
- ▶  $p_{tV}, p_{tH} > 200$  GeV
- ▶  $|\eta_V|, |\eta_H| < 2.5$
- ▶ Assume real/fake  $b$ -tag rates of 0.7/0.01.
- ▶ Some extra cuts in  $HW$  channels to reject  $t\bar{t}$ .
- ▶ Assume  $m_H = 115$  GeV.

At  $5.9\sigma$  for  $30 \text{ fb}^{-1}$  this looks like a competitive channel for light Higgs discovery. **Deserves serious exp. study!**

- ▶ Know the algorithms and make sure they're fully described somewhere
- ▶ Be sure they're infrared & collinear safe (there's plenty of choice)
- ▶ LHC's a very different environment from what we've had before
- ▶ One should explore a range of jet definitions for any given analysis
- ▶ Both to optimize the analysis
- ▶ And to verify robustness of conclusions wrt jet-def
- ▶ Jets are *not just partons*
- ▶ Tell you lots about UE, pileup → model it better, subtract it better
- ▶ Have substructure, which can help reveal underlying physics
- ▶ As data arrives at LHC, new ideas in jet-definition will arise, geared to actual LHC conditions. Keep your eyes open for best tools.

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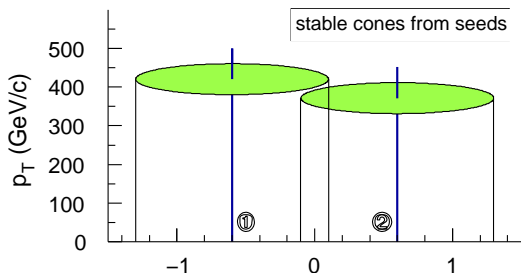


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# EXTRA MATERIAL

## Use of seeds is *dangerous*



Extra soft particle adds new seed  $\rightarrow$  changes final jet configuration.

This is **IR unsafe**.

Divergences of real and virtual contributions do not cancel at  $\mathcal{O}(\alpha_s^4)$

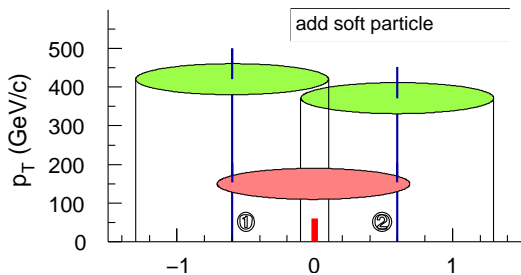
Kilgore & Giele '97

Solution: add extra seeds at midpoints of all pairs, triplets, ... of stable cones.

Seymour '97 (?)

**NB:** only in past 3-4 years has this fix appeared in CDF and D0 analyses...

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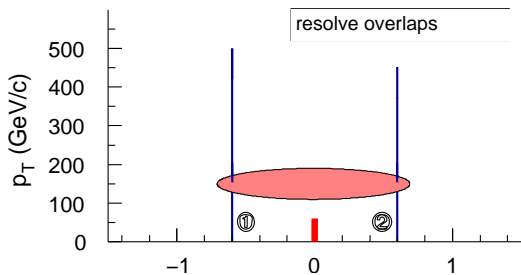
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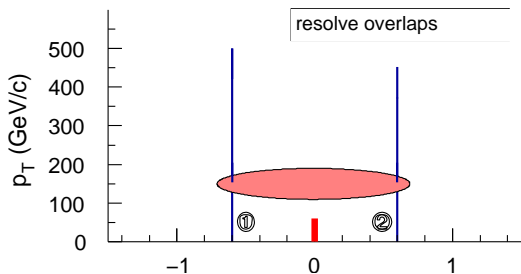
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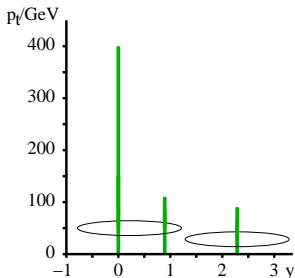
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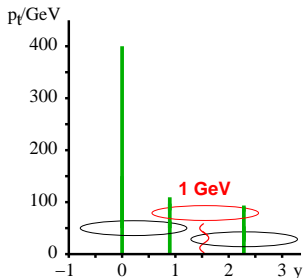


Stable cones  
with midpoint:

$\{1,2\}$  &  $\{3\}$

Jets with  
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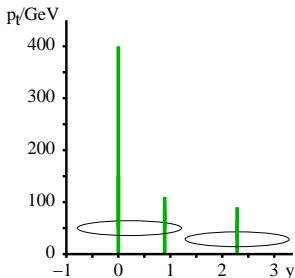
Midpoint cone alg. misses some stable cones; extra soft particle  $\rightarrow$  extra starting point  $\rightarrow$  extra stable cone found

**MIDPOINT IS INFRARED UNSAFE**

Or collinear unsafe with seed threshold



# Midpoint IR problem

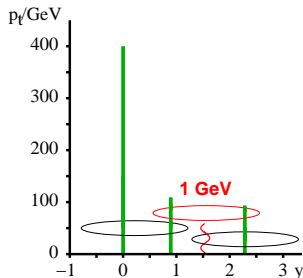


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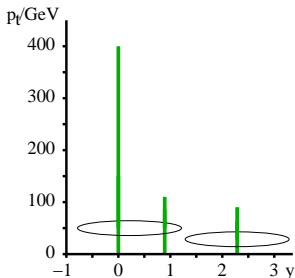
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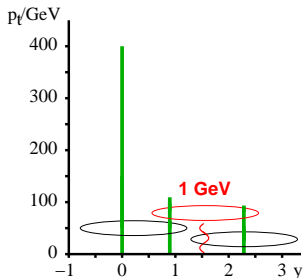
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# How much does IR safety *really* matter?

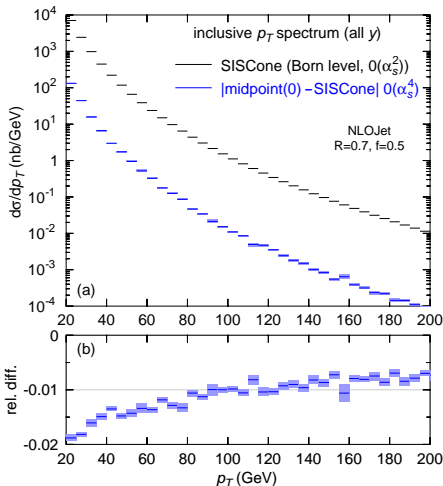
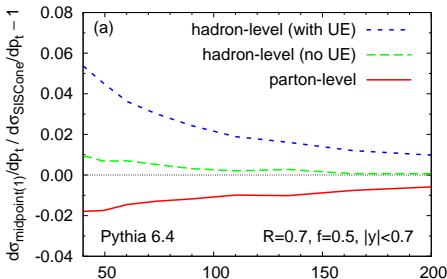
## Compare midpoint and SISCone

Result depends on observable:

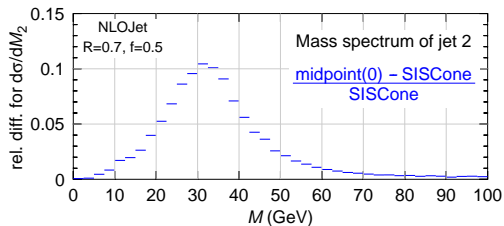
- ▶ inclusive jet spectrum is the least sensitive (affected at NNLO)
- ▶ larger differences (5 – 10%) at hadron level

seedless reduces UE effect

$p\bar{p}$   $\sqrt{s} = 1.96$  TeV



Look at jet masses in multijet events. **NB: Jet masses reconstruct boosted  $W/Z/H/top$  in BSM searches**



Select 3-jet events

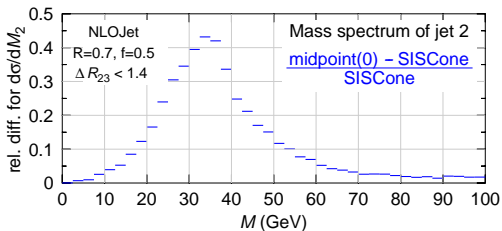
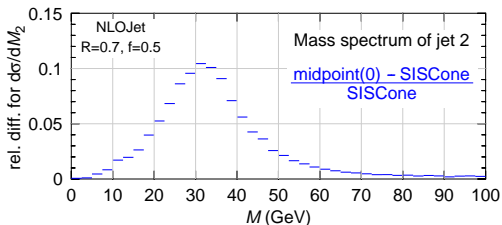
$$p_{t1,2,3} > \{120, 60, 20\} \text{ GeV,}$$

Calculate LO jet-mass spectrum for jet 2, compare midpoint with SISCone.

- ▶ 10% differences by default
- ▶ **40% differences** with extra cut  $\Delta R_{2,3} < 1.4$   
e.g. for jets from common decay chain

In complex events, IR safety matters

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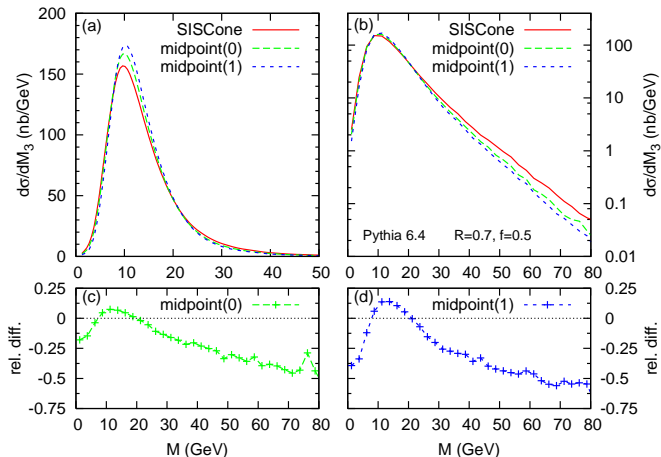
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**In complex events, IR safety matters**

Showering puts in many extra seeds: missing stable cones (in midpoint)  
should be less important?

Look at 3rd jet mass distribution (no  $\Delta R_{23}$  cut):



**Missing stable cones  $\rightarrow$  50% effects even after showering**

There are  $N(N - 1)/2$  distances  $d_{ij}$  — surely we have to calculate them all in order to find smallest?

$k_t$  distance measure is partly *geometrical*:

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*In words*: if  $i, j$  form smallest  $d_{ij}$  then  $j$  is geometrical nearest neighbour (GNN) of  $i$ .

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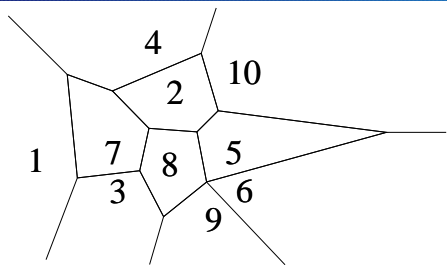
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Given a set of vertices on plane (1...10) a *Voronoi diagram* partitions plane into cells containing all points closest to each vertex

Dirichlet '1850, Voronoi '1908

A vertex's nearest other vertex is always in an adjacent cell.

E.g. GNN of point 7 will be found among 1,4,2,8,3 (it turns out to be 3)

Construction of Voronoi diagram for  $N$  points:  $N \ln N$  time      Fortune '88

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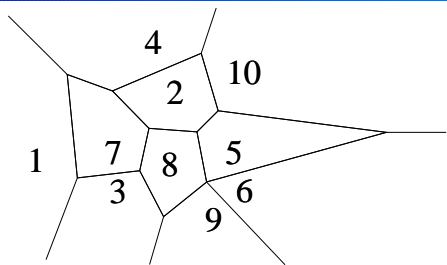
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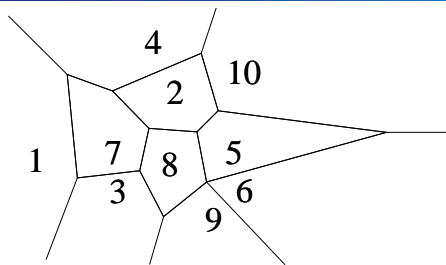
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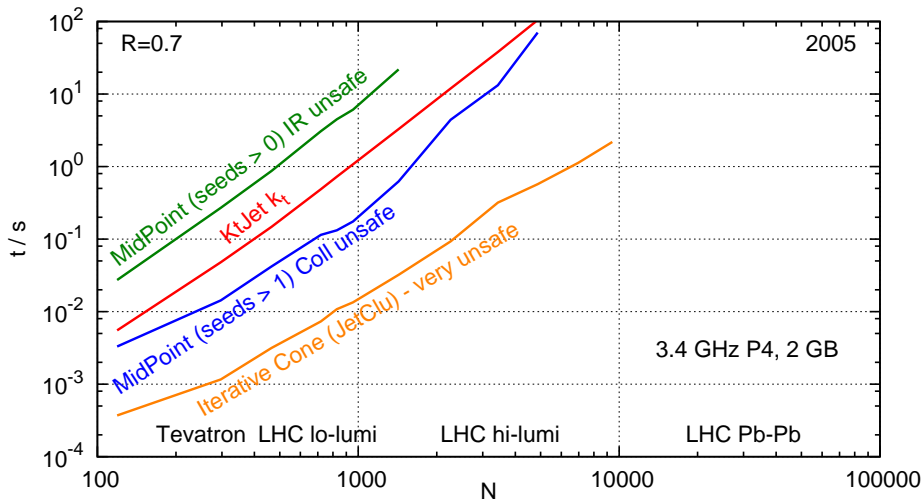
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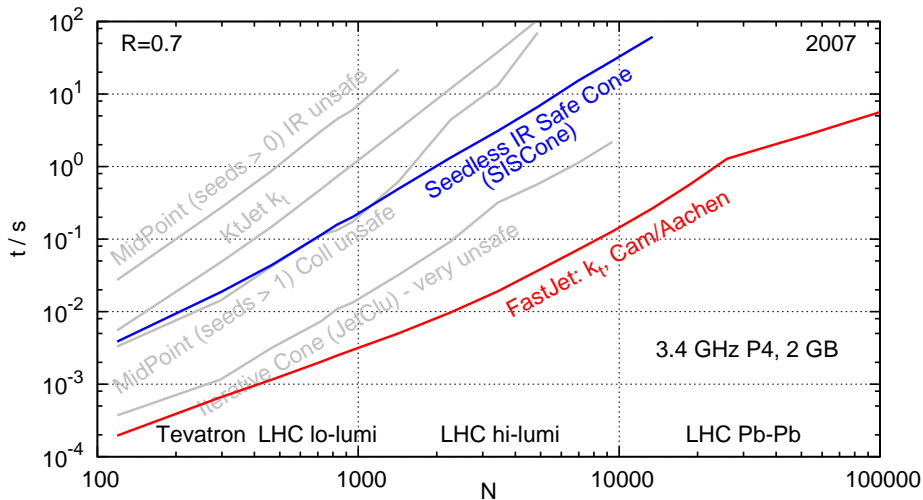
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Here: sum of squared means

Better still: calculate fluctuations

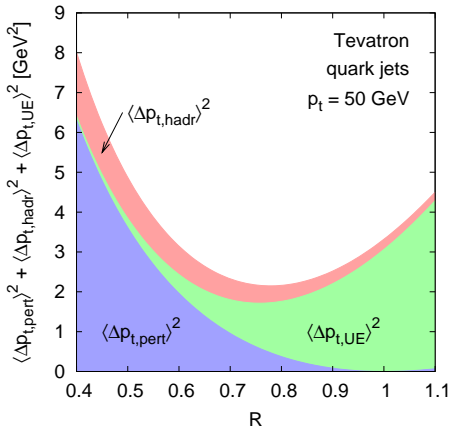
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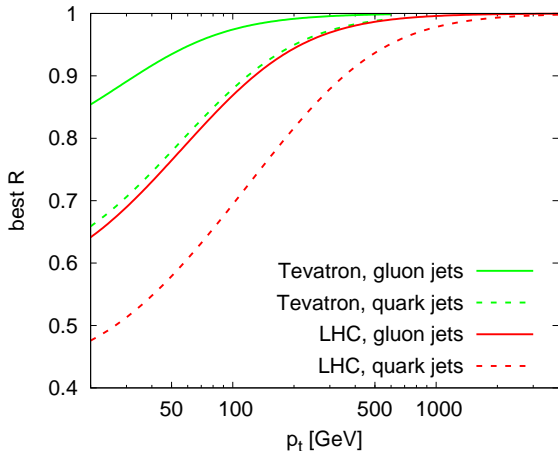
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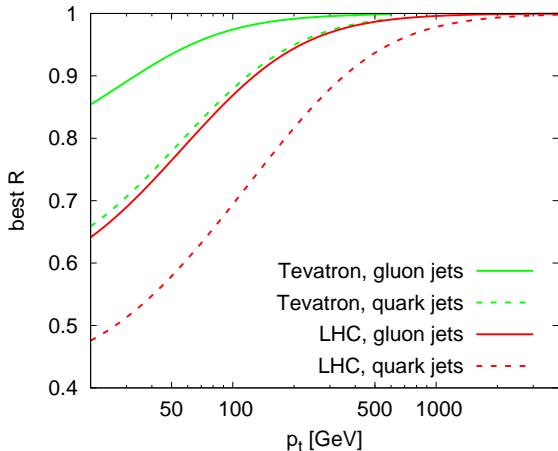




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- ▶ smaller  $R$  at LHC than Tevatron  
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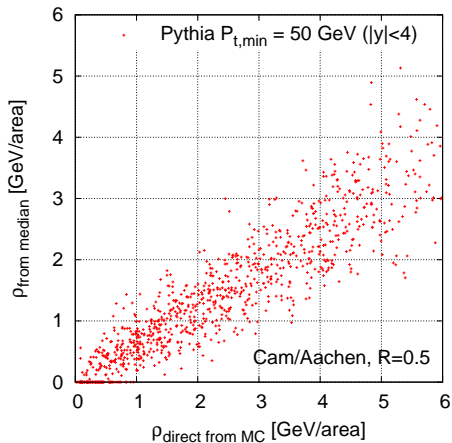
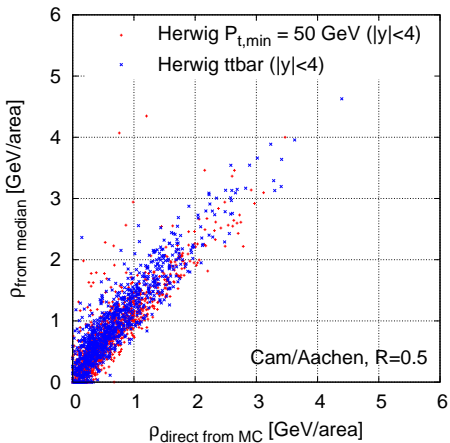


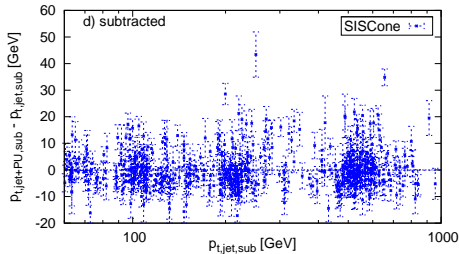
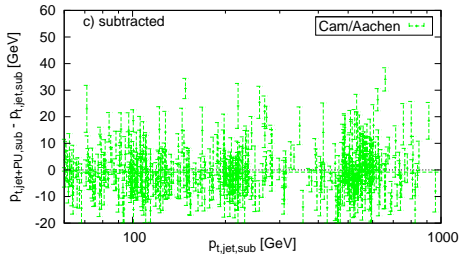
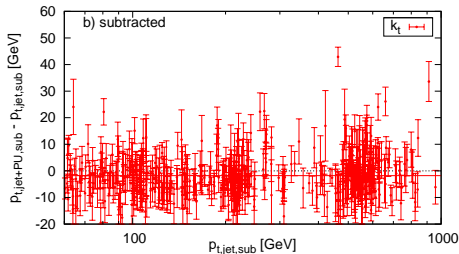
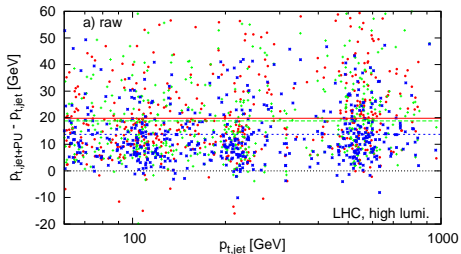
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# Correlate MC & measured UE (LHC)





How can high- $p_t$  HV do better than low  $p_t$  search?

- ▶ Acceptance improves at high- $p_t$ 
  - ▶ H and V both more likely to be central
  - ▶ H and V, since boosted, have central decay products
  - ▶ Decay products have high  $p_t$ , so you always see them
- ▶ Backgrounds fall faster than signal
  - ▶ for  $Vb\bar{b}$ , not gigantic effect (factor 2–3), but there
  - ▶ for  $t\bar{t}$ , top mass  $\rightarrow$   $b$ 's with  $p_t = 60 - 70$  GeV, just like light Higgs at rest; much harder to fake high- $p_t$  H
- ▶ At high- $p_t$  capture most of gluon radiation from  $H \rightarrow b\bar{b} + X$ , without too much UE; much harder at low  $p_t$ .  $\rightarrow$  Better mass resolution.

# $b$ -tagging, mass-dep for high- $p - t$ HV

