

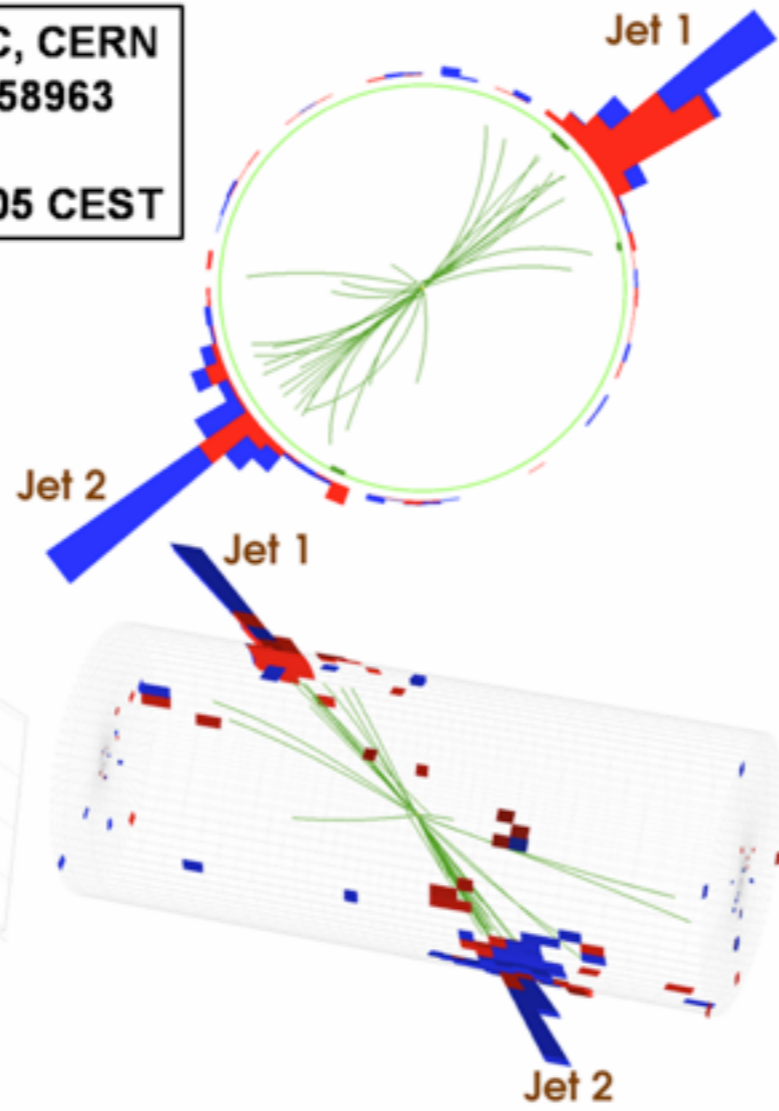
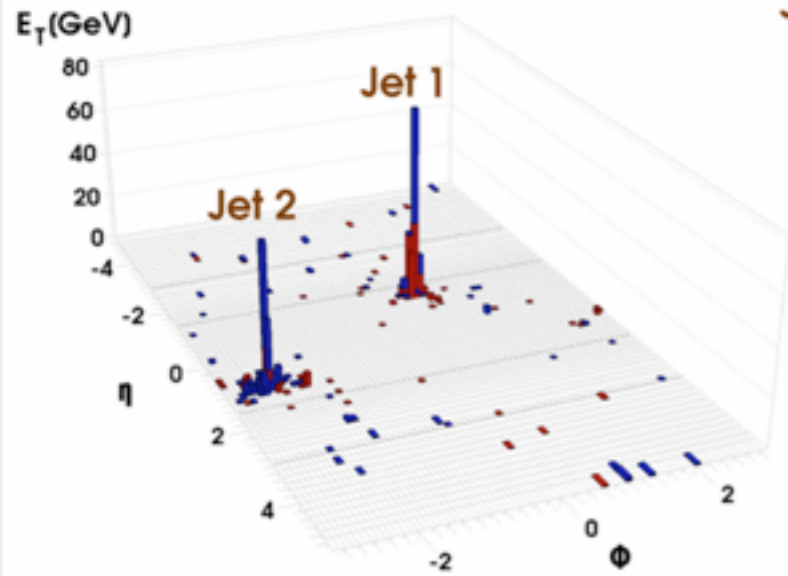
QCD lecture 8: jets

Gavin Salam, Oxford, February 2021
as part of Claire Gwenlan's QCD PhD course
(<http://www-pnp.physics.ox.ac.uk/~gwenlan/teaching/qcd.html>)

(with extensive use of material by
Matteo Cacciari and Gregory Soyez)

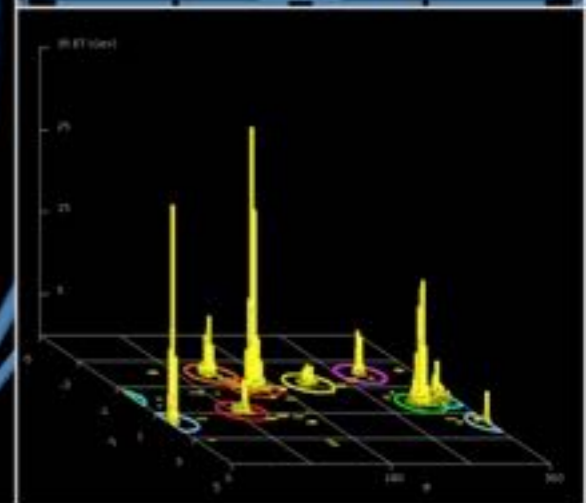
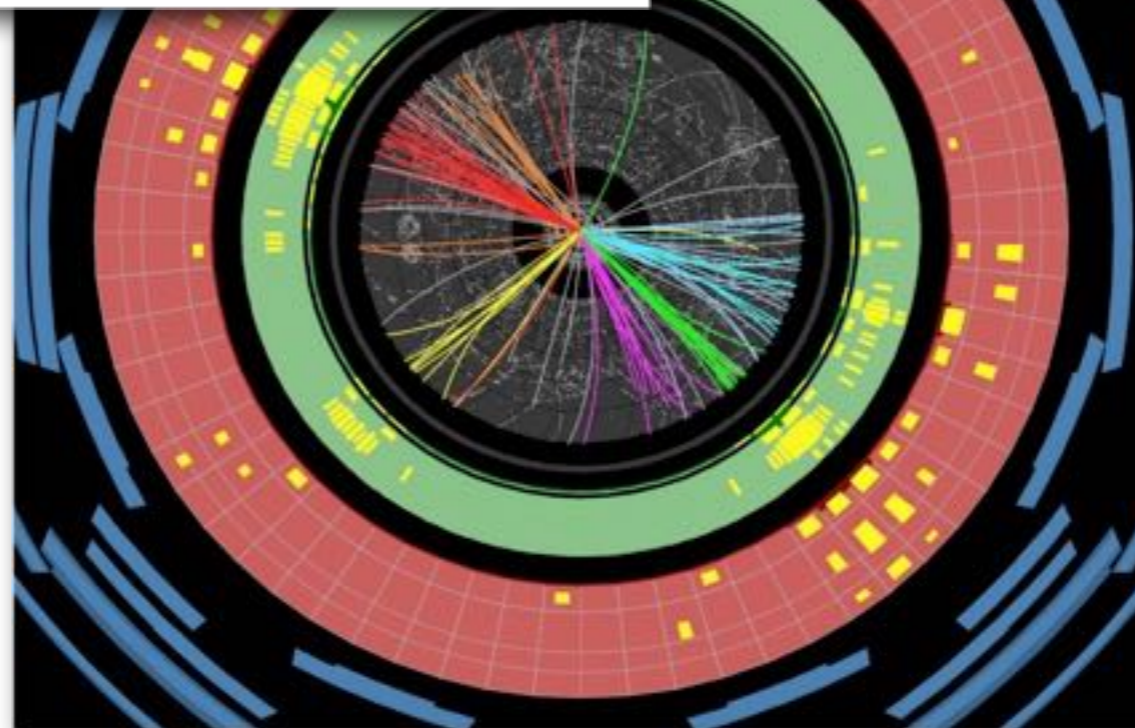
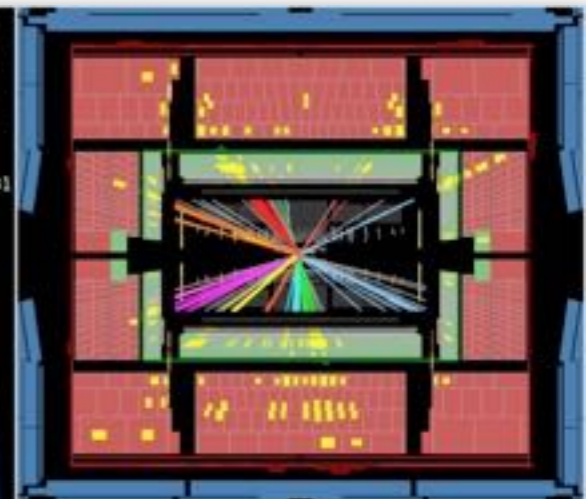


CMS Experiment at LHC, CERN
Run 133450 Event 16358963
Lumi section: 285
Sat Apr 17 2010, 12:25:05 CEST



JETS

Collimated, energetic bunches of particles








Find all papers by ATLAS and CMS




2252 records found

literature ▾ (collaboration:ATLAS or collaboration:CMS) and reportnumber:CERN*

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Emulating the impact of additional proton-proton interactions in the ATLAS simulation by pre-sampling sets of inelastic Monte Carlo events #1
ATLAS Collaboration • [Georges Aad \(Marseille, CPPM\)](#) et al. (Feb 18, 2021)
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Measurements of angular distance and momentum ratio distributions in three-jet and Z + two-jet final states in pp collisions #2
CMS Collaboration • [Albert M Sirunyan \(Yerevan Phys. Inst.\)](#) et al. (Feb 17, 2021)
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Pull out those that refer to one widely used jet-alg
1403 records found

literature ▾ (collaboration:ATLAS or collaboration:CMS) and reportnumber:CERN* and refersto:recid:779080

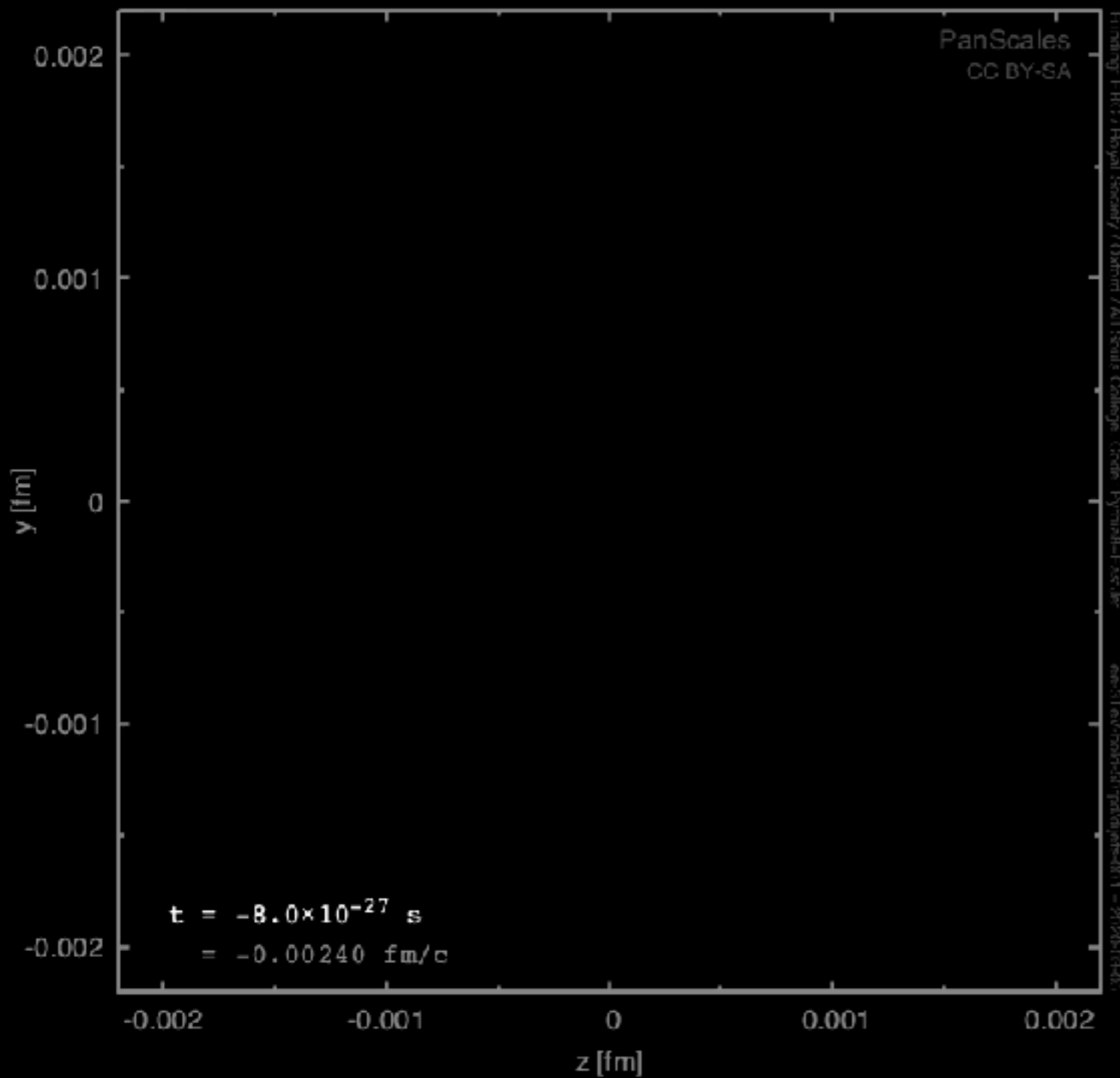
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Search for a heavy vector resonance decaying to a Z boson and a Higgs boson in proton-proton collisions at $\sqrt{s} = 13$ TeV #2
CMS Collaboration • Albert M Sirunyan (Yerevan Phys. Inst.) et al. (Feb 16, 2021)
e-Print: [2102.08198](#) [hep-ex]
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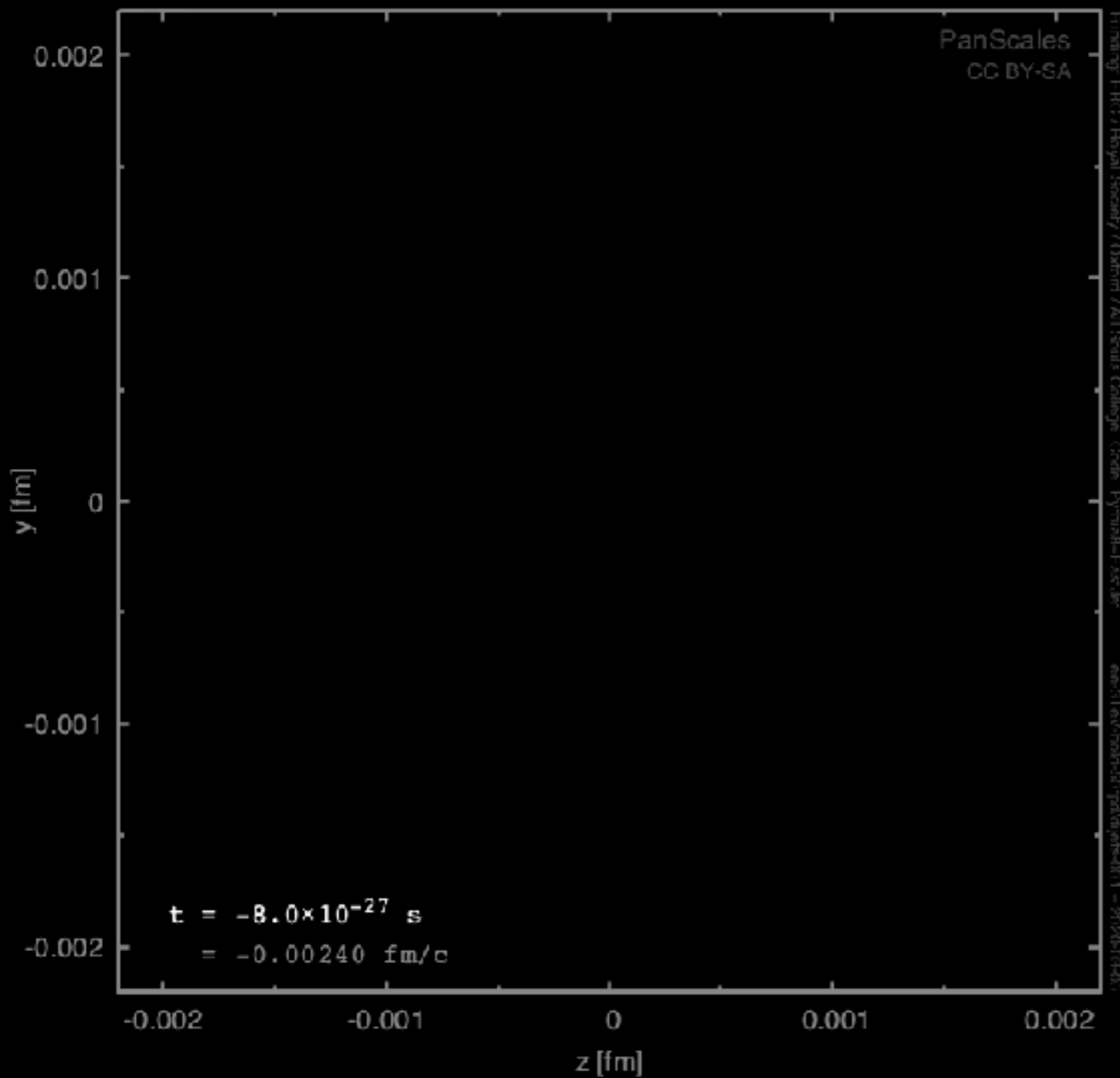
> 60% of papers use jets!



- incoming beam particle
- intermediate particle (quark or gluon)
- final particle (hadron)

Event evolution spans 7 orders of magnitude in space-time

<http://panscales.org/videos.html>



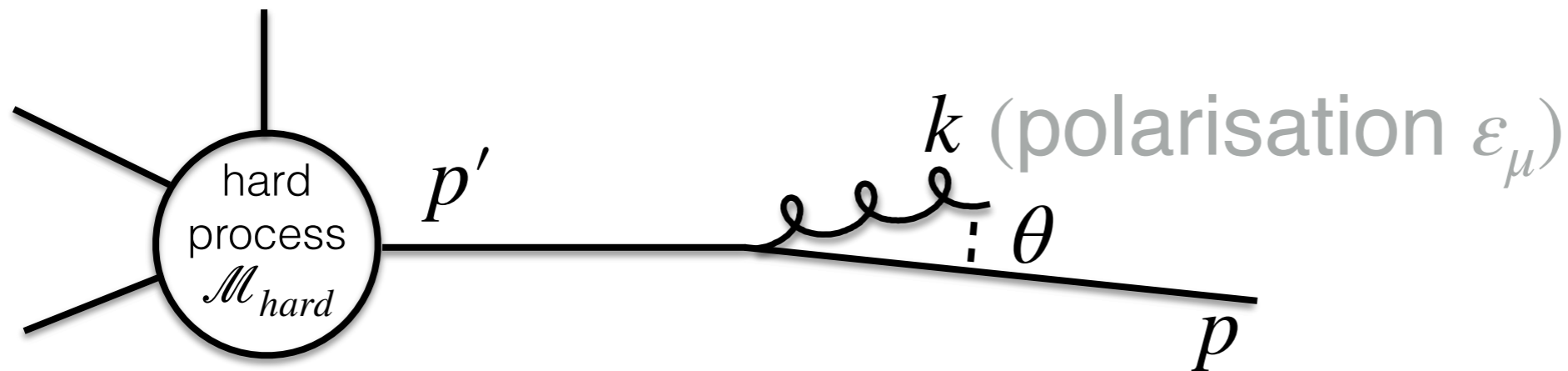
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why do quarks and
gluons fragment into jets?

Soft & collinear gluon emission



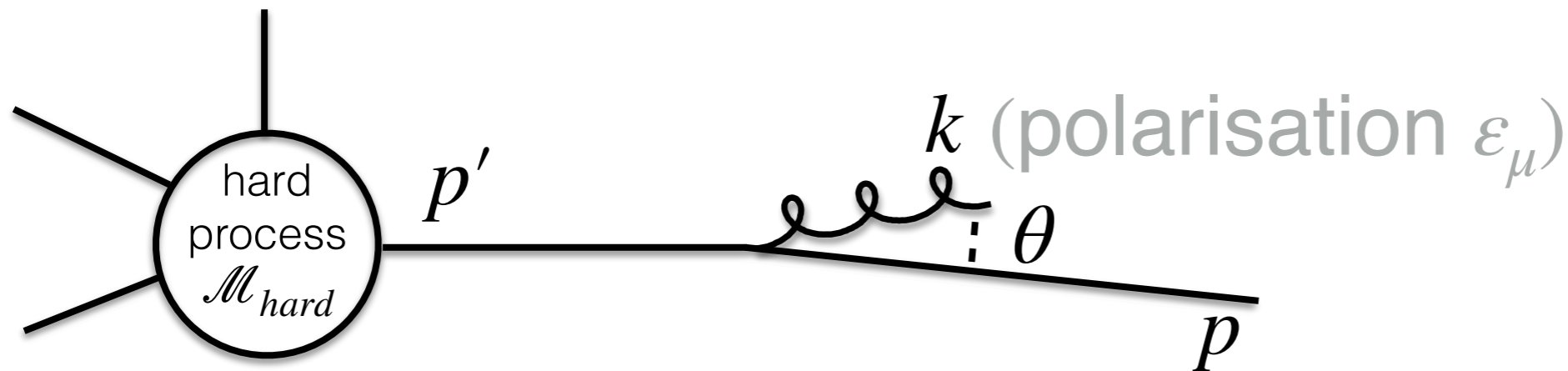
- **soft limit:** low-energy gluon emission: $E_k \ll E_p$
- **collinear limit:** small-angle emission, $\theta \ll 1$
- amplitude: (leaving out colour factors)

$$\mathcal{M} \propto \mathcal{M}_{hard} \cdot g_s \frac{p \cdot \epsilon}{p \cdot k} \simeq \mathcal{M}_{hard} \cdot g_s \frac{\sin \theta}{E_k (1 - \cos \theta)} \simeq \mathcal{M}_{hard} \cdot g_s \frac{2}{E_k \theta}$$

- phase space:

$$d\Phi \simeq d\Phi_{hard} \frac{E_k^2 dE_k}{(2\pi)^3 2E_k} d(\cos \theta) d\phi \simeq d\Phi_{hard} \cdot E_k dE_k \frac{\theta d\theta}{4\pi^2} \frac{d\phi}{2\pi}$$

Soft & collinear gluon emission



- Full result:

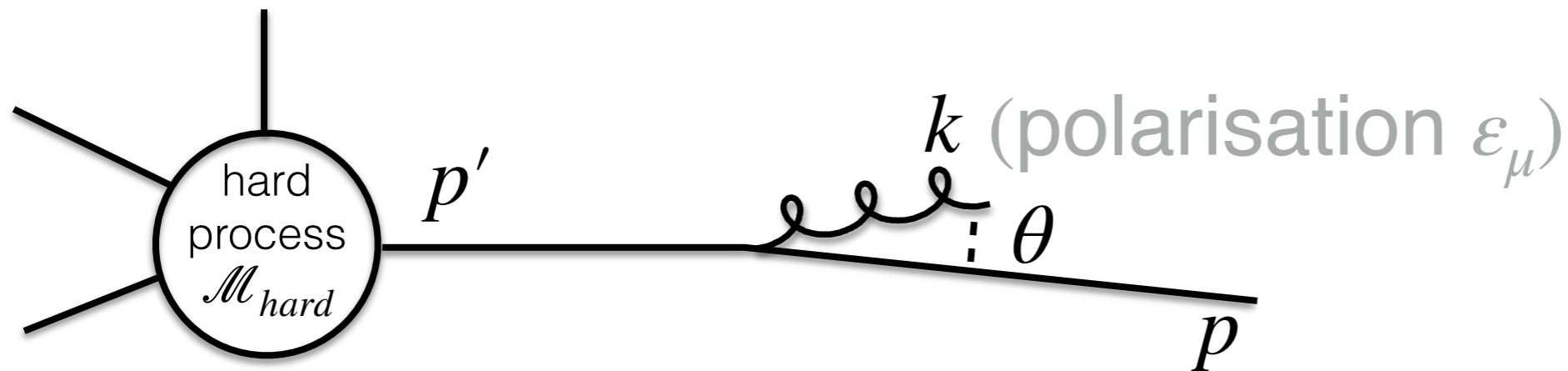
$$|\mathcal{M}|^2 d\Phi \simeq |\mathcal{M}_{hard}|^2 d\Phi_{hard} \cdot \frac{2C_F \alpha_s}{\pi} \frac{dE_k}{E_k} \frac{d\theta}{\theta}$$

- **factorises** into product of original hard matrix element and an additional soft-gluon emission probability

$$dP_{soft-gluon-emission} = \frac{2C_F \alpha_s}{\pi} \frac{dE_k}{E_k} \frac{d\theta}{\theta}$$

- **diverges** in soft ($E_k \rightarrow 0$) and collinear ($\theta \rightarrow 0$) limits

total probability of gluon emission



Total probability of gluon emission:

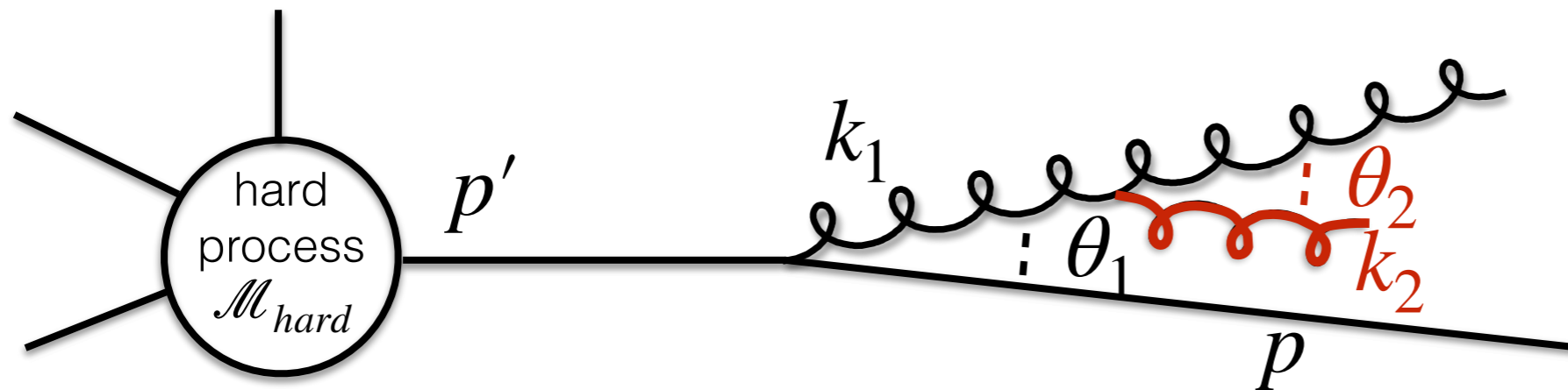
$$\langle N_{gluon} \rangle \simeq P_{gluon-emission} = \int dP = \frac{2\alpha_s C_F}{\pi} \int_{\Lambda_{QCD}/E_p}^1 \frac{d\theta}{\theta} \int_{\Lambda_{QCD}/\theta}^{E_p} \frac{dE_k}{E_k} = \frac{\alpha_s C_F}{\pi} \ln^2 \frac{E_p}{\Lambda_{QCD}}$$

Suppose (~wrongly!) that scale of the coupling is given by E_p , i.e. use $\alpha_s(E_p) = (2b_0 \ln E_p / \Lambda_{QCD})^{-1}$, then

$$\langle N_{gluon} \rangle \simeq P_{gluon-emission} \simeq \frac{C_F}{\pi b_0} \ln \frac{E_p}{\Lambda_{QCD}} \sim \frac{1}{\alpha_s} \gg 1$$

i.e. gluon emission is bound to happen and the average number of gluons is large

emission of gluon from gluon?



Emission of gluon 1 from gluon 2 also **factorises**, with colour factor $C_A = 3$, instead of $C_F = 4/3$:

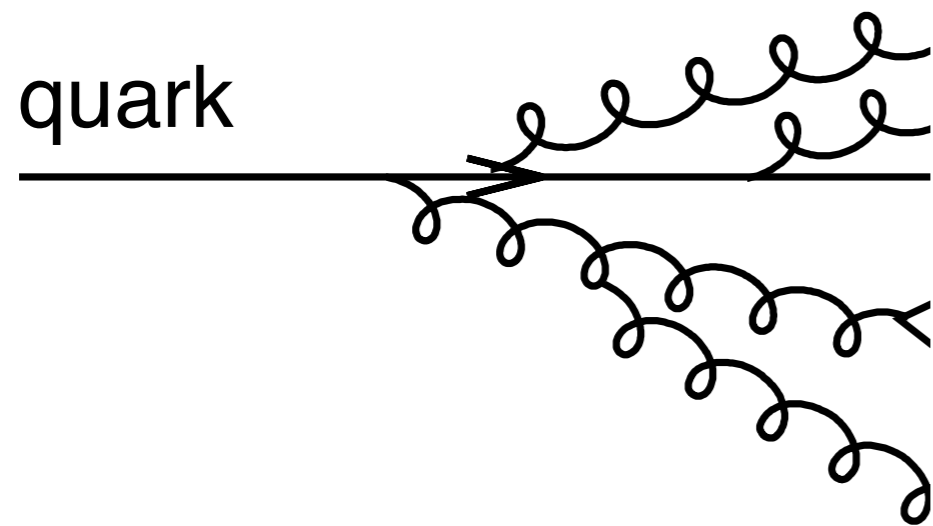
$$dP_{\text{gluon-2-from-gluon-1}} = \frac{2C_A \alpha_s}{\pi} \frac{dE_2}{E_2} \frac{d\theta_2}{\theta_2}$$

Additional gluon radiation due to emission from gluon 1 is confined within a cone of angle $\theta_2 \lesssim \theta_1$ (a.k.a. **Angular Ordering**) and will have energy $E_2 \lesssim E_1$

Total extra number of gluons emitted from gluon 1:

$$\frac{\alpha_s C_A}{\pi} \ln^2 \frac{E_1 \theta_1}{\Lambda_{QCD}}$$

Why do we see jets?

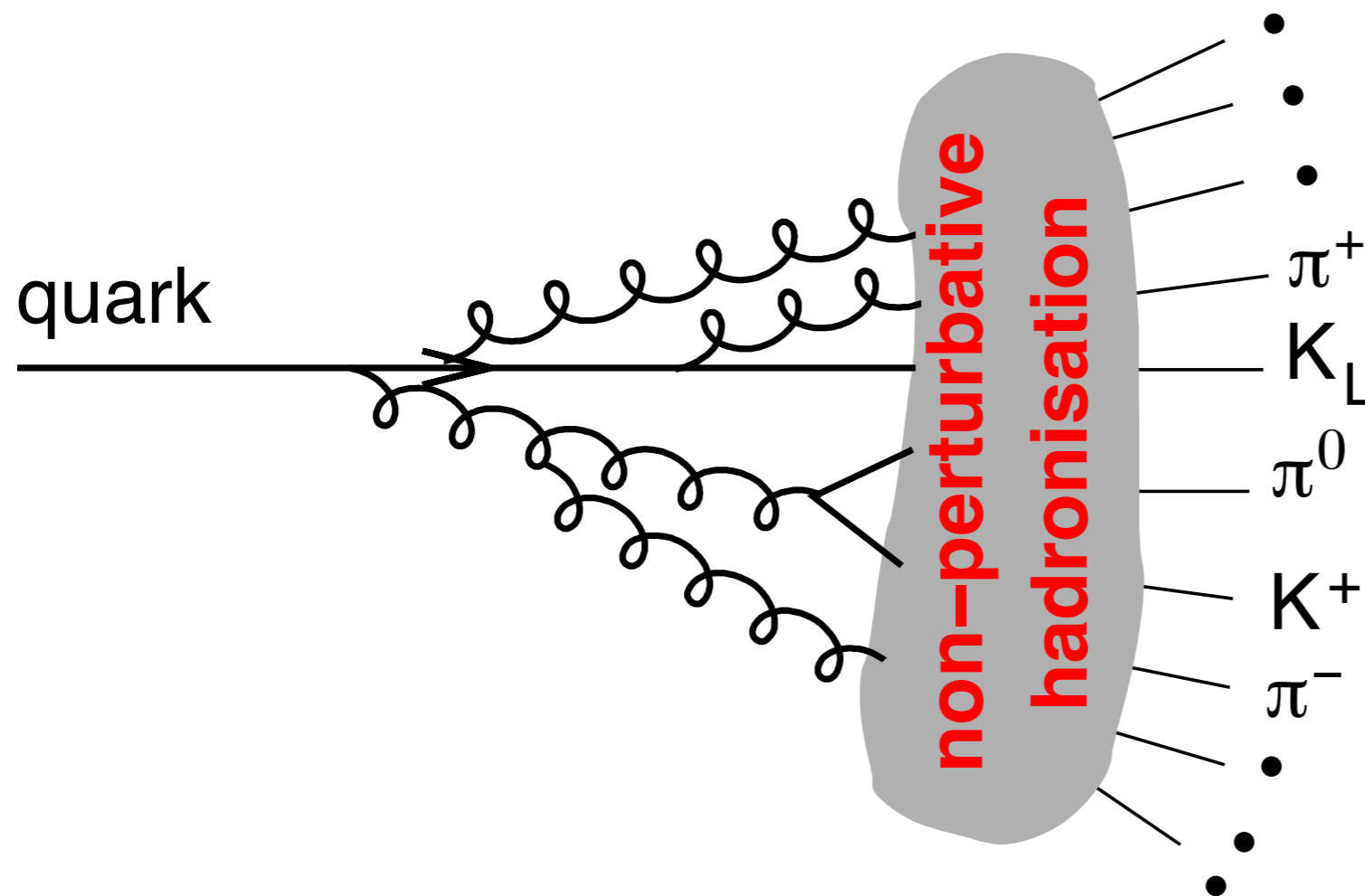


Starting from energetic quark, emit a cascade of many low-energy (soft) and small-angle (collinear) gluons

$$\frac{\alpha_s C_{F/A}}{\pi} \frac{dE}{E} \frac{d\theta}{\theta}$$

giving a collimated jet of partons

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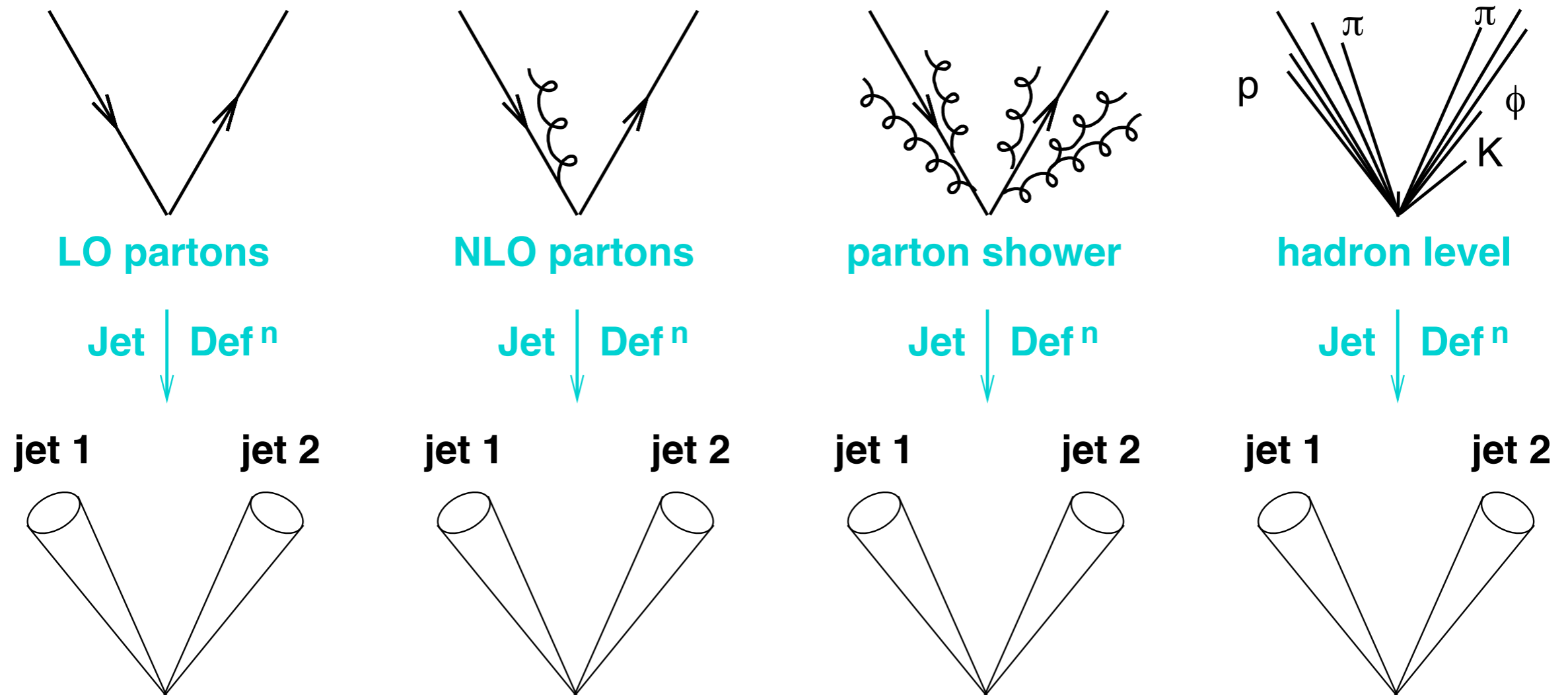
giving a collimated jet of partons

When the partons become separated by a distance $\sim 1/\Lambda_{QCD}$ they **confine** into hadrons.

The hadrons go in similar directions to the partons

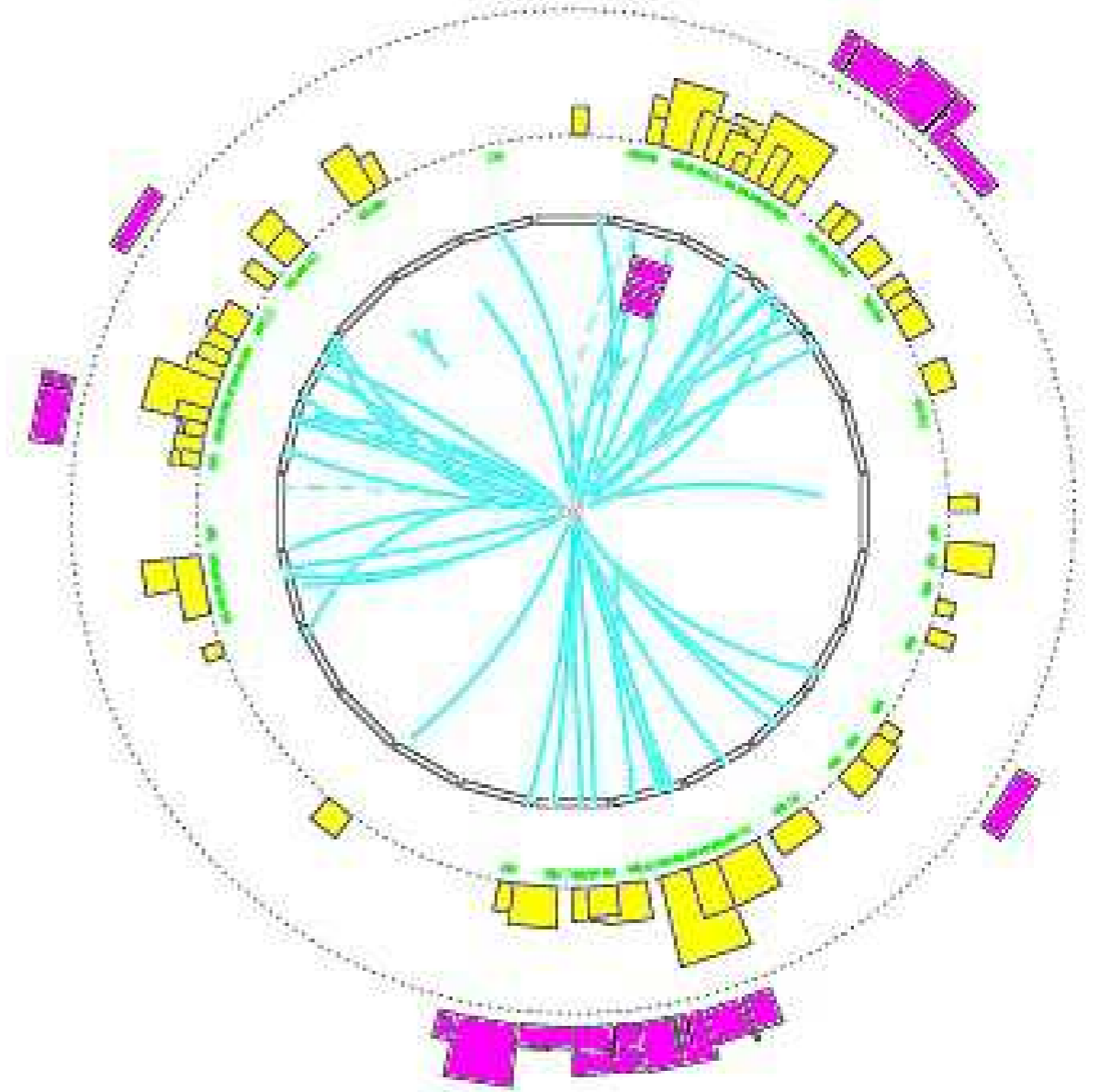
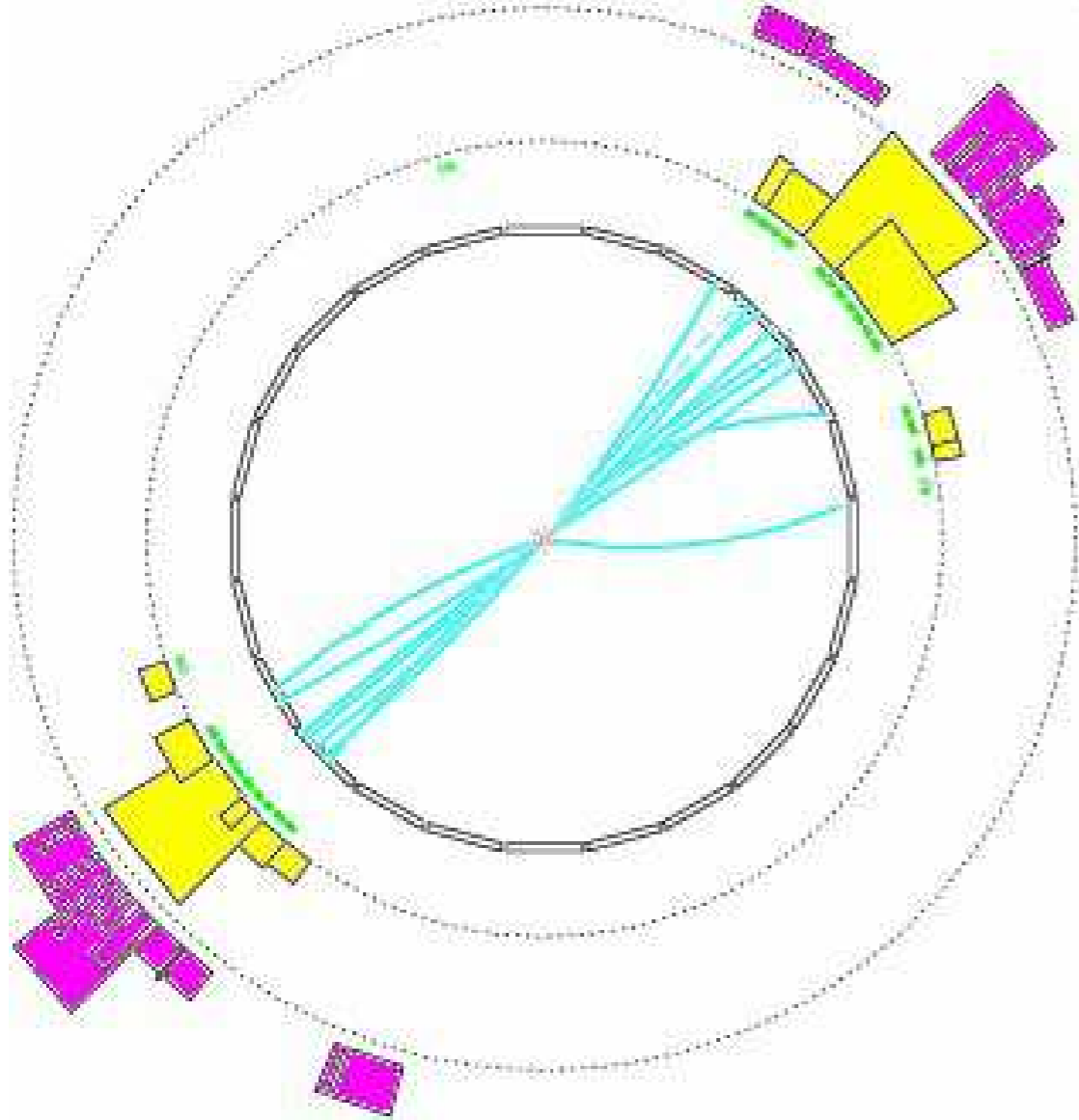
core ideas in jet reconstruction

Jet finding as a form of projection

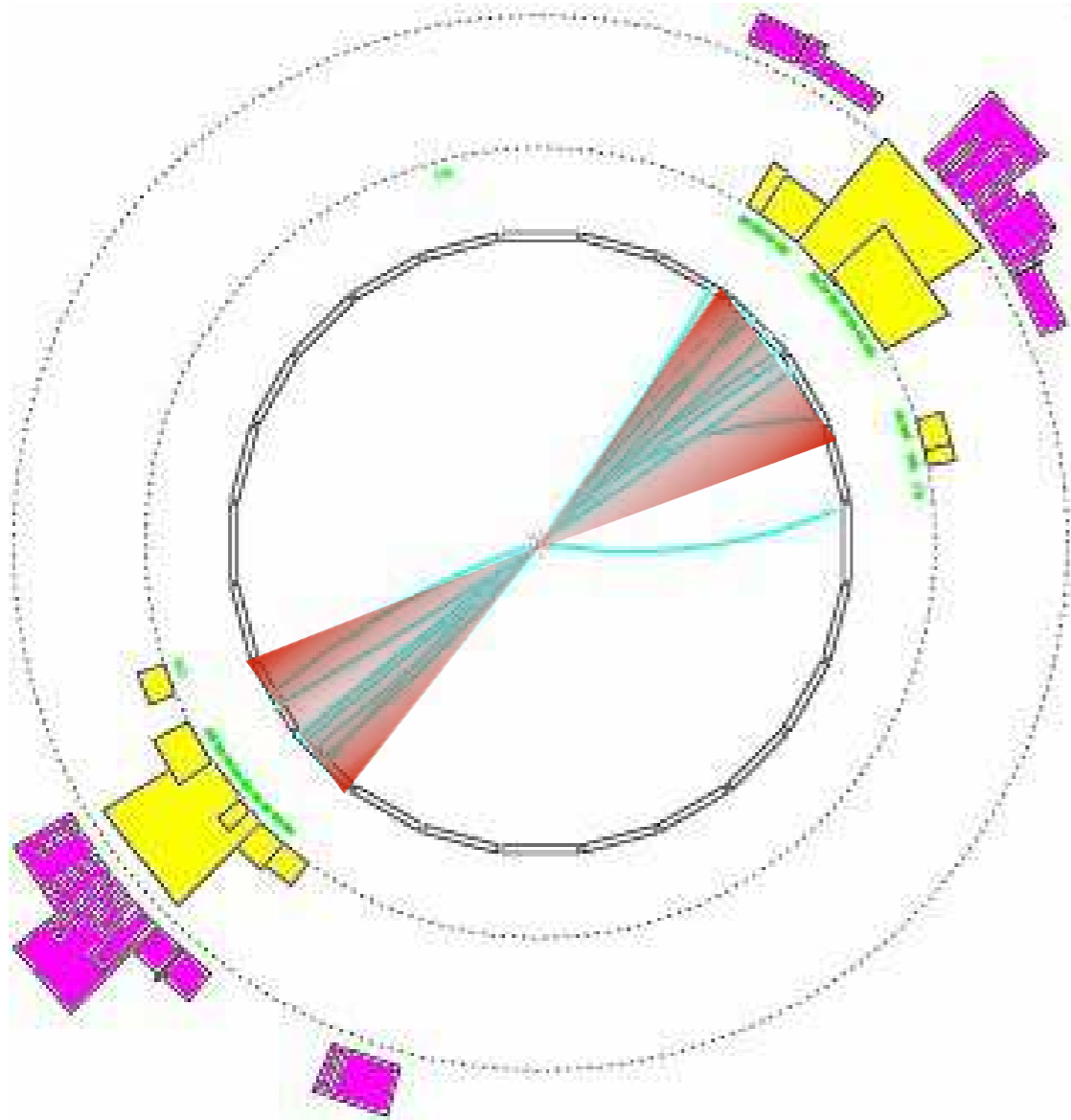


Projection to jets should be resilient to QCD effects

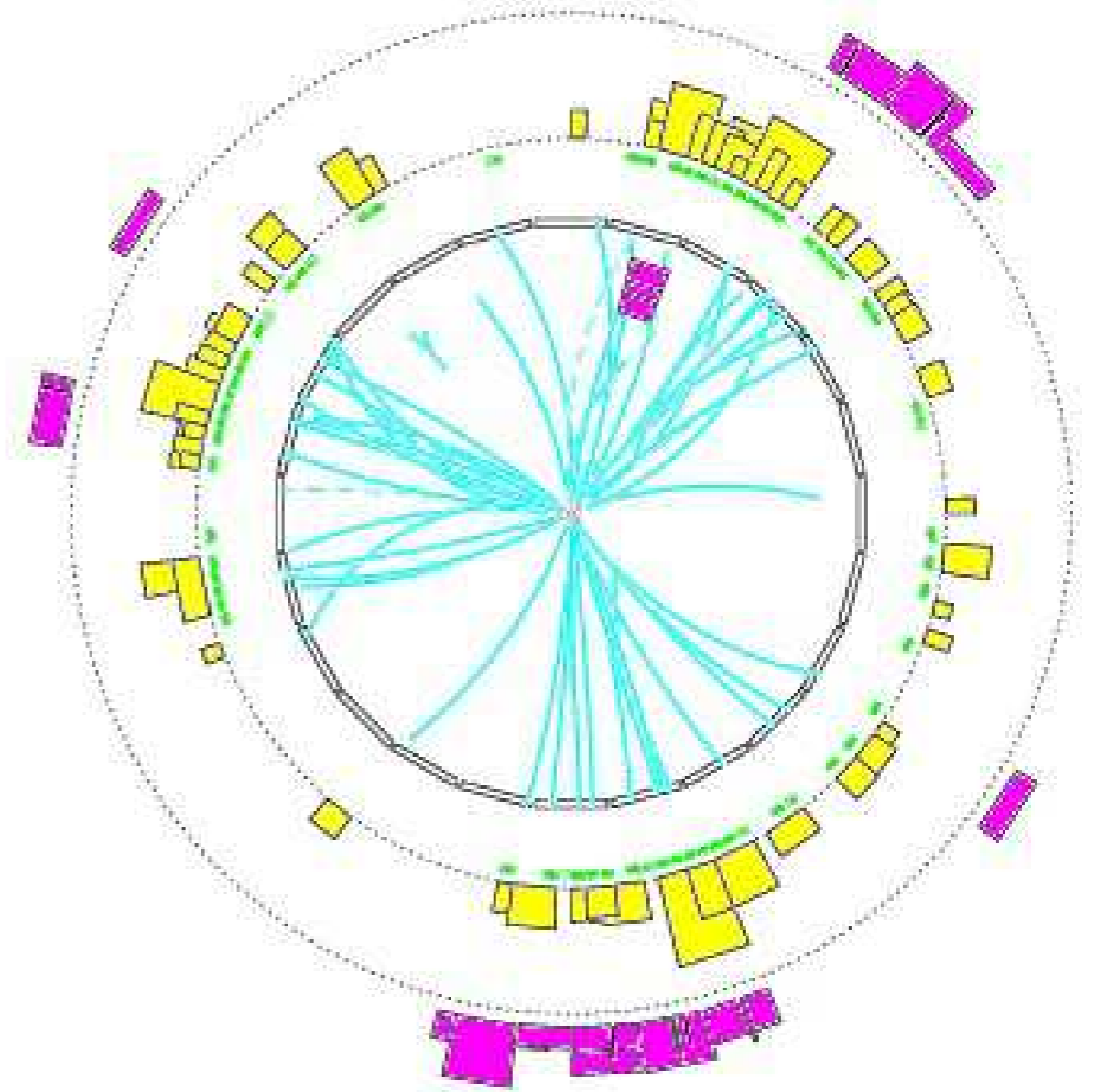
Reconstructing jets is an ambiguous task



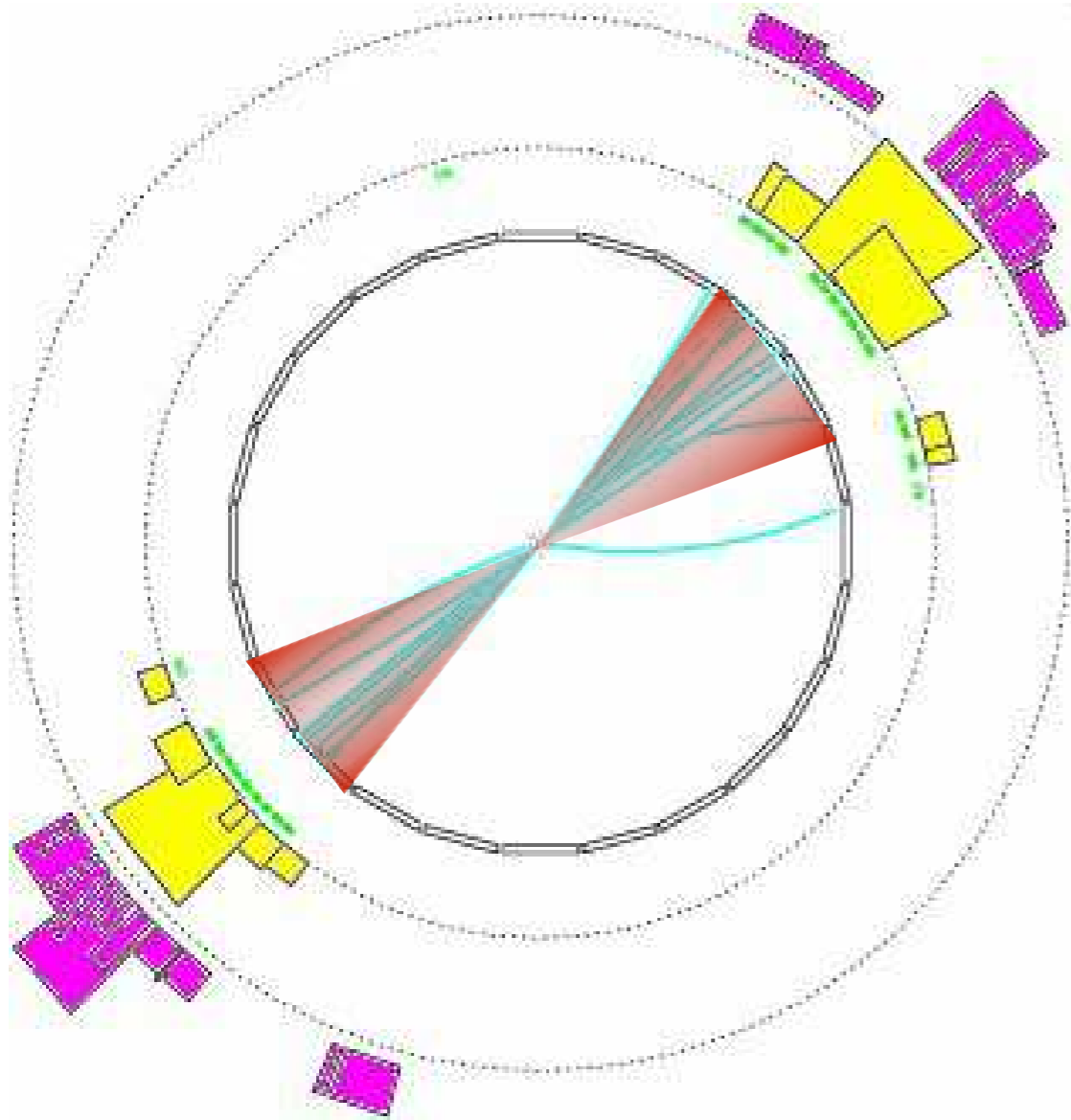
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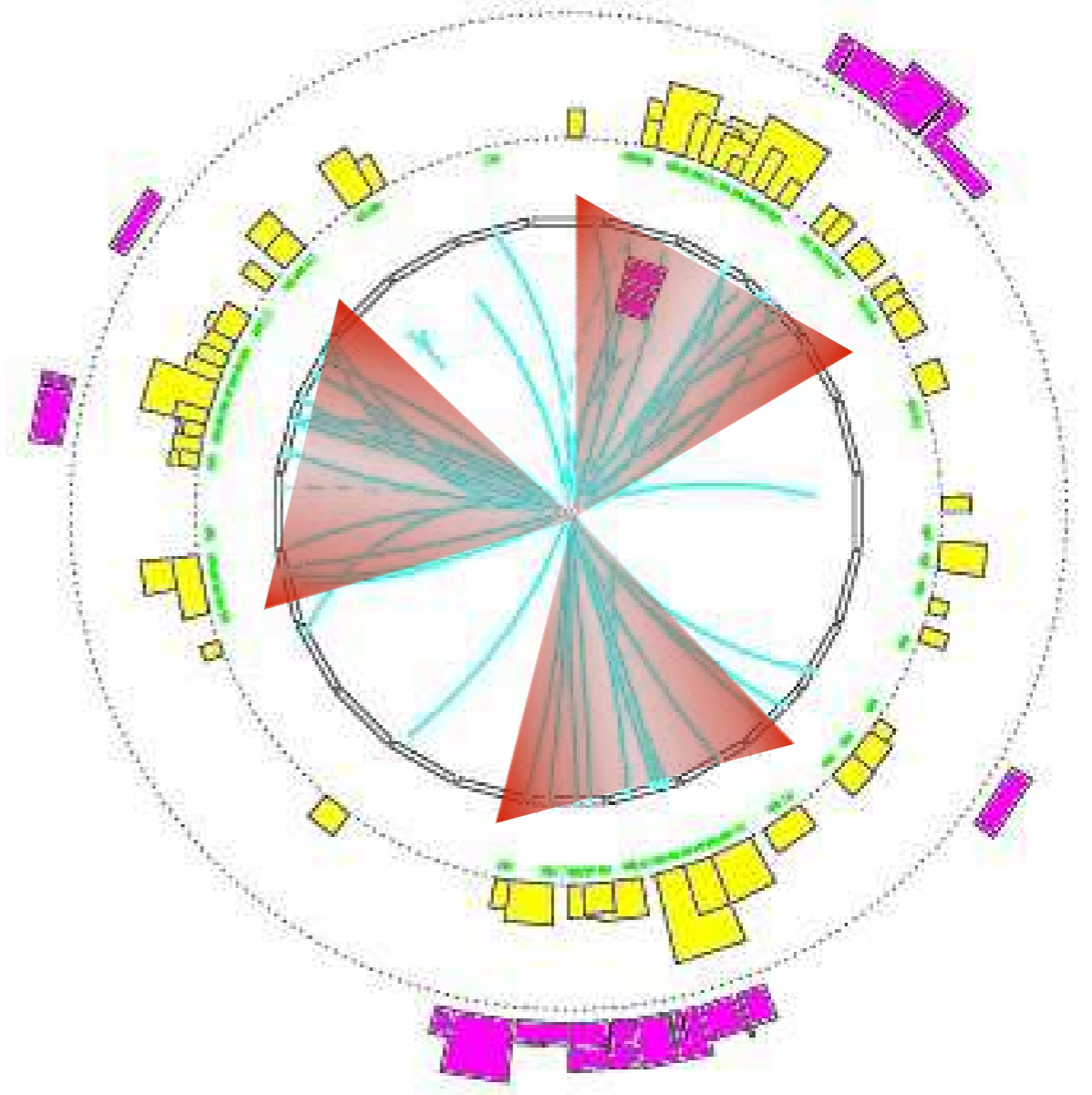
2 clear jets



Reconstructing jets is an ambiguous task

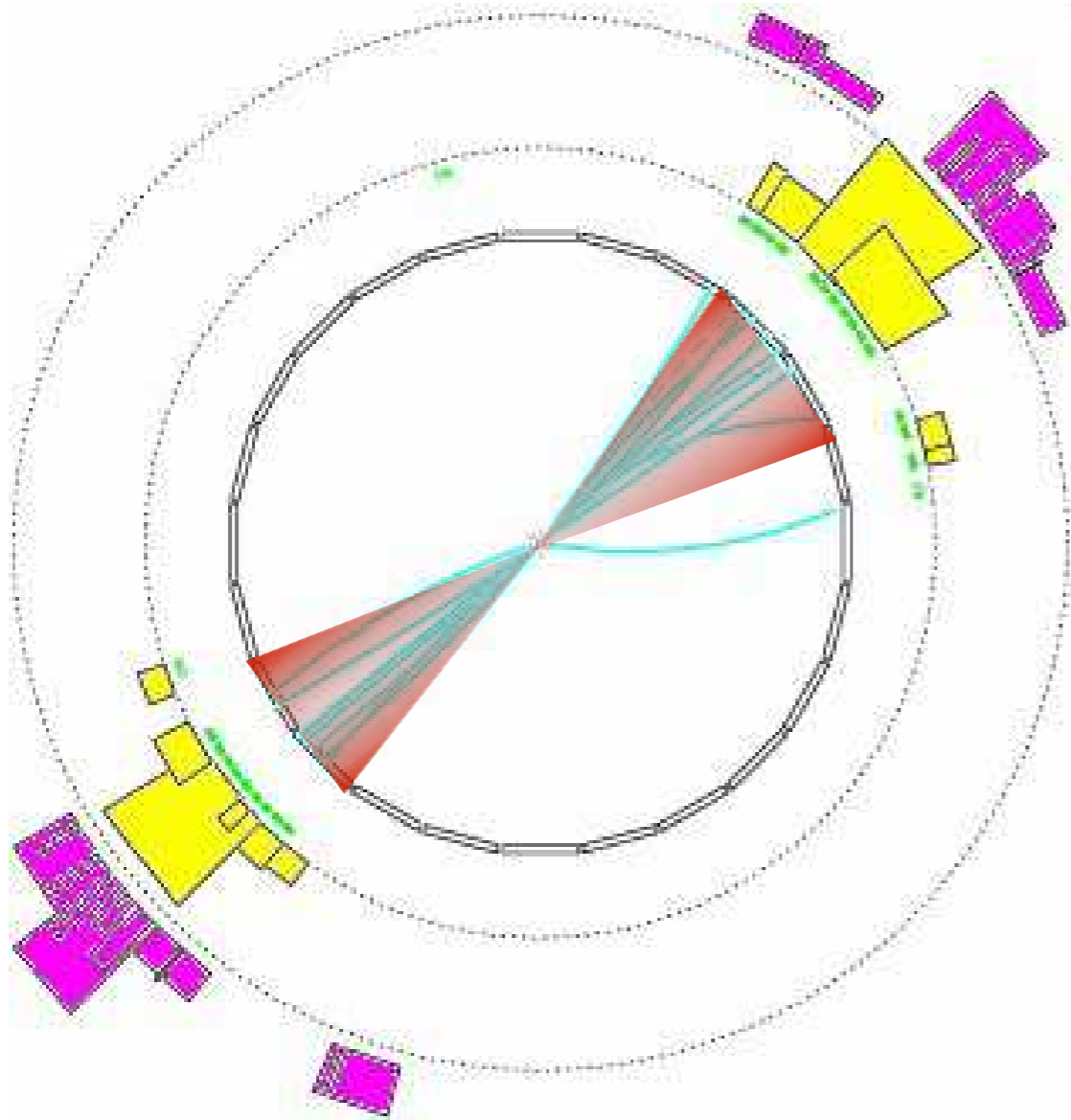


2 clear jets

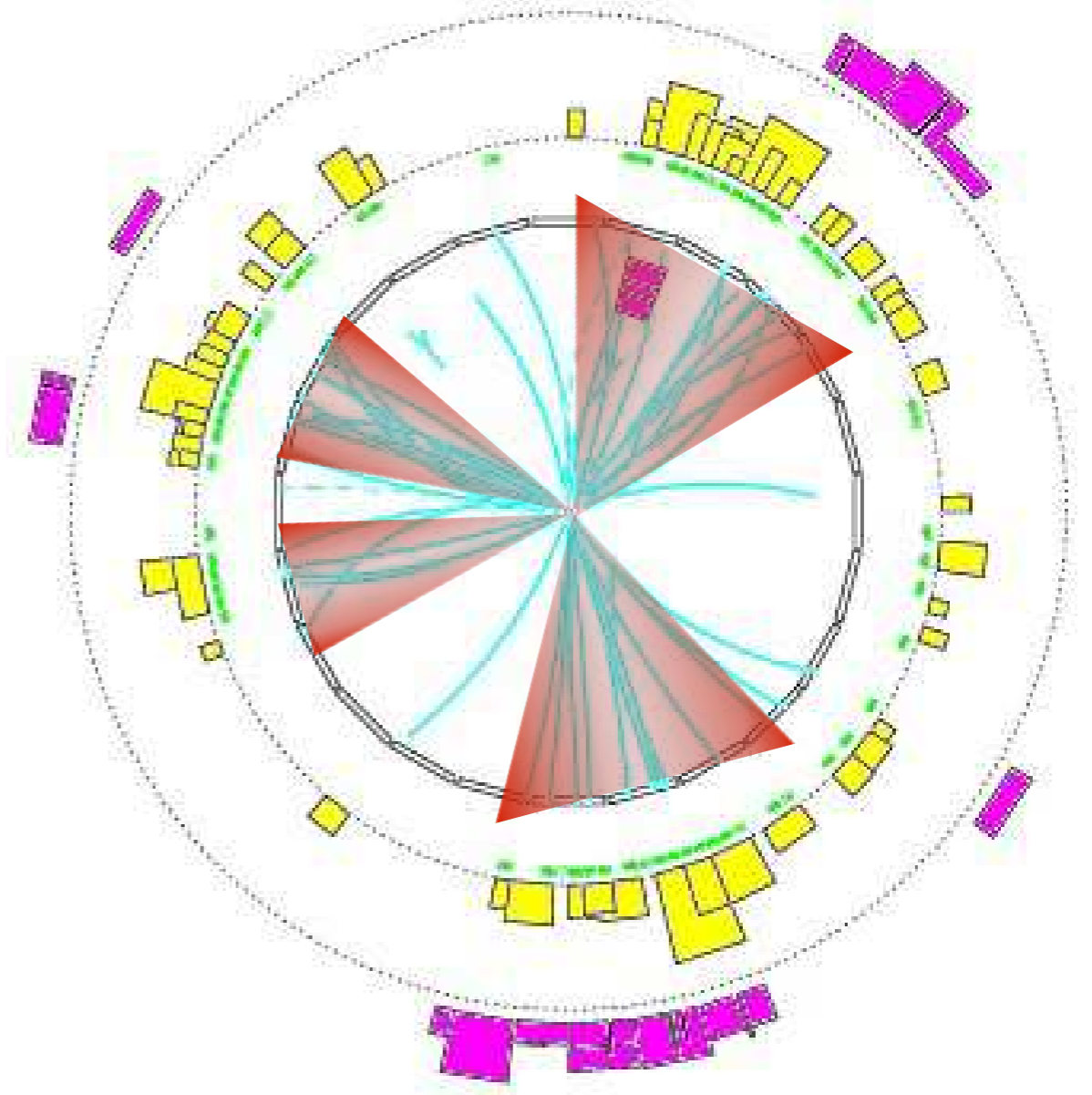


3 jets?

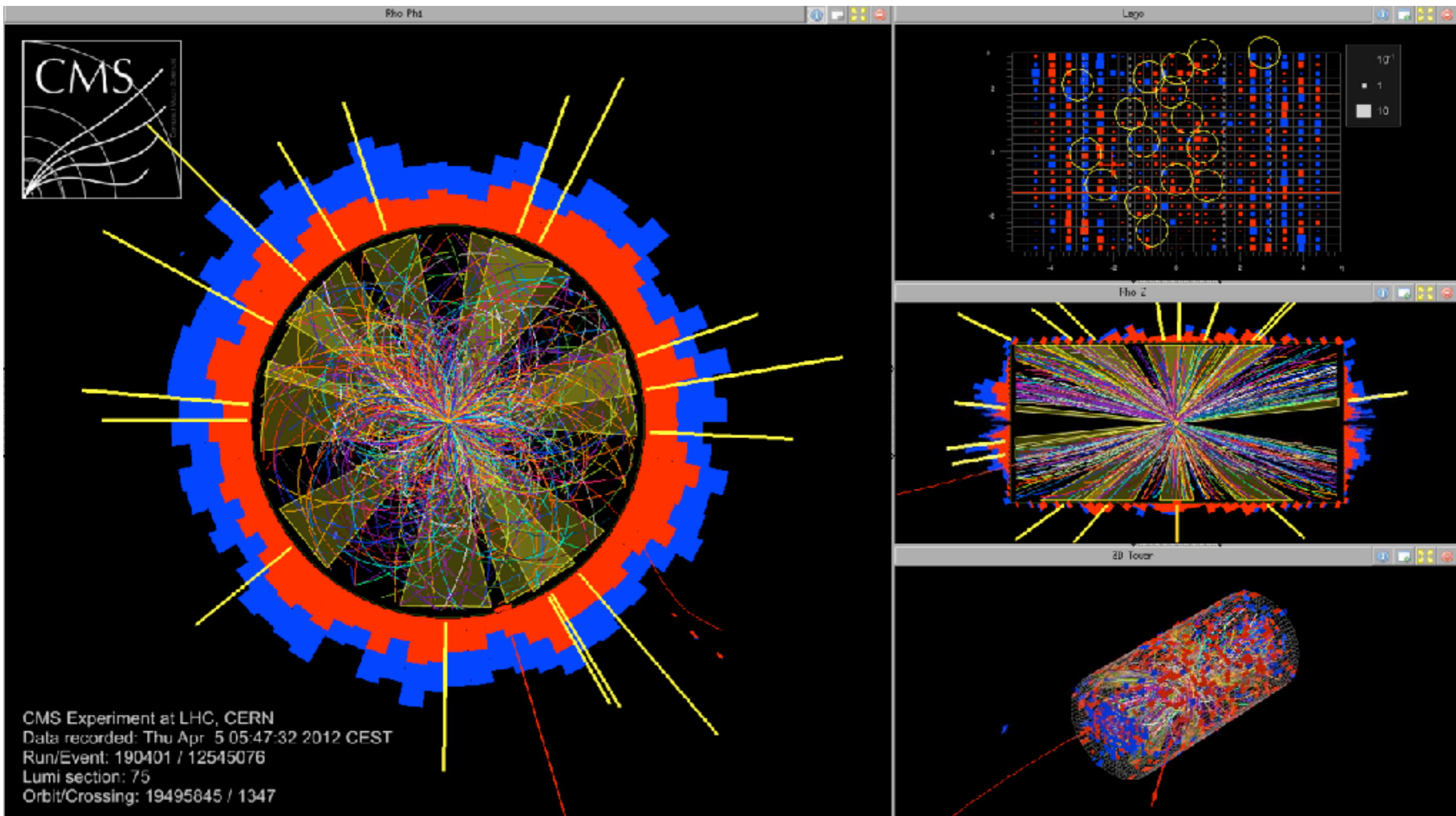
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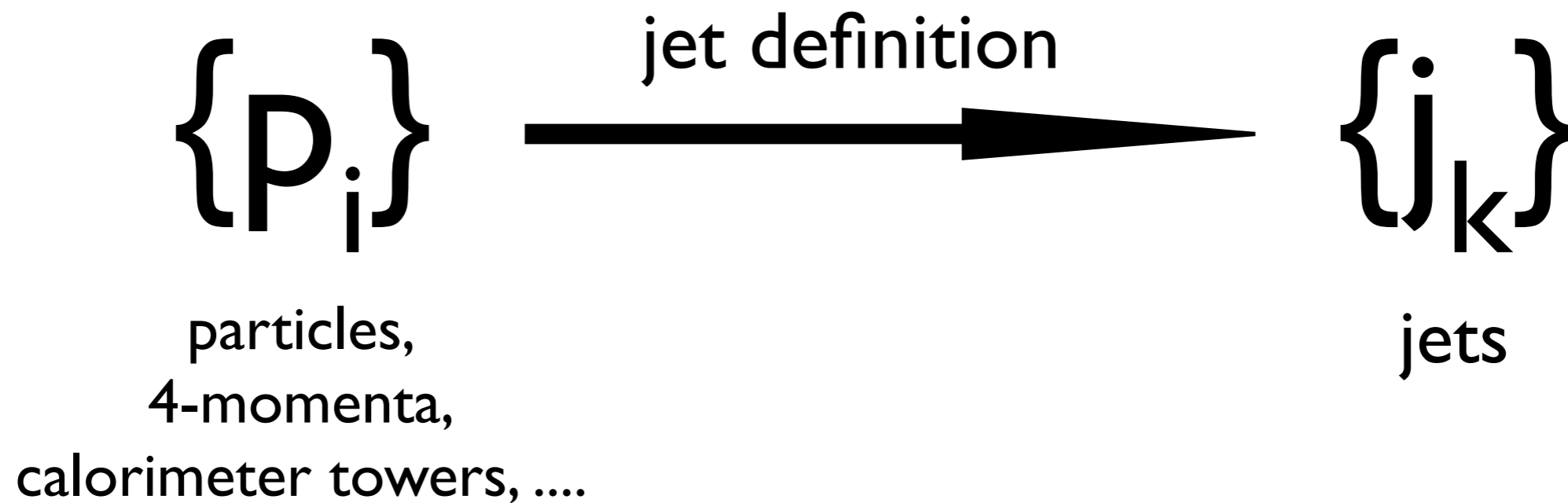
2 clear jets



3 jets?
or 4 jets?



Make a choice: specify a jet definition



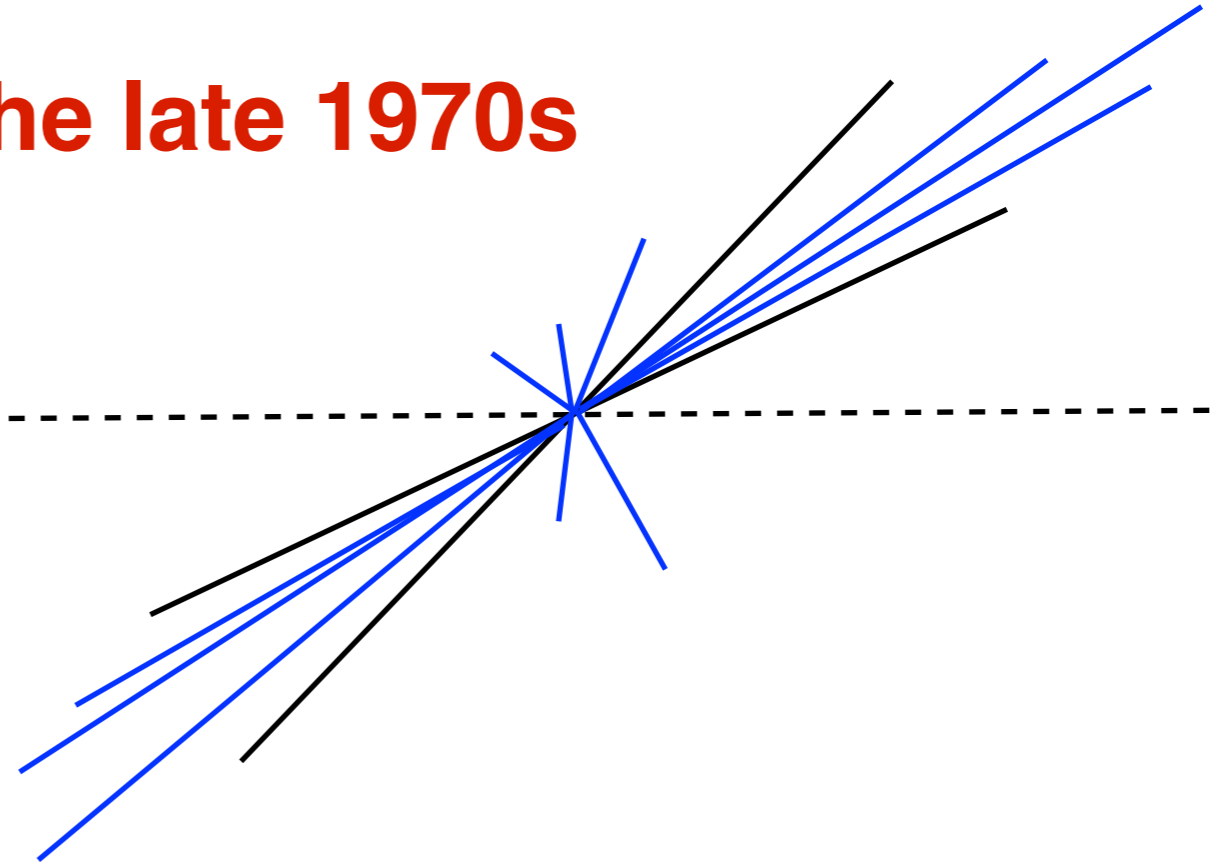
- Which particles do you put together into a same jet?
- How do you recombine their momenta (4-momentum sum is the obvious choice, right?)

“Jet [definitions] are legal contracts between theorists and experimentalists”
-- MJ Tannenbaum

They're also a way of organising the information in an event
1000's of particles per events, up to 40.000,000 events per second

Jet definitions date back to the late 1970s

Sterman and Weinberg,
Phys. Rev. Lett. 39, 1436 (1977):



To study jets, we consider the partial cross section

$\sigma(E, \theta, \Omega, \epsilon, \delta)$ for e^+e^- hadron production events, in which all but

a fraction $\epsilon \ll 1$ of the total e^+e^- energy E is emitted within

some pair of oppositely directed cones of half-angle $\delta \ll 1$,

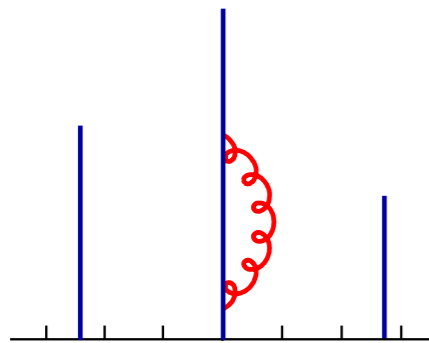
lying within two fixed cones of solid angle Ω (with $\pi\delta^2 \ll \Omega \ll 1$)

at an angle θ to the e^+e^- beam line. We expect this to be measur-

$$\sigma(E, \theta, \Omega, \epsilon, \delta) = (d\sigma/d\Omega)_0 \Omega \left[1 - (g_E^2/3\pi^2) \left\{ 3\ln \delta + 4\ln \delta \ln 2\epsilon + \frac{\pi^3}{3} - \frac{5}{2} \right\} \right]$$

Key requirement: infrared and collinear safety

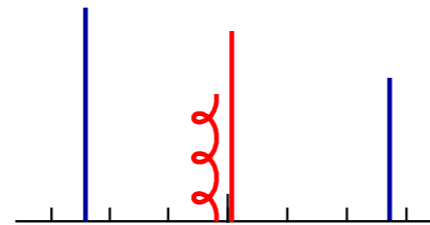
Collinear Safe



jet 1

$$\alpha_s^n \times (-\infty)$$

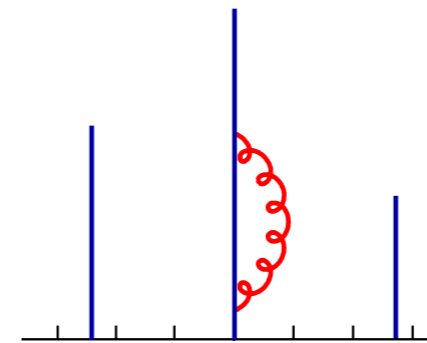
Infinities cancel



jet 1

$$\alpha_s^n \times (+\infty)$$

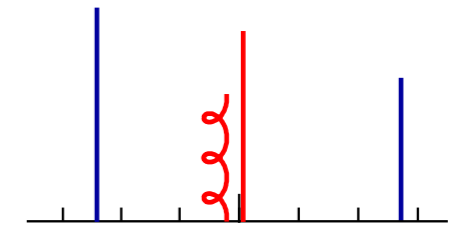
Collinear Unsafe



jet 1

$$\alpha_s^n \times (-\infty)$$

Infinities do not cancel



jet 1

jet 2

$$\alpha_s^n \times (+\infty)$$

Invalidates perturbation theory

hadron collider jet algorithms

Two parameters, R (jet opening angle) **and $p_{t,min}$**
(These are the two parameters in essentially every widely used hadron-collider jet algorithm)

Sequential recombination algorithm

$$d_{ij} = \min(p_{ti}^2, p_{tj}^2) \frac{\Delta R_{ij}^2}{R^2}, \quad d_{iB} = p_{ti}^2, \quad \Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$$

Define d_{ij} distance between every pair of particles as squared transverse momentum (k_t) of softer of i and j relative to harder one.

It will be small when particles are collinear ($\Delta R_{ij} \ll 1$) or if one of the particles is soft (p_{ti} or $p_{tj} \ll$ hard scale).

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1. Find smallest of d_{ij}, d_{iB}

2. If ij , recombine particles i and j

3. If iB , call i a jet and remove from list of particles

4. repeat from step 1 until no particles left

Only use jets with $p_t > p_{t,min}$

Inclusive k_t algorithm

S.D. Ellis & Soper, 1993

Catani, Dokshitzer, Seymour & Webber, 1993

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1. Find smallest of d_{ij} , d_{iB}
2. If ij , recombine particles i and j
3. If iB , call i a jet and remove from list of particles

If a particle i has no neighbours j within a distance $\Delta R_{ij} \leq R$, then $d_{iB} < \text{all } d_{ij}$, and i becomes a jet.

k_t alg.: Find smallest of

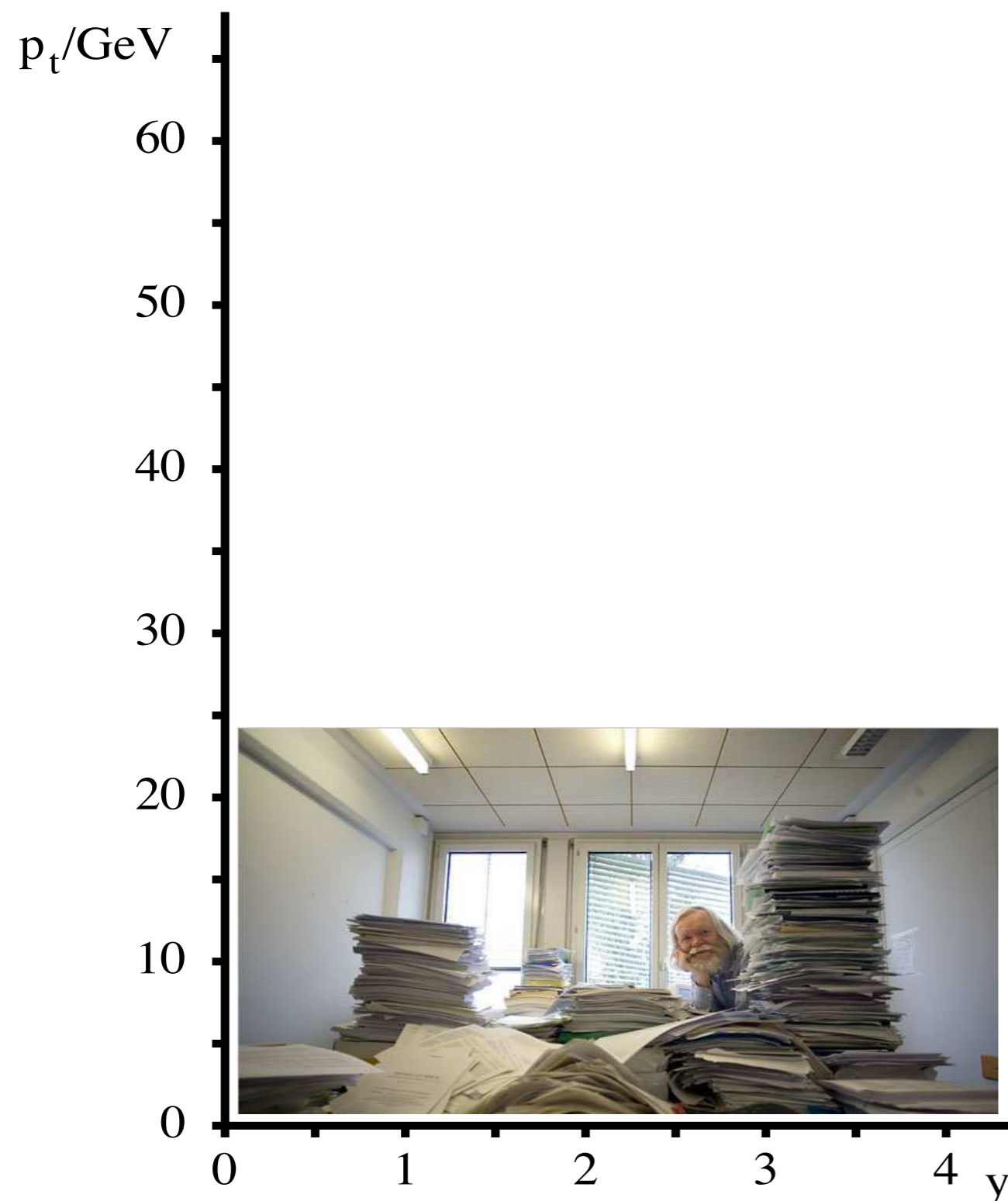
$$d_{ij} = \min(k_{ti}^2, k_{tj}^2) \frac{\Delta R_{ij}^2}{R^2}, \quad d_{iB} = k_{ti}^2$$

- ▶ If d_{ij} recombine
- ▶ if d_{iB} , i is a jet

Example clustering with k_t algorithm, $R = 1.0$

ϕ assumed 0 for all towers





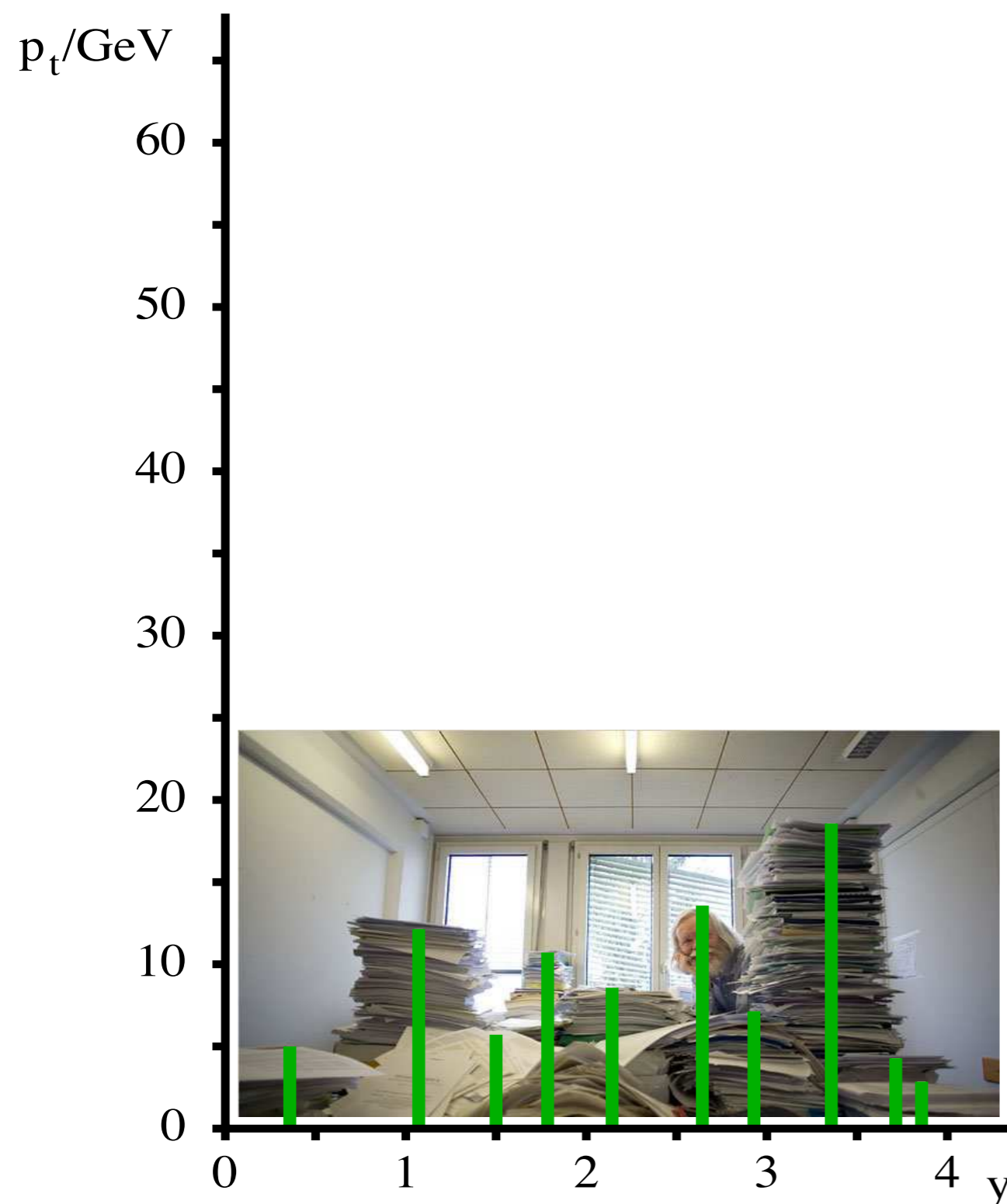
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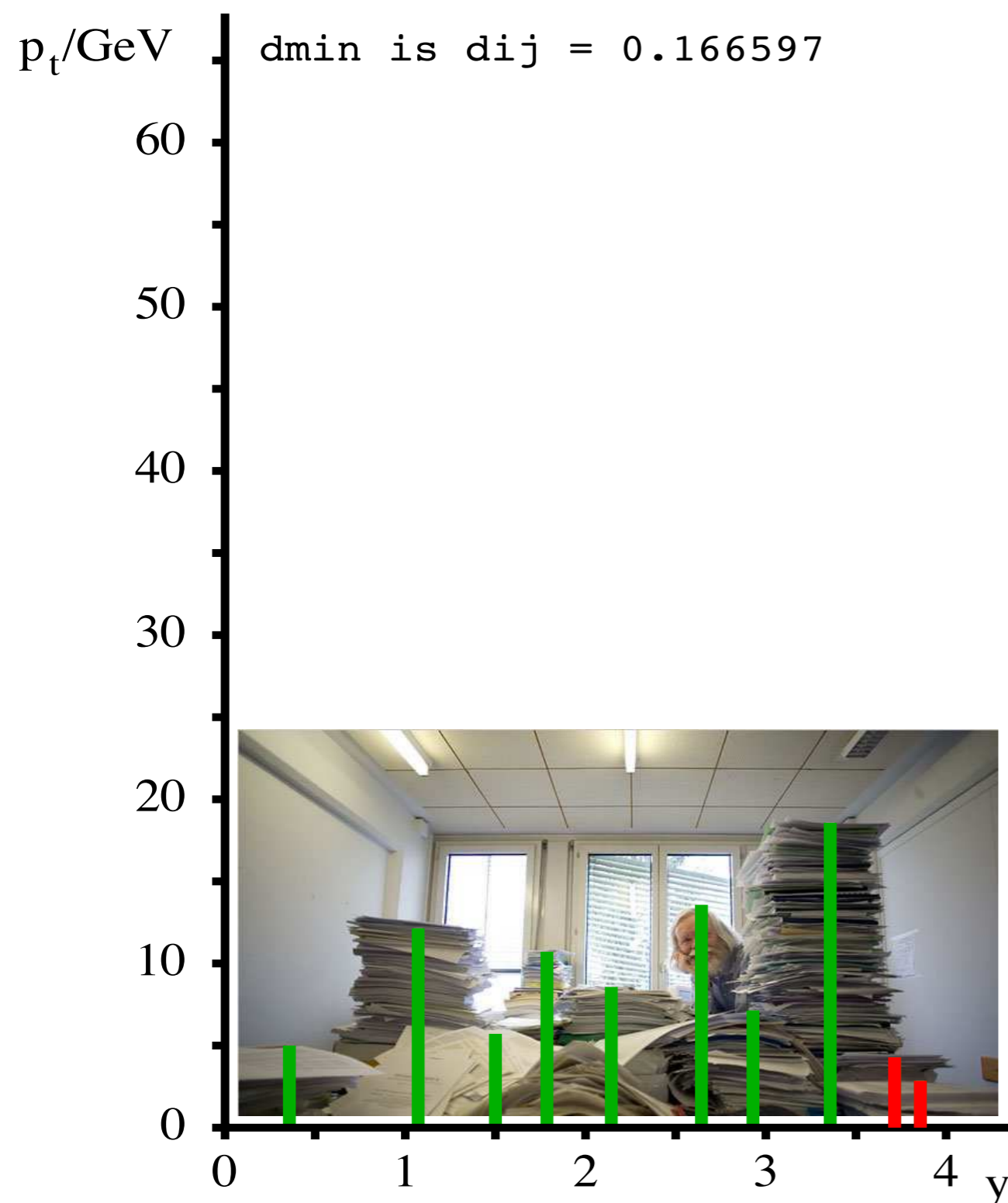
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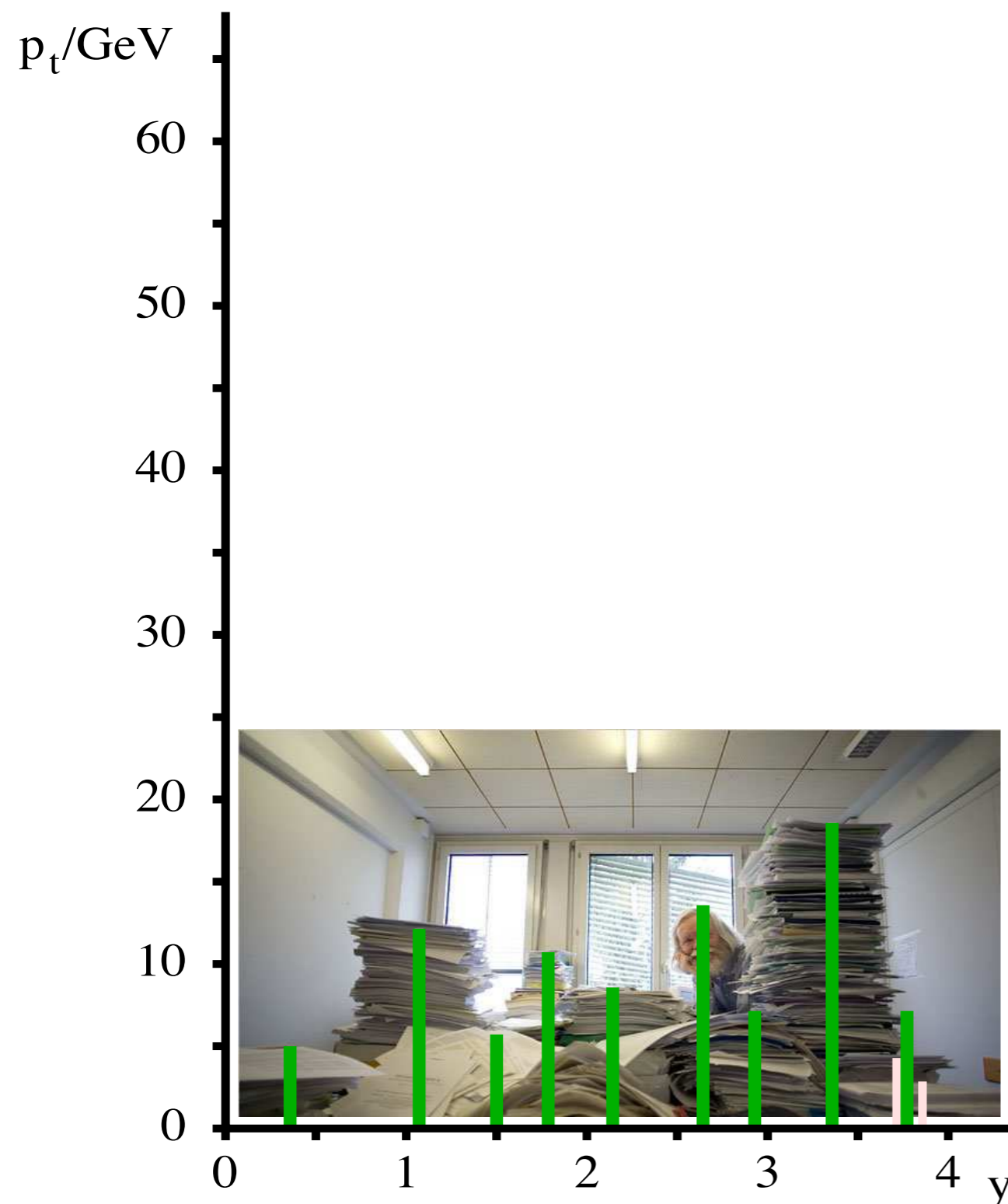
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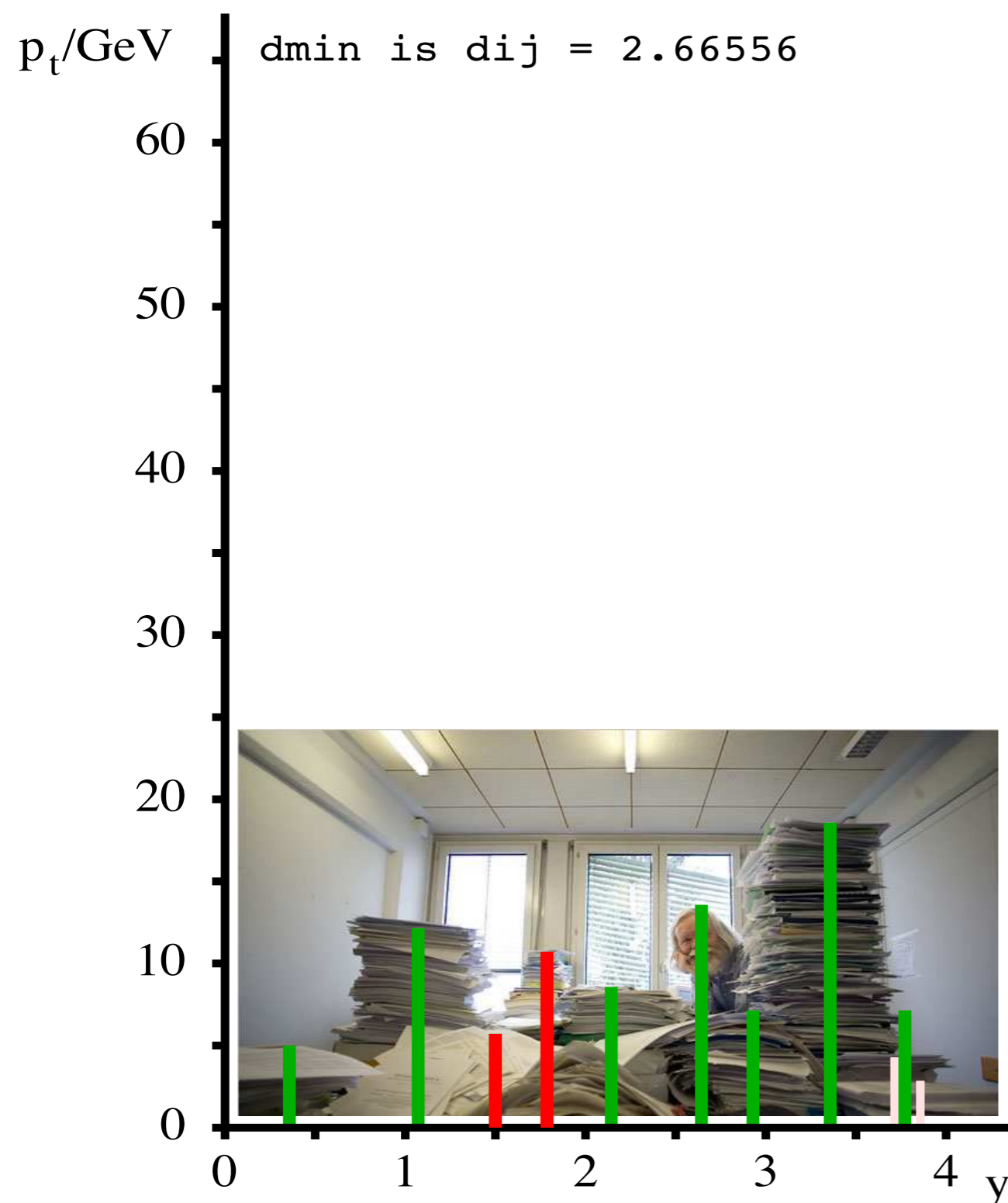
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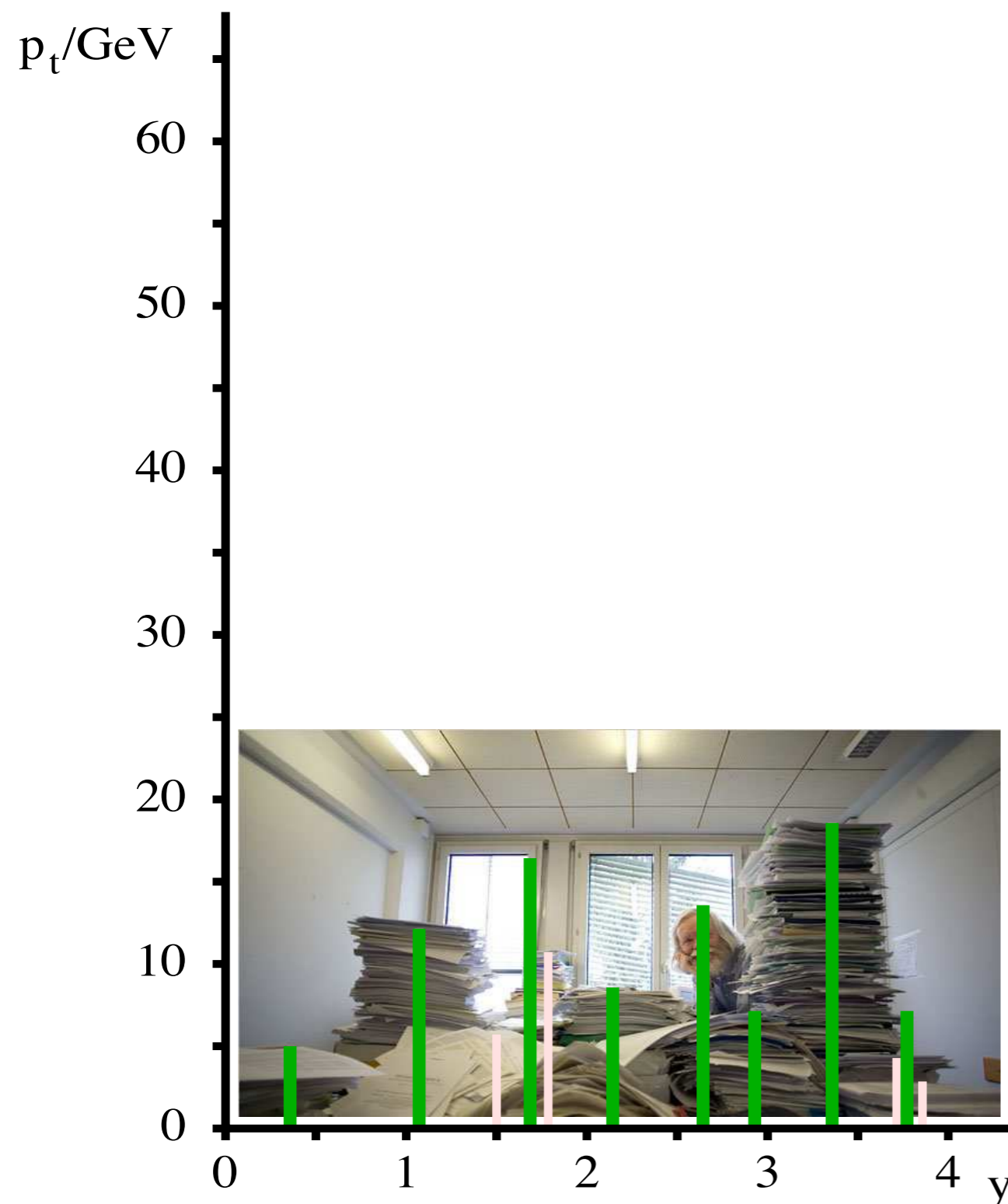
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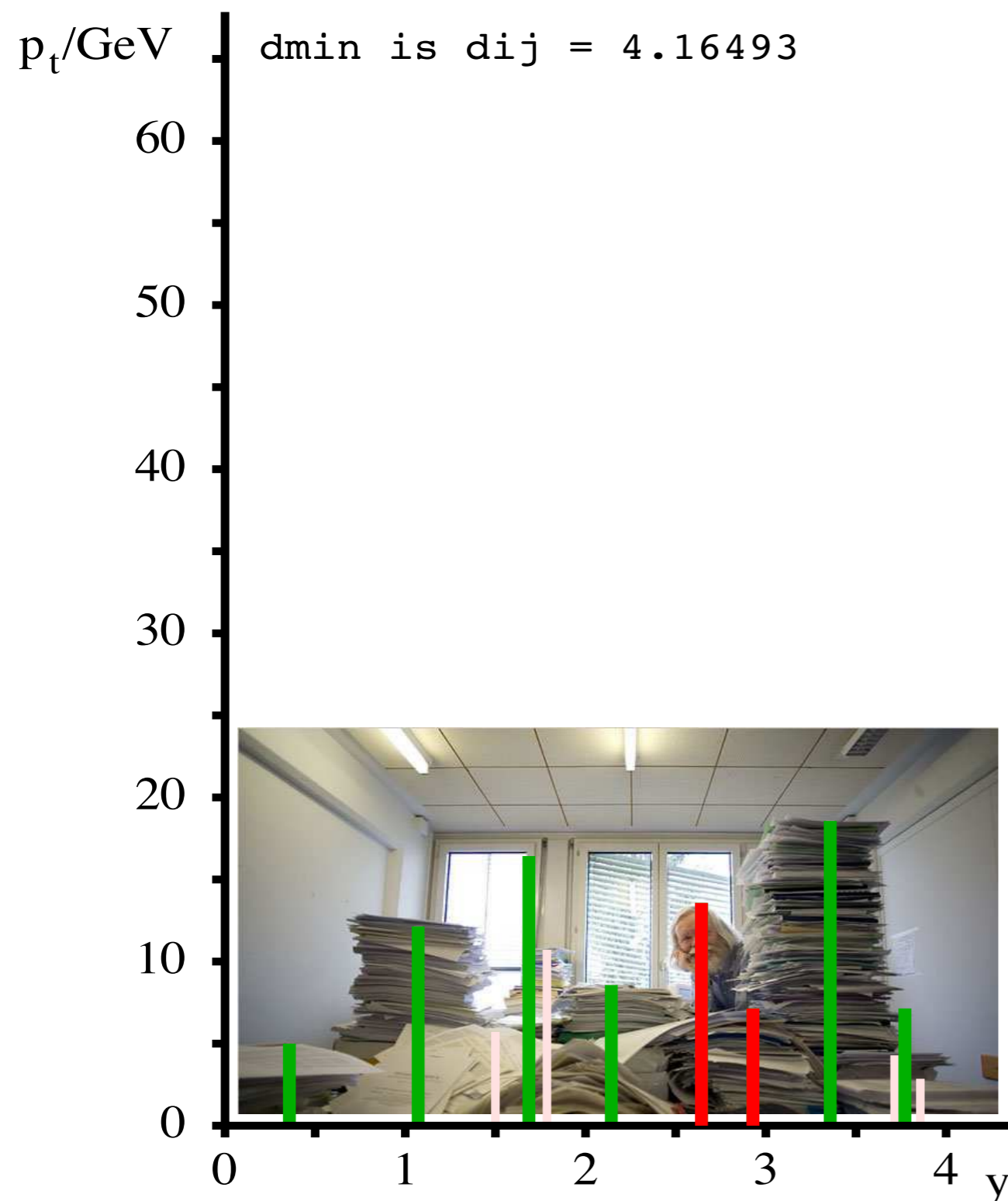
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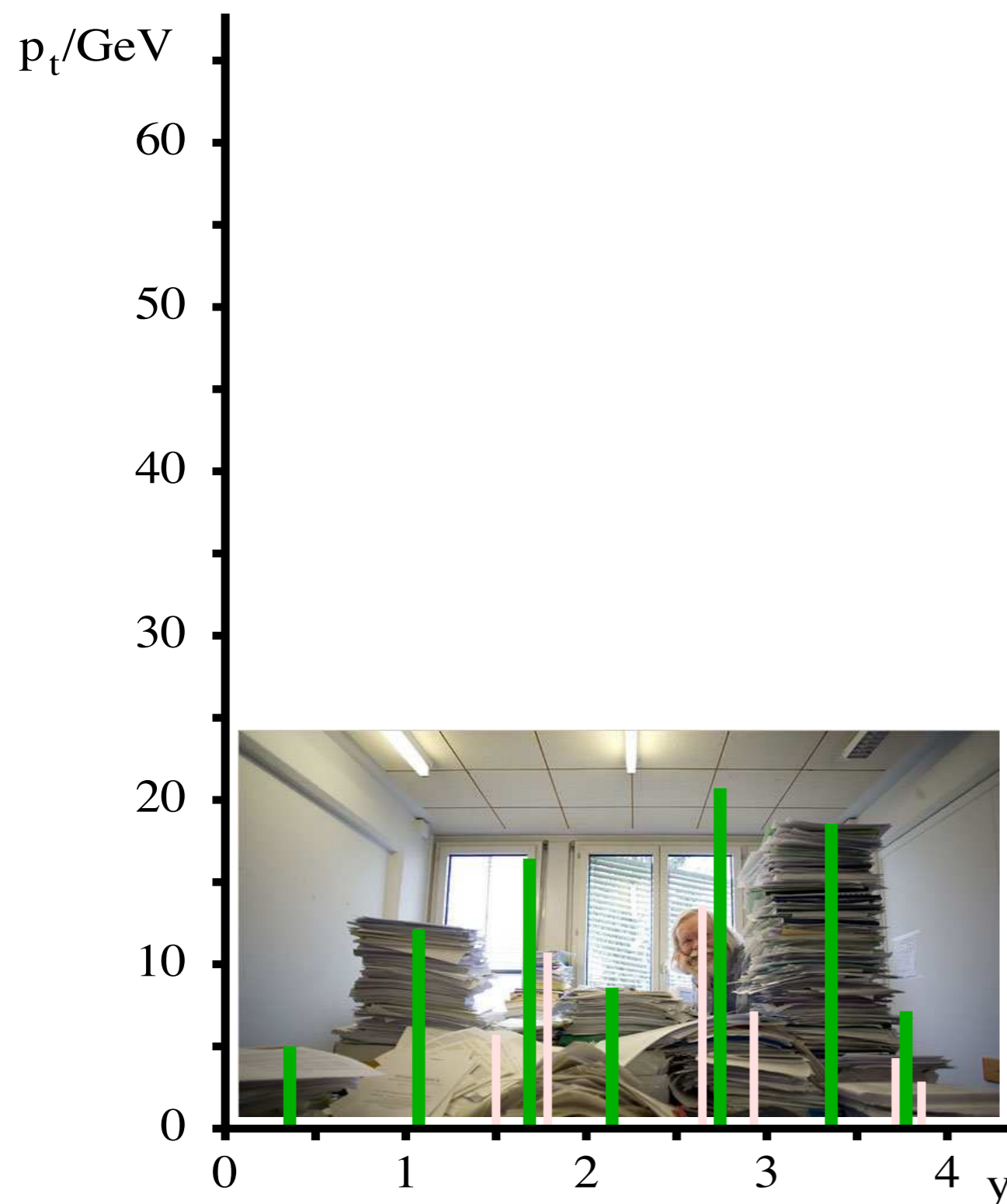
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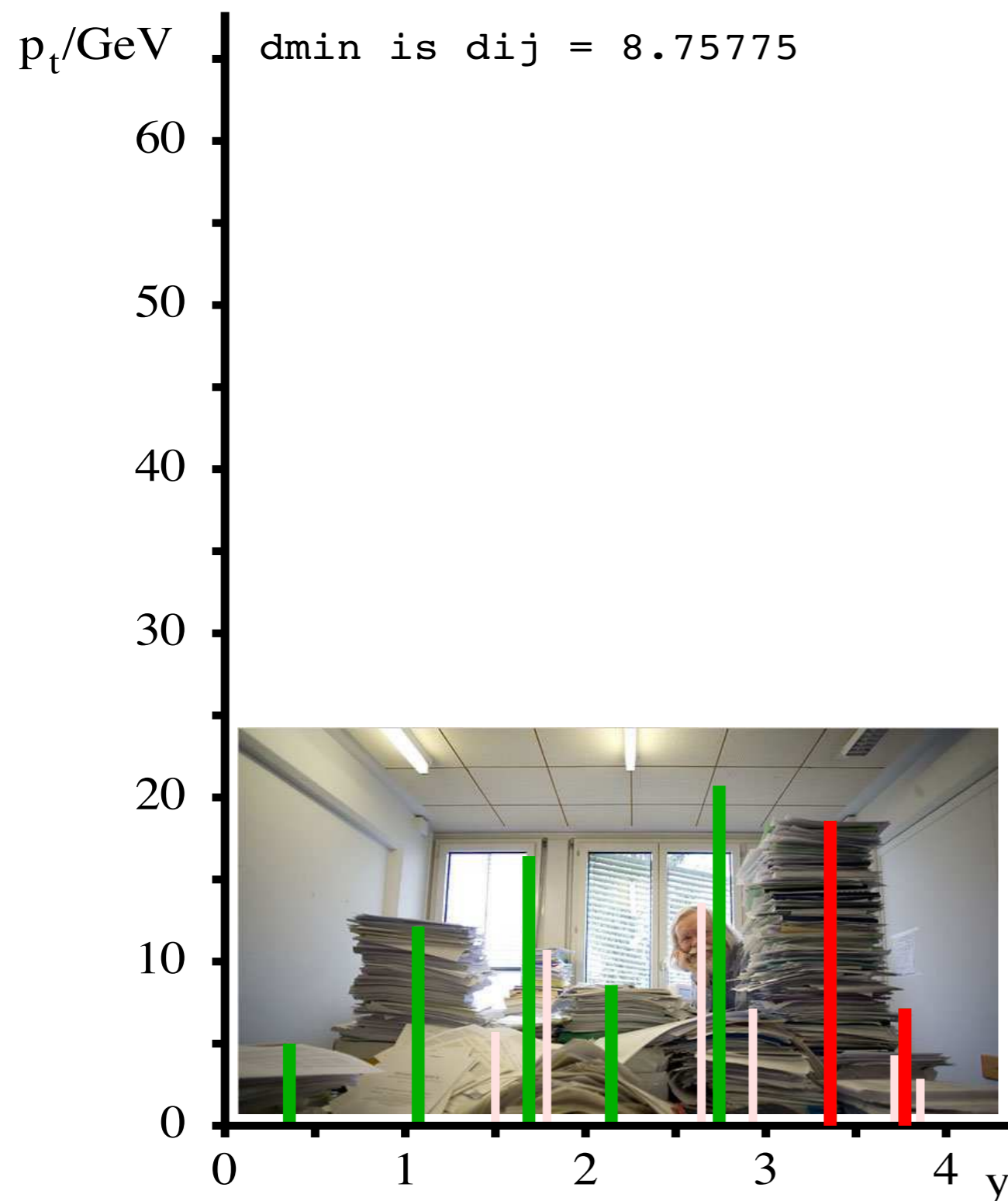
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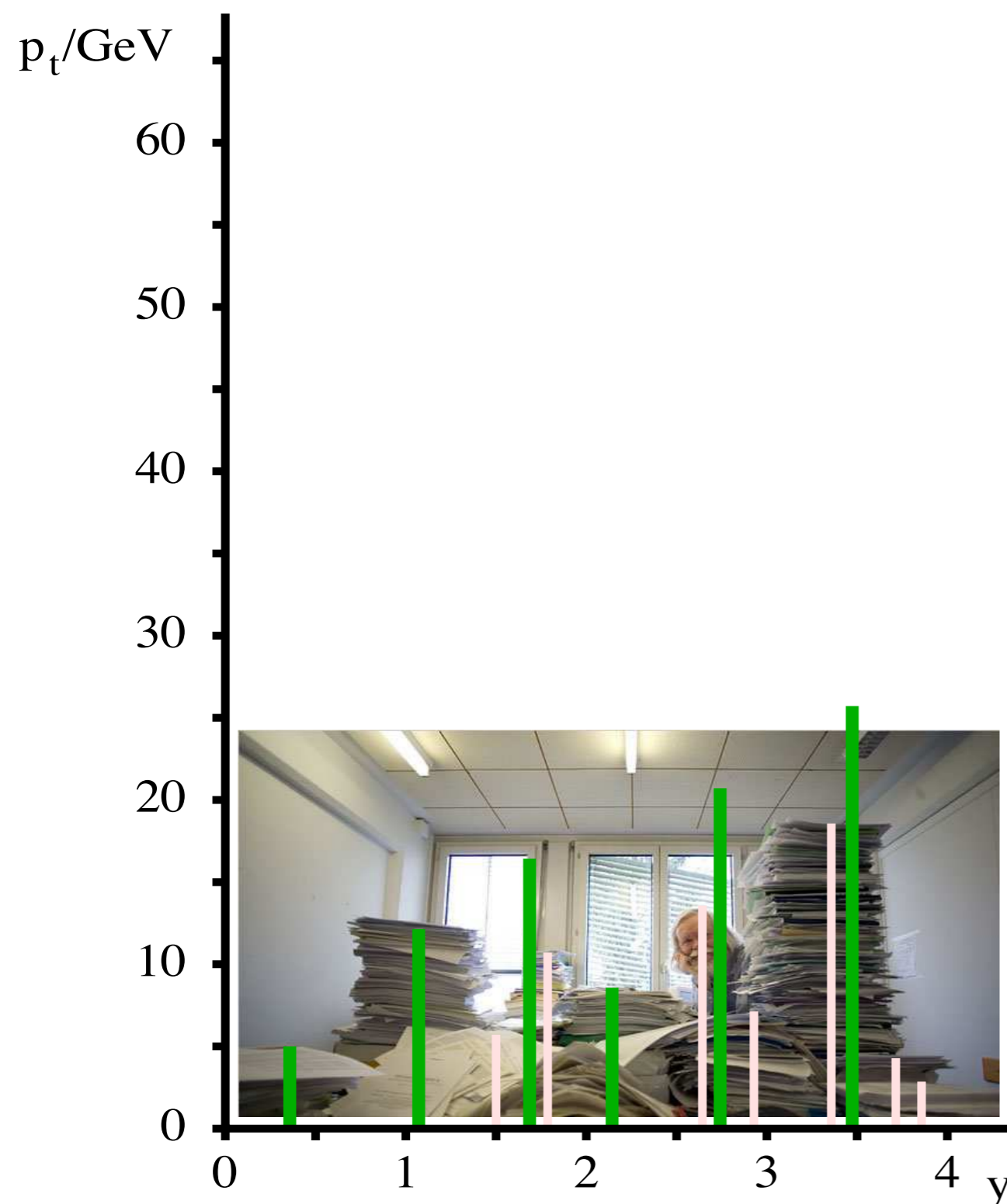
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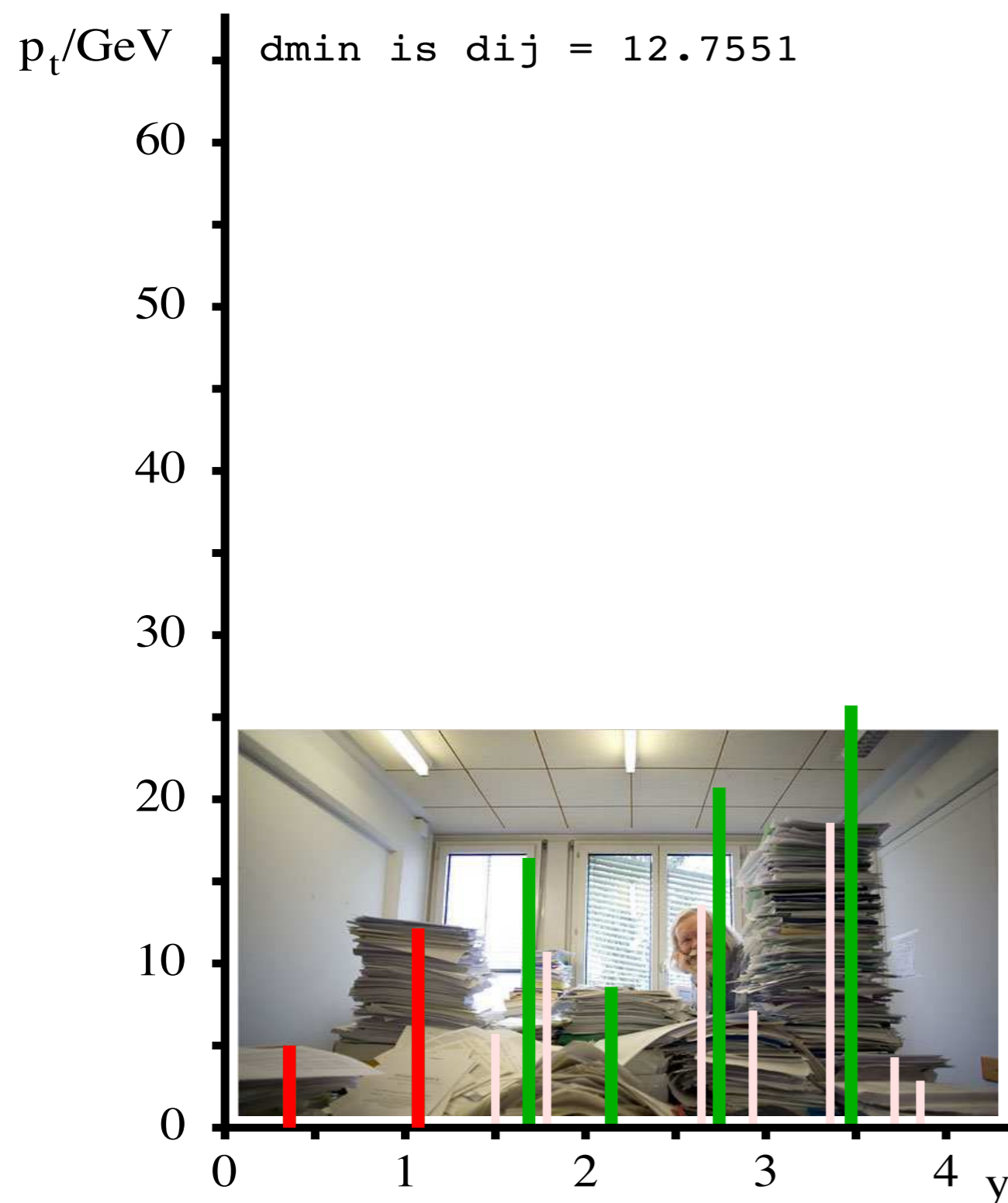
k_t alg.: Find smallest of

$$d_{ij} = \min(k_{ti}^2, k_{tj}^2) \frac{\Delta R_{ij}^2}{R^2}, \quad d_{iB} = k_{ti}^2$$

- ▶ If d_{ij} recombine
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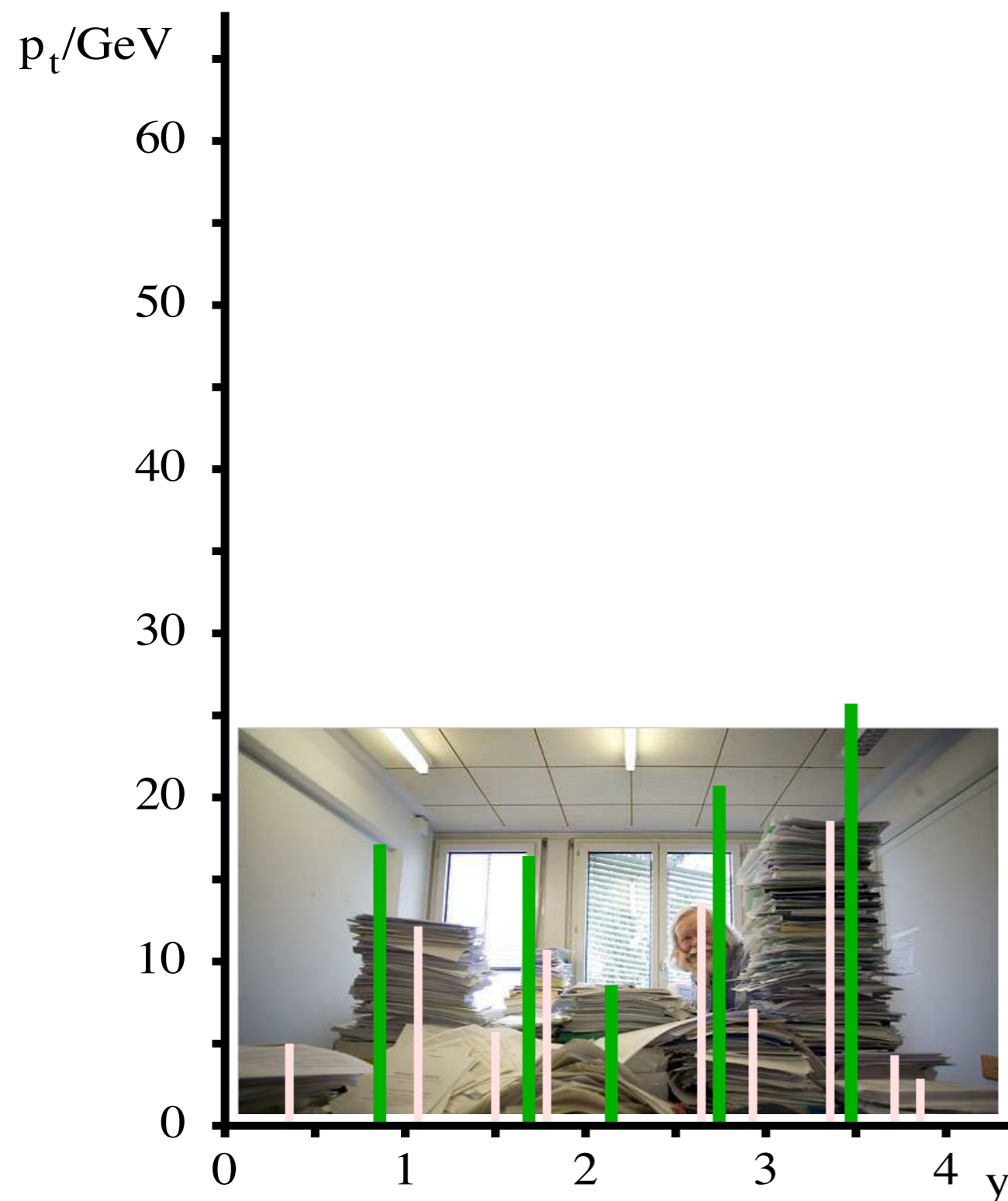
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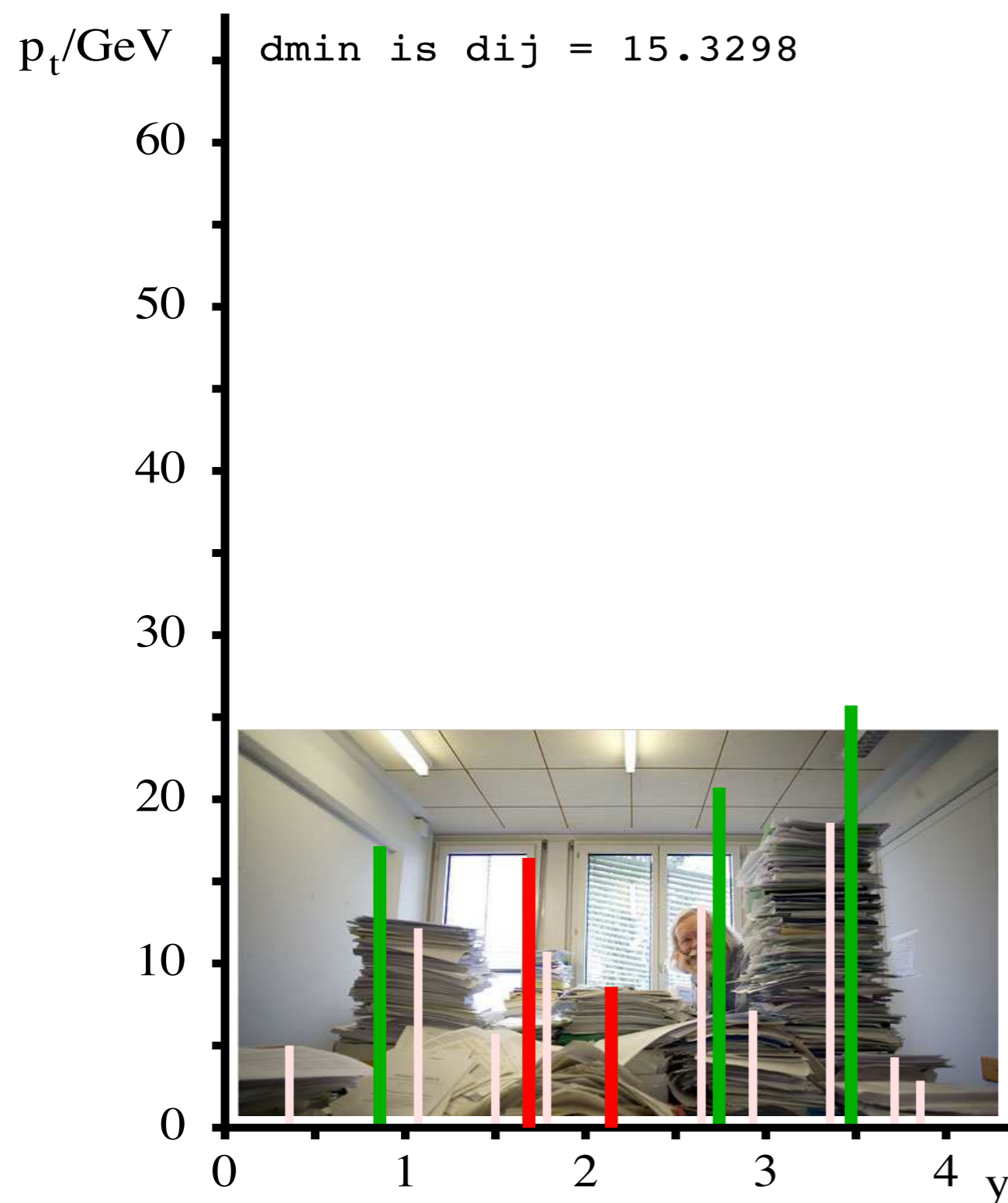
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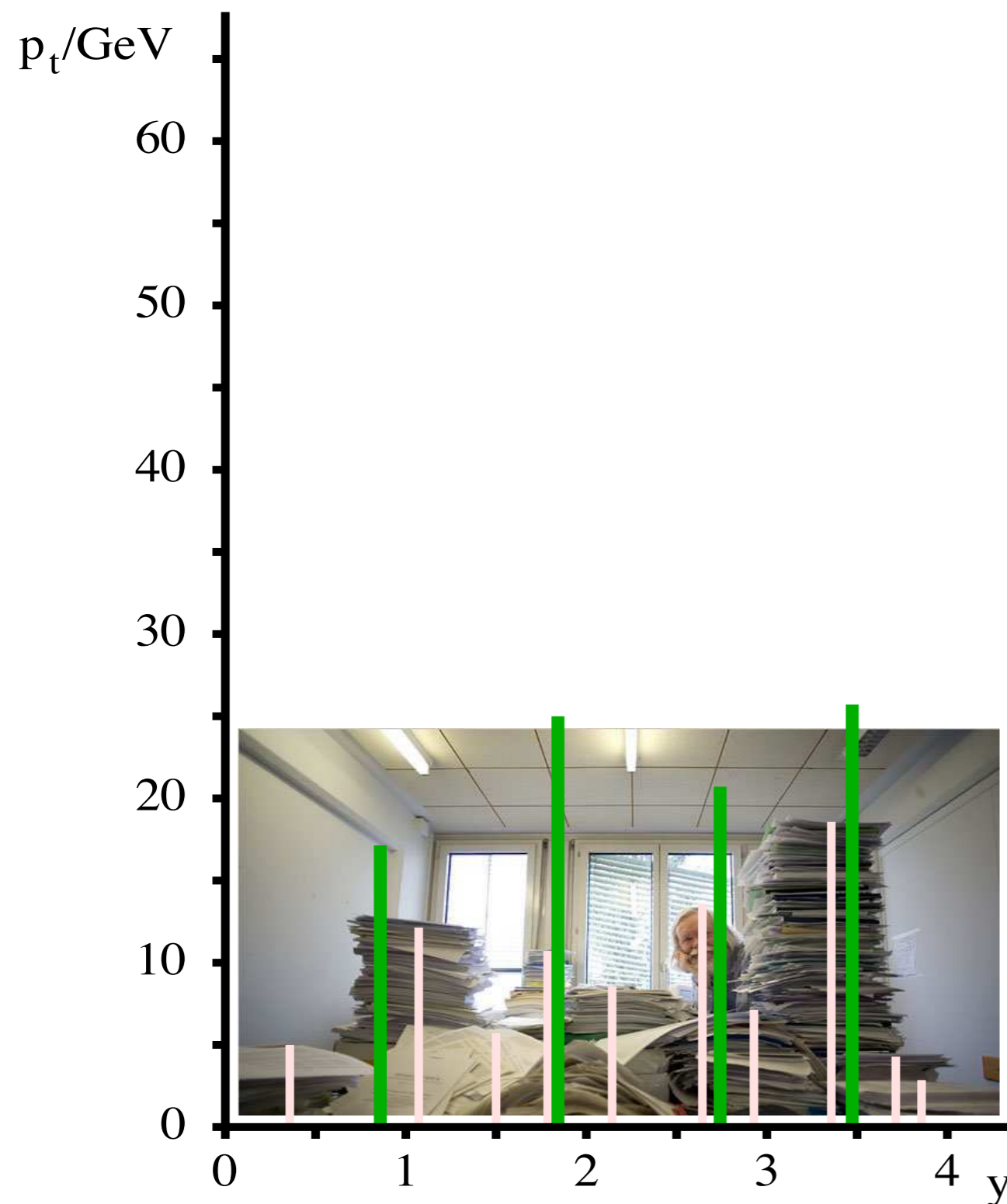
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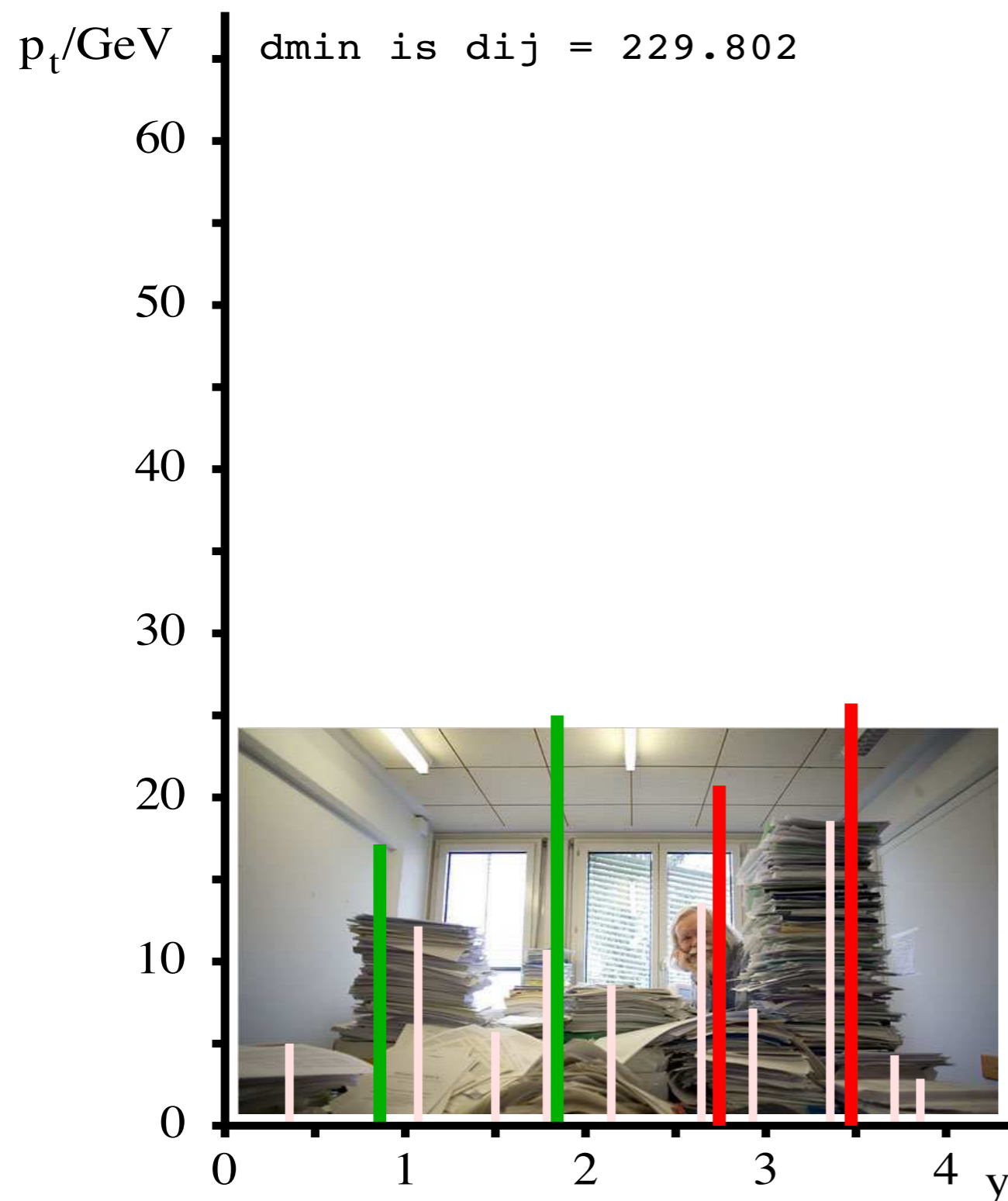
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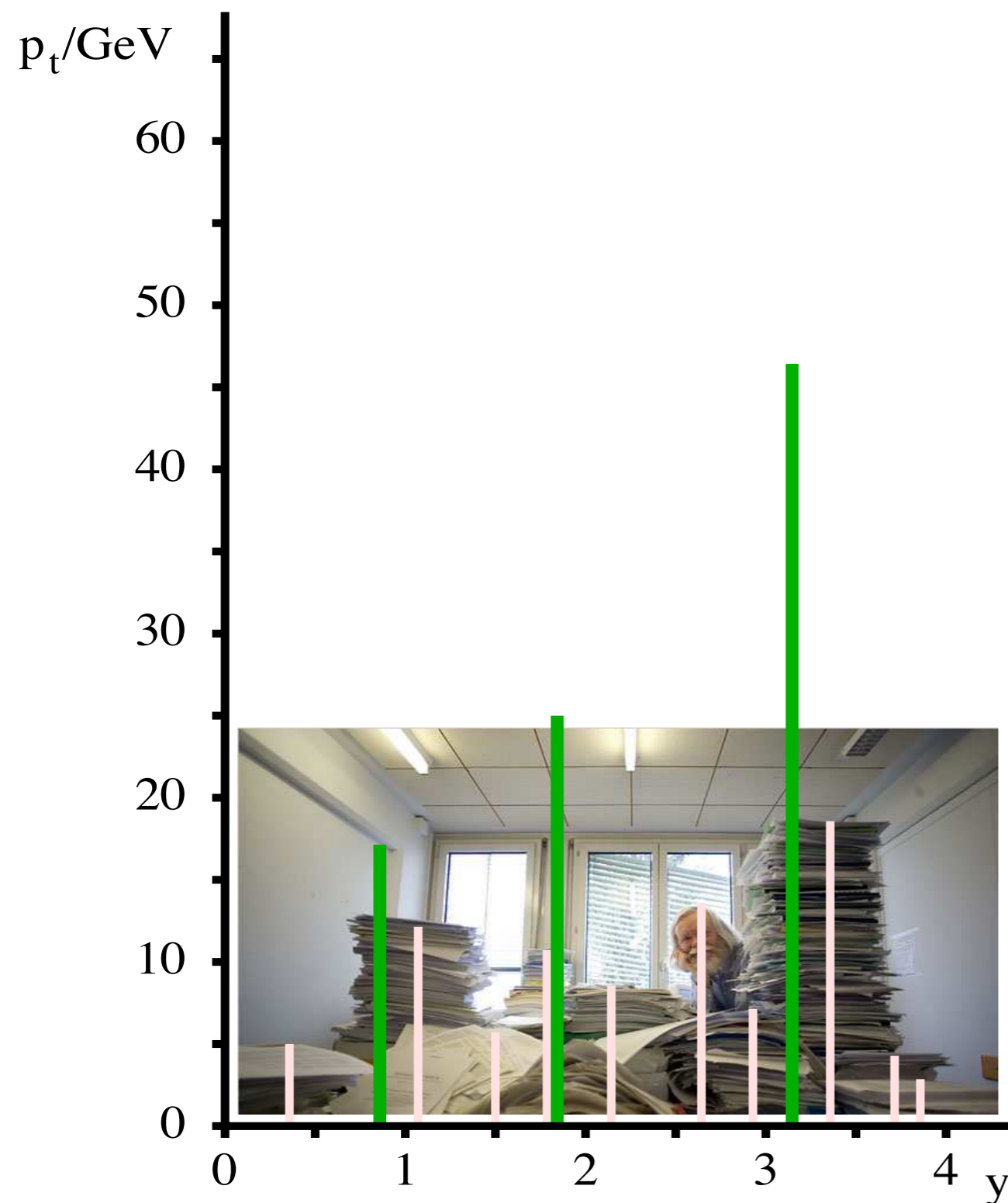
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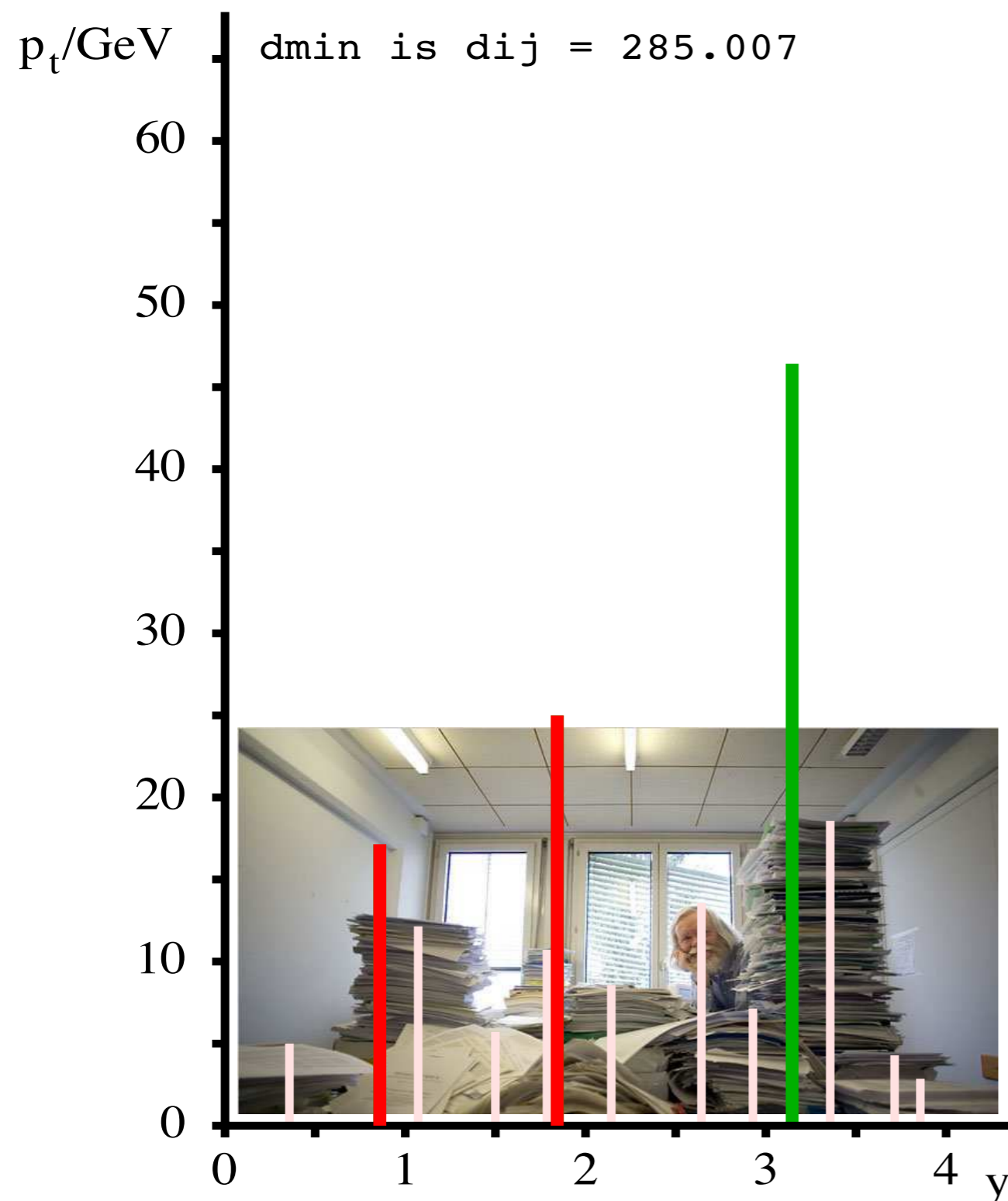
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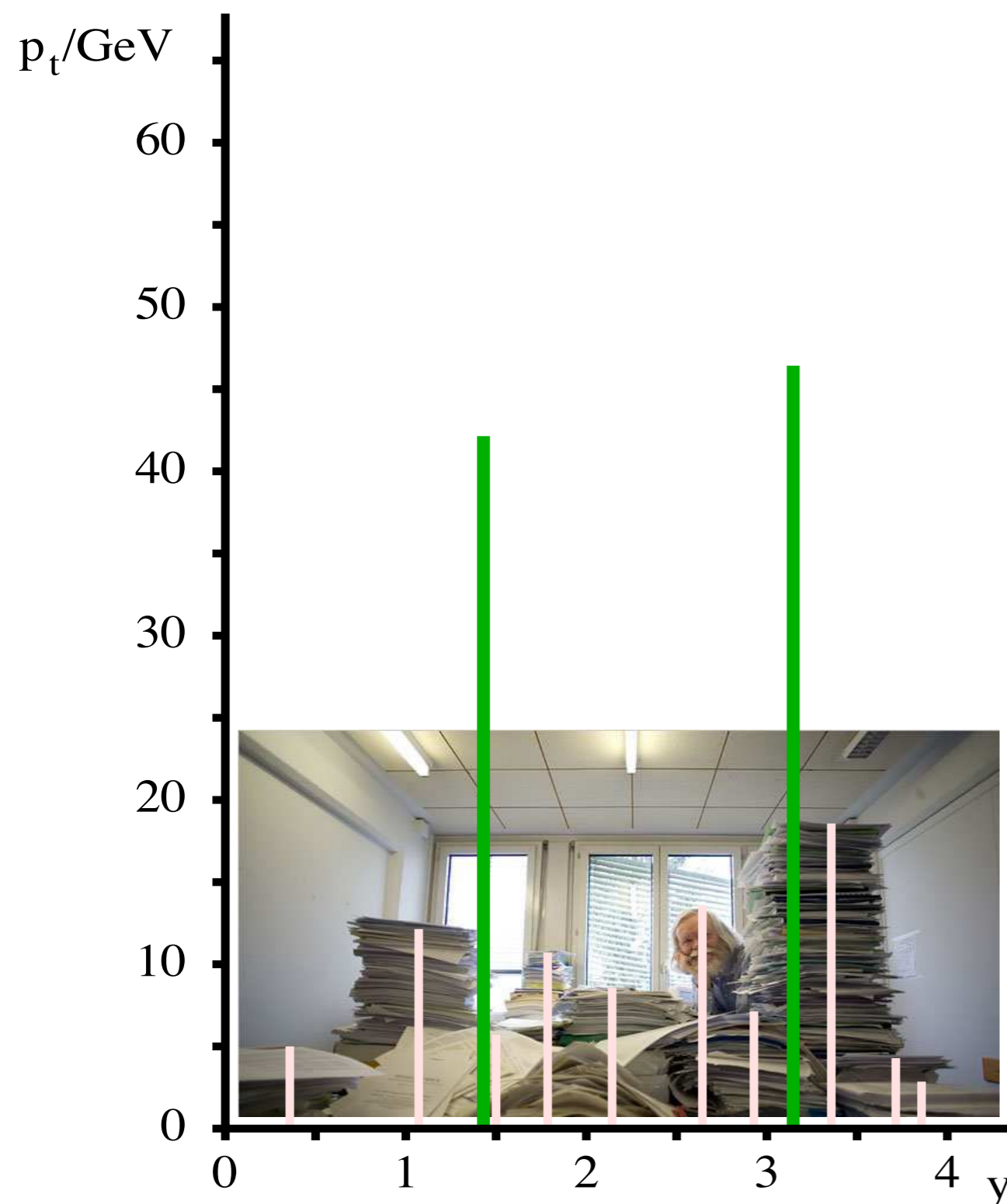
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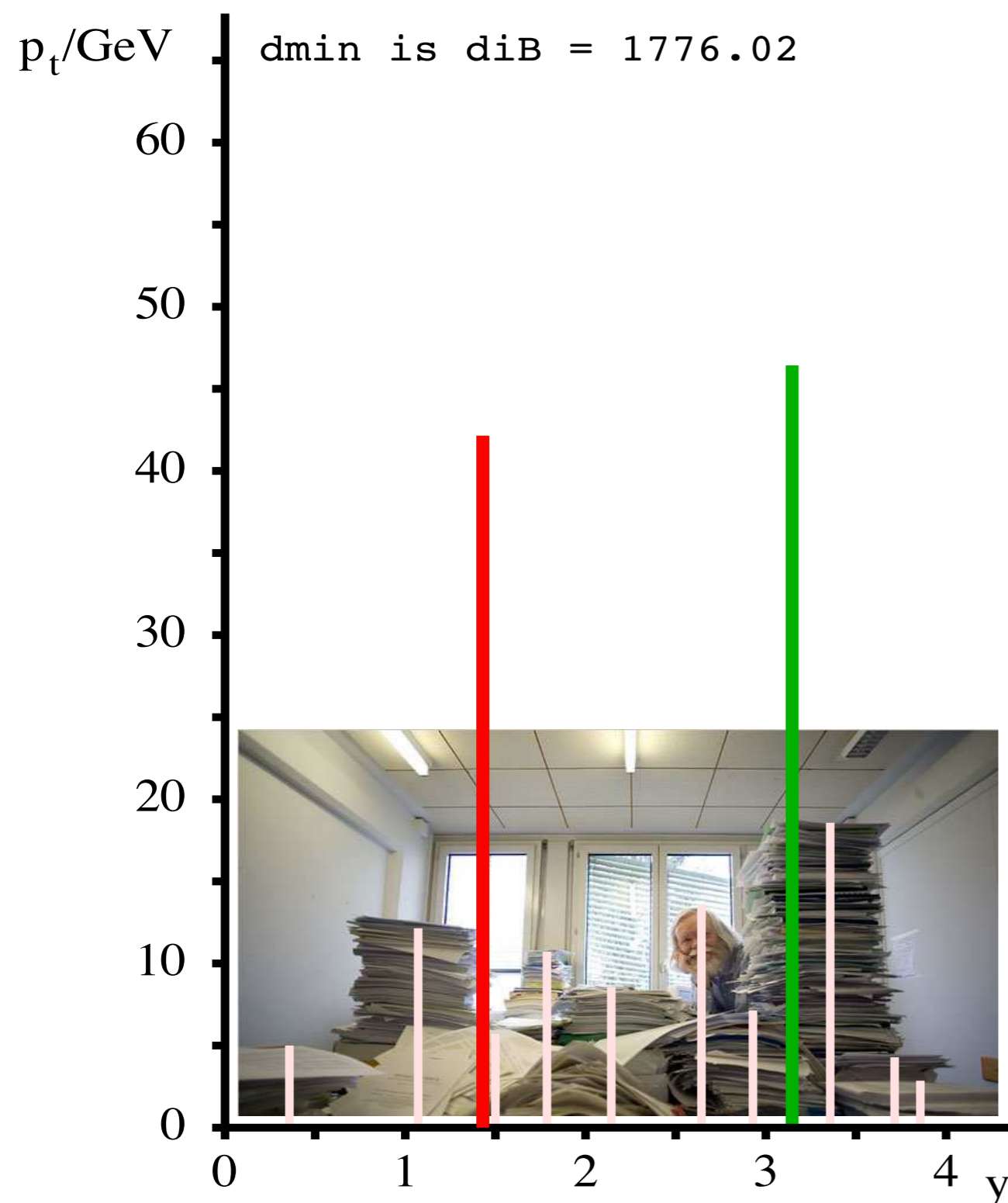
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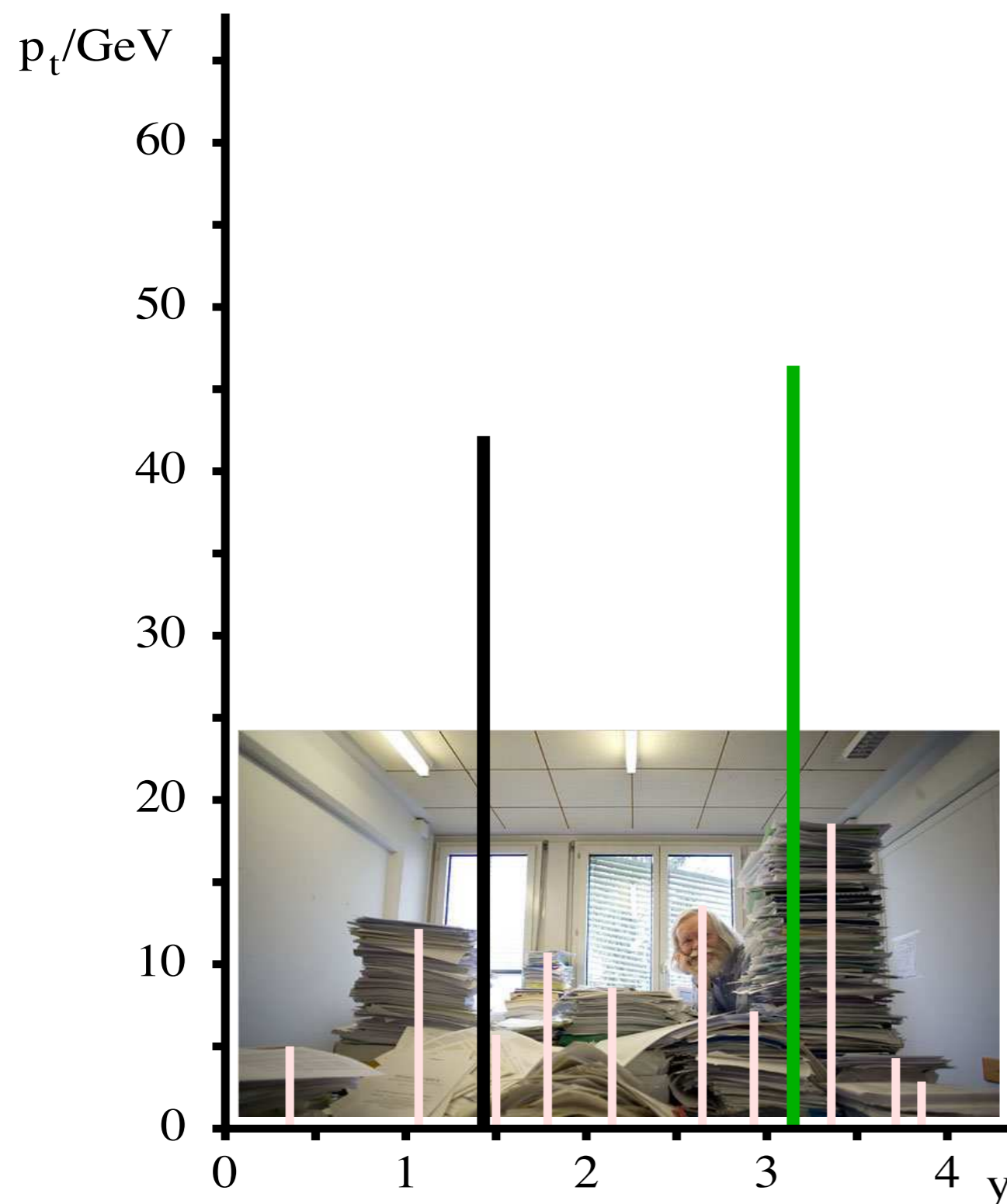
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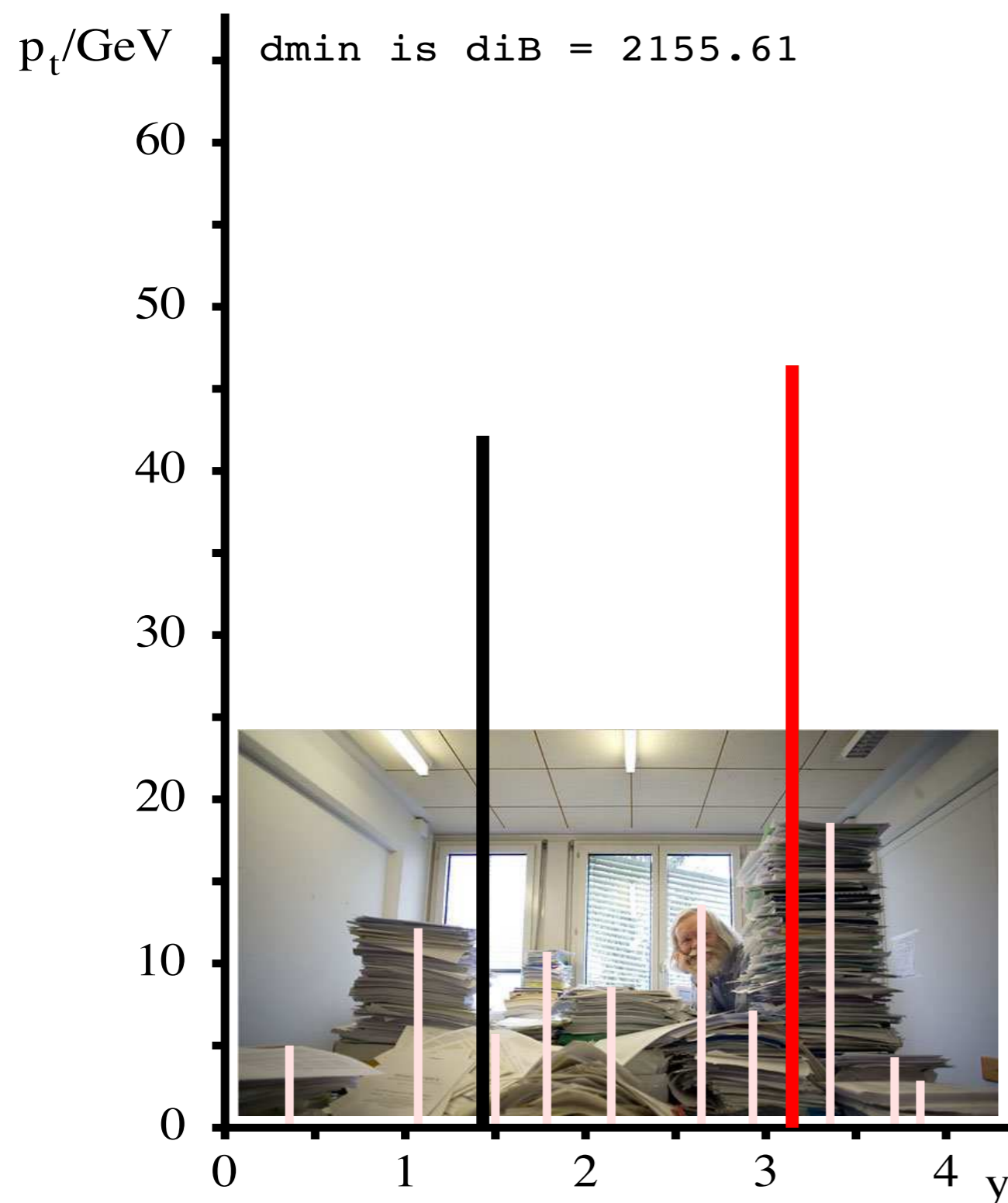
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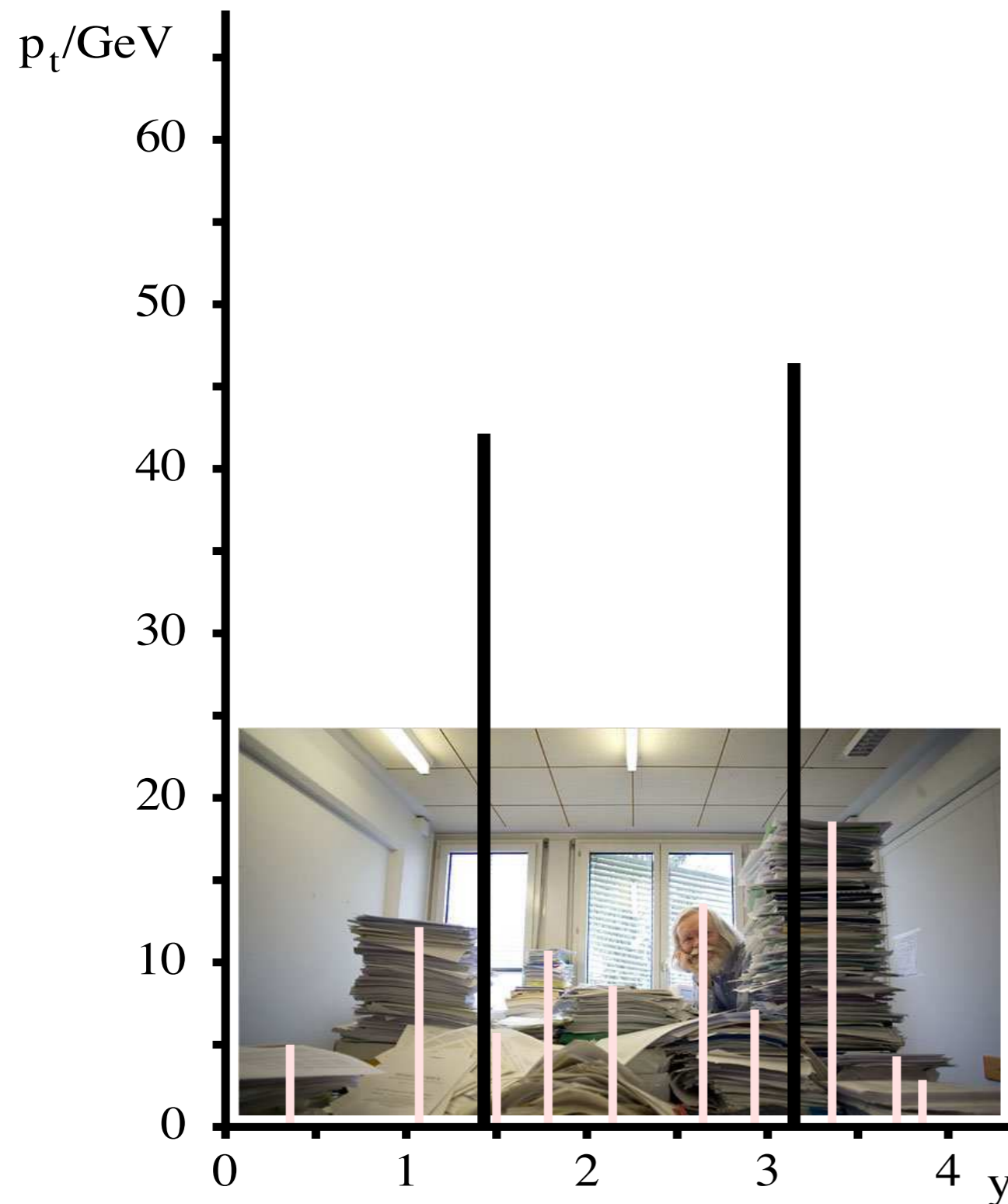
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Sequential recombination variants

Cambridge/Aachen: the simplest of hadron-collider algorithms

- Recombine pair of objects closest in ΔR_{ij}
- Repeat until all $\Delta R_{ij} > R$ — remaining objects are jets

Dokshitzer, Leder, Moretti, Webber '97 (Cambridge): more involved e^+e^- form

Wobisch & Wengler '99 (Aachen): simple inclusive hadron-collider form

One still applies a $p_{t,\min}$ cut to the jets, as for inclusive k_t

C/A privileges the collinear divergence of QCD;
it 'ignores' the soft one

Anti- k_t : formulated similarly to inclusive k_t , but with

$$d_{ij} = \frac{1}{\max(p_{ti}^2, p_{tj}^2)} \frac{\Delta R_{ij}^2}{R^2}, \quad d_{iB} = \frac{1}{p_{ti}^2}$$

Cacciari, GPS & Soyez '08 [+Delsart unpublished]

Anti- k_t privileges the collinear divergence of QCD and disfavours clustering between pairs of soft particles

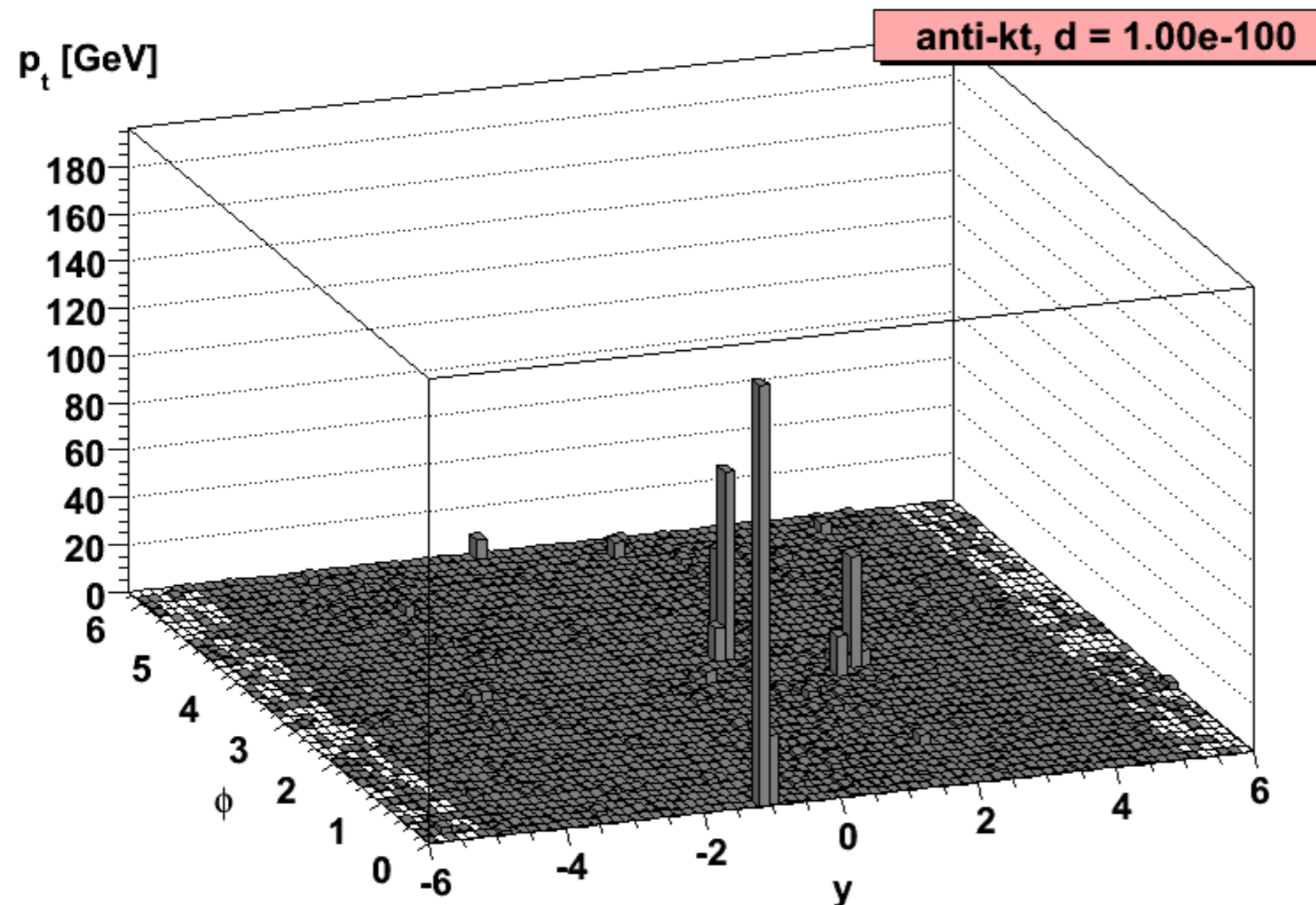
Most pairwise clusterings involve at least one hard particle

Clustering grows
around hard cores

$$d_{ij} = \frac{1}{\max(p_{ti}^2, p_{tj}^2)} \frac{\Delta R_{ij}^2}{R^2}, \quad d_{iB} = \frac{1}{p_{ti}^2}$$

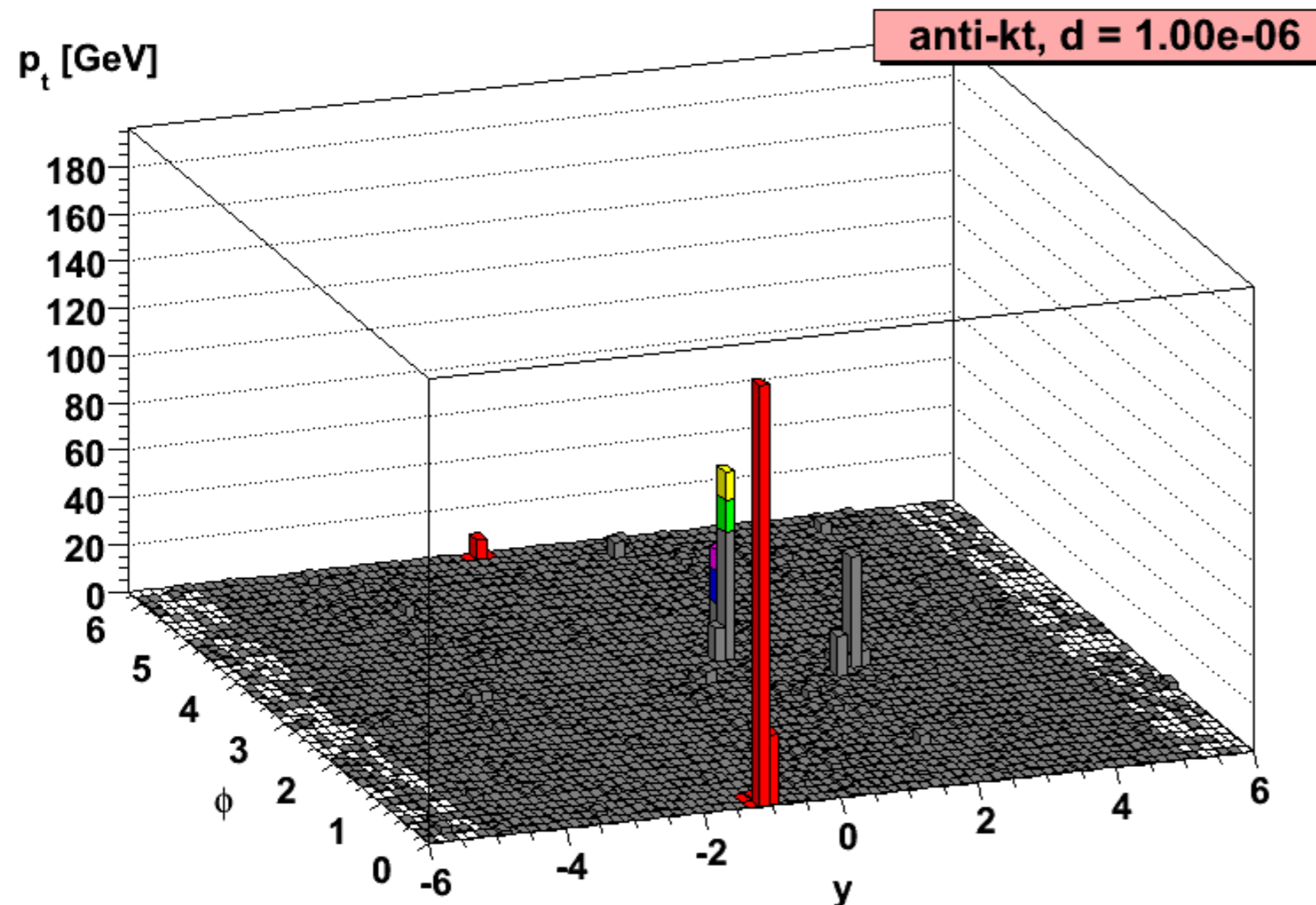
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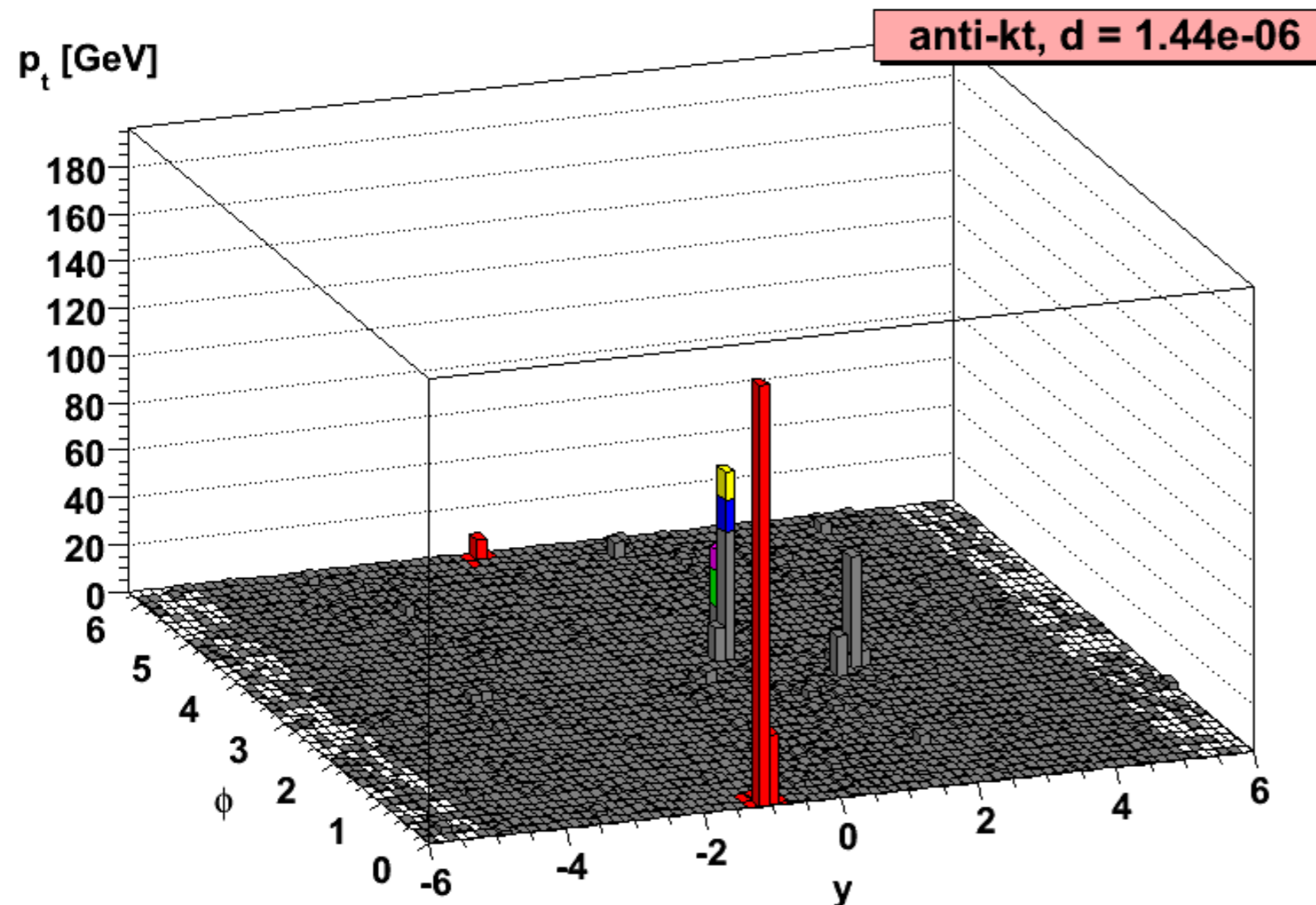
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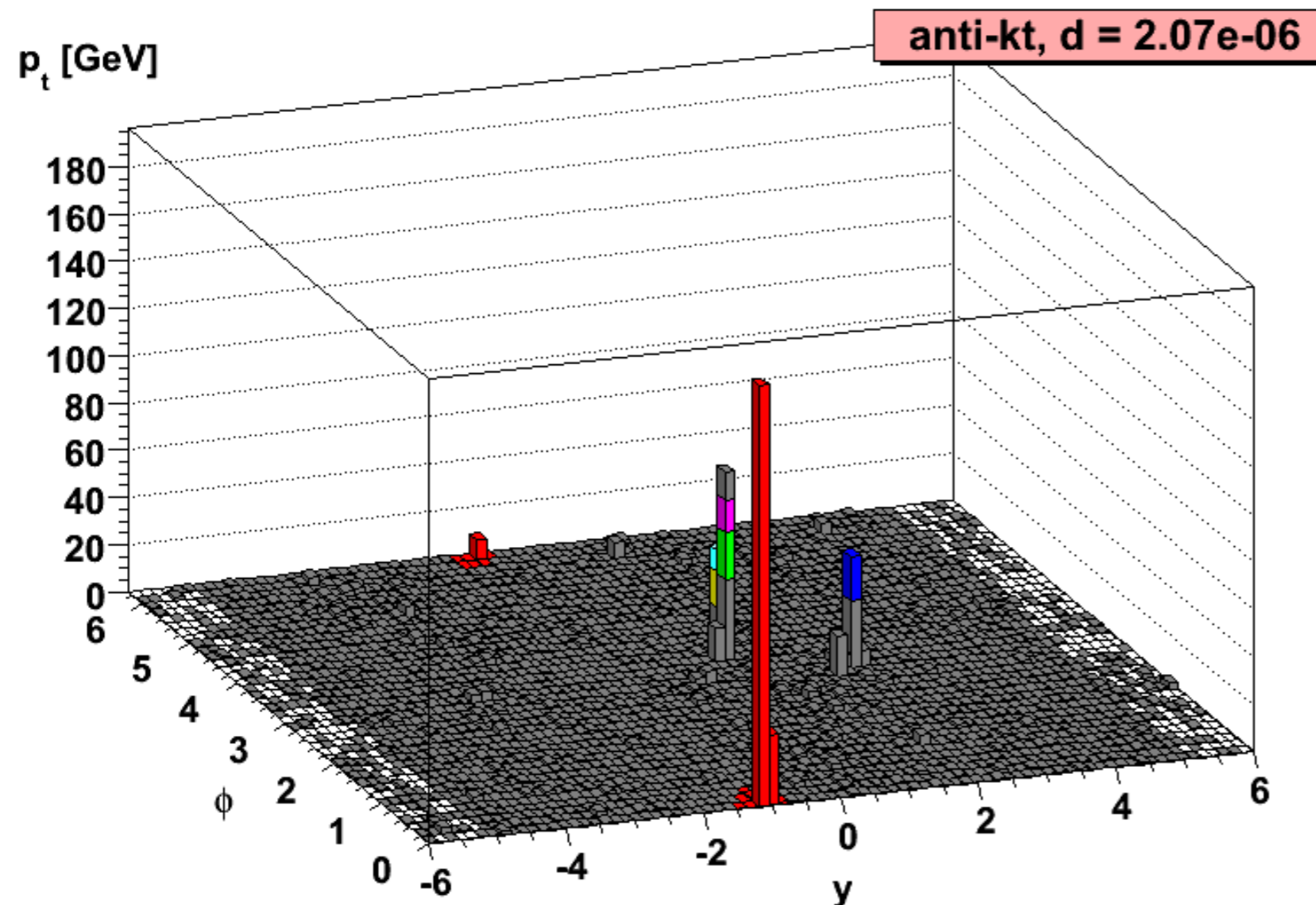
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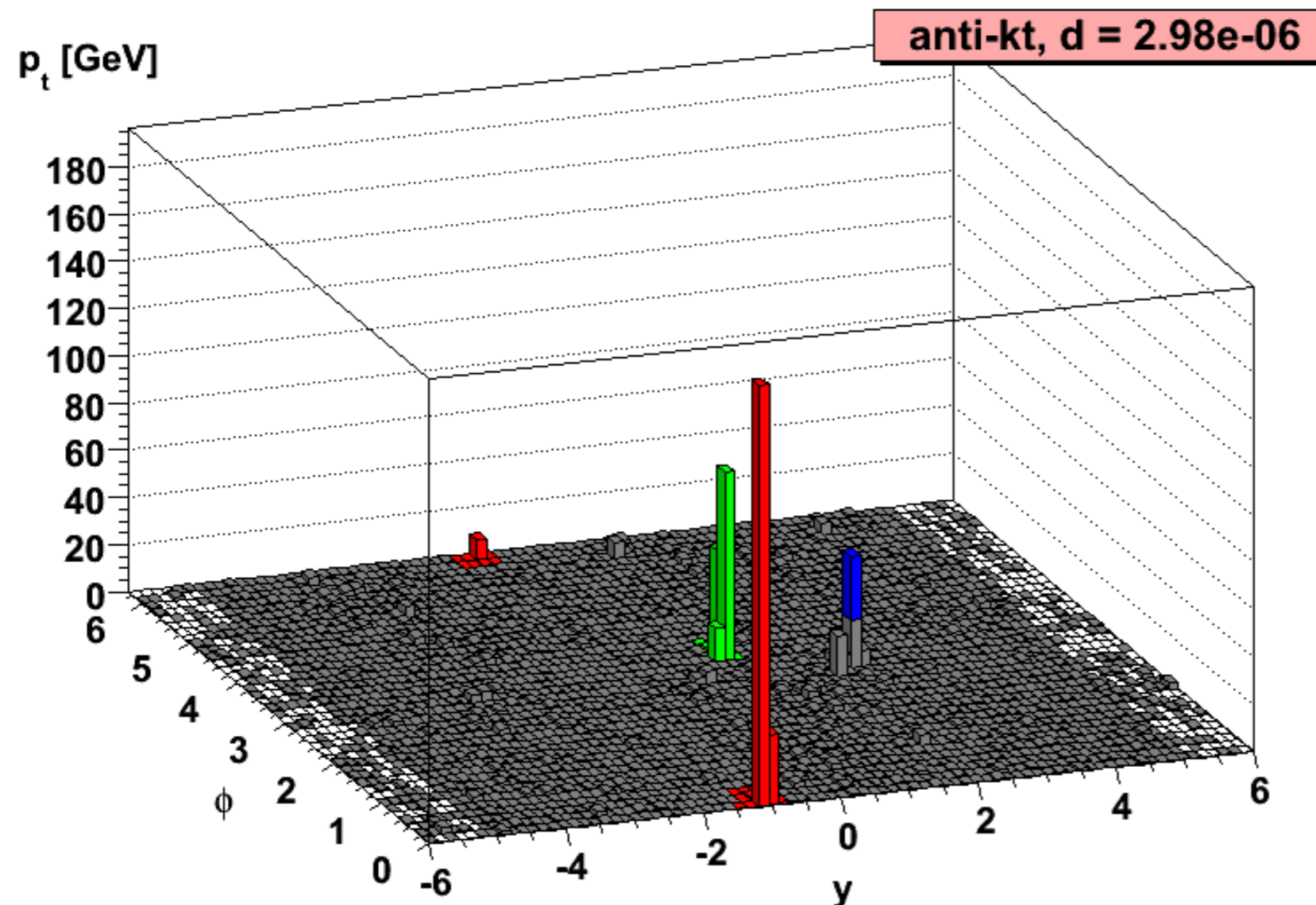
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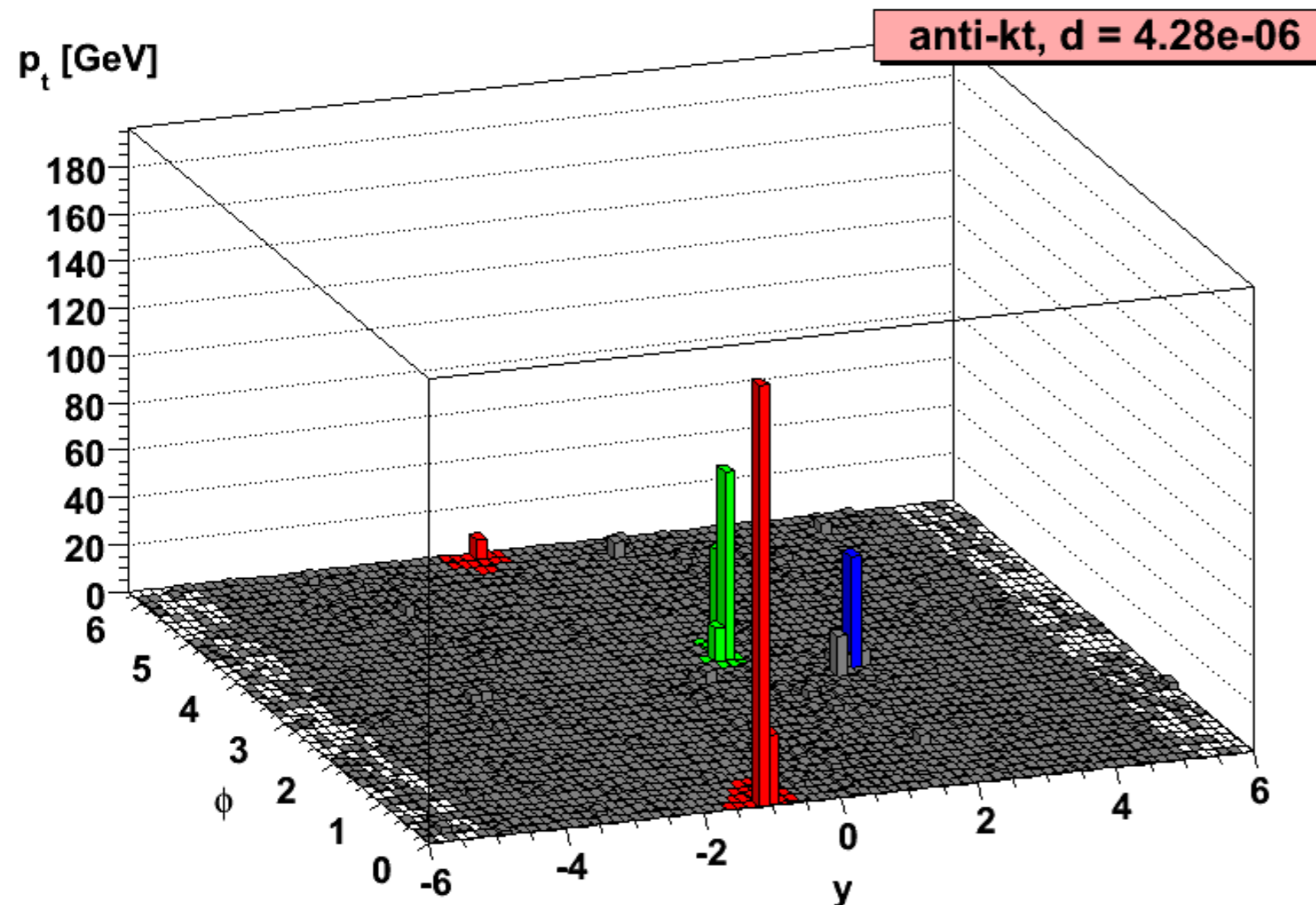
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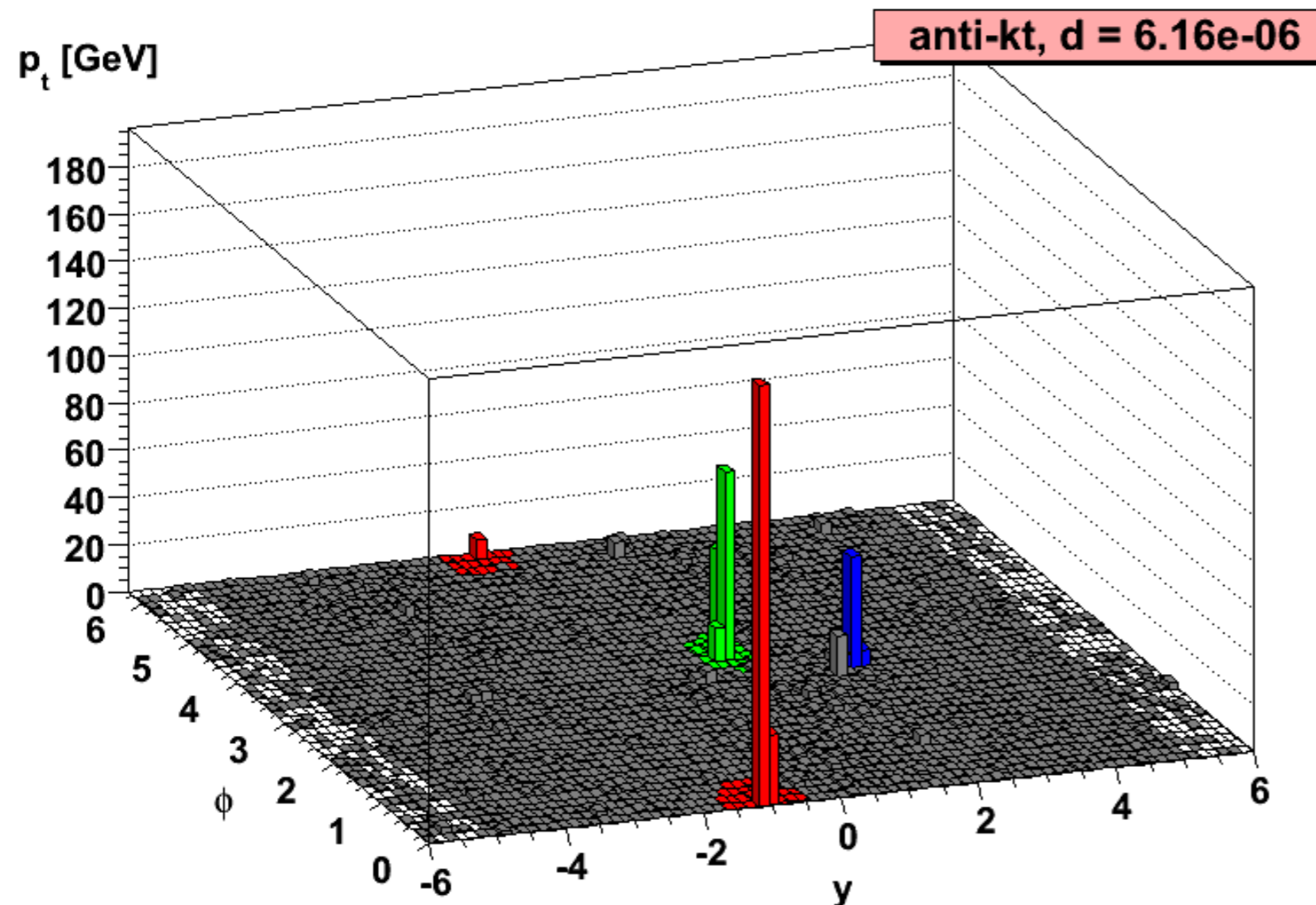
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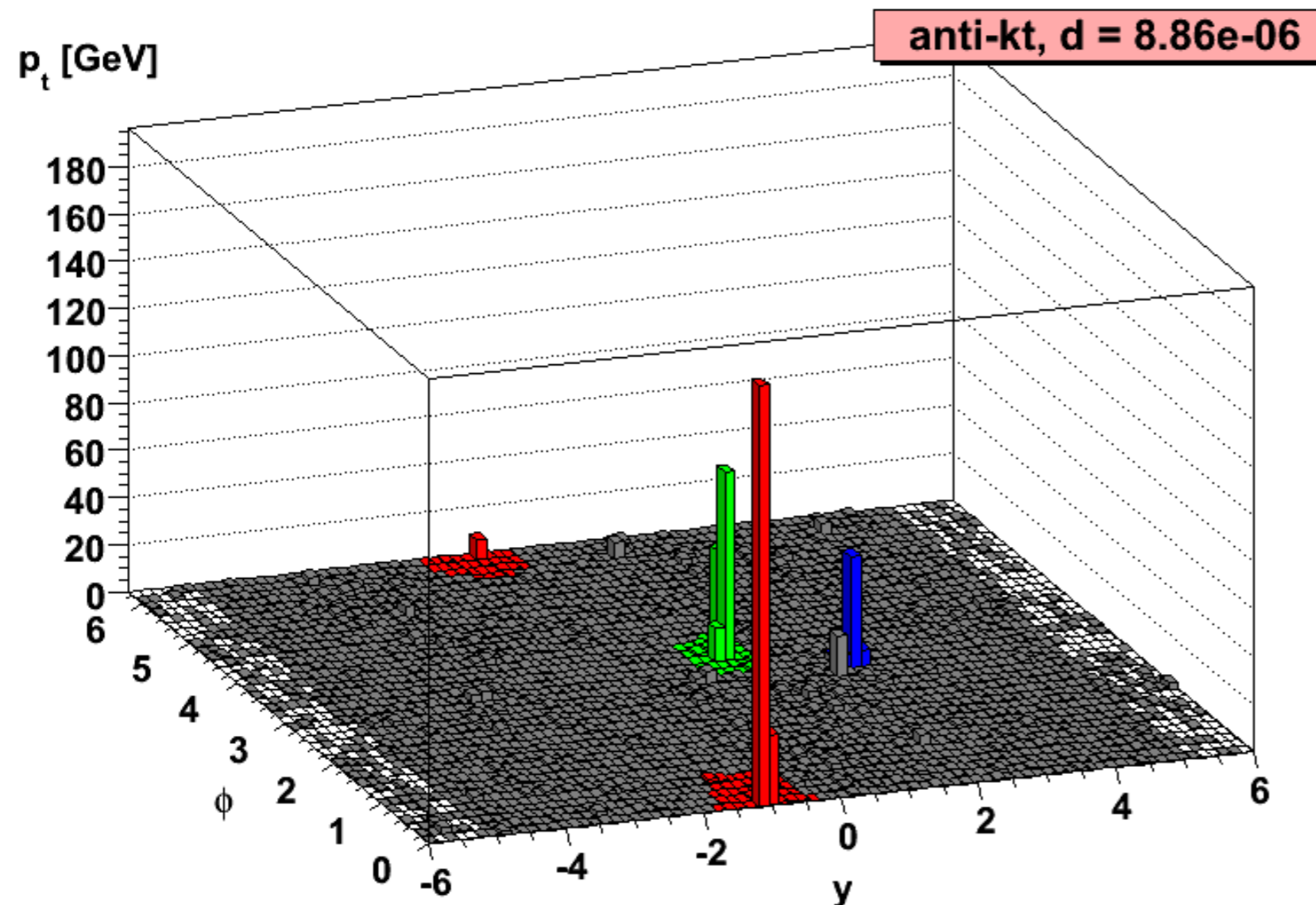
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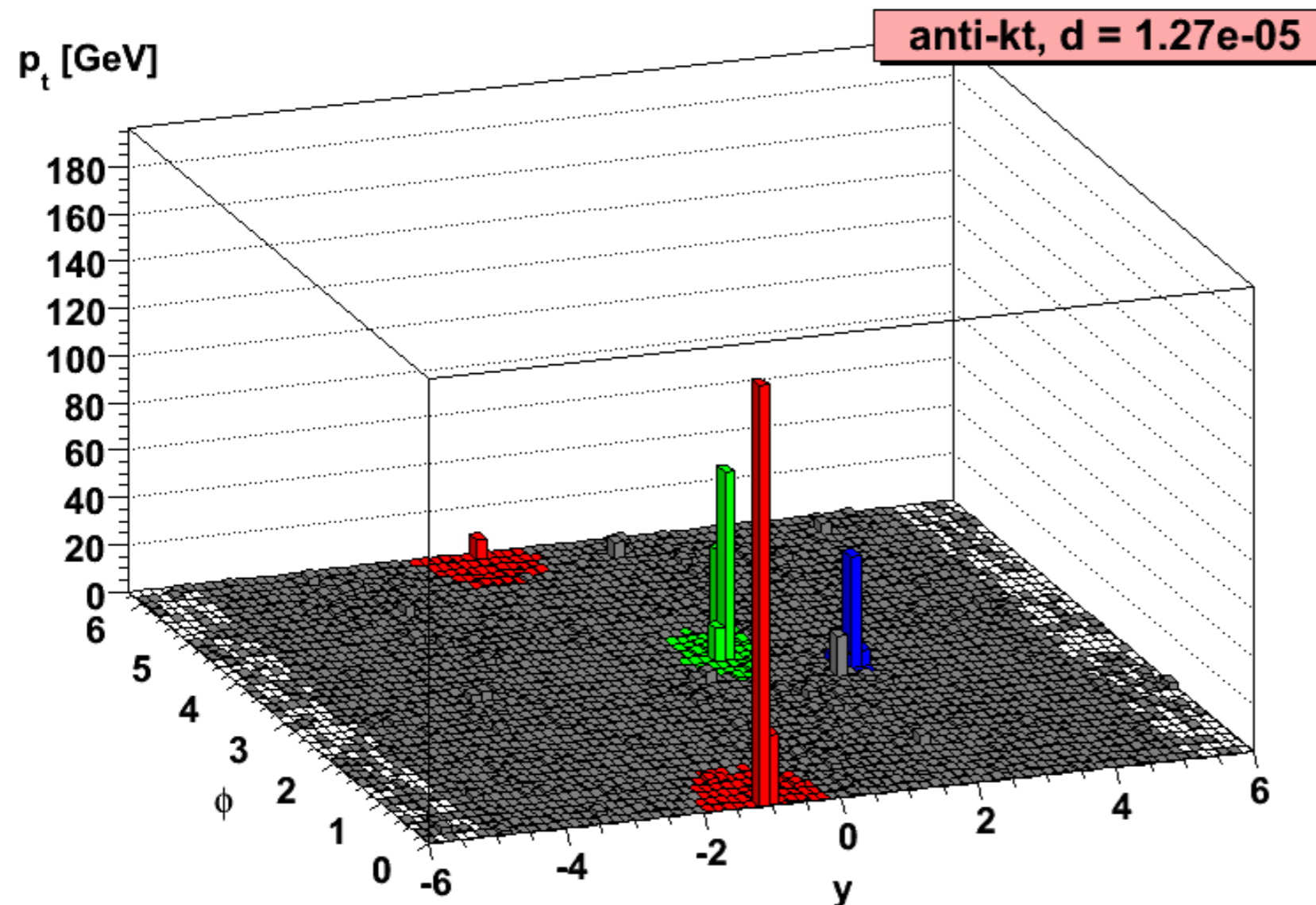
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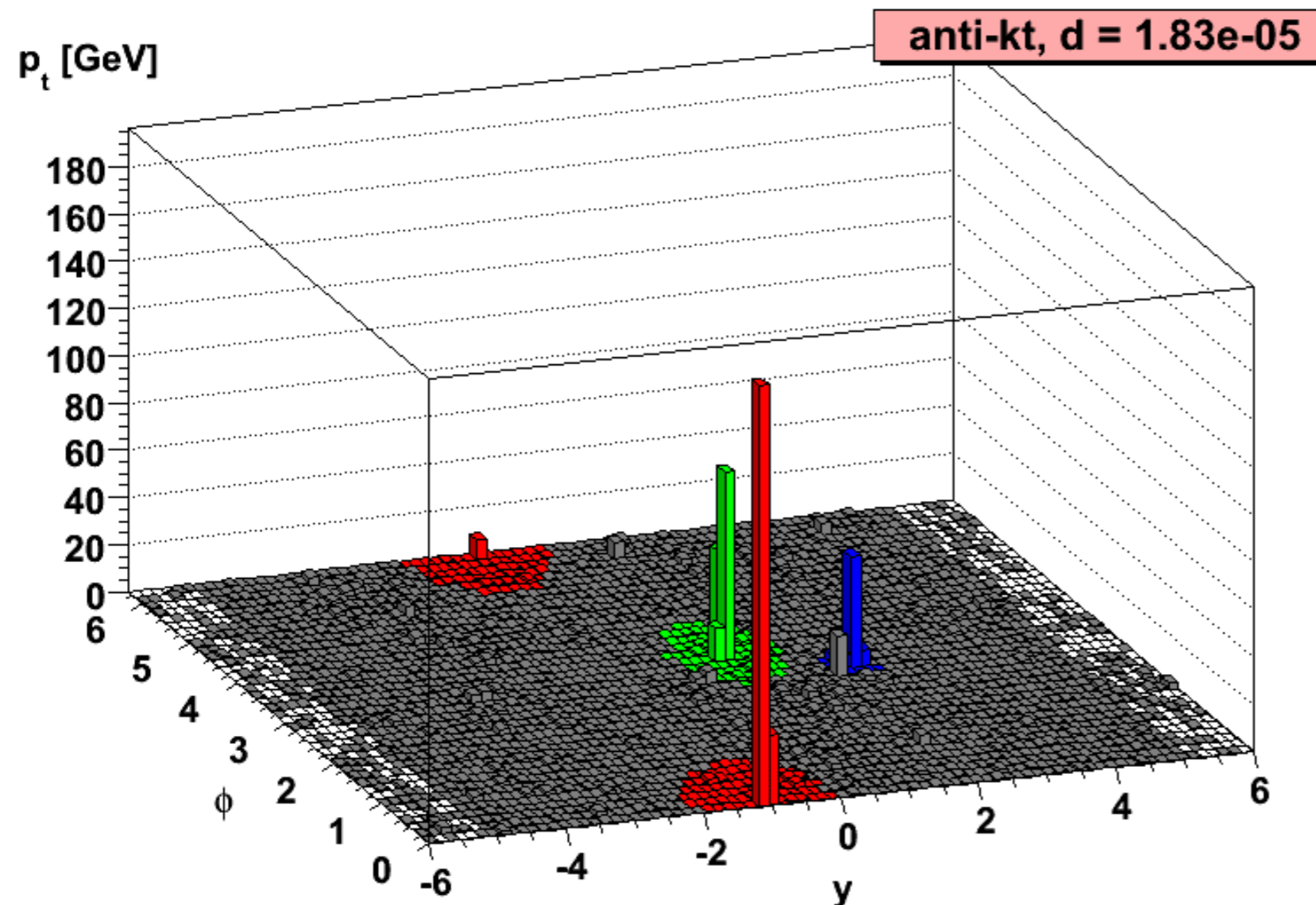
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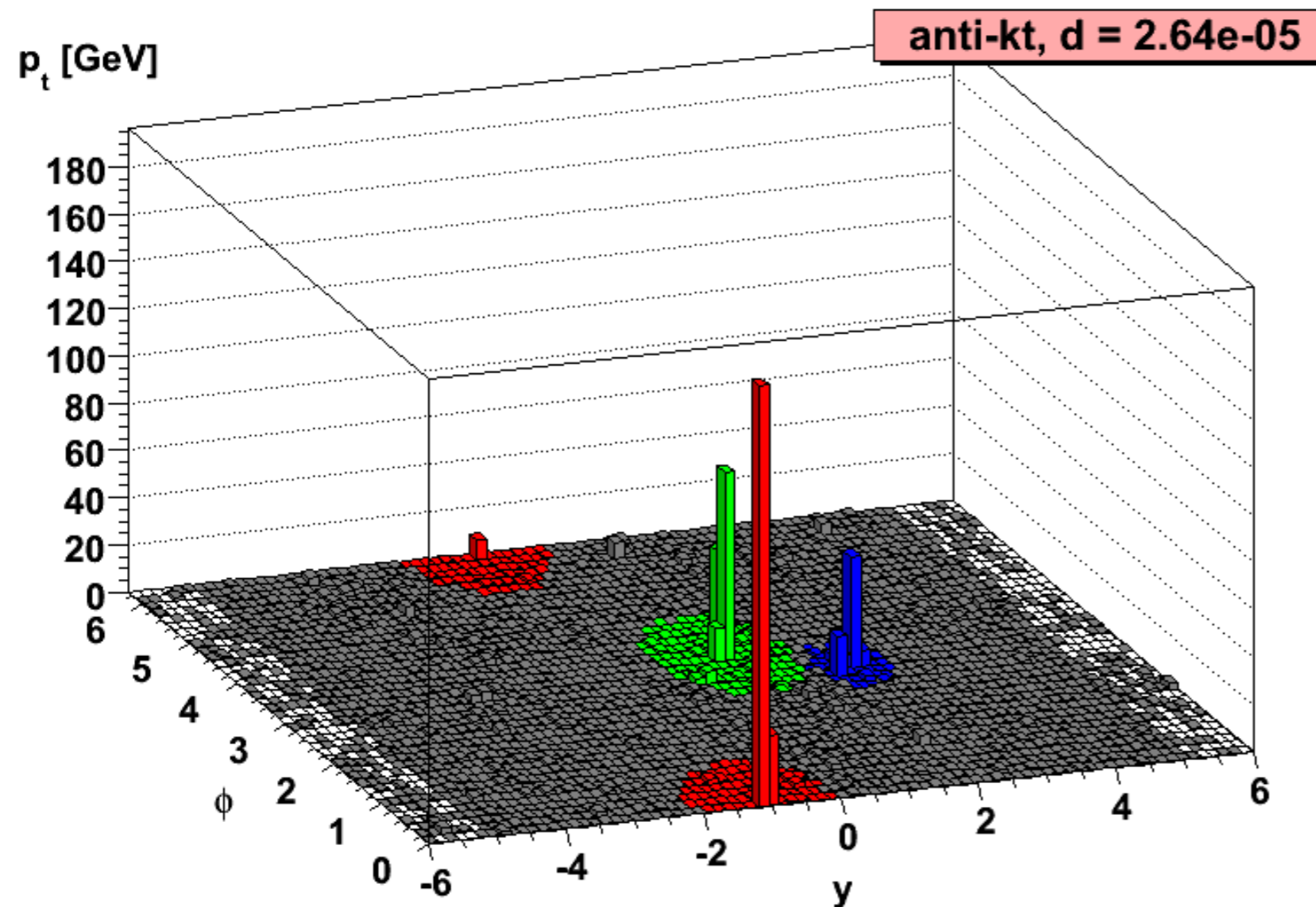
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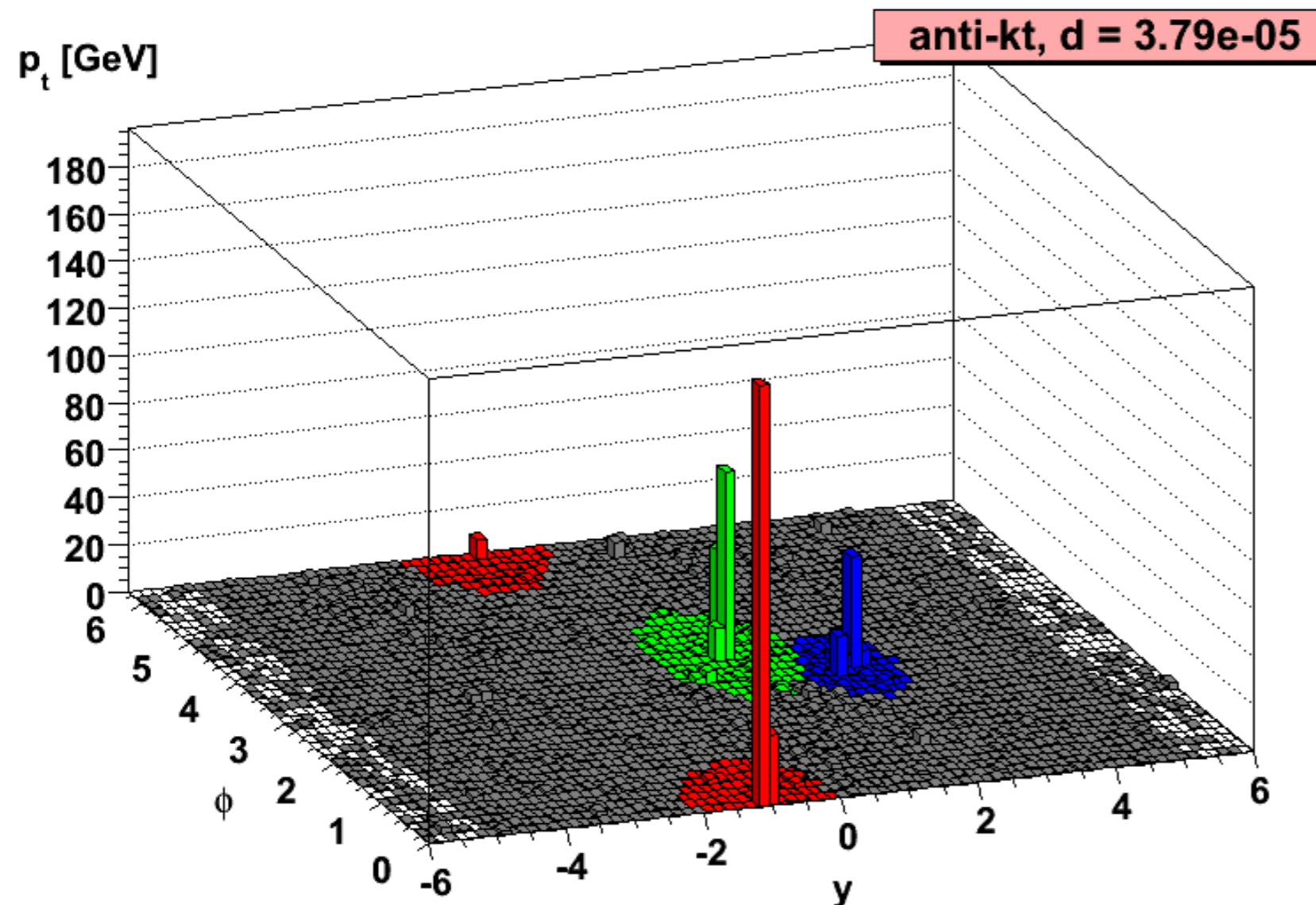
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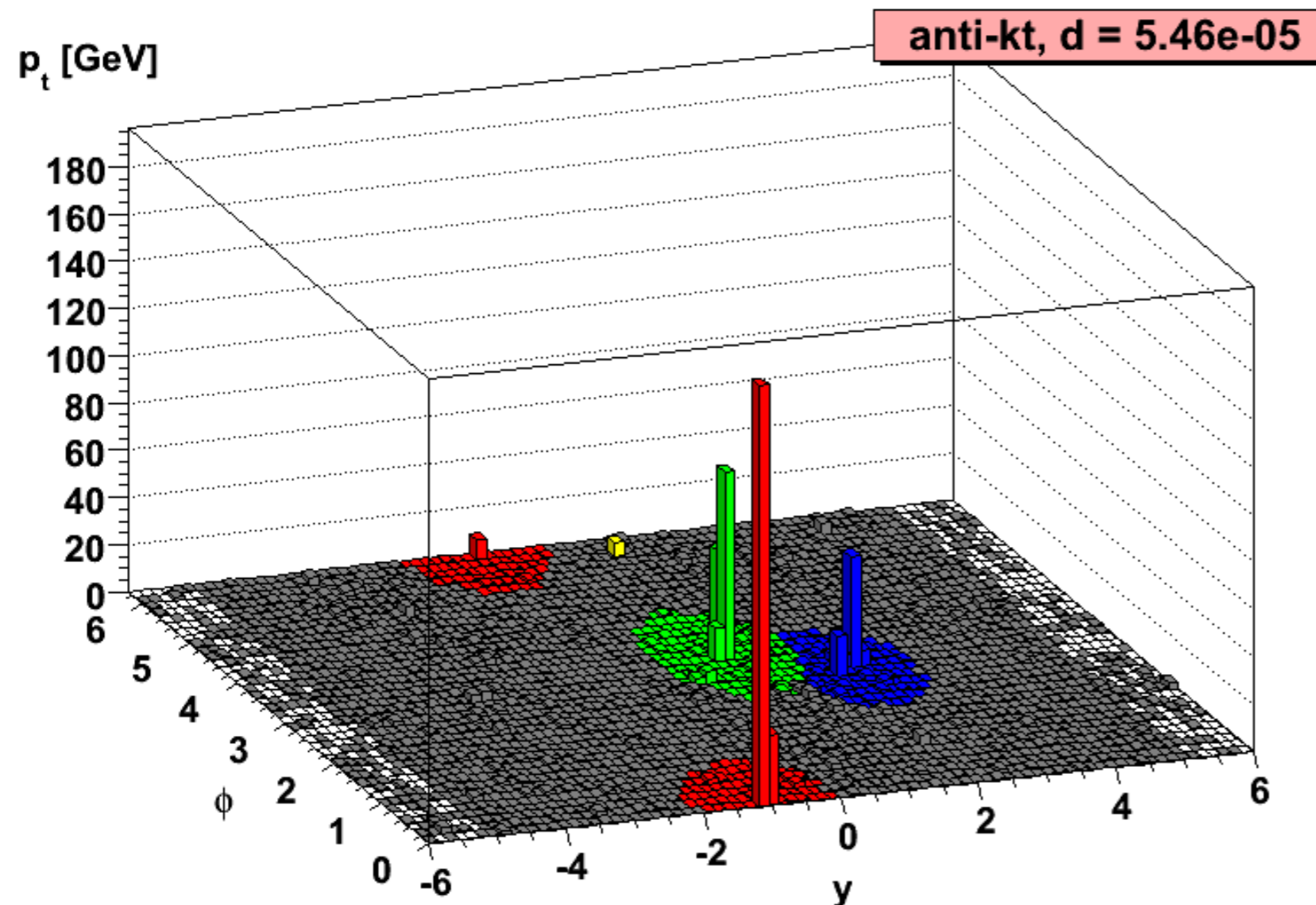
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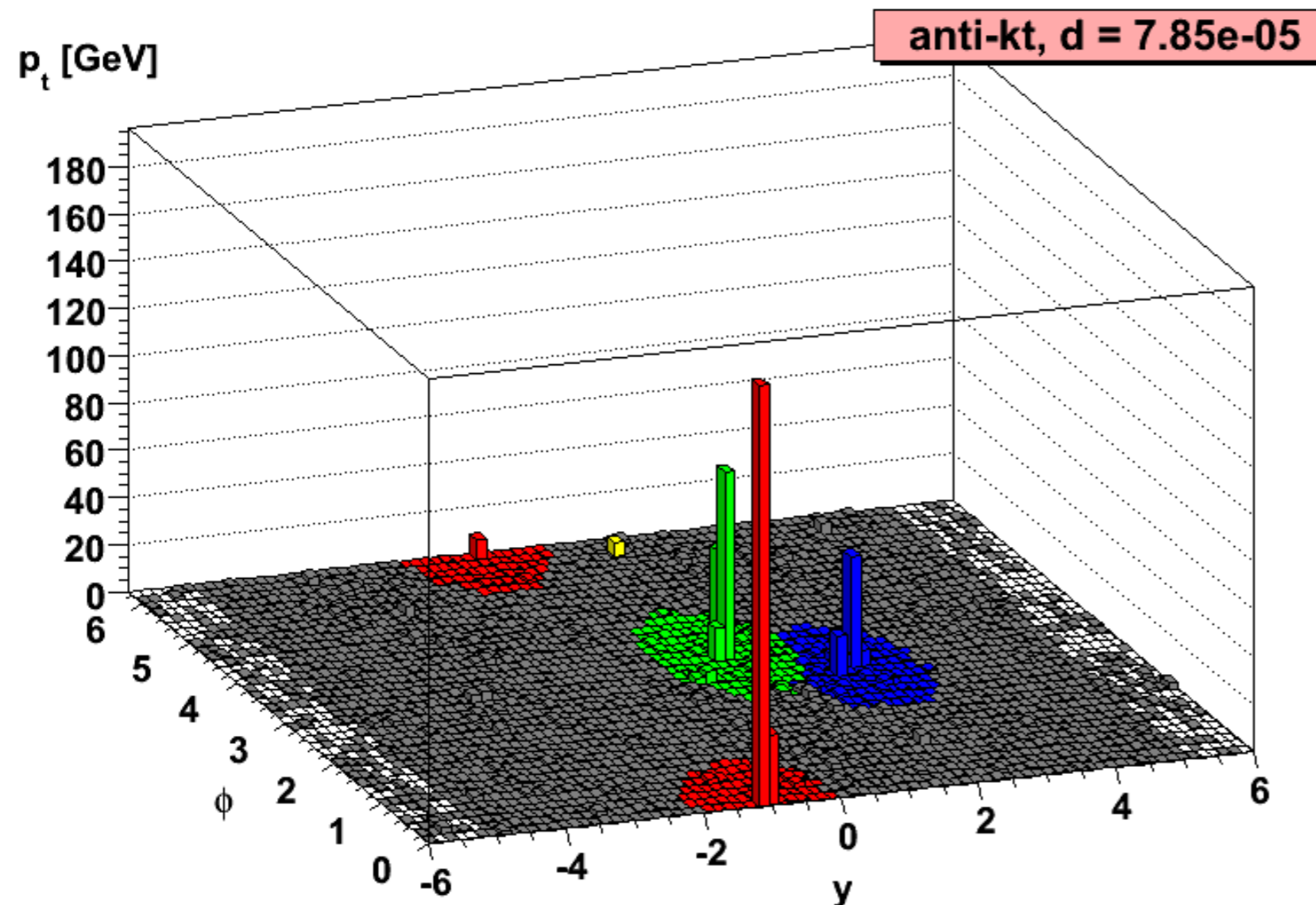
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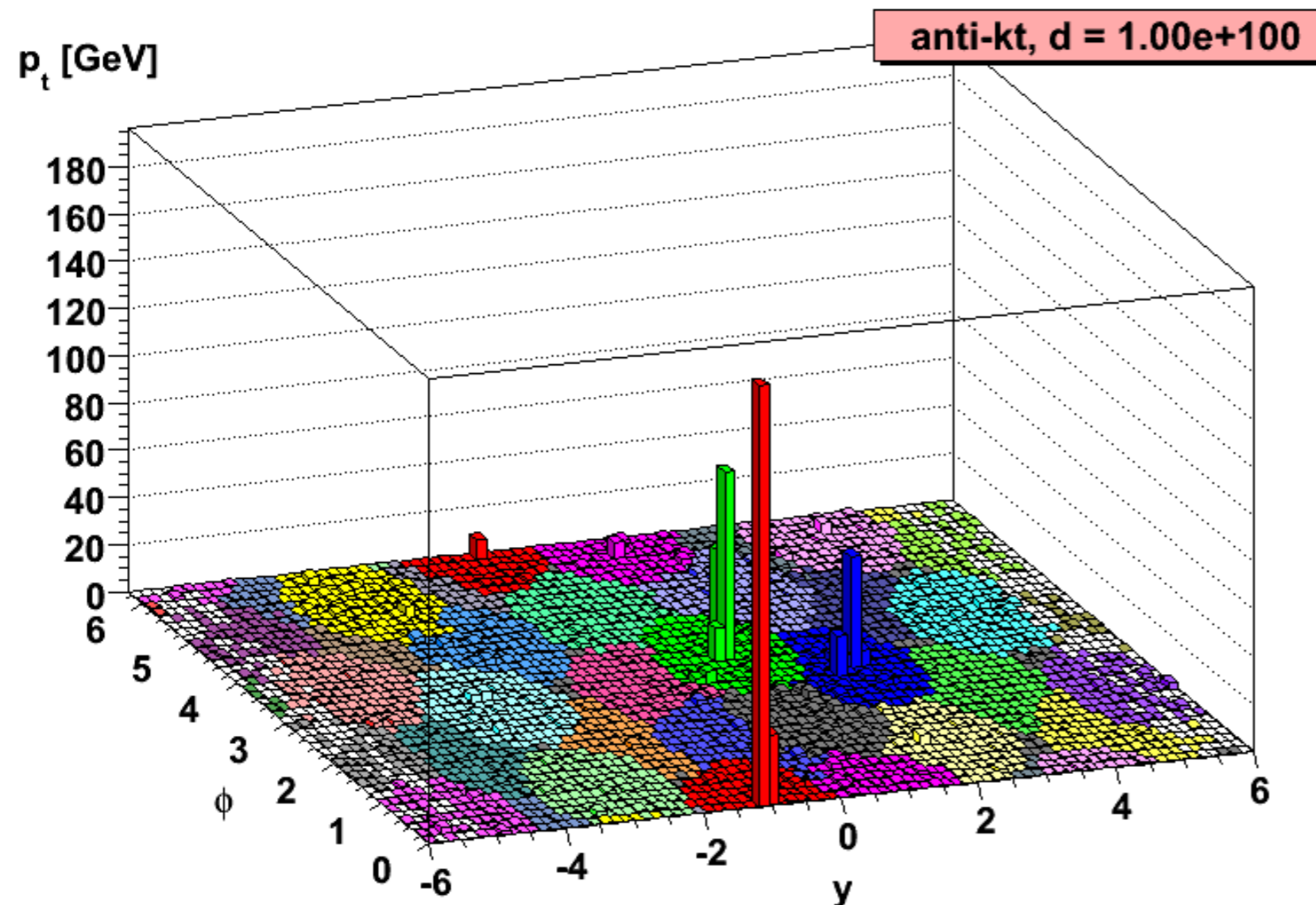
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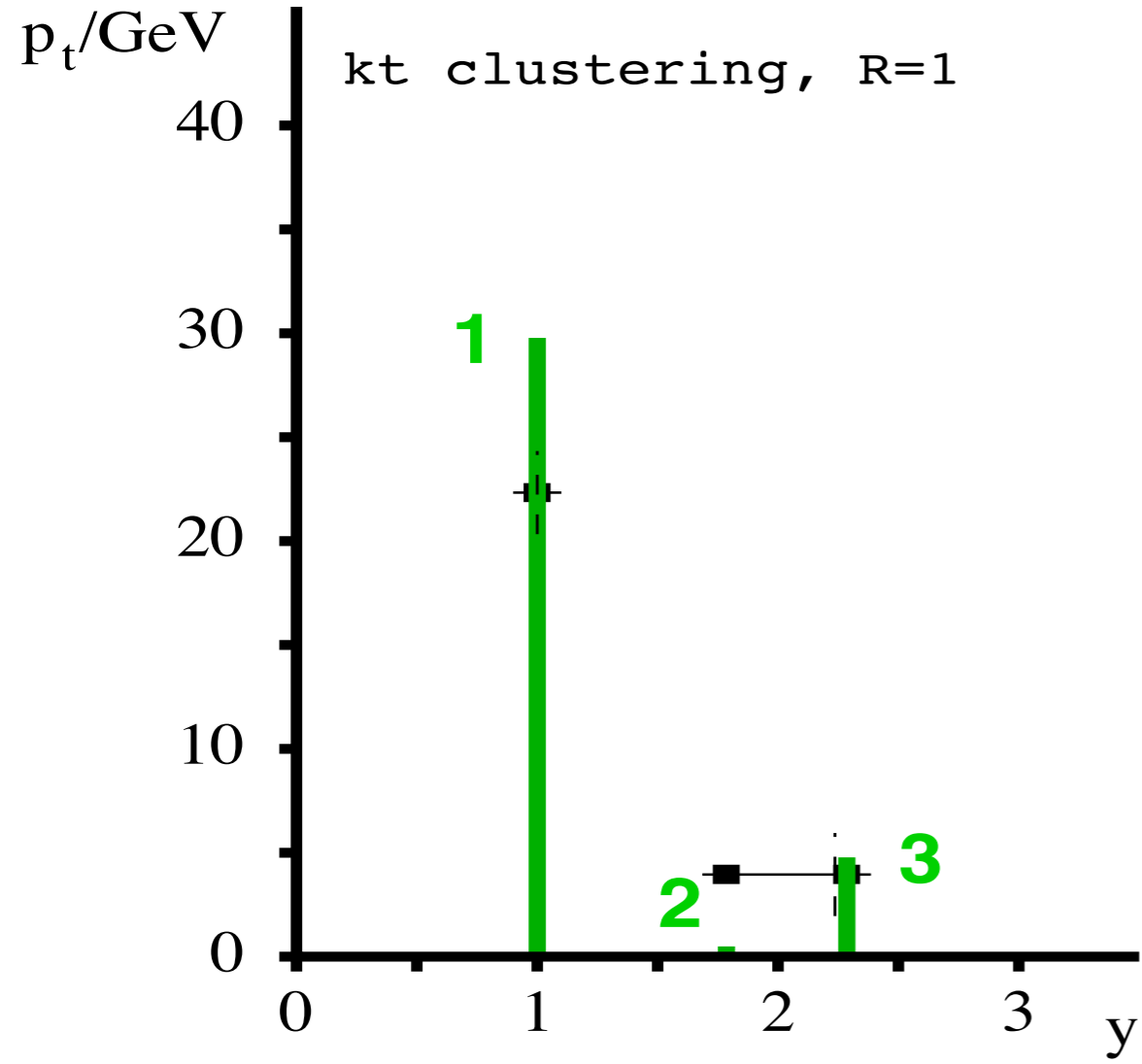


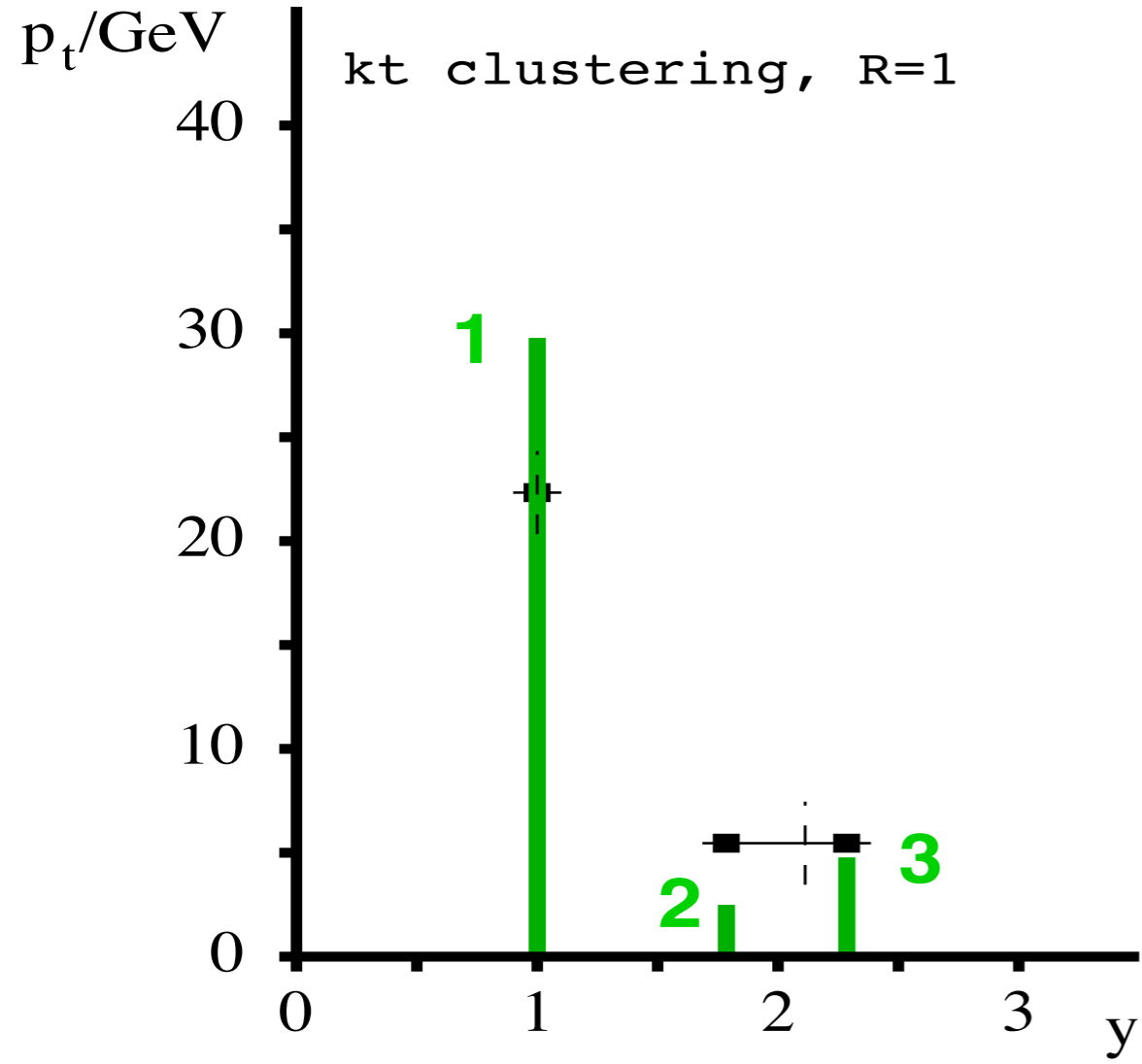
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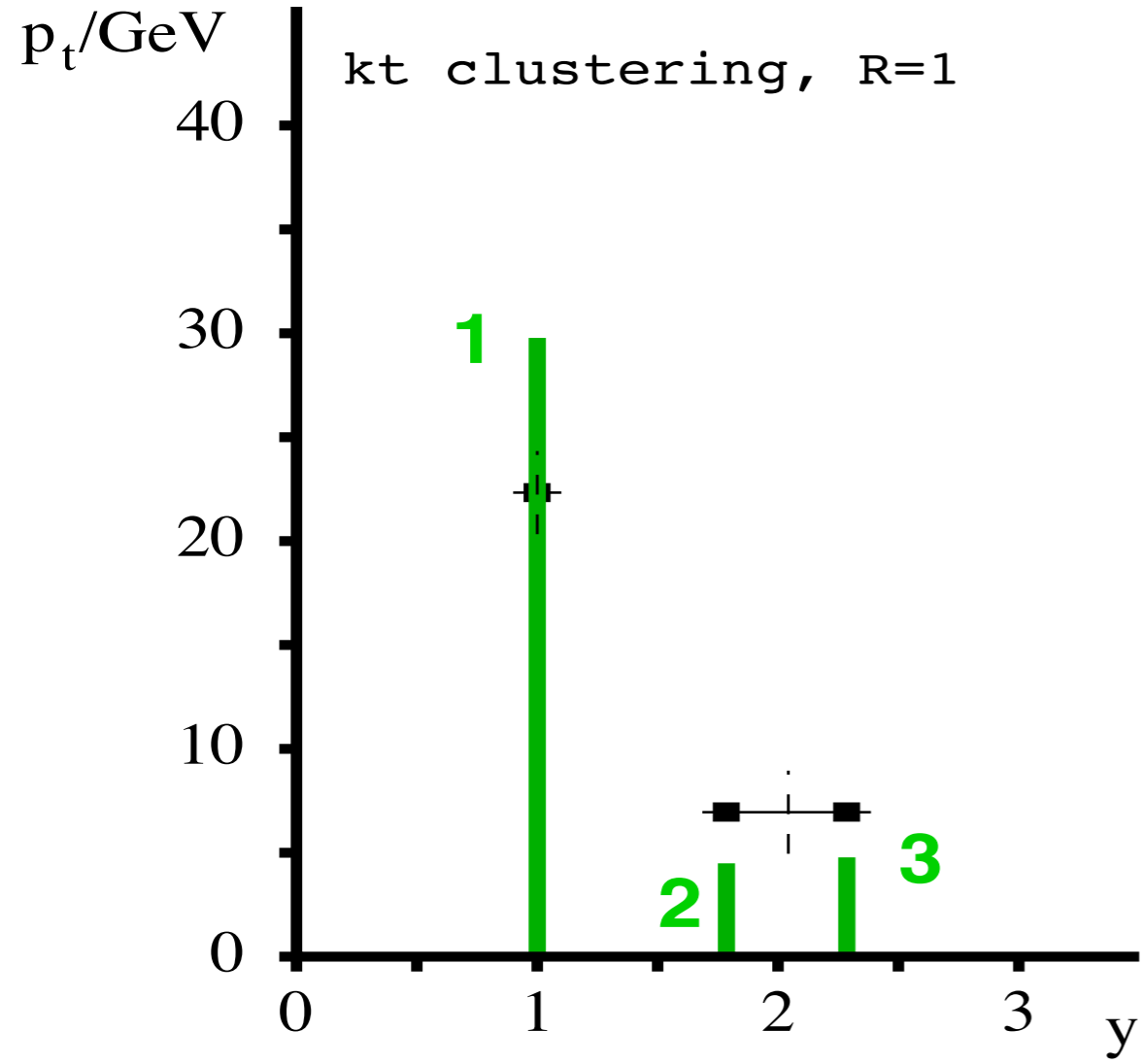
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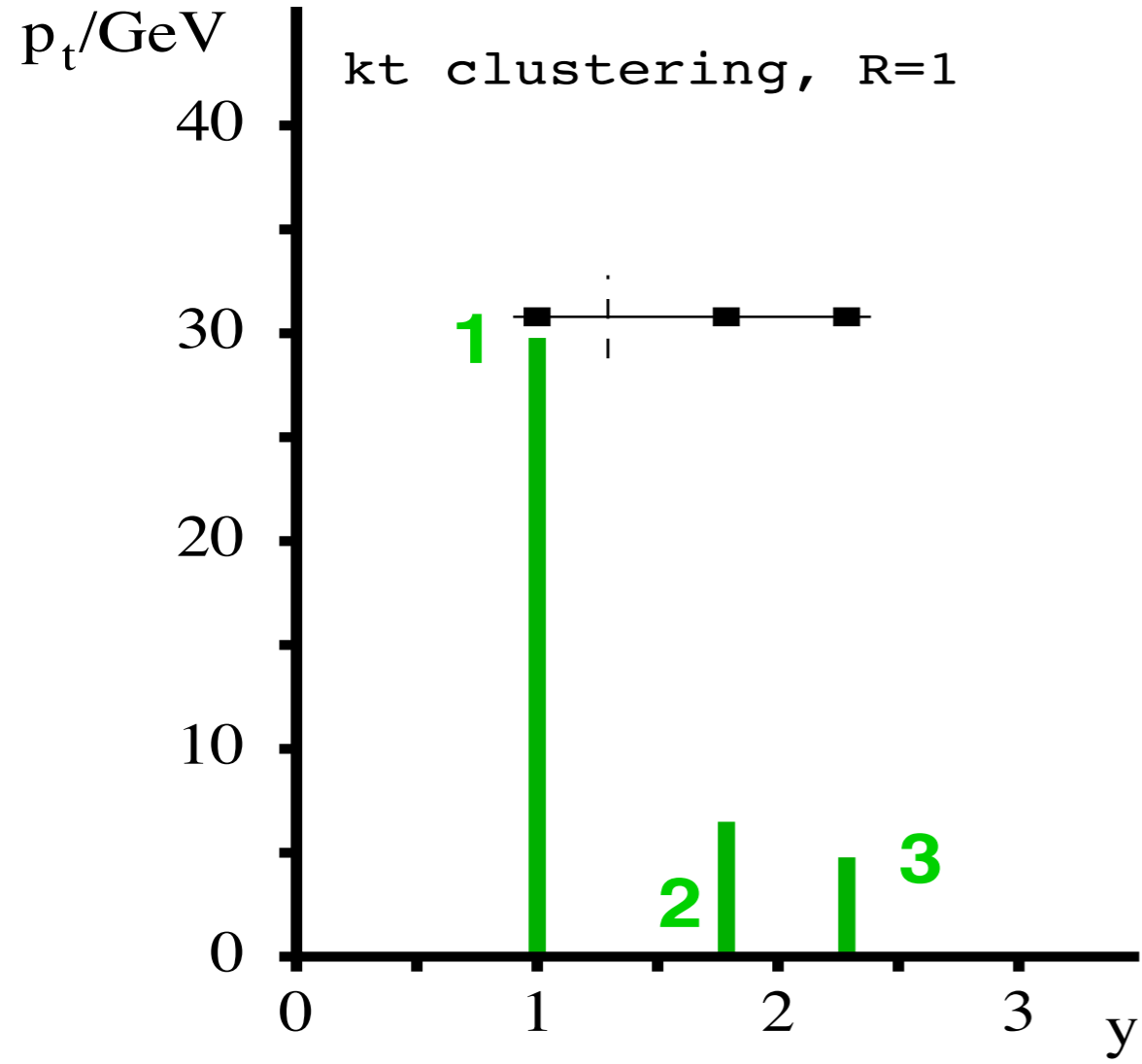


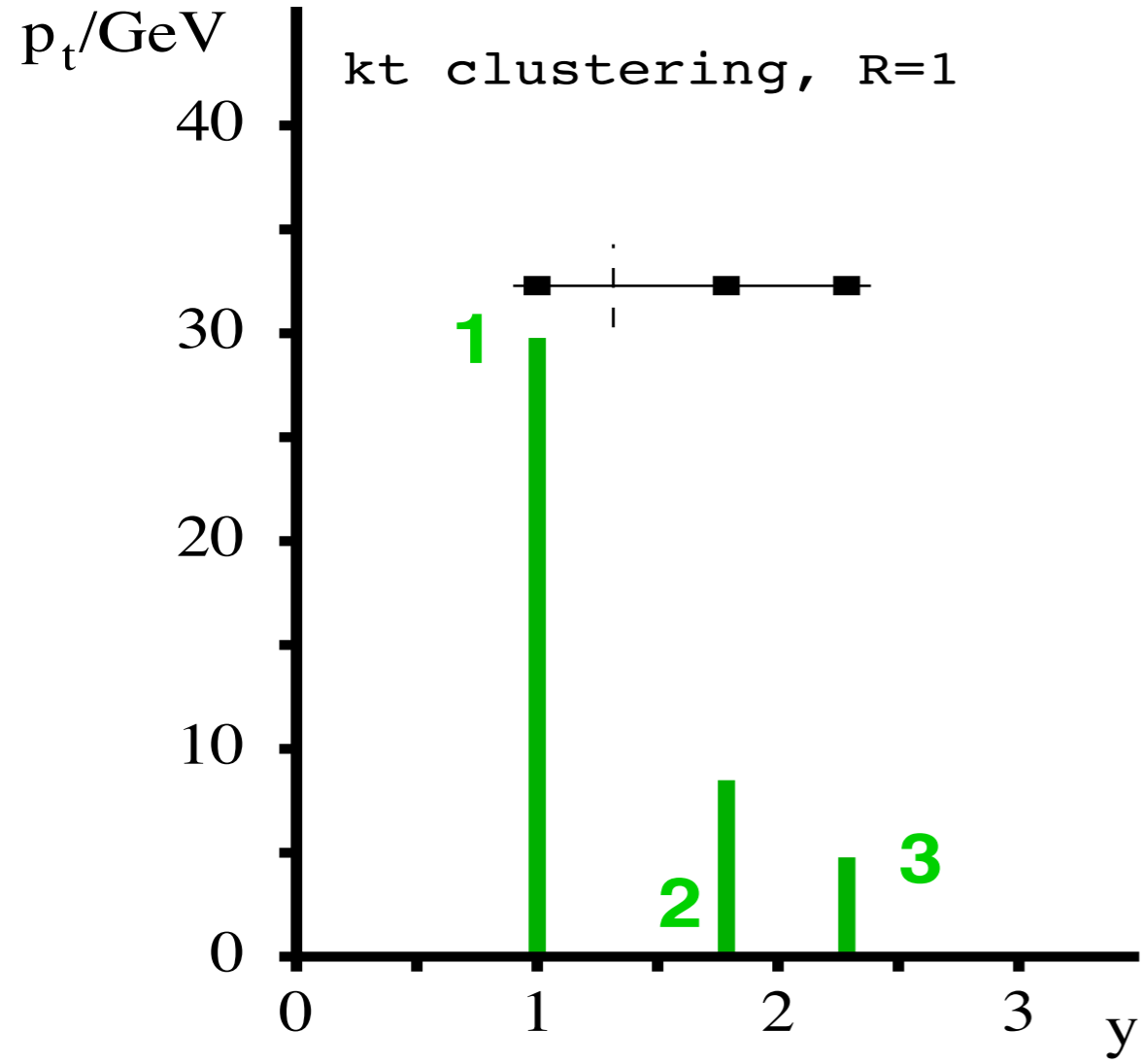
Anti- k_t gives circular jets (“cone-like”) in a way that’s infrared safe

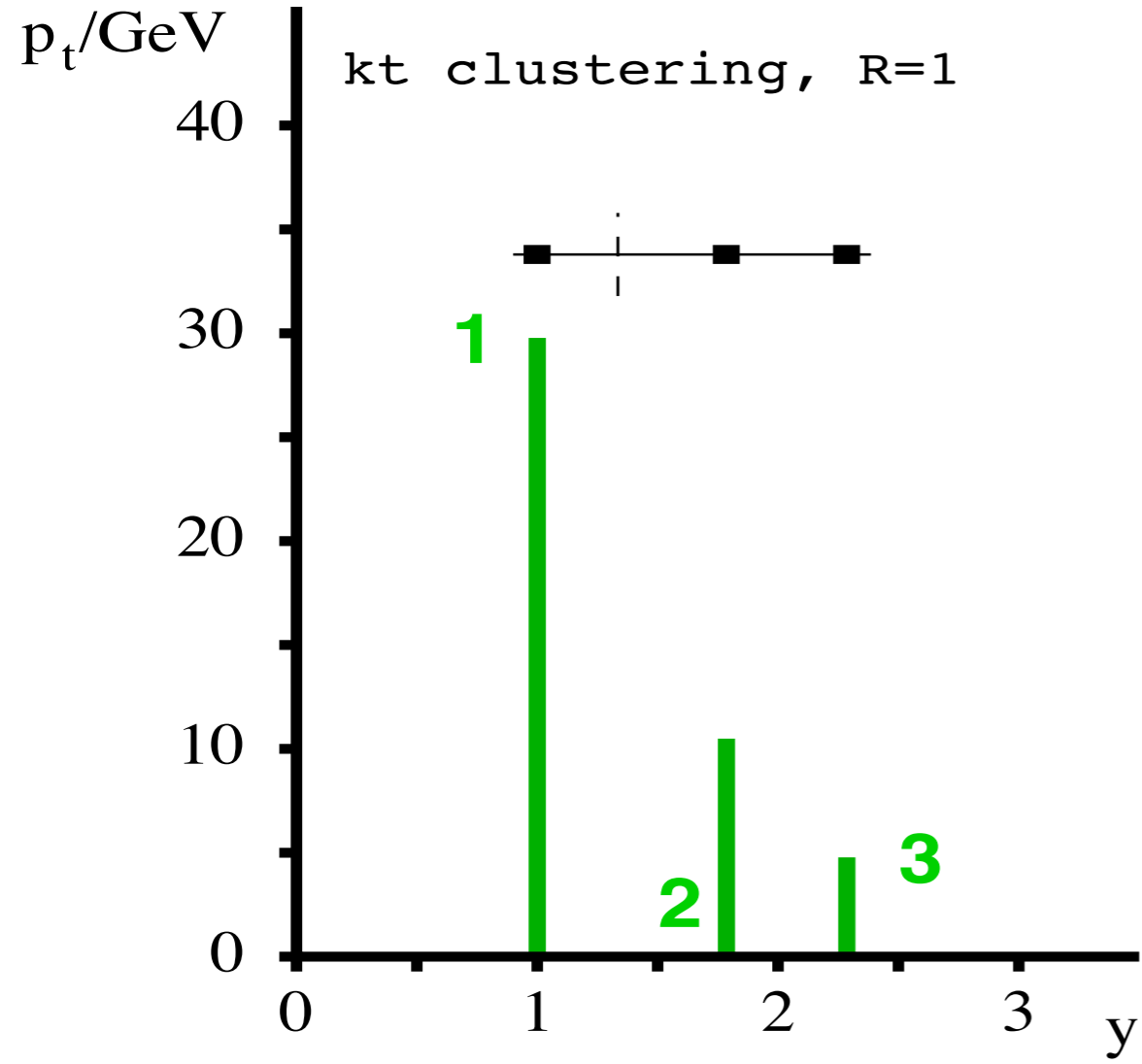


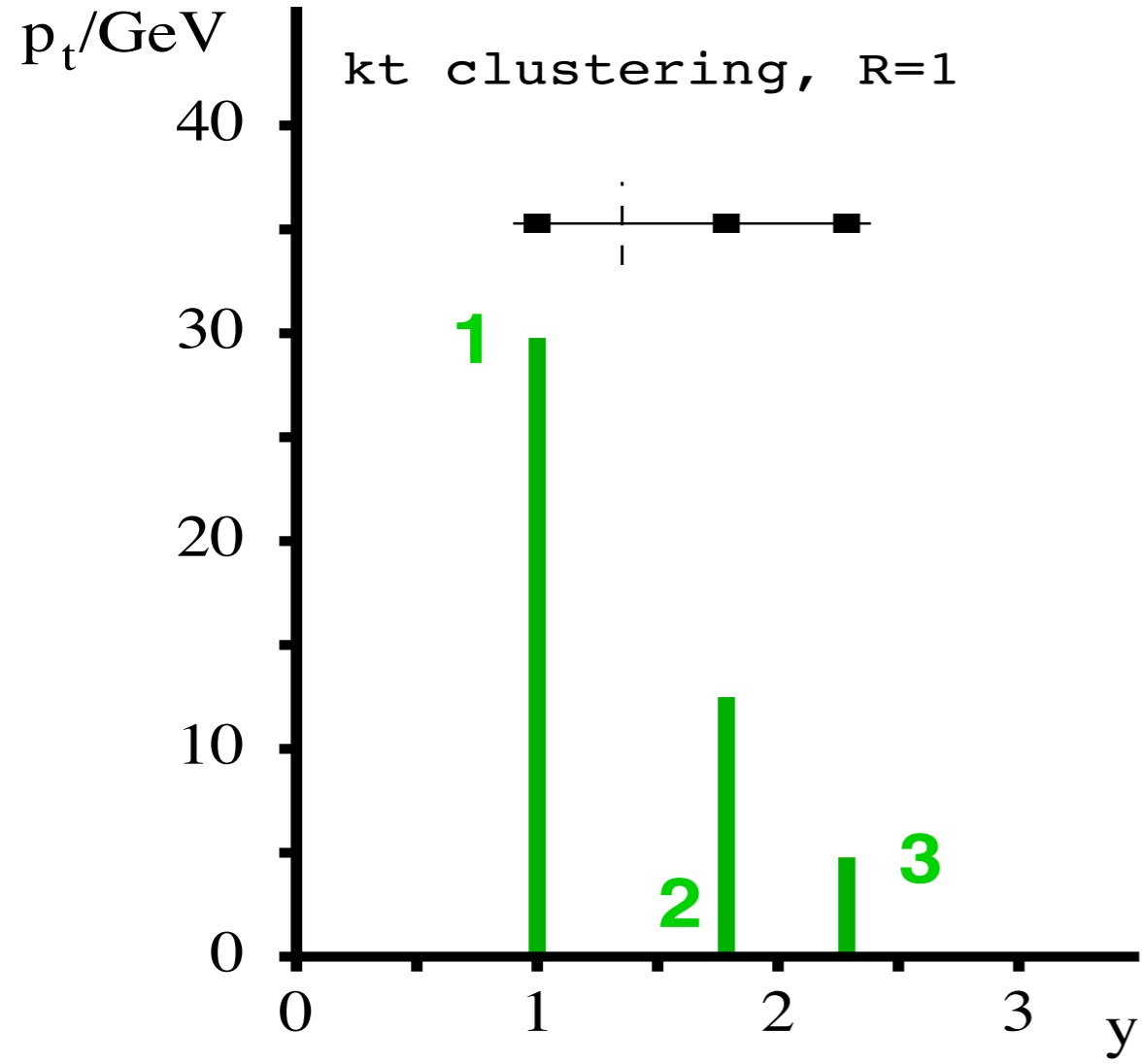


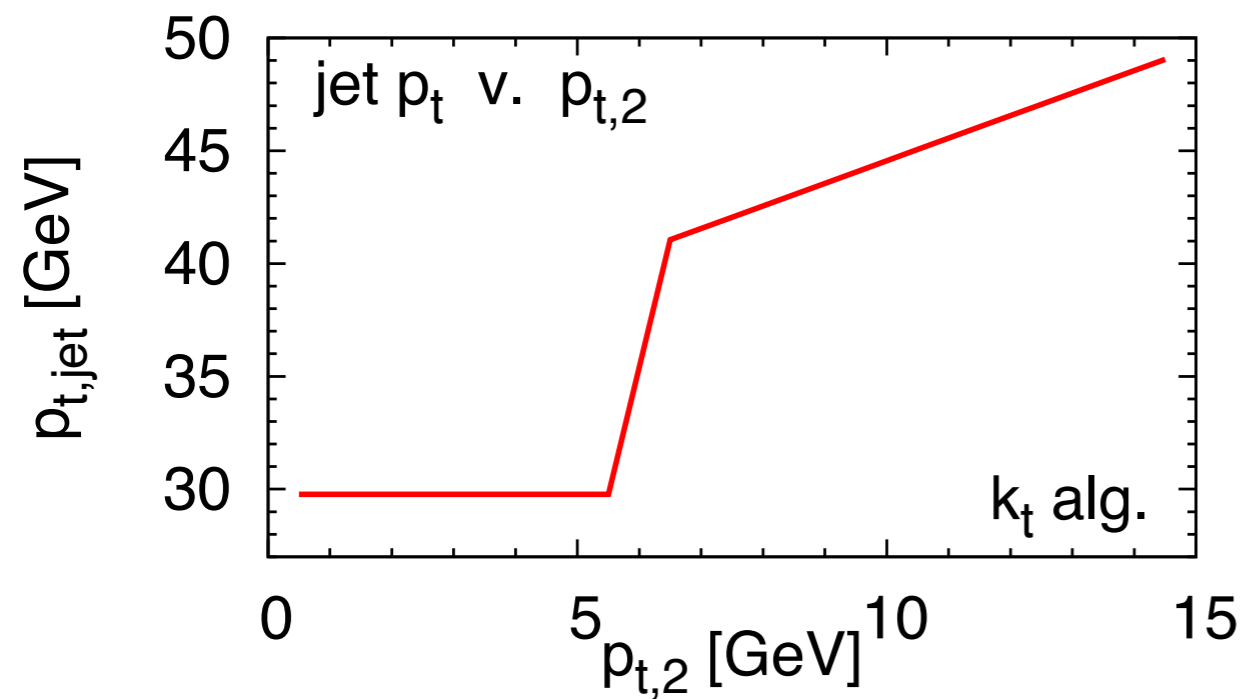
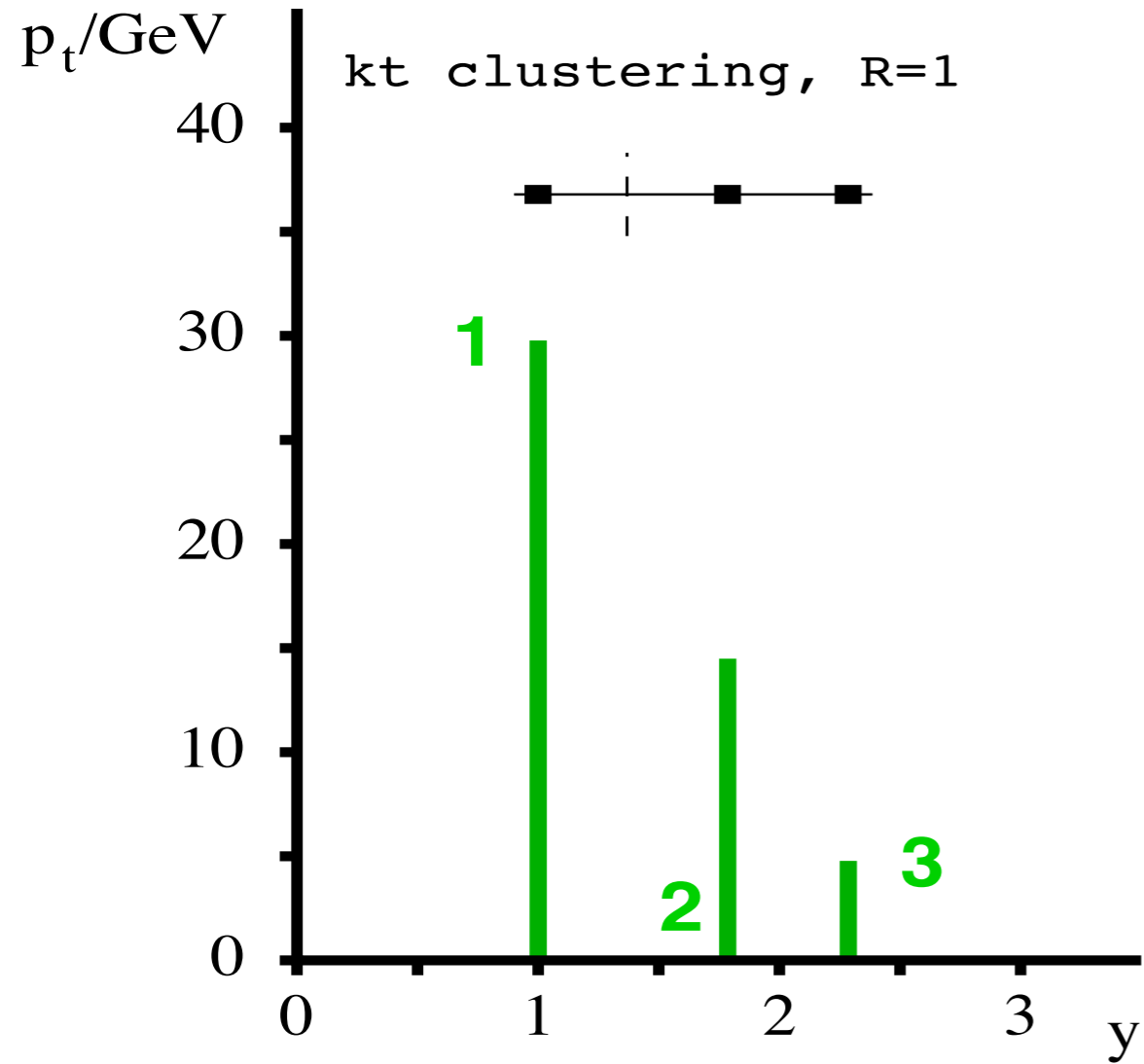




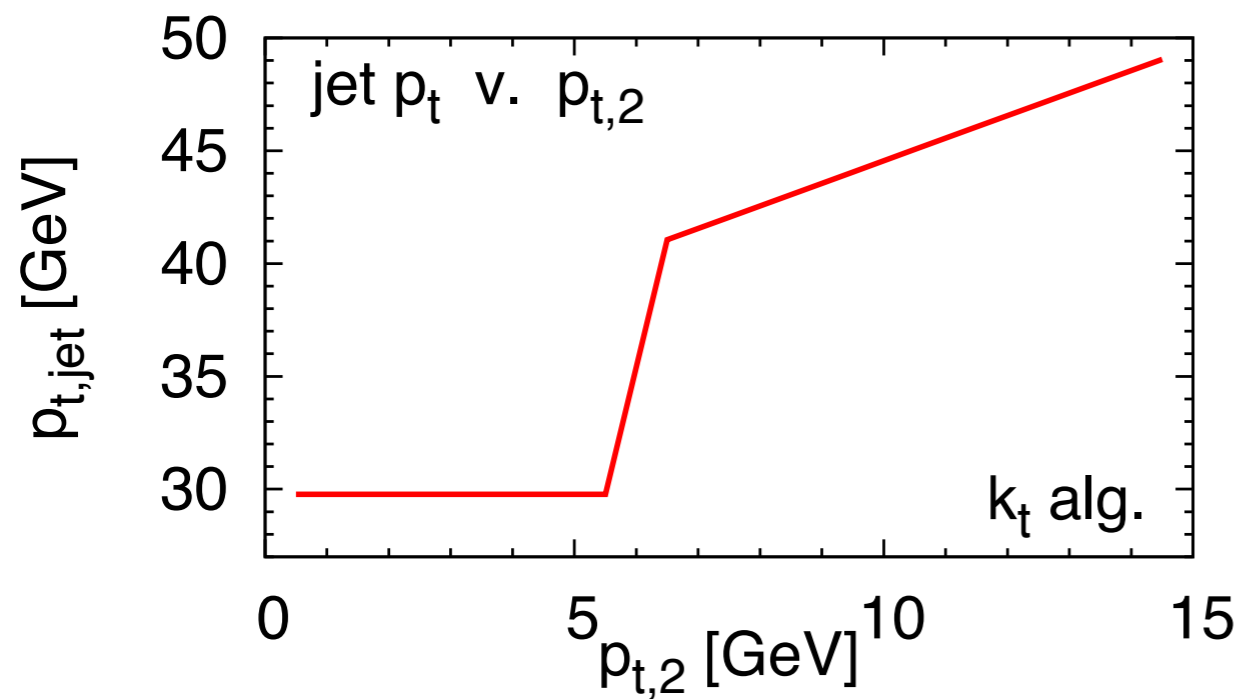
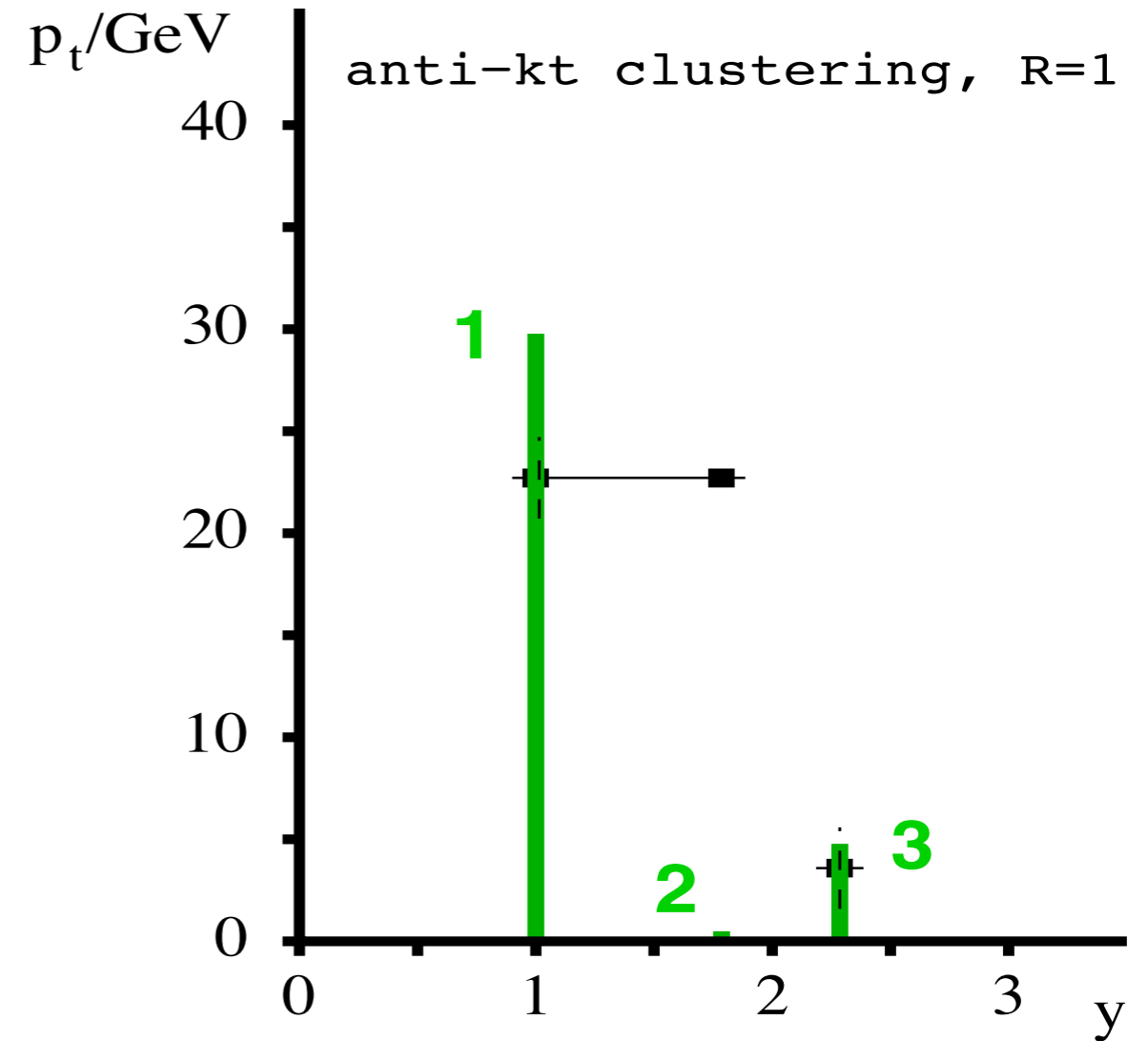
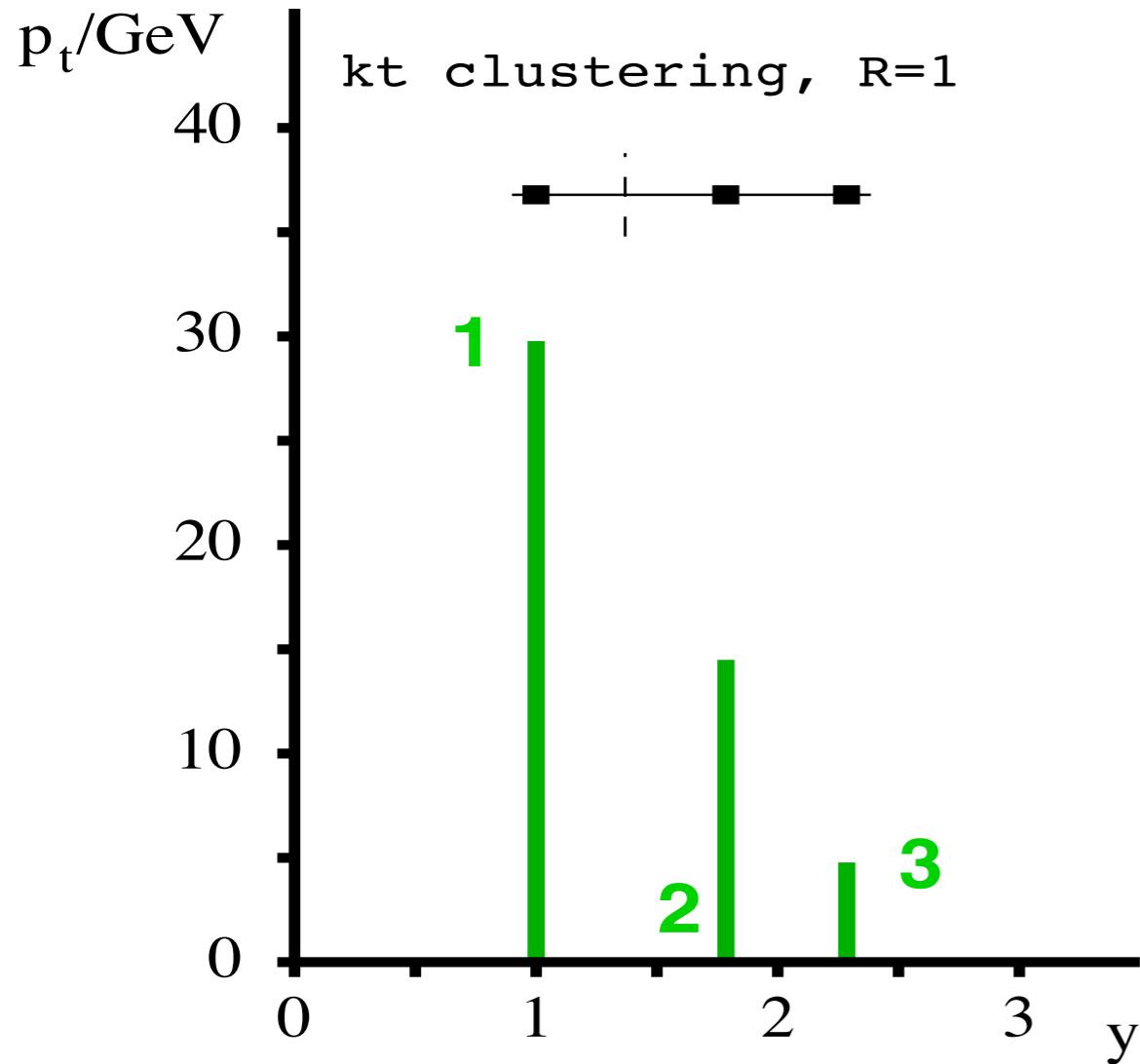




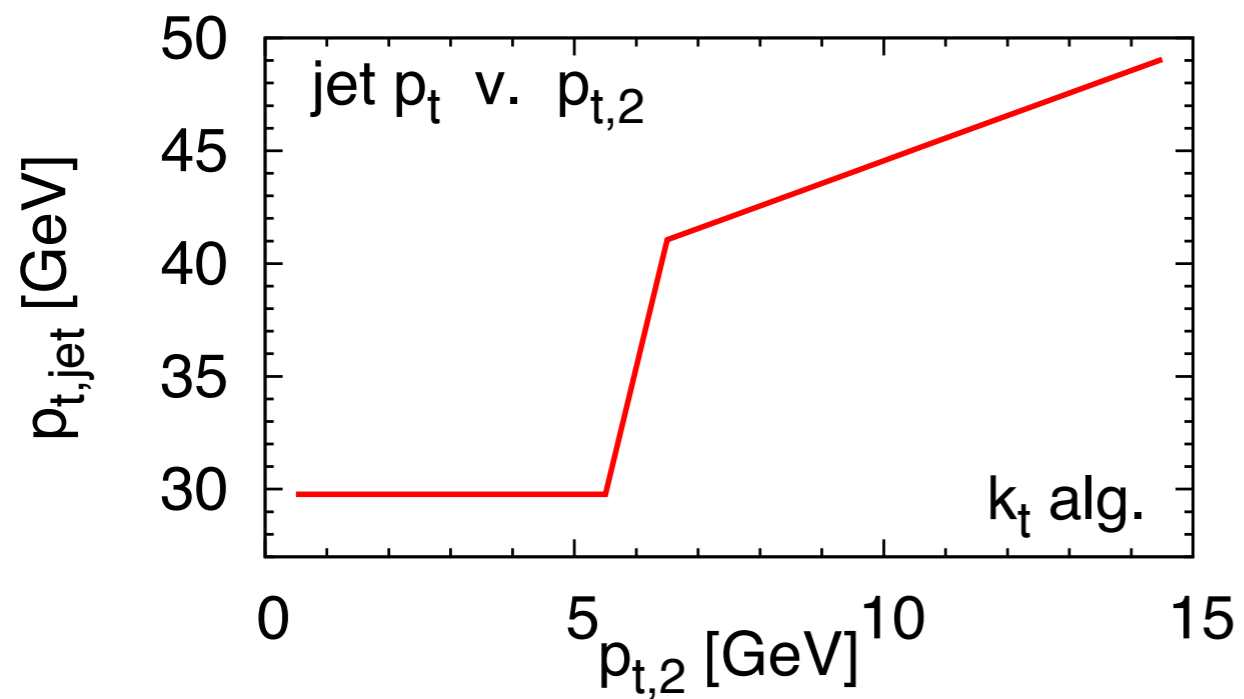
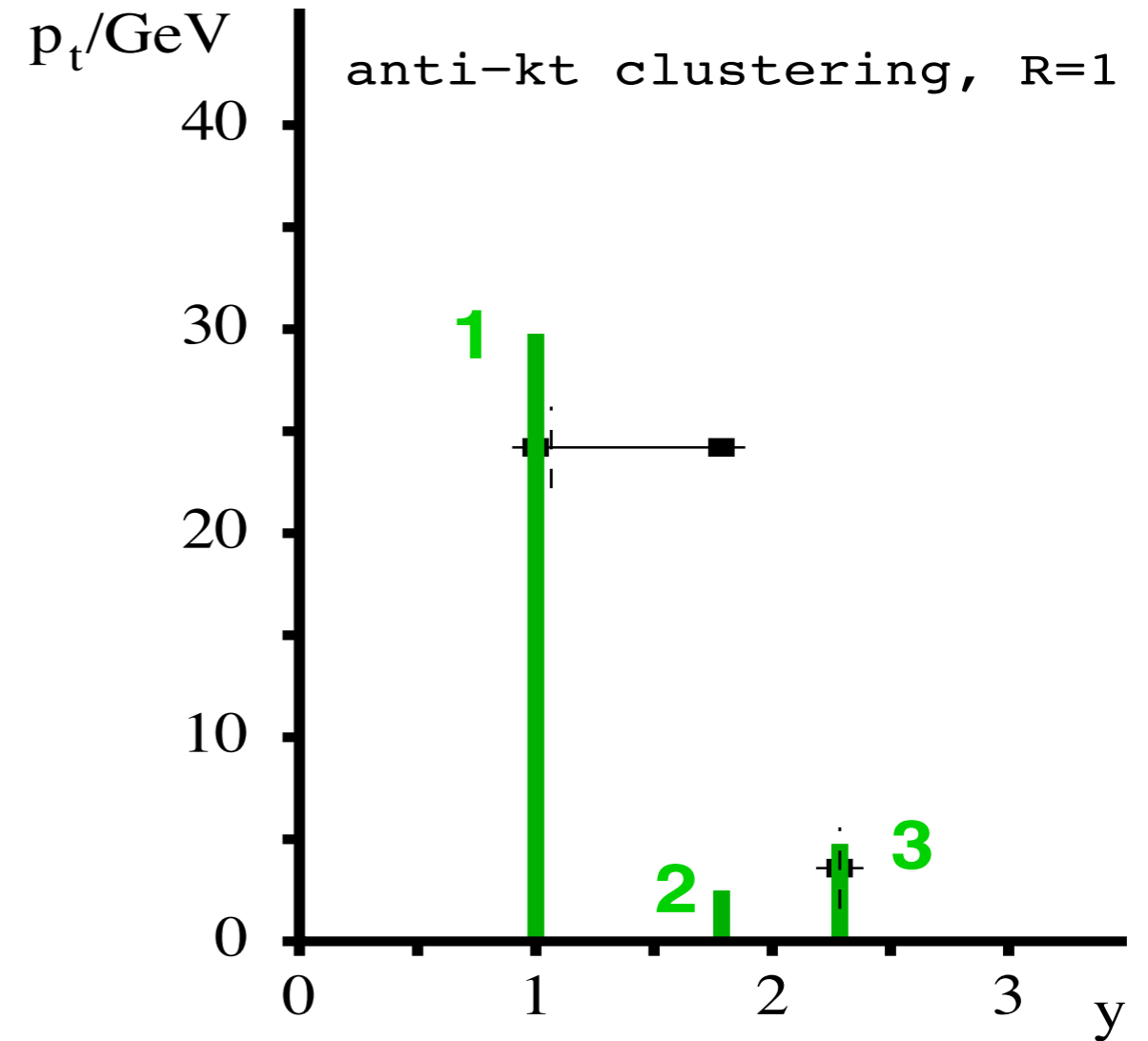
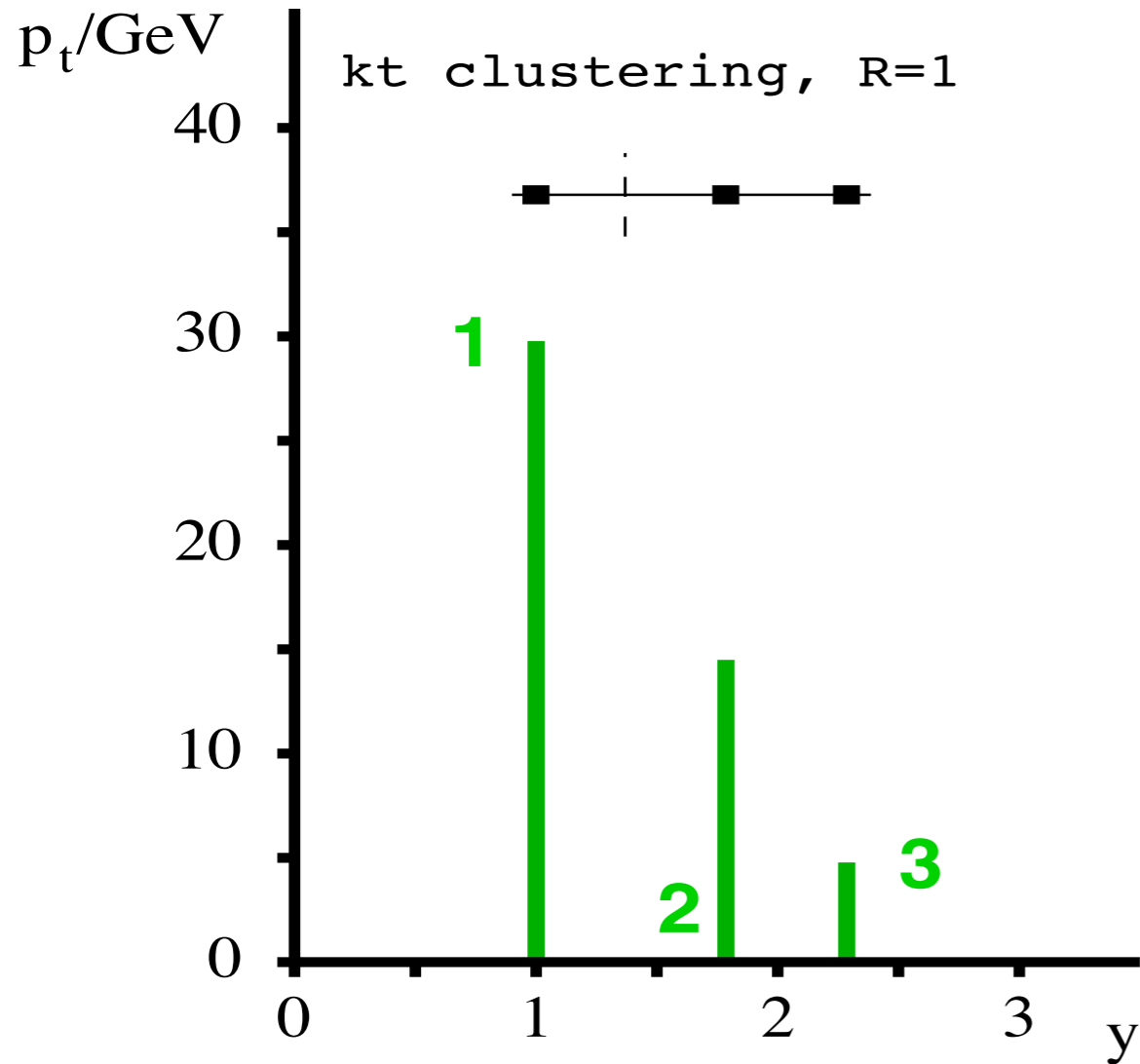




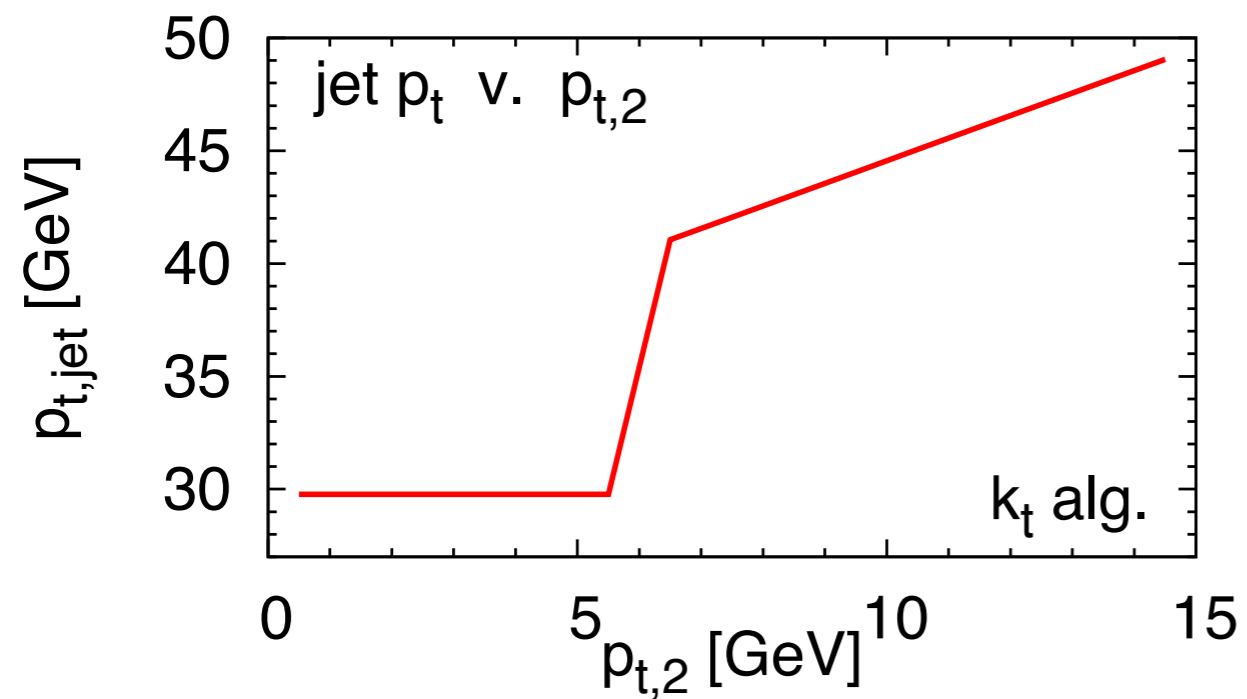
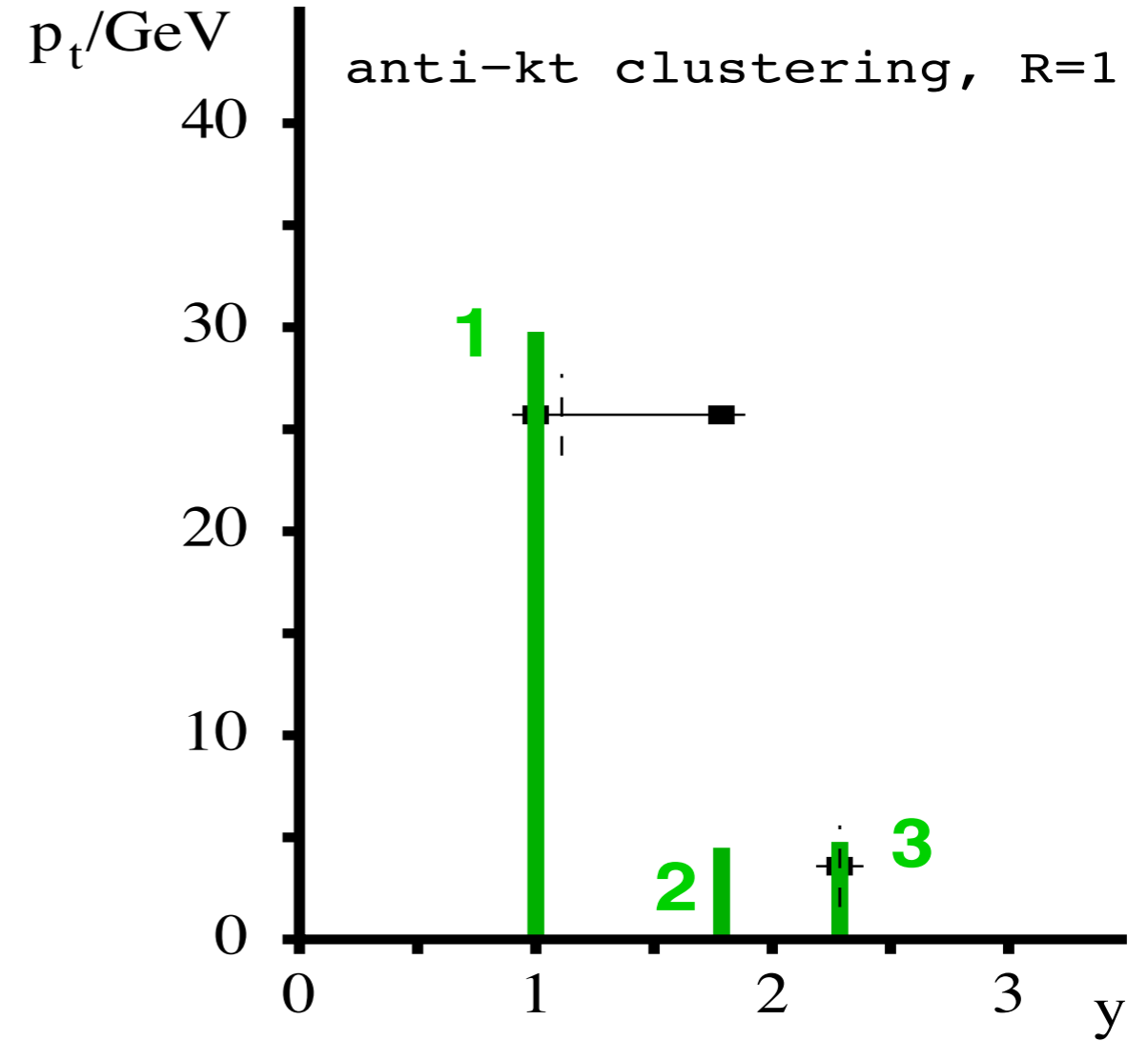
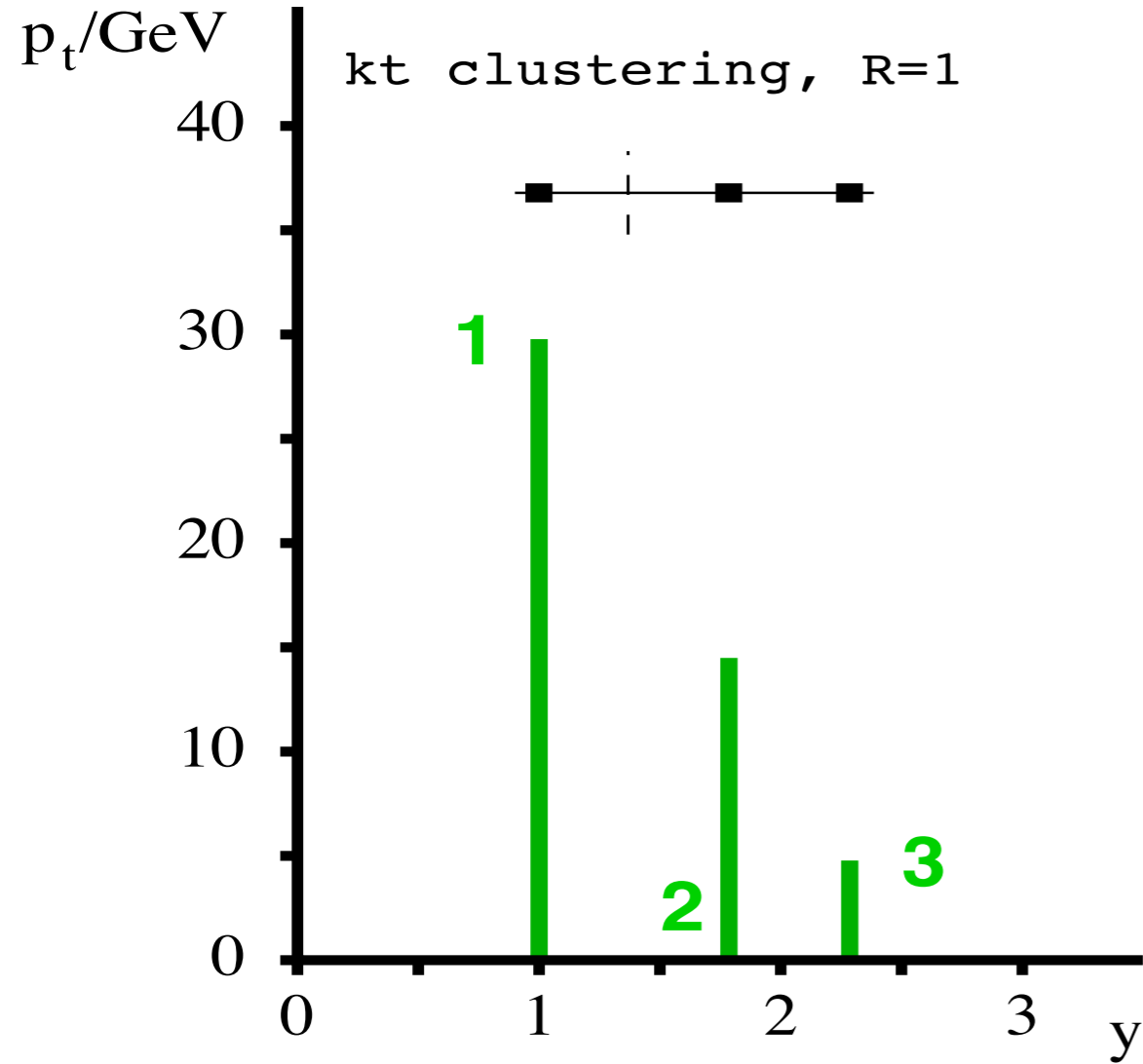
Linearity: k_t v. anti- k_t



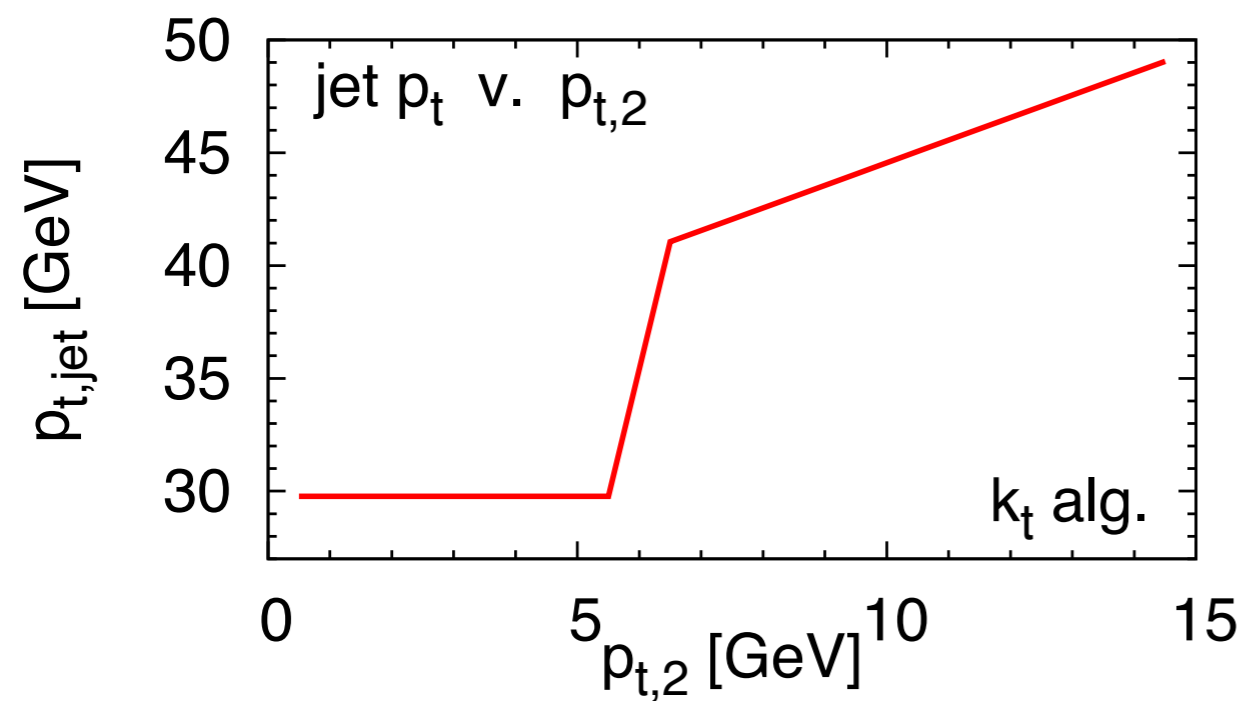
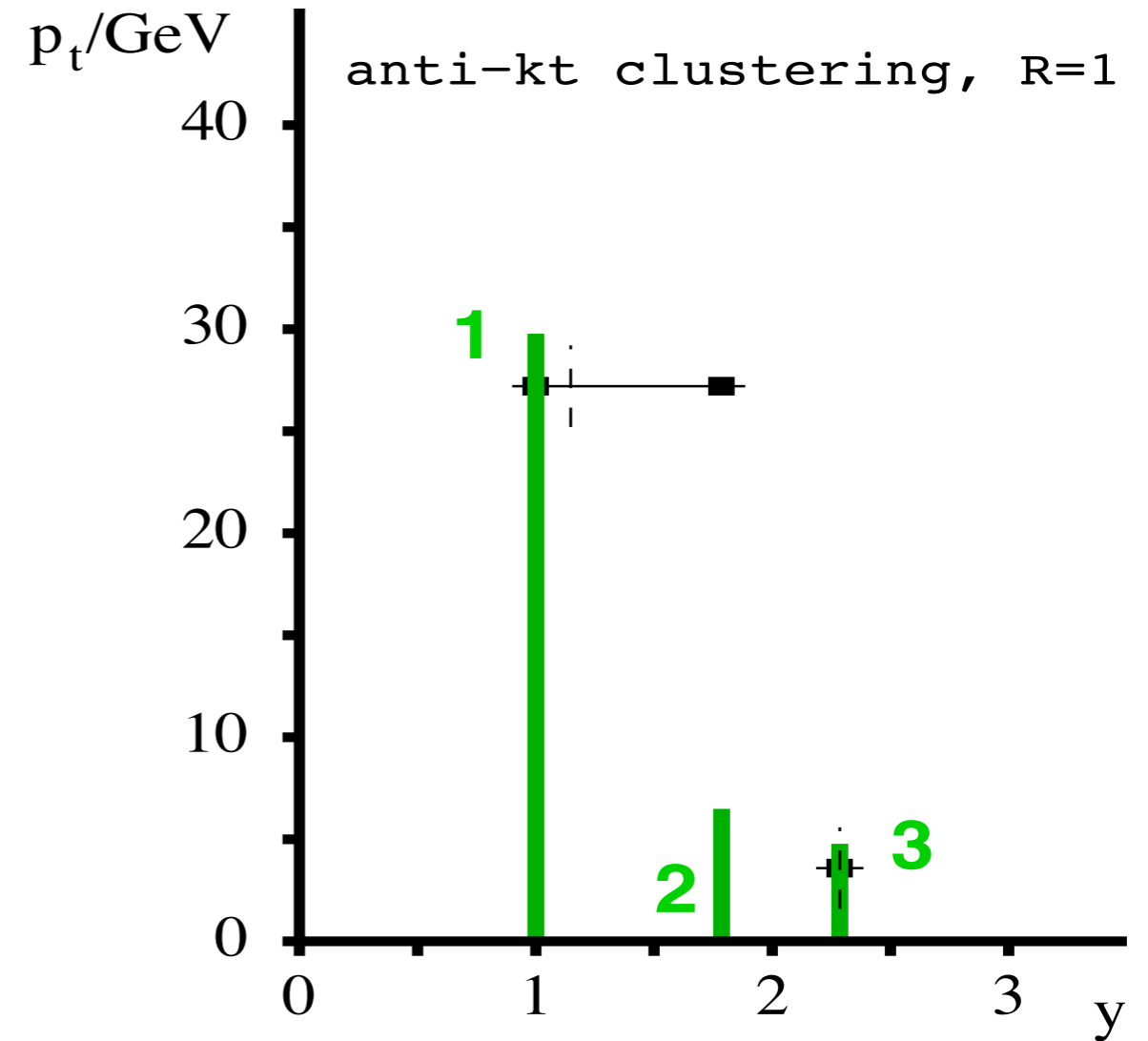
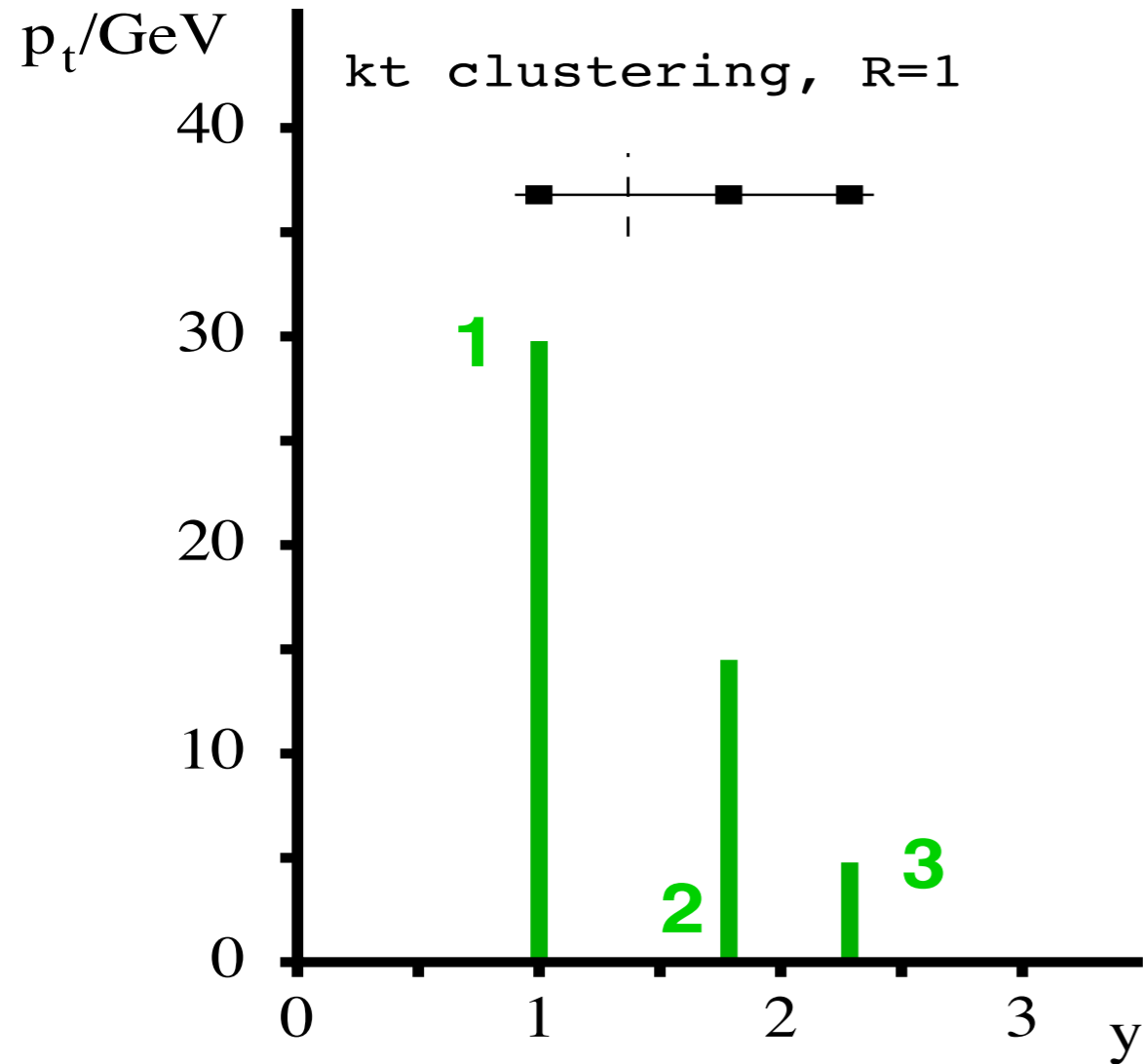
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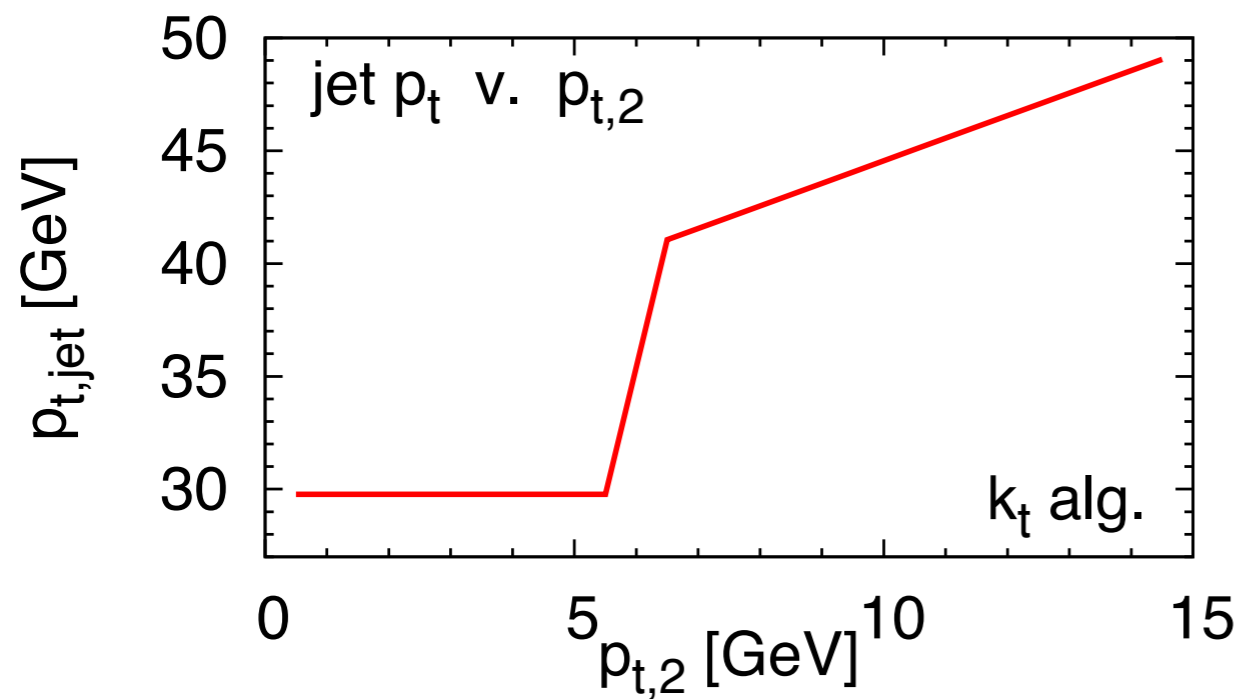
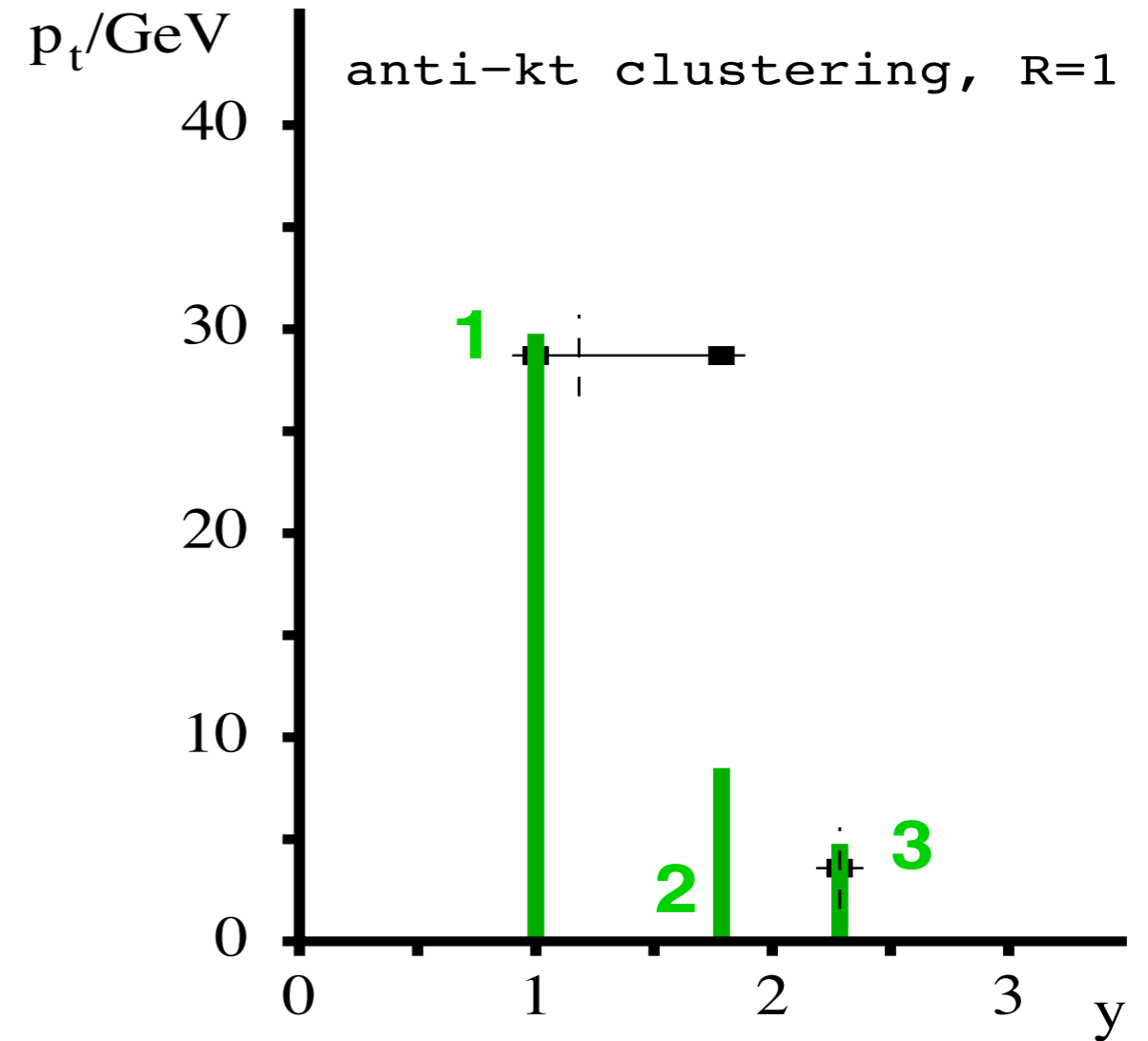
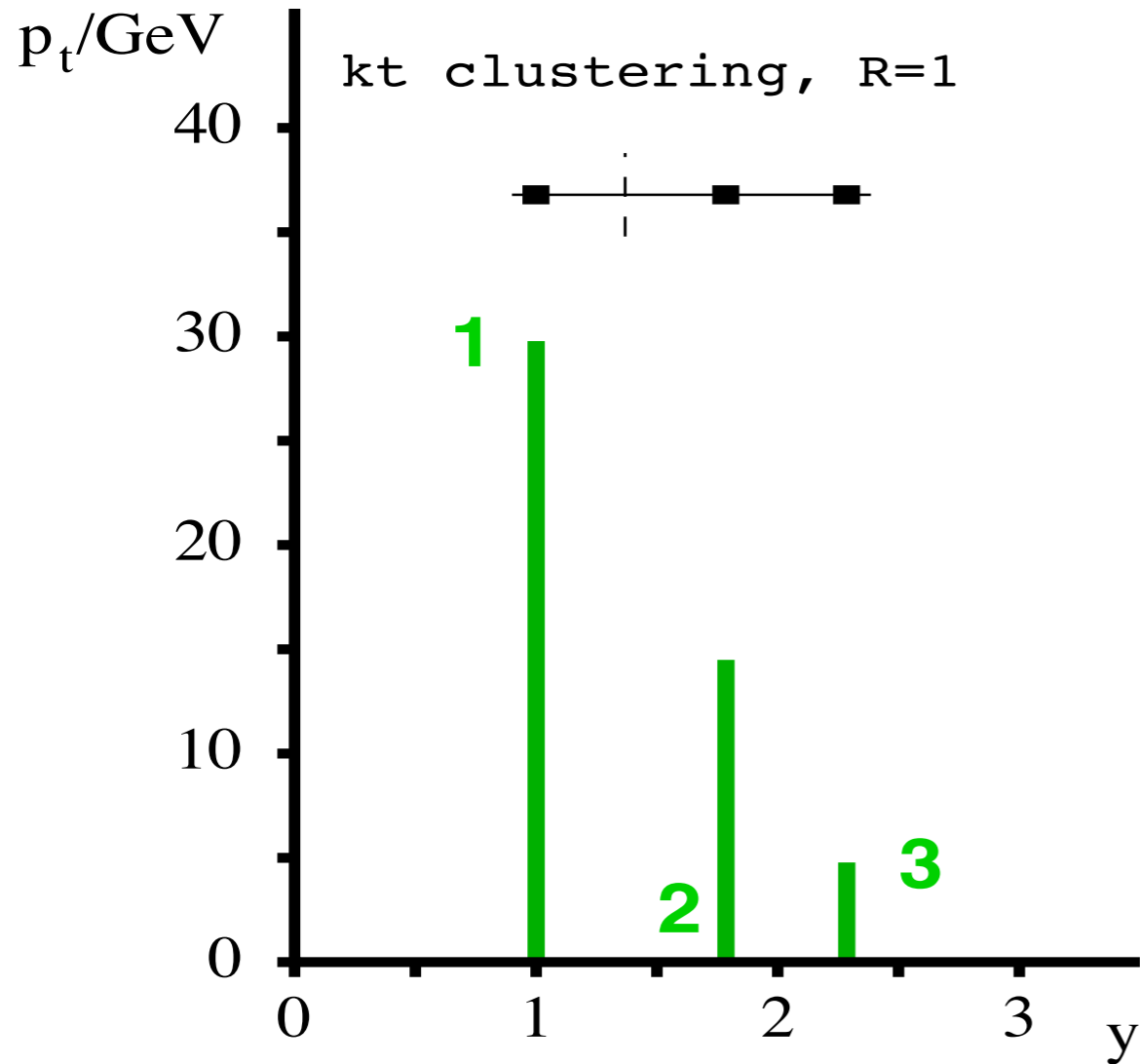
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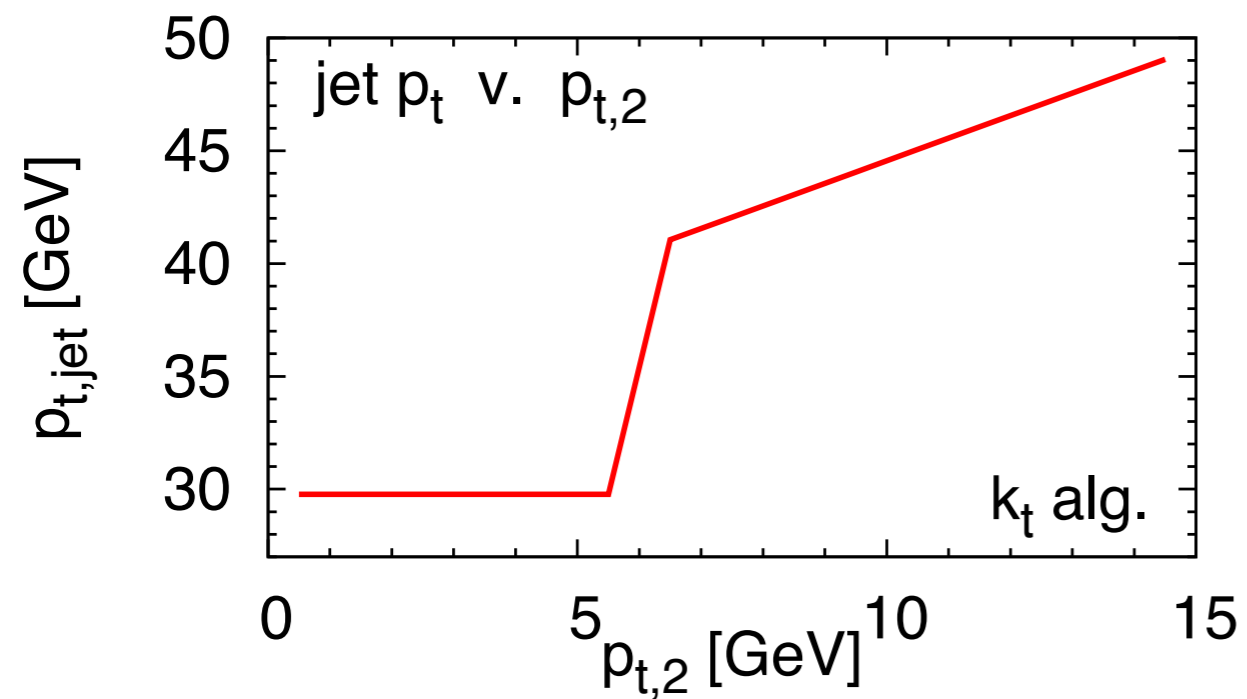
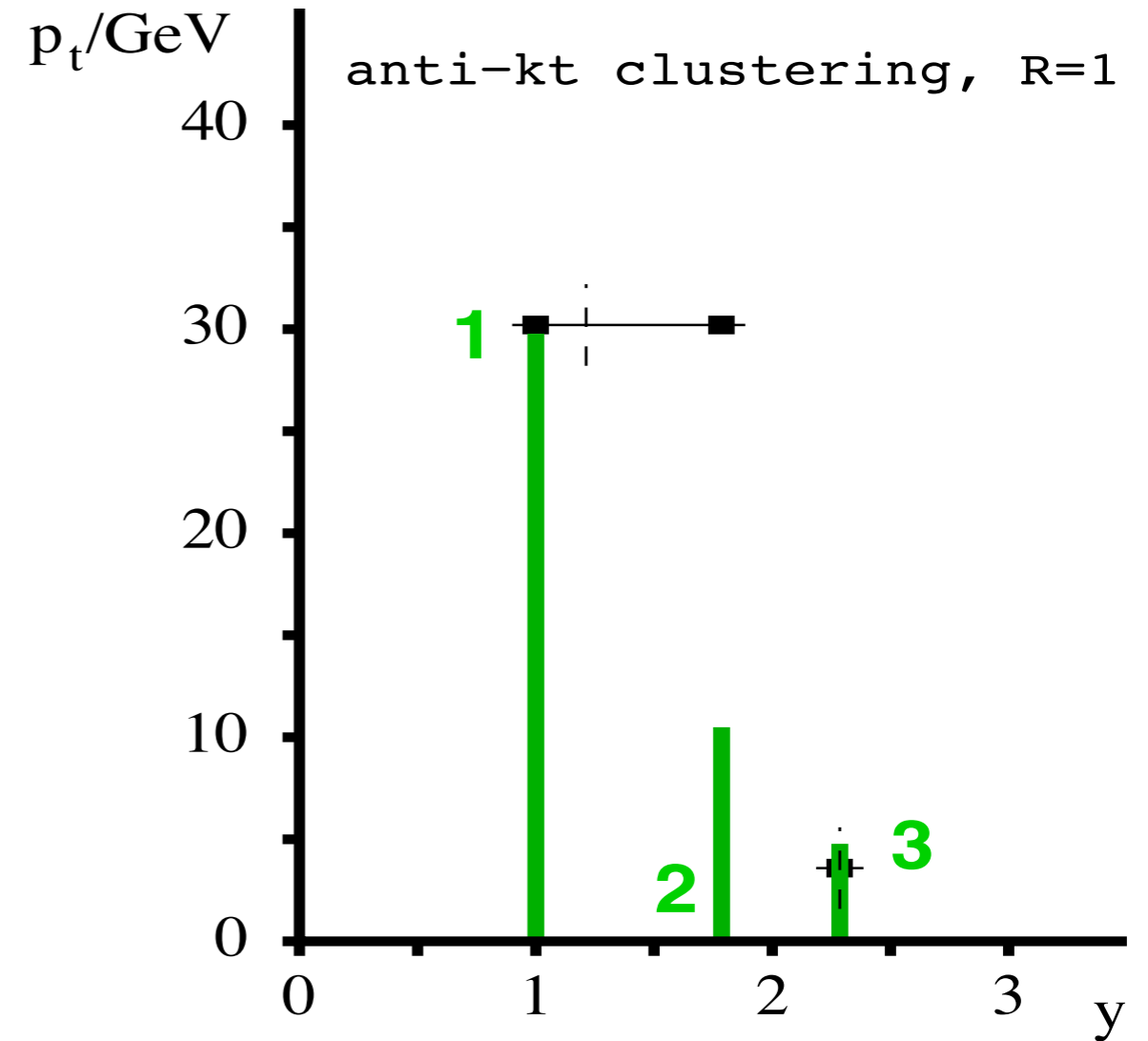
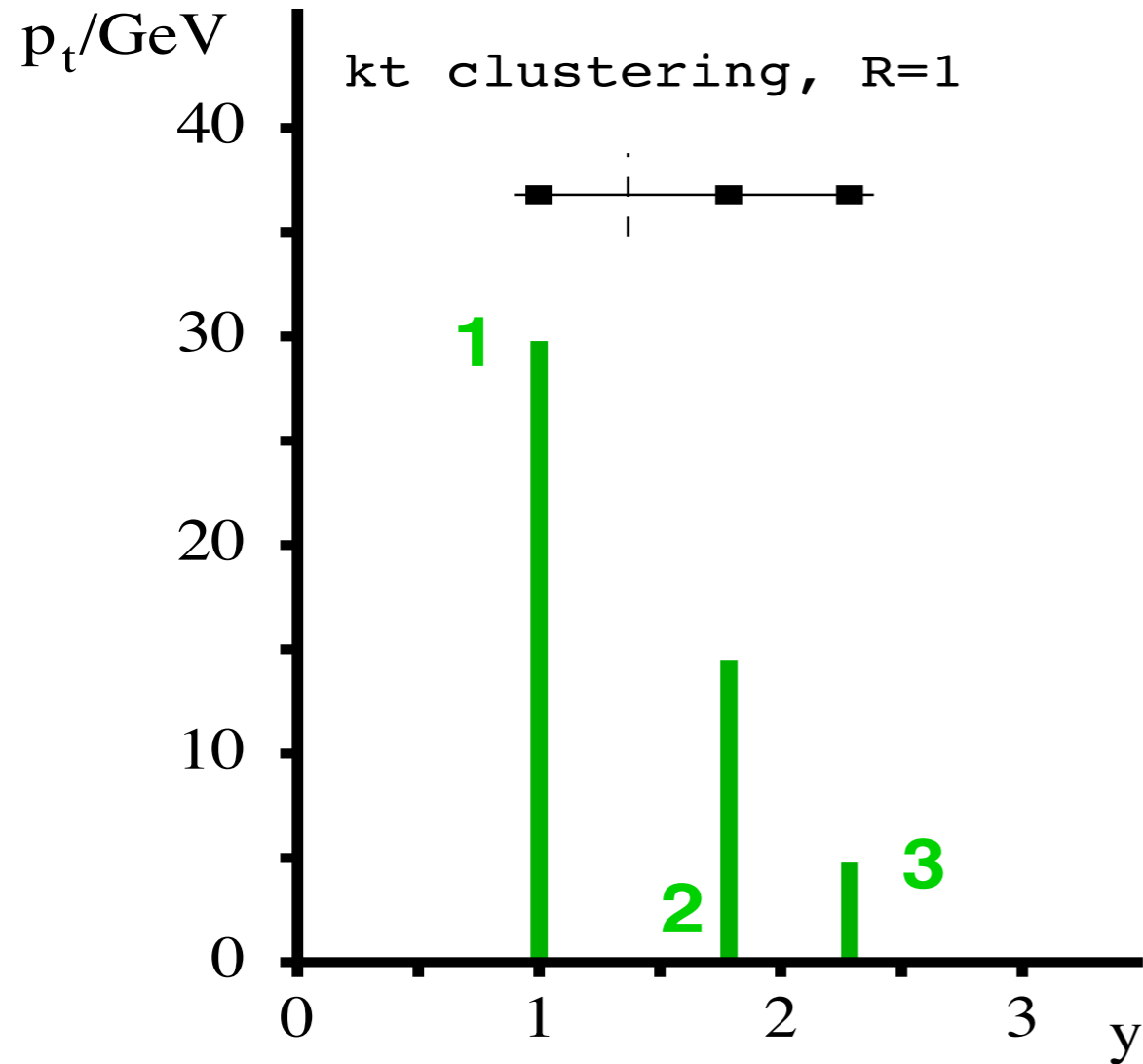
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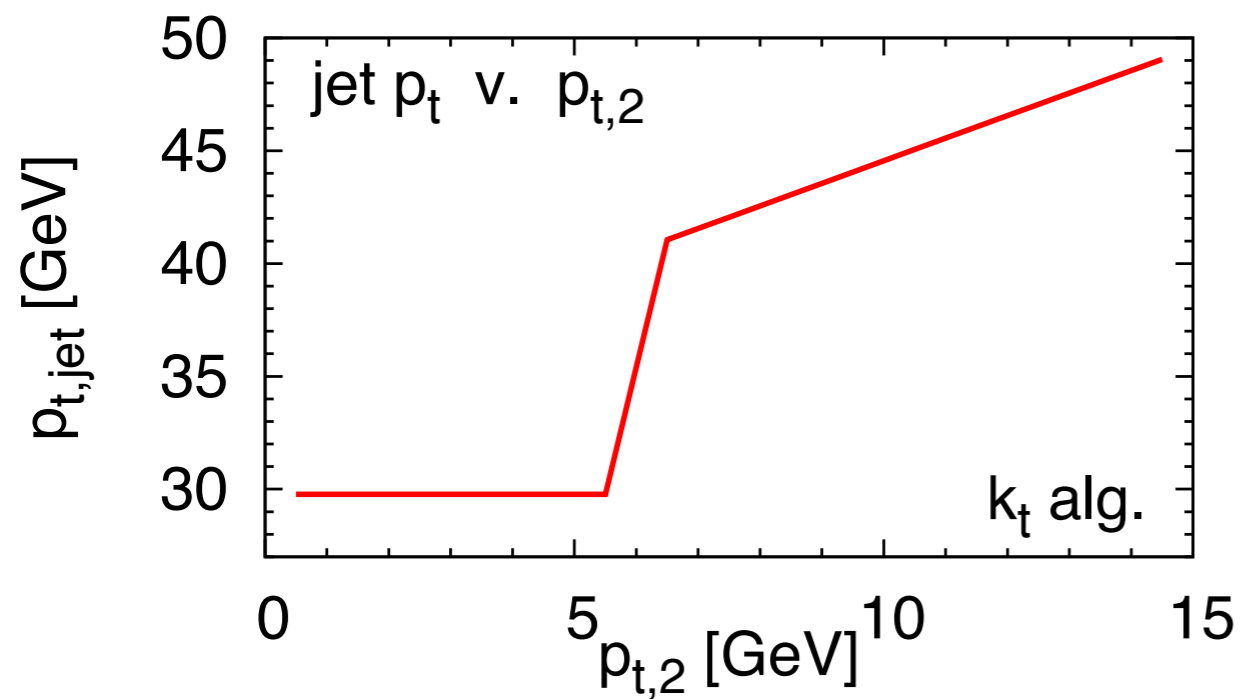
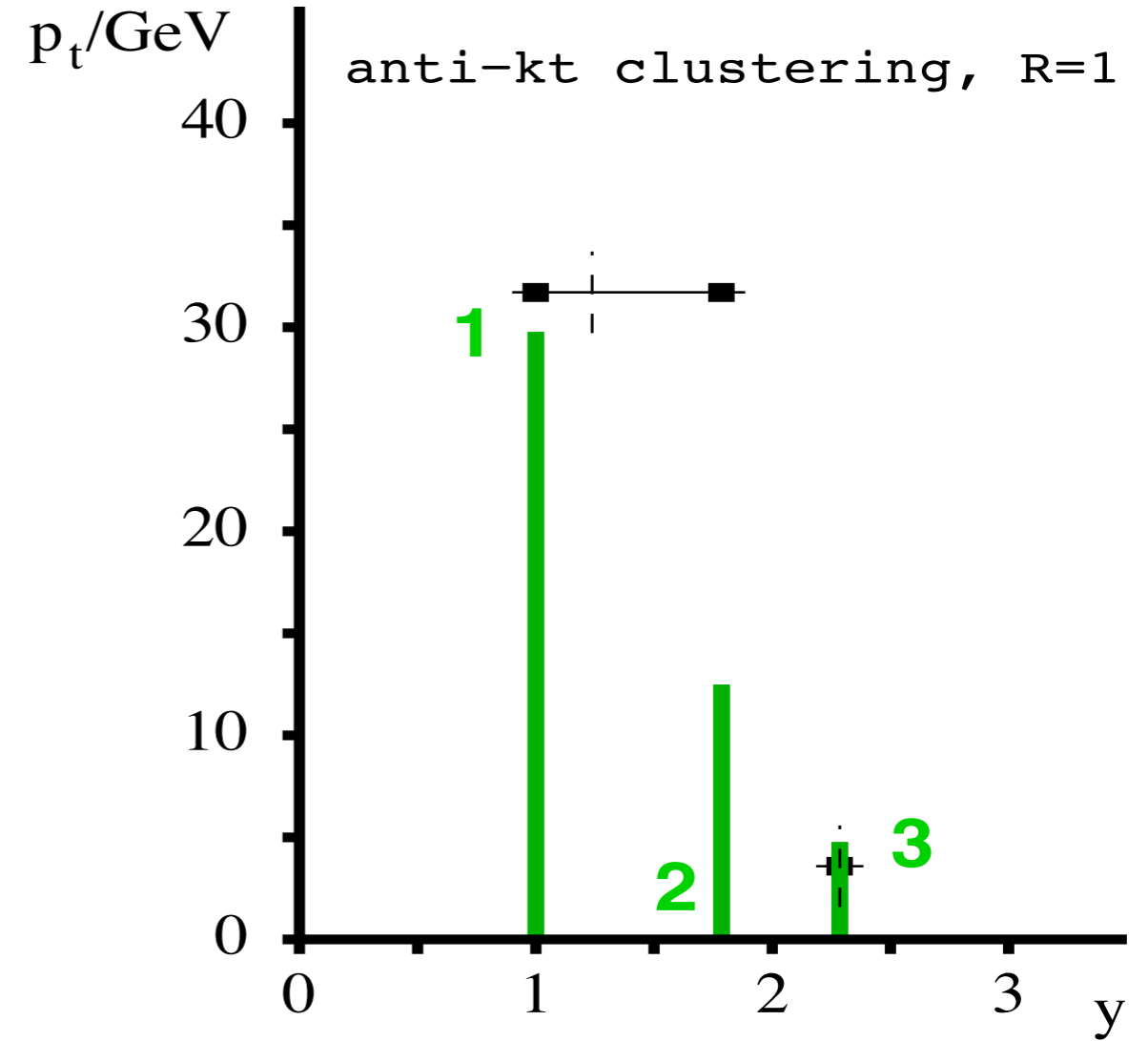
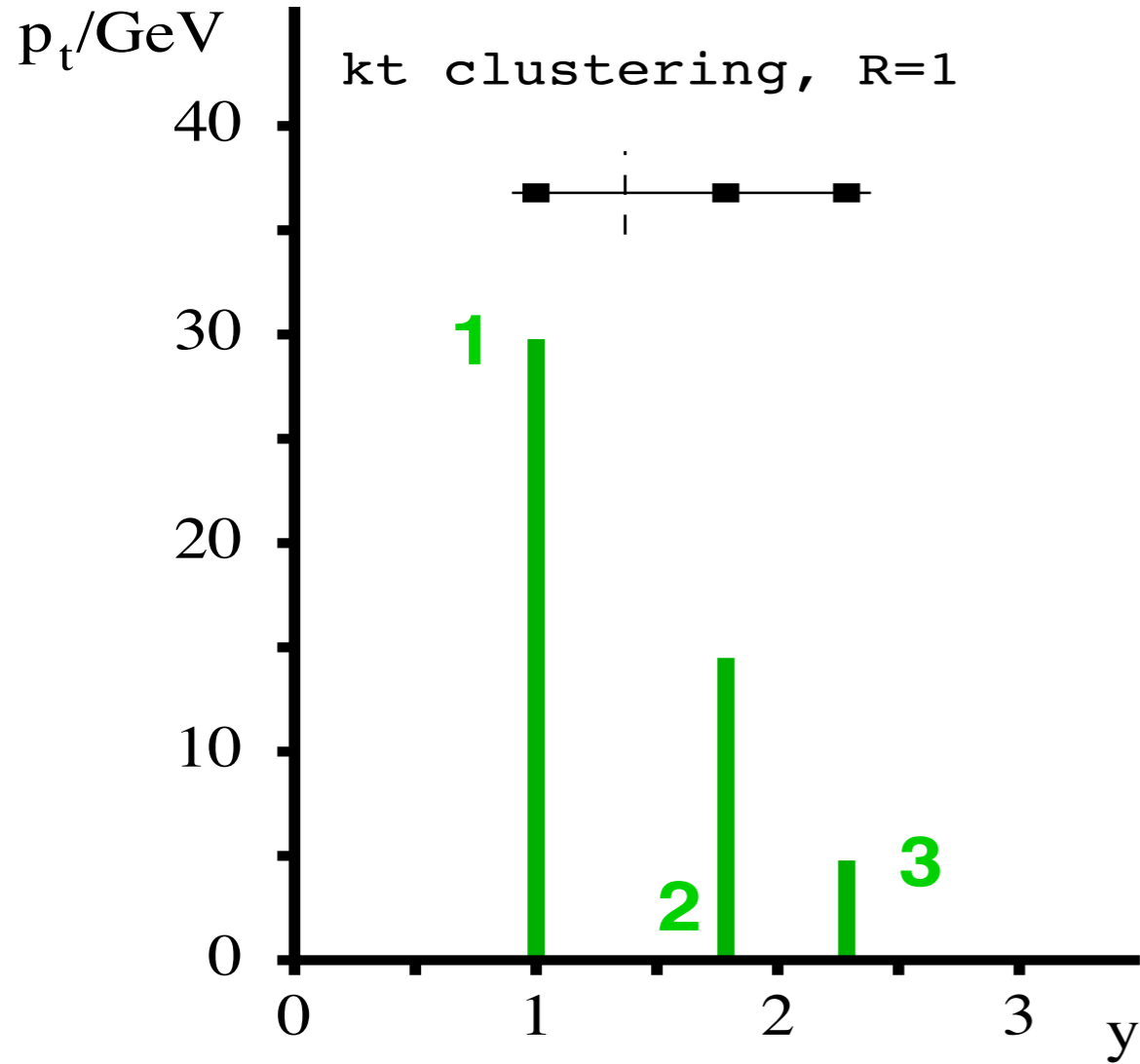
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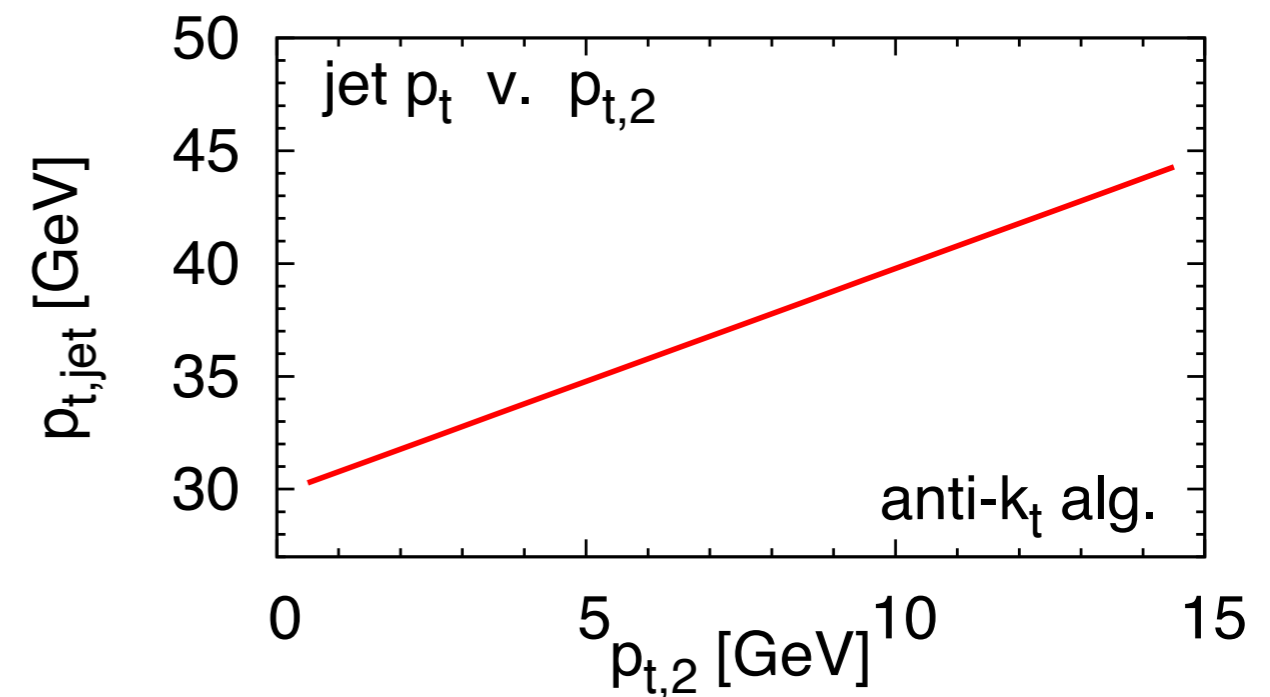
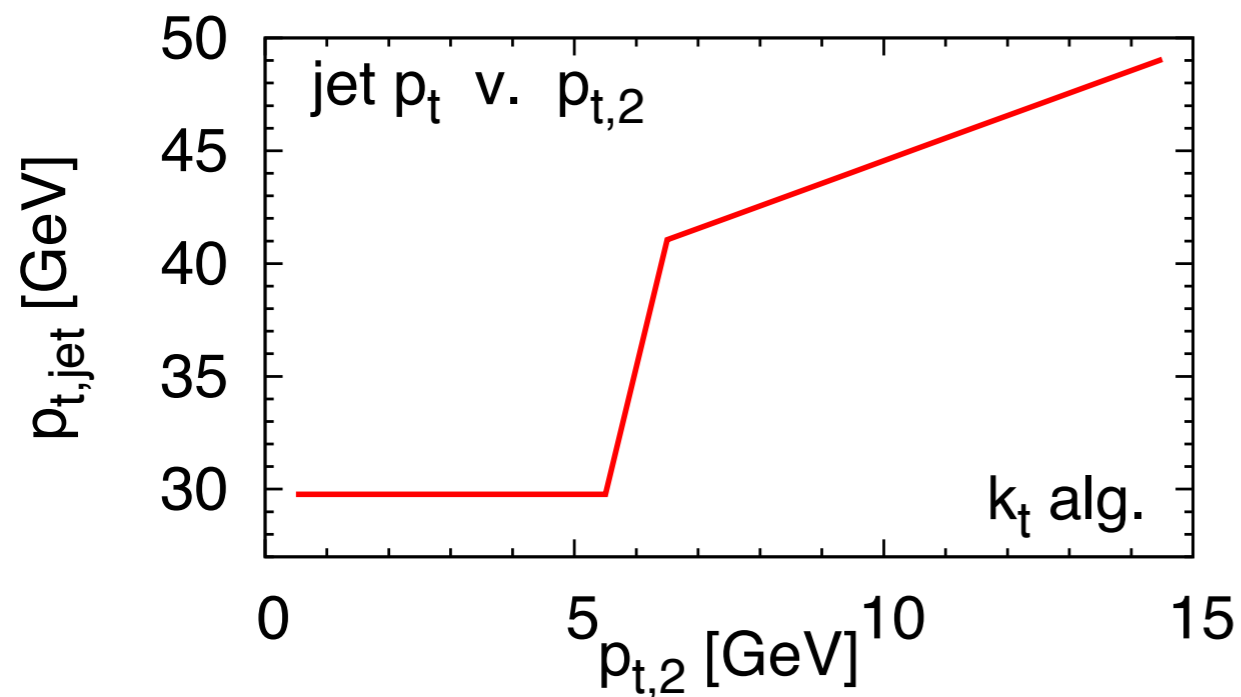
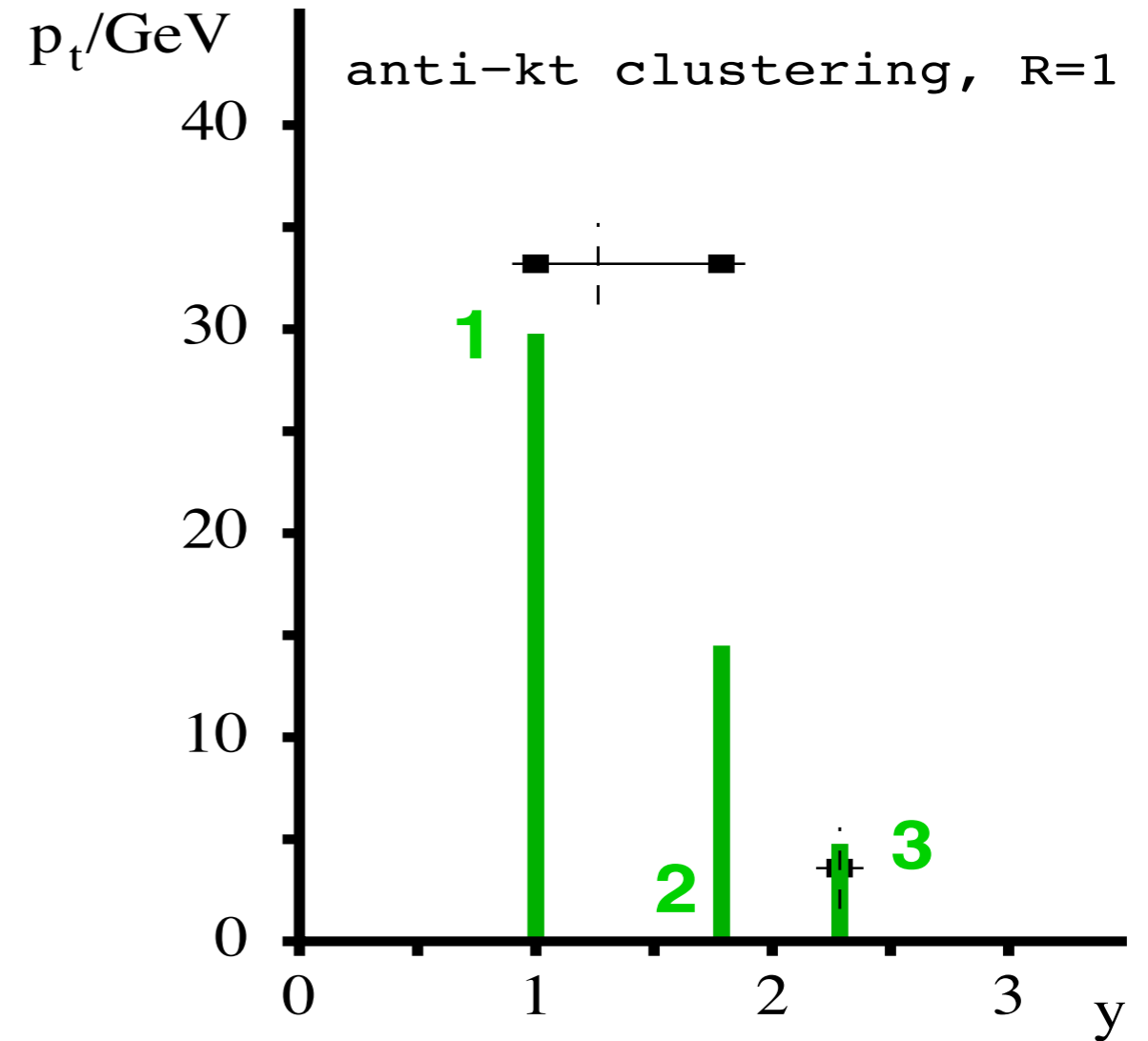
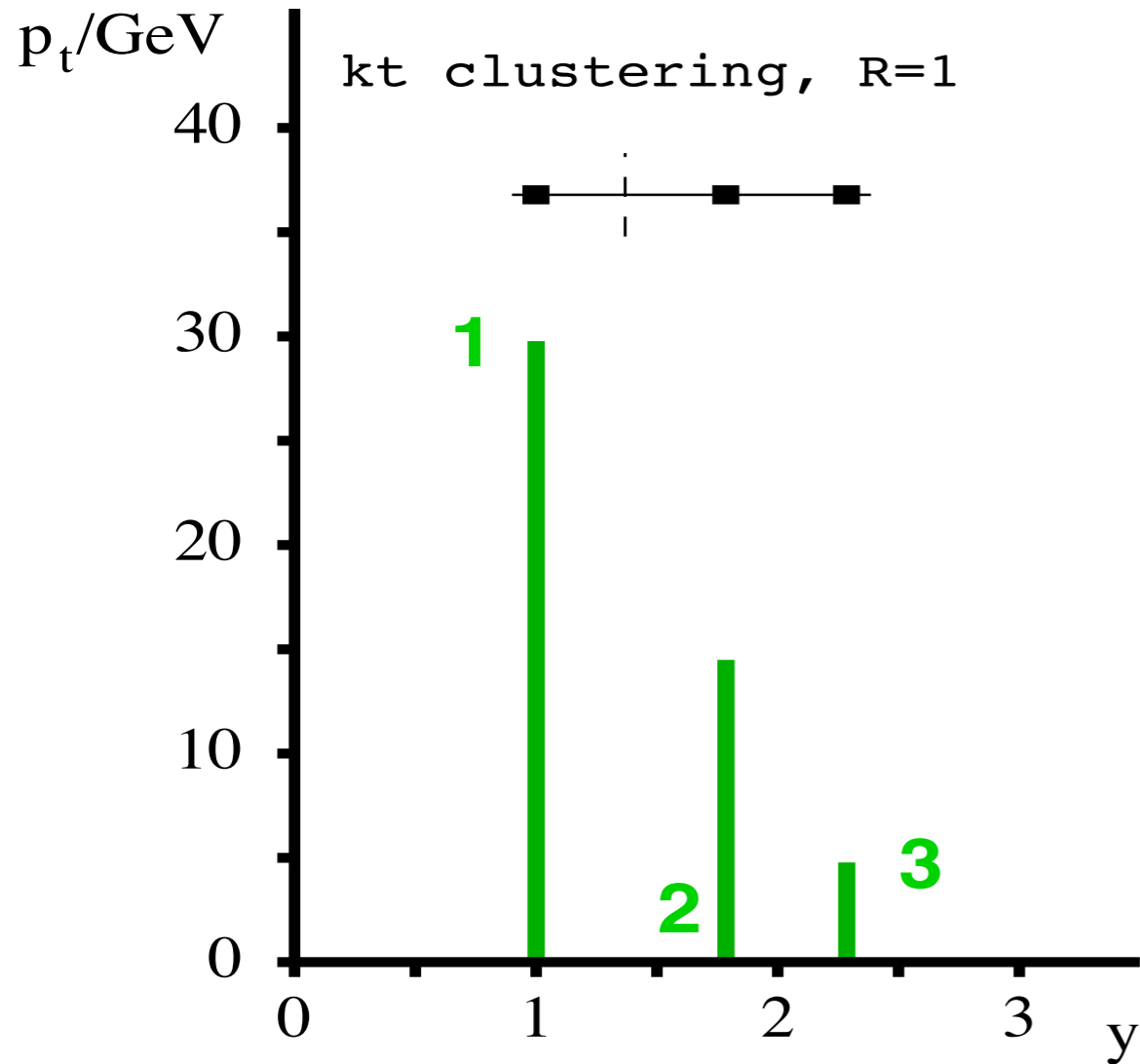
Linearity: k_t v. anti- k_t

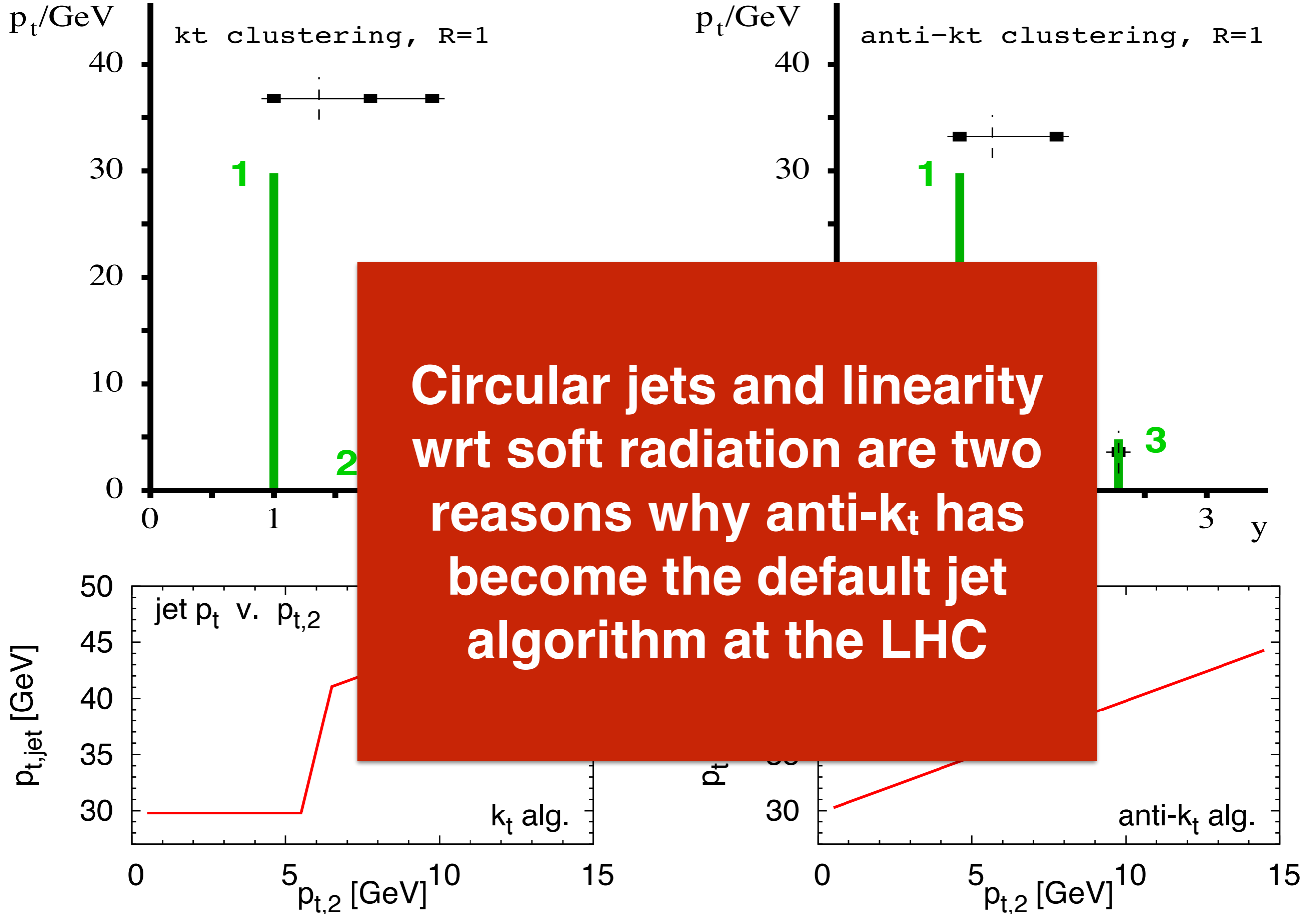


Linearity: k_t v. anti- k_t



Linearity: k_t v. anti- k_t





```
// specify a jet definition  
double R = 0.4  
JetDefinition jet_def(antikt_algorithm, R);
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jet_algorithm can be any one of the four IRC safe pp-collider algorithms, or also a variety of e⁺e⁻ algorithms, both native and plugins

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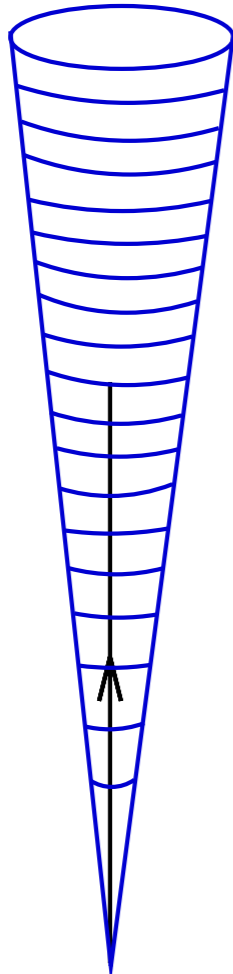
```
// extract the jets
vector<PseudoJet> jets = jet_def(input_particles);

// pt of hardest jet
double pt_hardest = jets[0].pt();

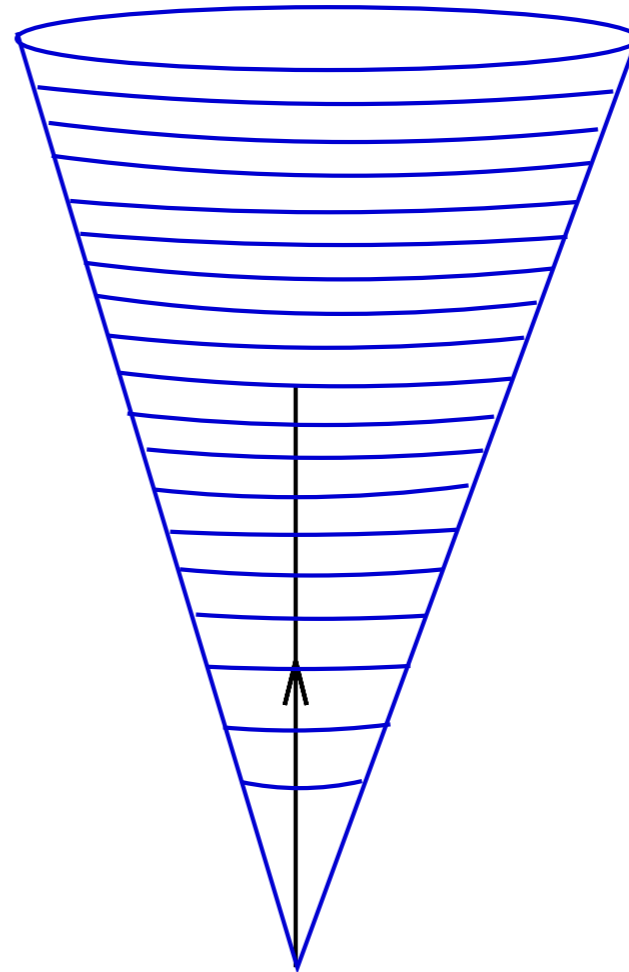
// constituents of hardest jet
vector<PseudoJet> constituents = jets[0].constituents();
```

hadron collider jet reconstruction parameters

Small jet radius

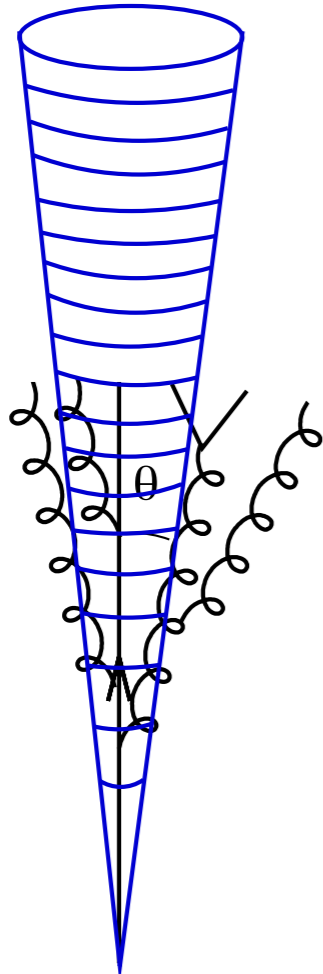


Large jet radius

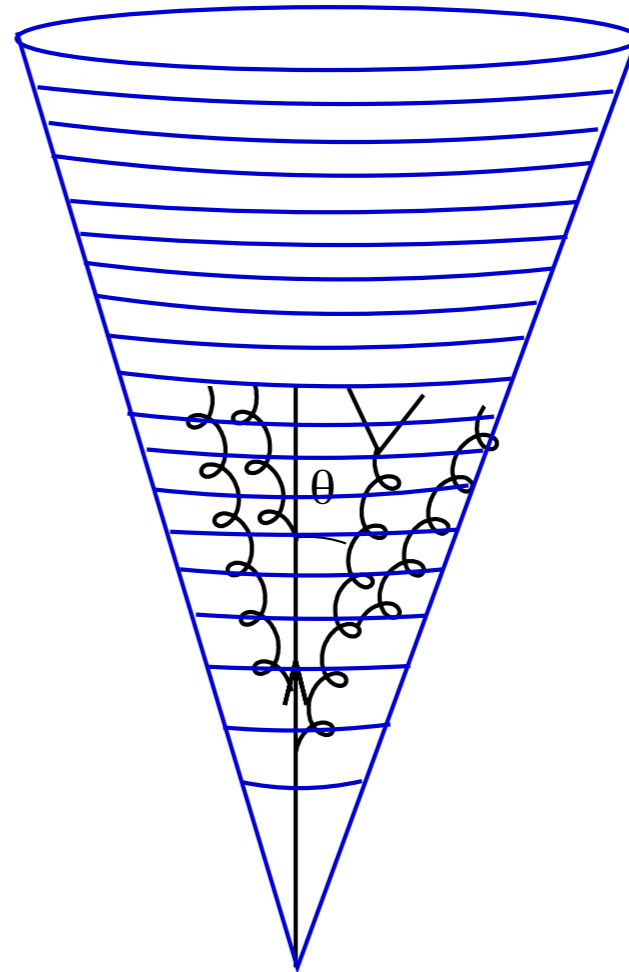


single parton @ LO: **jet radius irrelevant**

Small jet radius

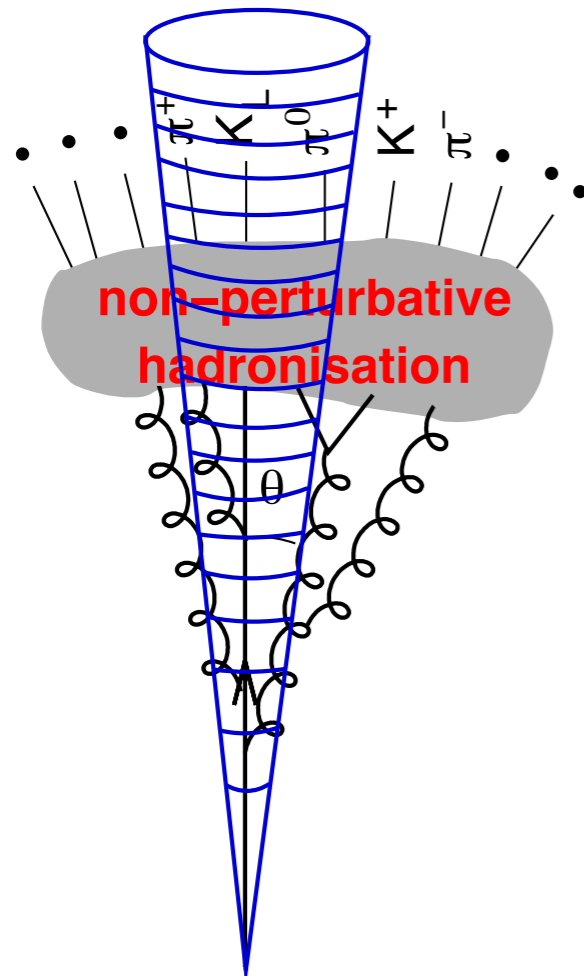


Large jet radius

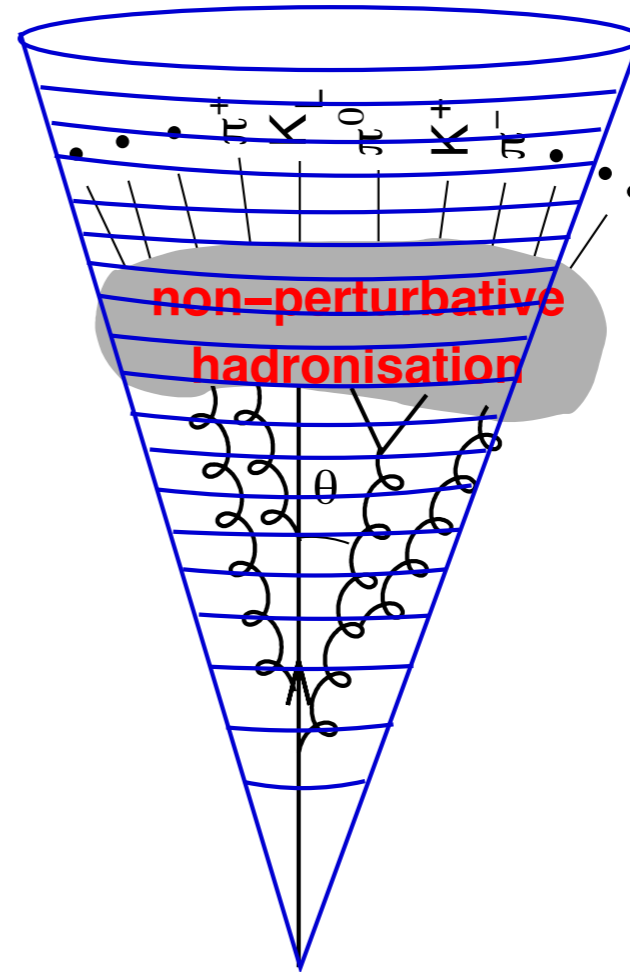


perturbative fragmentation: **large jet radius better**
(it captures more)

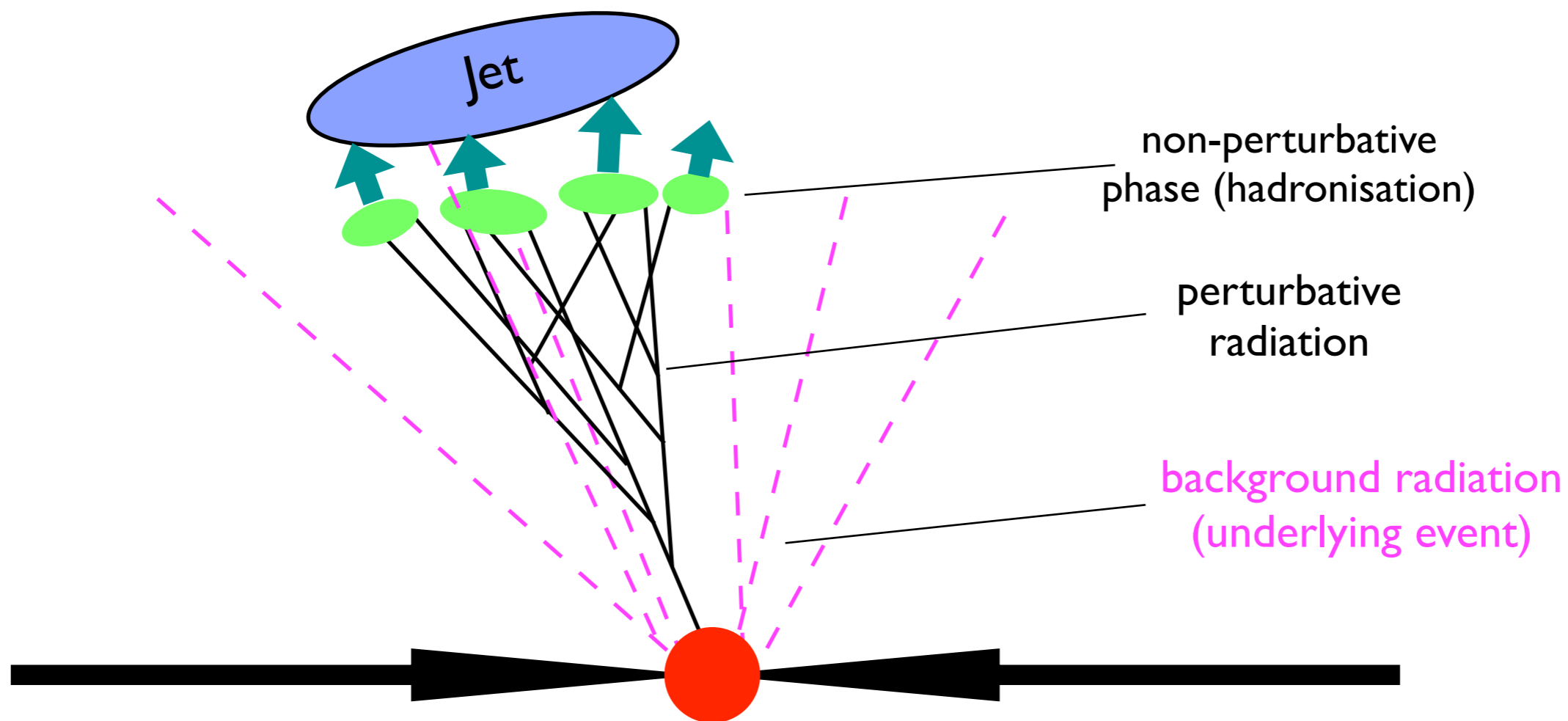
Small jet radius

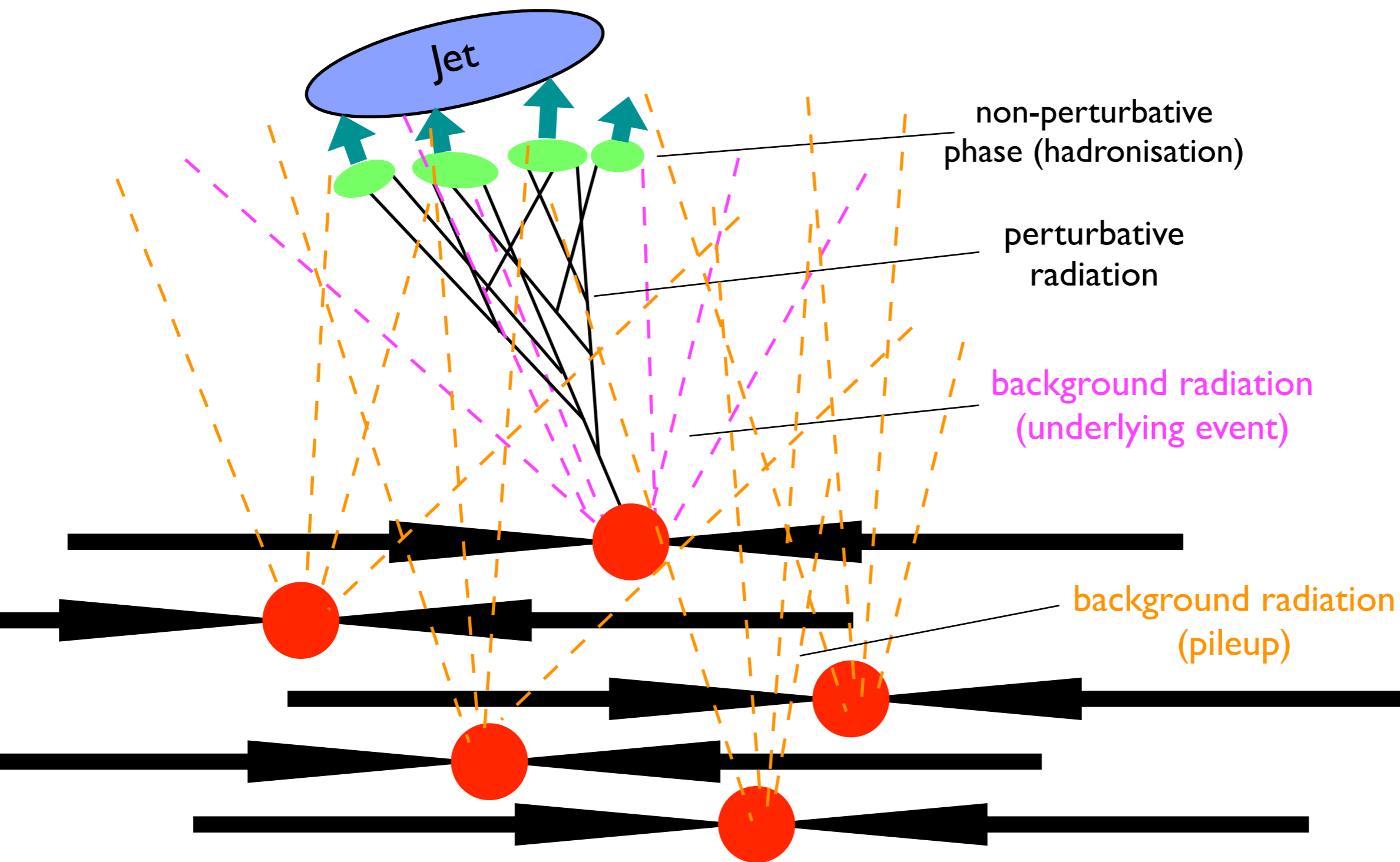


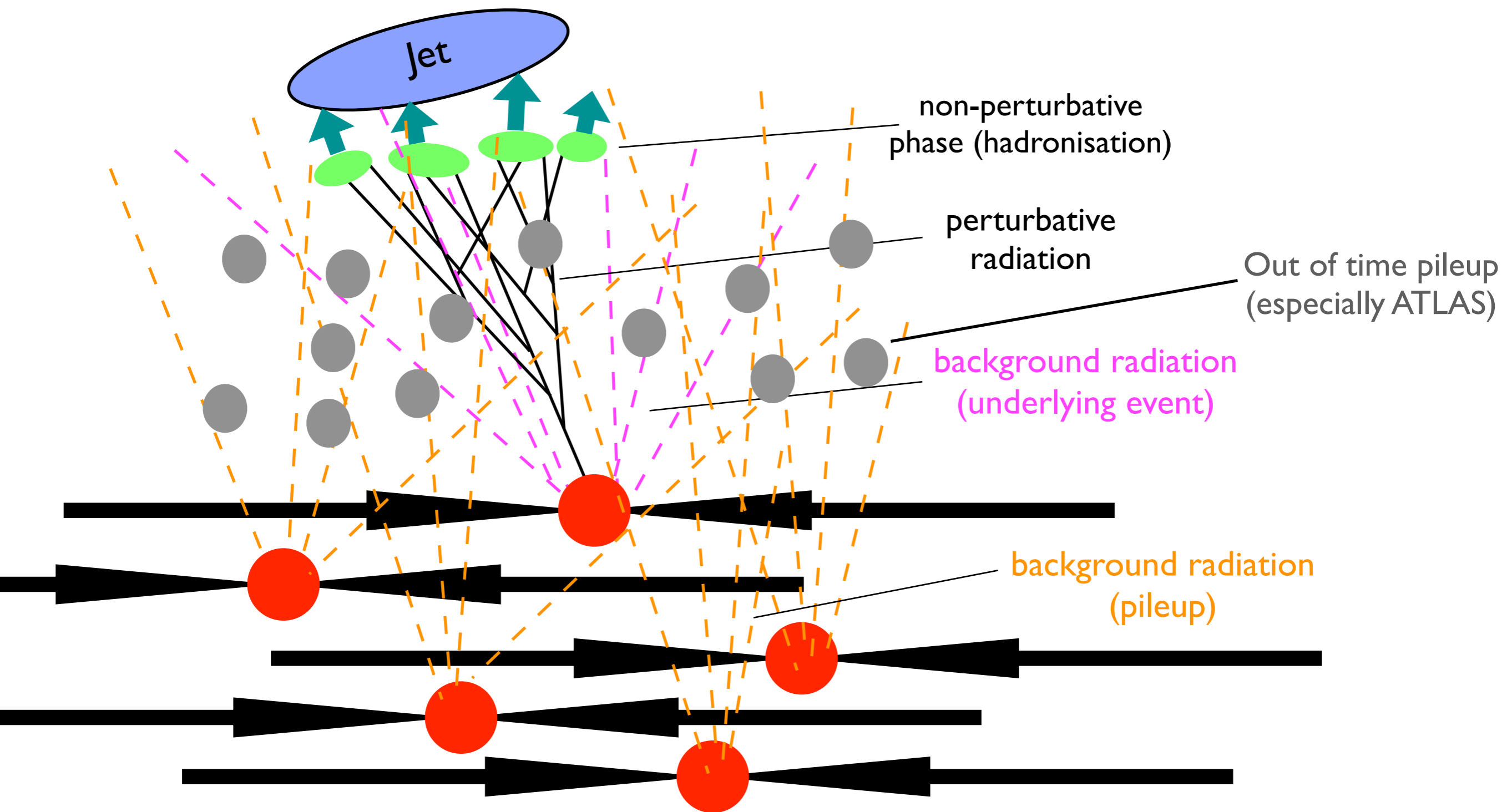
Large jet radius



non-perturbative fragmentation: **large jet radius better**
(it captures more)



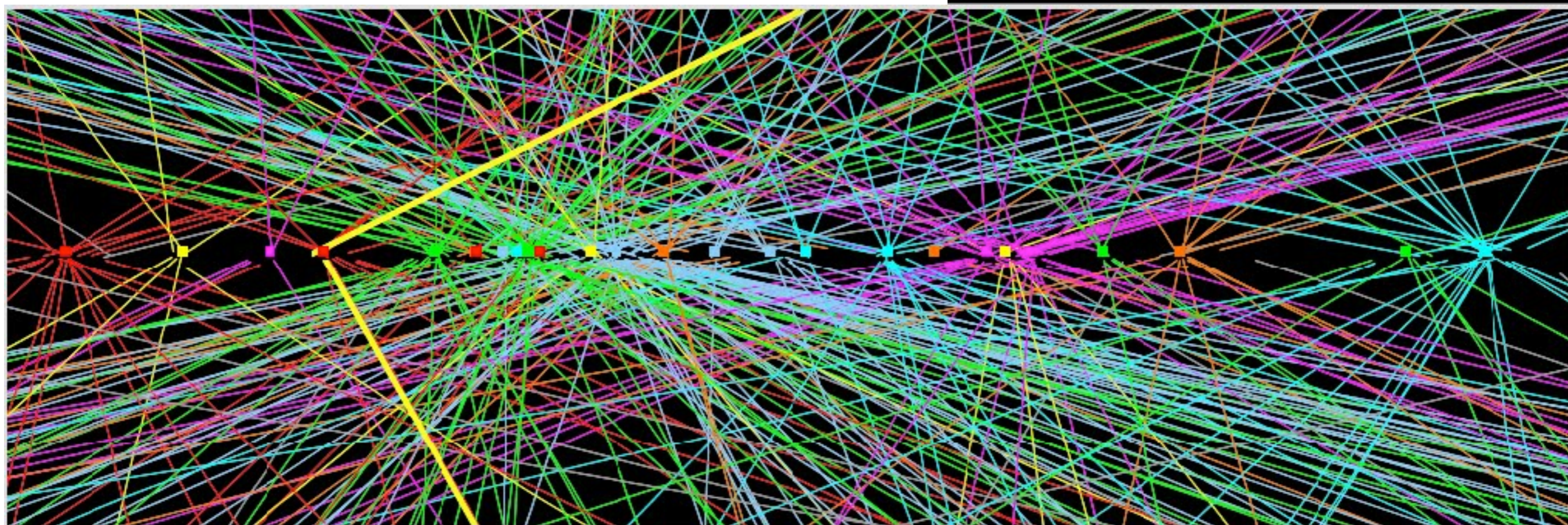
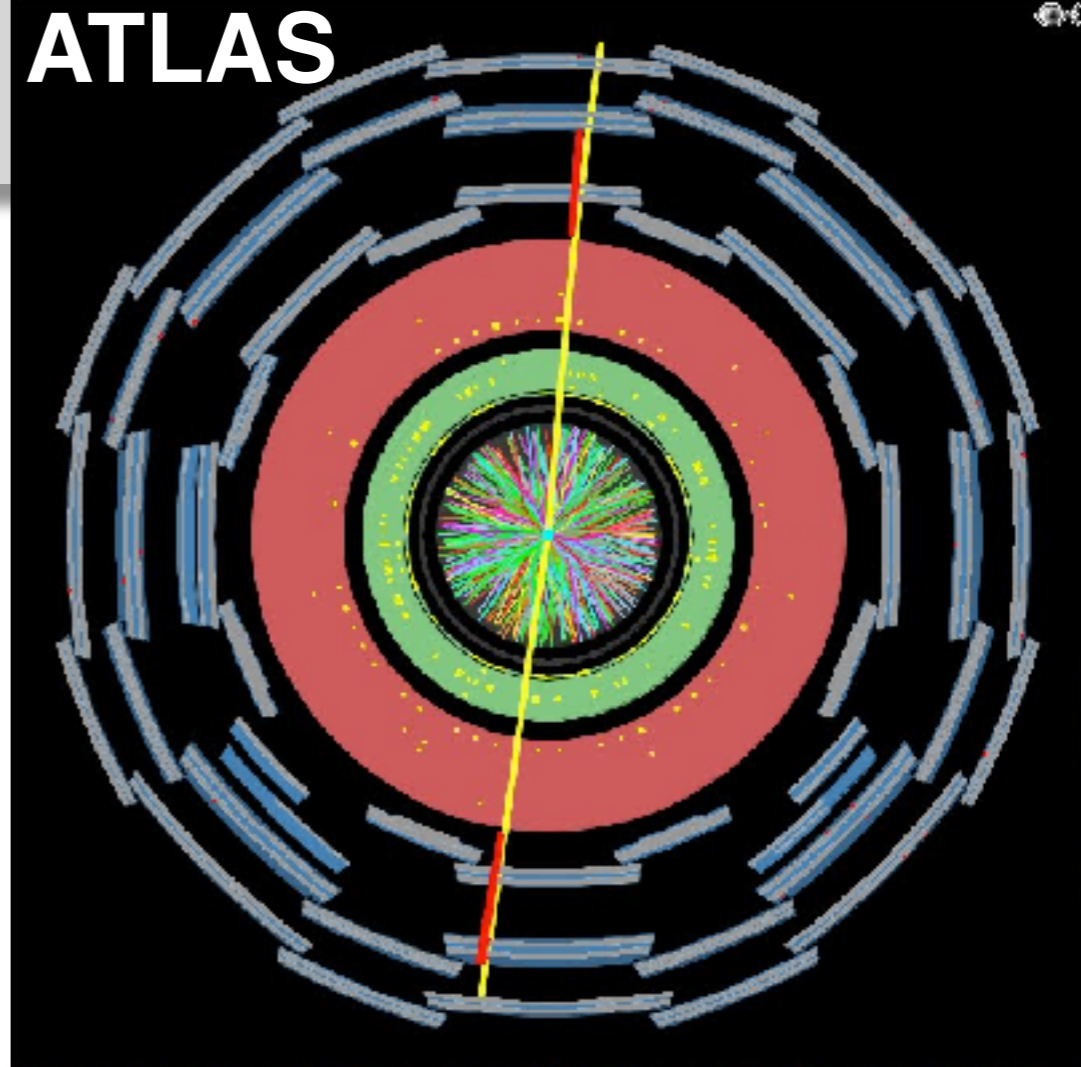




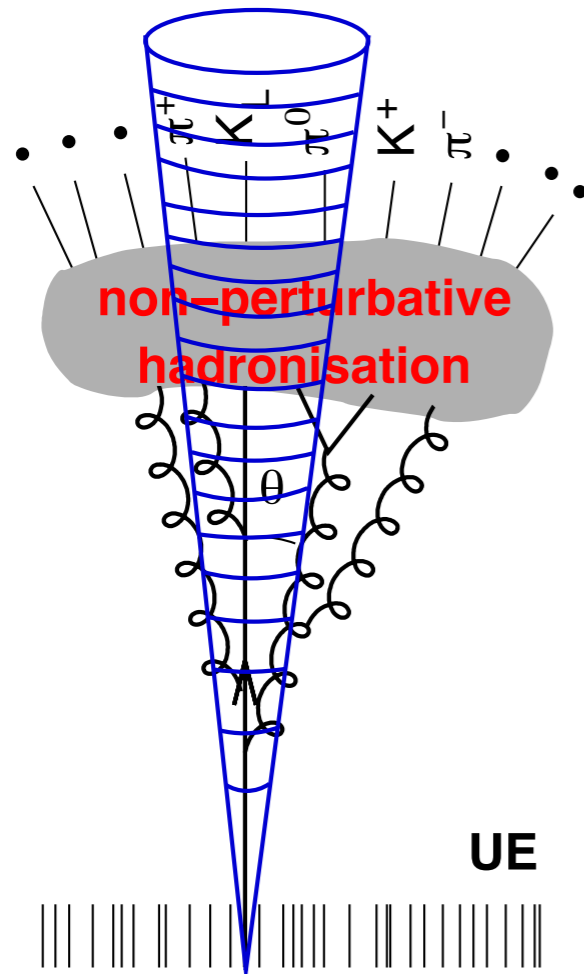
Pileup for real

a few cm

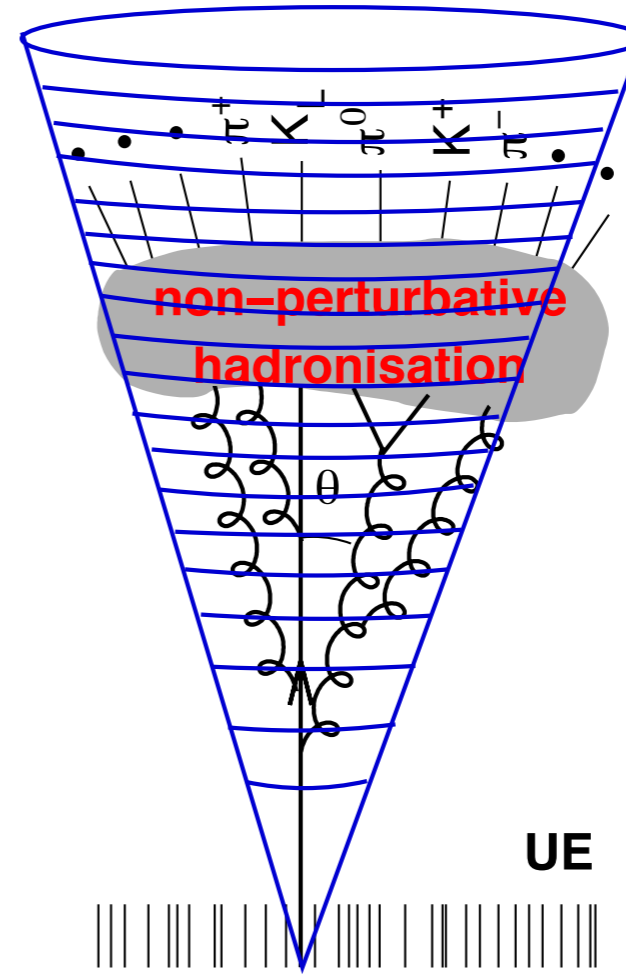
~ 20 m



Small jet radius

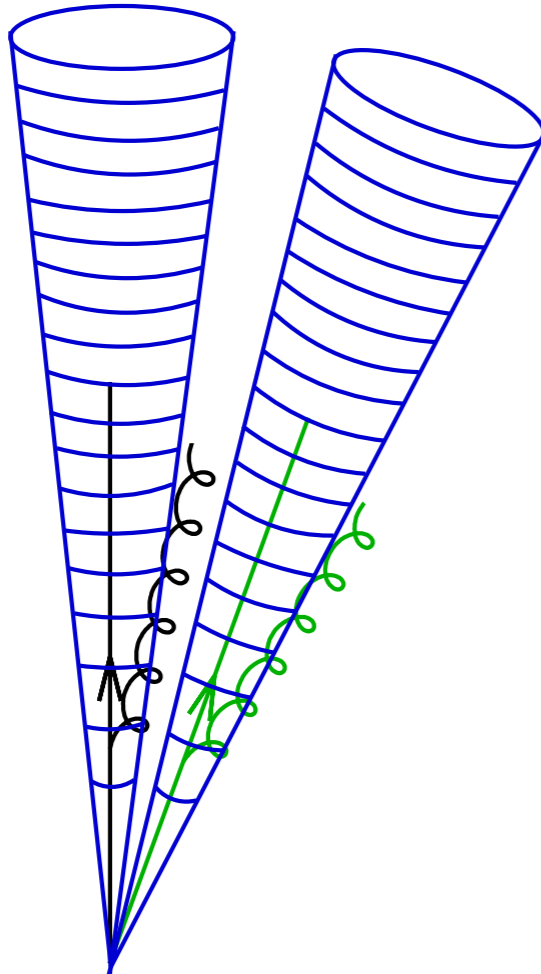


Large jet radius

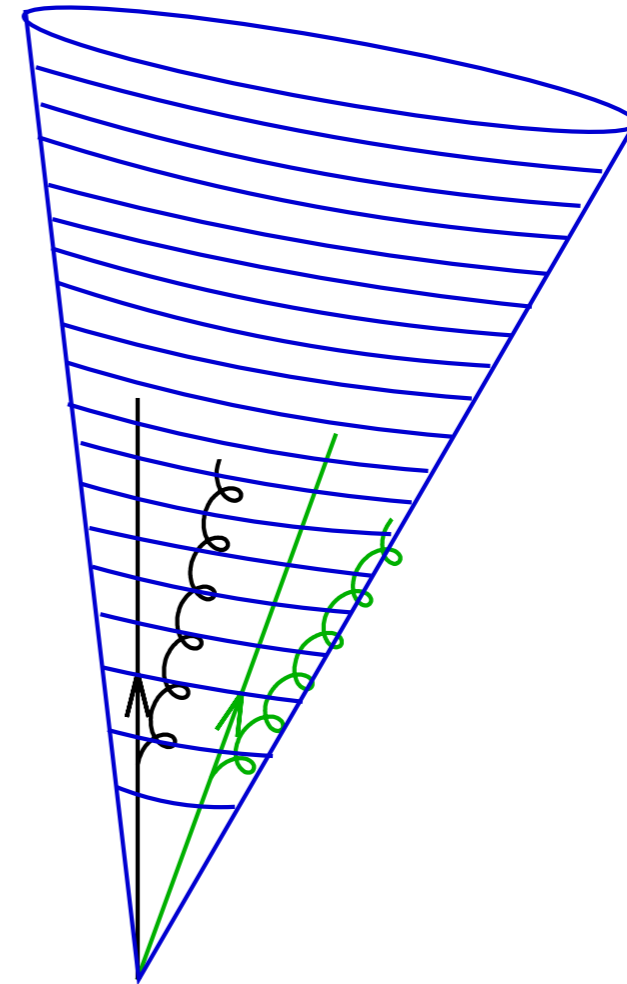


underlying ev. & pileup “noise”: **small jet radius better**
(it captures less)

Small jet radius



Large jet radius



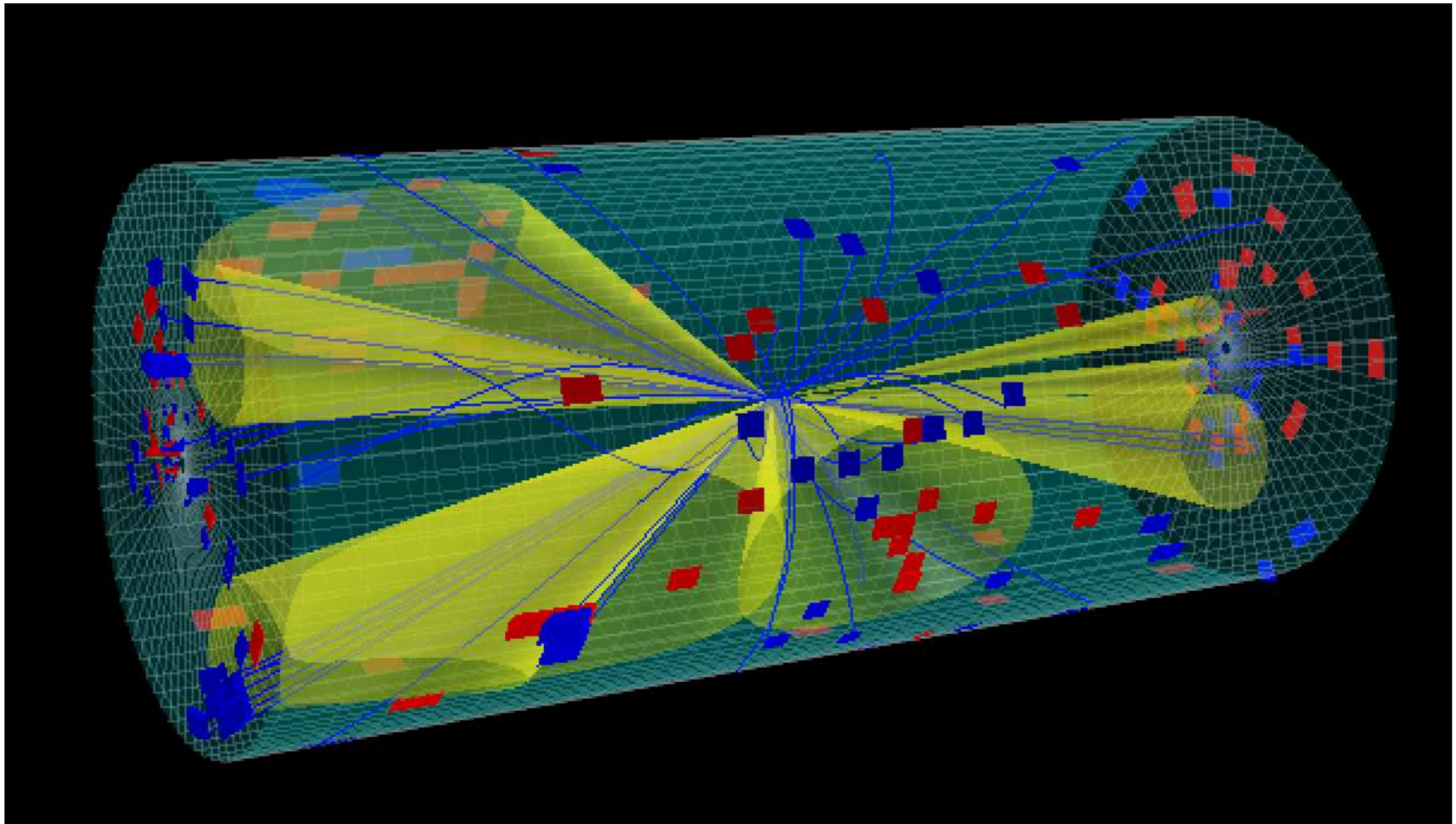
multi-hard-parton events: **small jet radius better**
(it resolves partons more effectively)

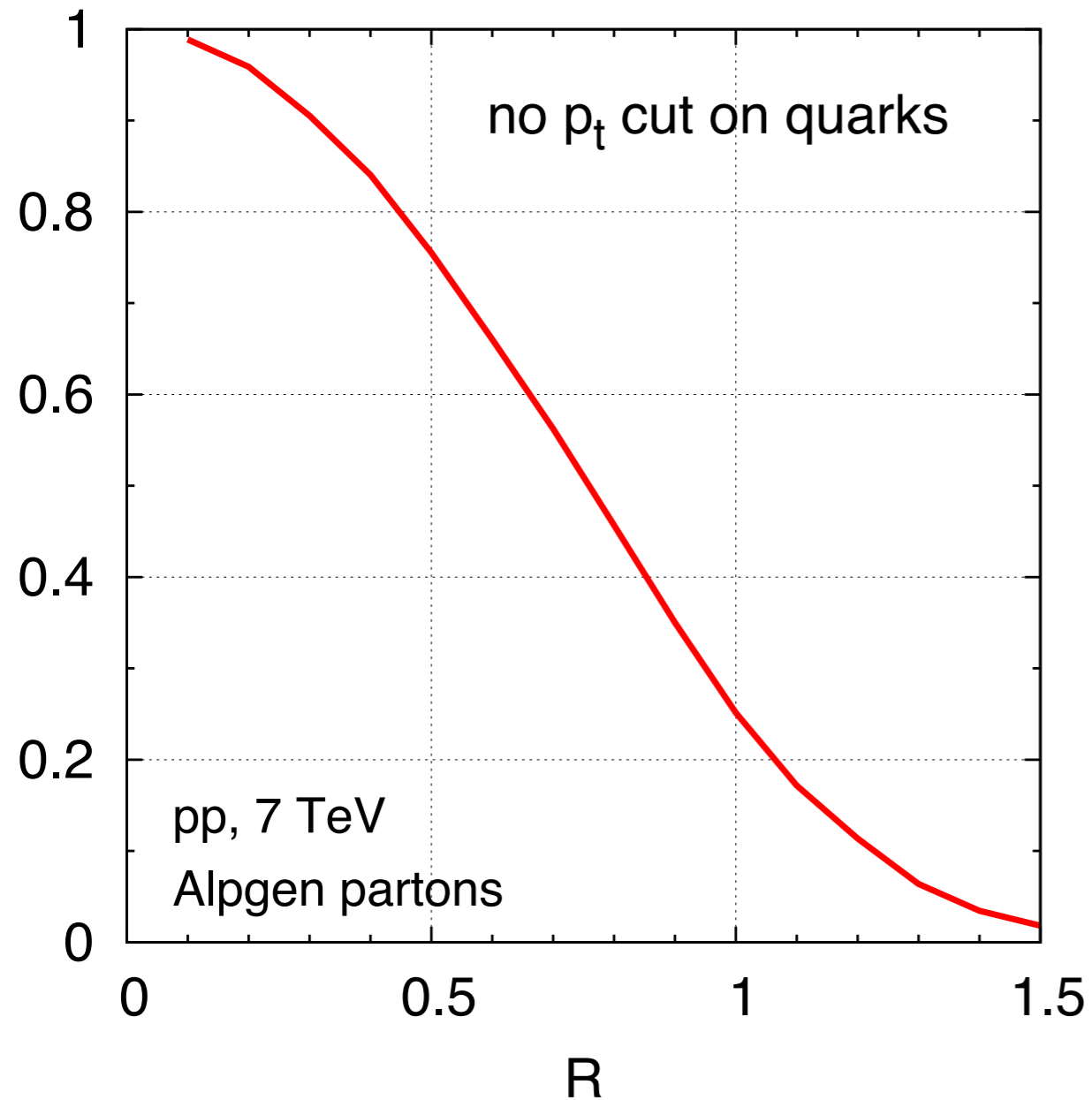
Can we capture all quarks and gluons?

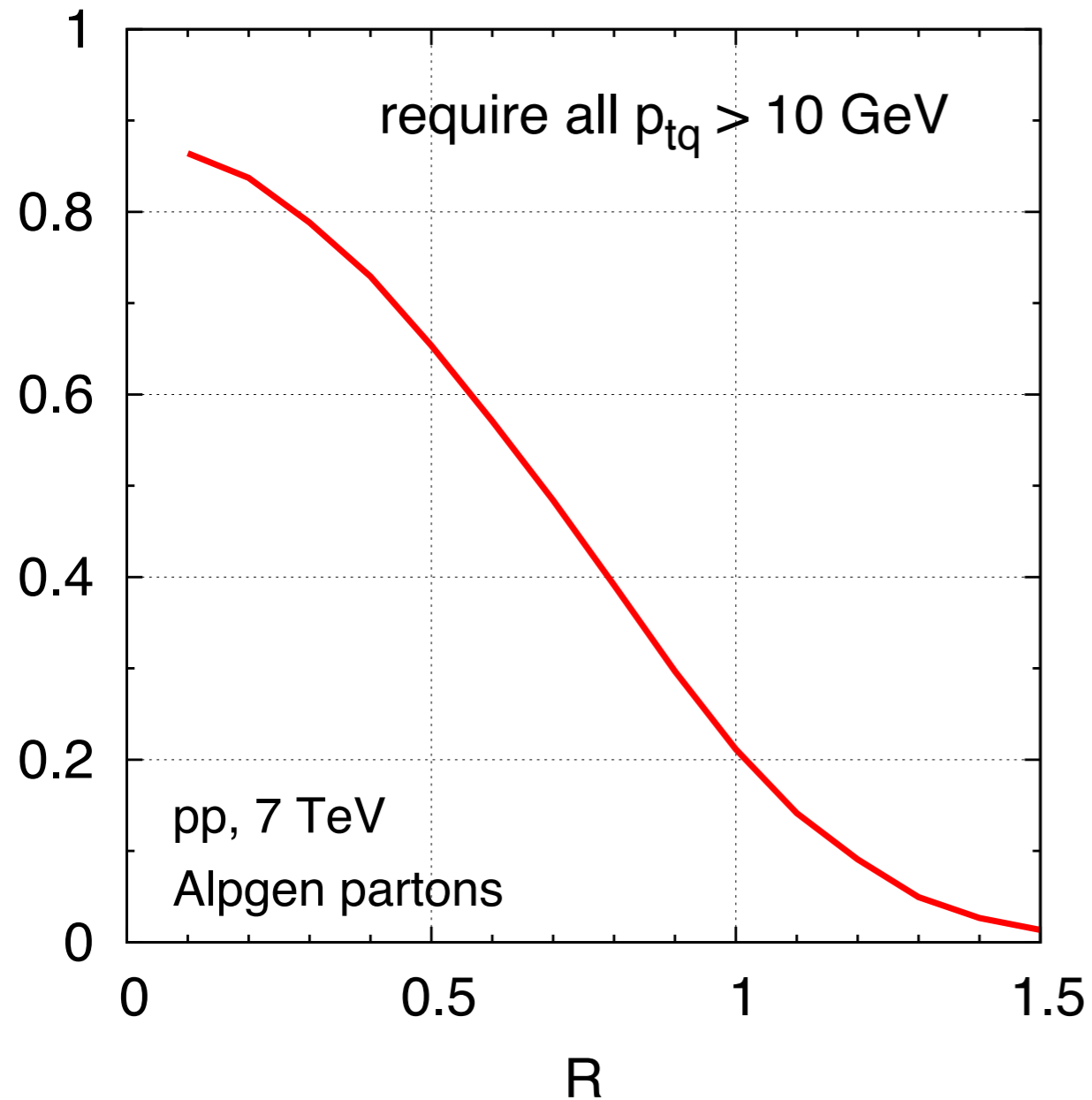
Should we capture all quarks and gluons?

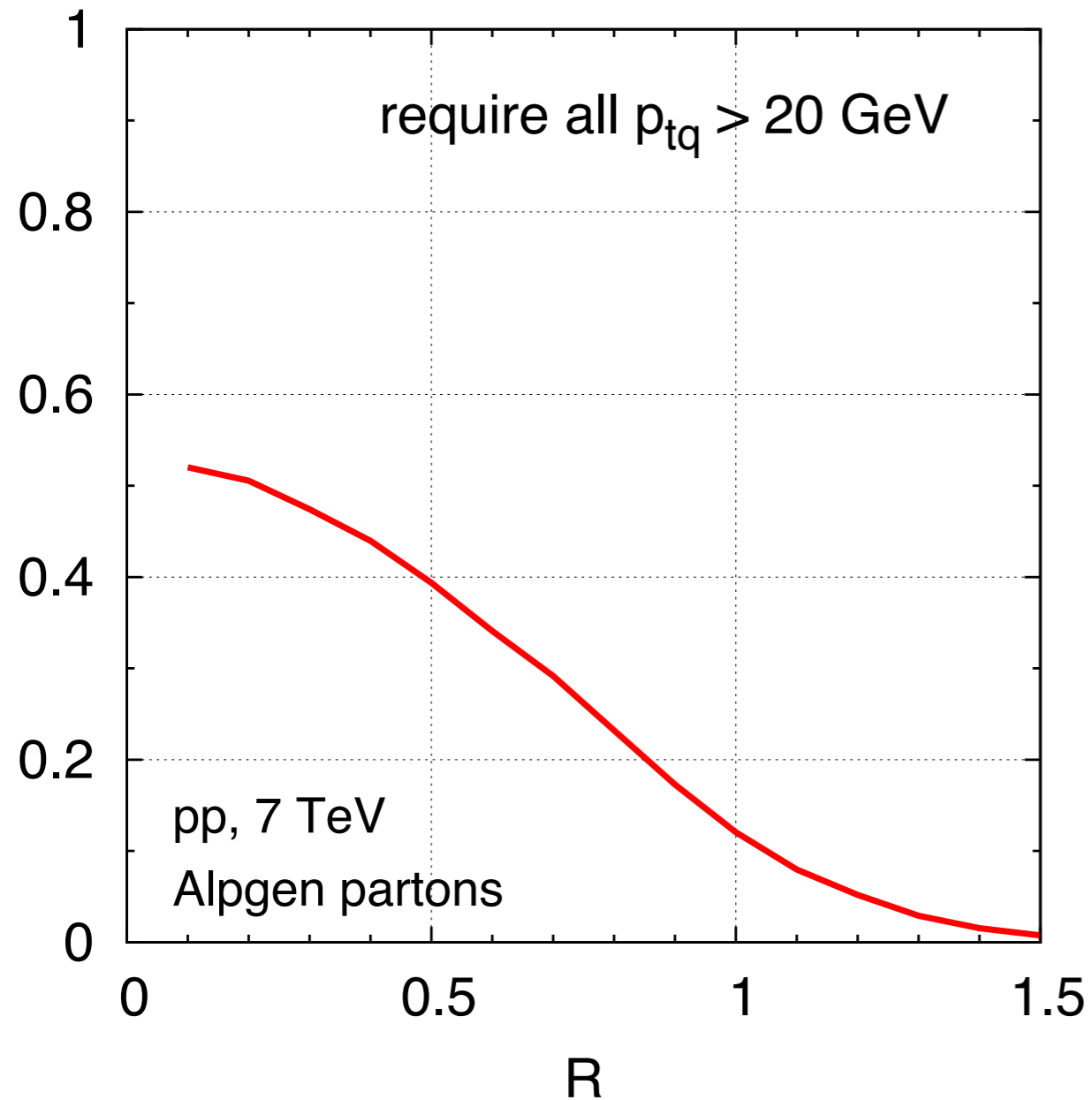
$$pp \rightarrow t\bar{t}$$

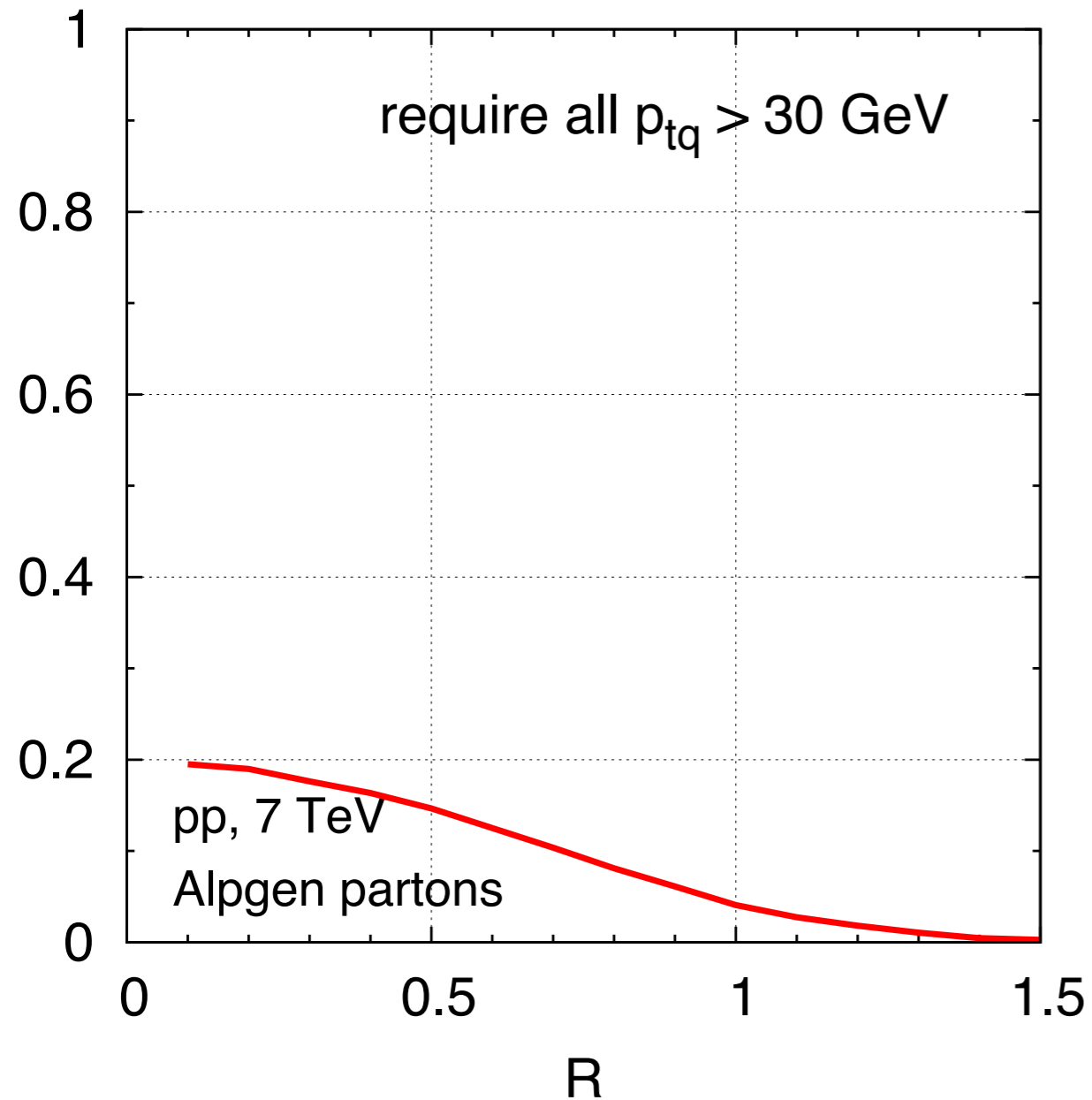
simulated with Pythia, displayed with Delphes

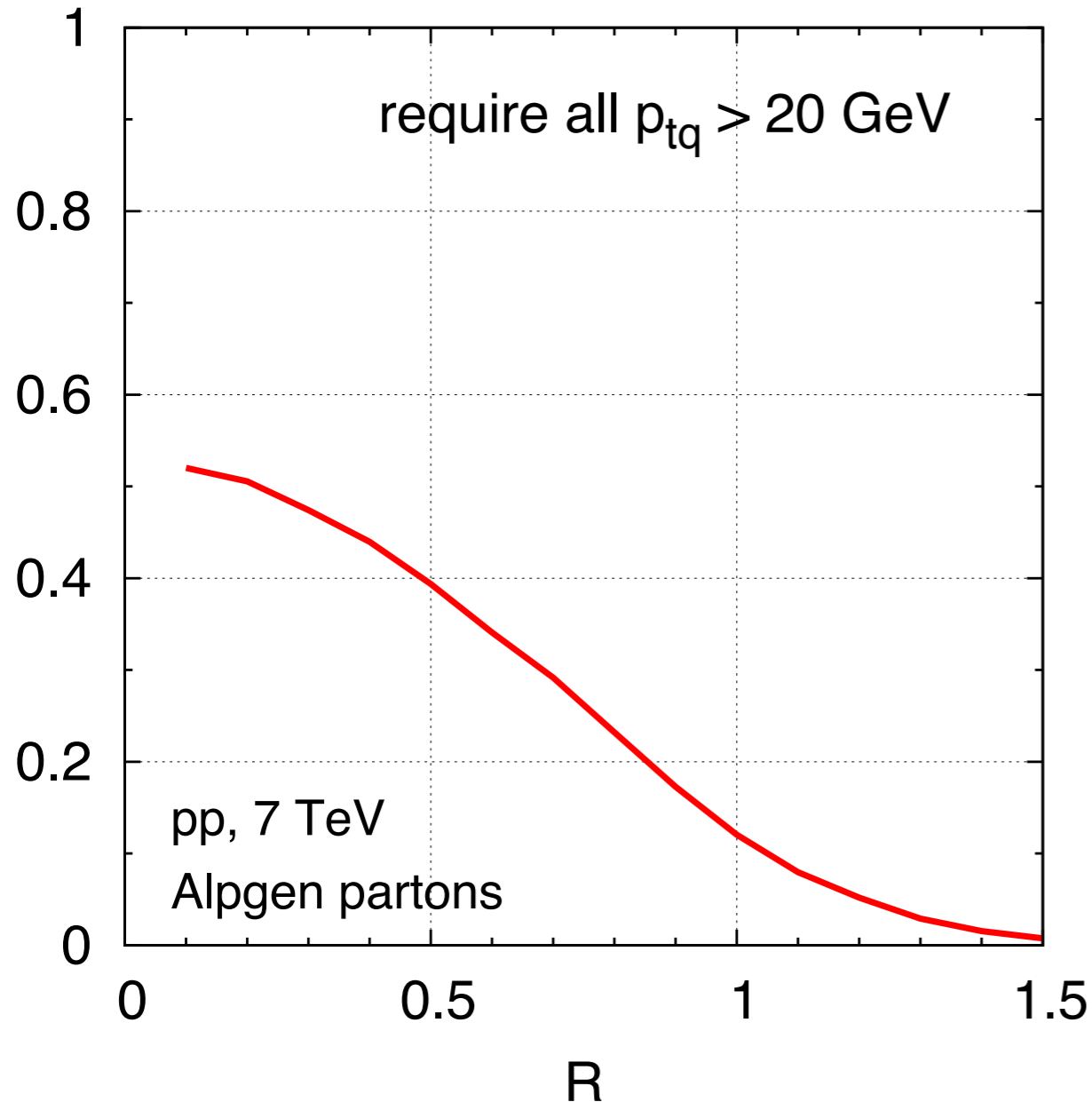


Alpgen $pp \rightarrow t\bar{t} \rightarrow 6q$ fraction of $pp \rightarrow t\bar{t} \rightarrow 6q$ events with all $R_{qq} > R$ 

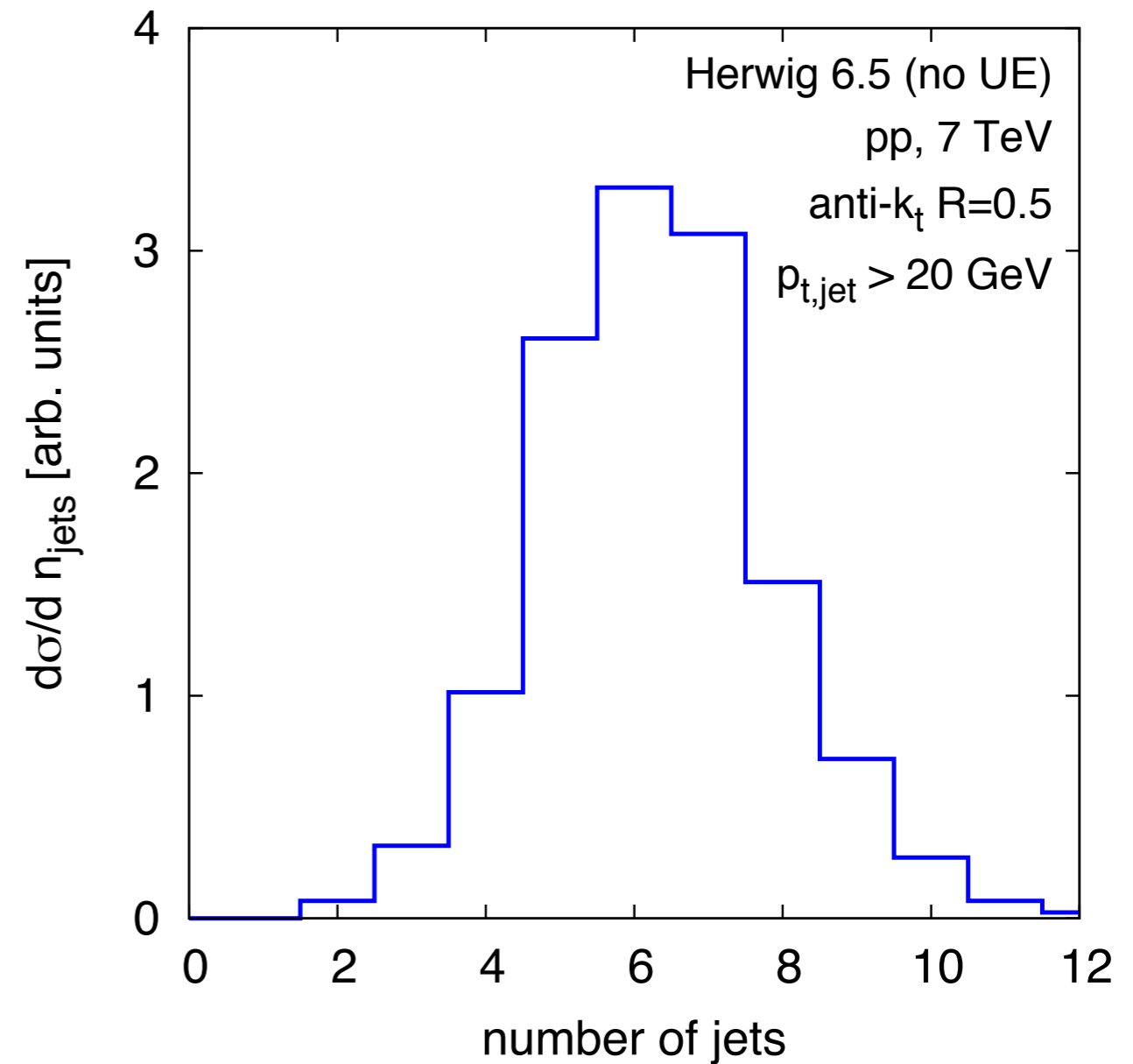
Alpgen $pp \rightarrow t\bar{t} \rightarrow 6q$ fraction of $pp \rightarrow t\bar{t} \rightarrow 6q$ events with all $R_{qq} > R$ 

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Alpgen $pp \rightarrow t\bar{t} \rightarrow 6q$ fraction of $pp \rightarrow t\bar{t} \rightarrow 6q$ events with all $R_{qq} > R$ **Herwig $pp \rightarrow t\bar{t} \rightarrow$ hadrons**

Distribution of number of jets



- Uses the anti- k_t algorithm
- Uses a jet radius $R=0.4$
- Uses a transverse momentum threshold that is typically at least 20 GeV (exact value depends on the analysis)

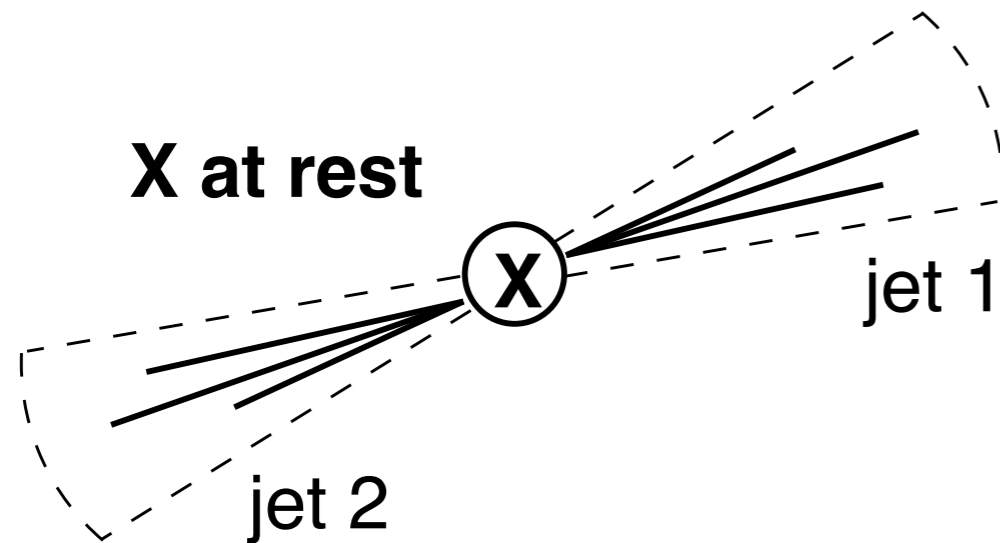
Radius and p_T threshold choices give a good compromise between

- ability to resolve multi-jet physics
- loss of radiation from jets
- additional spurious jets
- contamination from pileup

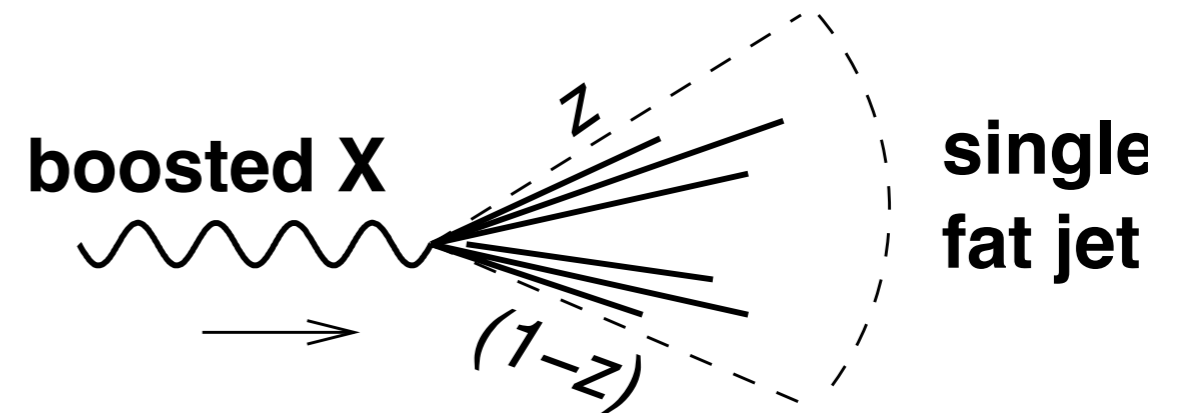
boosted object reconstruction

Boosted EW scale objects

Normal analyses: two quarks from $X \rightarrow q\bar{q}$ reconstructed as two jets



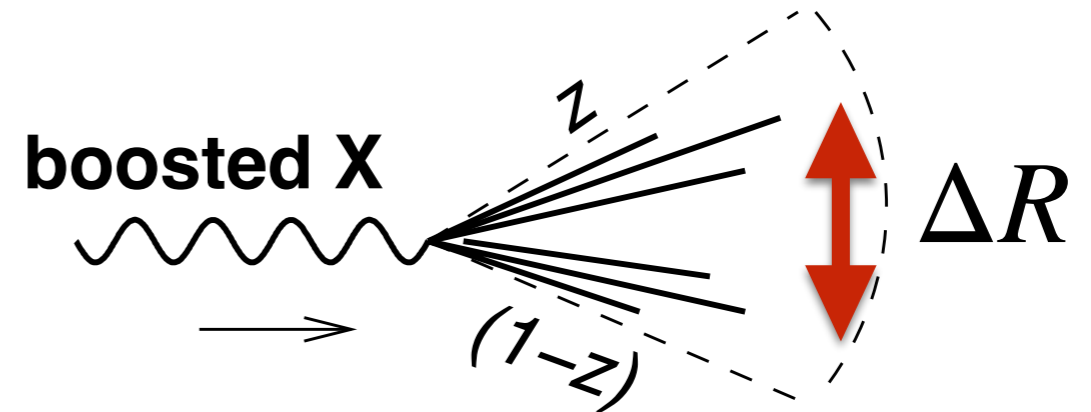
High- p_t regime: EW object X is boosted, decay is collimated, $q\bar{q}$ both in same jet



Boosted EW scale objects

High- p_t regime: EW object X is boosted, decay is collimated, $q\bar{q}$ both in same jet

$$\begin{aligned} m_X^2 &\simeq E_X^2 \cdot z(1-z) \cdot 2(1-\cos\theta) \\ &\simeq p_{tX}^2 \cdot z(1-z) \cdot \Delta R^2 \end{aligned}$$



The two prongs end up in a single jet if

$$\Delta R \simeq \frac{m}{p_t} \frac{1}{\sqrt{z(1-z)}} \sim \frac{2m}{p_t} < R \quad \text{or} \quad p_t \gtrsim \frac{2m}{R}$$

E.g. W-boson with $p_t > 400 \text{ GeV}$ ends up collimated in a single jet.

Two widely used terms
though there's not a
consensus about
what they mean

Tagging

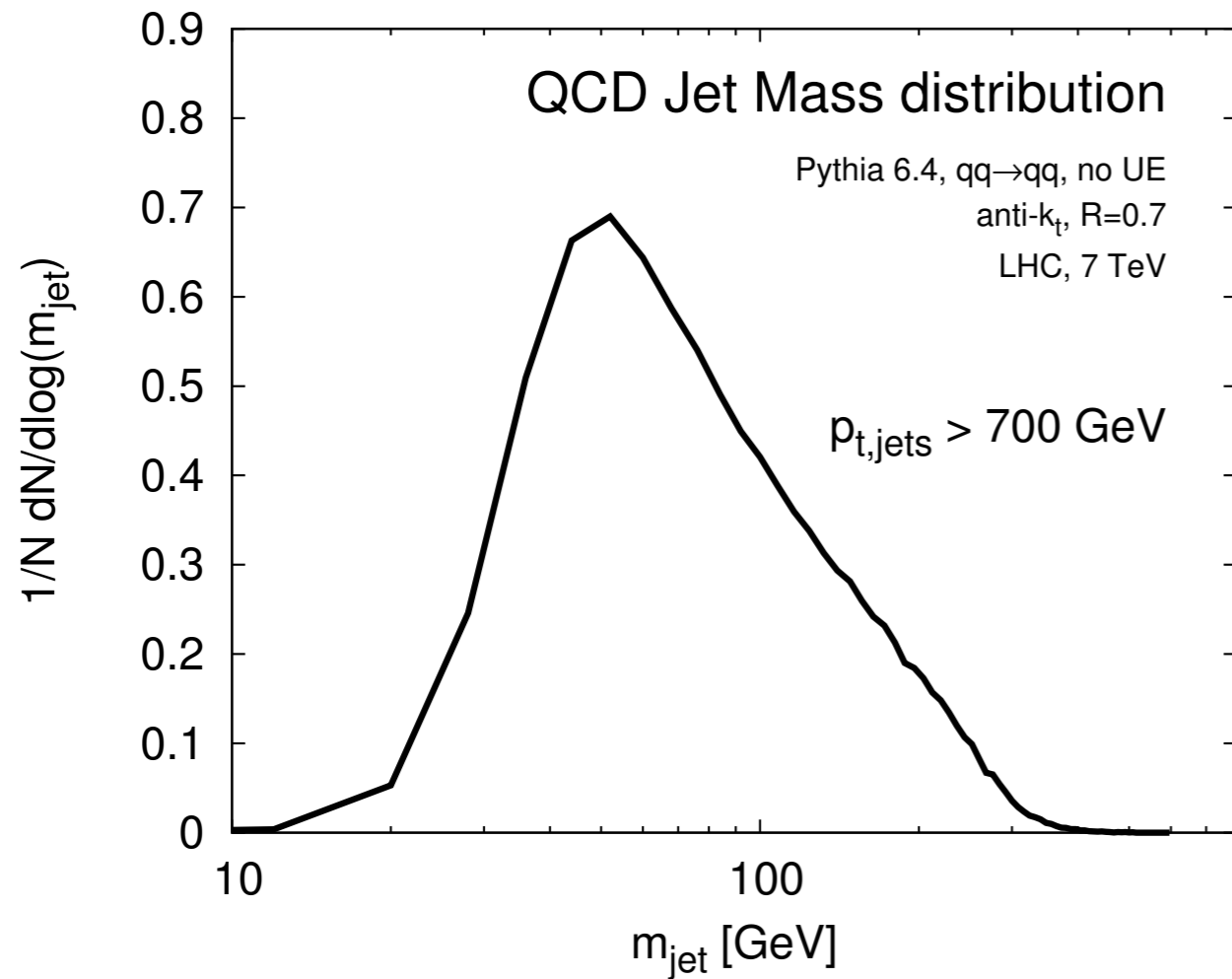
- reduces the background, leaves much of signal
- you can tag with underlying hard n-prong structure and based on radiation pattern

Grooming

- improves signal mass resolution (removing pileup, etc.), without significantly changing background & signal event numbers

One core idea for
tagging

Inside the jet mass

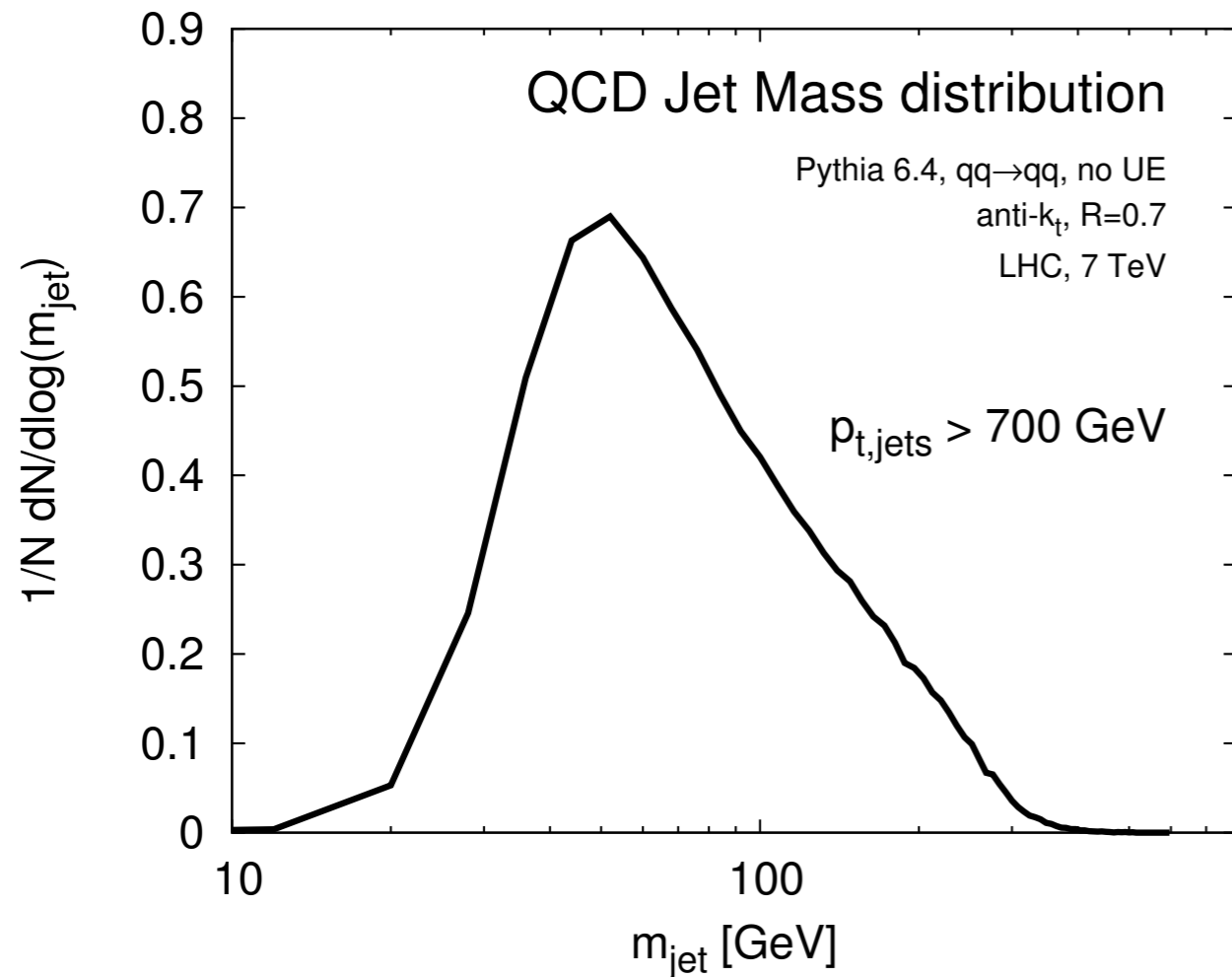


QCD jet mass distribution has the approximate

$$\frac{dN}{d \ln m} \sim \alpha_s \ln \frac{p_t R}{m} \times \text{Sudakov}$$

Work from '80s and '90s

Inside the jet mass



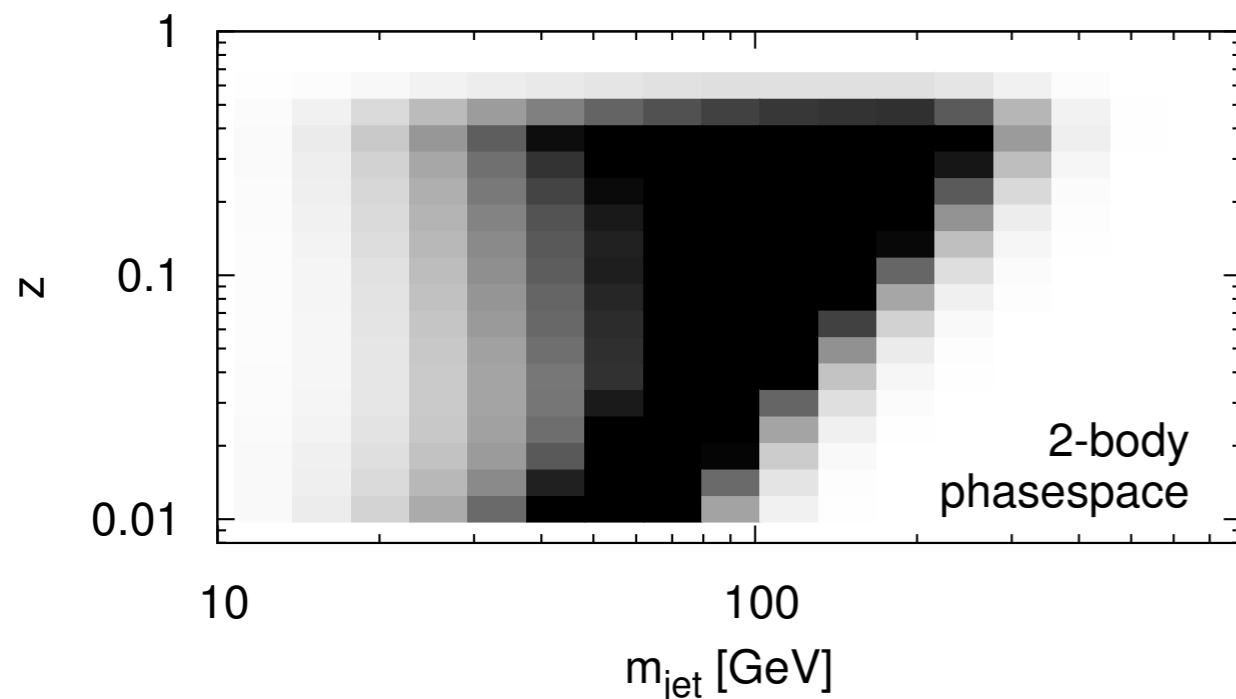
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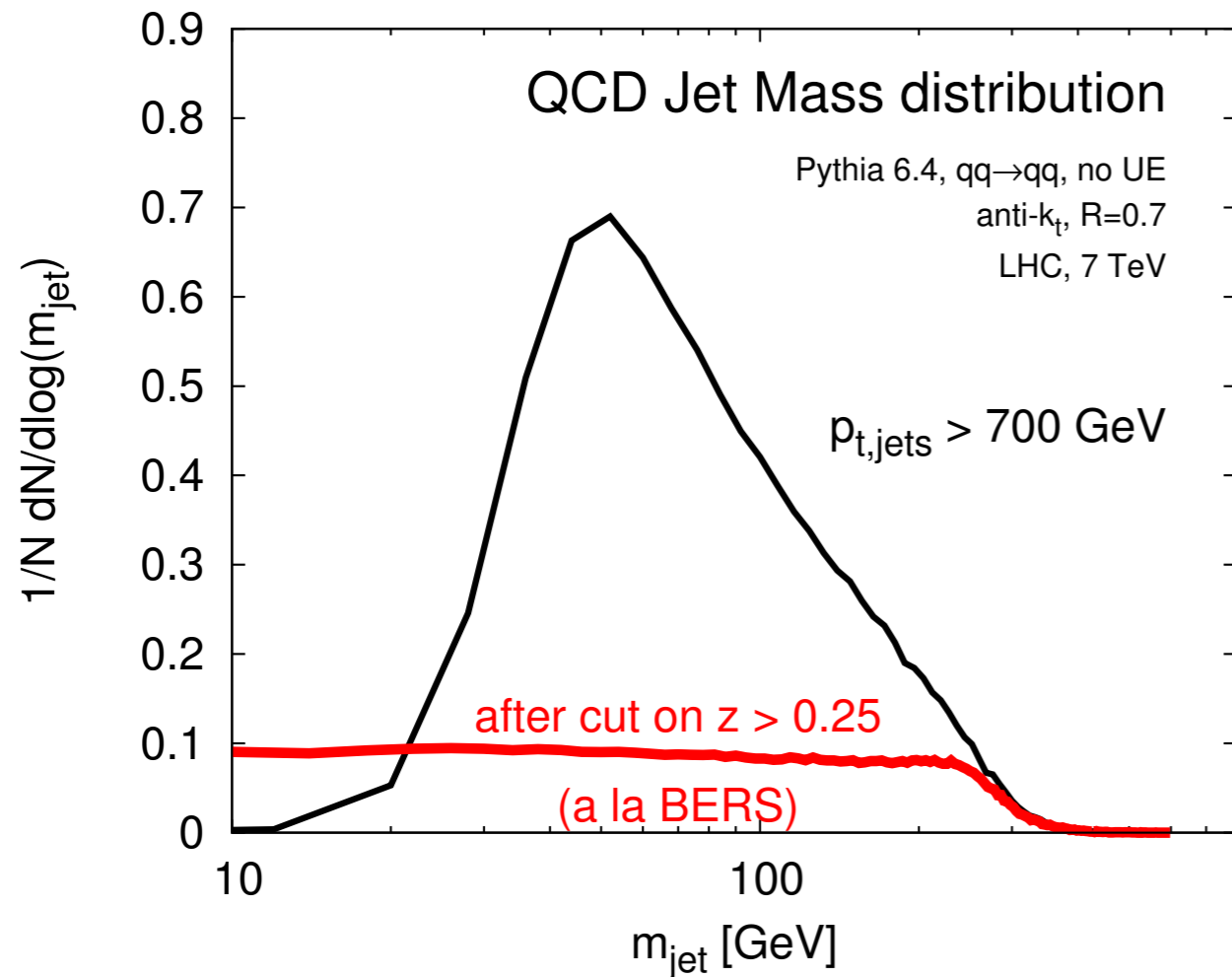
Work from '80s and '90s

The logarithm comes from integral over soft divergence of QCD:

$$\int_{\frac{m^2}{p_t^2 R^2}}^{\frac{1}{2}} \frac{dz}{z}$$



Inside the jet mass



QCD jet mass distribution has the approximate

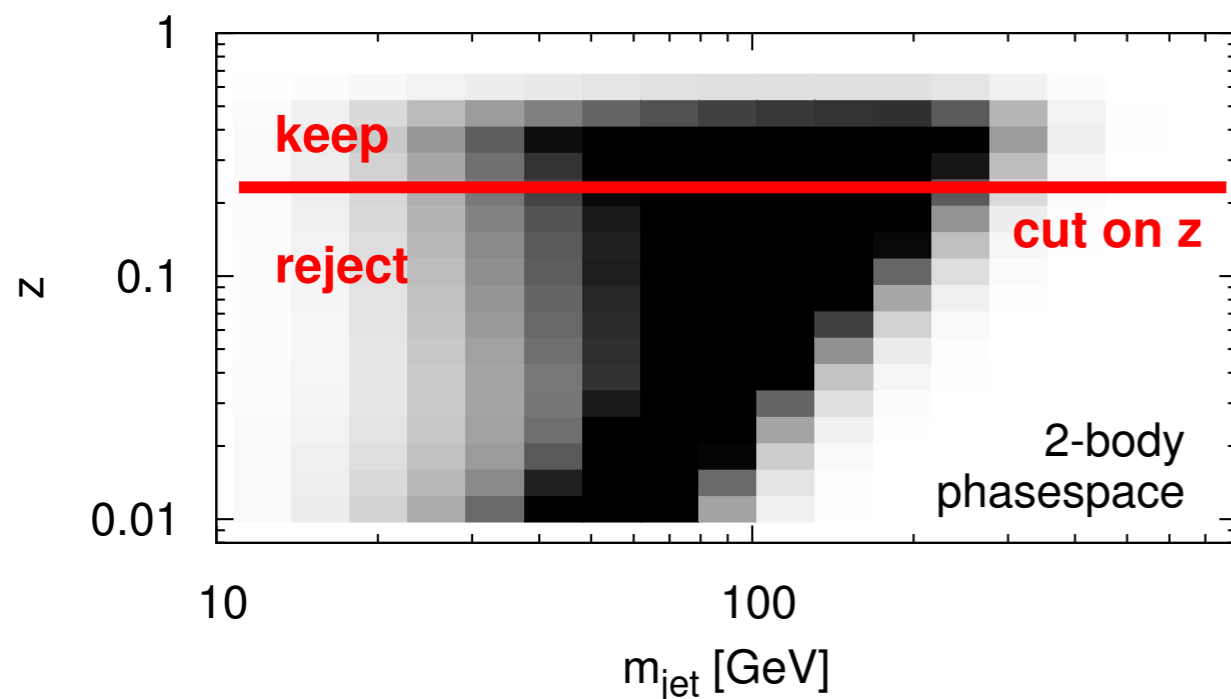
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Work from '80s and '90s

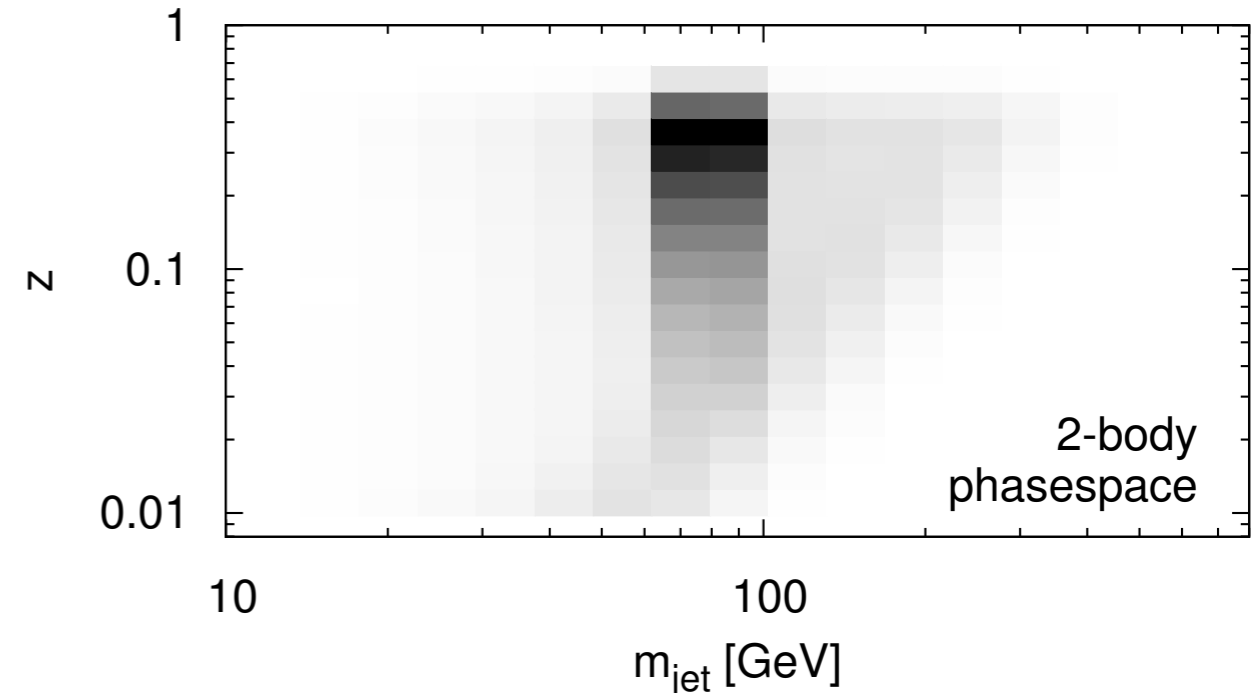
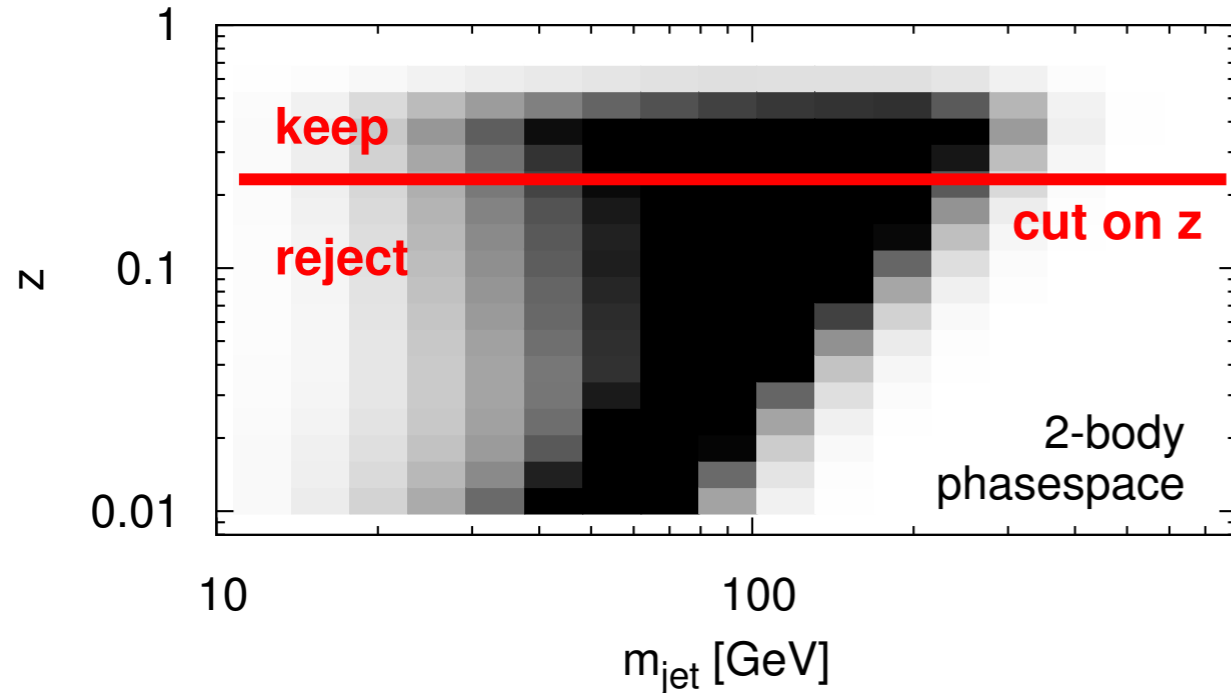
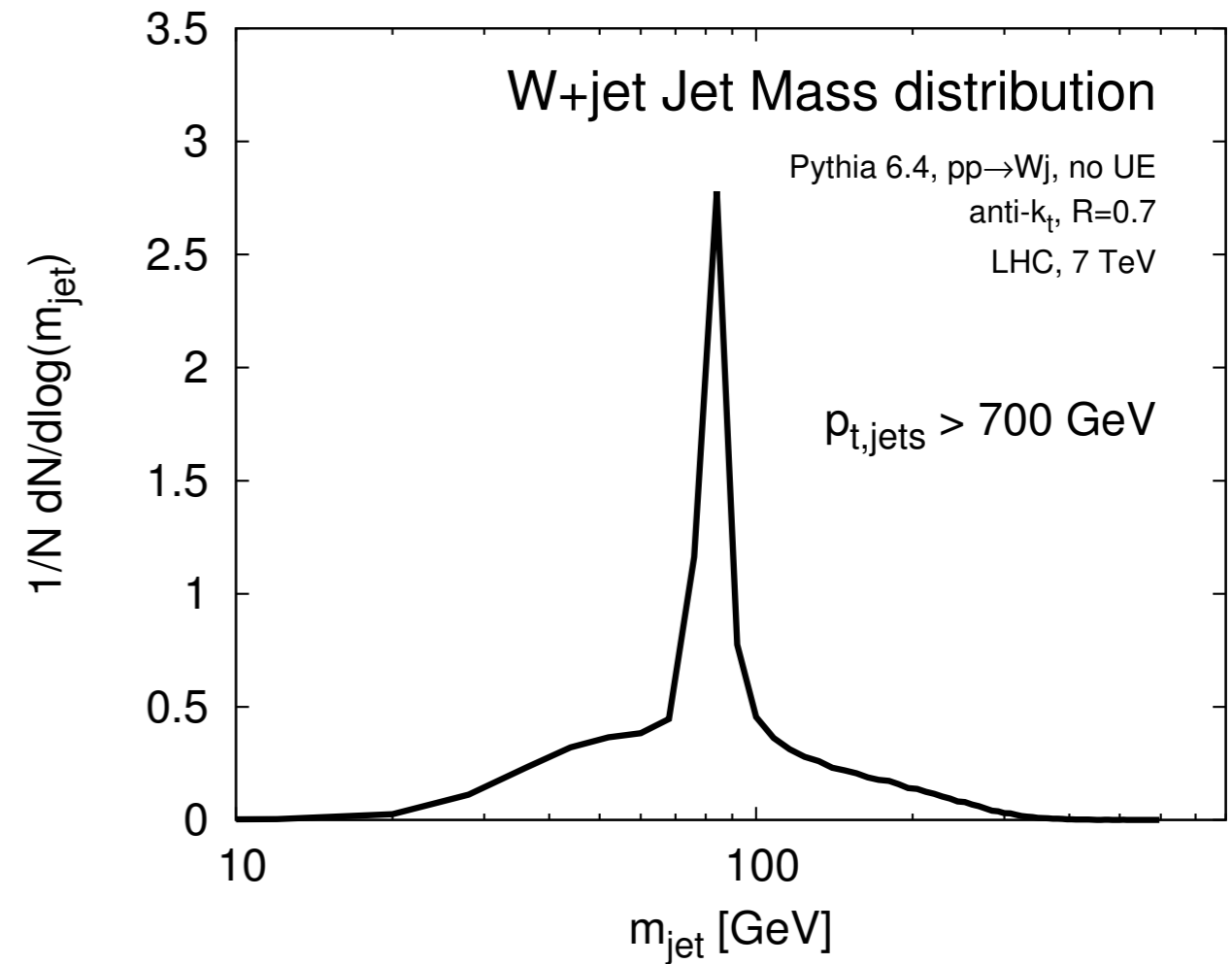
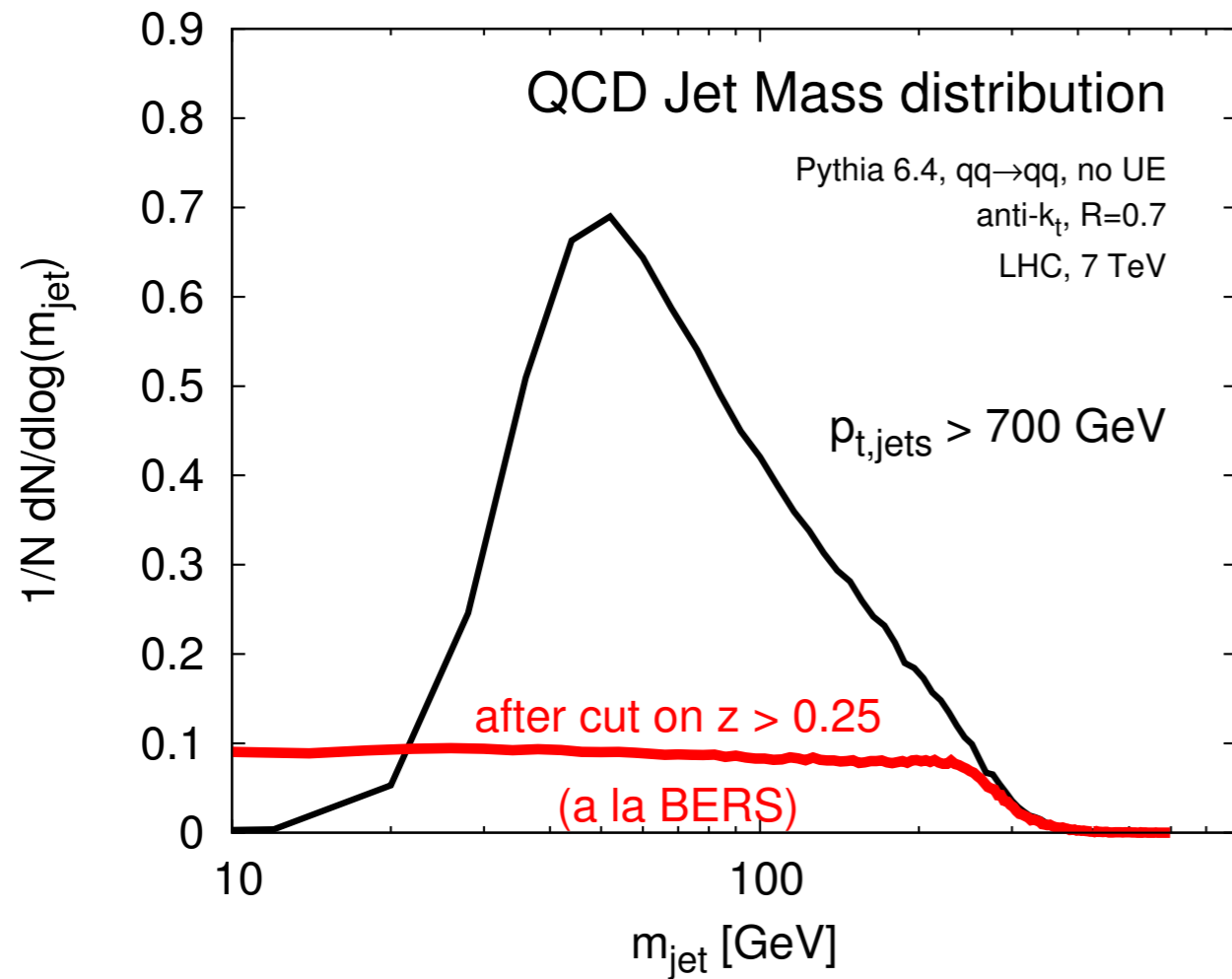
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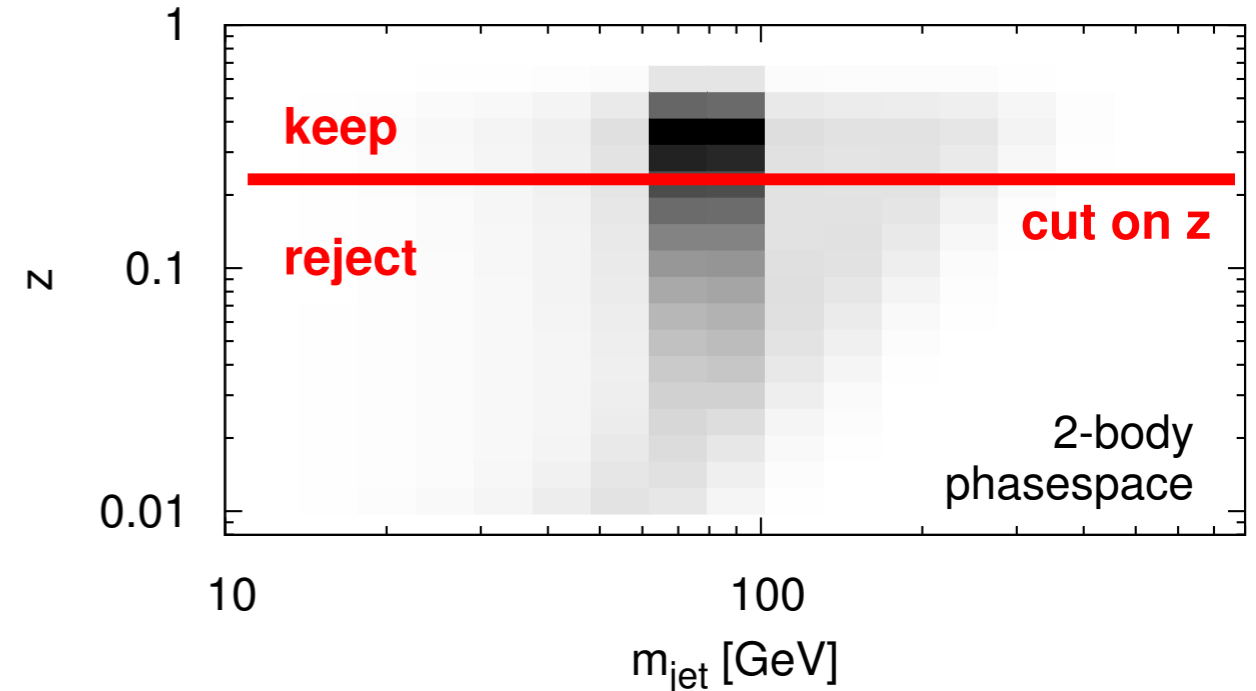
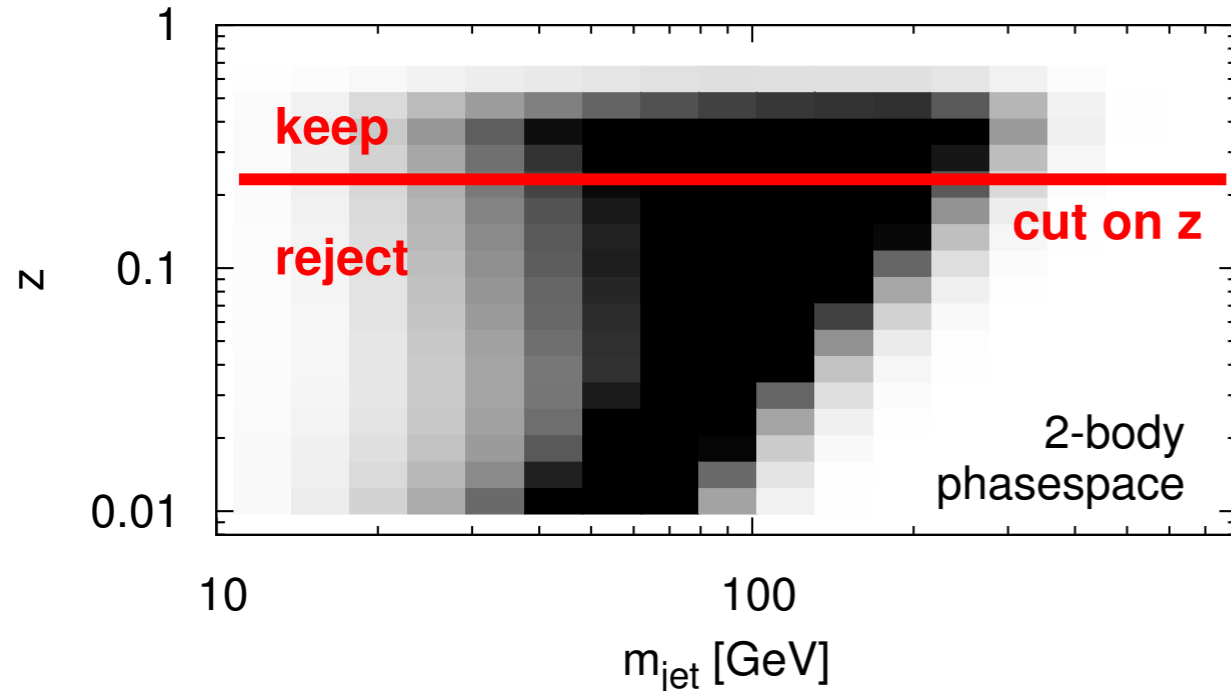
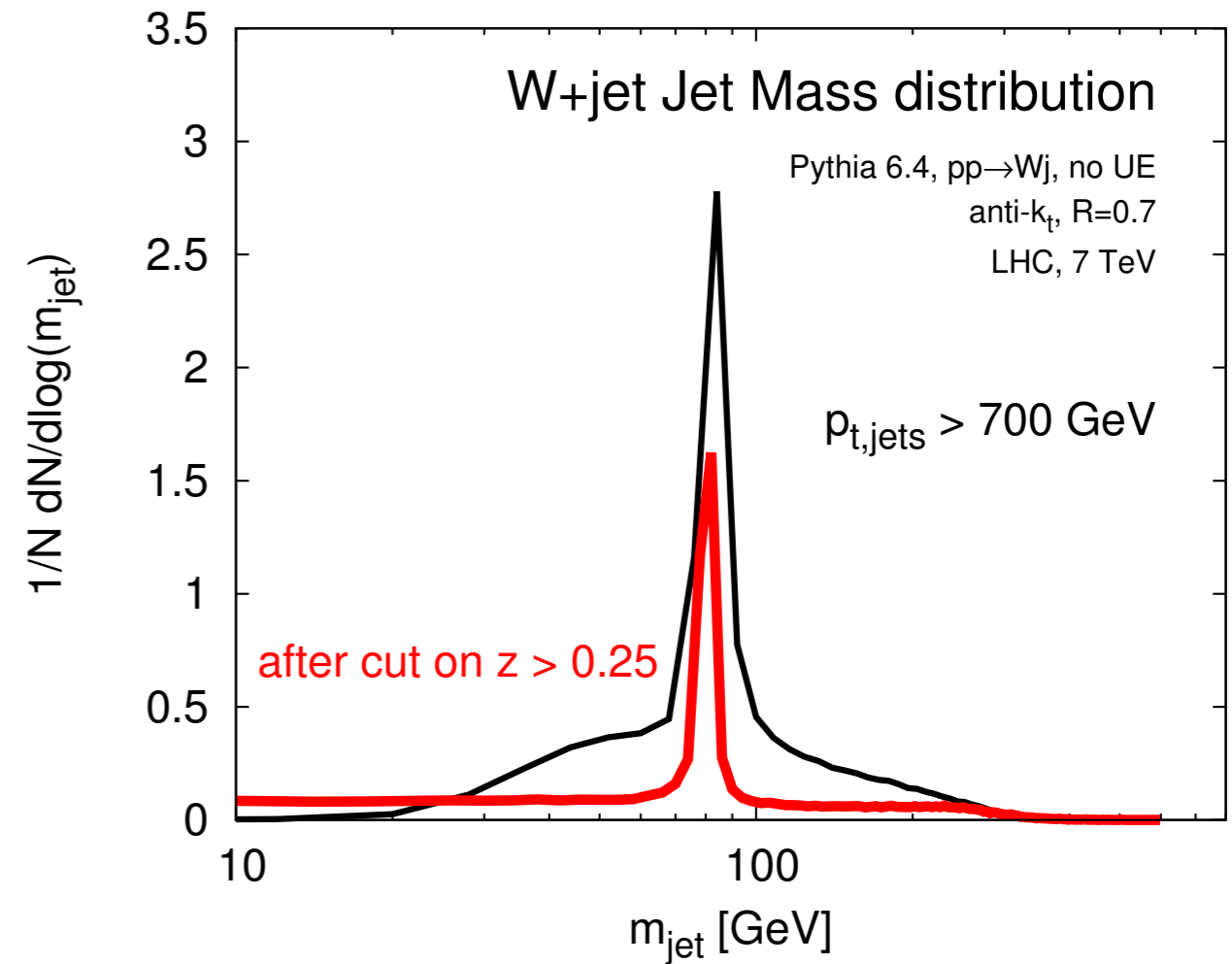
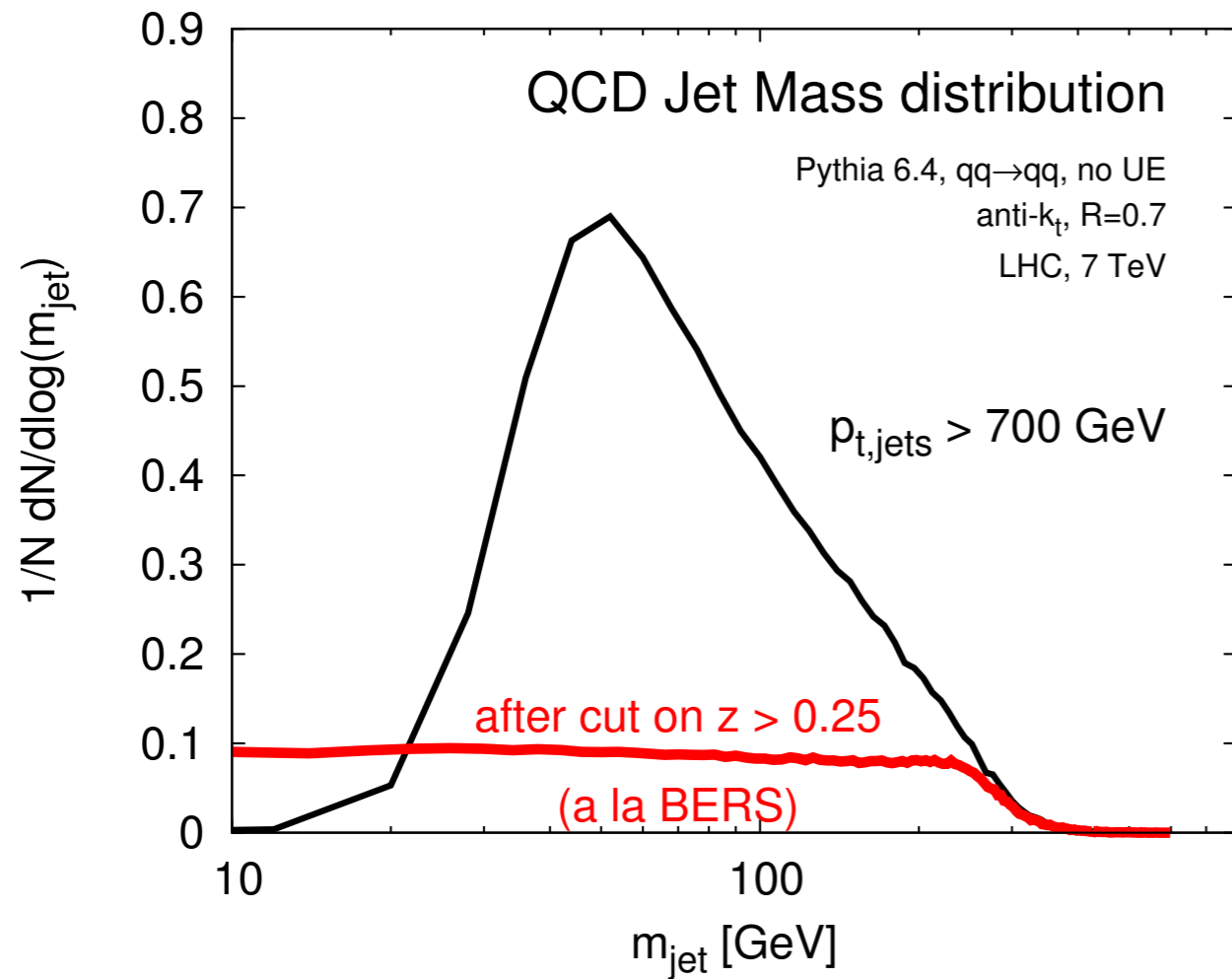
A hard cut on z reduces QCD background & simplifies its shape

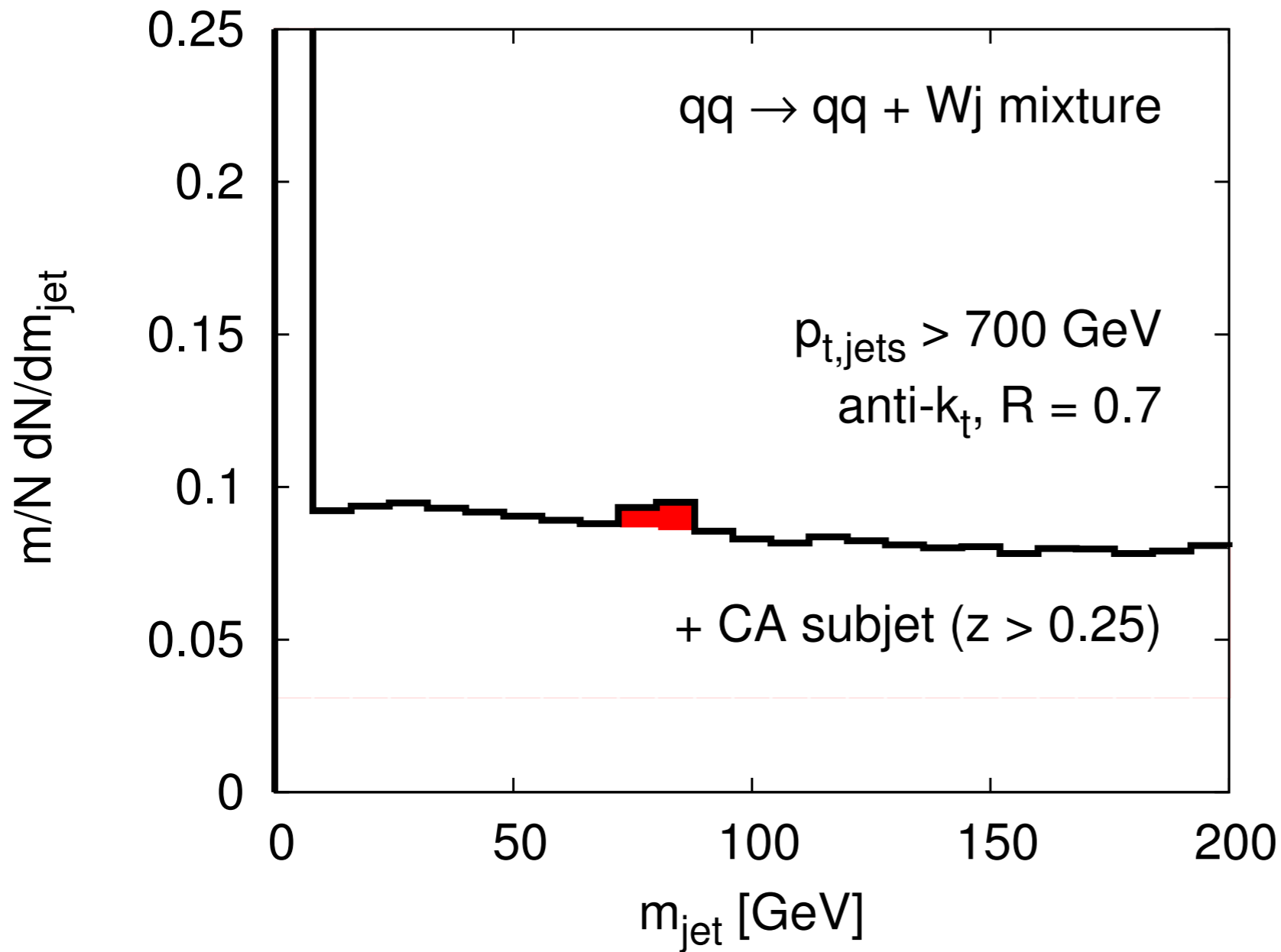


Inside the jet mass



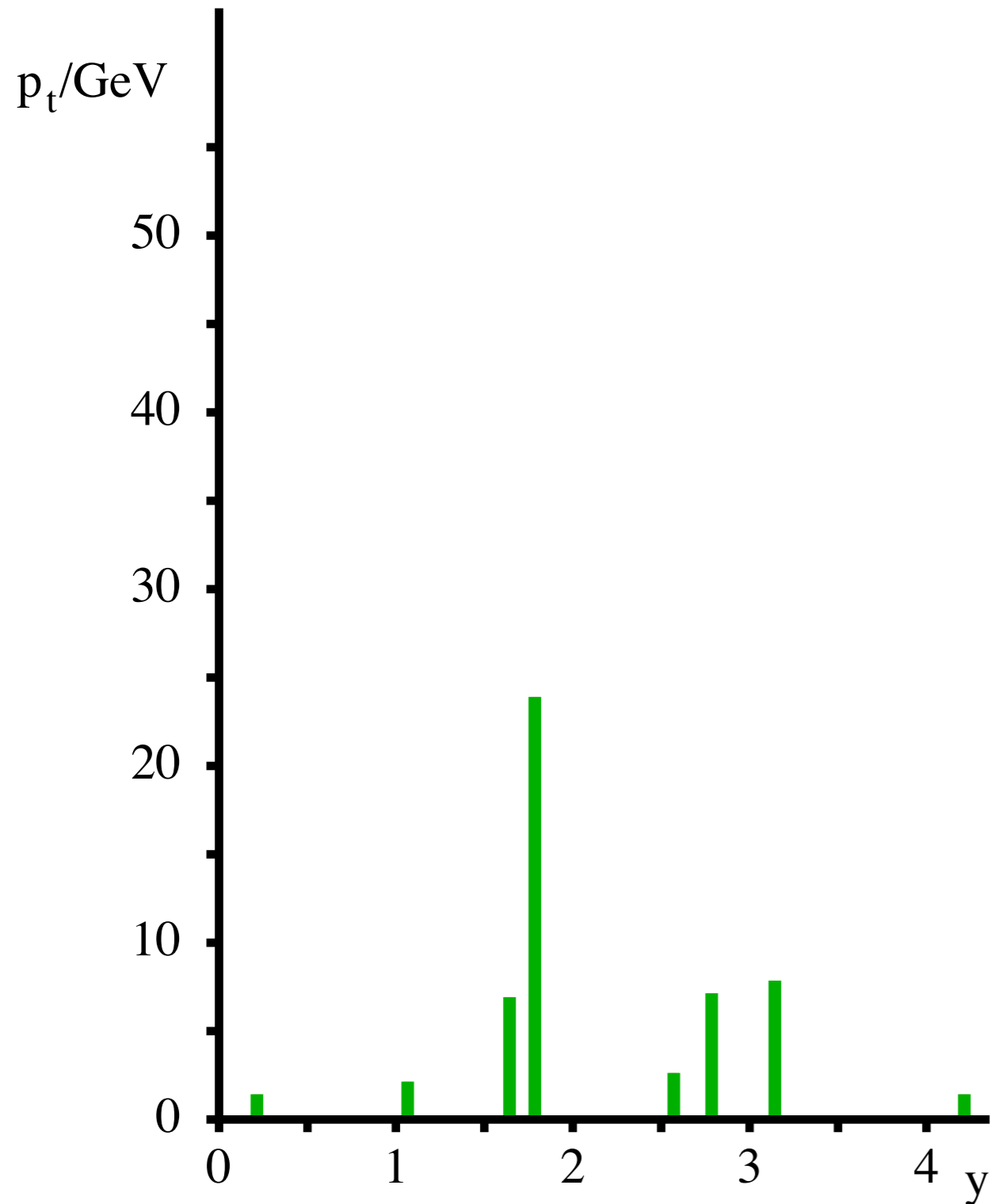
Inside the jet mass





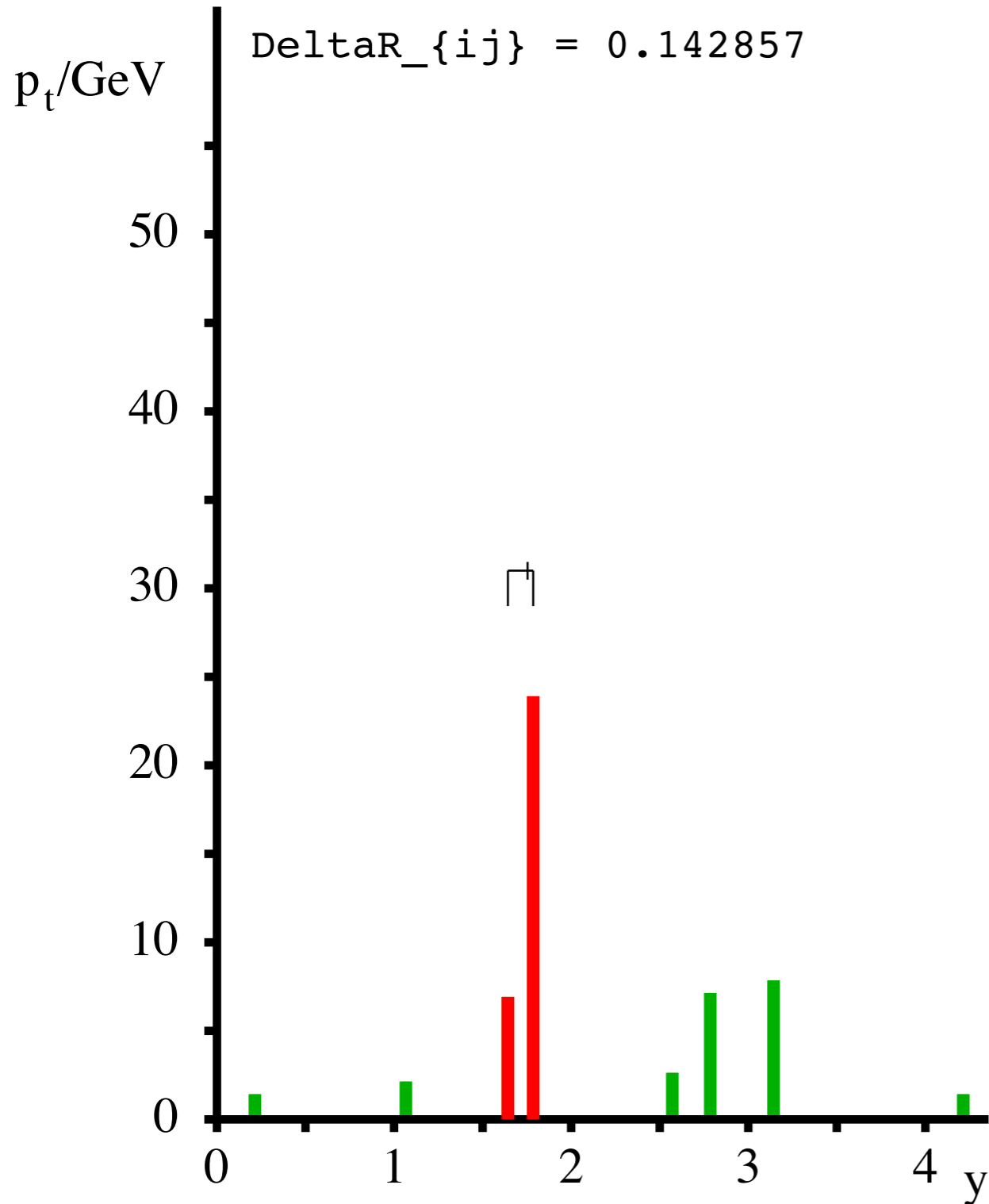
Signal + bkgd
after cut on z

Cambridge/Aachen algorithm



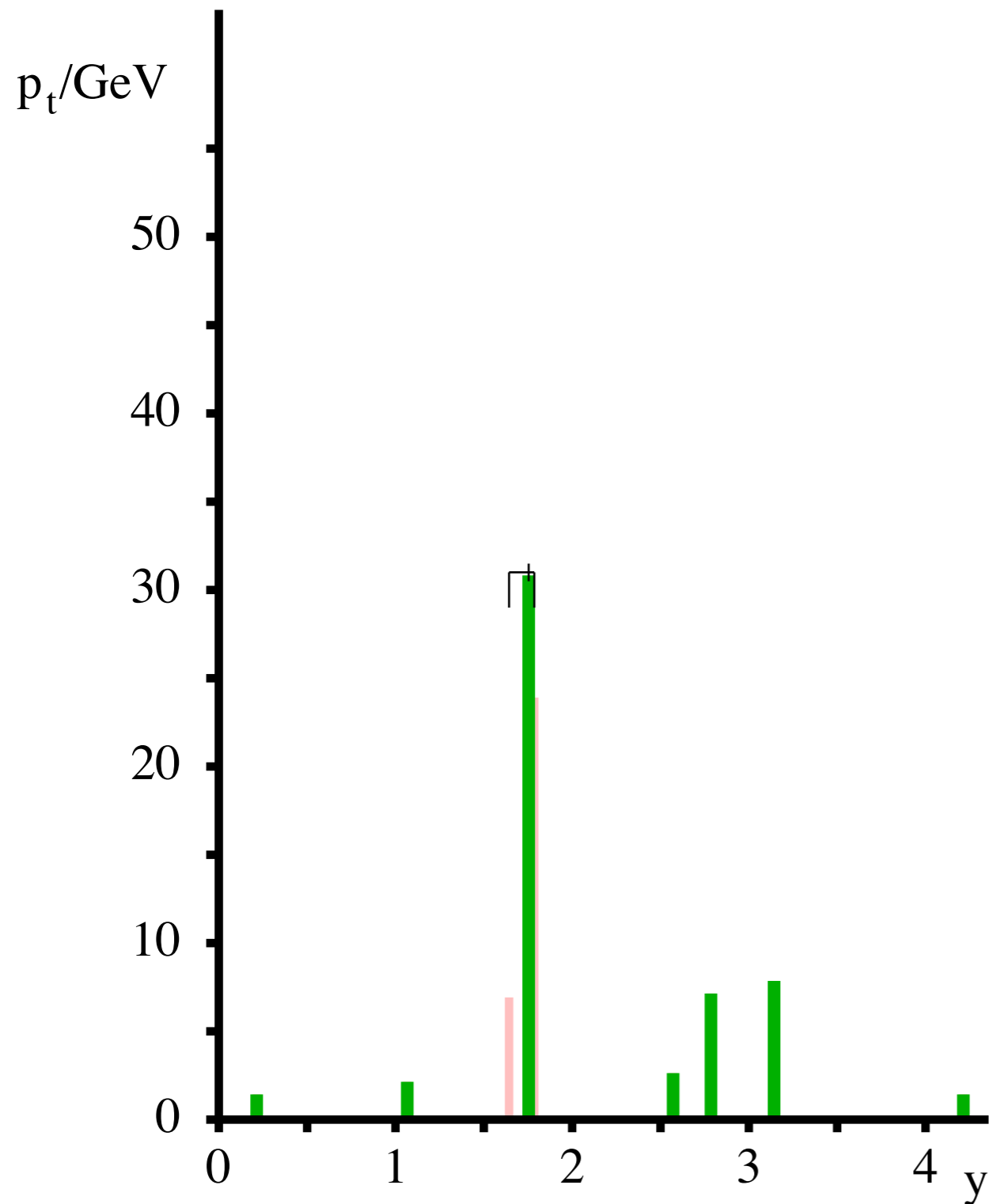
How well can an algorithm identify the “blobs” of energy inside a jet that come from different partons?

Cambridge/Aachen algorithm



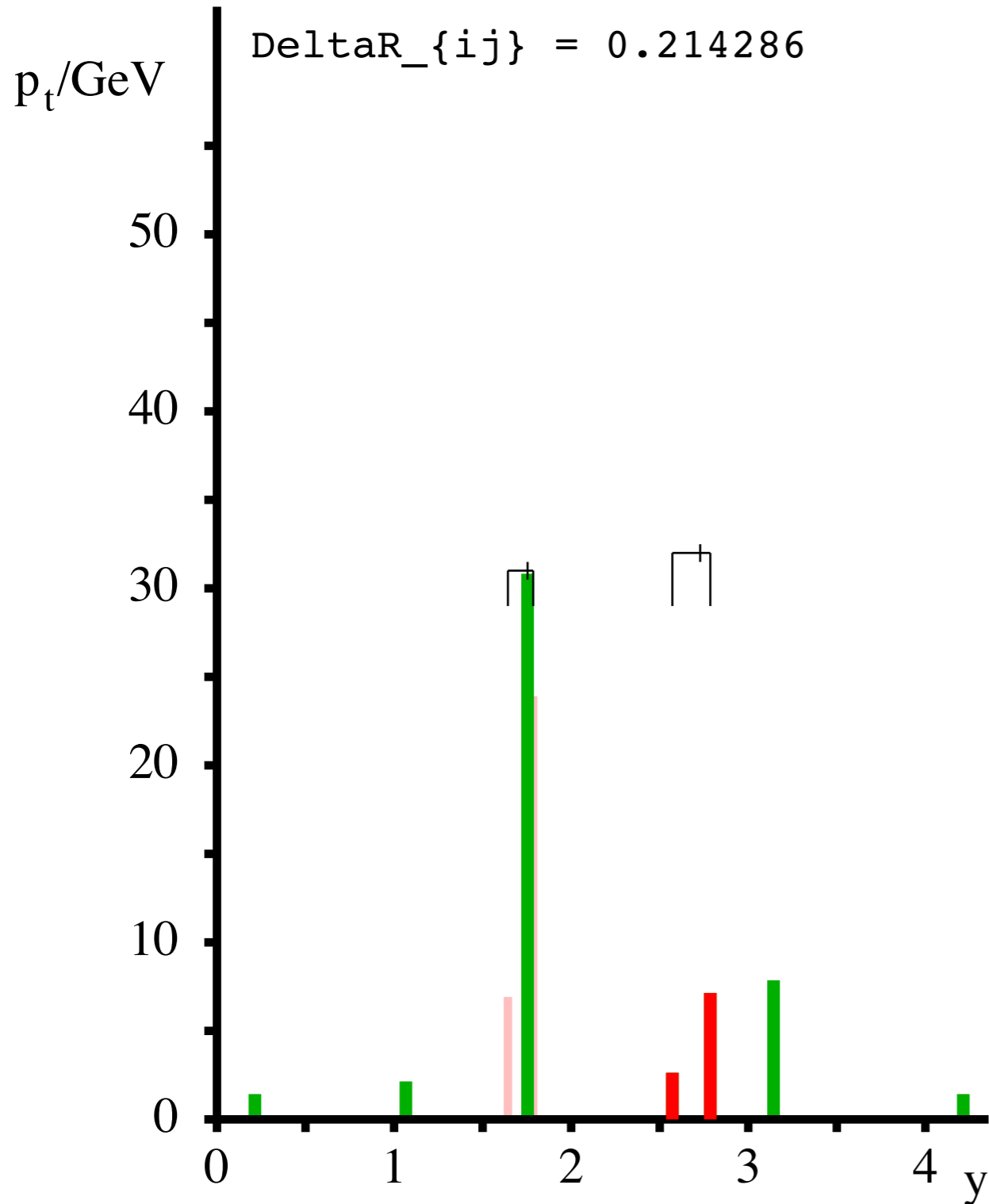
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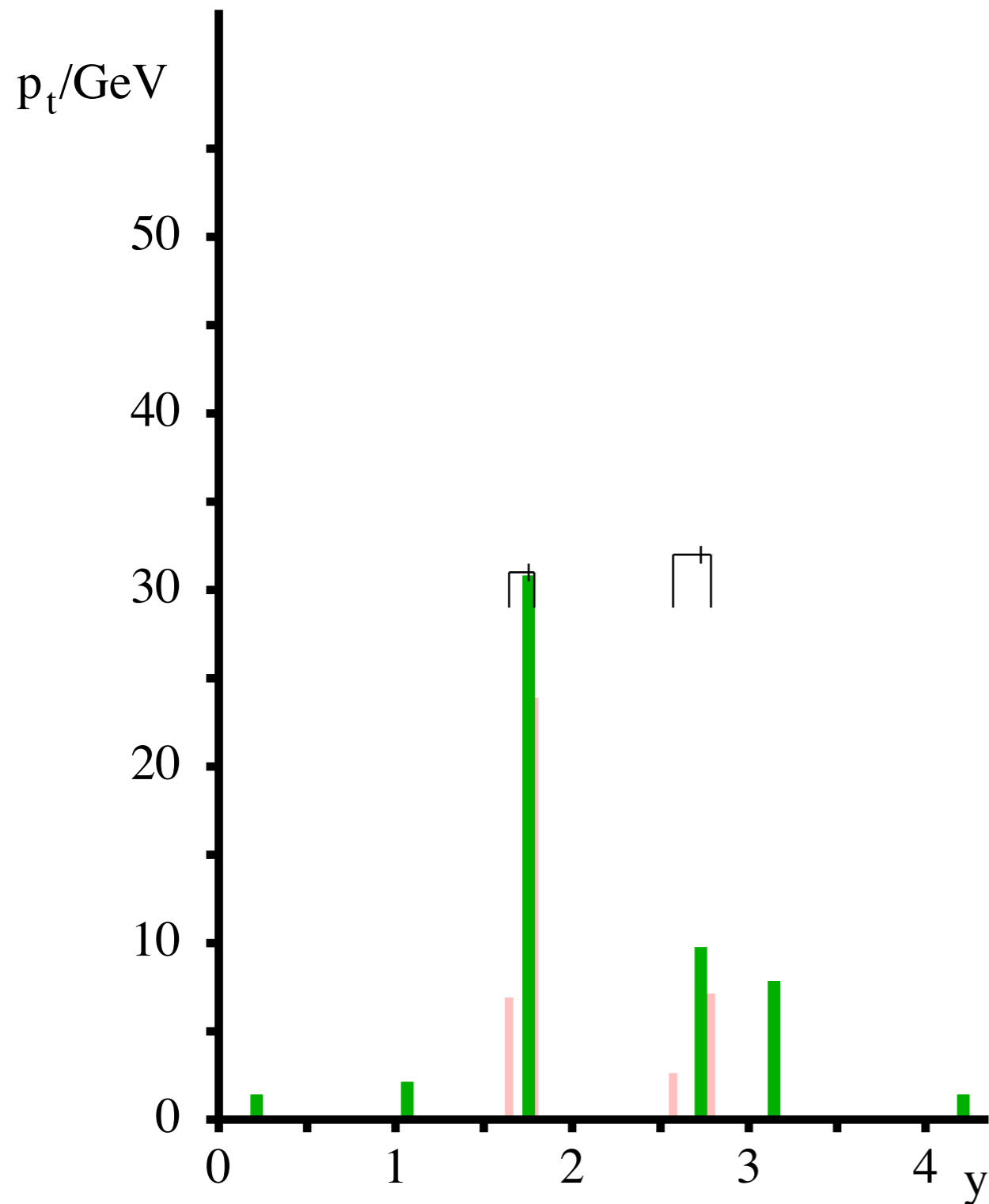
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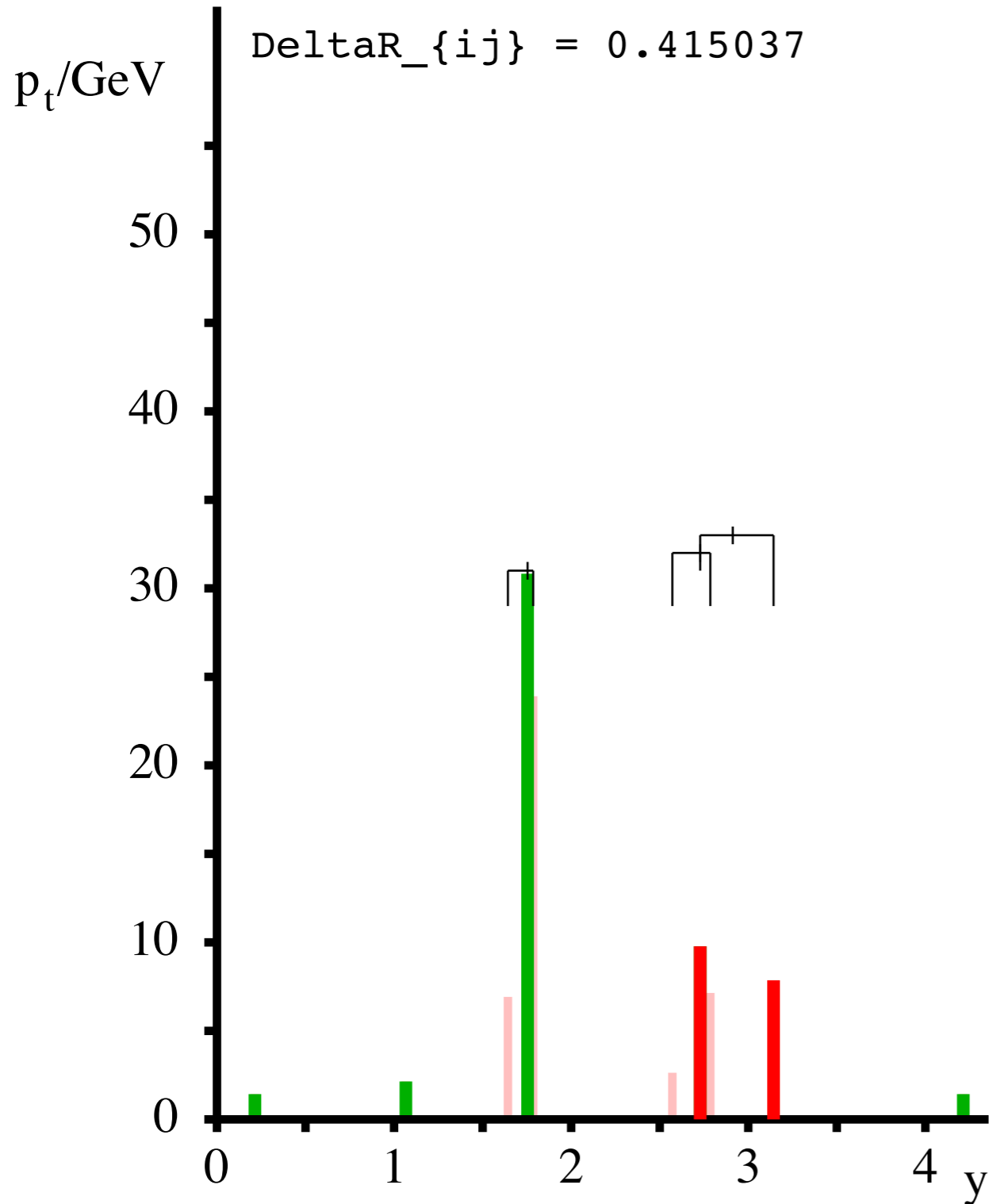
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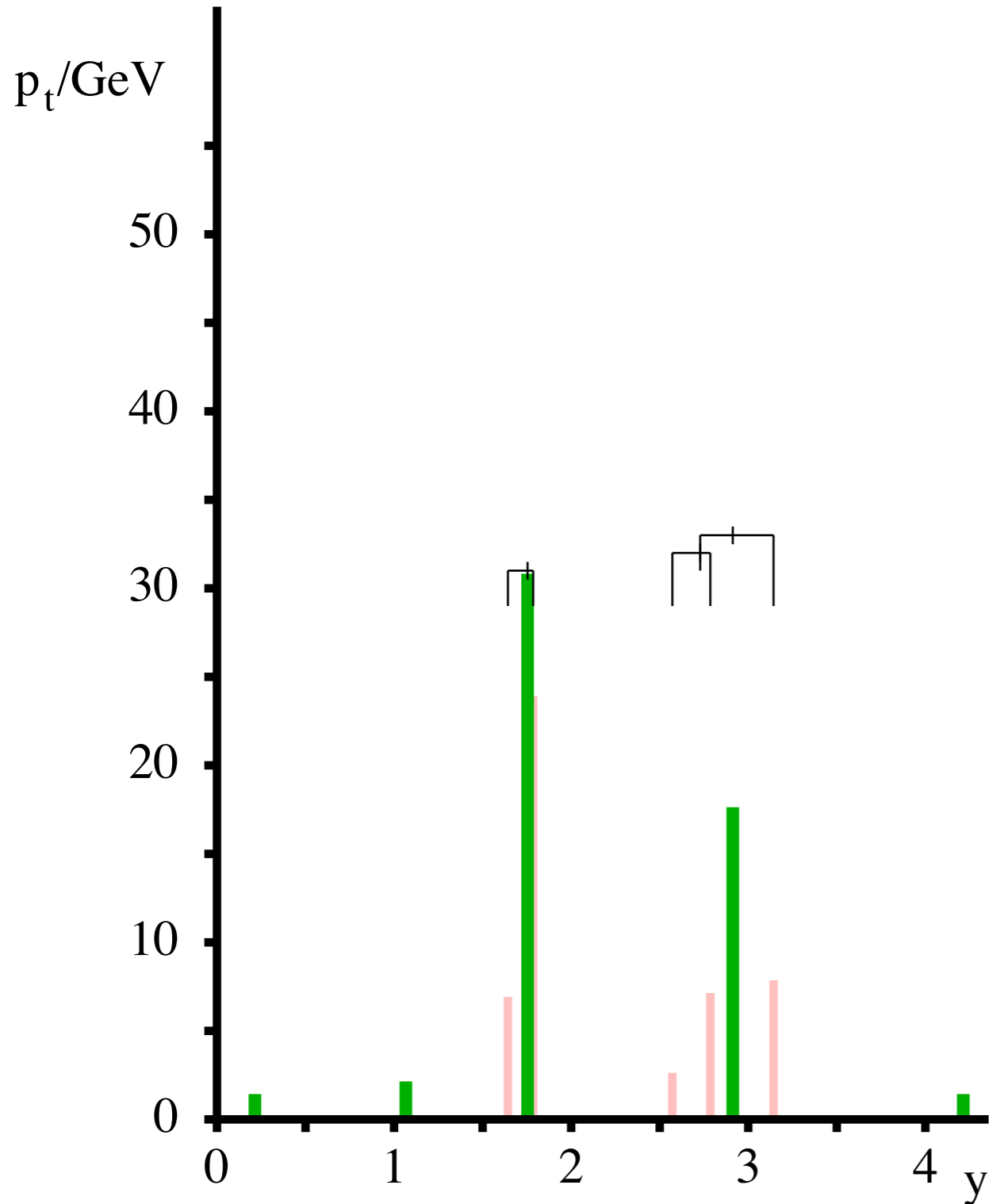
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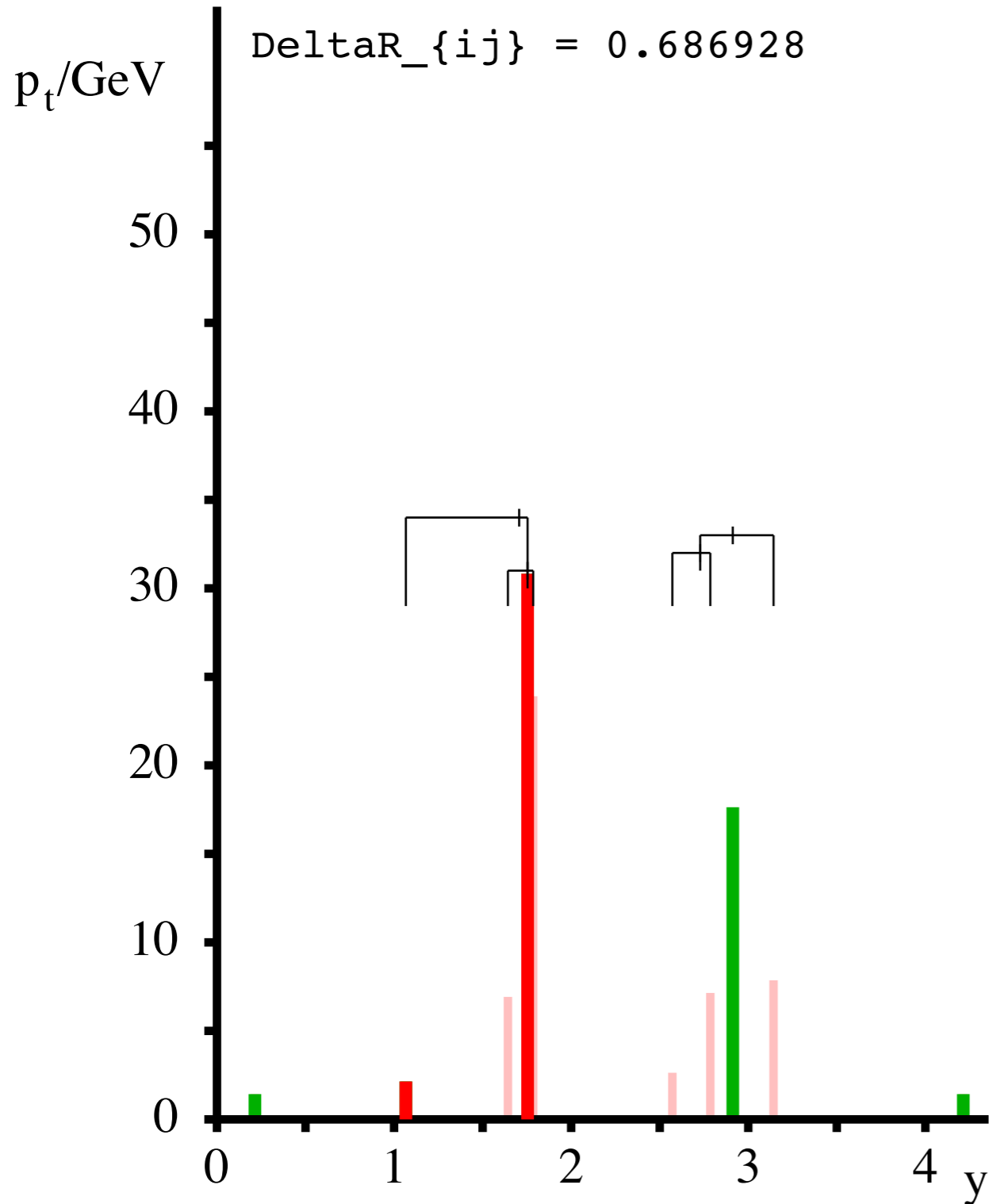
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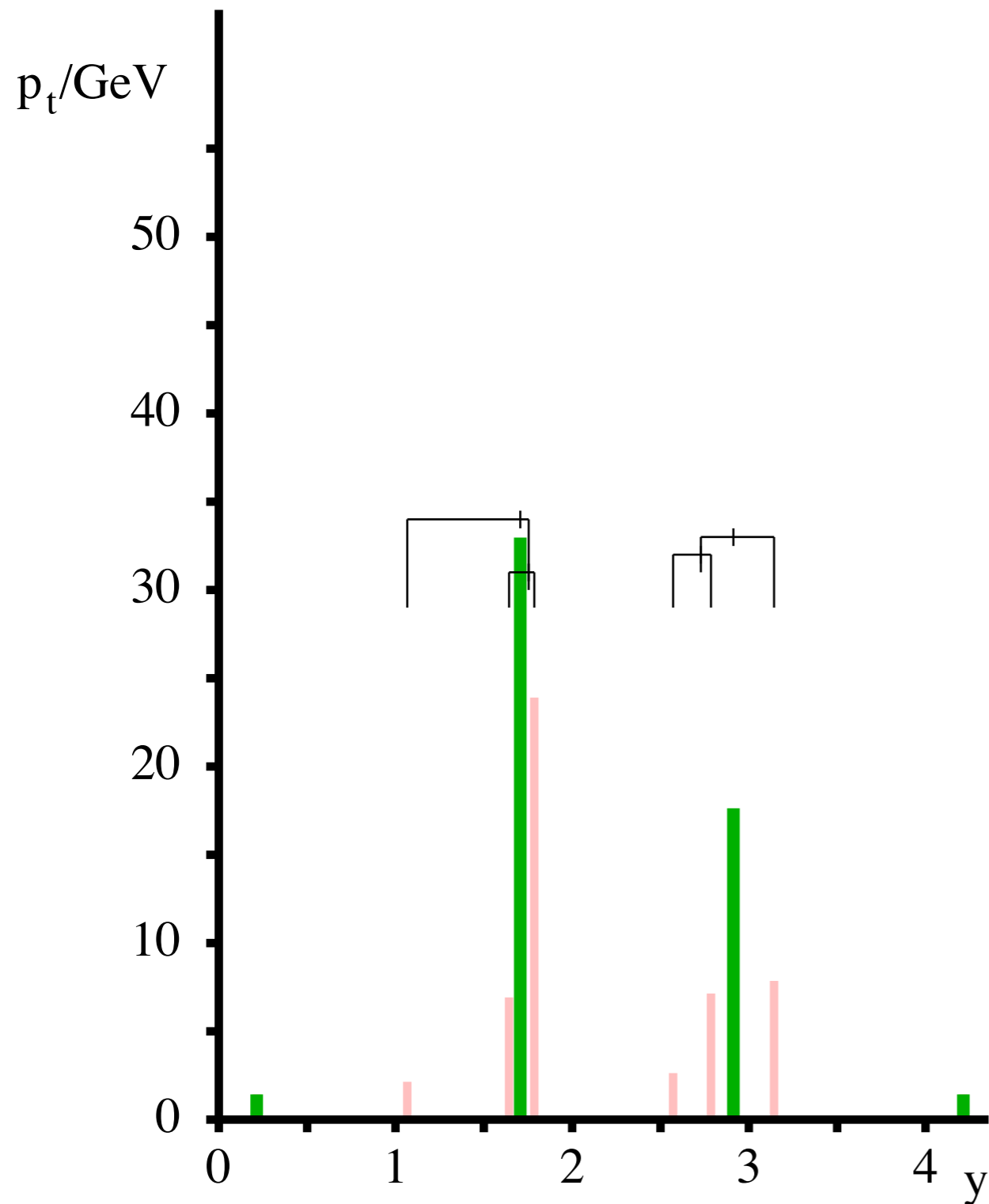
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Cambridge/Aachen algorithm



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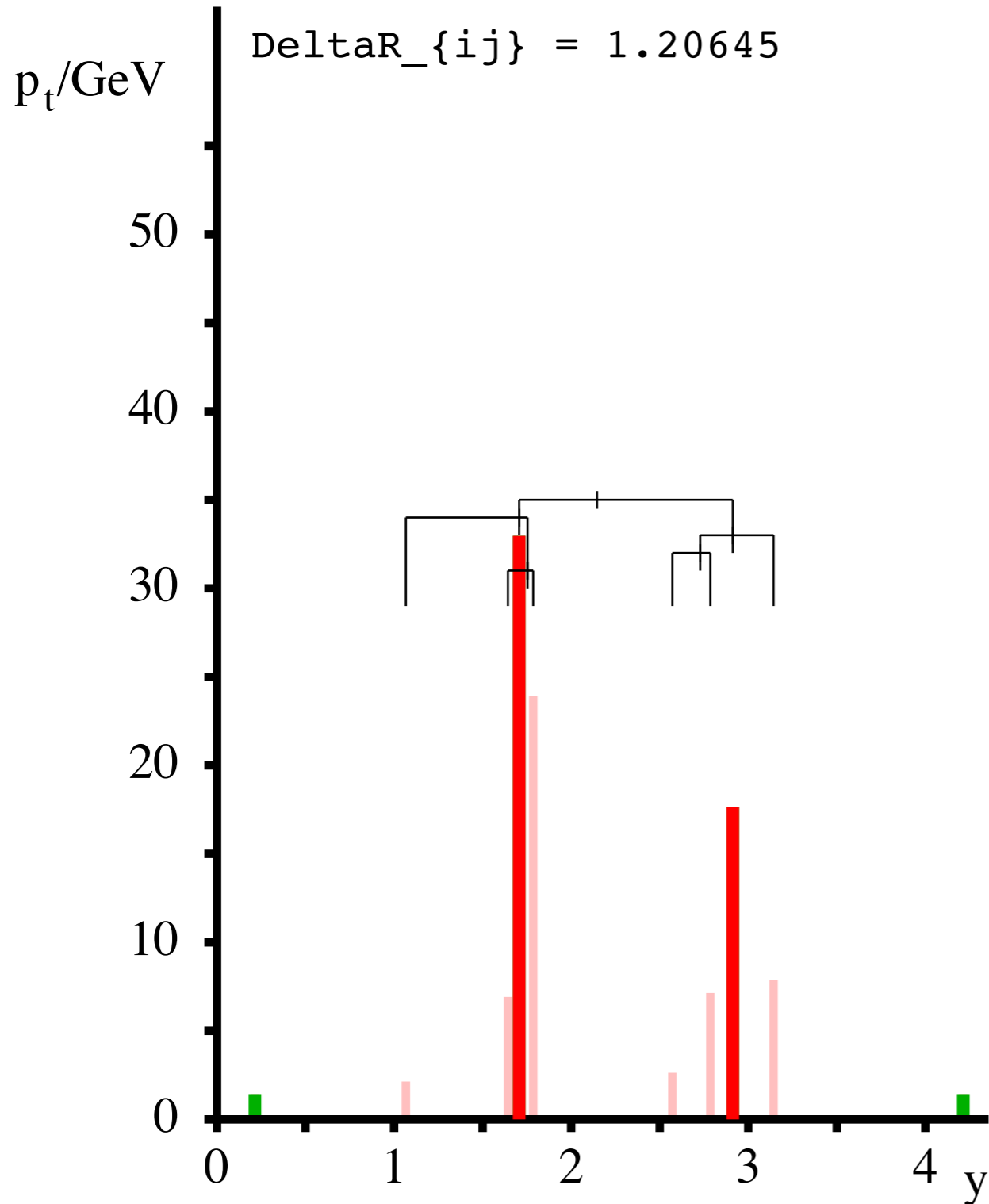
Cambridge/Aachen algorithm



How well can an algorithm identify the “blobs” of energy inside a jet that come from different partons?

C/A identifies two hard blobs with limited soft contamination

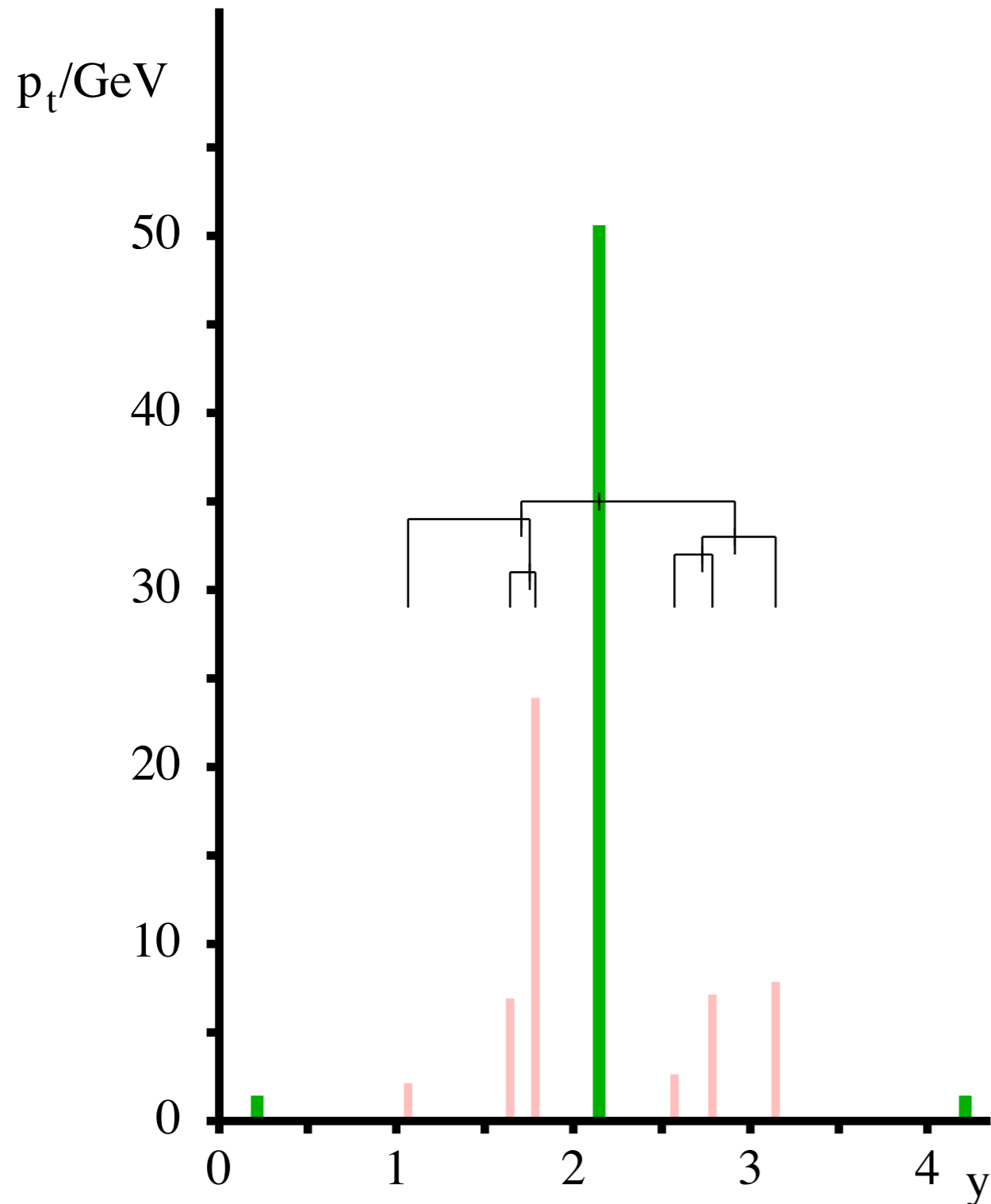
Cambridge/Aachen algorithm



How well can an algorithm identify the “blobs” of energy inside a jet that come from different partons?

C/A identifies two hard blobs with limited soft contamination, **joins them**

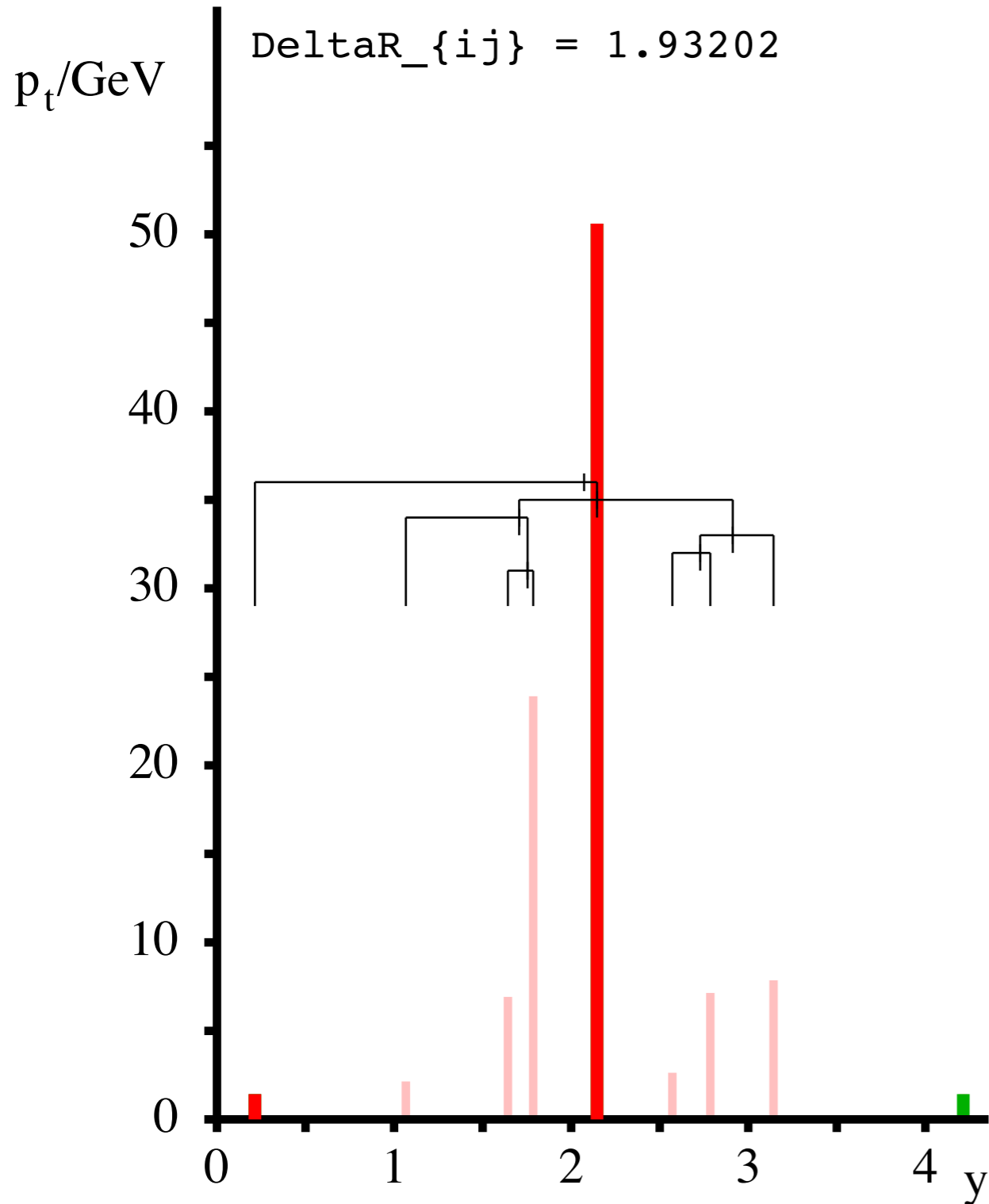
Cambridge/Aachen algorithm



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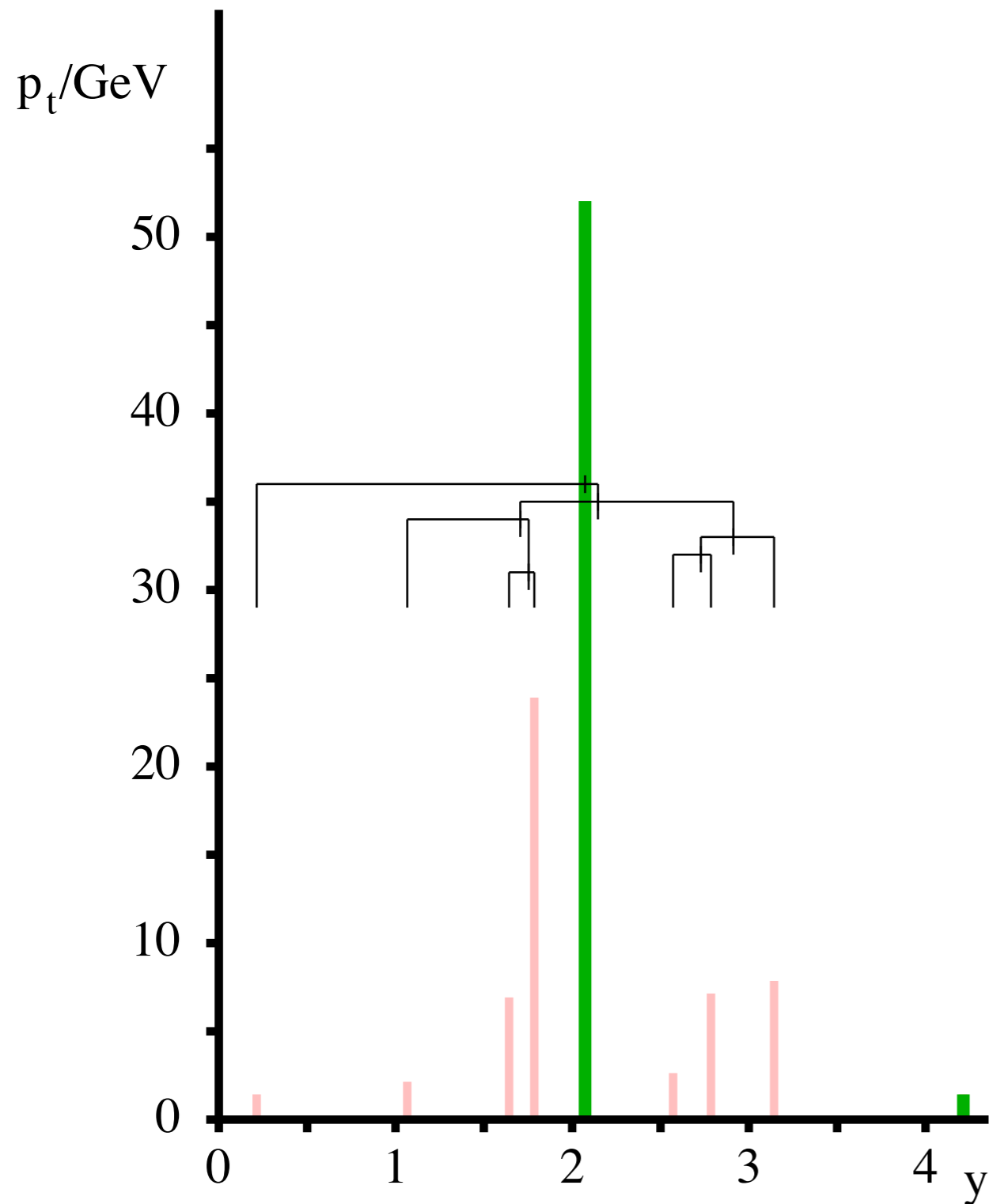
Cambridge/Aachen algorithm



How well can an algorithm identify the “blobs” of energy inside a jet that come from different partons?

C/A identifies two hard blobs with limited soft contamination, joins them, and then adds in remaining soft junk

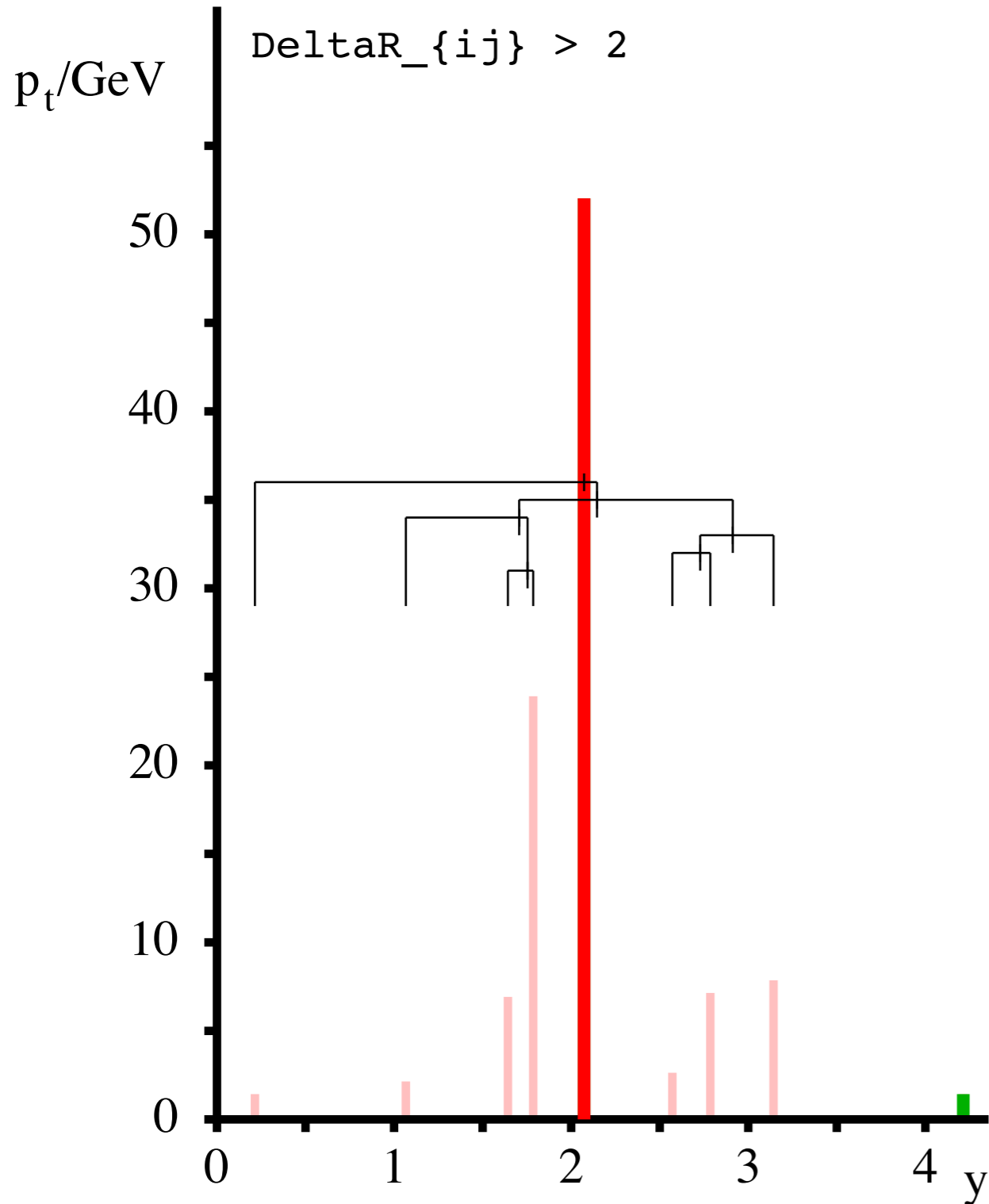
Cambridge/Aachen algorithm



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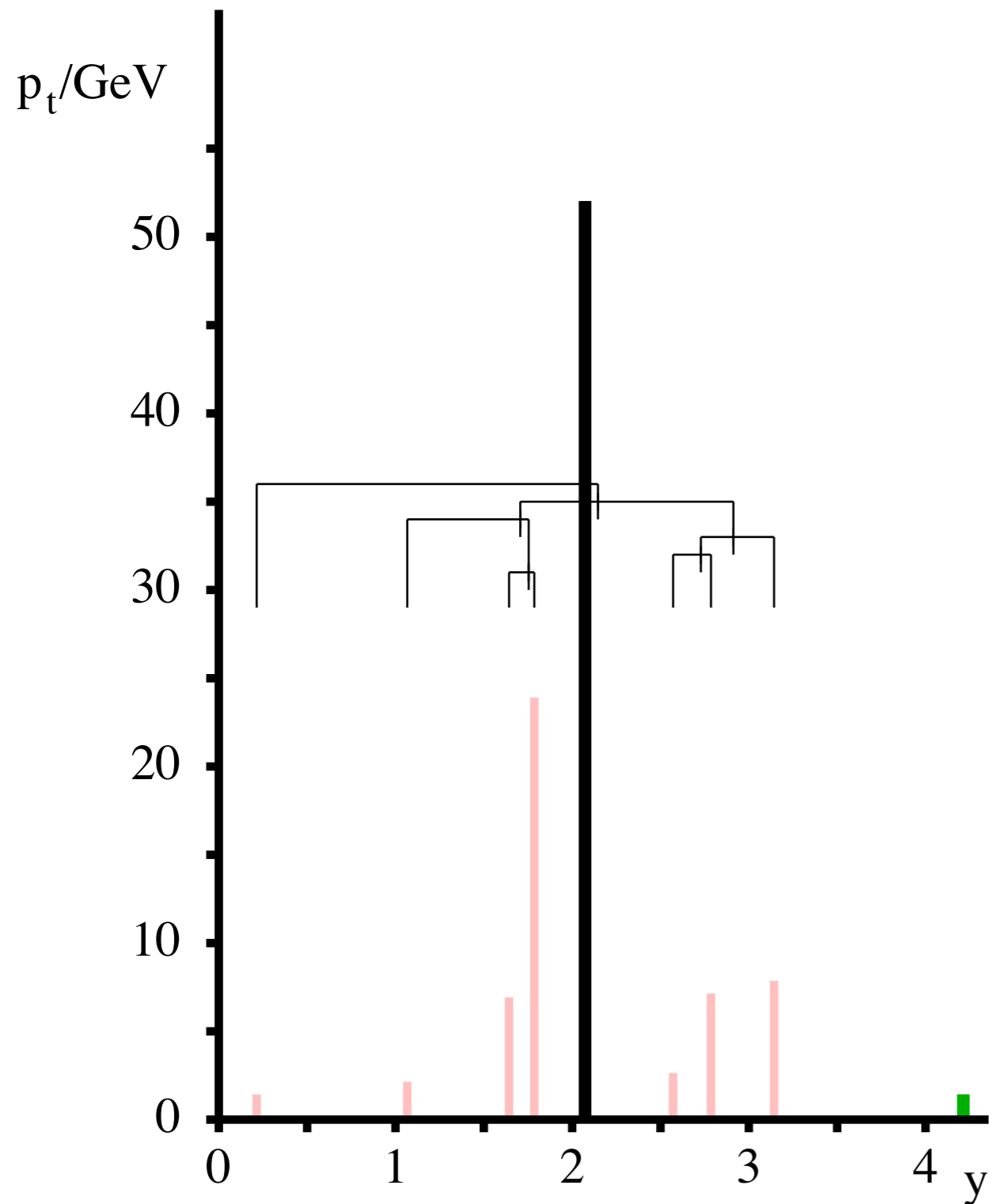
Cambridge/Aachen algorithm



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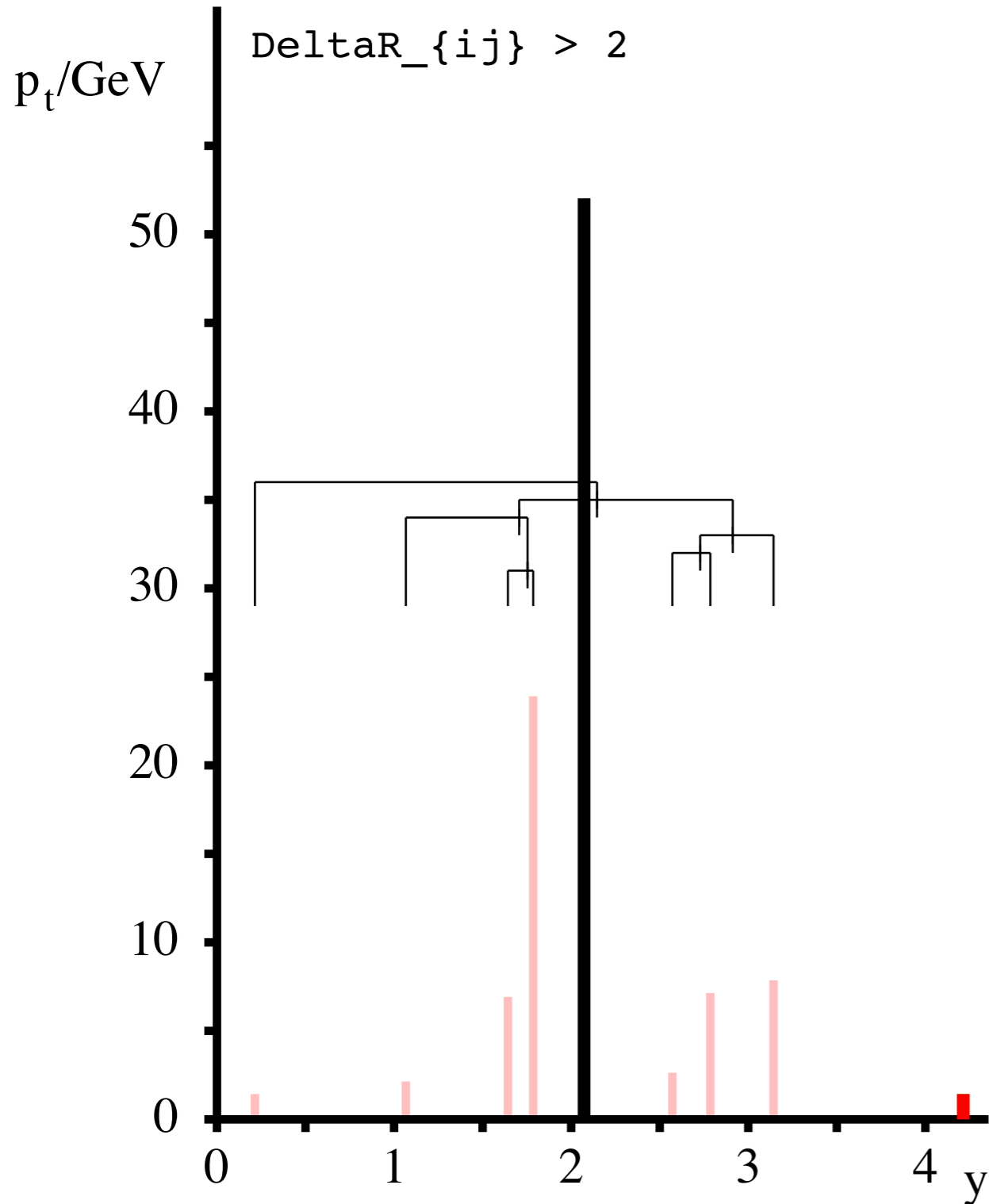
Cambridge/Aachen algorithm



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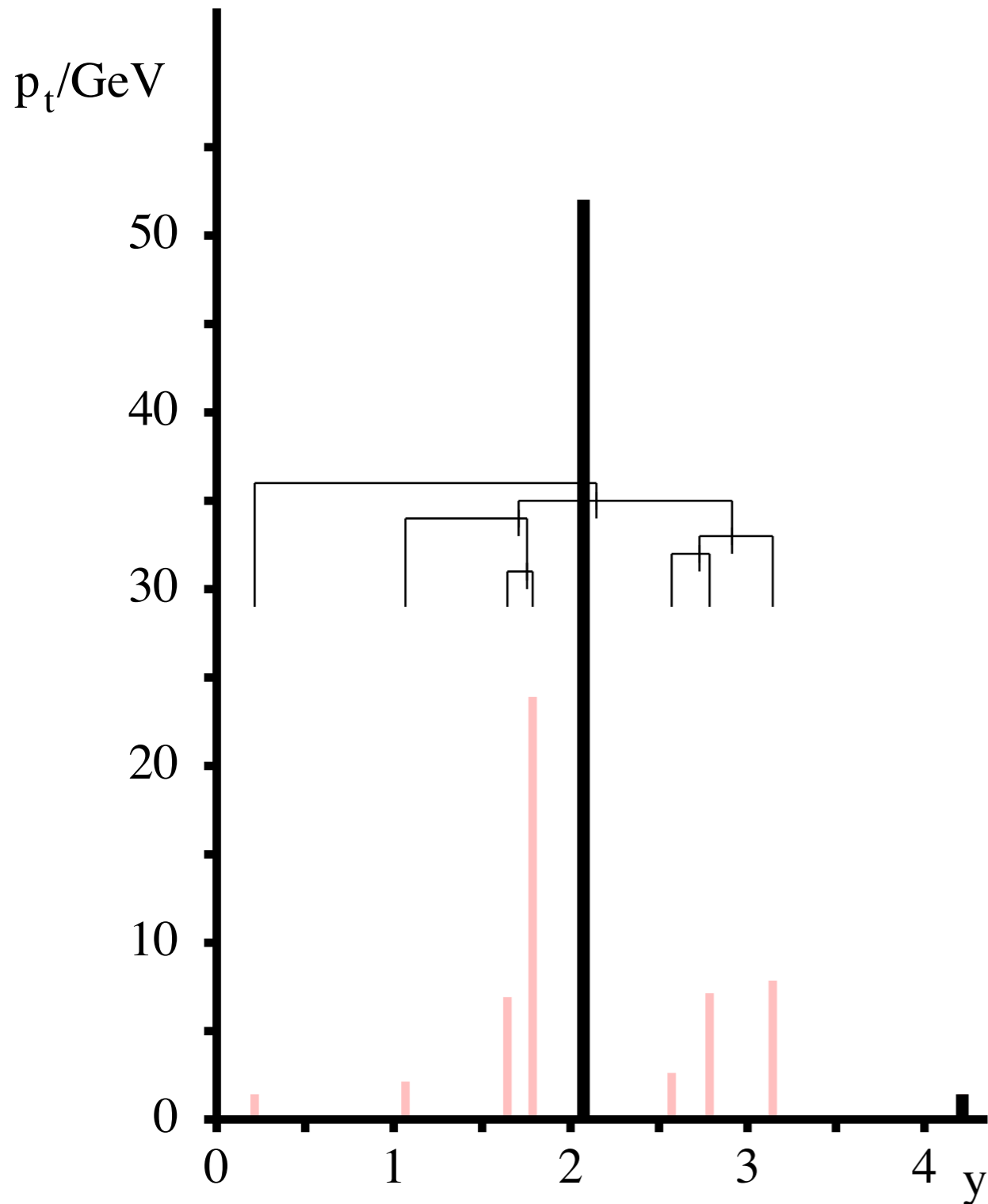
Cambridge/Aachen algorithm



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Cambridge/Aachen algorithm



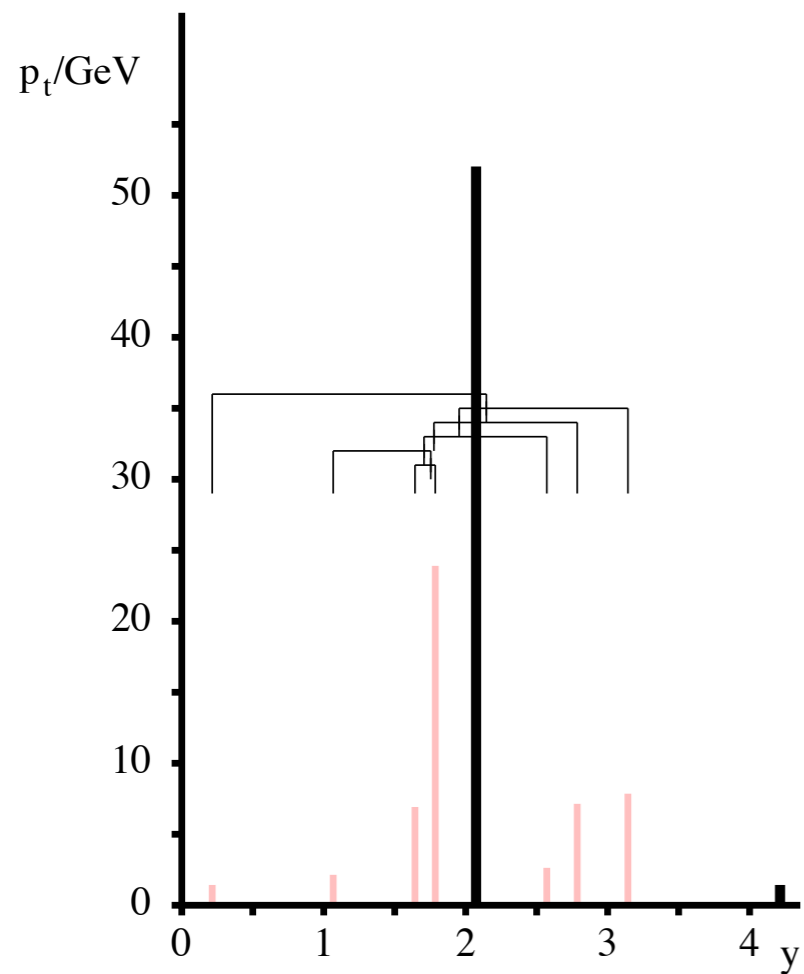
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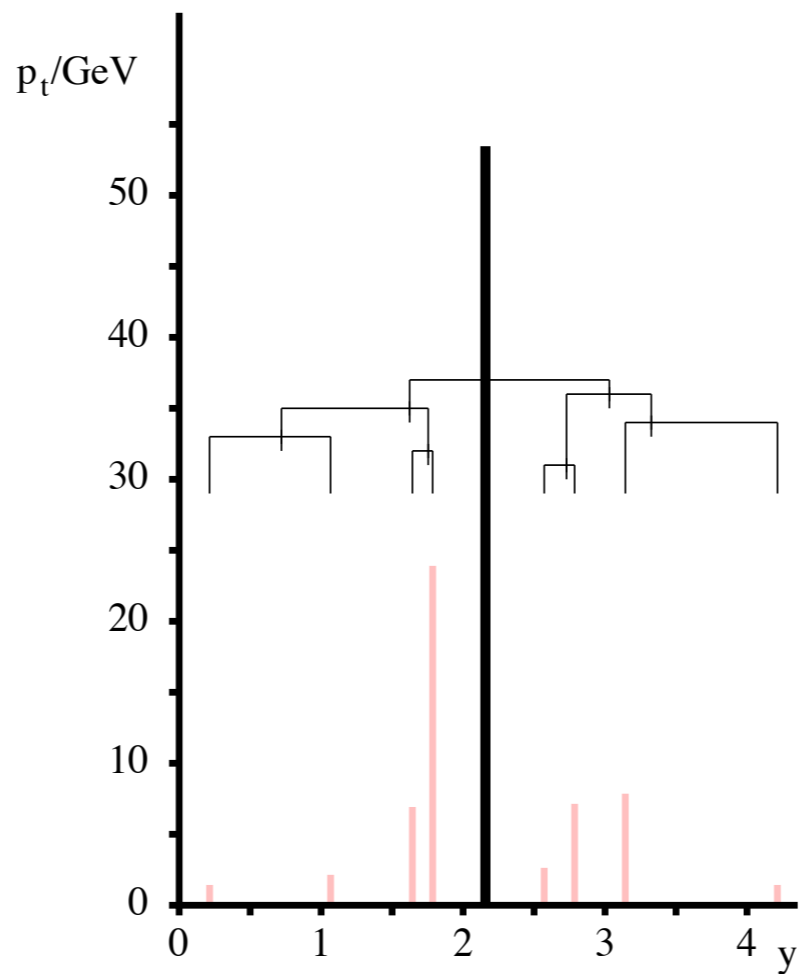
The interesting substructure is buried inside the clustering sequence — **it's less contaminated by soft junk, but needs to be pulled out with special techniques**

Butterworth, Davison, Rubin & GPS '08
Kaplan, Schwartz, Reherman & Tweedie '08
Butterworth, Ellis, Rubin & GPS '09
Ellis, Vermilion & Walsh '09

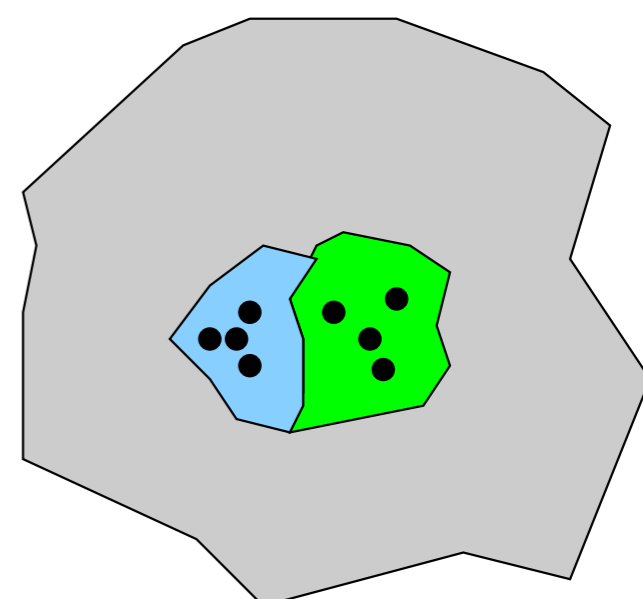
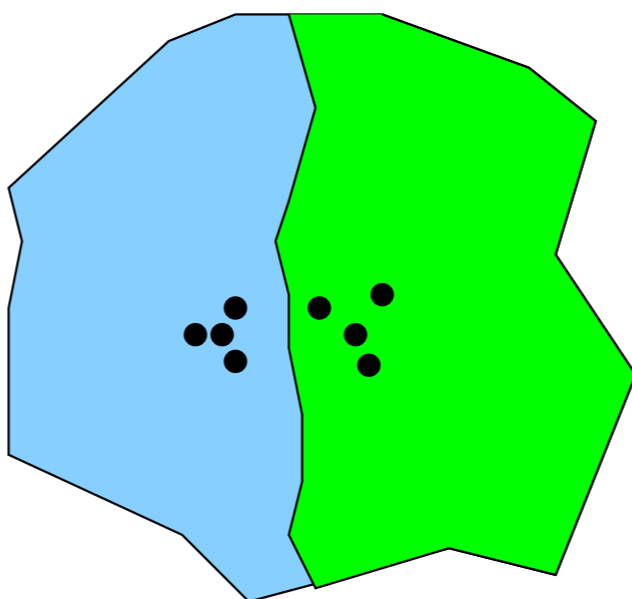
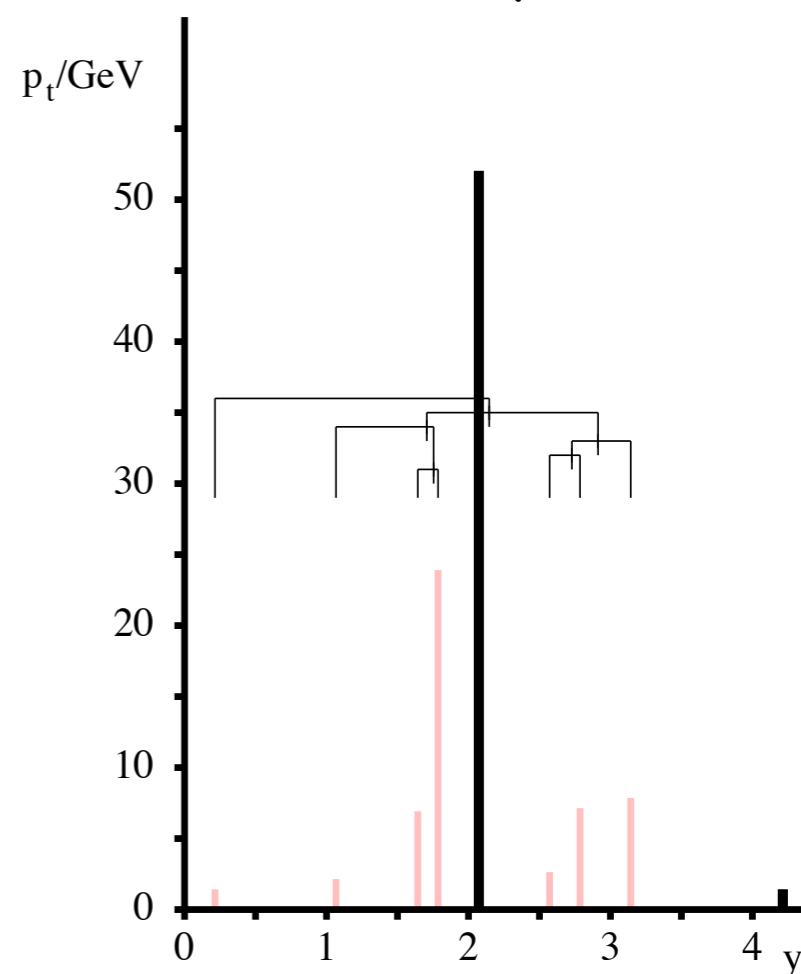
anti- k_t algorithm



k_t algorithm

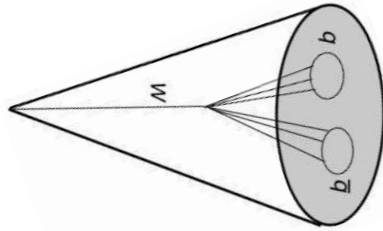


Cambridge/Aachen

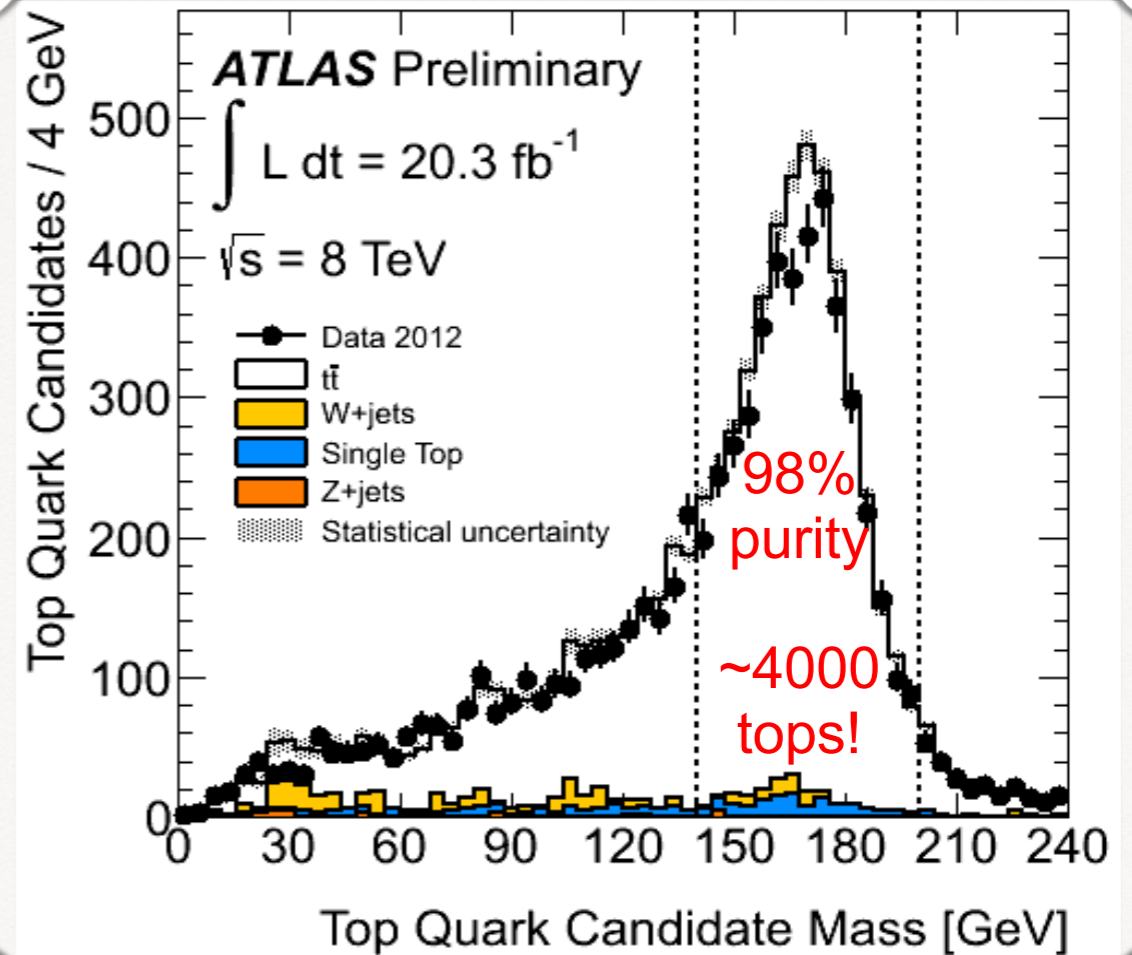
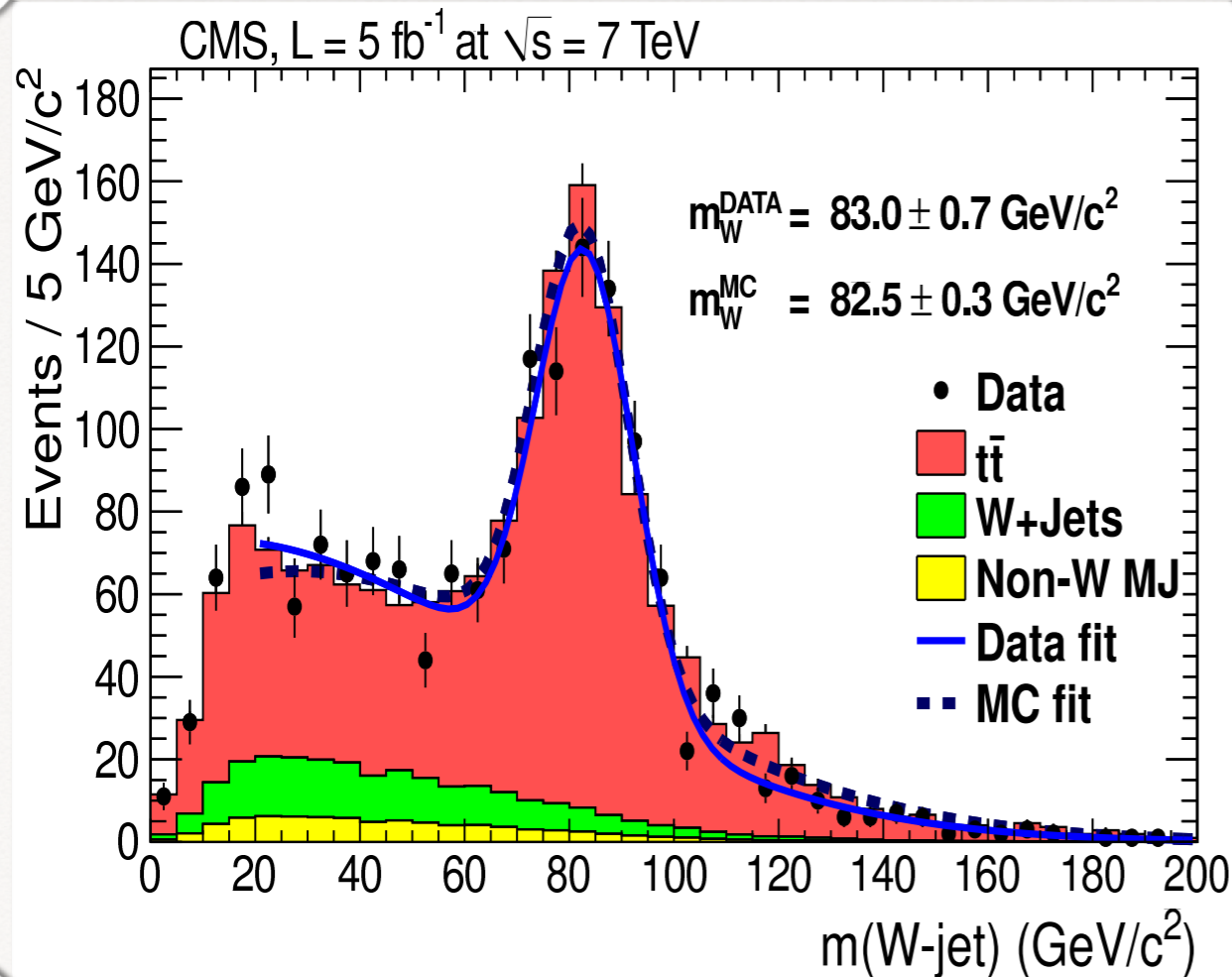
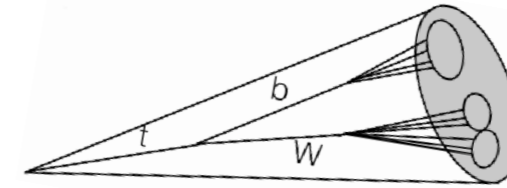


Seeing W's and tops in a single jet

W's in a single jet

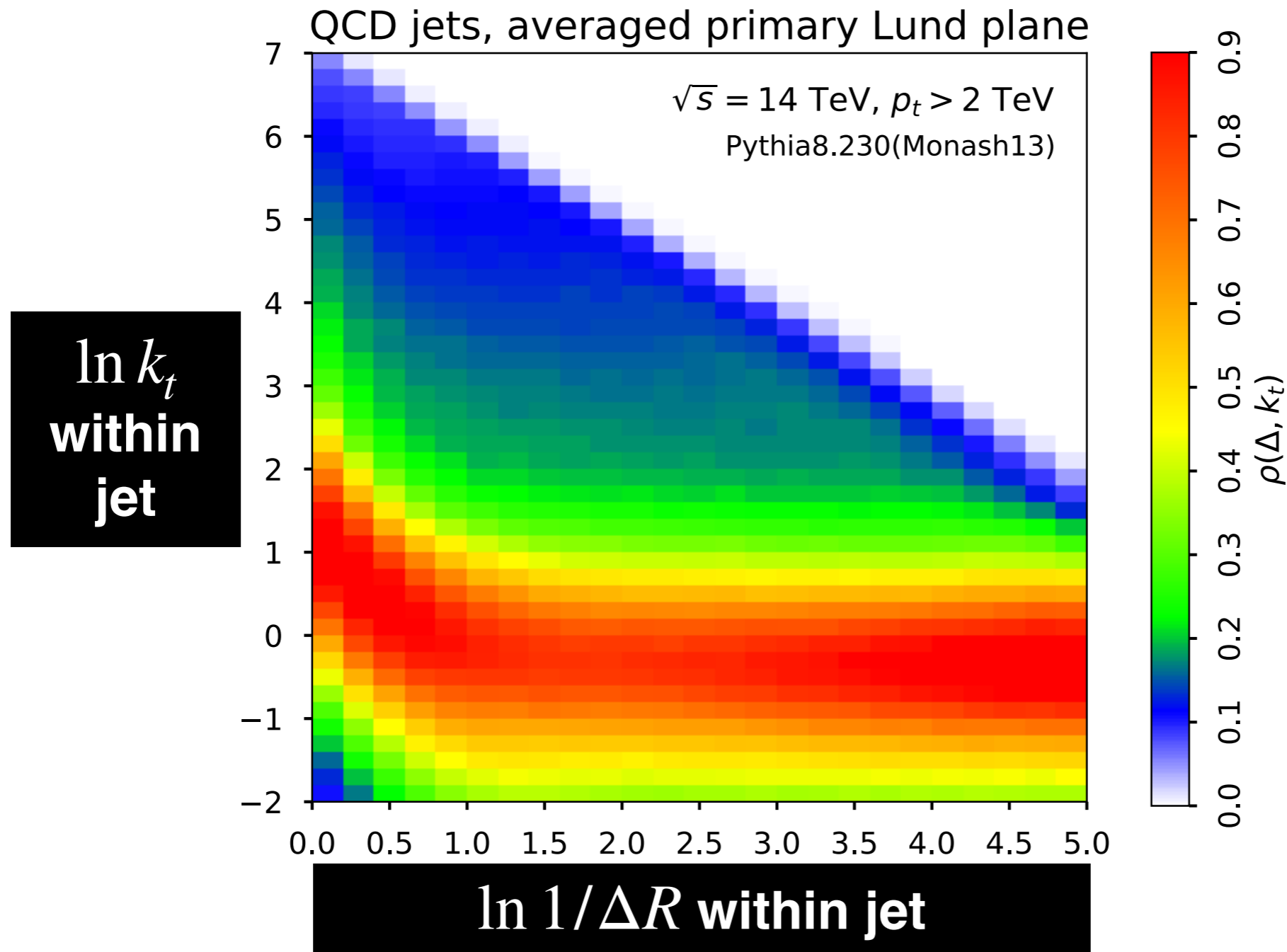


tops in a single jet



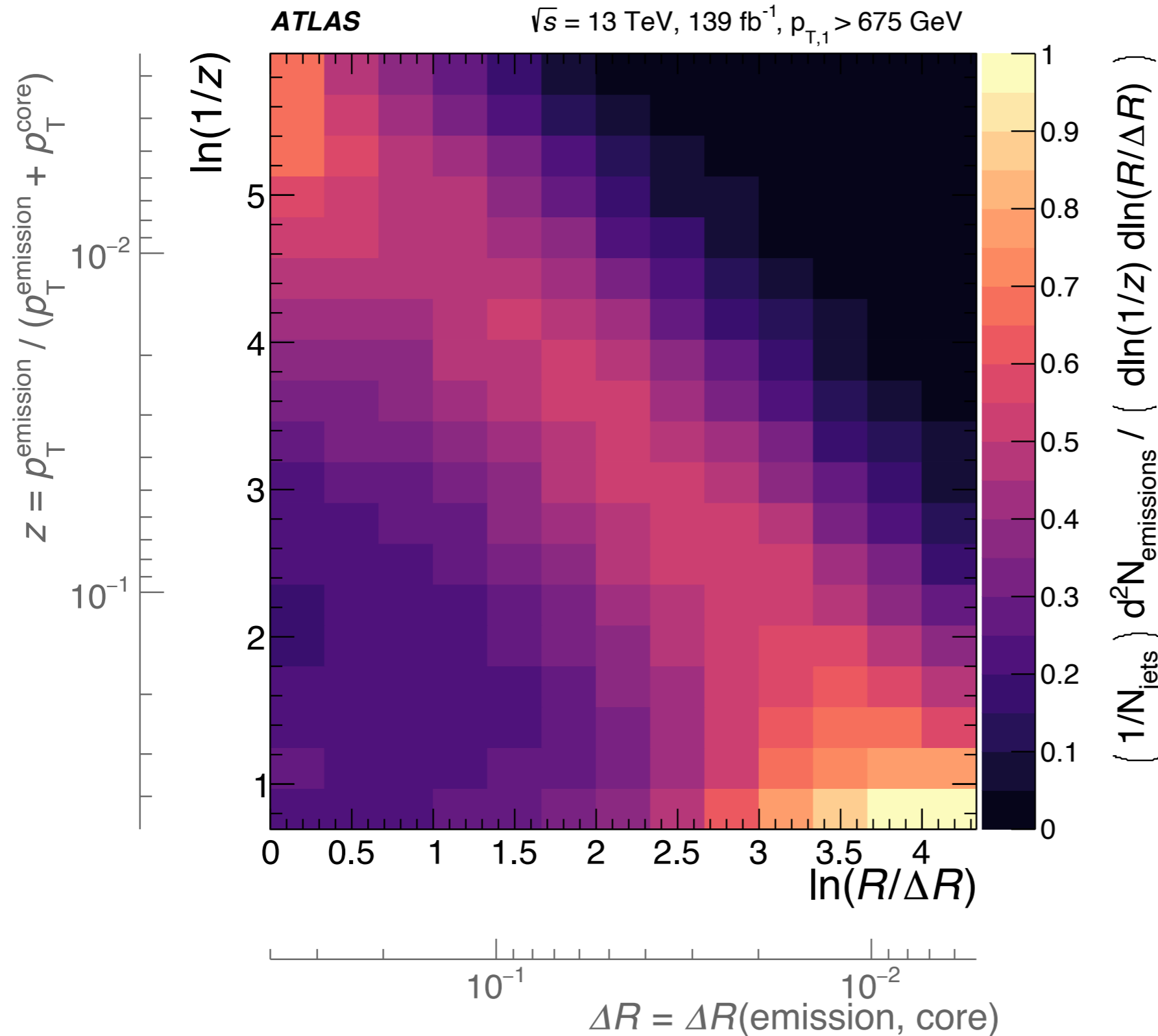
- **SoftDrop:** uses the same key ideas of C/A declustering, but with better theoretical properties and more flexibility in phasespace
- **Subjettiness / energy-energy-correlations / energy-flow polynomials / Lund Plane structure:** all try to measure the energy flow around the core n -prong structure of a jet (e.g. 2-prong for Higgs decay)
- **Machine learning:** jet substructure is one of the most dynamic playgrounds for ML, with large gains to be had in pulling out all info from jets

density of intrajet emissions in QCD jets



Dreyer, GPS & Soyez, [1807.04758](#); Lifson, GPS & Soyez, [2007.06578](#)

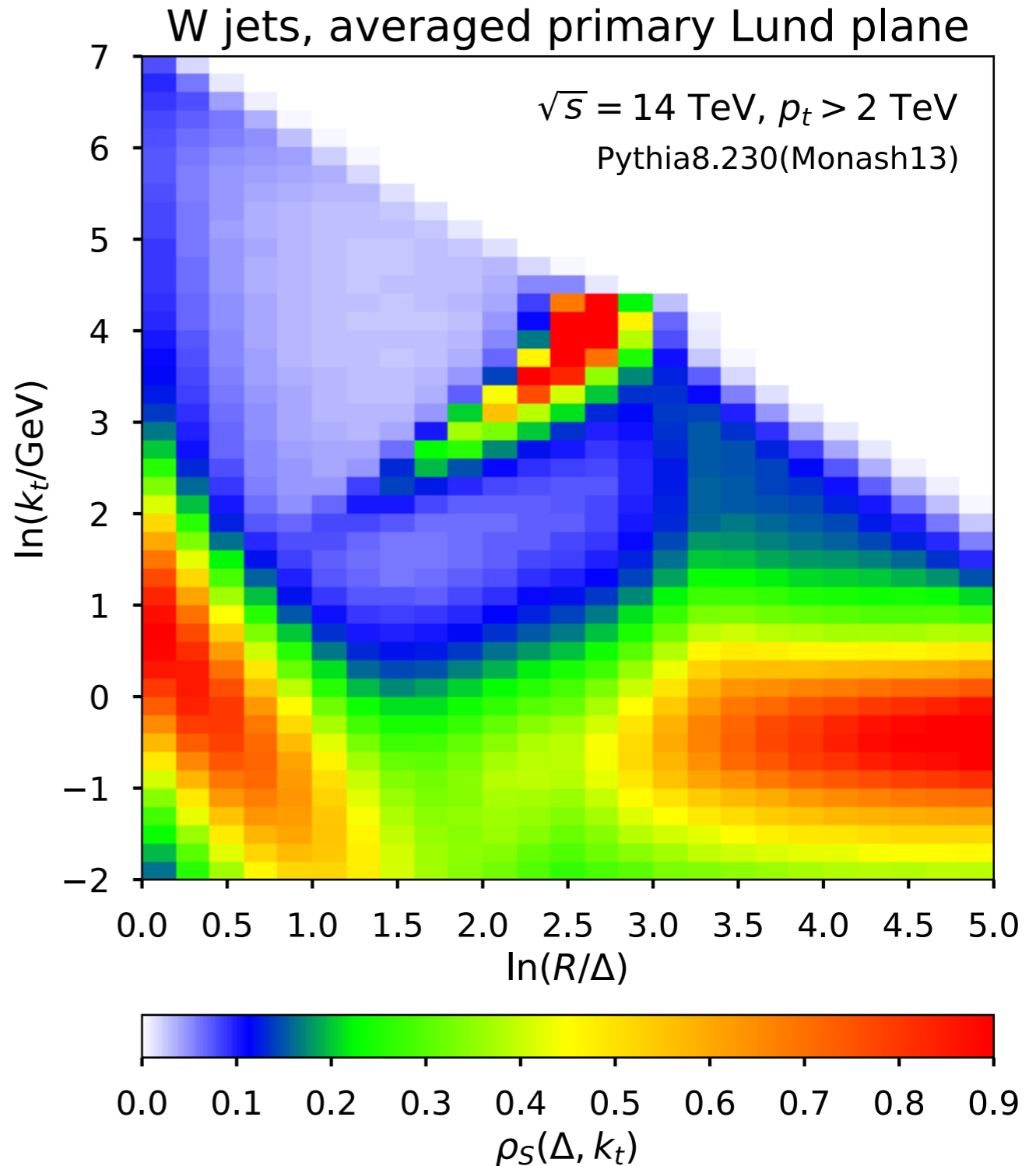
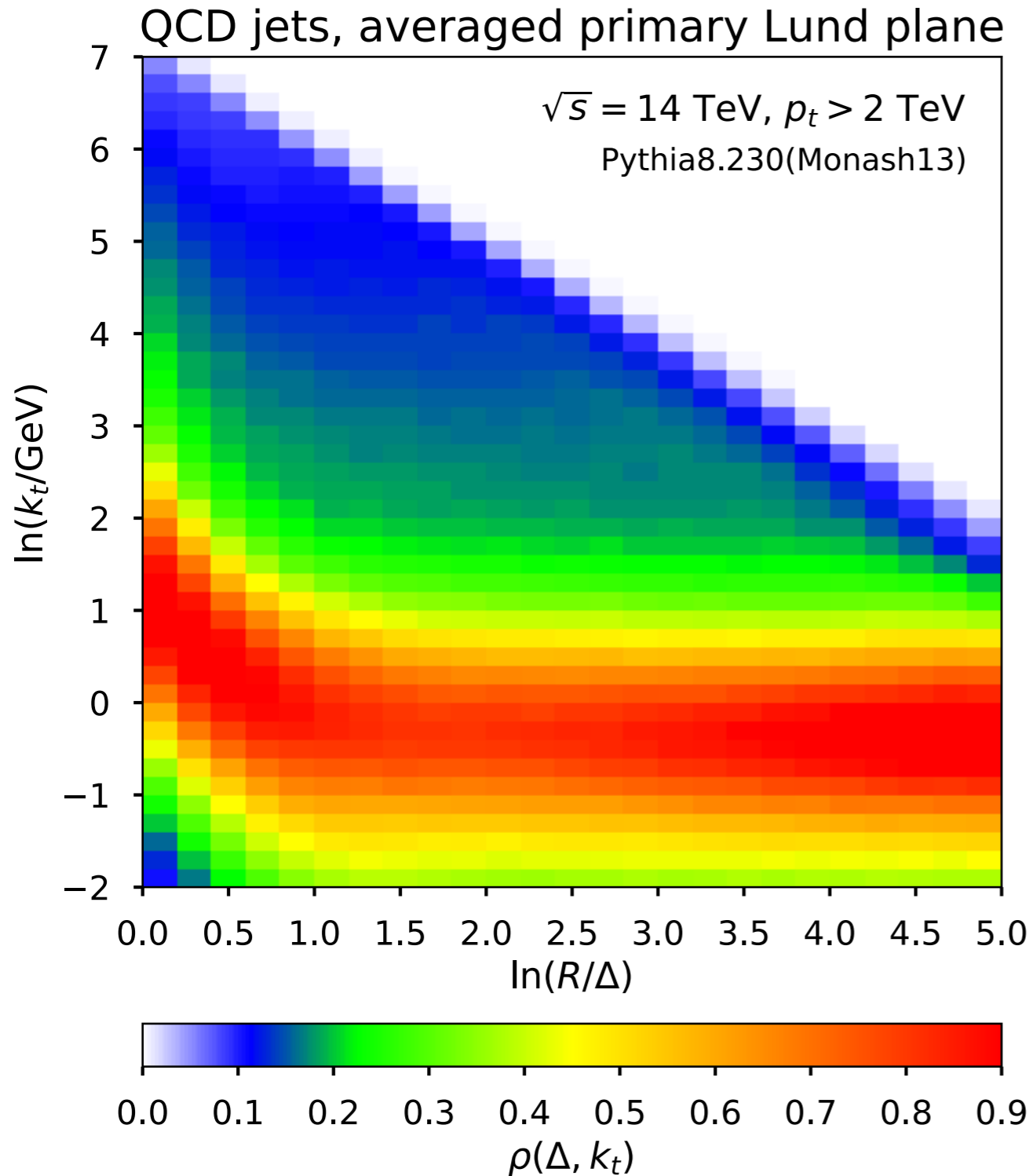
ATLAS measurement of Lund jet plane



ATLAS
2004.03540

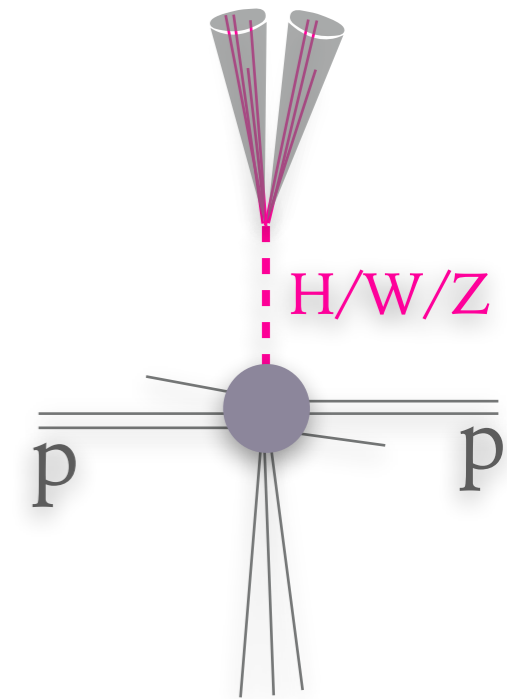
NB: vertical axis is $\ln z$ rather than $\ln k_t$

intrajet energy flow for QCD jets & W jets

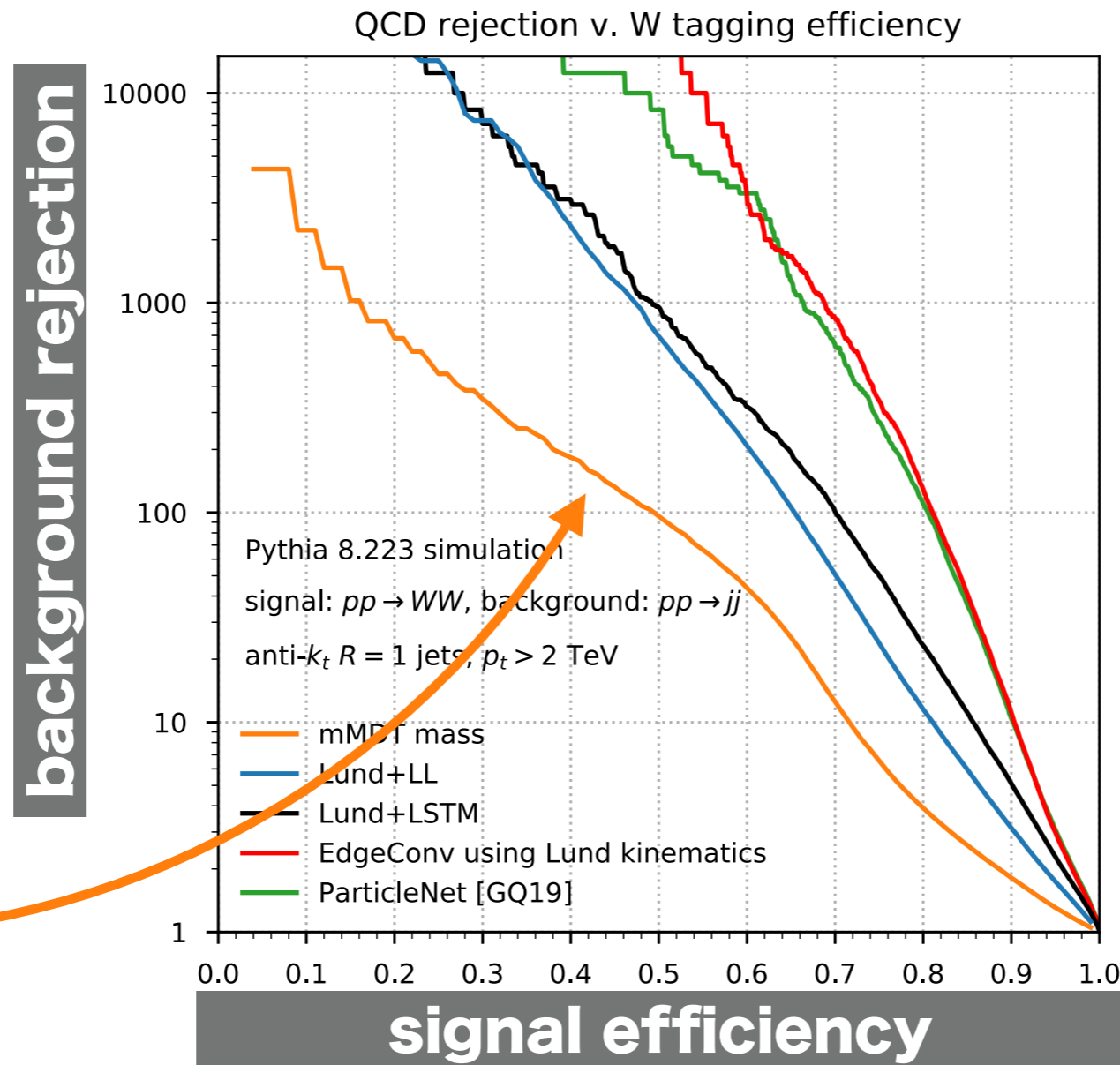


using full jet/event information for H/W/Z-boson

F. Dreyer & H. Qu, 2012.08526

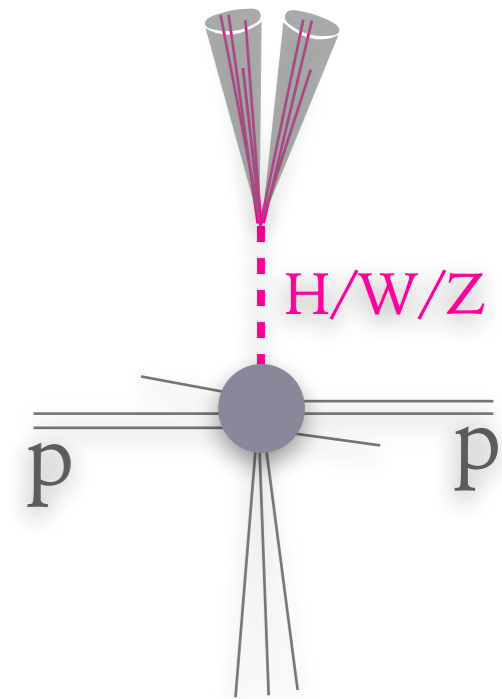


*QCD rejection
with
just jet mass
(SD/mMDT)
i.e. 2008 tools
& their
2013/14
descendants*

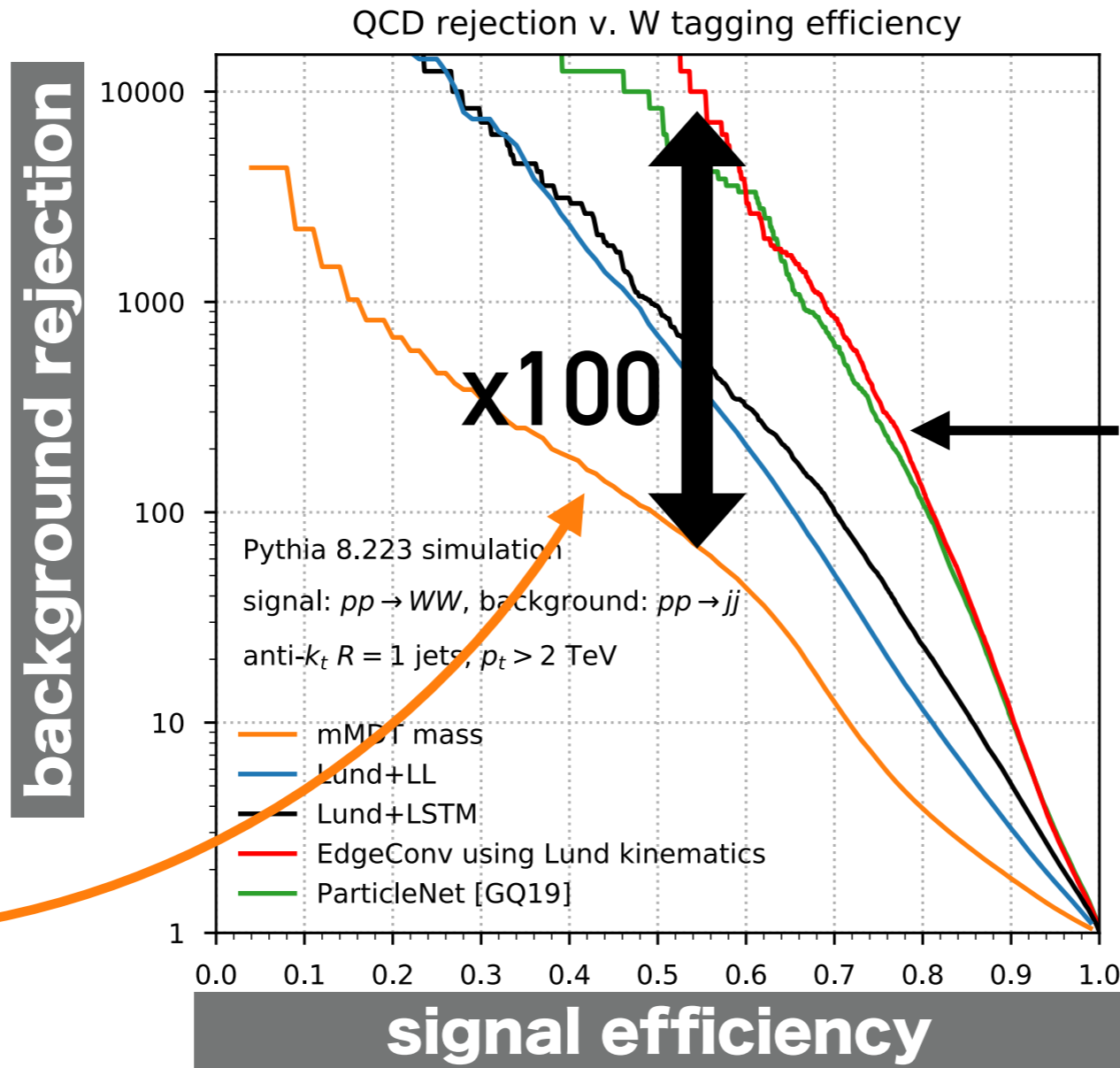


using full jet/event information for H/W/Z-boson

F. Dreyer & H. Qu, 2012.08526



QCD rejection with just jet mass (SD/mMDT) i.e. 2008 tools & their 2013/14 descendants



QCD rejection with use of full jet substructure (2019 tools) 100x better

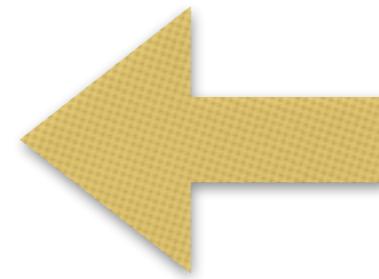
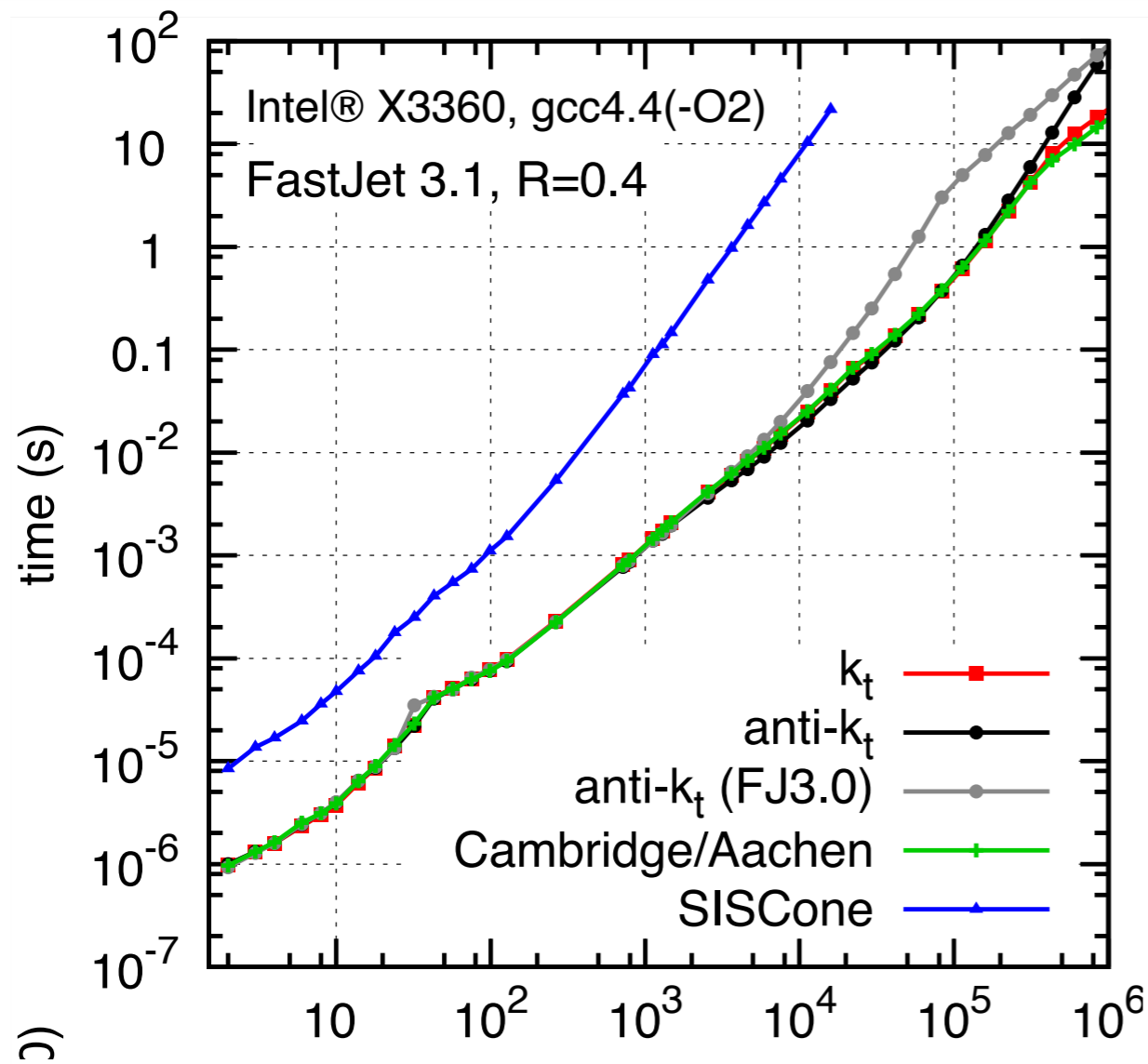
First started to be exploited by Thaler & Van Tilburg with “N-subjettiness” (2010/11)

- Jets are a consequence of the soft & collinear enhancements of gluon emission (even at small coupling), followed by hadronisation
- There are myriad approaches to jet finding
- For applications with a single moderately hard scale (e.g. $t\bar{t}$), anti-kt, $R=0.4$, with a p_t cut of a few tens of GeV is often a good default
- For problems with multiple hard scales (e.g. highly boosted top / W / H / etc.) one needs to look at events on multiple angular scales: jet substructure

- Towards Jetography, *GPS*, [0906.1833](#)
- Jet Substructure at the Large Hadron Collider: A Review of Recent Advances in Theory and Machine Learning, *Larkoski, Moult and Nachman*, [1709.04464](#)
- Jet Substructure at the Large Hadron Collider: Experimental Review, *L. Asquith et al.*, [1803.06991](#)
- Looking inside jets: an introduction to jet substructure and boosted-object phenomenology, *Marzani, Soyez and Spannowsky*, [1901.10342](#)

EXTRAS

Time to cluster N particles in FastJet



Time to cluster N particles

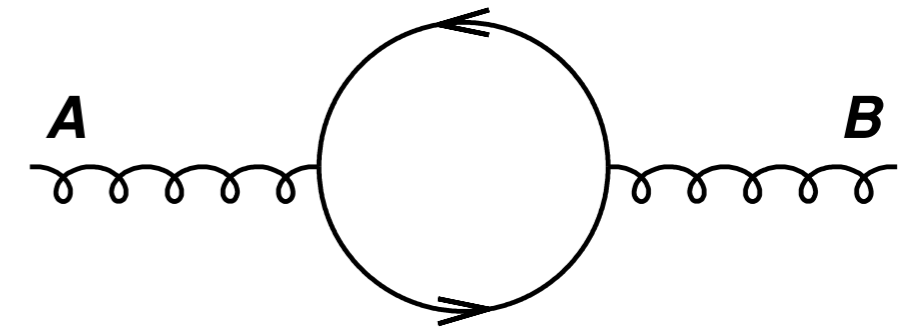
FJContrib packages

Version 1.045 of FastJet Contrib is distributed with the following packages

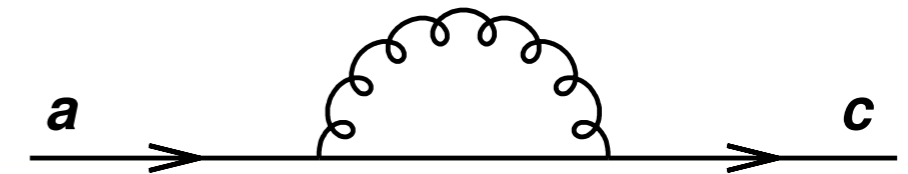
Package	Version	Release date	Information
Centauro	1.0.0	2020-08-04	README NEWS
ClusteringVetoPlugin	1.0.0	2015-05-04	README NEWS
ConstituentSubtractor	1.4.5	2020-02-23	README NEWS
EnergyCorrelator	1.3.1	2018-02-10	README NEWS
FlavorCone	1.0.0	2017-09-07	README NEWS
GenericSubtractor	1.3.1	2016-03-30	README NEWS
JetCleanser	1.0.1	2014-08-16	README NEWS
JetFFMoments	1.0.0	2013-02-07	README NEWS
JetsWithoutJets	1.0.0	2014-02-22	README NEWS
LundPlane	1.0.3	2020-02-23	README NEWS
Nsubjettiness	2.2.5	2018-06-06	README NEWS
QCDAwarePlugin	1.0.0	2015-10-08	README NEWS
RecursiveTools	2.0.0	2020-03-03	README NEWS
ScJet	1.1.0	2013-06-03	README NEWS
SoftKiller	1.0.0	2014-08-17	README NEWS
SubjetCounting	1.0.1	2013-09-03	README NEWS
ValenciaPlugin	2.0.2	2018-12-22	README NEWS
VariableR	1.2.1	2016-06-01	README NEWS

more details on soft
emission

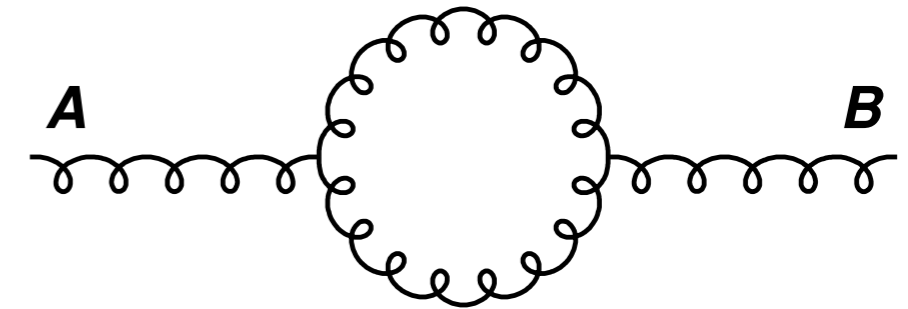
$$\text{Tr}(t^A t^B) = T_R \delta^{AB}, \quad T_R = \frac{1}{2}$$



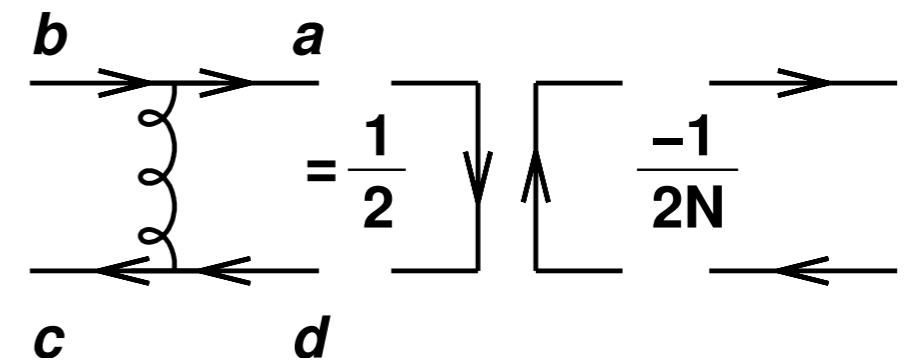
$$\sum_A t_{ab}^A t_{bc}^A = C_F \delta_{ac}, \quad C_F = \frac{N_c^2 - 1}{2N_c} = \frac{4}{3}$$



$$\sum_{C,D} f^{ACD} f^{BCD} = C_A \delta^{AB}, \quad C_A = N_c = 3$$



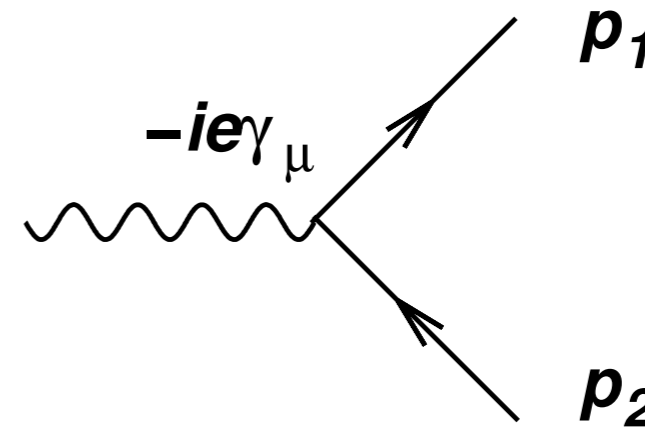
$$t_{ab}^A t_{cd}^A = \frac{1}{2} \delta_{bc} \delta_{ad} - \frac{1}{2N_c} \delta_{ab} \delta_{cd} \quad (\text{Fierz})$$



$N_c \equiv$ number of colours = 3 for QCD

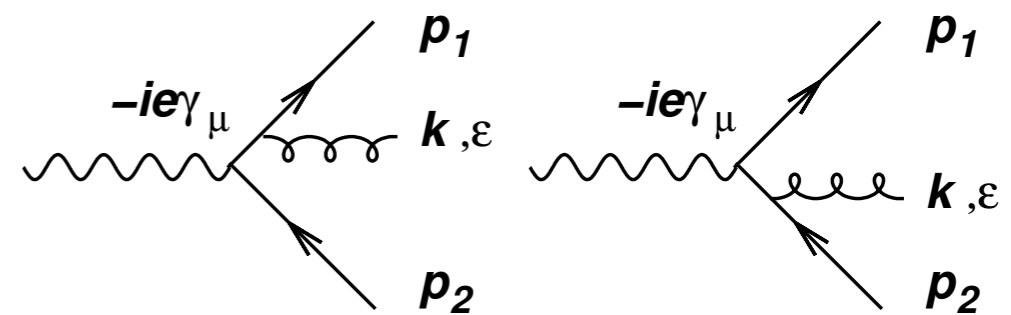
Start with $\gamma^* \rightarrow q\bar{q}$:

$$\mathcal{M}_{q\bar{q}} = -\bar{u}(p_1)ie_q\gamma_\mu v(p_2)$$



Emit a gluon:

$$\begin{aligned} \mathcal{M}_{q\bar{q}g} = & \bar{u}(p_1)ig_s\not{\epsilon}t^A \frac{i}{\not{p}_1 + \not{k}} ie_q\gamma_\mu v(p_2) \\ & - \bar{u}(p_1)ie_q\gamma_\mu \frac{i}{\not{p}_2 + \not{k}} ig_s\not{\epsilon}t^A v(p_2) \end{aligned}$$



Make gluon *soft* $\equiv k \ll p_{1,2}$; ignore terms suppressed by powers of k :

$$\mathcal{M}_{q\bar{q}g} \simeq \bar{u}(p_1)ie_q\gamma_\mu t^A v(p_2) g_s \left(\frac{p_1 \cdot \epsilon}{p_1 \cdot k} - \frac{p_2 \cdot \epsilon}{p_2 \cdot k} \right) \quad \left| \begin{array}{l} \not{p}v(p) = 0, \\ \not{p}k + k\not{p} = 2p \cdot k \end{array} \right.$$

Start with $\gamma^* \rightarrow q \bar{q}$:

$$\bar{u}(p_1) i g_s \not{\epsilon} t^A \frac{i}{\not{p}_1 + \not{k}} i e_q \gamma_\mu v(p_2) = -i g_s \bar{u}(p_1) \not{\epsilon} \frac{\not{p}_1 + \not{k}}{(p_1 + k)^2} e_q \gamma_\mu t^A v(p_2)$$

Use $\not{A} \not{B} = 2A \cdot B - \not{B} \not{A}$:

$$= -i g_s \bar{u}(p_1) [2\epsilon \cdot (p_1 + k) - (\not{p}_1 + \not{k}) \not{\epsilon}] \frac{1}{(p_1 + k)^2} e_q \gamma_\mu t^A v(p_2)$$

Use $\bar{u}(p_1) \not{p}_1 = 0$ and $k \ll p_1$ (p_1, k massless)

$$\simeq -i g_s \bar{u}(p_1) [2\epsilon \cdot p_1] \frac{1}{(p_1 + k)^2} e_q \gamma_\mu t^A v(p_2)$$

$$= -i g_s \frac{p_1 \cdot \epsilon}{p_1 \cdot k} \underbrace{\bar{u}(p_1) e_q \gamma_\mu t^A v(p_2)}_{\text{pure QED spinor structure}}$$

$$\begin{aligned}
 |M_{q\bar{q}g}^2| &\simeq \sum_{A,\text{pol}} \left| \bar{u}(p_1) i e_q \gamma_\mu t^A v(p_2) g_s \left(\frac{p_1 \cdot \epsilon}{p_1 \cdot k} - \frac{p_2 \cdot \epsilon}{p_2 \cdot k} \right) \right|^2 \\
 &= -|M_{q\bar{q}}^2| C_F g_s^2 \left(\frac{p_1}{p_1 \cdot k} - \frac{p_2}{p_2 \cdot k} \right)^2 = |M_{q\bar{q}}^2| C_F g_s^2 \frac{2p_1 \cdot p_2}{(p_1 \cdot k)(p_2 \cdot k)}
 \end{aligned}$$

Include phase space:

$$d\Phi_{q\bar{q}g} |M_{q\bar{q}g}^2| \simeq (d\Phi_{q\bar{q}} |M_{q\bar{q}}^2|) \frac{d^3 \vec{k}}{2E(2\pi)^3} C_F g_s^2 \frac{2p_1 \cdot p_2}{(p_1 \cdot k)(p_2 \cdot k)}$$

Note property of **factorisation** into hard $q\bar{q}$ piece and **soft-gluon emission piece, dS** .

$$dS = EdE d\cos\theta \frac{d\phi}{2\pi} \cdot \frac{2\alpha_s C_F}{\pi} \frac{2p_1 \cdot p_2}{(2p_1 \cdot k)(2p_2 \cdot k)}$$

$$\begin{aligned}
 \theta &\equiv \theta_{p_1 k} \\
 \phi &= \text{azimuth}
 \end{aligned}$$

