CAESAR: Computer automated resummations

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work done with Andrea Banfi & Giulia Zanderighi

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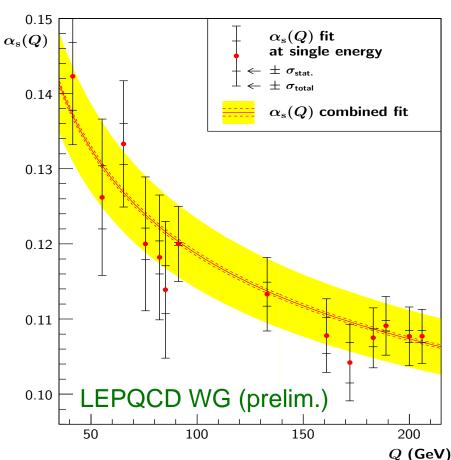
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Measurements of the coupling α_s and its renormalization group running



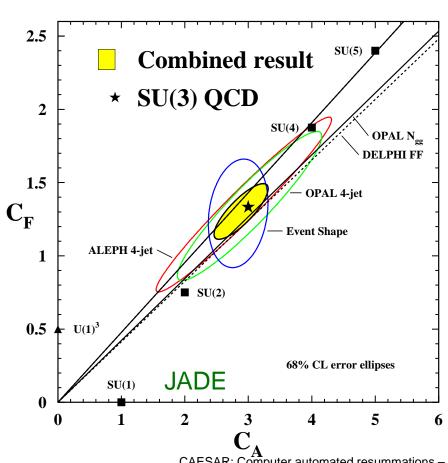
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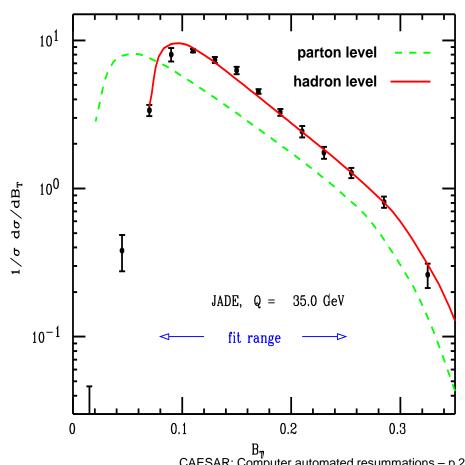


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Provide a wealth of information, e.g.:

- Measurements of the coupling α_s and its renormalization group running
- Measurements/cross checks of the values of the colour factors of QCD
- Studies of connection between parton-level (QCD calculations of quarks and gluons) and (the real) hadron-level



Extending scope of these studies

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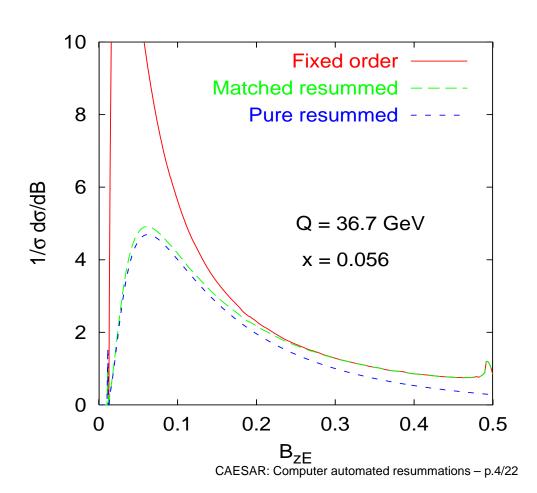
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- Studies of event-shapes for three (2+1) jet events is vital for testing non-perturbative (NP) models
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- A wide range of studies are possible in hadron-hadron colliders (dijet events, W,Z + jet)
 - ullet analogous studies to those in e^+e^- and DIS
 - purely analytical studies of underlying event and minium bias properties

Need the theoretical tools

- Account for hard gluon/quark radiation (fixed order)
 - ✓ Programs calculate NLO predictions given subroutine for observable
 - ✓ Many processes: Event2, Disent, Disaster++, NLOJET++, MCFM, . . .
 - distributions diverge in soft/collinear region
- Account for multiple soft/collinear radiation (NLL resummation)
 - Ensures sensible distribution
 - Has to be done by hand for each observable & process



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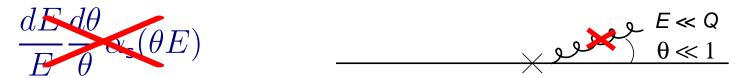
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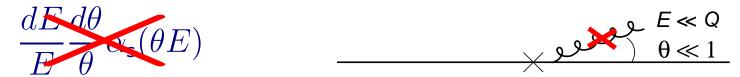
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$$\frac{dE}{E}\frac{d\theta}{\theta} \propto (\theta E) \qquad \qquad \times \underbrace{\theta \in \mathbb{Q}}_{\theta \ll 1}$$

virtual corrections are unaffected

$$-\frac{dE}{E}\frac{d\theta}{\theta}\alpha_{s}(\theta E) \qquad \qquad \underbrace{E \ll Q}_{\theta \ll 1}$$

Imbalance leads to large logarithms in distribution of event shape:

$$\operatorname{Prob}(1 - T < \tau) \simeq 1 - \frac{\alpha_{\mathsf{s}} C_F}{2\pi} \ln^2 \tau + \dots \tag{} \tau \ll 1)$$

Resummation: Large Logarithms at all orders

There is a soft and a collinear divergence (\mapsto logs) for each emitted gluon.

At all orders, probability of event being two-jet like has *poorly convergent perturbation series:*

$$P(1 - T < \tau) \equiv \Sigma(\tau) = 1 + \sum_{n=1}^{\infty} R_{n,2n} \alpha_s^n \ln^{2n} \tau + \dots$$

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Today's state of the art involves *exponentiation* and resummation of *Leading Logs* (LL) and *Next-to-Leading Logs* (NLL):

$$\Sigma(\tau) \simeq \exp\left[\sum_{n=1}^{\infty} \left(\underbrace{G_{n,n+1} \alpha_{s}^{n} \ln^{n+1} \tau}_{\text{LL}} + \underbrace{G_{n,n} \alpha_{s}^{n} \ln^{n} \tau}_{\text{NLL}} + \cdots\right)\right]$$

NB: $\alpha_s^n \ln^{2n} \tau$ in Σ , but only $\alpha_s^n \ln^{n+1} \tau$ in exponent.

The analytical resummation industry...

$e^+e^- \rightarrow 2 \text{ jets}$

- S. Catani, G. Turnock, B. R. Webber and L. Trentadue, *Thrust distribution* in e^+e^- annihilation, Phys. Lett. B **263** (1991) 491.
- S. Catani, G. Turnock and B. R. Webber, *Heavy jet mass distribution in* e^+e^- *annihilation*, Phys. Lett. B **272** (1991) 368.
- S. Catani et al., New clustering algorithm for multi-jet cross-sections in e^+e^- annihilation, Phys. Lett. B **269** (1991) 432.
- S. Catani, L. Trentadue, G. Turnock and B. R. Webber, *Resummation of large logarithms* [...] , Nucl. Phys. B **407** (1993) 3.
- S. Catani, G. Turnock and B. R. Webber, *Jet broadening measures in* e^+e^- *annihilation*, Phys. Lett. B **295** (1992) 269.
- G. Dissertori and M. Schmelling, An Improved theoretical prediction for the two jet rate in e^+e^- annihilation, Phys. Lett. B **361** (1995) 167.
- Y. L. Dokshitzer, A. Lucenti, G. Marchesini and GPS, *On the QCD analysis of jet broadening*, JHEP **9801** (1998) 011
- S. Catani and B. R. Webber, Resummed C-parameter distribution in e^+e^- annihilation, Phys. Lett. B **427** (1998) 377
- S. J. Burby and E. W. Glover, *Resumming the light hemisphere mass and* [...], JHEP **0104** (2001) 029
- M. Dasgupta and GPS, Resummation of non-global QCD observables, Phys. Lett. B **512** (2001) 323
- C. F. Berger, T. Kucs and G. Sterman, *Event shape / energy flow correlations*, Phys. Rev. D **68** (2003) 014012
- E. Gardi and J. Rathsman, *Renormalon resummation and exponentiation of soft and collinear gluon radiation in the thrust distribution*, Nucl. Phys. B **609** (2001) 123
- E. Gardi and J. Rathsman, *The thrust and heavy-jet mass distributions in the two-jet region*, Nucl. Phys. B **638** (2002) 243

E. Gardi and L. Magnea, *The C parameter distribution in e+ e- annihilation*, JHEP **0308** (2003) 030

F. Krauss and G. Rodrigo, Resummed jet rates for e^+e^- annihilation into massive quarks, Phys. Lett. B **576** (2003) 135

DIS 1+1 jet

V. Antonelli, M. Dasgupta and GPS, *Resummation of thrust distributions in DIS*, JHEP **0002** (2000) 001

M. Dasgupta and GPS, Resummation of the jet broadening in DIS, Eur. Phys. J. C **24** (2002) 213

M. Dasgupta and GPS, Resummed event-shape variables in DIS, JHEP **0208** (2002) 032

e^+e^- , DY, DIS 3 jets

A. Banfi , G. Marchesini, Y. L. Dokshitzer and G. Zanderighi, *QCD analysis of near-to-planar 3-jet events*, JHEP **0007** (2000) 002

A. Banfi, Y. L. Dokshitzer, G. Marchesini and G. Zanderighi, *Near-to-planar 3-jet events in and beyond QCD perturbation theory*, Phys. Lett. B **508** (2001) 269

A. Banfi, Y. L. Dokshitzer, G. Marchesini and G. Zanderighi, *QCD analysis* of *D-parameter in near-to-planar three-jet events*, JHEP **0105** (2001) 040

A. Banfi , G. Marchesini, G. Smye and G. Zanderighi, *Out-of-plane QCD radiation in hadronic Z0 production*, JHEP **0108** (2001) 047

A. Banfi , G. Marchesini, G. Smye and G. Zanderighi, *Out-of-plane QCD radiation in DIS with high p(t) jets*, JHEP **0111** (2001) 066

A. Banfi , G. Marchesini and G. Smye, *Azimuthal correlation in DIS*, JHEP **0204** (2002) 024

Average: 1 observable per article

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 - Only know if program is suitable for observable if you've already done most of the resummation...
 - Matching with fixed order (LO, NLO, NNLO) is complex

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Computer Automated Expert Semi-Analytical Resummation (CAESAR)

$$V(\lbrace p\rbrace, k) = d_{\ell} \left(\frac{k_{t}}{Q}\right)^{a_{\ell}} e^{-b_{\ell}\eta} g_{\ell}(\phi) .$$

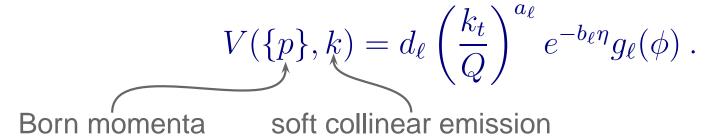
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 Born momenta soft collinear emission

Observable must have standard functional form for soft & collinear gluon emission

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- Require continuous globalness, i.e. uniform dependence on k_t independently of emission direction ($a_1 = a_2 = \cdots = a$)

Example: thrust in e^+e^-

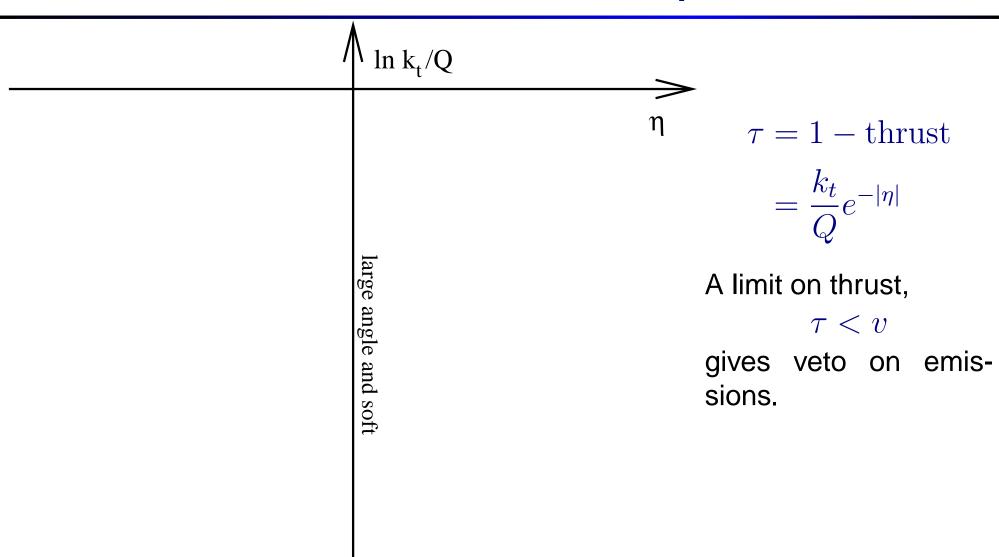
$$\tau = 1 - \text{thrust}$$
$$= \frac{k_t}{Q} e^{-|\eta|}$$

A limit on thrust,

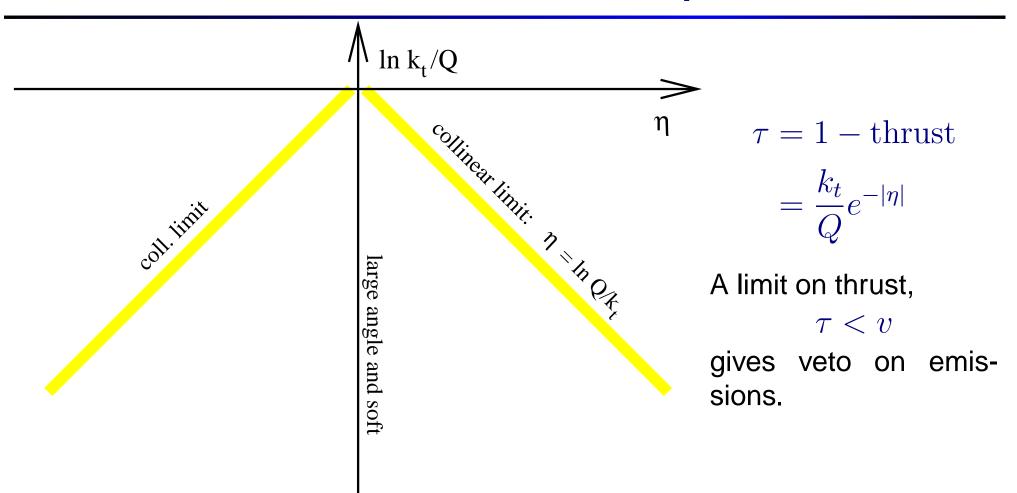
$$\tau < v$$

gives veto on emissions.

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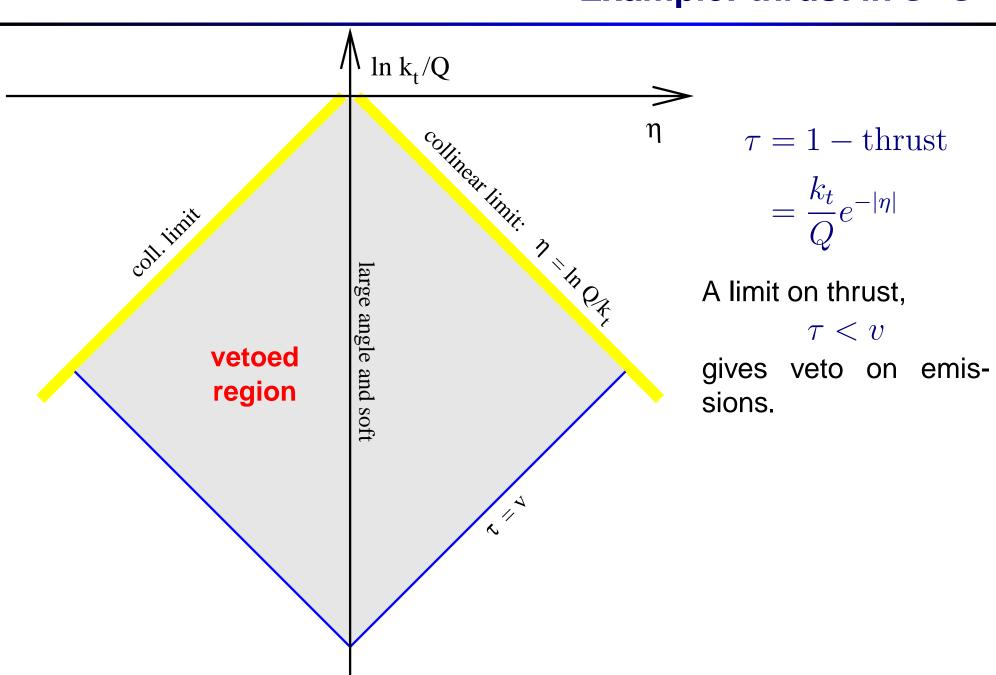


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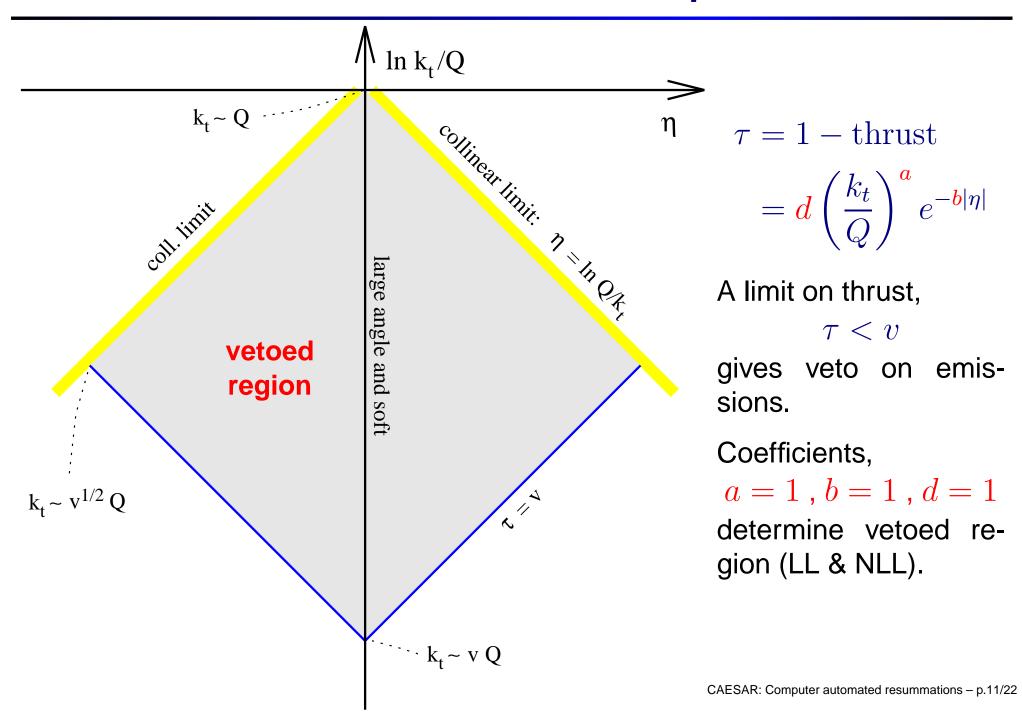


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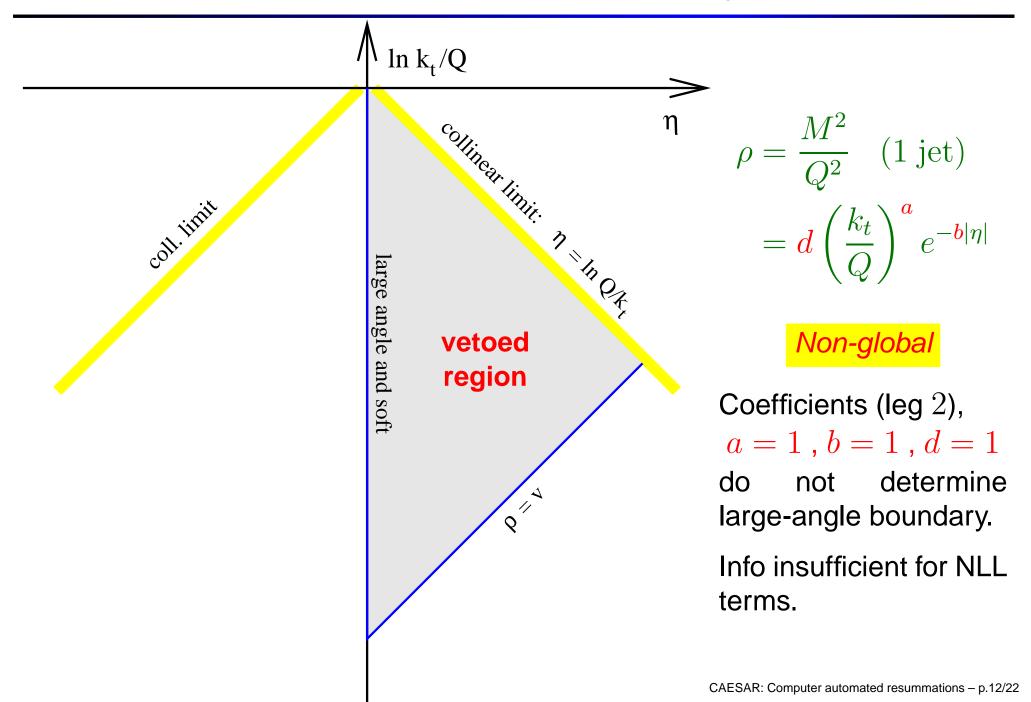
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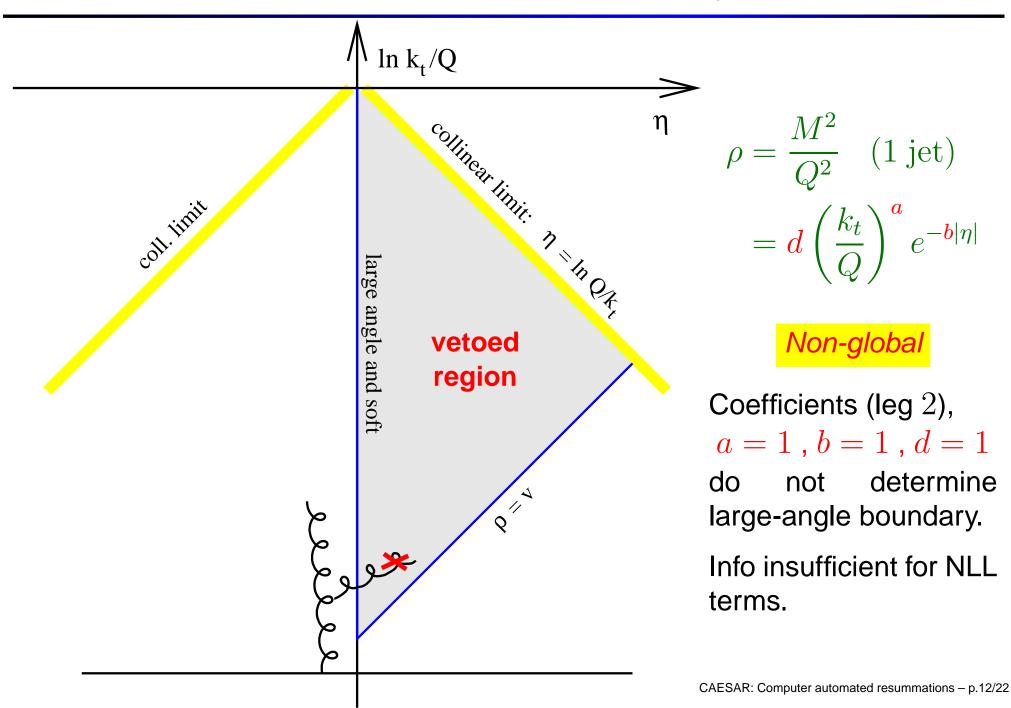
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Not 'supported': invariant jet mass in e^+e^-



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 $\kappa(v)$ is any momentum such that $V(\{p\},\kappa(v))=v$

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$$\lim_{v \to 0} \frac{1}{v} V(\{p\}, \kappa_1(\zeta_1 v), \kappa_2(\zeta_2 v), \ldots) = f(\zeta_1, \zeta_2, \ldots)$$

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• Require *variant* of infrared-collinear safety for $f(\zeta_1, \zeta_2, \ldots)$:

$$\lim_{\substack{\zeta_n \to 0}} f(\zeta_1, \zeta_2, \dots, \zeta_{n-1}, \underline{\zeta_n}) = f(\zeta_1, \zeta_2, \dots, \zeta_{n-1})$$

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Or:
$$\left[\lim_{v\to 0}, \lim_{\zeta_n\to 0}\right] \frac{1}{v} V(\{p\}, \kappa_1(\zeta_1 v), \kappa_2(\zeta_2 v), \dots, \kappa_n(\zeta_n v)) = 0$$

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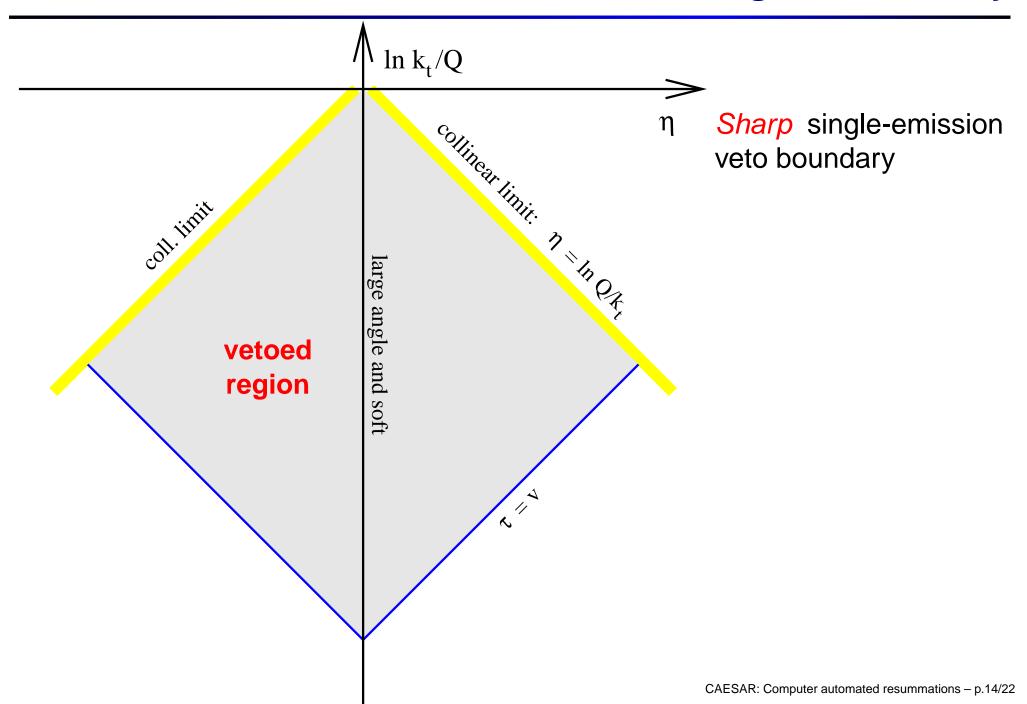
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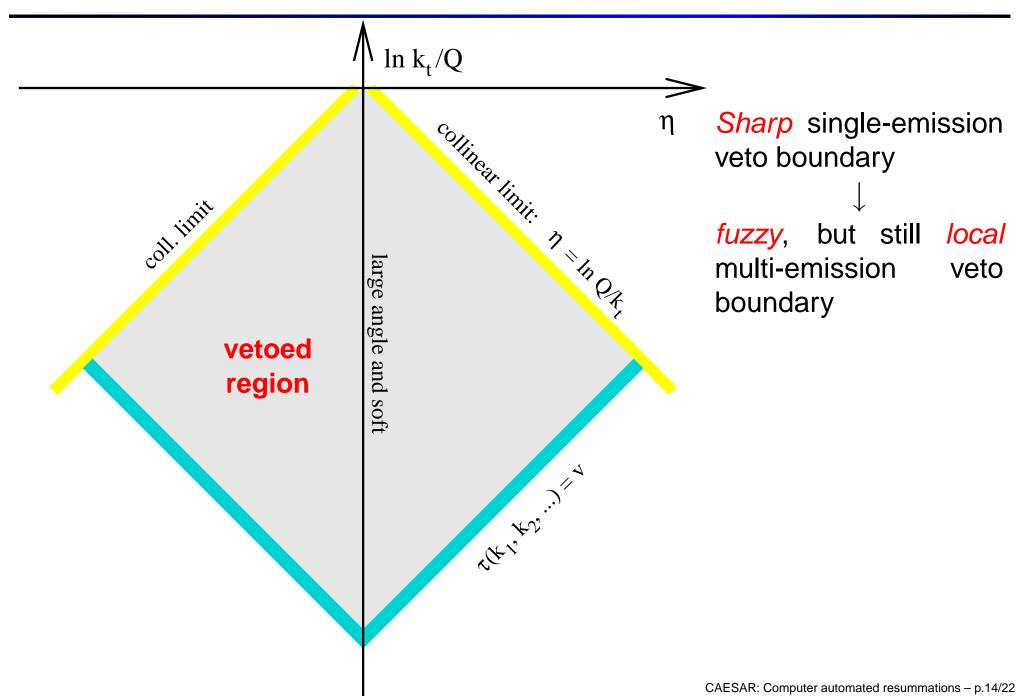
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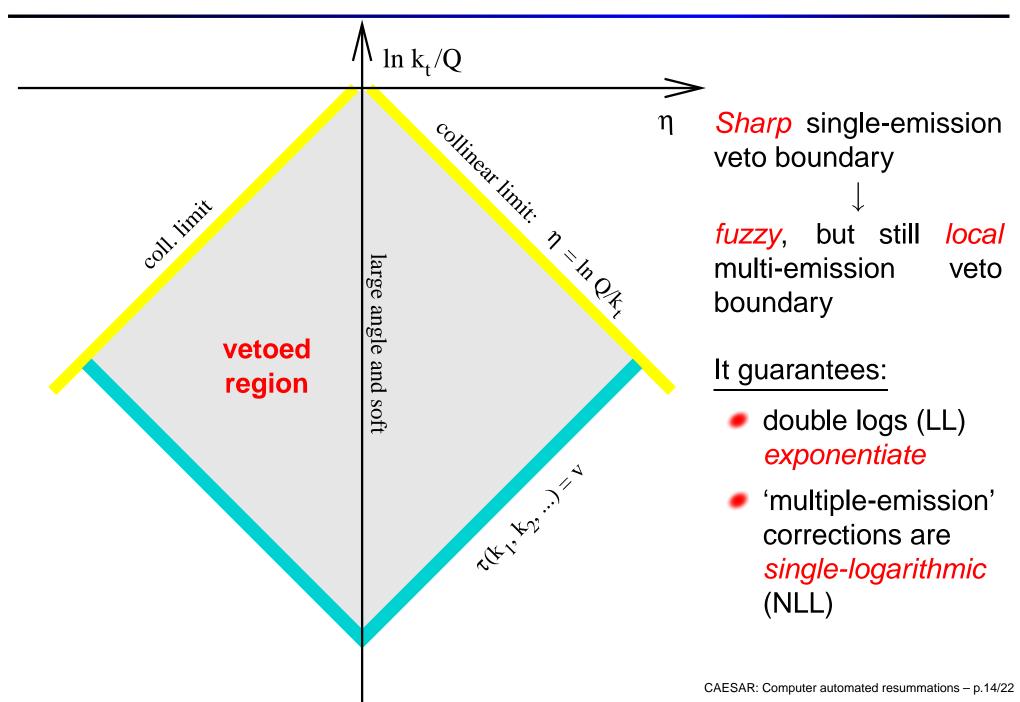
Meaning of rIRC safety



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Multi-emission single logs

Determine function \mathcal{F} that accounts for average extra veto to be applied to exponentiated double logs.

Given by:

$$\mathcal{F} = \lim_{\epsilon \to 0} \frac{\epsilon^{R'}}{R'} \sum_{m=0}^{\infty} \frac{1}{m!} \left(\prod_{i=1}^{m+1} \sum_{\ell_i=1}^n C_{\ell} r'_{\ell_i} \int_{\epsilon}^1 \frac{d\zeta_i}{\zeta_i} \int_{0}^{2\pi} \frac{d\phi_i}{2\pi} \right) \delta(\ln \zeta_1) \times \exp\left(-R' \ln \lim_{\bar{v} \to 0} \frac{V(\{\tilde{p}\}, \kappa_1(\zeta_1 \bar{v}), \dots, \kappa_{m+1}(\zeta_{m+1} \bar{v}))}{\bar{v}} \right).$$

- Result explicitly depends on single-logarithmic quantity $r'_{\ell} \sim \alpha_{\rm s} \ln 1/v$ (roughly: length of boundaries of vetoed region, for each leg).
- Evaluated by Monte Carlo integration.
- Potentially divergent for non-rIRC safe observables.

Given info from previous pages, final answer is analytical:

$$\ln \Sigma(v) = -\sum_{\ell=1}^{n} C_{\ell} \left[r_{\ell}(v) + r'_{\ell}(v) \left(\ln \bar{d}_{\ell} - b_{\ell} \ln \frac{2E_{\ell}}{Q} \right) + B_{\ell} T \left(\frac{\ln 1/v}{a + b_{\ell}} \right) \right] + \sum_{\ell=1}^{n_{i}} \ln \frac{f_{\ell}(x_{\ell}, v^{\frac{2}{a + b_{\ell}}} \mu_{f}^{2})}{f_{\ell}(x_{\ell}, \mu_{f}^{2})} + \ln S \left(T \left(\frac{\ln 1/v}{a} \right) \right) + \ln \mathcal{F}(C_{1}r'_{1}, \dots, C_{n}r'_{n}),$$

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$$+ B_{\ell} T \left(\frac{\ln 1/v}{a + b_{\ell}} \right) + \sum_{\ell=1}^{n_{i}} \ln \frac{f_{\ell}(x_{\ell}, v^{\frac{2}{a + b_{\ell}}} \mu_{f}^{2})}{f_{\ell}(x_{\ell}, \mu_{f}^{2})}$$

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$$C_\ell = {
m colour\ factor\ } (C_F\ {
m or\ } C_A), \qquad f_\ell(x_\ell,\mu_f^2) = {
m parton\ distributions}$$

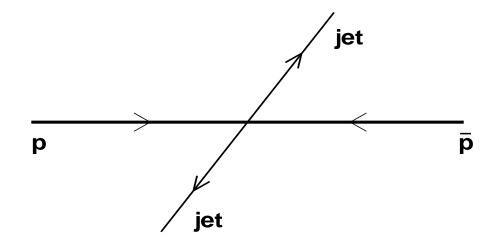
$$r_{\ell}(L) = \int_{v^{\frac{2}{a+b_{\ell}}}Q^{2}}^{v^{\frac{2}{a+b_{\ell}}}Q^{2}} \frac{dk_{t}^{2}}{k_{t}^{2}} \frac{\alpha_{s}(k_{t})}{\pi} \ln\left(\frac{k_{t}}{v^{1/a}Q}\right)^{a/b_{\ell}} + \int_{v^{\frac{2}{a+b_{\ell}}}Q^{2}}^{Q^{2}} \frac{dk_{t}^{2}}{k_{t}^{2}} \frac{\alpha_{s}(k_{t})}{\pi} \ln\frac{Q}{k_{t}},$$

 $S(T(\frac{1}{a}\ln 1/v)) =$ large-angle logarithms (process dependence)

. . .

Example: global thrust in hadronic dijet production

In hadronic-dijet scattering, invent observables to measure *deviation* from Born event.

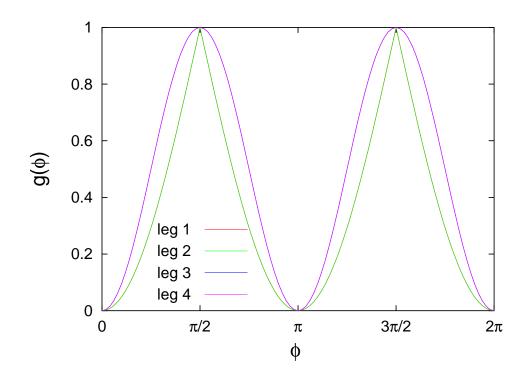


E.g.: global transverse thrust

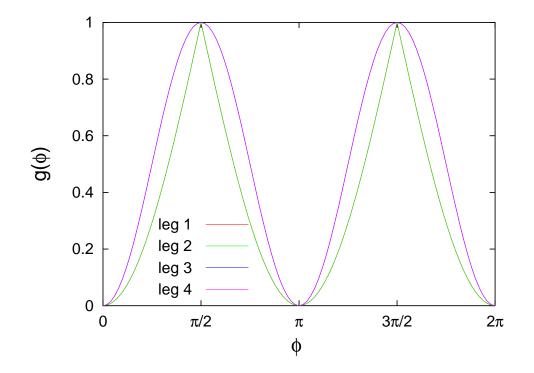
$$1 - T_{t,g} \equiv 1 - \max_{\vec{n}_t} \frac{\sum_i |\vec{p}_{ti} \cdot \vec{n}_t|}{\sum_i p_{ti}},$$

$\log \ell$	a_{ℓ}	b_ℓ	$g_\ell(\phi)$	d_ℓ	$\langle \ln g_{\ell}(\phi) \rangle$
1	1.000	0.000	tabulated	1.02062	-1.859
2	1.000	0.000	tabulated	1.02062	-1.859
3	1.000	1.000	$\sin^2 \phi$	1.042	$-2\ln(2)$
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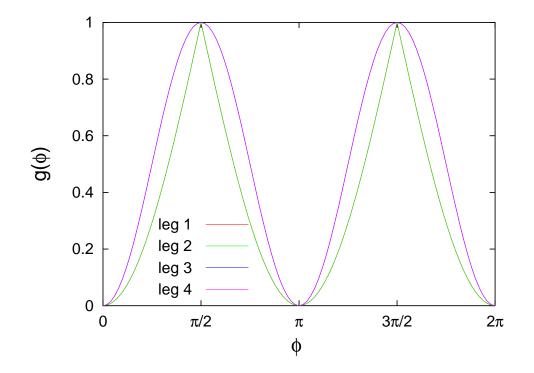


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Test	result
continuously global	Т
rec. IRC safe (cond. 1)	Т
rec. IRC safe (cond. 2a)	Т
rec. IRC safe (cond. 2b)	Т
additivity	Т

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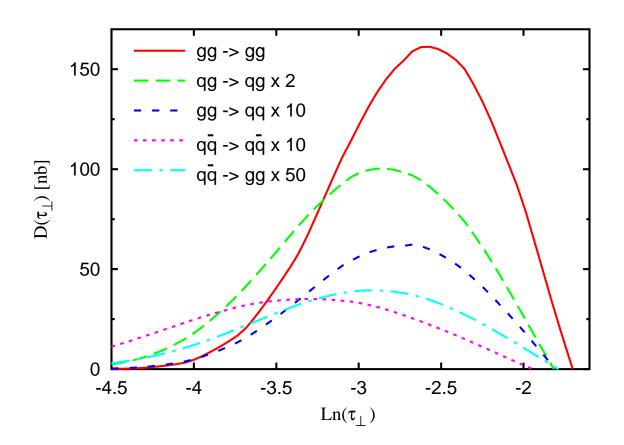


Test	result
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$\mathcal{F} = e^{-\gamma_E R'}$	
$\Gamma = \frac{\Gamma}{\Gamma(1+R')}$	

Resummed thrust for Tevatron

ullet run II regime $\sqrt{s}=1.96\,{\rm TeV}$

- ullet cut on jet-rapidity $|\eta| < 1$
- cut on jet transverse energy $E_T > 50 \text{GeV}$

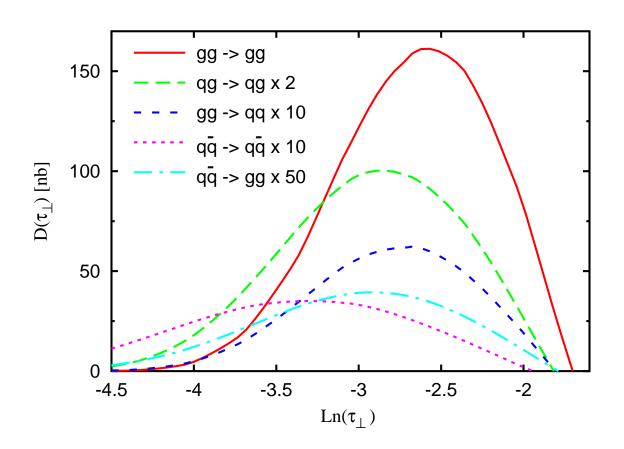


PRELIMINARY!

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Conclusions/Outlook

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Longer-term Outlook

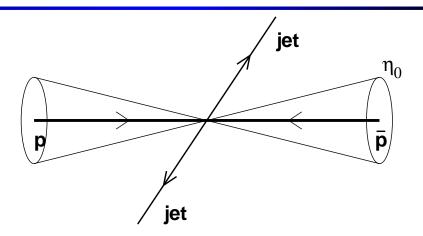
- Matching with fixed order (in progress) → Phenomenology
- Extending scope (e.g. non-global observables?)

EXTRA SLIDES

Observables in hadronic dijet production

Cut around the beam $|\eta| < \eta_0$

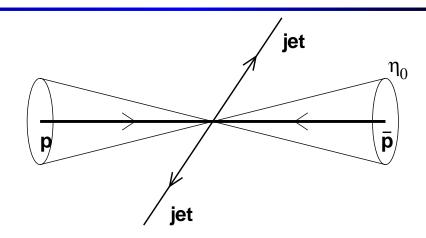
→ Problems with globalness



Observables in hadronic dijet production

Cut around the beam $|\eta| < \eta_0$

→ Problems with globalness



Directly global observables: $\eta_0 \gg 1$

Transverse thrust

$$T_T = \frac{1}{E_T} \max_{\vec{n}_T} \sum_{i} |\vec{p}_{ti} \cdot \vec{n}_T|$$

Thrust minor

$$T_m = \frac{1}{E_T} \sum_{i} |p_i^{out}|$$

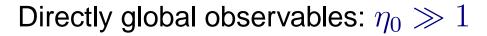
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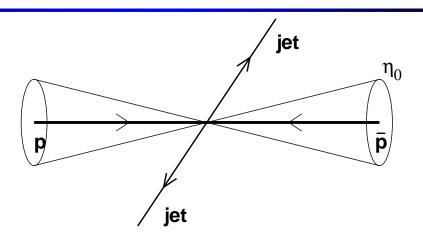
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Predictions valid as long as $\log 1/v < (a+b_{\ell})|\eta_0|$



Indirectly global observables: $\eta_0 = \mathcal{O}(1)$

Transverse thrust

$$T_T = \frac{1}{E_{T,\eta_0}} \left(\max_{\vec{n}_T} \sum_{|\eta_i| < \eta_0} |\vec{p}_{ti} \cdot \vec{n}_T| - \left| \sum_{|\eta_i| < \eta_0} \vec{p}_{ti} \right| \right)$$

Thrust minor

$$T_{m} = \frac{1}{E_{T,\eta_{0}}} \left(\sum_{|\eta_{i}| < \eta_{0}} |p_{i}^{out}| + |\sum_{|\eta_{i}| < \eta_{0}} \vec{p}_{ti}| \right)$$

Predictions valid as usual but \mathcal{F} diverges for $R'=R'_c$