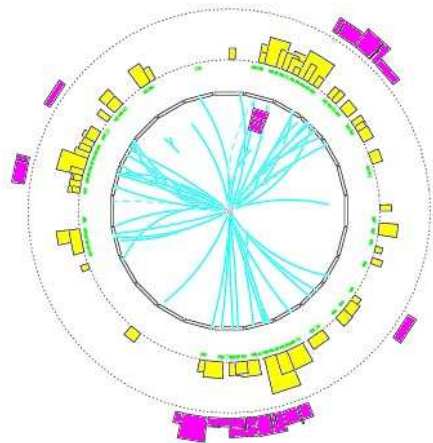
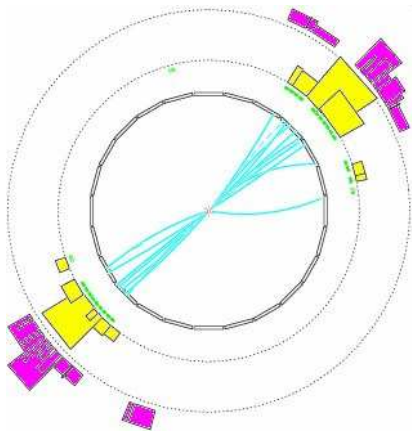


Jets

Gavin Salam

LPTHE, CNRS and UPMC (Univ. Paris 6)

CTEQ MCNet school
Debrecen, Hungary, August 2008



Jets are everywhere in QCD
Our *window on partons*

But *not* the same as partons:
Partons ill-defined; jets *well-definable*

Perturbatively

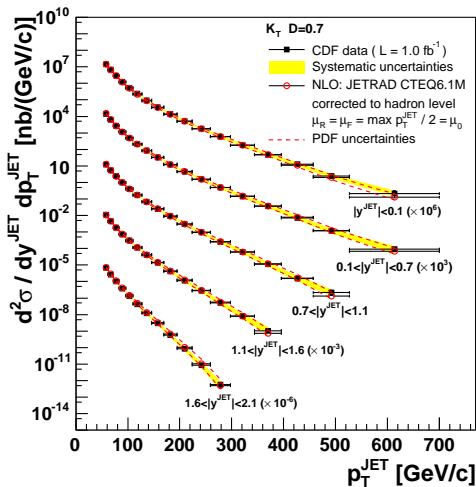
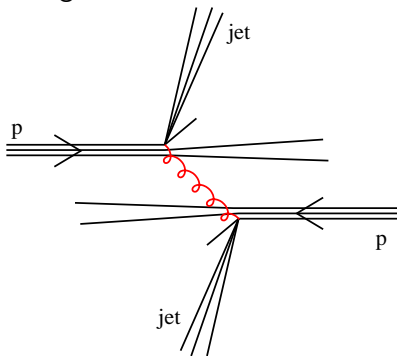
- ▶ Quarks fragment: soft & collinear divergences for gluon emission
- ▶ Gluons fragment: soft & collinear divergences for gluon emission
collinear divergences for quark emission
- ▶ Even perturbative coupling is not so small

Non-perturbatively

- ▶ precise process long way from being understood, even by lattice
- ▶ good models contain many parameters — complex process

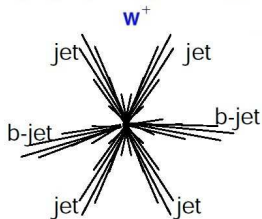
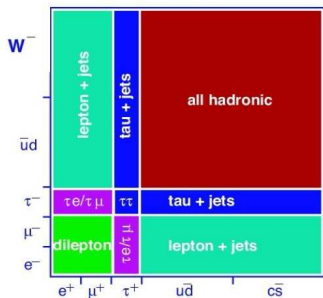
High-energy partons unavoidably lead to collimated bunches of hadrons.

Jets are unavoidable at hadron colliders, e.g. from parton scattering



Jet cross section: data and theory agree over many orders of magnitude \Leftrightarrow probe of underlying interaction

$t\bar{t}$ decay modes



All-hadronic

(BR~46%, huge bckg)

picture: Juste LP05

Heavy objects: multi-jet final-states

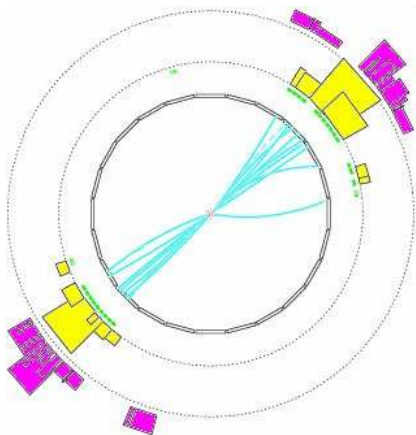
- ▶ 10^7 $t\bar{t}$ pairs for 10 fb^{-1}
- ▶ Vast # of QCD multijet events

# jets	# events for 10 fb^{-1}
3	$9 \cdot 10^8$
4	$7 \cdot 10^7$
5	$6 \cdot 10^6$
6	$3 \cdot 10^5$
7	$2 \cdot 10^4$
8	$2 \cdot 10^3$

Tree level

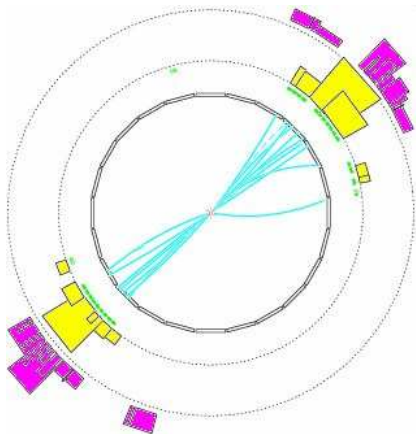
$p_t(\text{jet}) > 60 \text{ GeV}$, $\theta_{ij} > 30 \text{ deg}$, $|y_{ij}| < 3$

Draggiotis, Kleiss & Papadopoulos '02

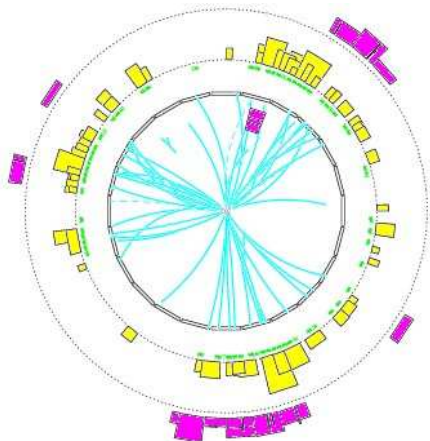


Jets are what we see.
Clearly(?) 2 jets here

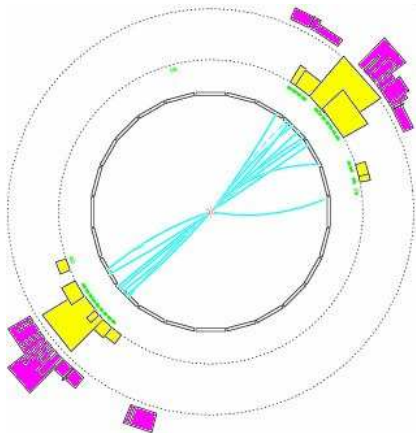
How many jets do you see?
Do you really want to ask yourself
this question for 10^8 events?



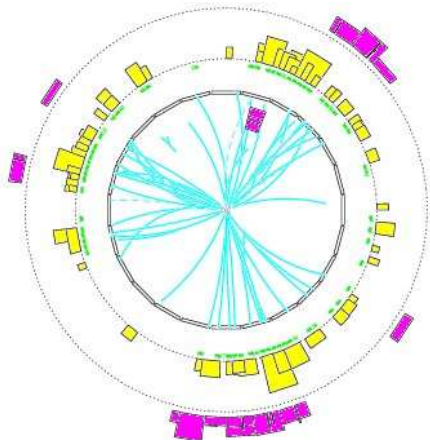
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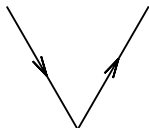
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How many jets do you see?
Do you really want to ask yourself
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- ▶ *A jet definition is a fully specified set of rules for projecting information from 100's of hadrons, onto a handful of parton-like objects:*
 - ▶ or project 1000's of calorimeter towers
 - ▶ or project dozens of (showered) partons
 - ▶ or project a handful of (unshowered) partons
- ▶ Resulting objects (jets) used for many things, e.g. :
 - ▶ reconstructing decaying massive particles
 - ▶ constraining proton structure
 - ▶ as a theoretical tool to attribute structure to an event

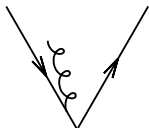
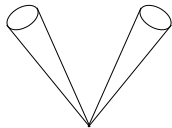
e.g. top → 3 jets
- ▶ You *lose much information* in projecting event onto jet-like structure:
 - ▶ Sometimes information you had no idea how to use
 - ▶ Sometimes information you may not trust, or of no relevance



LO partons

Jet ↓ Defⁿ

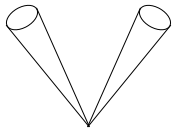
jet 1 jet 2



NLO partons

Jet ↓ Defⁿ

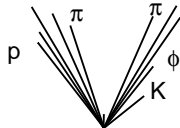
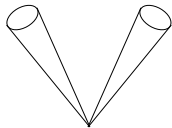
jet 1 jet 2



parton shower

Jet ↓ Defⁿ

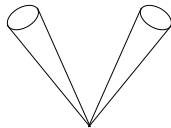
jet 1 jet 2



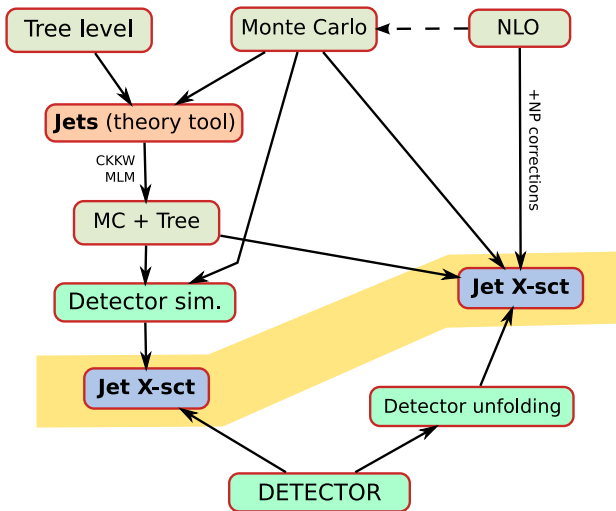
hadron level

Jet ↓ Defⁿ

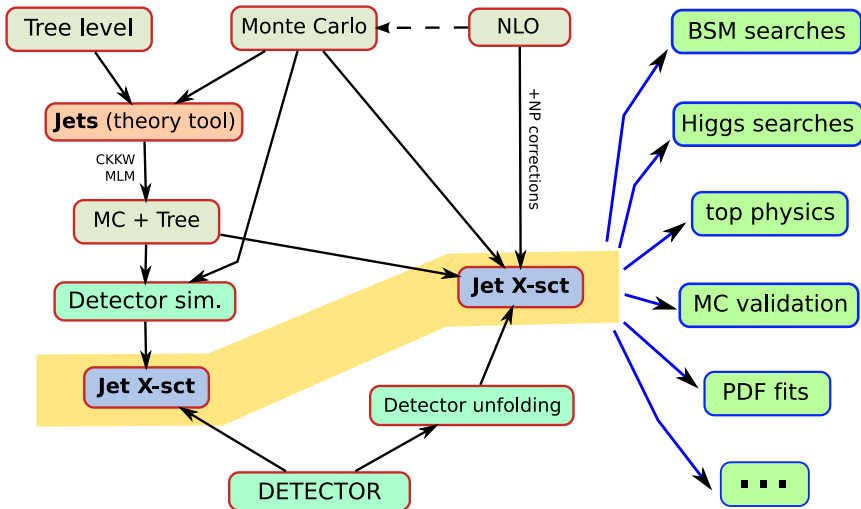
jet 1 jet 2



Projection to jets should be resilient to QCD effects



Jet (definitions) provide central link between expt., “theory” and theory
And jets are an input to almost all analyses



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And jets are an input to almost all analyses

Aims: to provide you with

- ▶ the “basics” needed to understand what goes into current jet-based measurements;
- ▶ some insight into the issues that are relevant when thinking about a jet measurement

Structure:

- ▶ General considerations
- ▶ Common jet definitions — we'll look at 2 broad classes
 - ▶ Sequential recombination
 - ▶ Cone
- ▶ The physics of jets [briefly]

today
today & tomorrow
tomorrow

There is no unique jet definition

The construction of a jet is unavoidably ambiguous. On at least two fronts:

1. which particles get put together into a common jet? Jet algorithm
+ parameters
2. how do you combine their momenta? Recombination scheme
Most commonly used: direct 4-vector sums (E -scheme)

Taken together, these different elements specify a choice of jet definition

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Taken together, these different elements specify a choice of jet definition

- ▶ Physical results (particle discovery, masses, PDFs, coupling) should be independent of your choice of jet definition
 - a bit like renormalisation scale/scheme invariance
 - Tests independence on modelling of radiation, hadronisation, etc.
- ▶ Except when there is a good reason for this not to be the case

Jetography, like photography



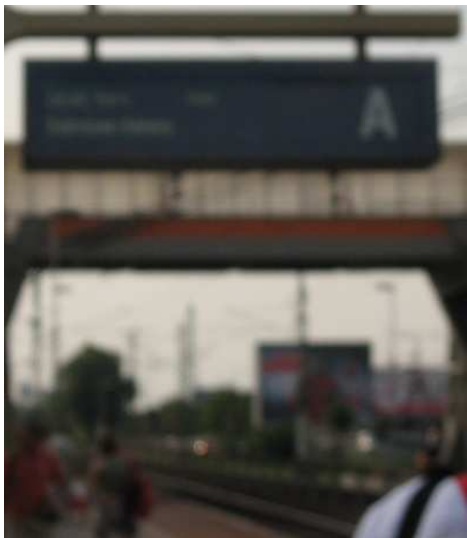
- ▶ Fine detail on bus ticket to train station — shoot from close up, focus = 40cm

[get to train station]

- ▶ Keep focus at 40cm
- ▶ Reset focus to 6m

Catch correct train

Jetography, like photography



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Catch correct train

Jets should be **invariant** with respect to certain modifications of the event:

- ▶ collinear splitting
- ▶ infrared emission

Why?

- ▶ Because otherwise lose real-virtual cancellation in NLO/NNLO QCD calculations → divergent results
- ▶ Hadron-level 'jets' fundamentally non-perturbative
- ▶ Detectors resolve neither full collinear nor full infrared event structure

Known as **infrared and collinear safety**

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Known as **infrared and collinear safety**

Sequential recombination (k_t , etc.)

- ▶ bottom-up
- ▶ successively undoes QCD branching

Cone

- ▶ top-down
- ▶ centred around idea of an 'invariant', directed energy flow

Sequential recombination jet algorithms

Majority of QCD branching is soft & collinear, with following divergences:

$$[dk_j] |M_{g \rightarrow g_i g_j}^2(k_j)| \simeq \frac{2\alpha_s C_A}{\pi} \frac{dE_j}{\min(E_i, E_j)} \frac{d\theta_{ij}}{\theta_{ij}}, \quad (E_j \ll E_i, \theta_{ij} \ll 1).$$

To invert branching process, take pair with strongest divergence between them — they're the most *likely* to belong together.

This is basis of **k_t /Durham algorithm** (e^+e^-):

1. Calculate (or update) distances between all particles i and j :

$$y_{ij} = \frac{2 \min(E_i^2, E_j^2)(1 - \cos \theta_{ij})}{Q^2}$$

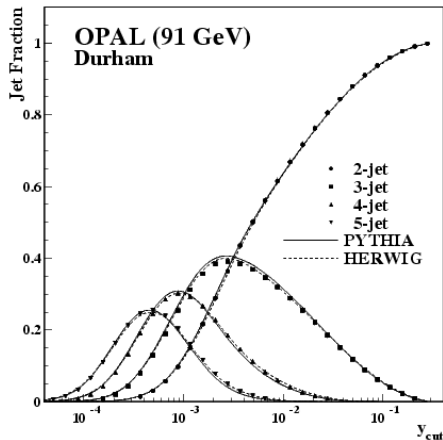
2. Find smallest of y_{ij}

NB: relative k_t between particles

- ▶ If $> y_{cut}$, stop clustering
- ▶ Otherwise recombine i and j , and repeat from step 1

Catani, Dokshitzer, Olsson, Turnock & Webber '91

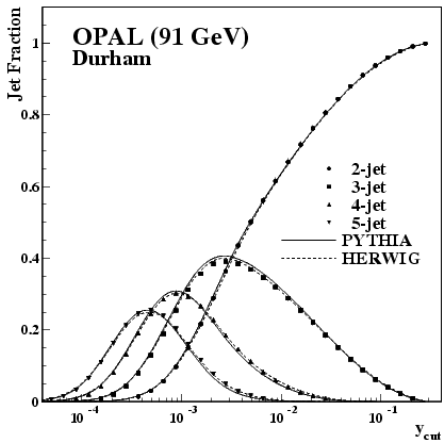
- ▶ Gives hierarchy to event and jets
 - Event can be specified by y_{23} , y_{34} , y_{45} .
- ▶ Resolution parameter related to minimal transverse momentum between jets



Most widely-used jet algorithm in e^+e^-

- ▶ Collinear safe: collinear particles recombined early on
- ▶ Infrared safe: soft particles have no impact on rest of clustering seq.

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Most widely-used jet algorithm in e^+e^-

- ▶ Collinear safe: collinear particles recombined early on
- ▶ Infrared safe: soft particles have no impact on rest of clustering seq.

1st attempt

- ▶ Lose absolute normalisation scale Q . So use unnormalised d_{ij} rather than y_{ij} :

$$d_{ij} = 2 \min(E_i^2, E_j^2)(1 - \cos \theta_{ij})$$

- ▶ Now also have *beam remnants* (go down beam-pipe, not measured)
Account for this with particle-beam distance

$$d_{iB} = 2E_i^2(1 - \cos \theta_{iB})$$

squared transv. mom. wrt beam

2nd attempt: make it longitudinally boost-invariant

- ▶ Formulate in terms of rapidity (y), azimuth (ϕ), p_t

$$d_{ij} = \min(p_{ti}^2, p_{tj}^2) \Delta R_{ij}^2, \quad \Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$$

NB: not η_i , E_{ti}

- ▶ Beam distance becomes

$$d_{iB} = p_{ti}^2$$

squared transv. mom. wrt beam

Catani, Dokshitzer, Seymour & Webber '93

Apart from measures, just like e^+e^- alg.

Known as **exclusive k_t algorithm**.

Problem: at hadron collider, no single fixed scale (as in Q in e^+e^-). So

how do you choose d_{cut} ?

See e.g. Seymour & Tevlin '06

3rd attempt: **inclusive k_t algorithm**

- ▶ Introduce angular radius R (NB: dimensionless!)

$$d_{ij} = \min(p_{ti}^2, p_{tj}^2) \frac{\Delta R_{ij}^2}{R^2}, \quad d_{iB} = p_{ti}^2$$

- ▶ 1. Find smallest of d_{ij} , d_{iB}
- 2. if ij , recombine them
- 3. if iB , call i a jet and remove from list of particles
- 4. repeat from step 1 until no particles left.

S.D. Ellis & Soper, '93; the simplest to use

Jets all separated by at least R on y, ϕ cylinder.

NB: number of jets not IR safe (soft jets near beam); number of jets above p_t cut **is** IR safe.

Fast Hierarchical Clustering and Other Applications of Dynamic Closest Pairs

David Eppstein
UC Irvine

We develop data structures for dynamic closest pair problems with arbitrary distance functions, that do not necessarily come from any geometric structure on the objects. Based on a technique previously used by the author for Euclidean closest pairs, we show how to insert and delete objects from an n -object set, maintaining the closest pair, in $O(n \log^2 n)$ time per update and $O(n)$ space. With quadratic space, we can instead use a quadtree-like structure to achieve an optimal time bound, $O(n)$ per update. We apply these data structures to hierarchical clustering, greedy matching, and TSP heuristics, and discuss other potential applications in machine learning, Gröbner bases, and local improvement algorithms for partition and placement problems. Experiments show our new methods to be faster in practice than previously used heuristics.

Categories and Subject Descriptors: F.2.2 [Analysis of Algorithms]: Nonnumeric Algorithms

General Terms: Closest Pair, Agglomerative Clustering

Additional Key Words and Phrases: TSP, matching, conga line data structure, quadtree, nearest neighbor heuristic

1. INTRODUCTION

Hierarchical clustering has long been a mainstay of statistical analysis, and clustering based methods have attracted attention in other fields: computational biology (reconstruction of evolutionary trees; tree-based multiple sequence alignment), scientific simulation (n -body problems), theoretical computer science (network design and nearest neighbor searching) and of course the web (hierarchical indices such as Yahoo). Many clustering methods have been devised and used in these applications, but less effort has gone into algorithmic speedups of these methods.

In this paper we identify and demonstrate speedups for a key subroutine used in several clustering algorithms, that of maintaining closest pairs in a dynamic set of objects. We also describe several other applications or potential applications of the

Idea behind k_t alg. is to be found over and over in many areas of (computer) science.

Sequential recombination

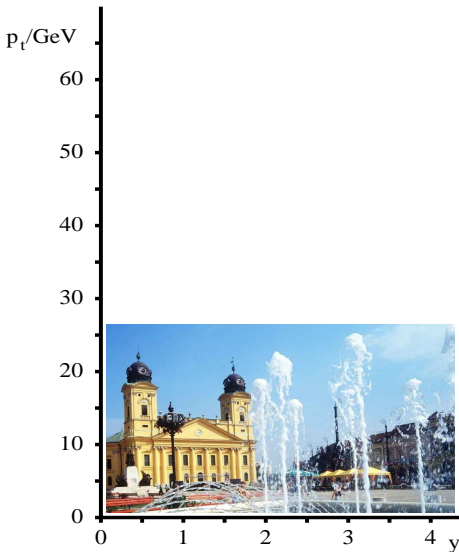
k_t alg.: Find smallest of

$$d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2, \quad d_{iB} = k_{ti}^2$$

If d_{ij} recombine; if d_{iB} , i is a jet
Example clustering with k_t algorithm, $R = 0.7$

ϕ assumed 0 for all towers



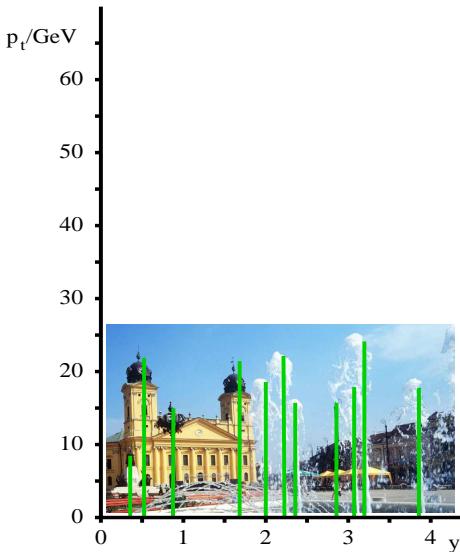


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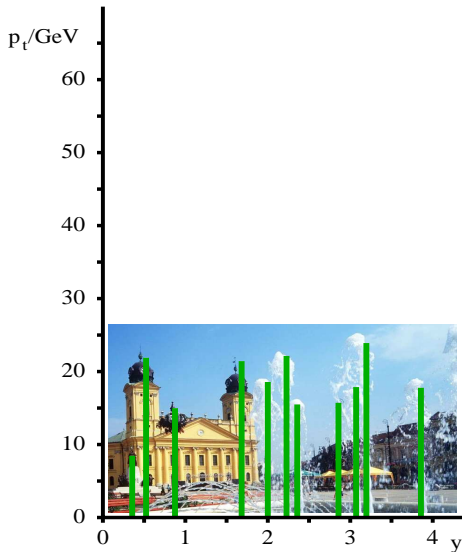


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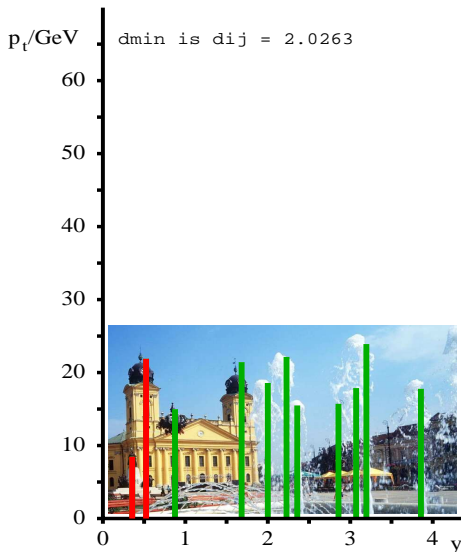


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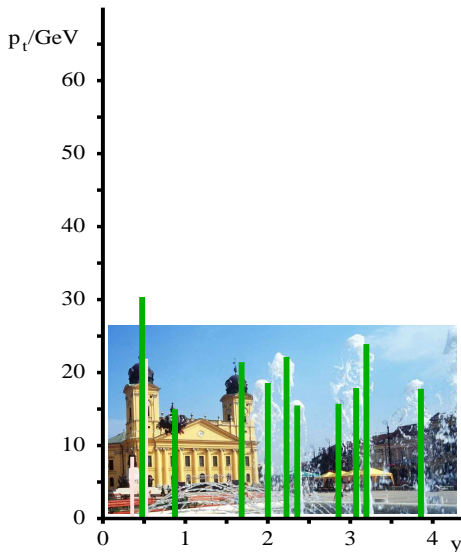


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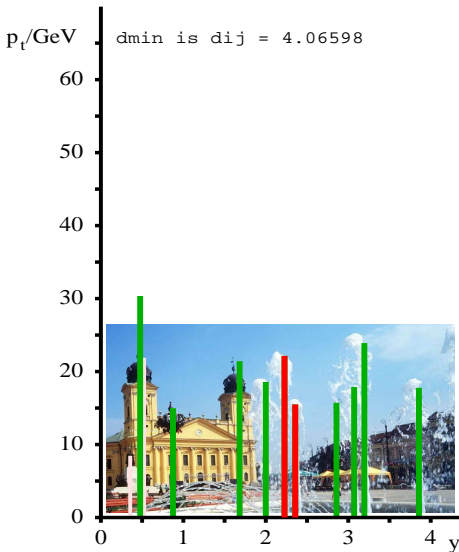


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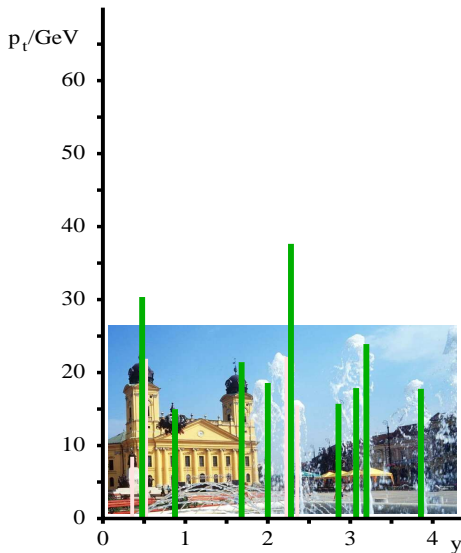


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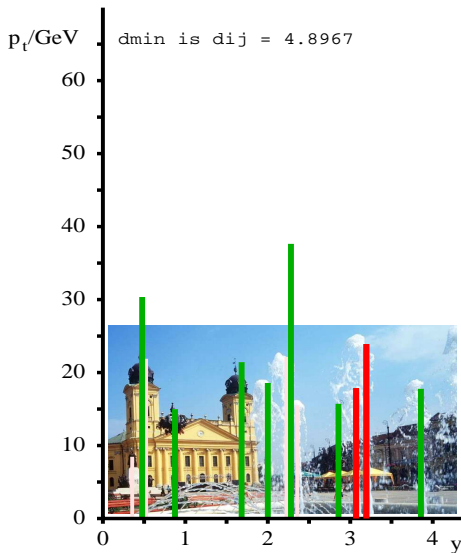


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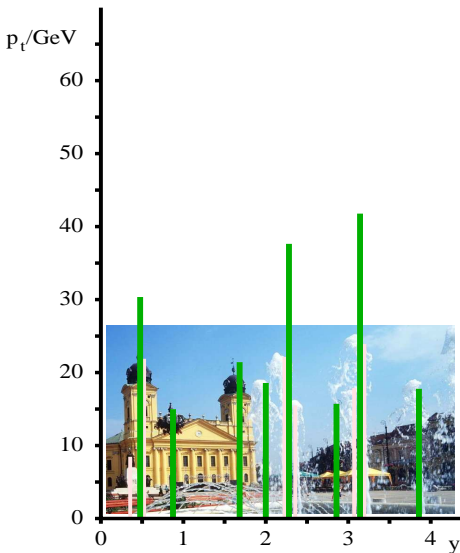


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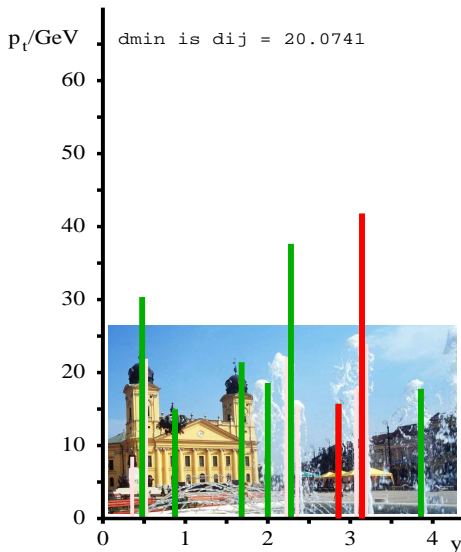


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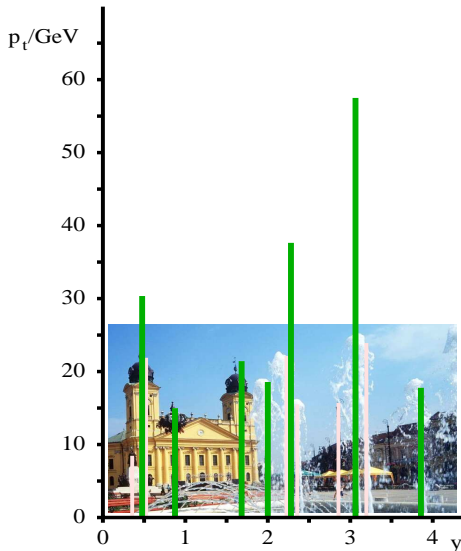


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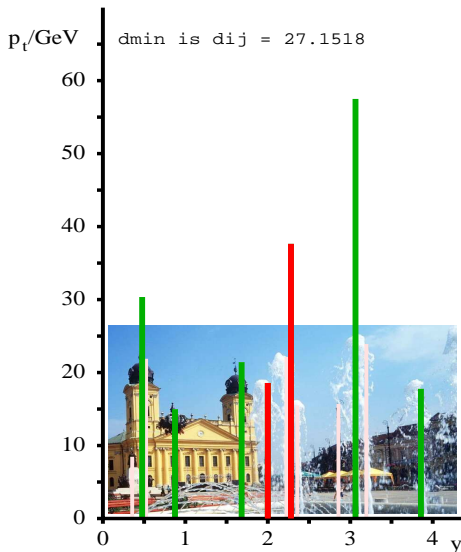


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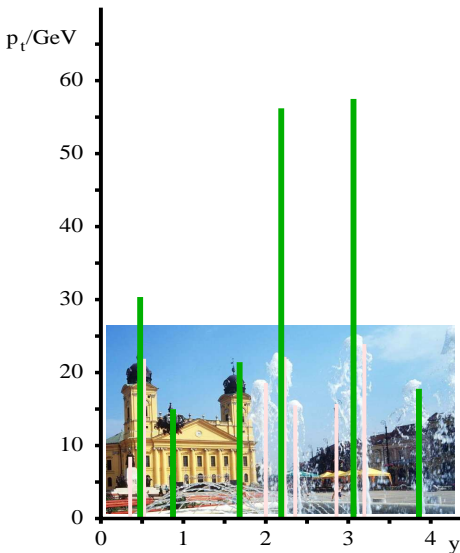


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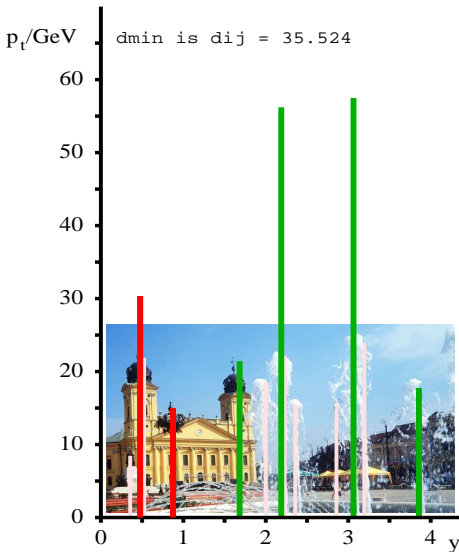


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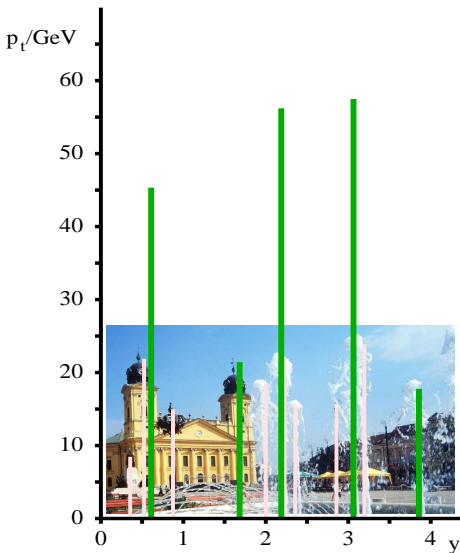


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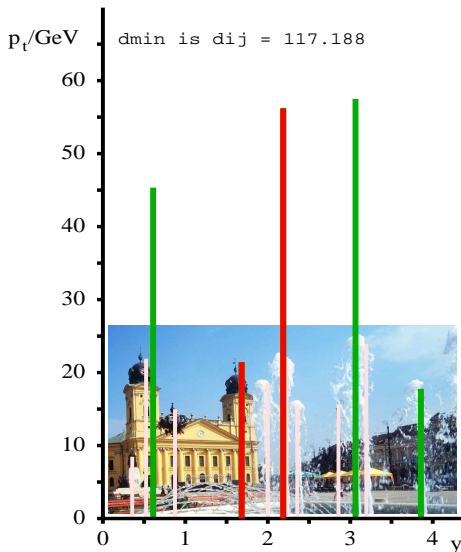


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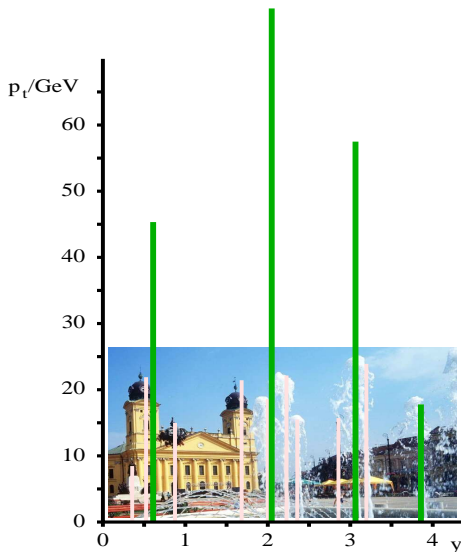


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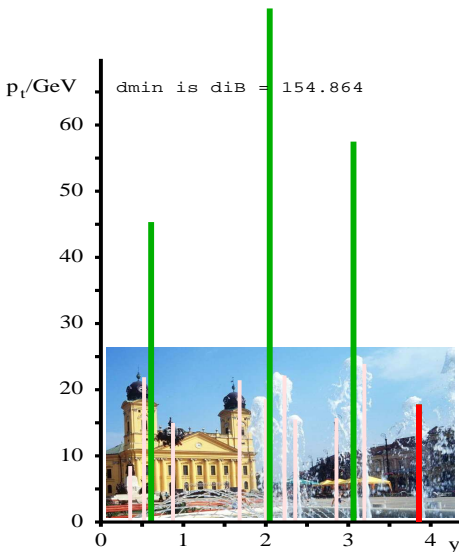


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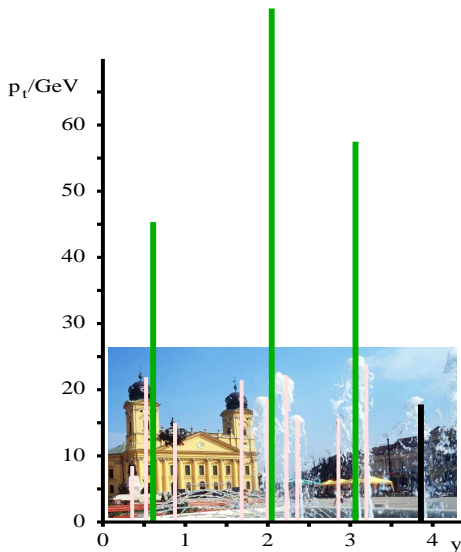


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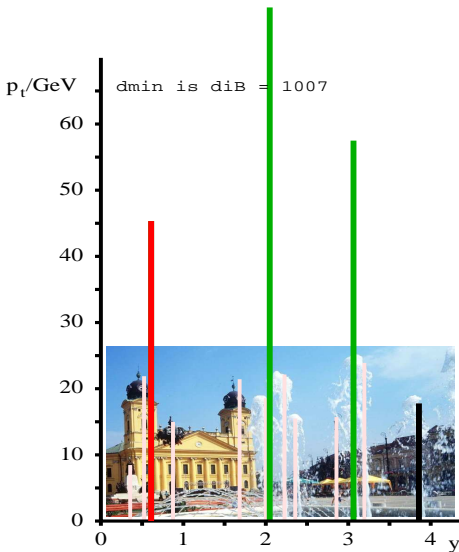


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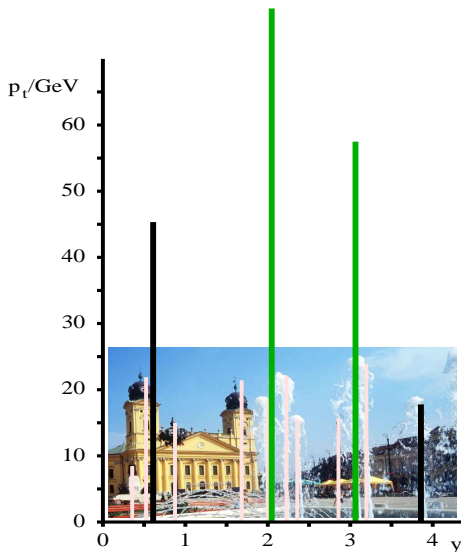


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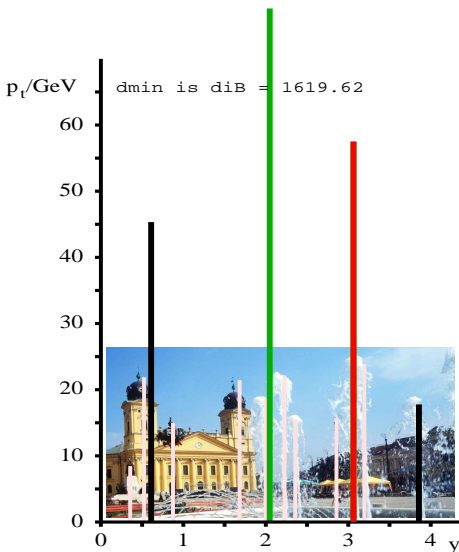


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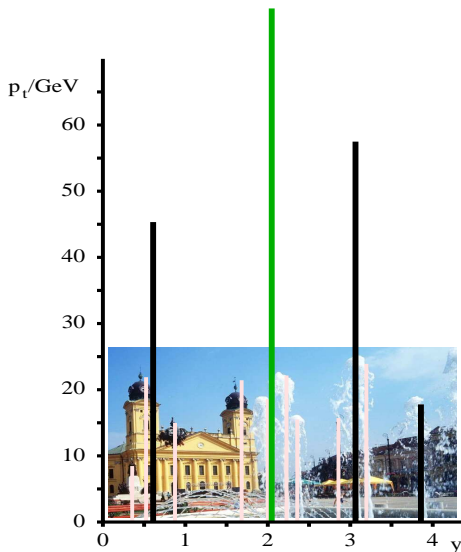


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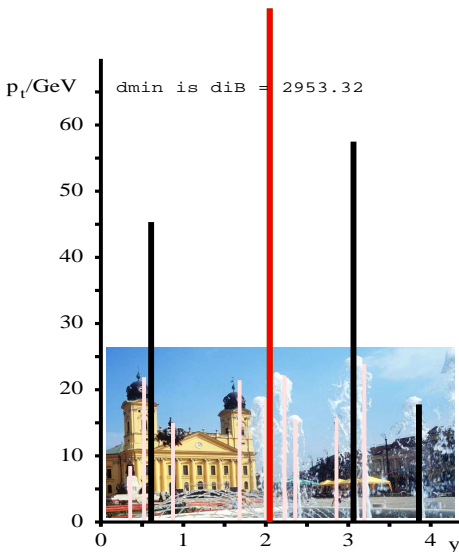


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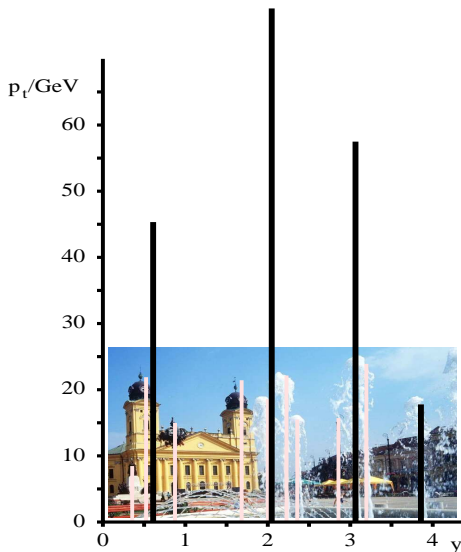


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Cambridge/Aachen: *the simplest of hadron-collider algorithms*

- ▶ Recombine pair of objects closest in ΔR_{ij}
- ▶ Repeat until all $\Delta R_{ij} > R$ — remaining objects are jets

Dokshitzer, Leder, Moretti, Webber '97 (Cambridge): more involved e^+e^- form

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Cacciari, GPS & Soyez, '08 [+ Delsart unpublished]
privileges clustering with *hard* particles first

Privileging different divergences \Leftrightarrow different jets; more later...

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Plenty more variants too, mostly in e^+e^- , e.g.

- ▶ JADE: $d_{ij} = m_{ij}^2/Q^2$
- ▶ Geneva $d_{ij} = 8E_i E_j (1 - \cos \theta_{ij})/9(E_i + E_j)^2$
- ▶ ARCLUS: perform $3 \rightarrow 2$ recombination

the original seq. rec. alg.

In pp , also have modifications of angular measure

- ▶ QCD-metric angular distance: $\Delta R_{ij}^2 \rightarrow 2(\cosh \Delta y_{ij} - \cos \Delta \phi_{ij})$

And beyond just momentum

- ▶ Flavour- k_t algorithm (e^+e^- and pp)

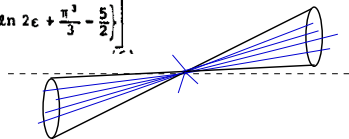
Cone algorithms

First 'jet algorithm' dates back to **Sterman and Weinberg (1977)** — the original infrared-safe cross section:

To study jets, we consider the partial cross section

$\sigma(E, \theta, \Omega, \epsilon, \delta)$ for e^+e^- hadron production events, in which all but a fraction $\epsilon \ll 1$ of the total e^+e^- energy E is emitted within some pair of oppositely directed cones of half-angle $\delta \ll 1$, lying within two fixed cones of solid angle Ω (with $\pi\delta^2 \ll \Omega \ll 1$) at an angle θ to the e^+e^- beam line. We expect this to be measur-

$$\sigma(E, \theta, \Omega, \epsilon, \delta) = (d\sigma/d\Omega)_0 \Omega \left[1 - (g_E^2/3\pi^2) \left\{ 3\ln \delta + 4\ln \delta \ln 2\epsilon + \frac{\pi^2}{3} - \frac{5}{2} \right\} \right]$$



Groundbreaking; good for 2 jets in e^+e^- ; but never widely generalised

Unifying idea: momentum flow within a cone only
marginally modified by QCD branching

But cones come in many variants

Finding cones \ Processing	Progressive Removal	Split-Merge	Split-Drop
Seeded, Fixed (FC)	GetJet CellJet		
Seeded, Iterative (IC)	CMS Cone	JetClu (CDF) [†] ATLAS cone	
Seeded, It. + Midpoints (IC _{mp})		CDF MidPoint D0 Run II cone	PxCone
Seedless (SC)		SISCone	

[†] JetClu also has “ratcheting”

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Common features in discussion of cones

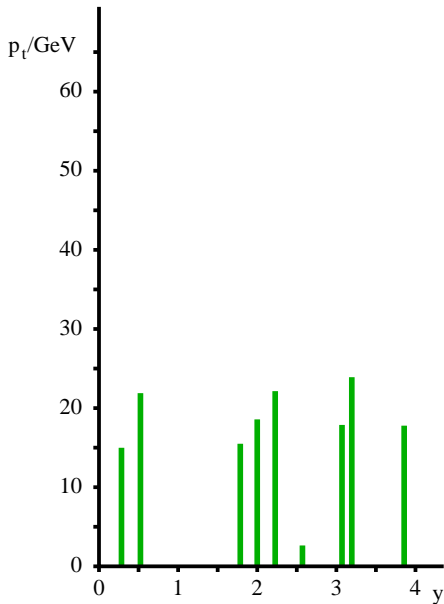
- ▶ Cones are always understood as circles in rapidity (y) and azimuth ϕ .
- ▶ A particle i is within the cone of radius R around the axis a if

$$\Delta R_{ia}^2 = (y_i - y_a)^2 + (\phi_i - \phi_a)^2 < R^2$$

The usual hadron collider variables

- ▶ We'll use $R = 0.7$ in the examples that follow
- ▶ And we'll use events all of whose particles are at $\phi = 0$, for simplicity

Fixed Cone, Prog Removal (FC-PR)



The simplest of the cones

PyCell, CellJet, GetJet

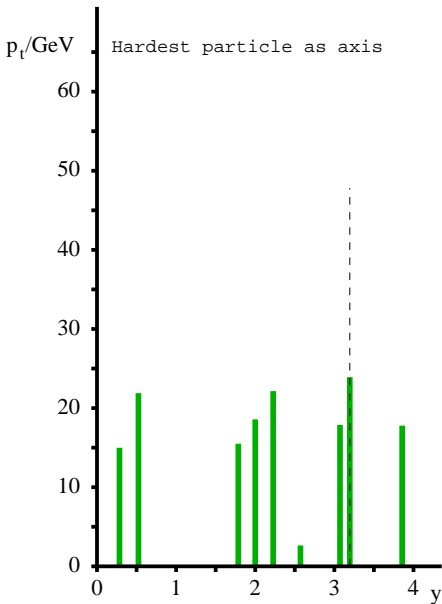
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Notes

- ▶ “Hardest particle” is collinear unsafe → more later...
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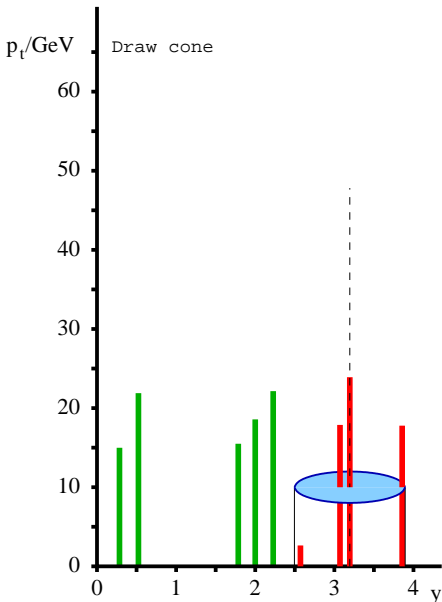
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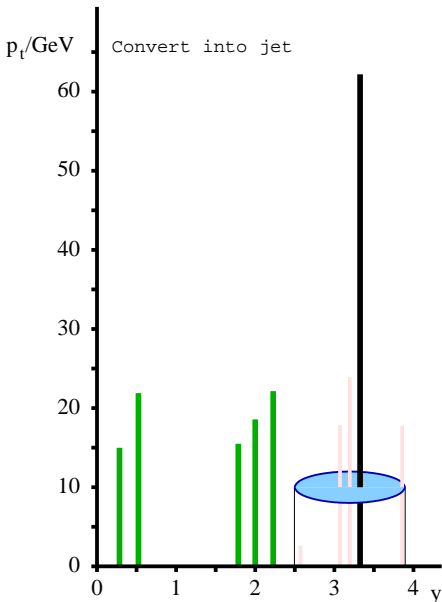
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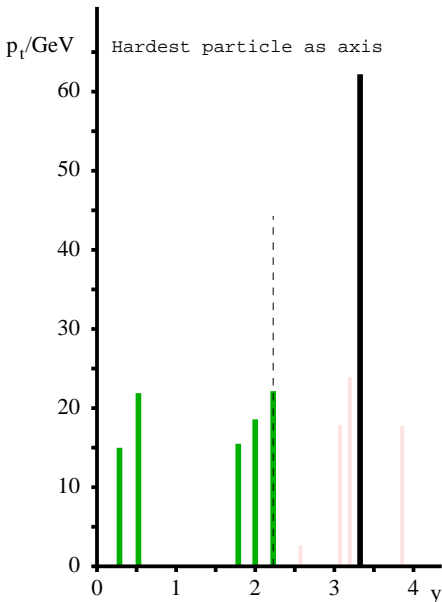
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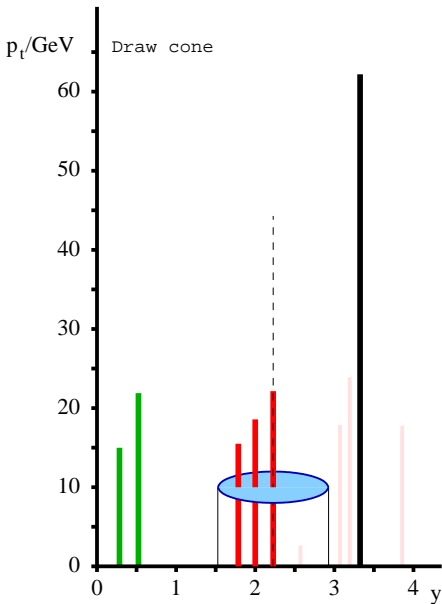
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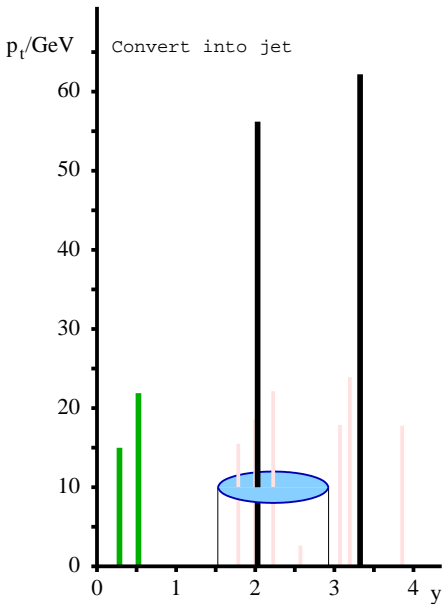
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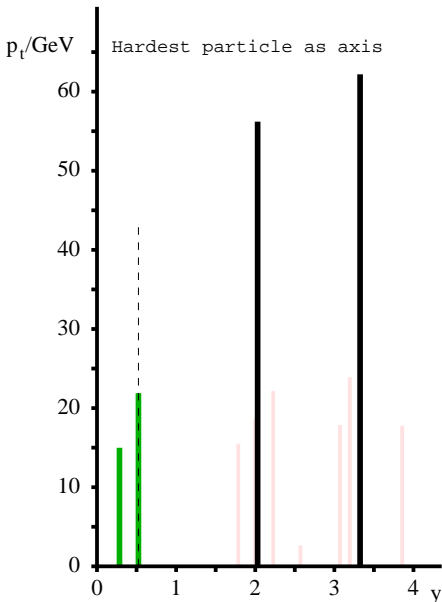
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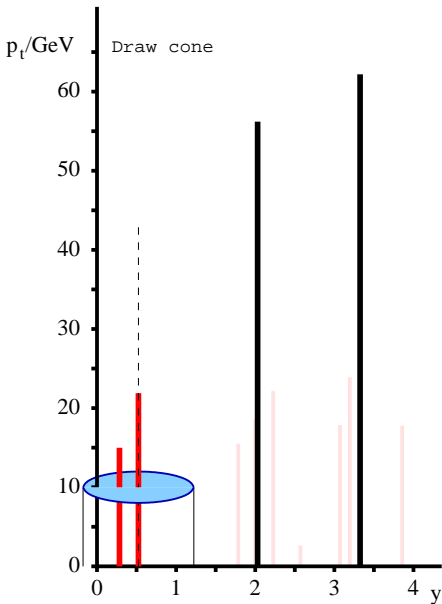
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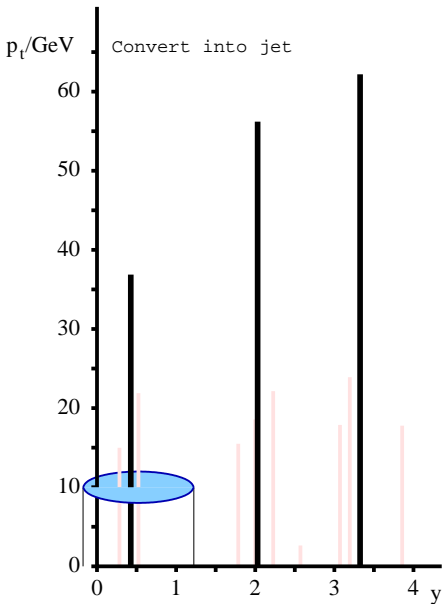
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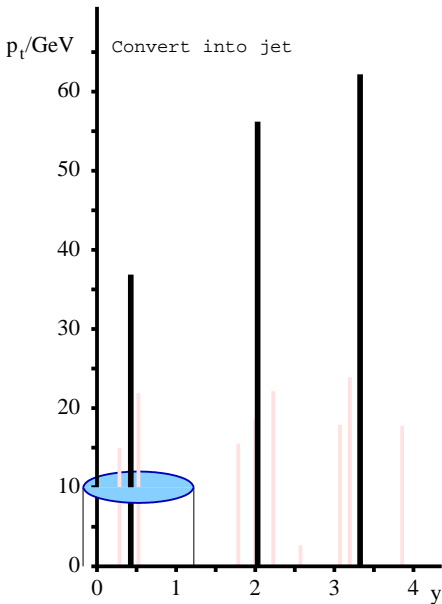
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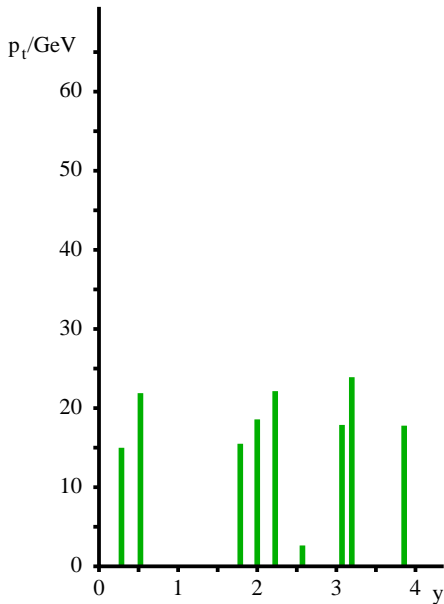
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Iterative Cone, Prog Removal (IC-PR)



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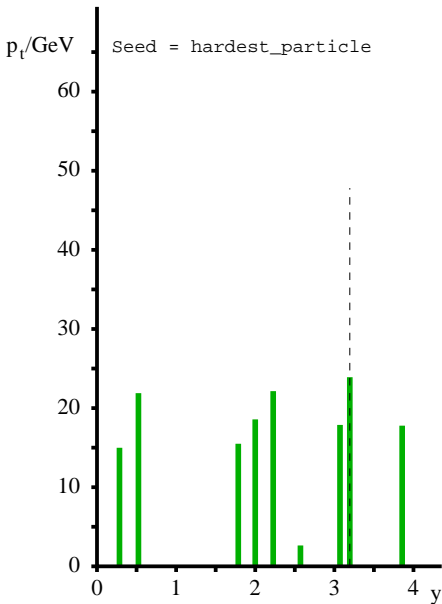
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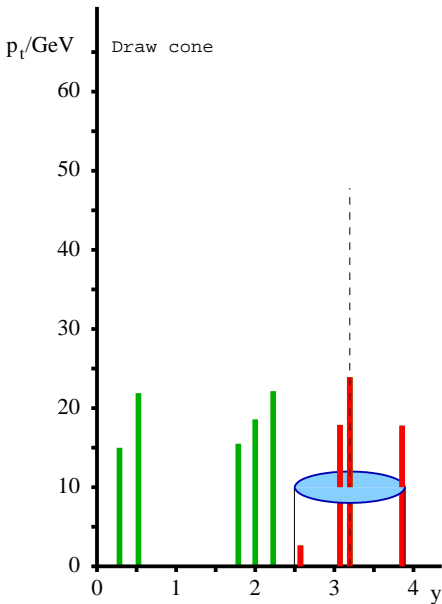
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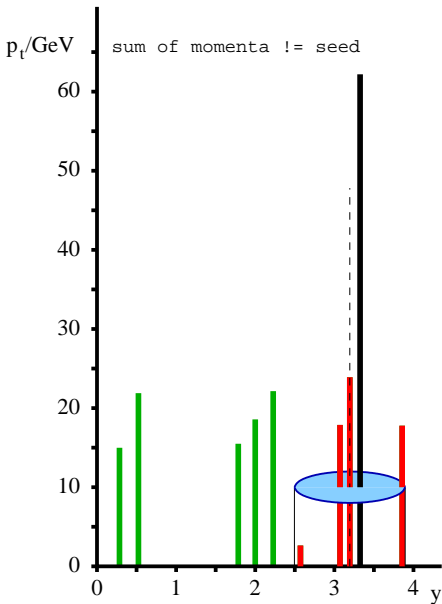
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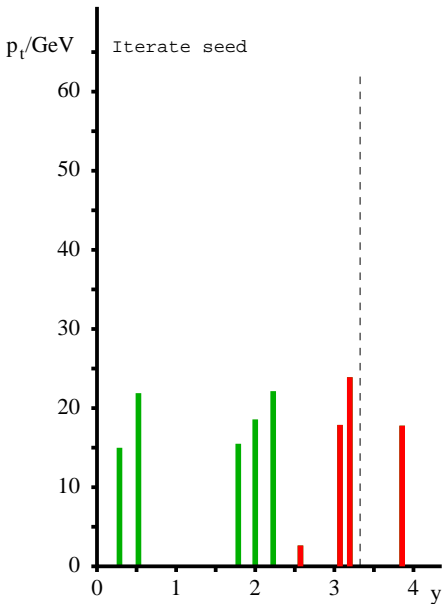
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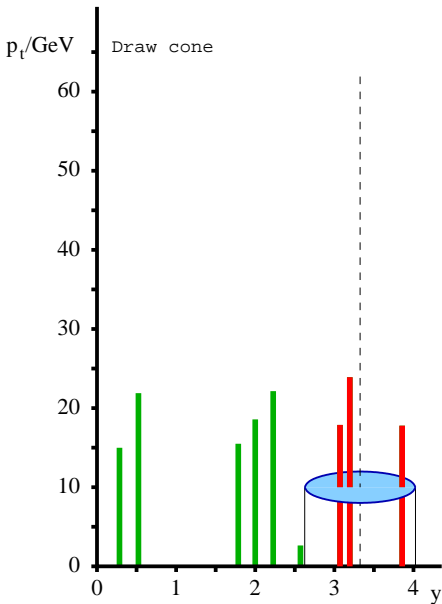
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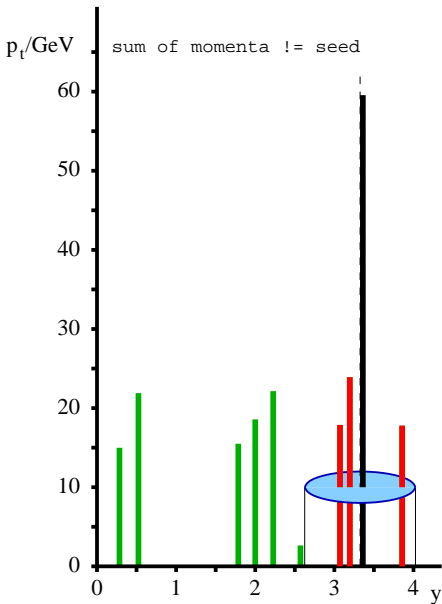
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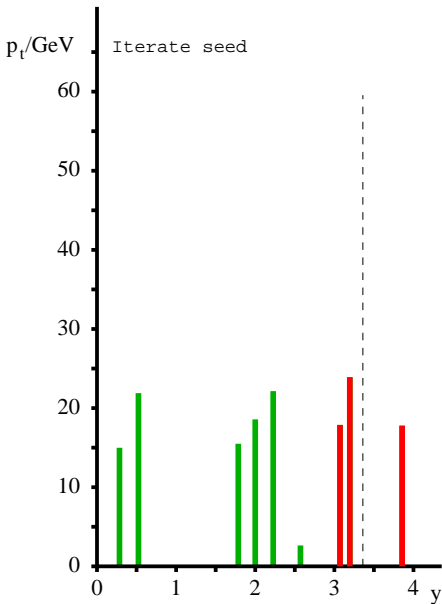
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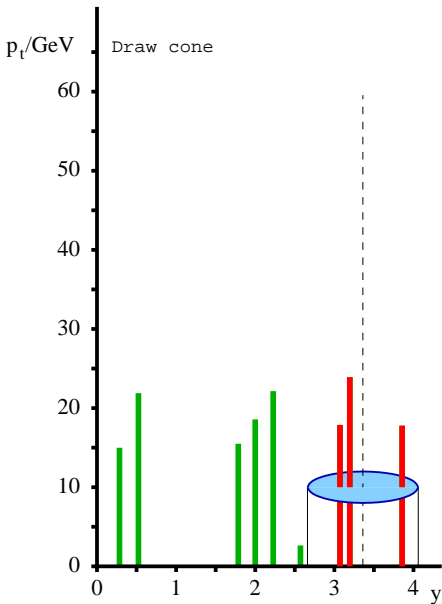
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- ▶ “Hardest particle” is collinear unsafe more right away...

Iterative Cone, Prog Removal (IC-PR)



Next-simplest of the cones

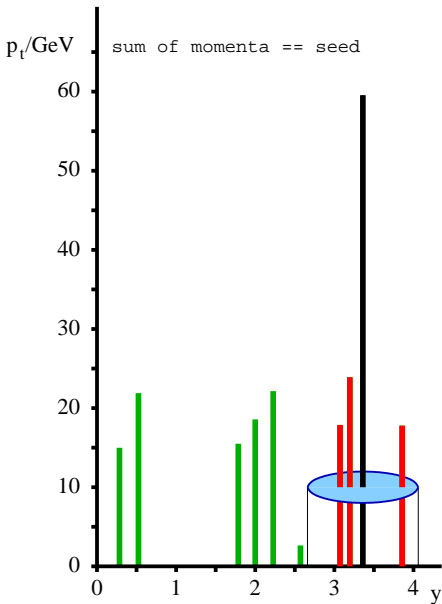
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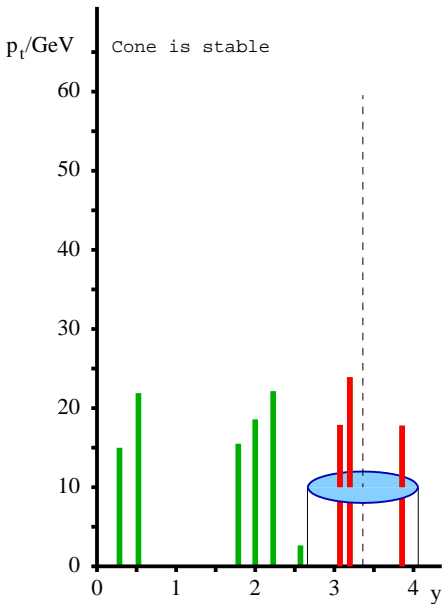
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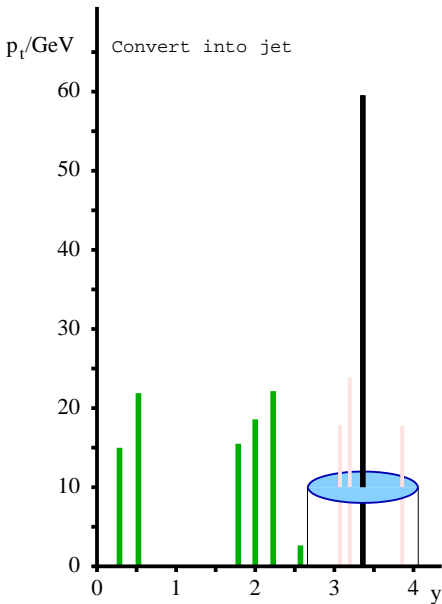
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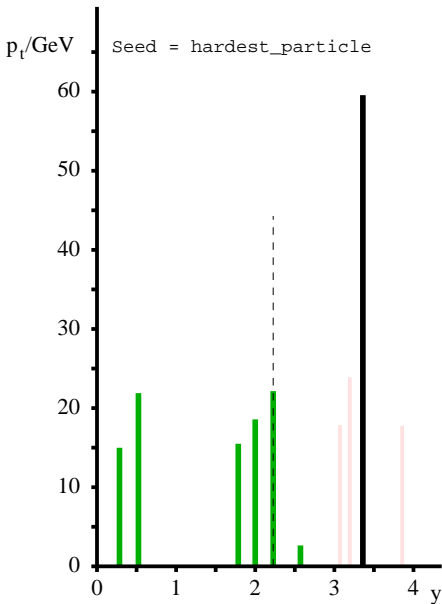
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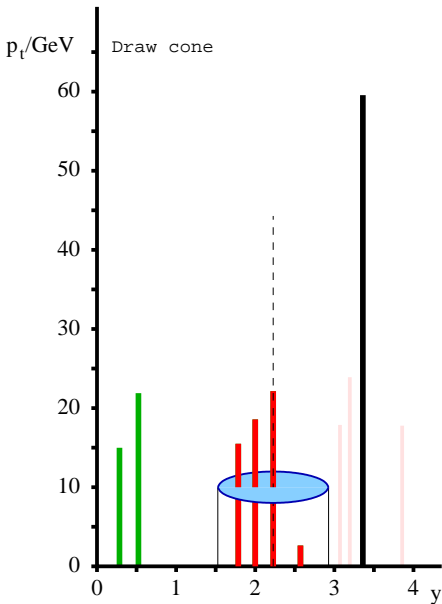
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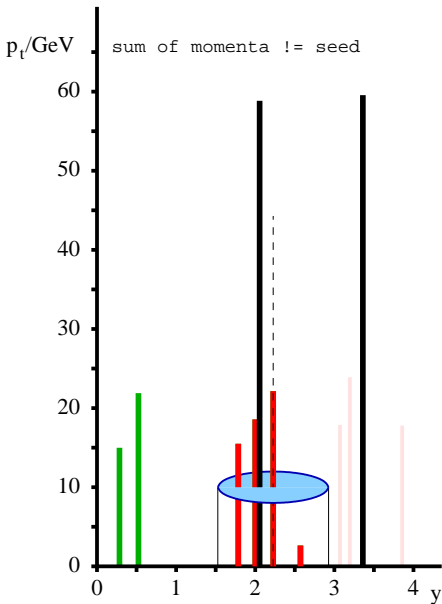
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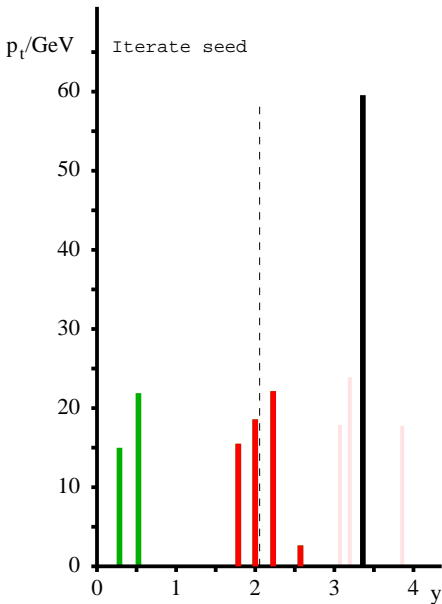
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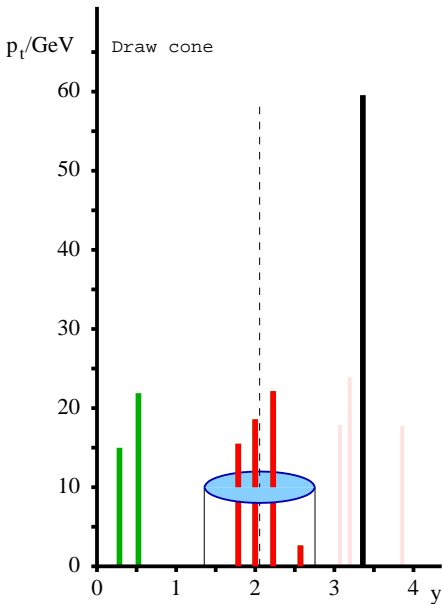
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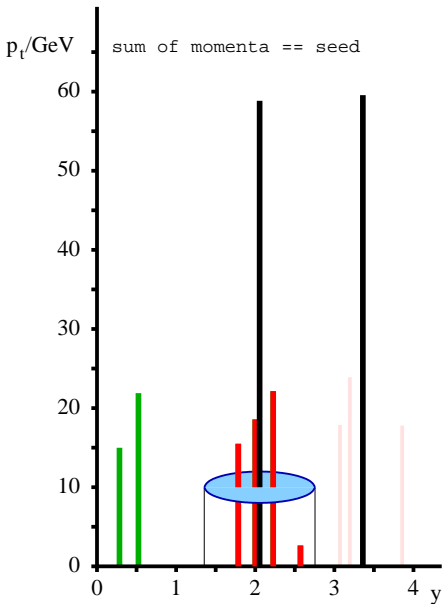
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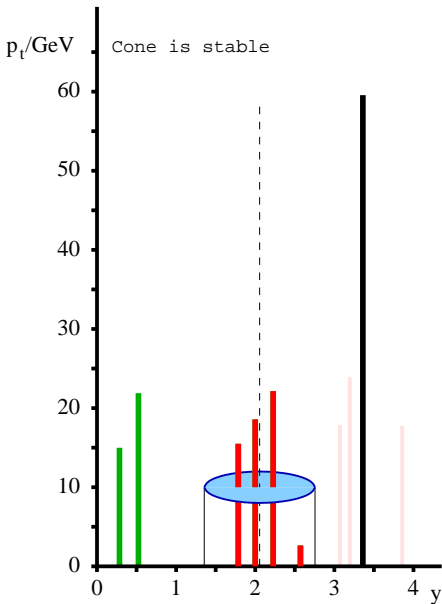
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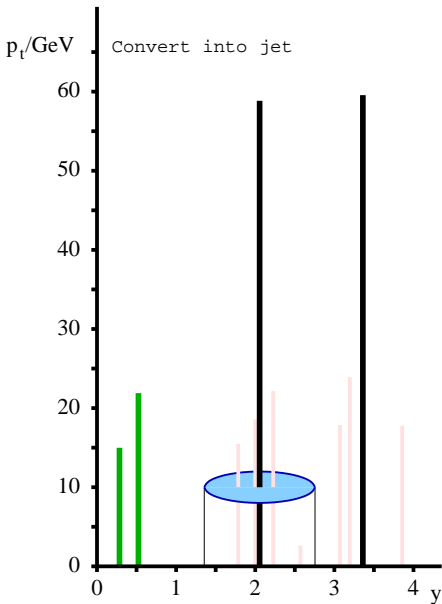
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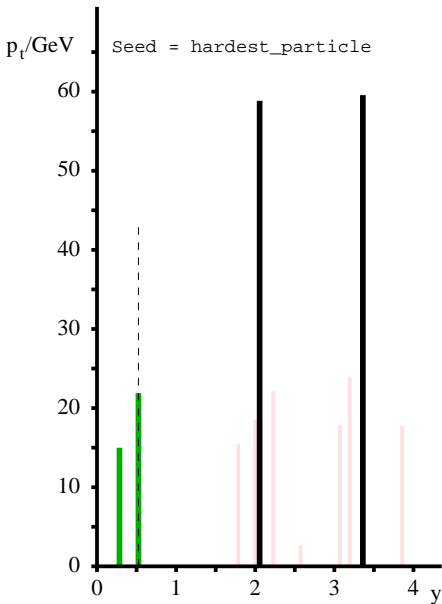
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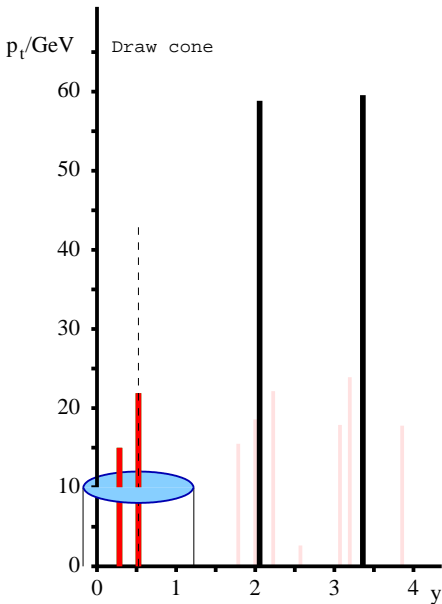
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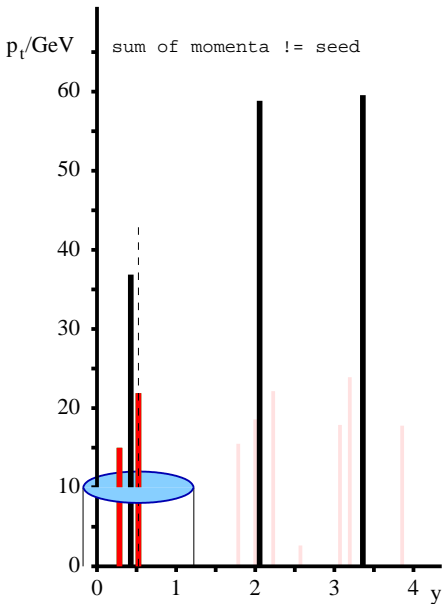
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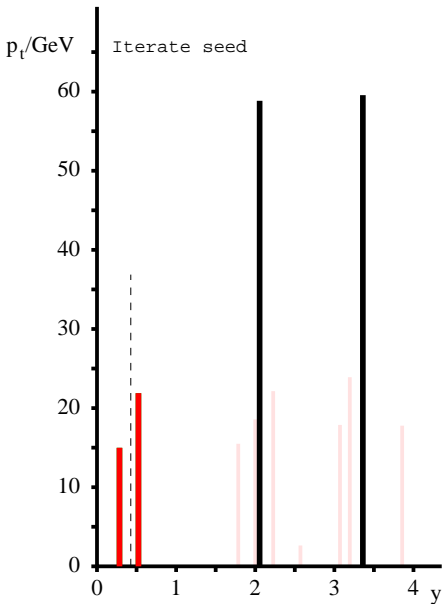
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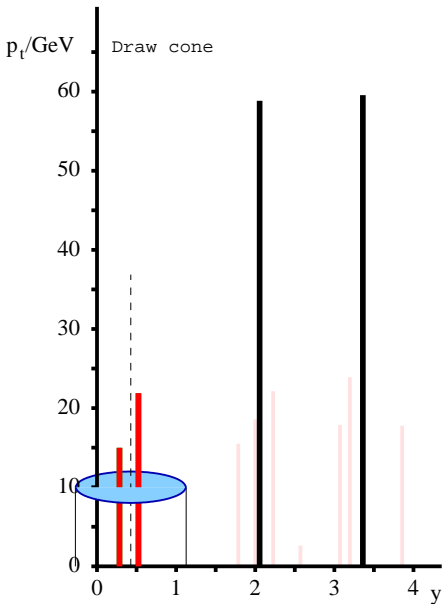
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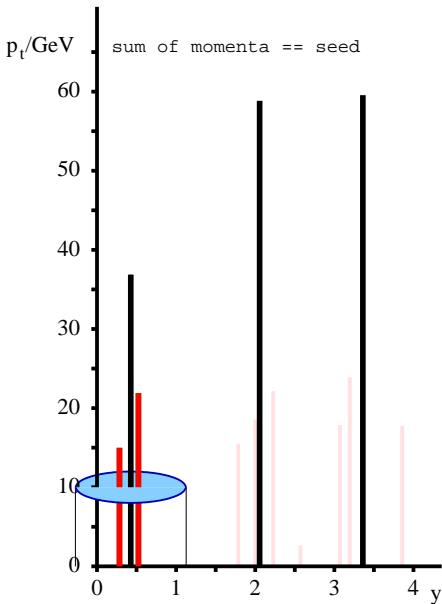
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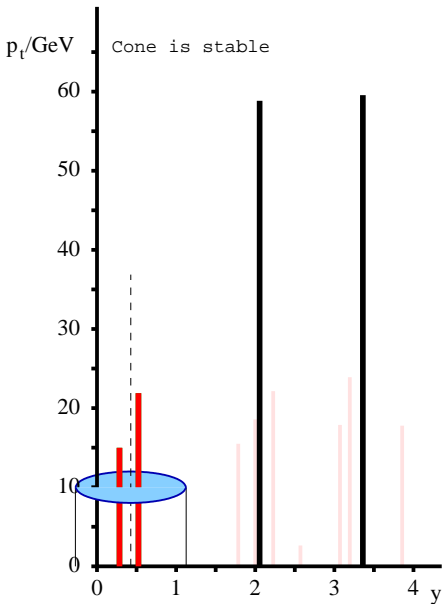
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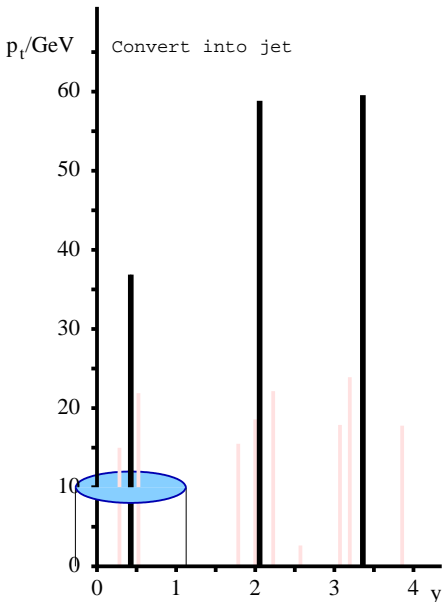
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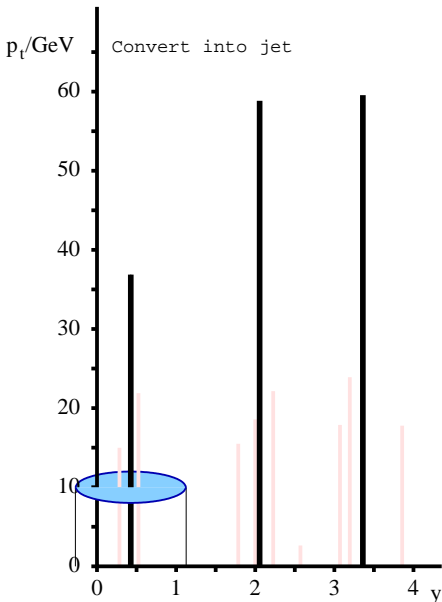
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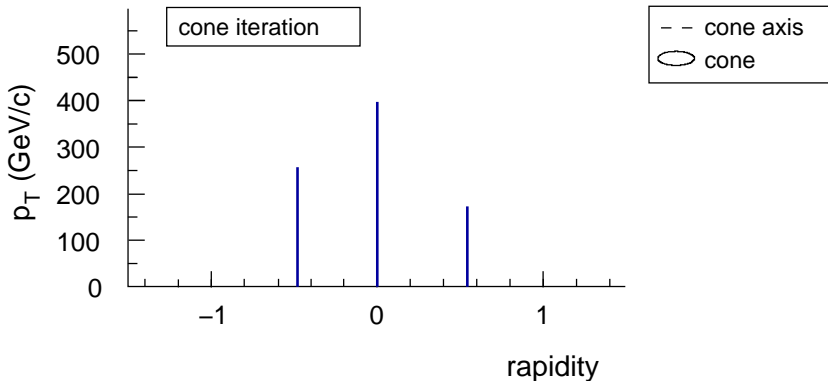
Next-simplest of the cones

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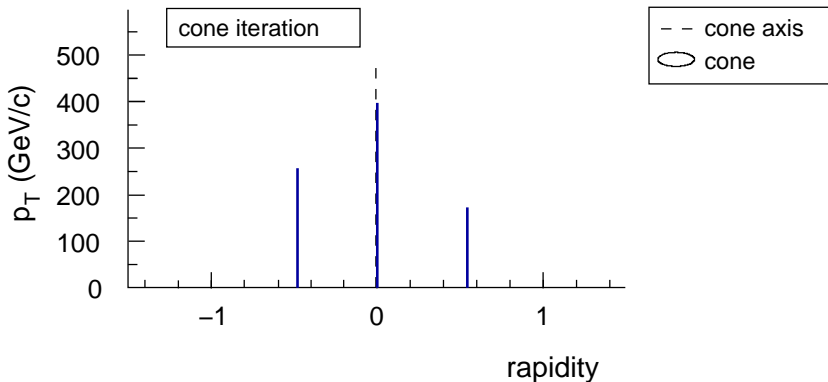
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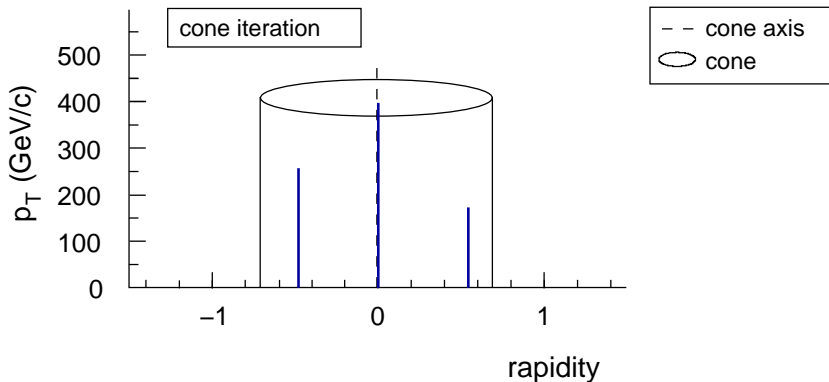
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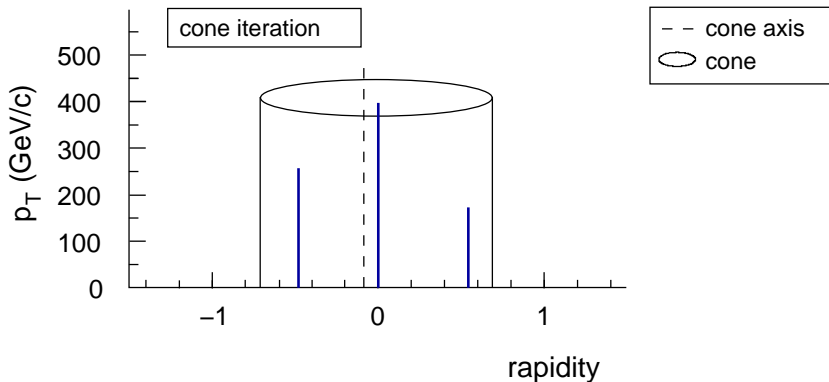
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



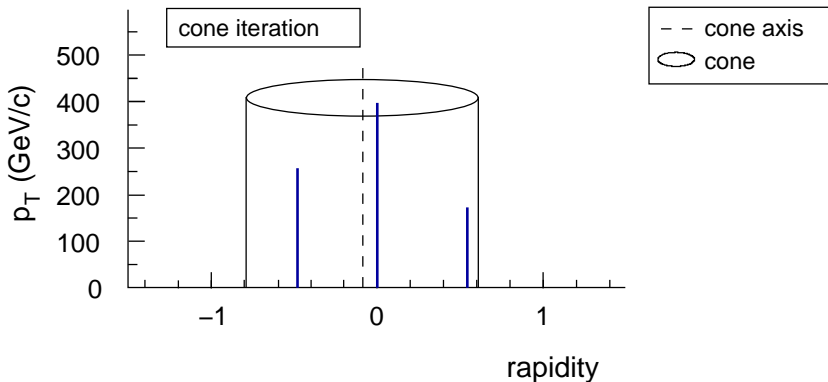
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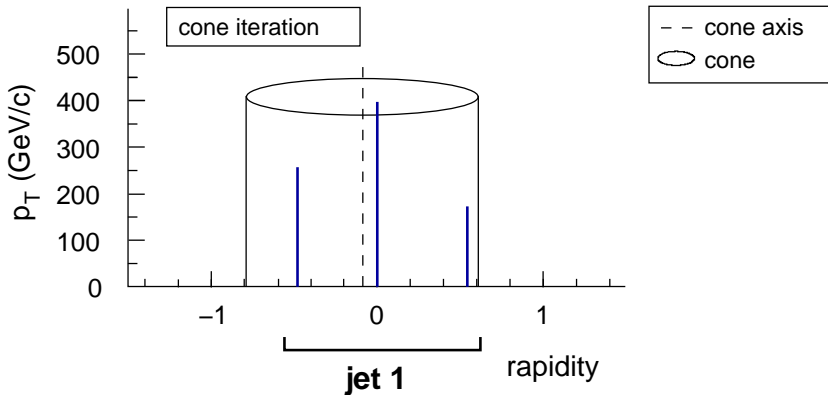
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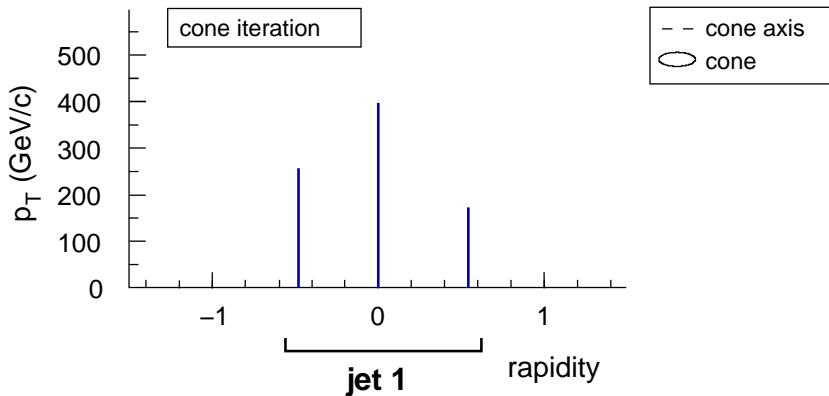
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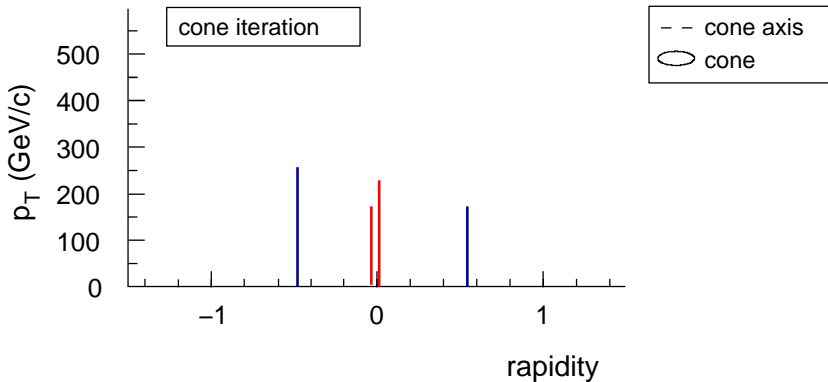
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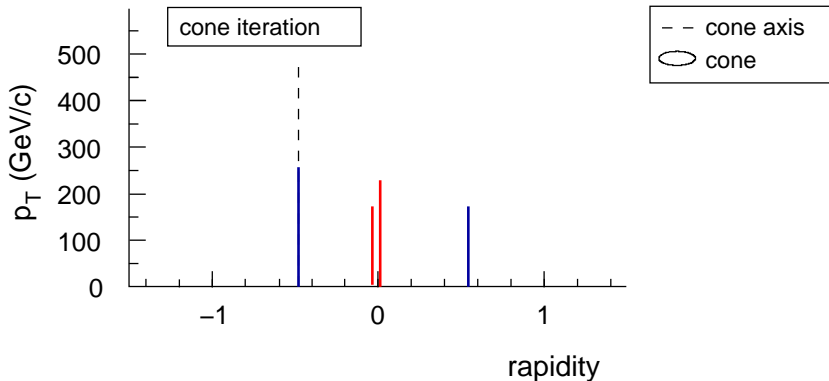
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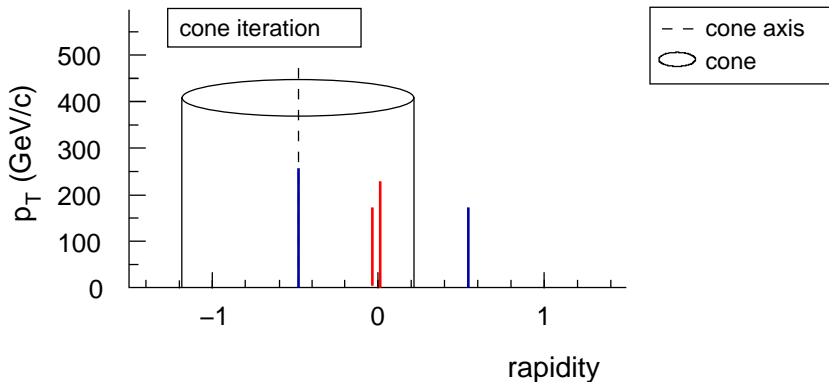
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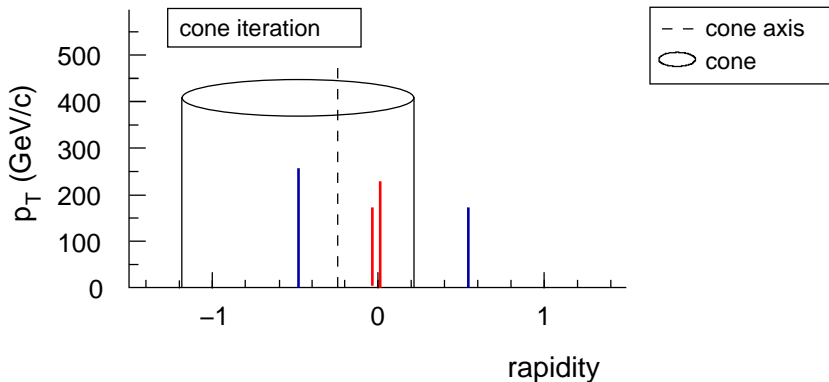
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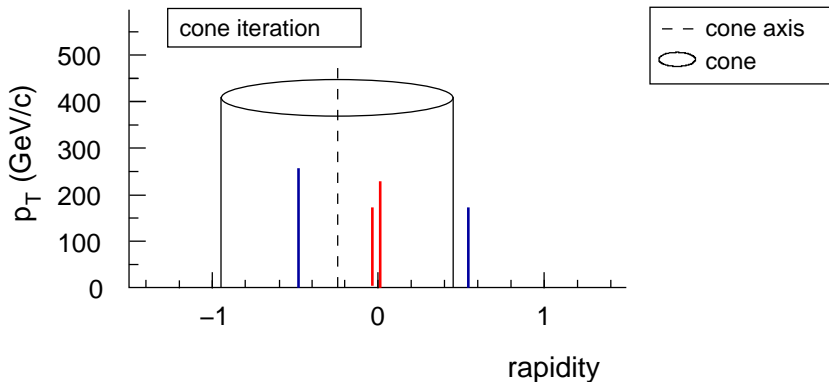
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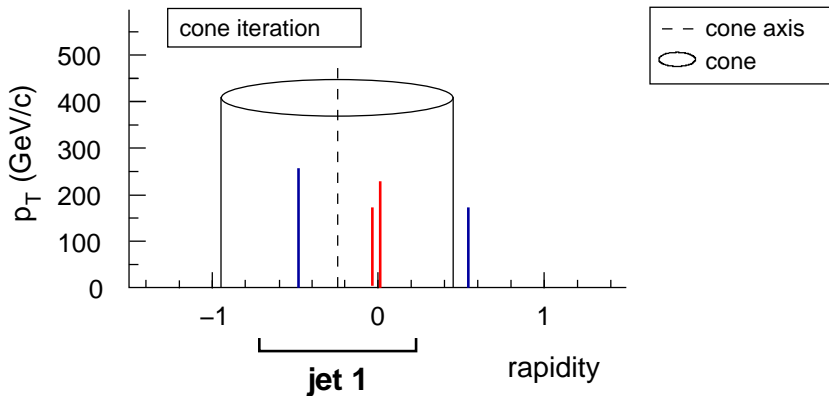
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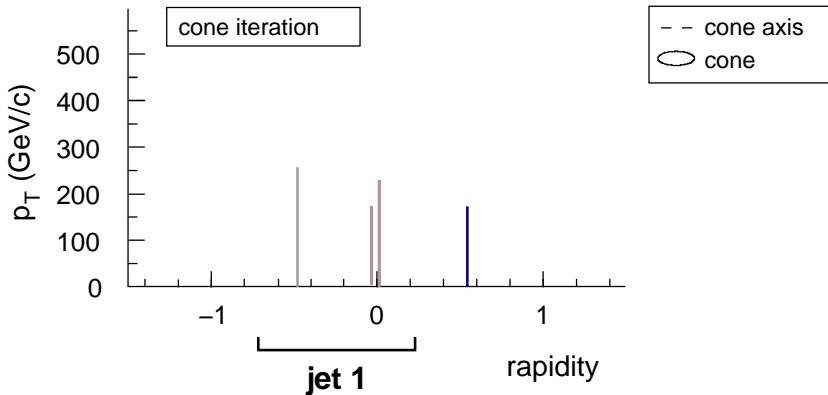
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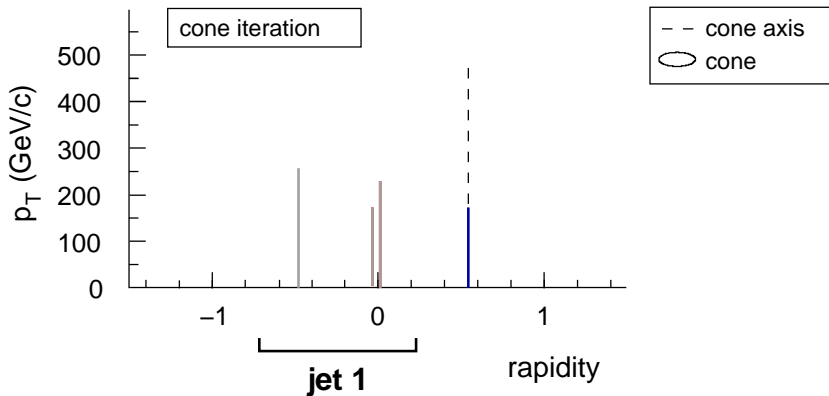
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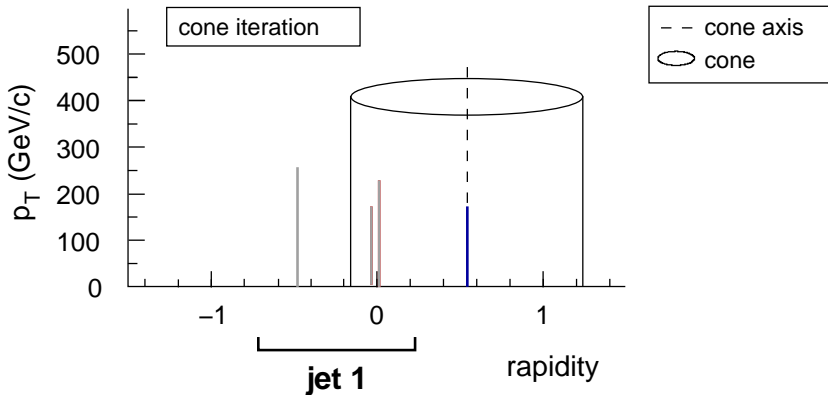
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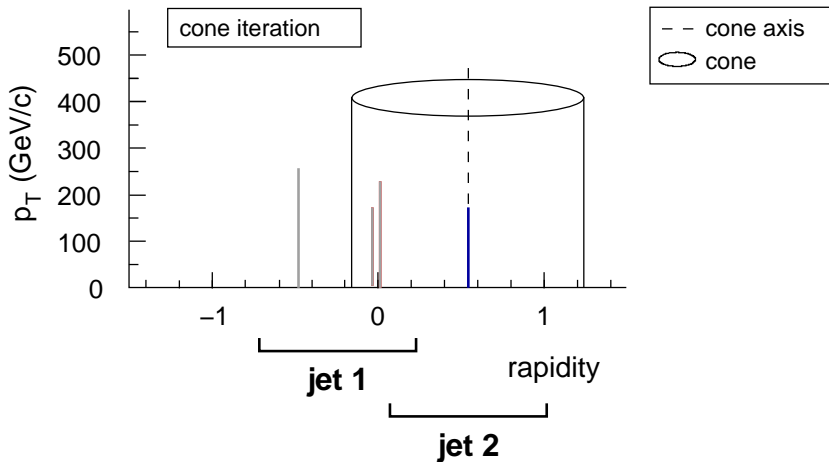
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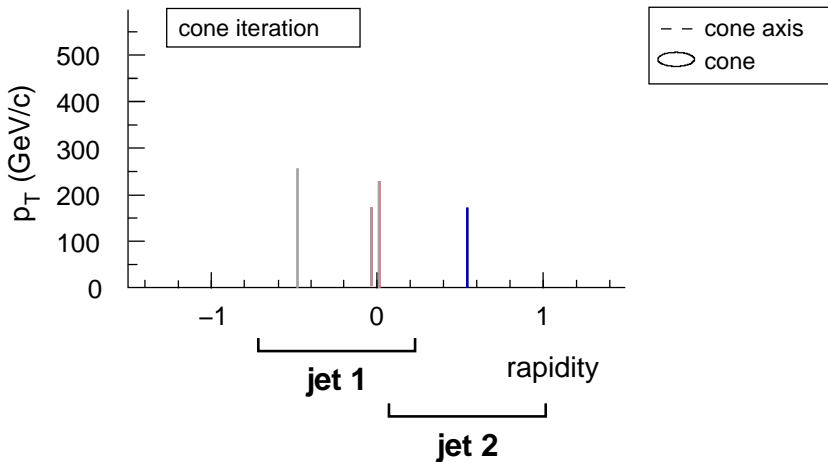
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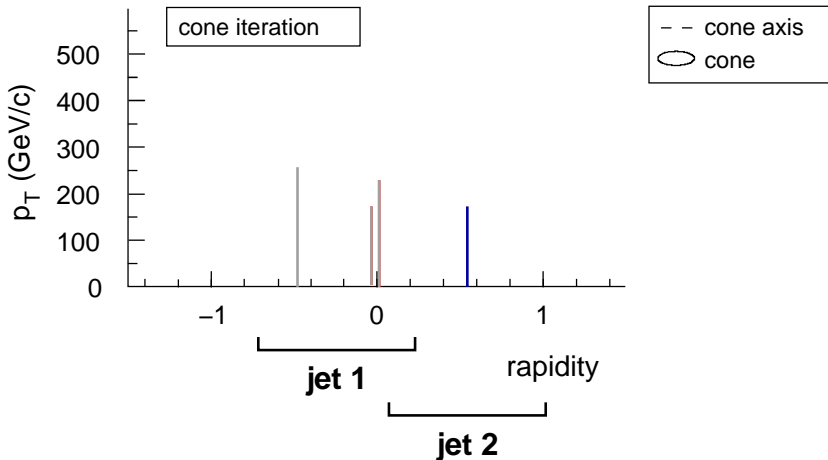
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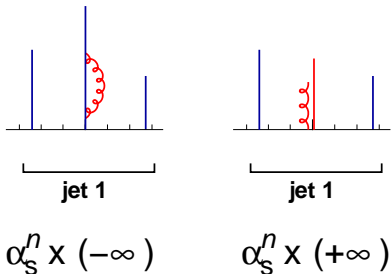


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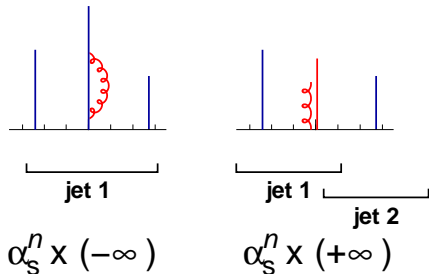
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Collinear Safe



Infinites cancel

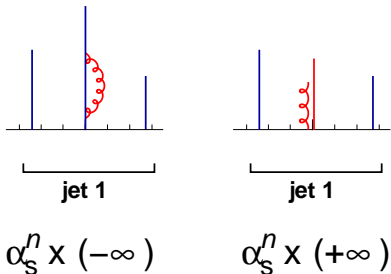
Collinear Unsafe



Infinites do not cancel

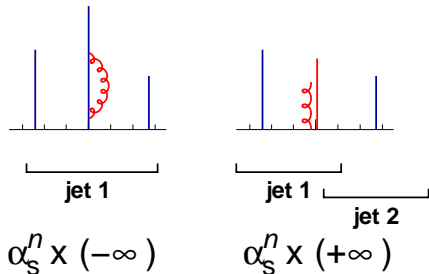
Invalidates perturbation theory

Collinear Safe



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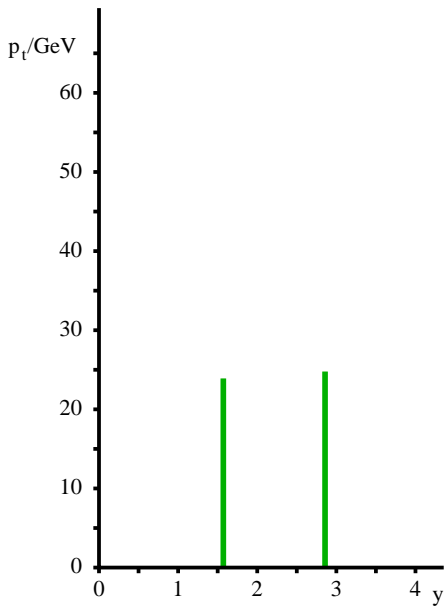
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So far

- ▶ We've seen sequential recombination jet algorithms
- ▶ And we've started looking at cone algorithms and run into problems

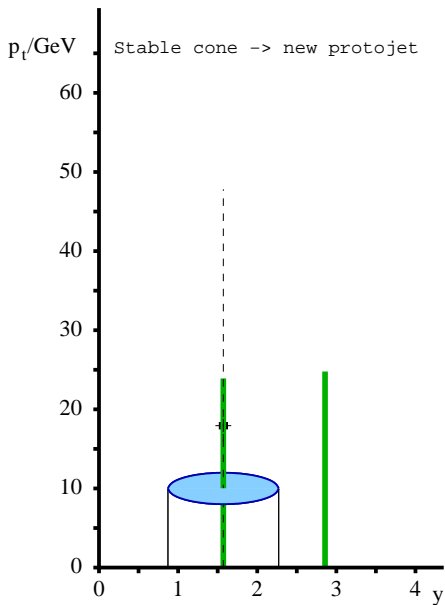
Tomorrow

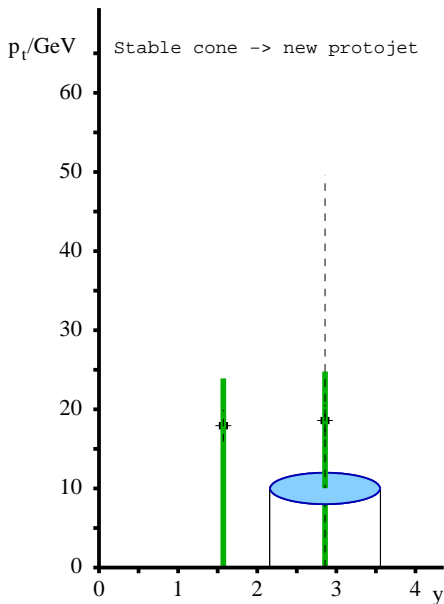
- ▶ Continue with the cones See more problems + some solutions
- ▶ Take a look at the physics of jet algorithms

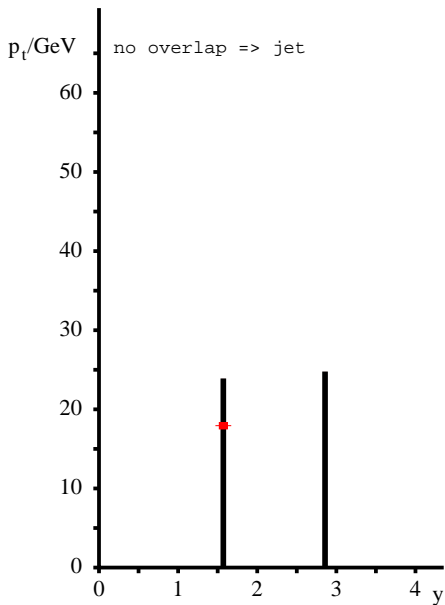


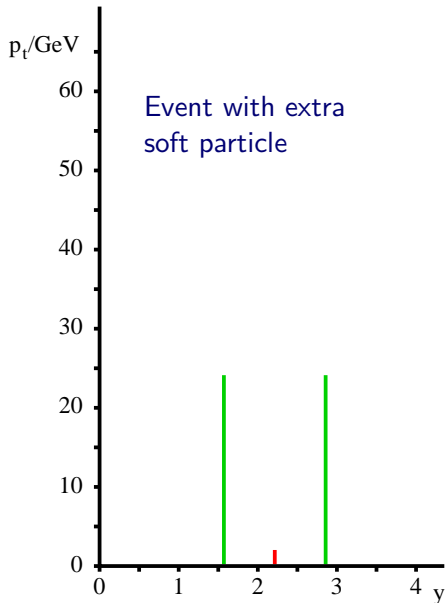
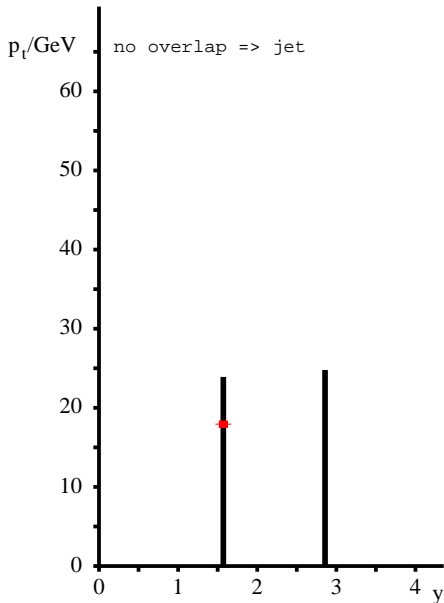
Event with extra

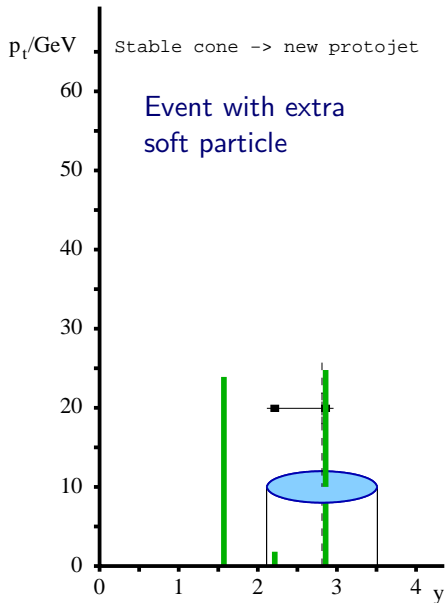
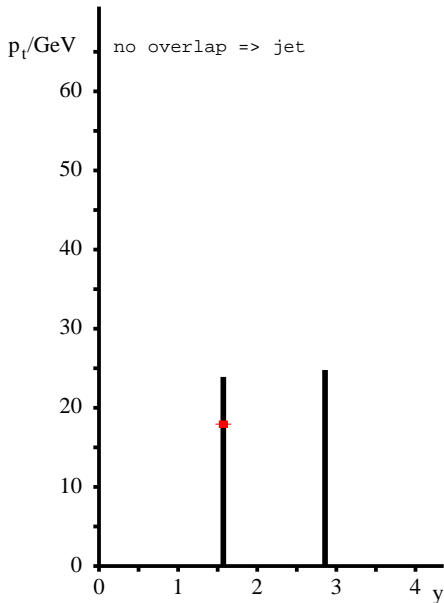
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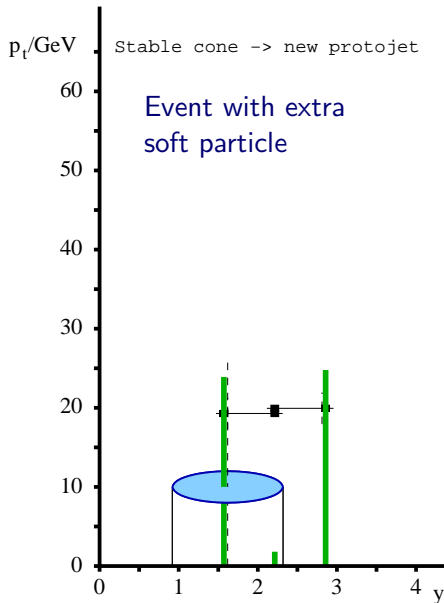
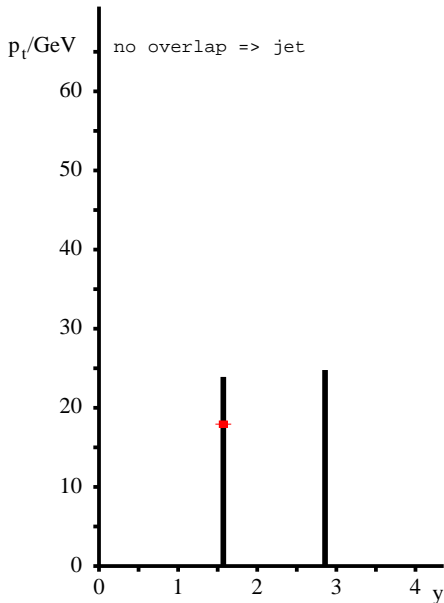


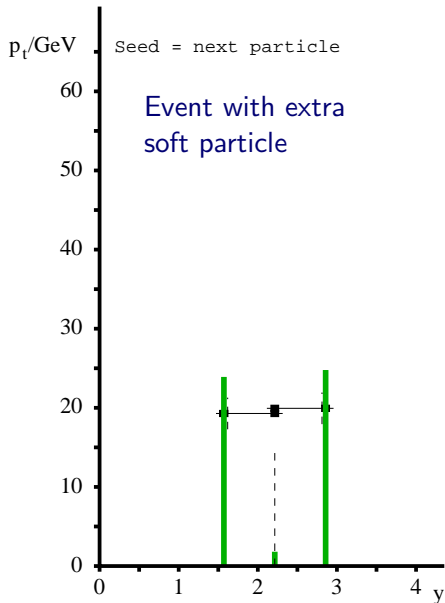
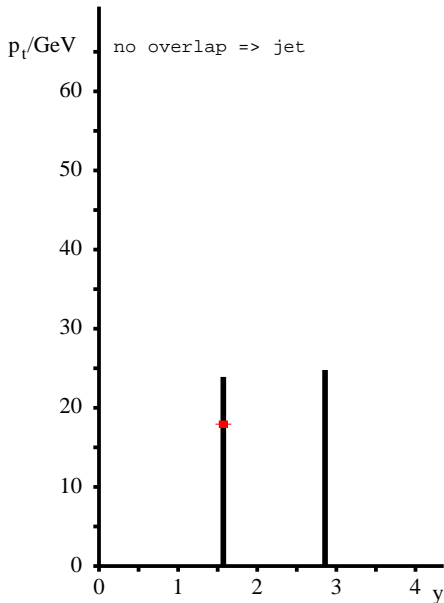


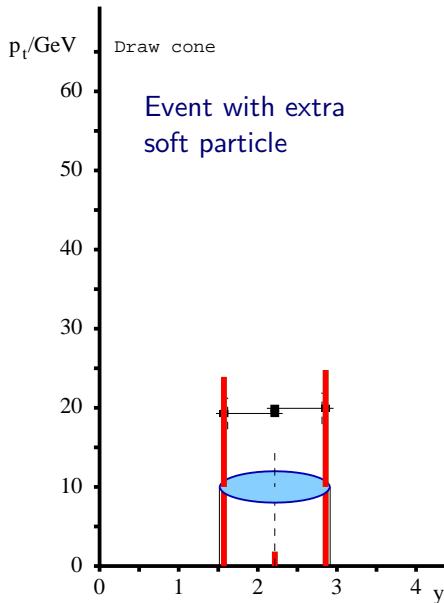
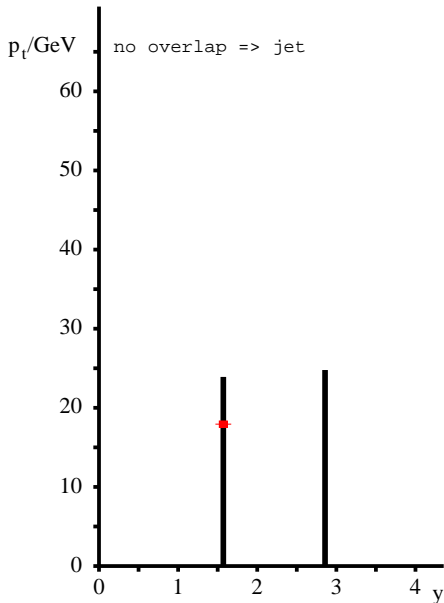


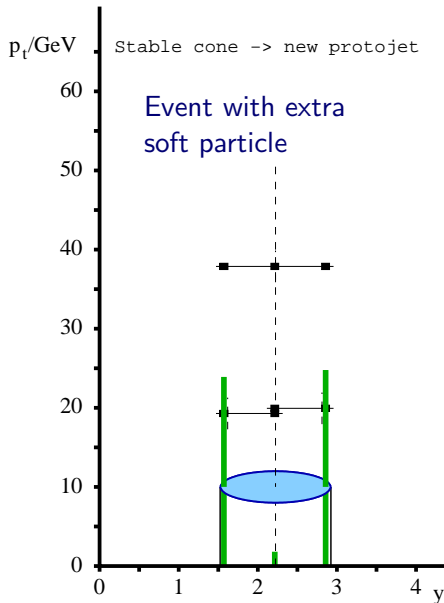
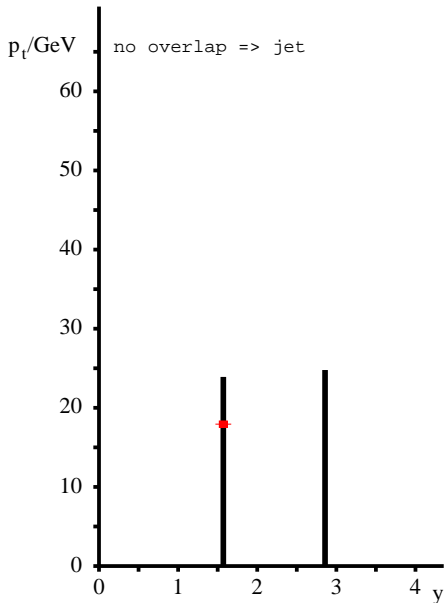


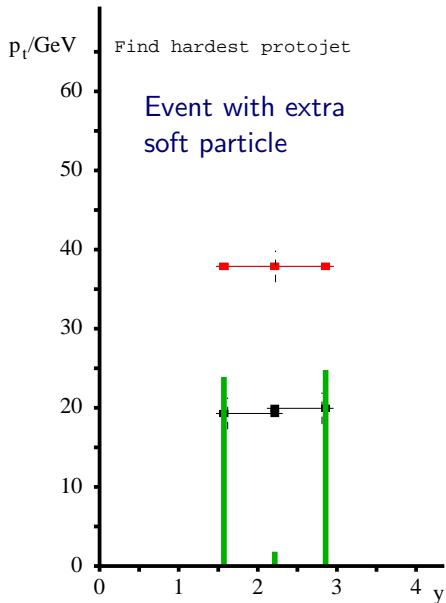
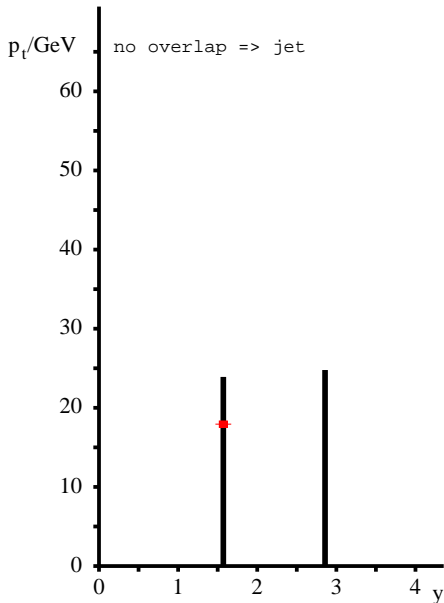


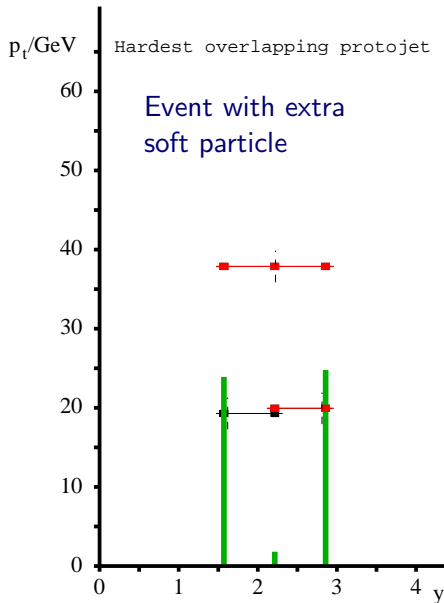
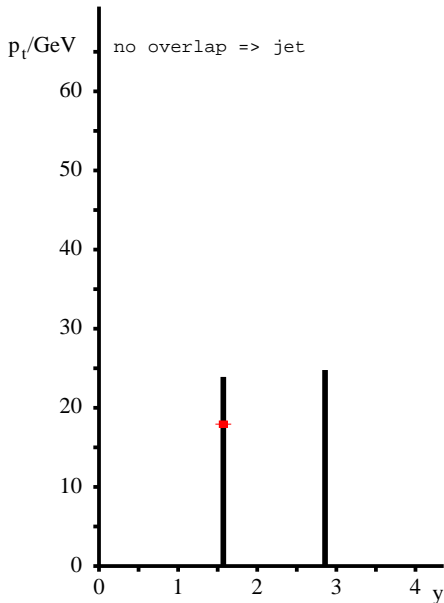


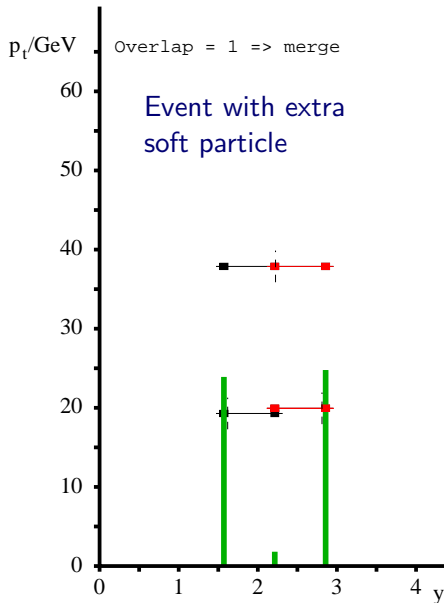
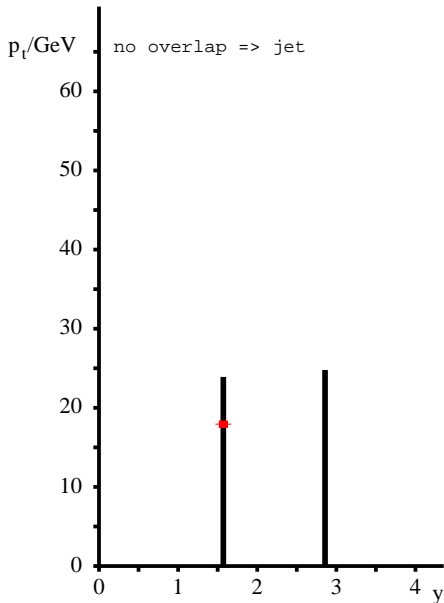


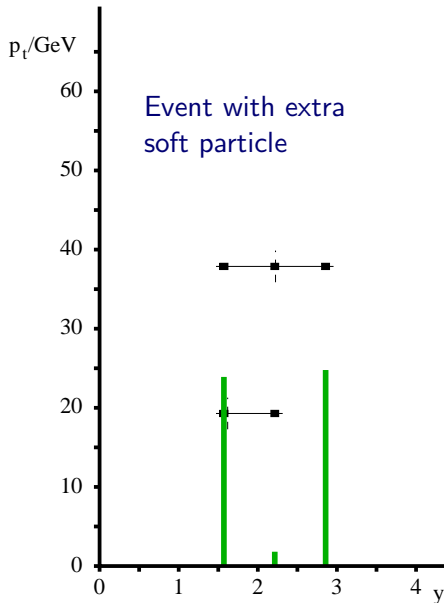
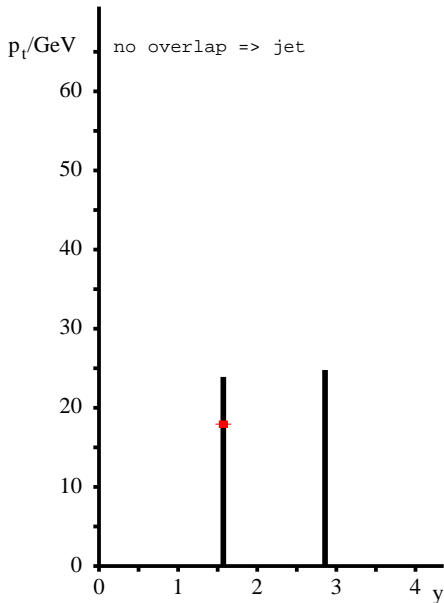


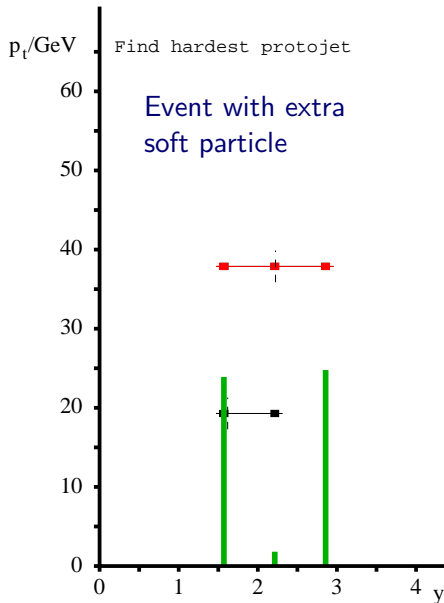
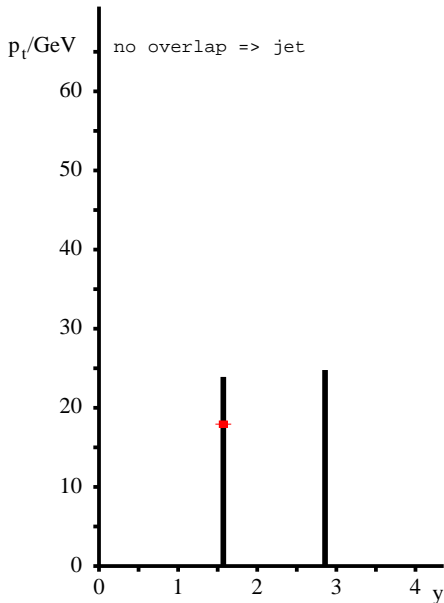


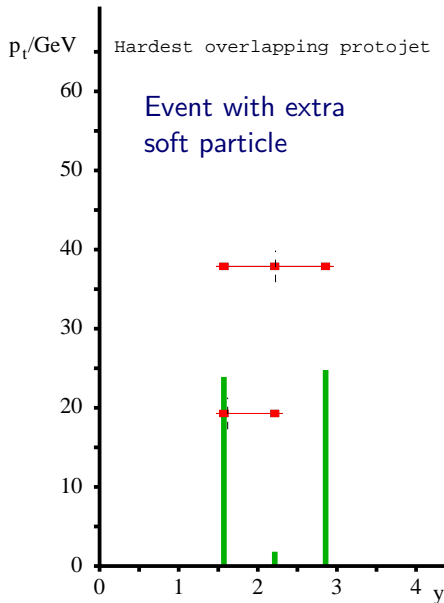
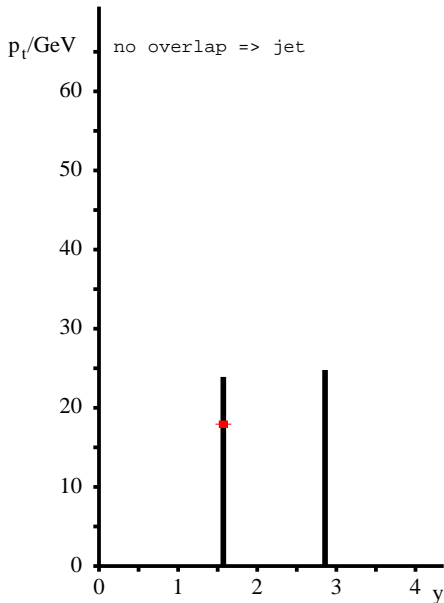


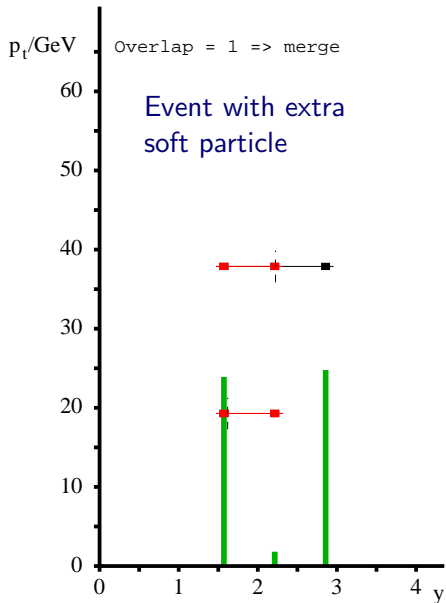
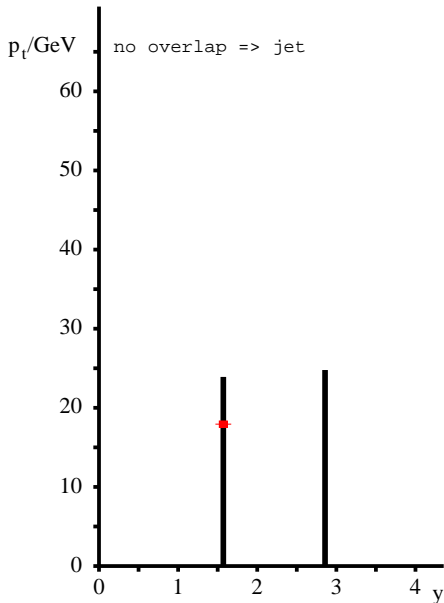


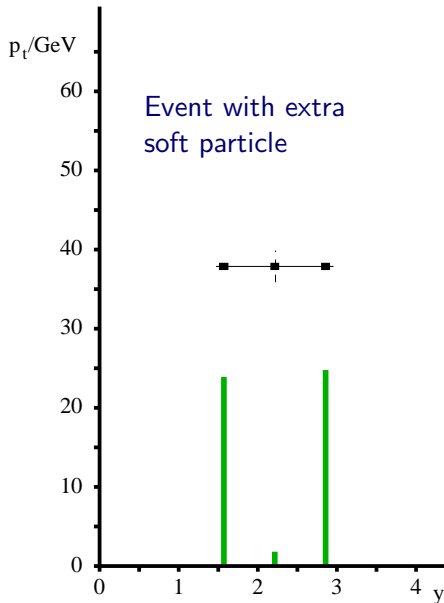
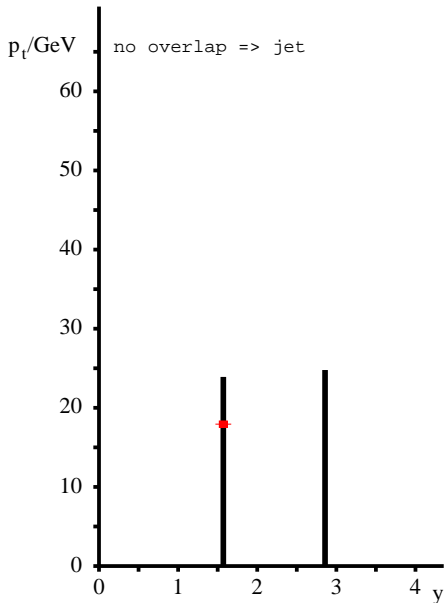


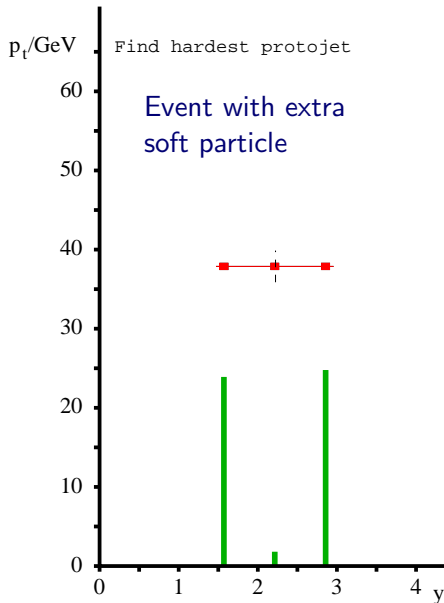
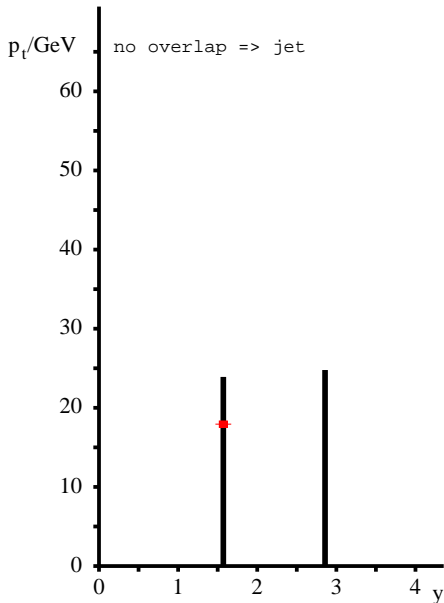


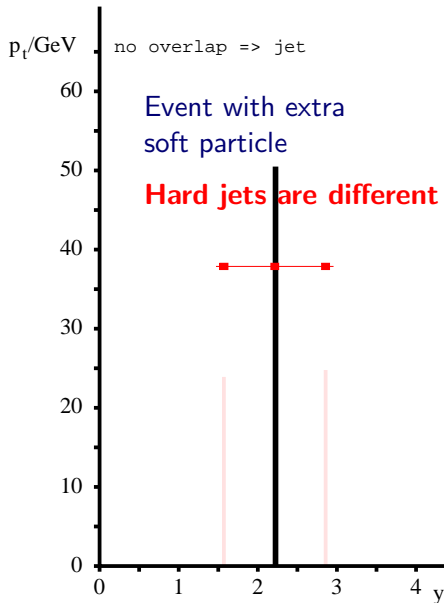
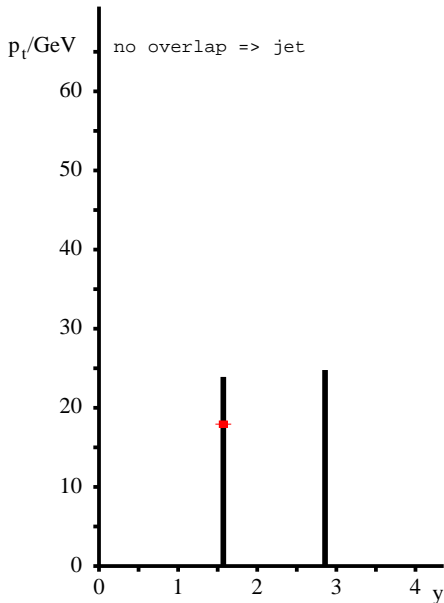












In lecture 1, we saw

- ▶ sequential recombination (k_t , etc.) algorithms
- ▶ the first of a series of cone-algorithms, those with “progressive removal” (xC-PR)
- ▶ and ran into collinear safety issues (from ordering of “seeds” for cone direction)

Today

- ▶ see the other series of cone-algorithms (with split-merge, xC-SM)
- ▶ look more at the physics of jet algs.

Unifying idea: momentum flow within a cone only
 marginally modified by QCD branching

But cones come in many variants

Finding cones \ Processing	Progressive Removal	Split-Merge	Split-Drop
Seeded, Fixed (FC)	GetJet CellJet		
Seeded, Iterative (IC)	CMS Cone	JetClu (CDF) [†] ATLAS cone	
Seeded, It. + Midpoints (IC _{mp})		CDF MidPoint D0 Run II cone	PxCone
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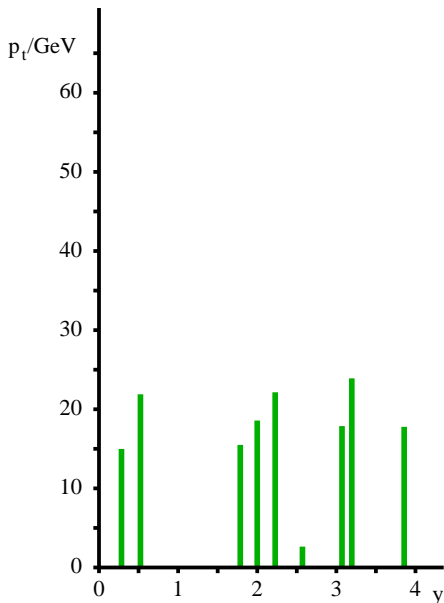
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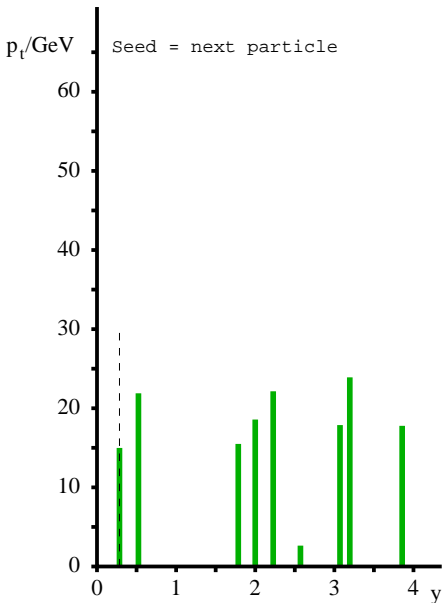
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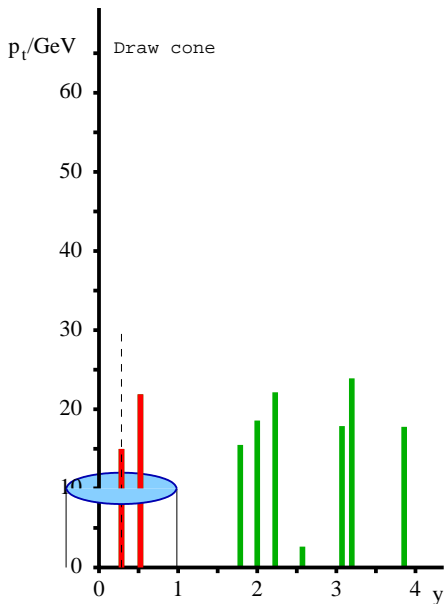
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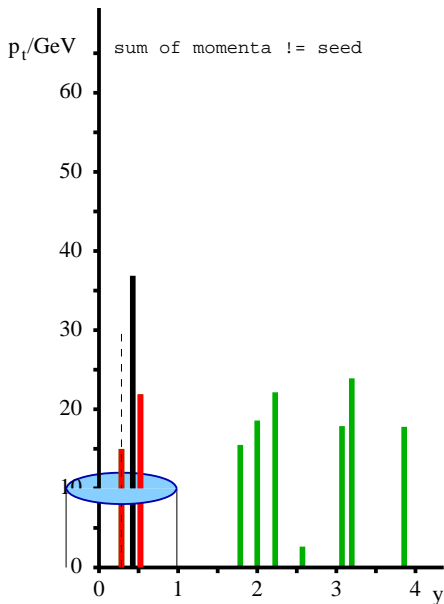
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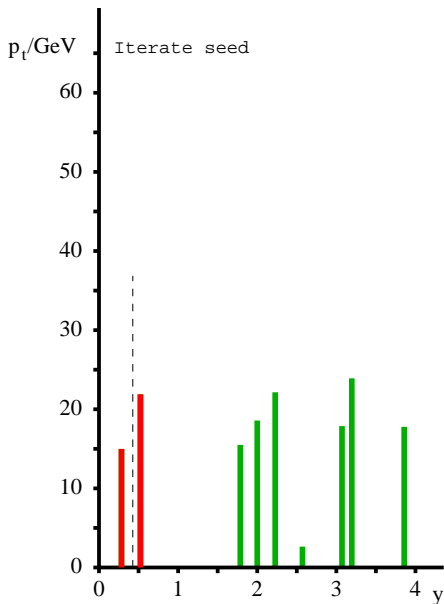
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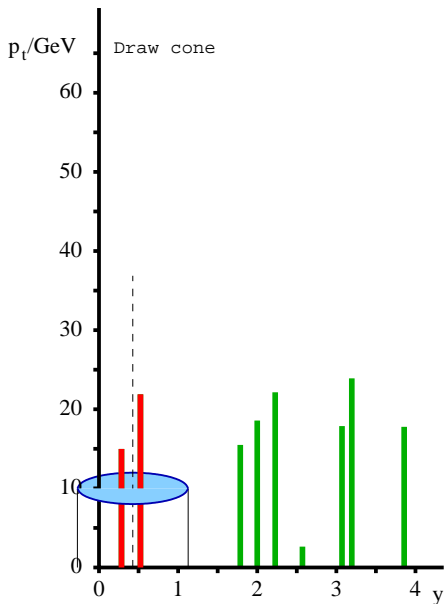
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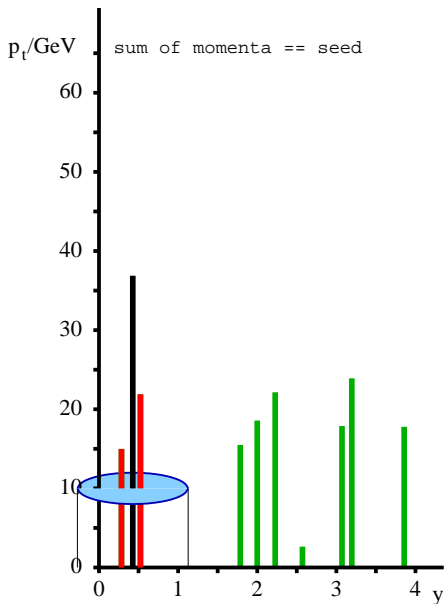
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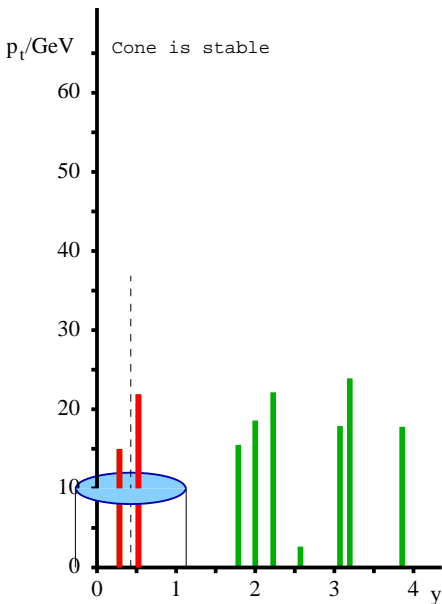
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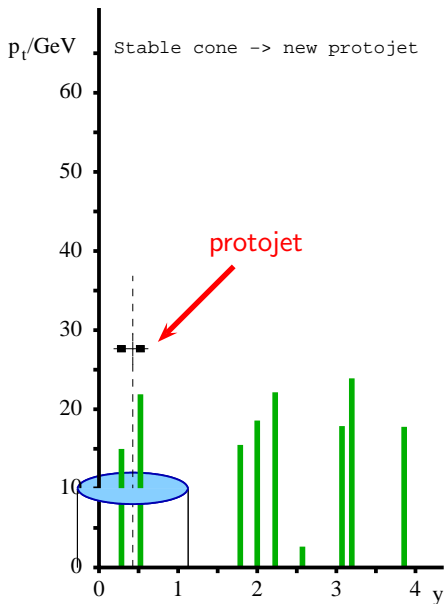
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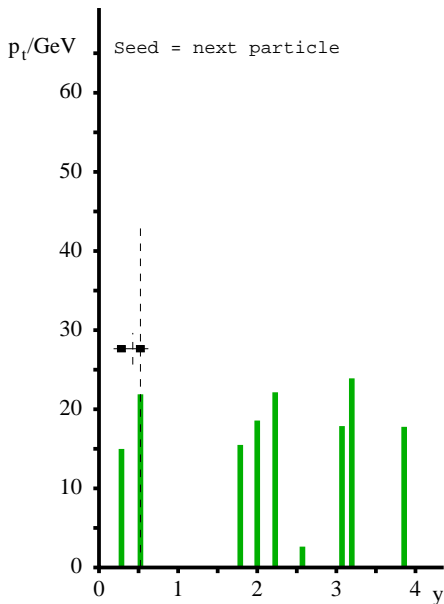
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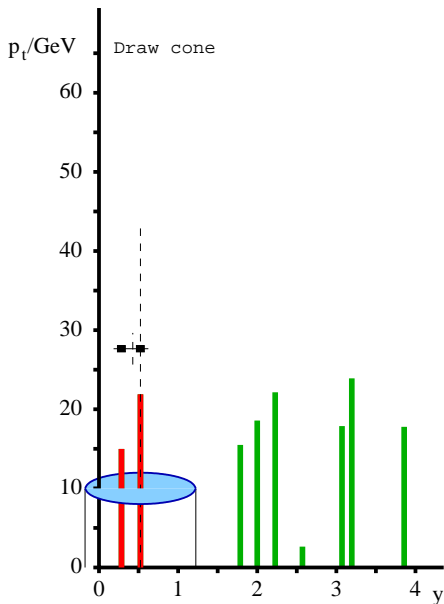
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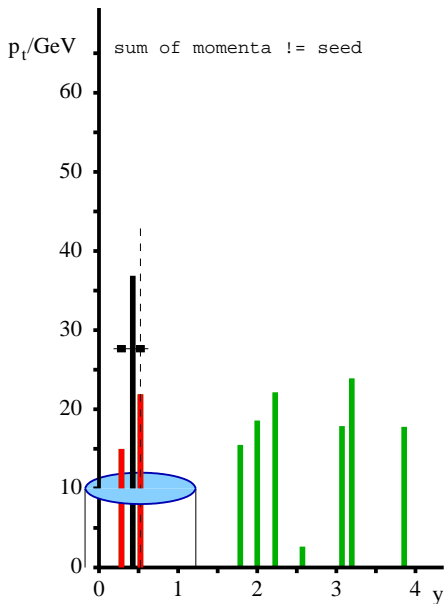
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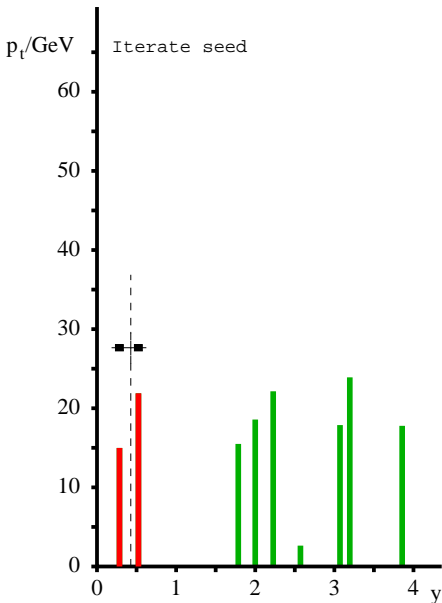
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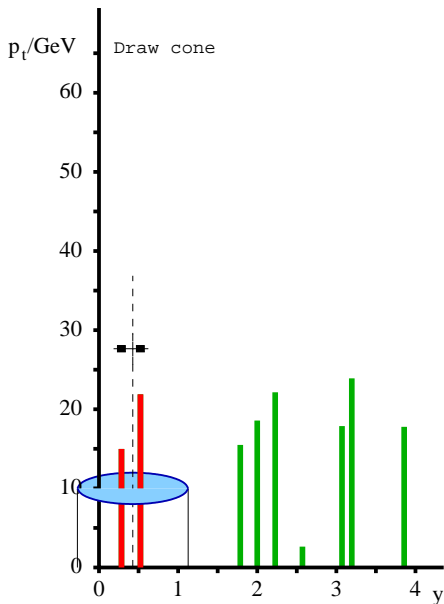
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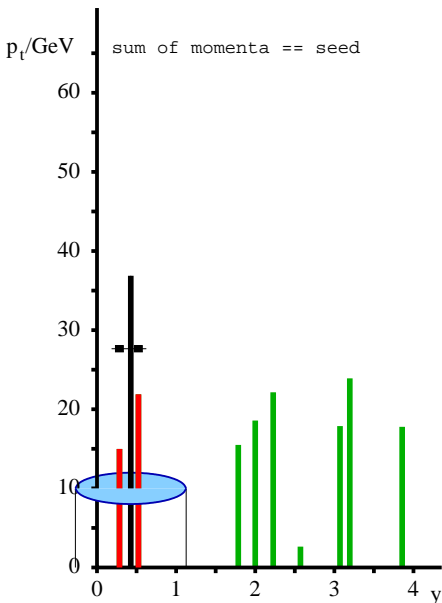
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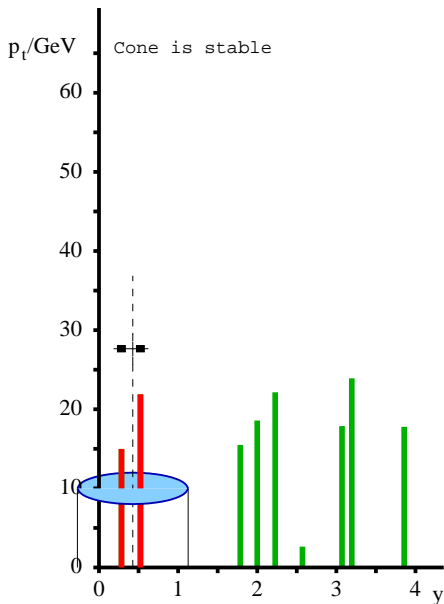
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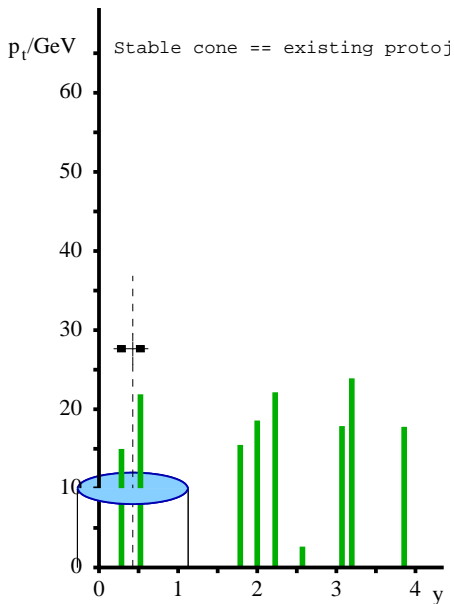
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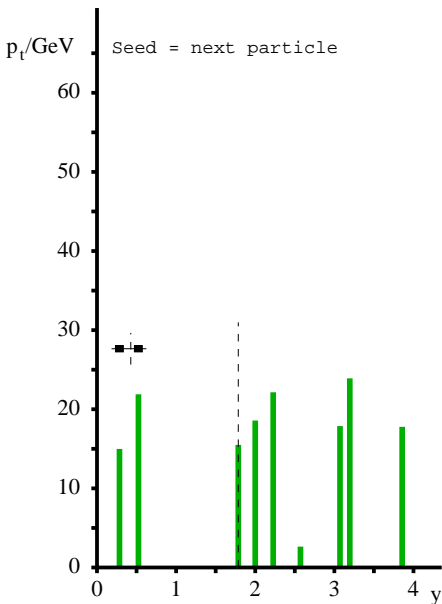
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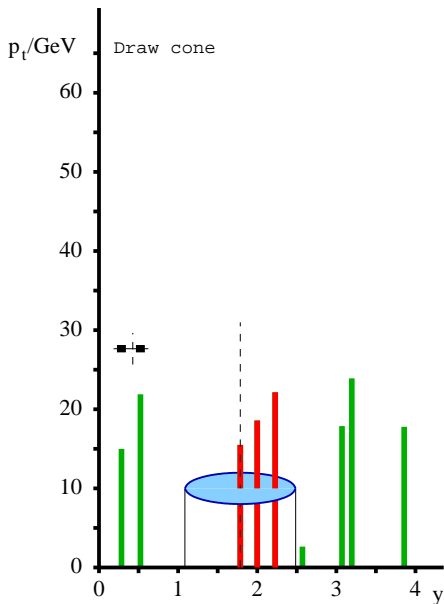
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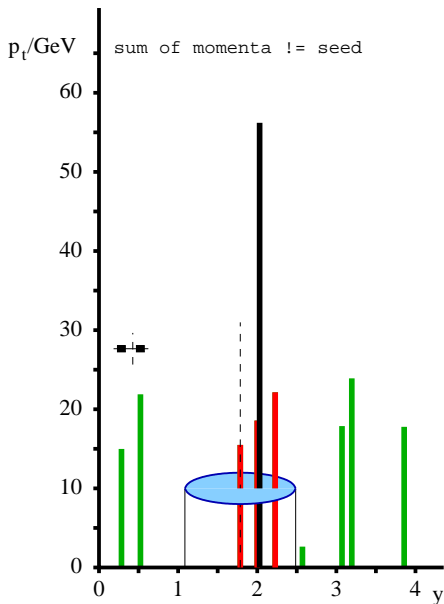
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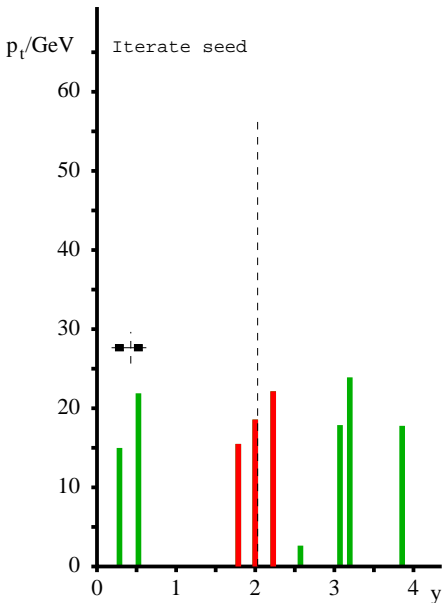
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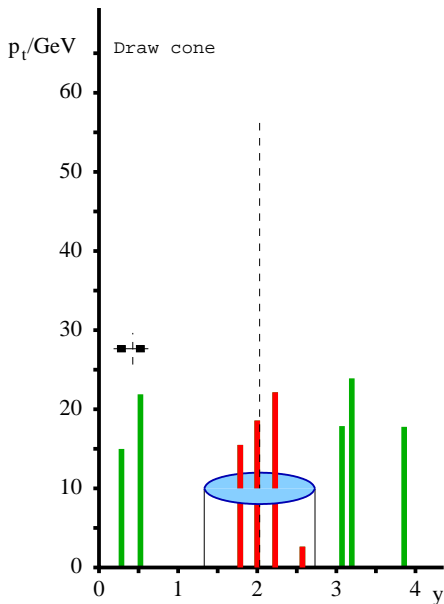
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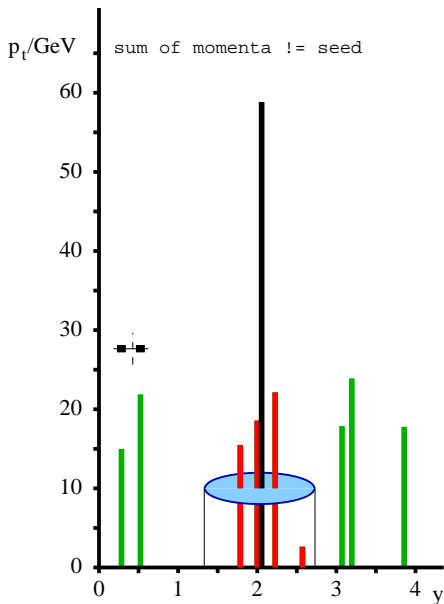
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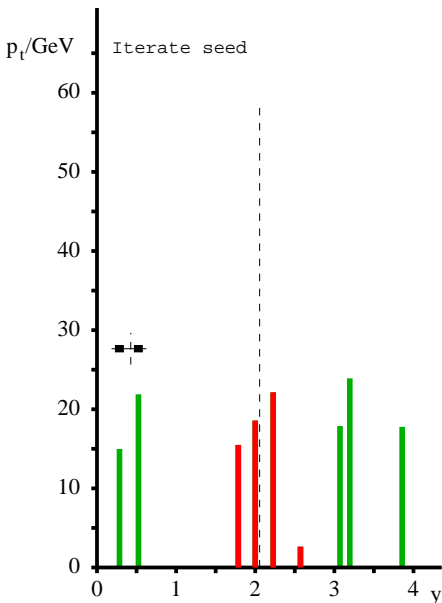
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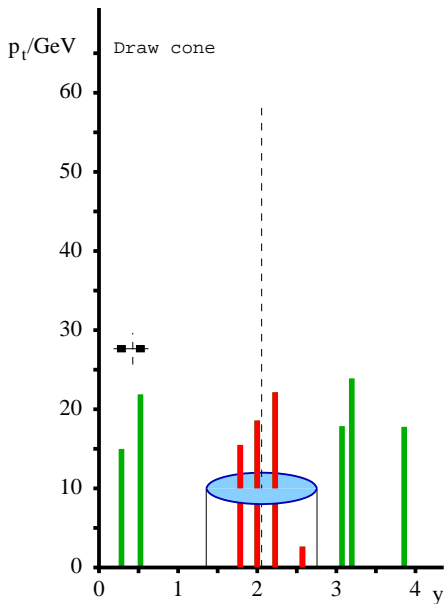
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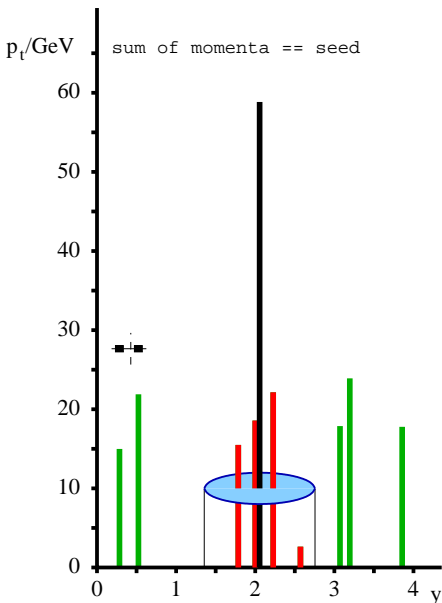
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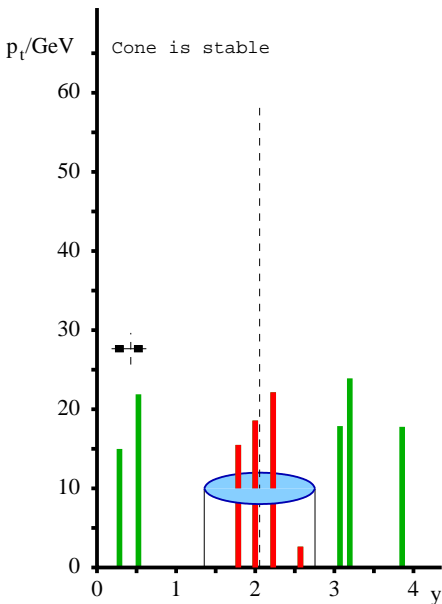
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It. Cone with Split-Merge (IC-SM)



Avoid ordering seeds (coll. unsafe)

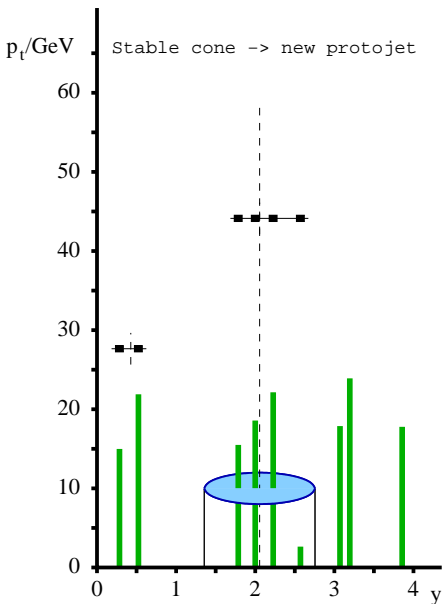
CDF JetClu[†] & ATLAS cones

- ▶ use every particle as possible seed (no particular order)
- ▶ **iterate until stable cone**
- ▶ add the stable cone to the list of **protojets** unless it's already there
- ▶ until all seeds done

Note: protojets **overlap**. Certain particles appear in many protojets
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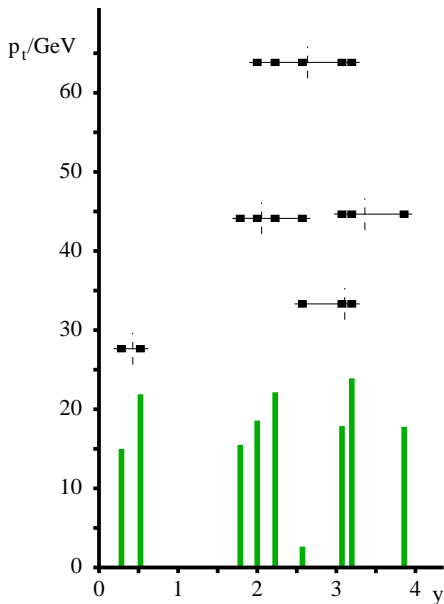
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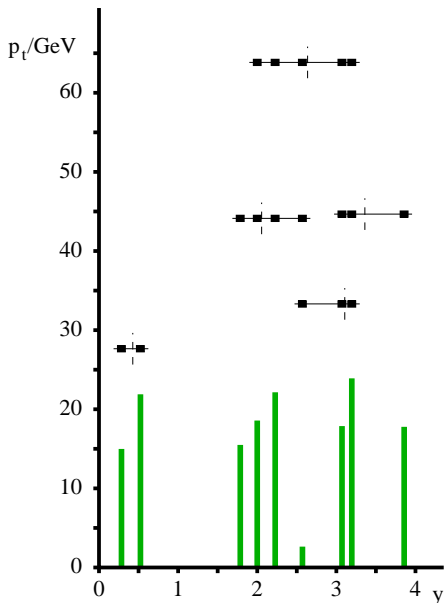
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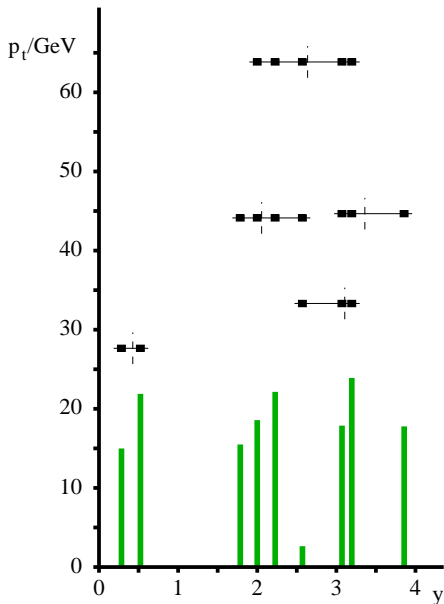
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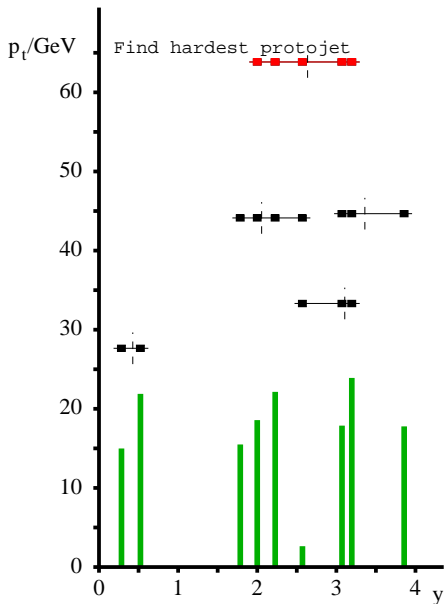


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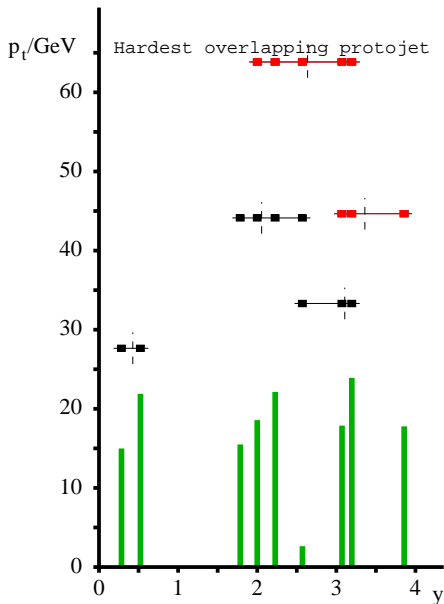


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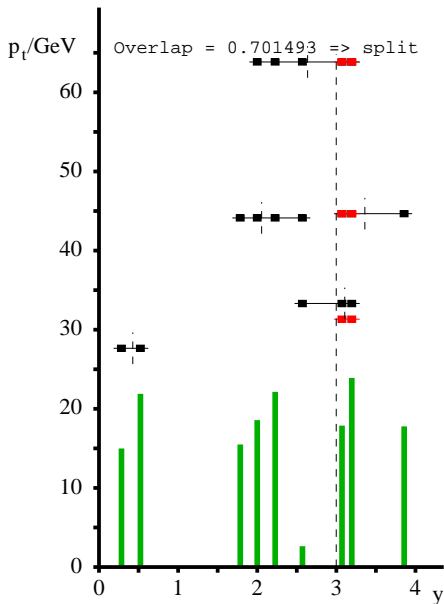


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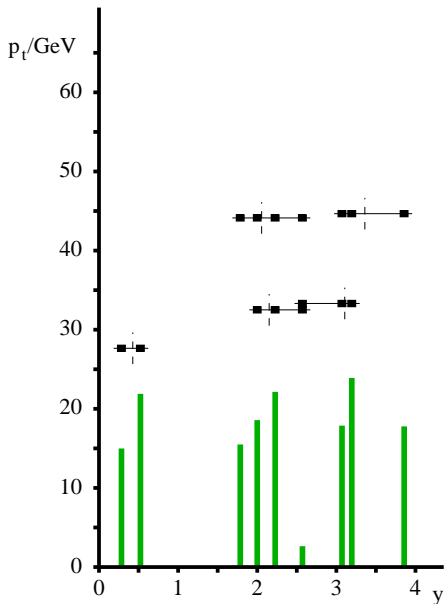


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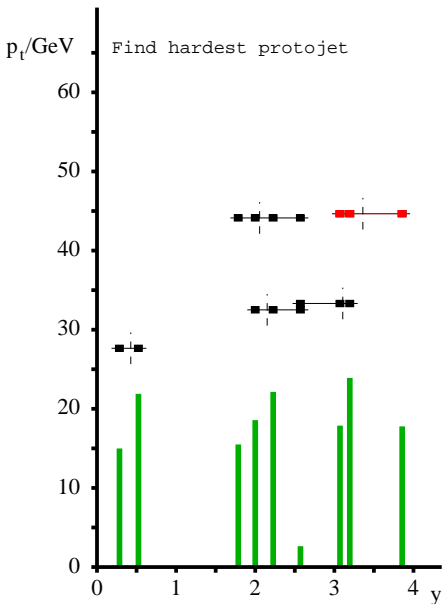


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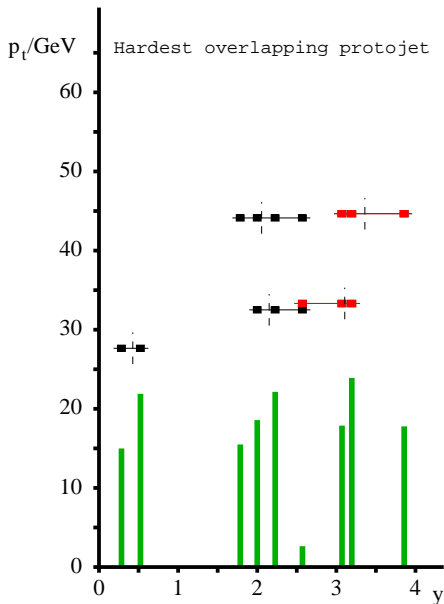


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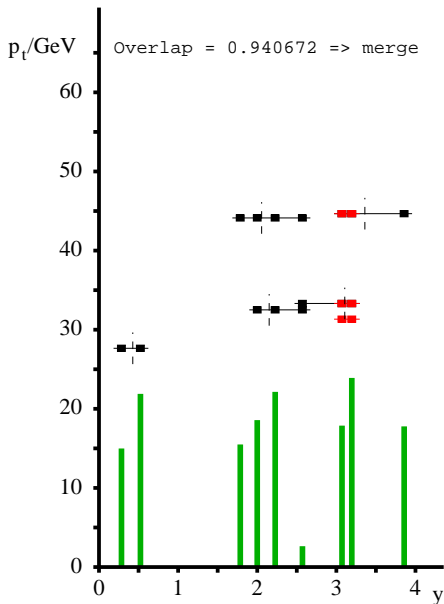


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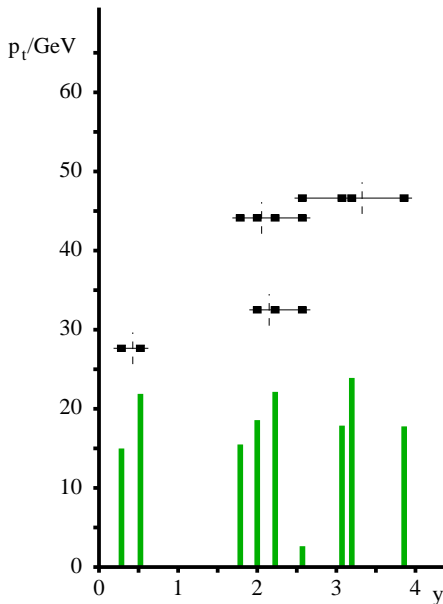


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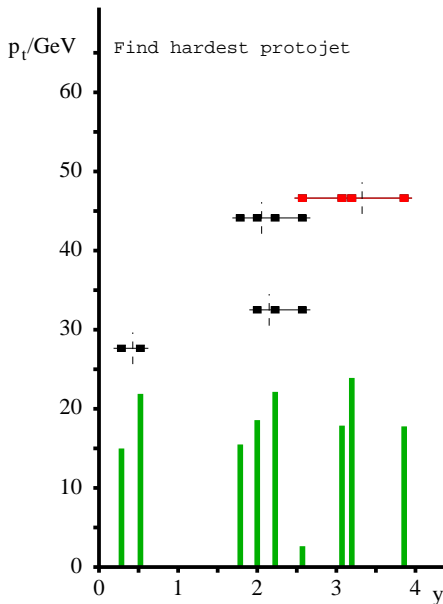


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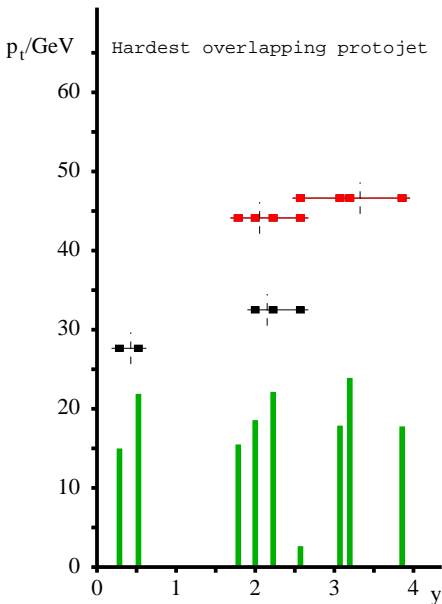


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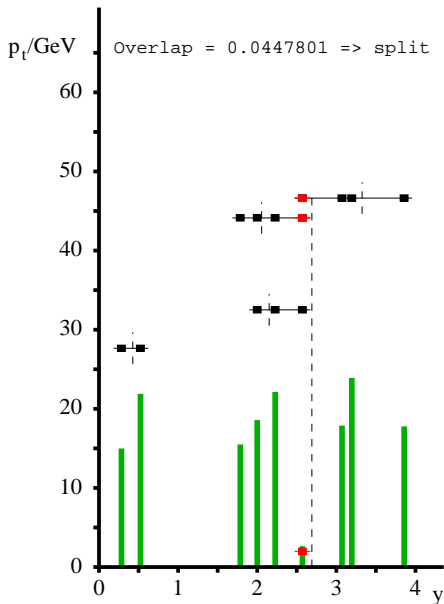


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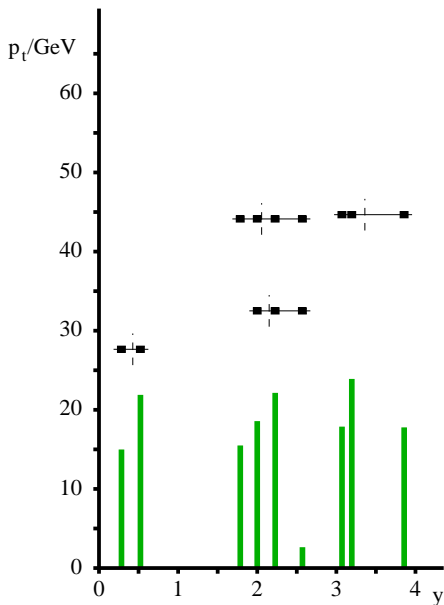


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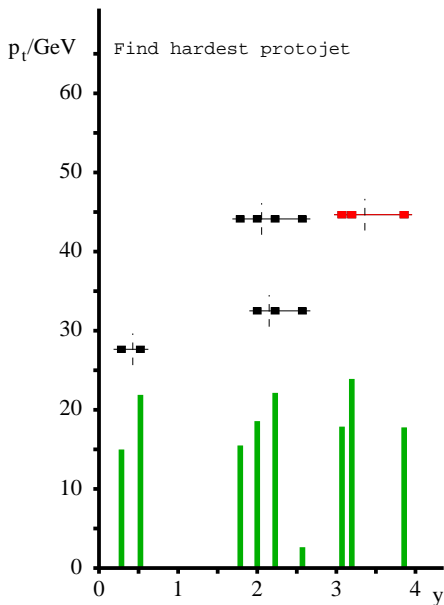


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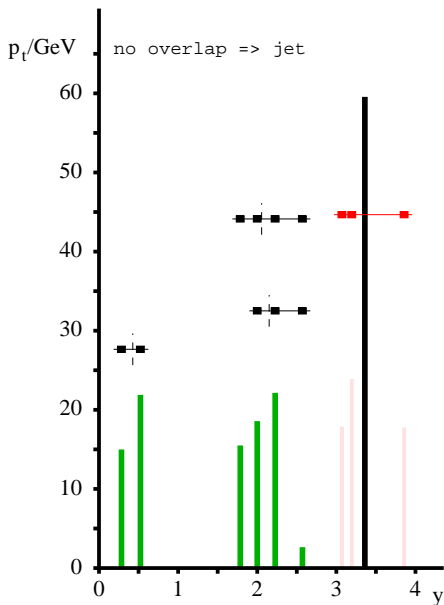


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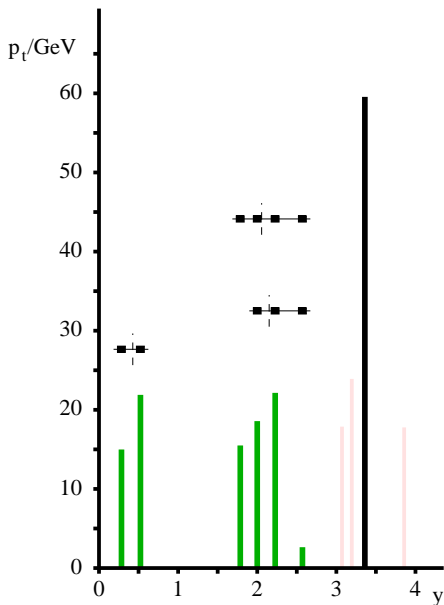


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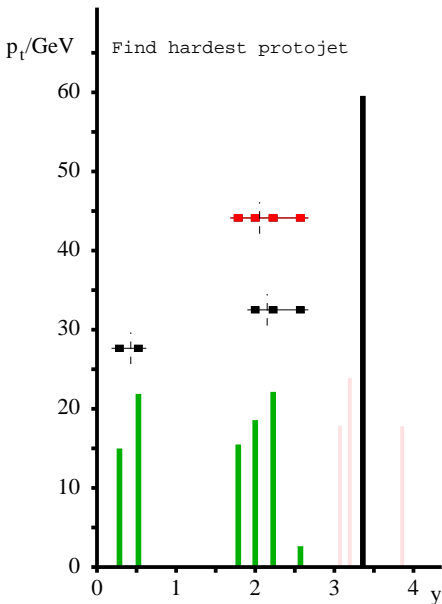


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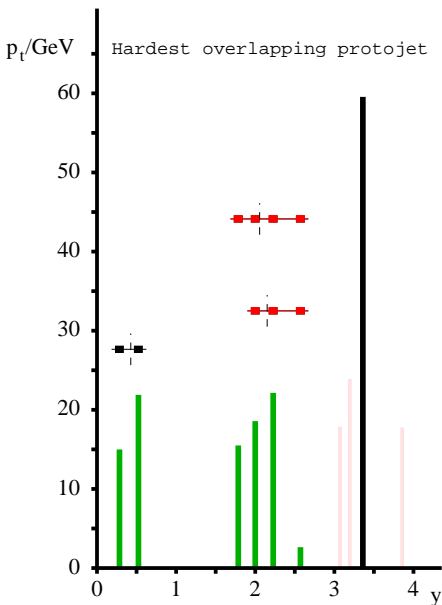


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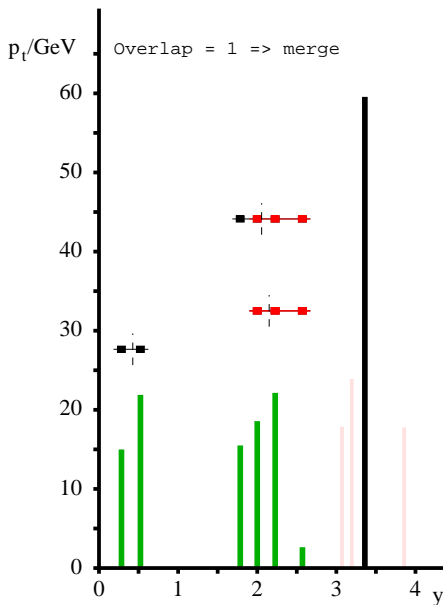


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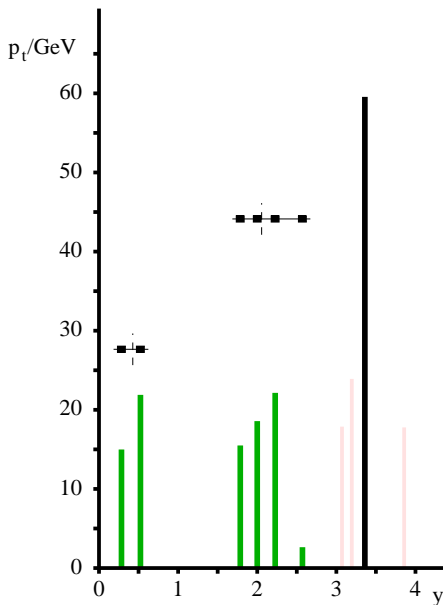


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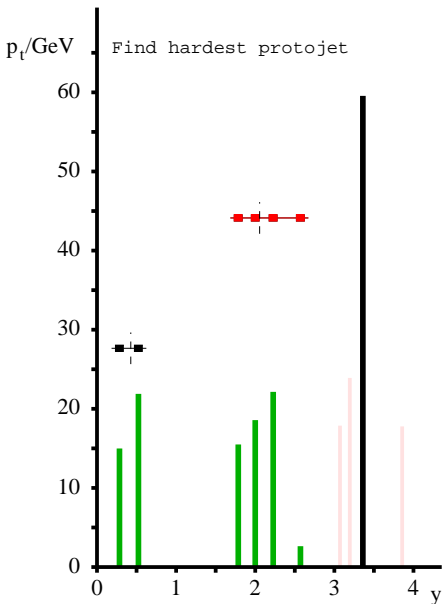


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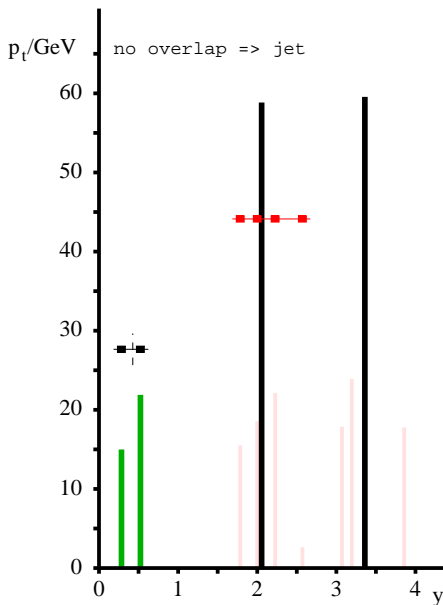


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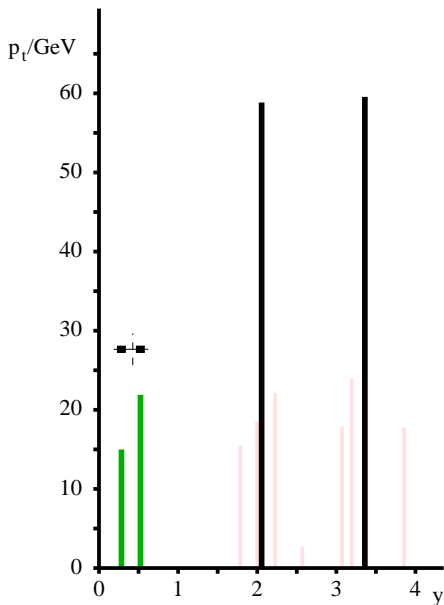


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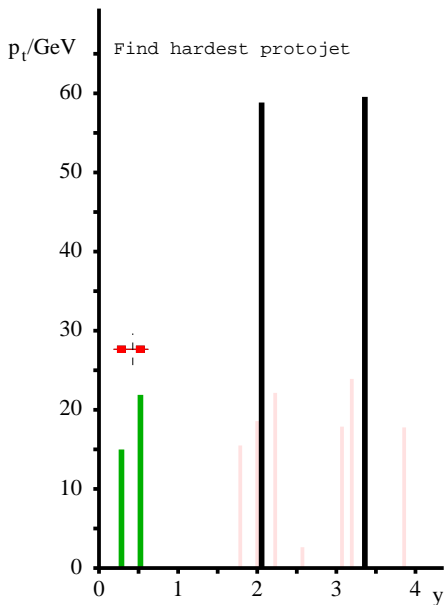


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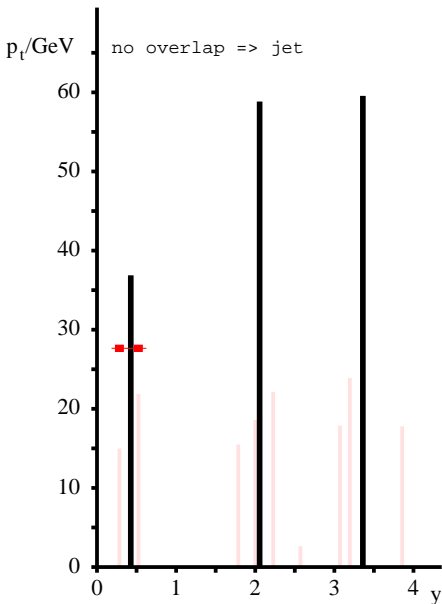


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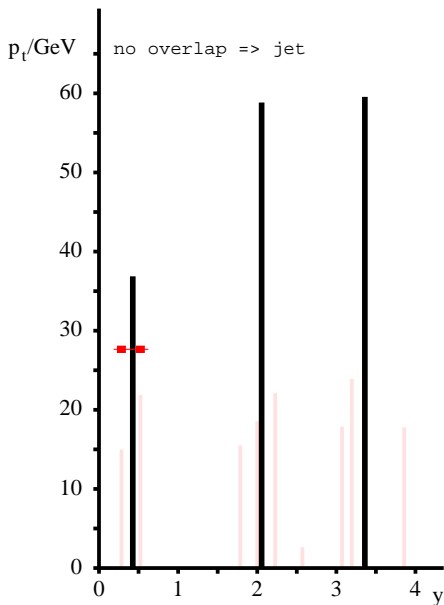


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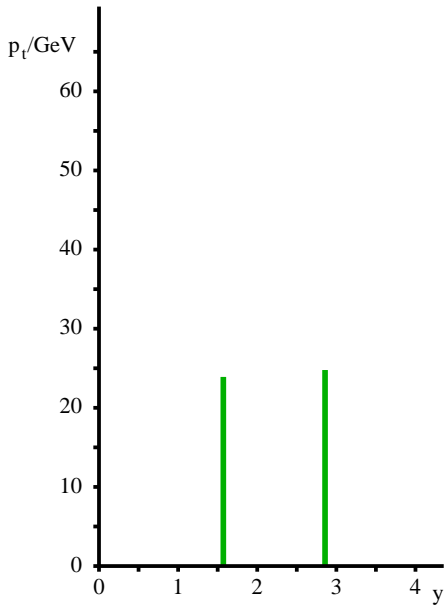


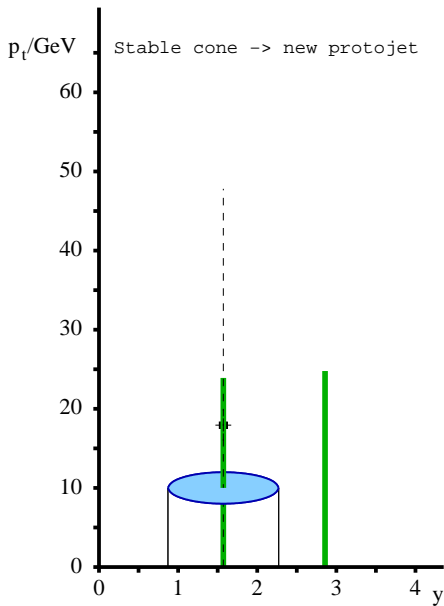
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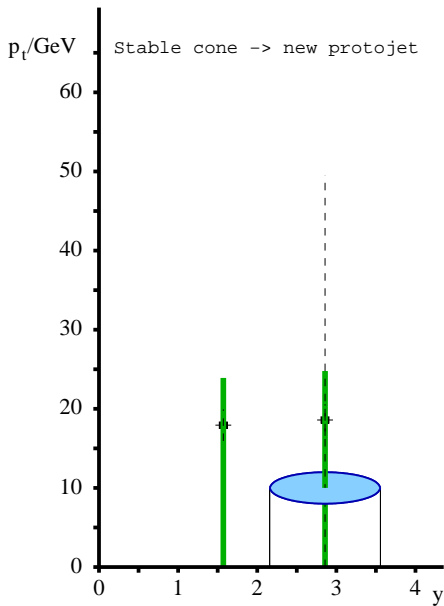
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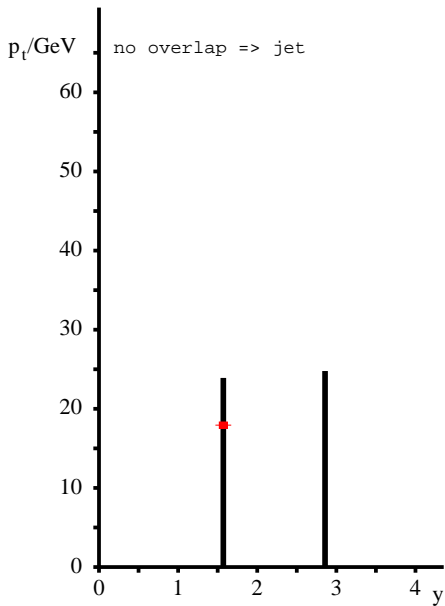
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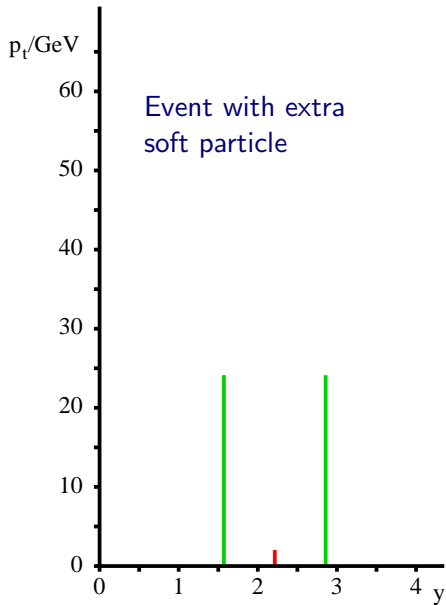
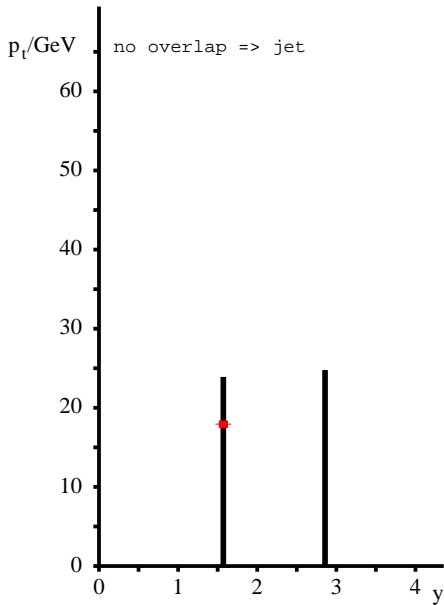
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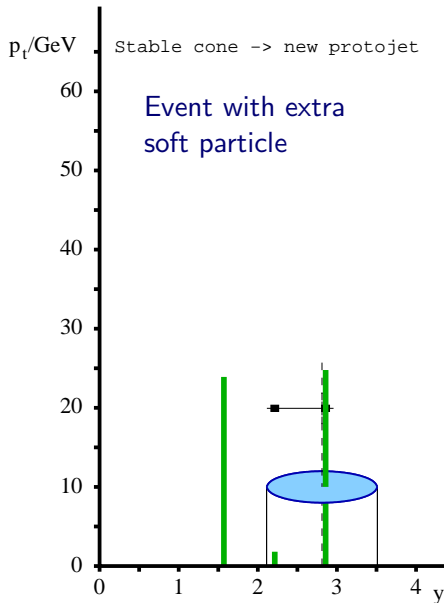
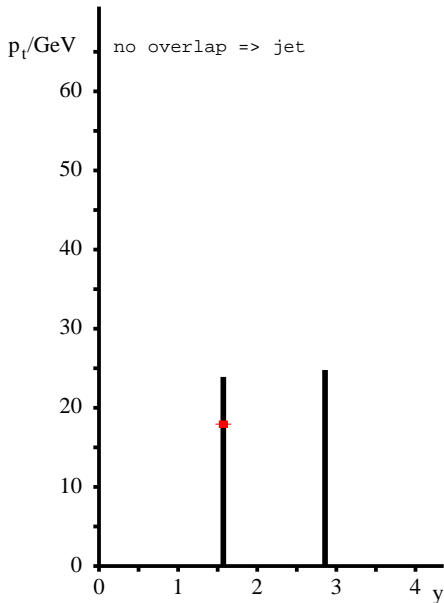


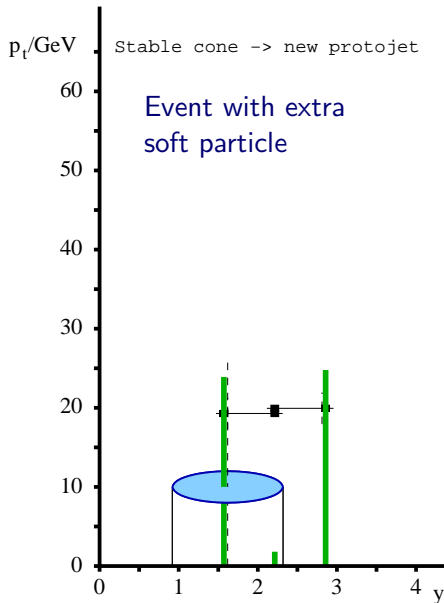
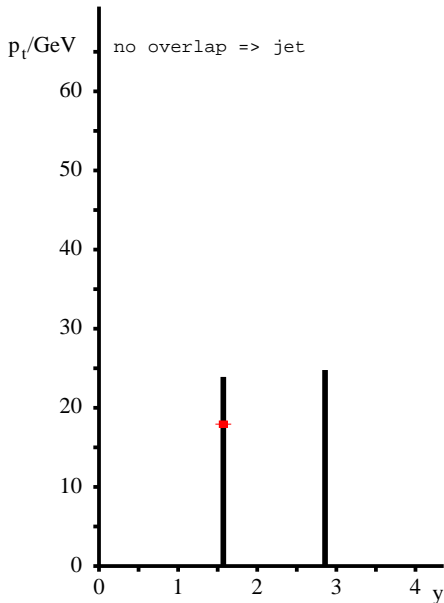


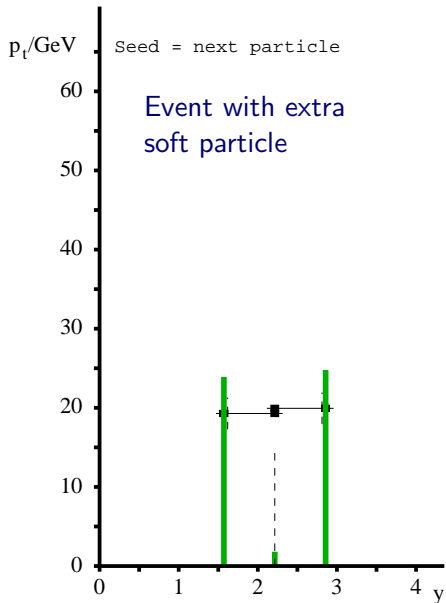
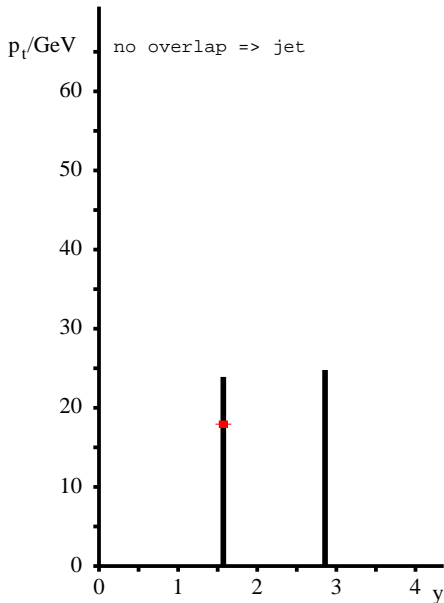


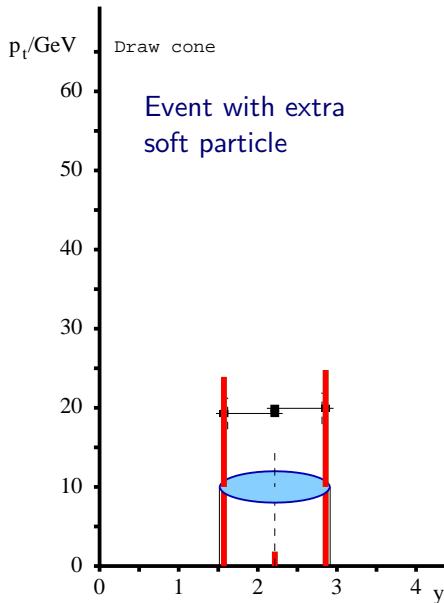
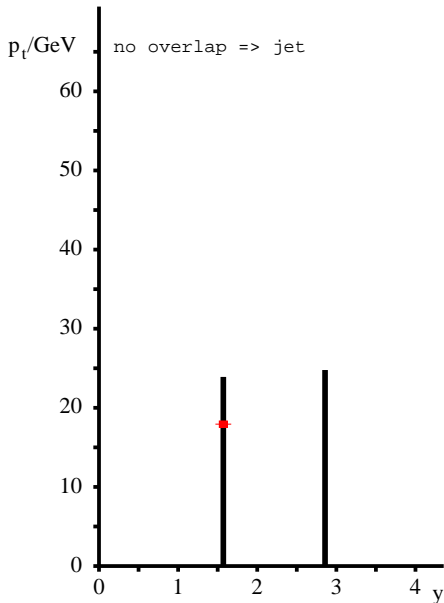


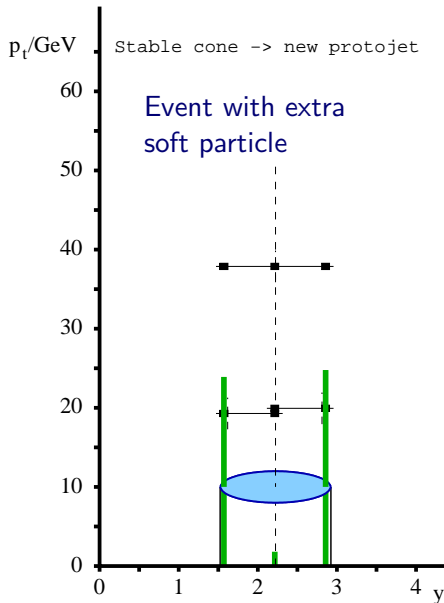
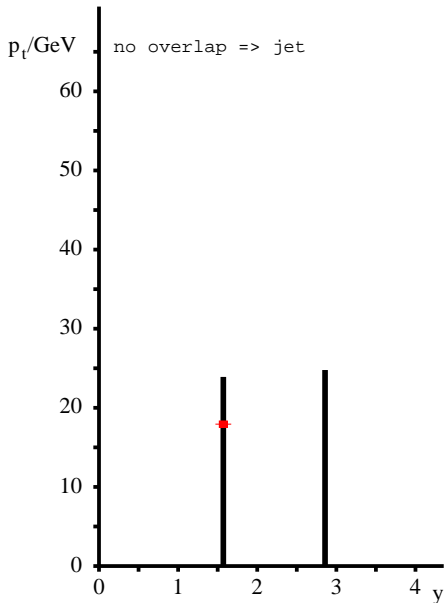


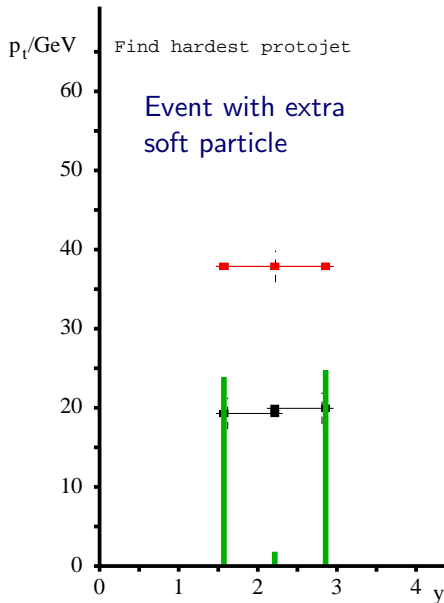
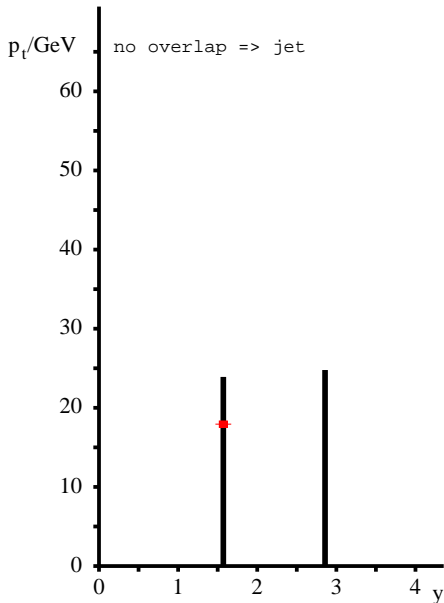


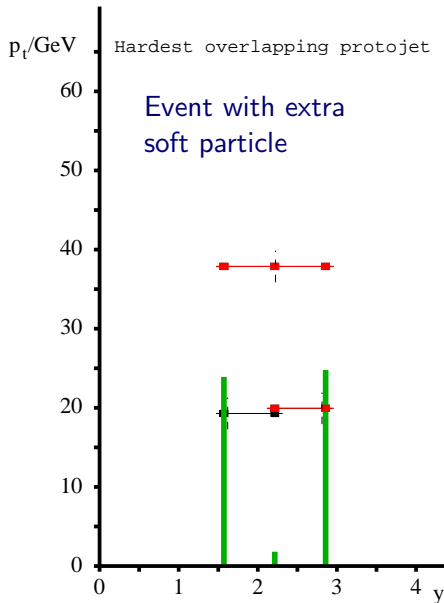
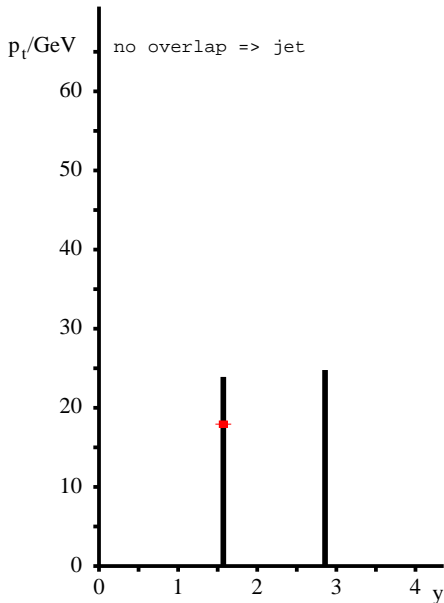


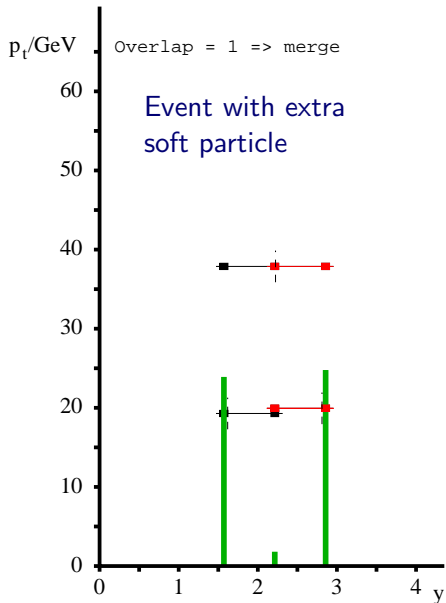
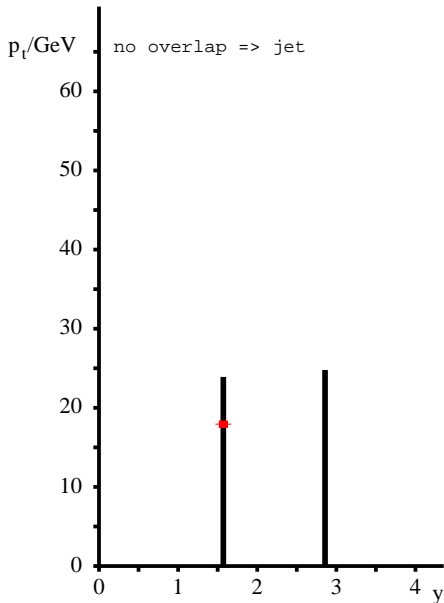


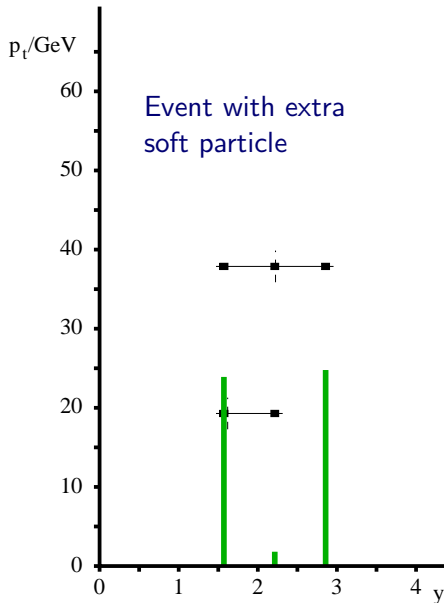
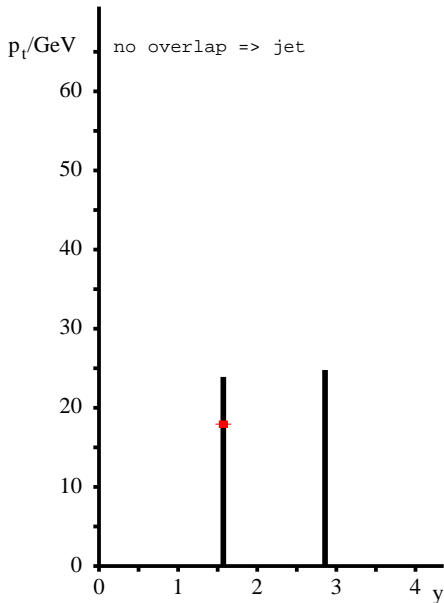


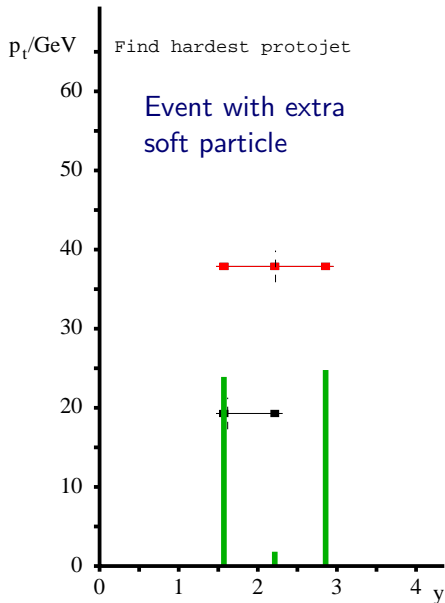
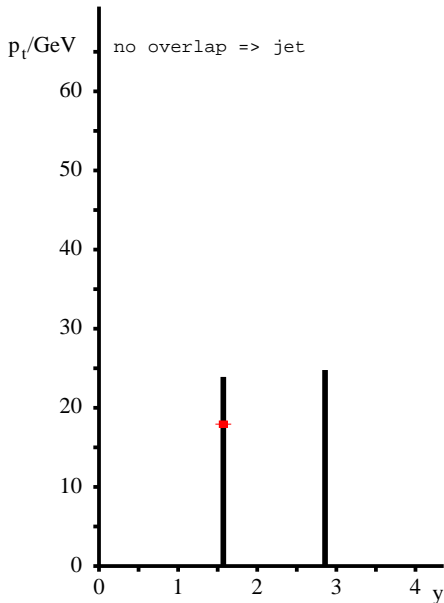


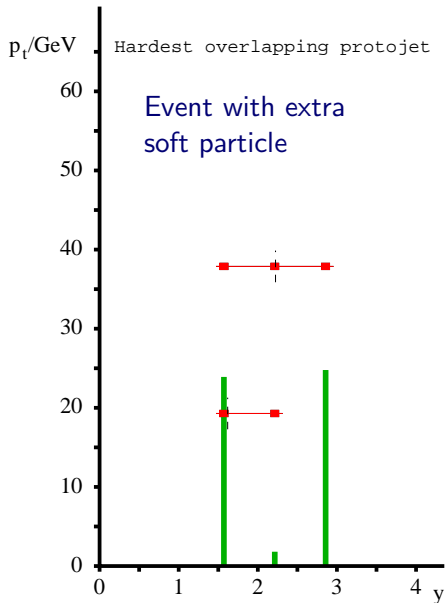
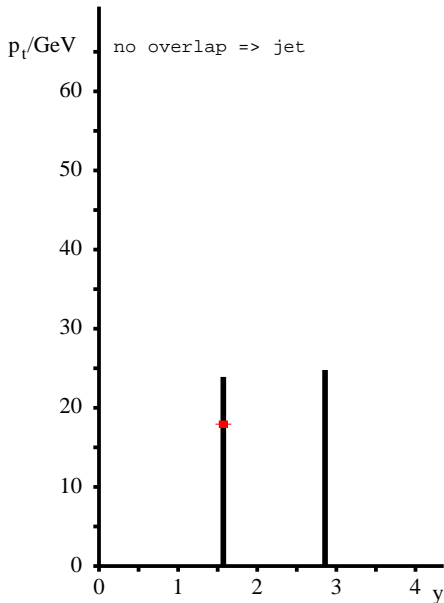


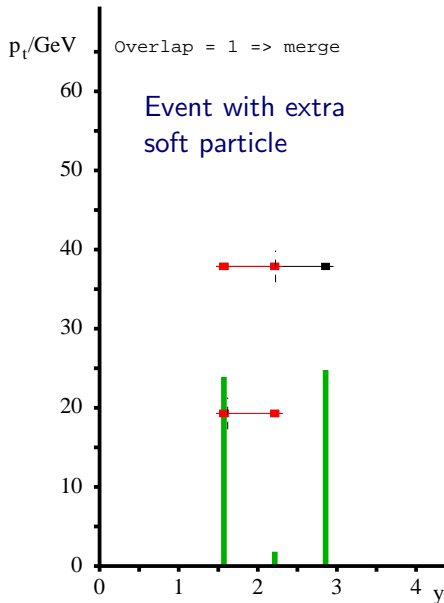
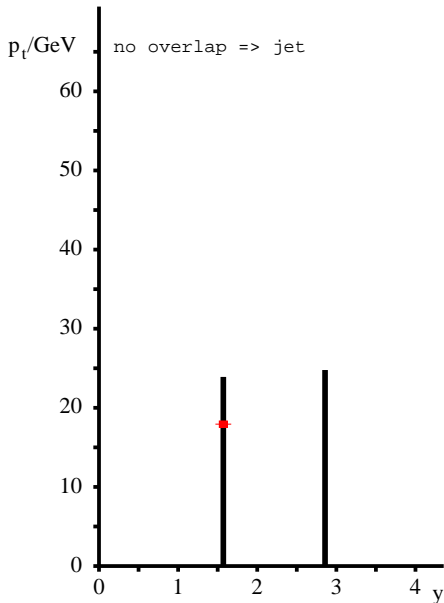


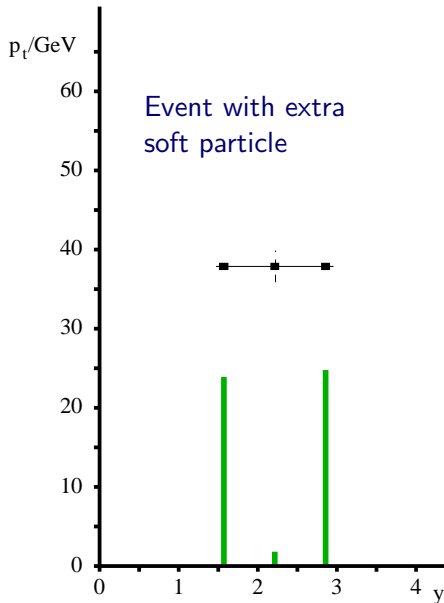
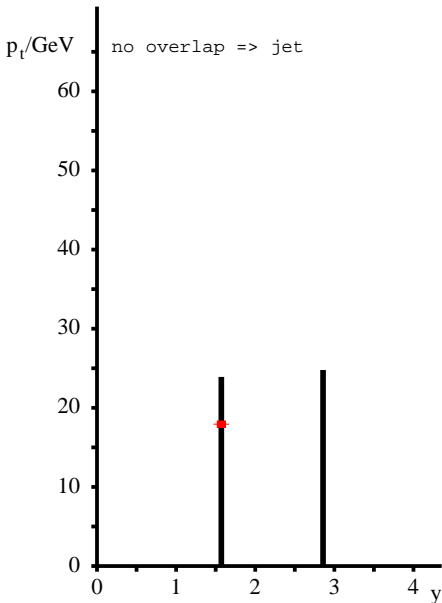


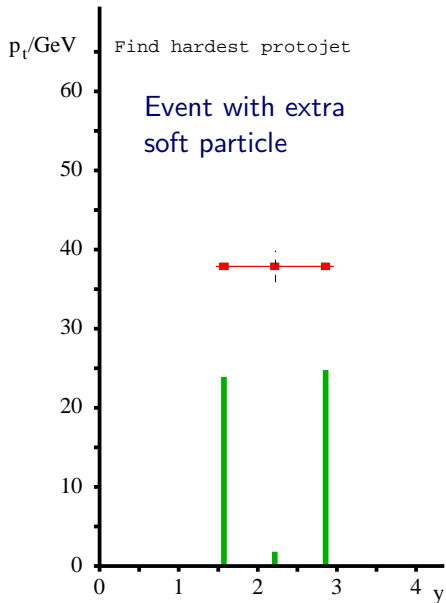
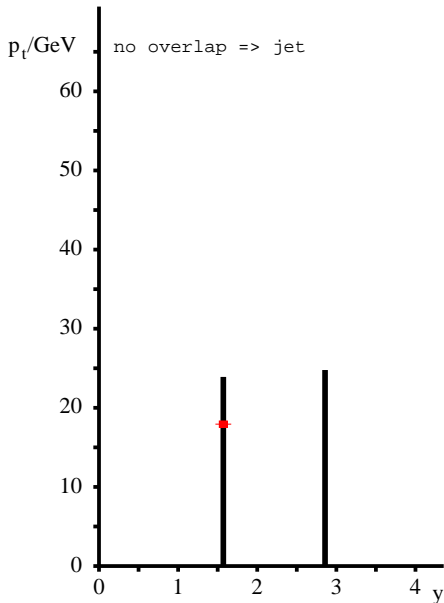


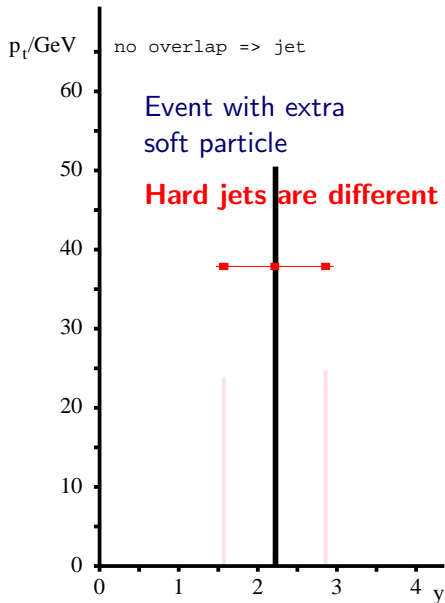
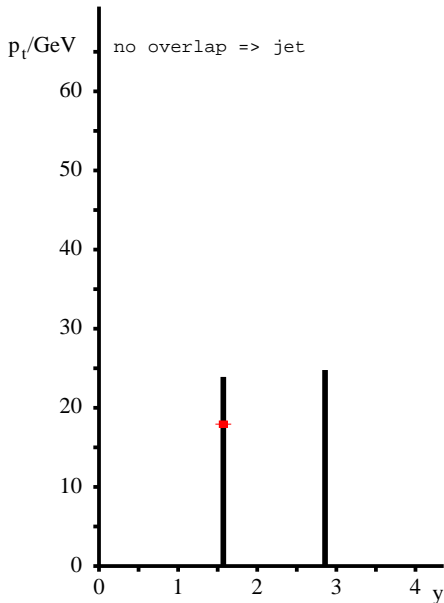






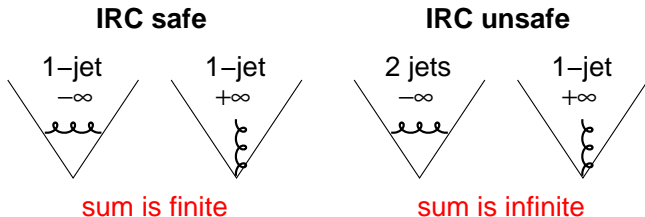




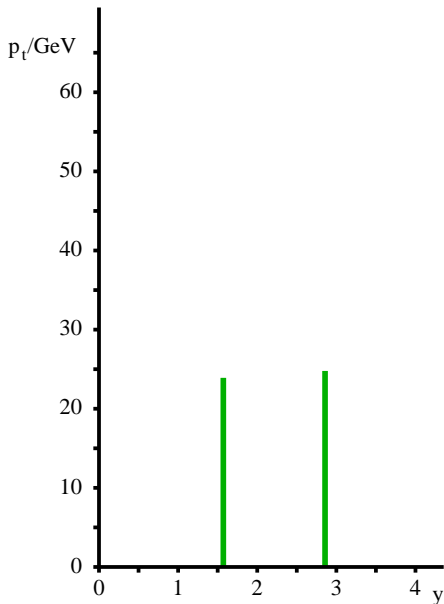


Soft emission, collinear splitting are both **infinite** in pert. QCD.

Infinites **cancel** with loop diagrams if jet-alg IRC safe



Some calculations simply become **meaningless**



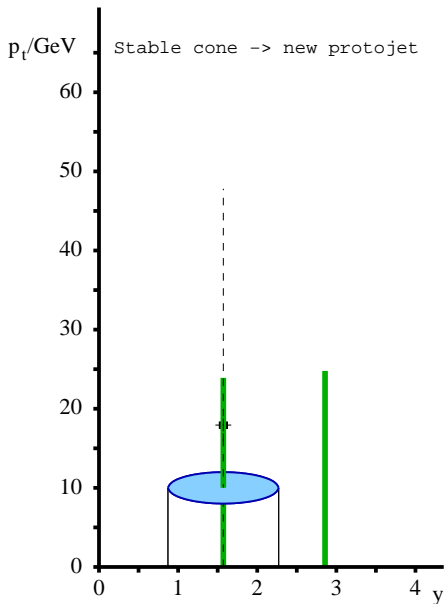
Looking for stable cones \simeq finding local minima of a potential.

Problem: set of iterative solution depends on set of starting points.

Patch: after 1st round of iteration, find midpoints between protojets, use as new seeds

CDF Midpoint algorithm
 D0 Run II algorithm

This solves problem for
 2-hard-particle configs.
 [But it persists for 3-hard]



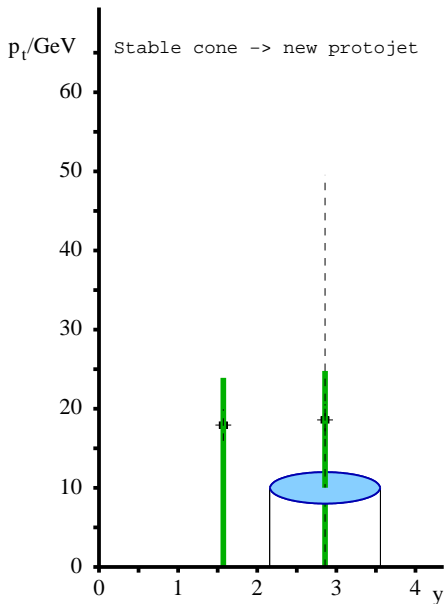
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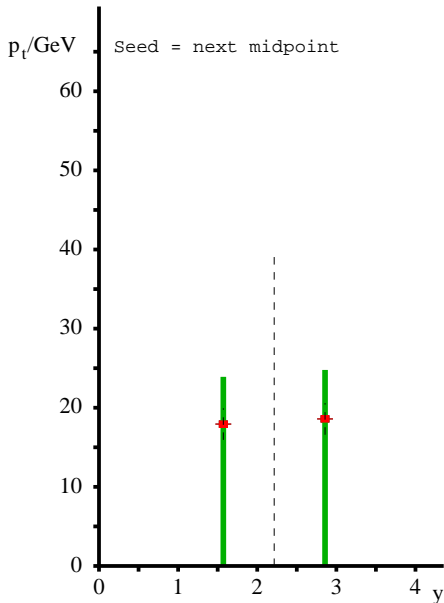
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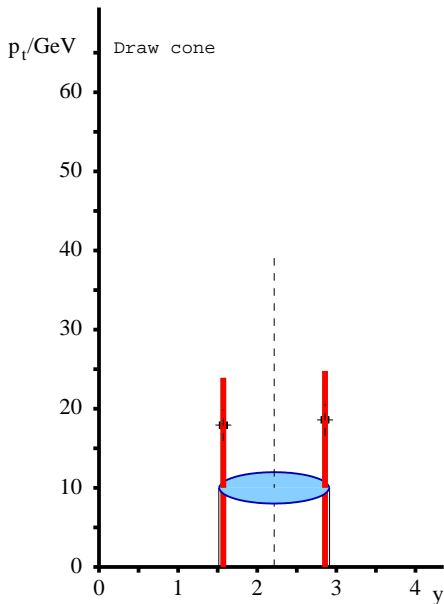
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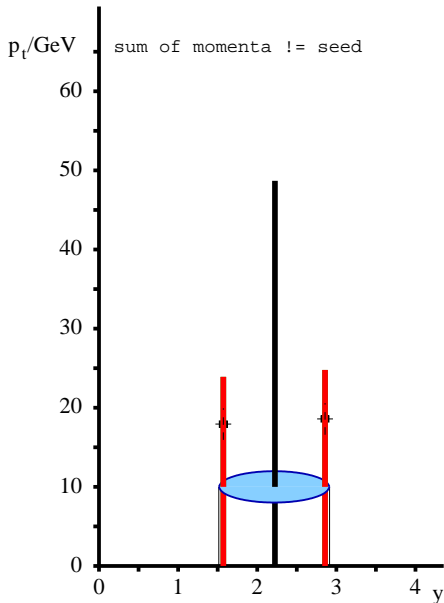
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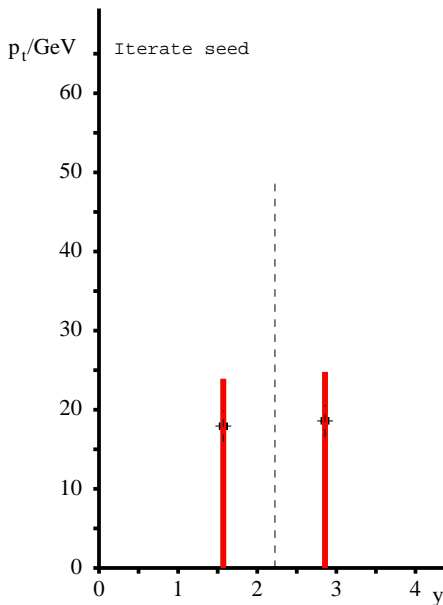
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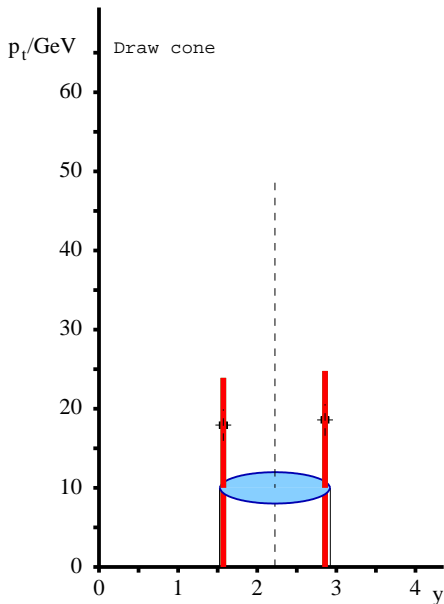
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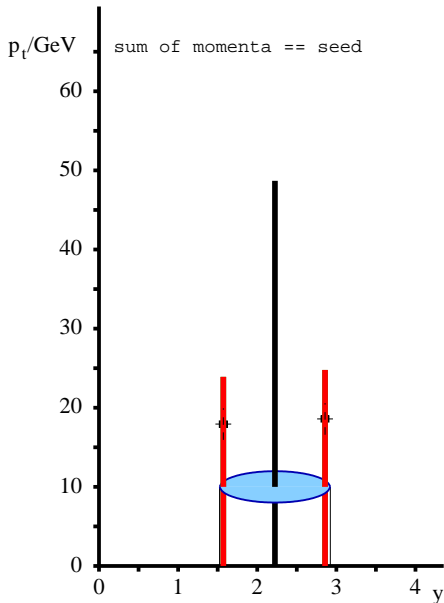
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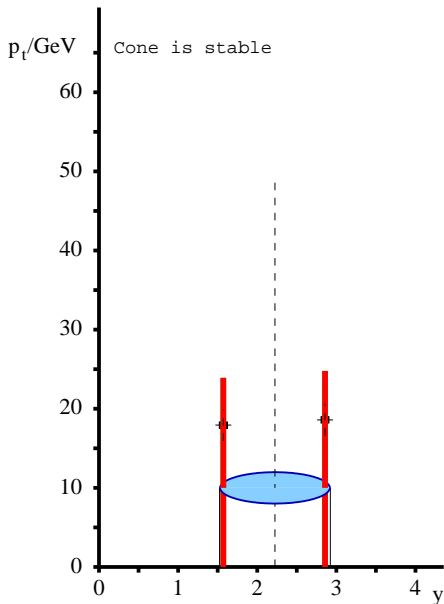
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Midpoint algorithm (IC_{mp} -SM)



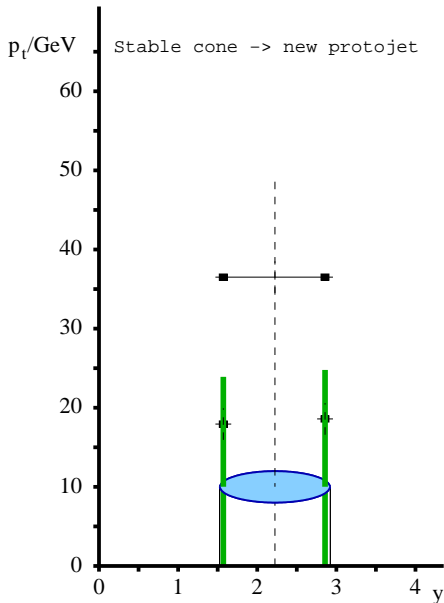
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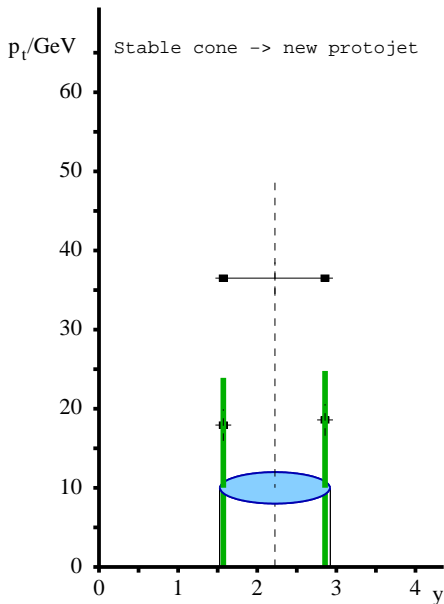
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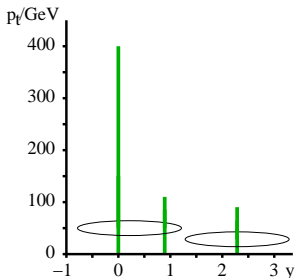
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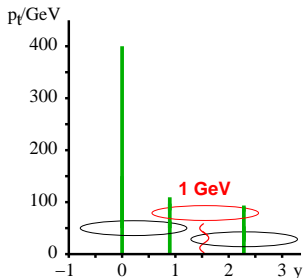


Stable cones
 with midpoint:

$\{1,2\}$ & $\{3\}$

Jets with
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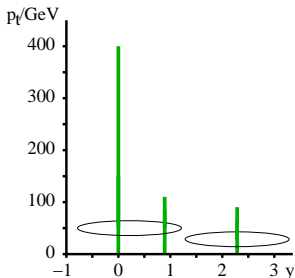
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Or collinear unsafe with seed threshold

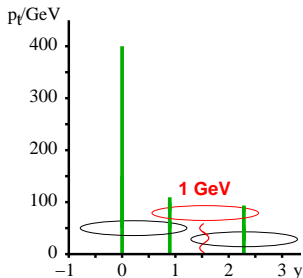


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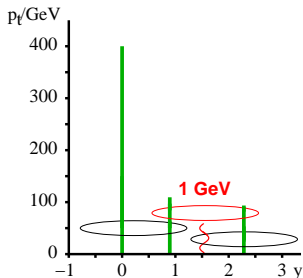
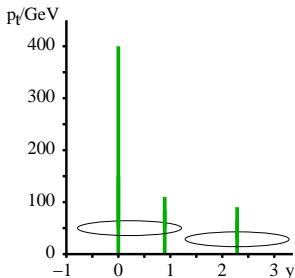
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Does IRC safety really matter?

Real life does not have infinities, but pert. infinity leaves a real-life trace

$$\alpha_s^2 + \alpha_s^3 + \alpha_s^4 \times \infty \rightarrow \alpha_s^2 + \alpha_s^3 + \alpha_s^4 \times \ln p_t/\Lambda \rightarrow \alpha_s^2 + \underbrace{\alpha_s^3 + \alpha_s^3}_{\text{BOTH WASTED}}$$

Among consequences of IR unsafety:

	<i>Last meaningful order</i>			Known at
	JetClu, ATLAS cone [IC-SM]	MidPoint [IC _{mp} -SM]	CMS it. cone [IC-PR]	
Inclusive jets	LO	NLO	NLO	NLO (→ NNLO)
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m _{jet} in 2j + X	none	none	none	LO

NB: \$30 – 50M investment in NLO

Multi-jet contexts much more sensitive: **ubiquitous at LHC**

And LHC will rely on QCD for background double-checks
 extraction of cross sections, extraction of parameters

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 - calorimeter thresholds
 - magnetic fields acting on charged particles
 - calorimeter noise
2. Detectors split/merge collinear particles
 - Two particles into single calo-tower
 - One particles showers into two calo-towers
3. High lumi adds lots of extra soft seeds

IRC safety provides resilience to these effects

1 & 3 shift energy scale, but don't change overall jet-structure

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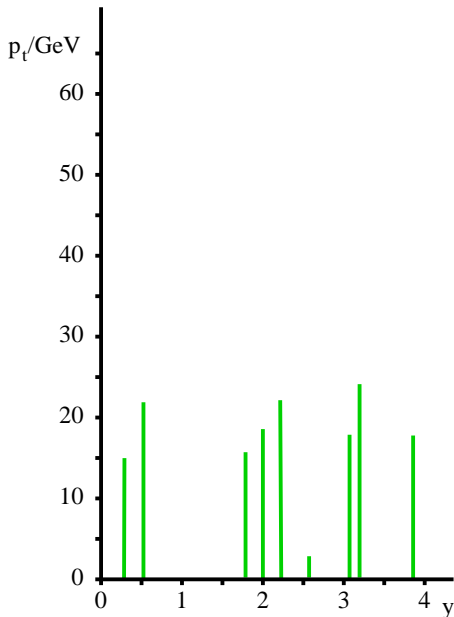
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Can we cure this IR safety
problem?

Seedless [Infrared Safe] cones (SC-SM)



Aim to identify *all* stable cones, independently of any seeds

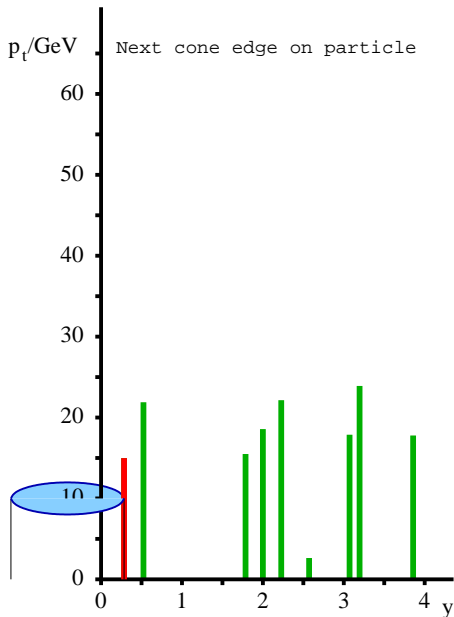
Procedure in 1 dimension (y):

- ▶ find all distinct enclosures of radius R by repeatedly sliding a cone sideways until edge touches a particle
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In 2 dimensions (y, ϕ) can design analogous procedure **SISCone**
 GPS & Soyez '07

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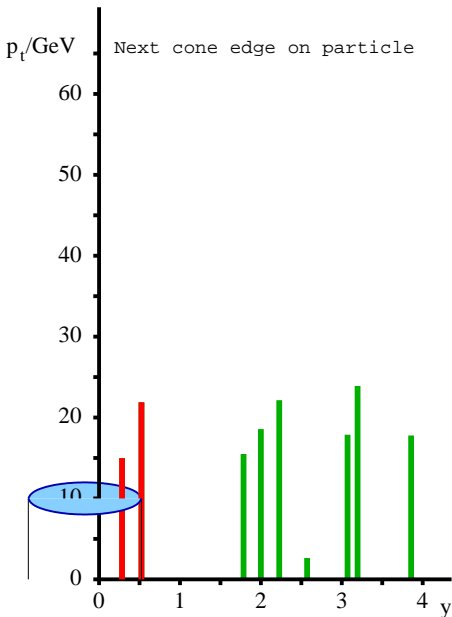
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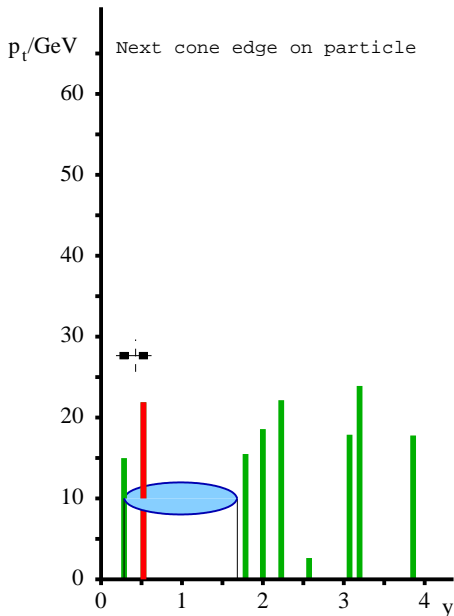
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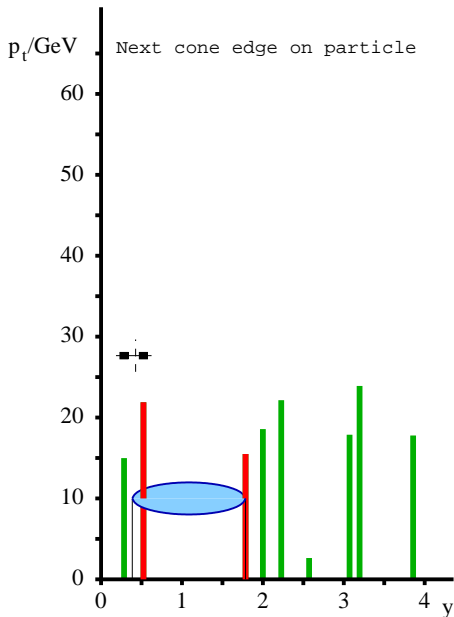
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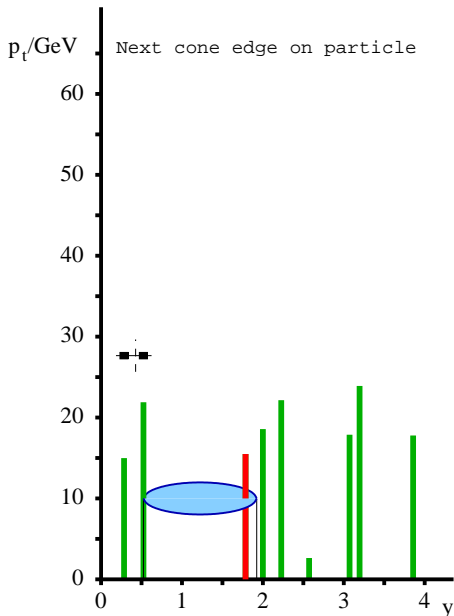
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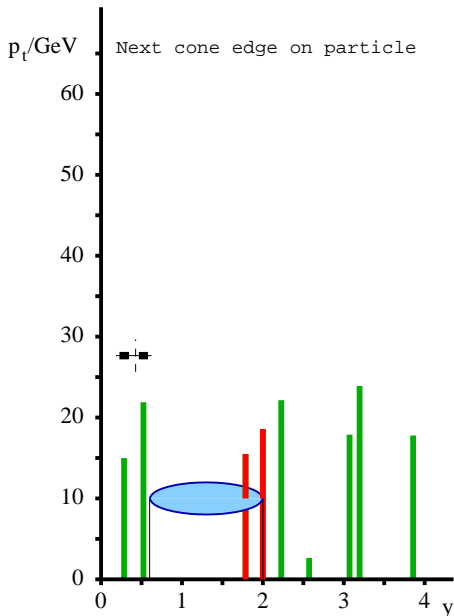
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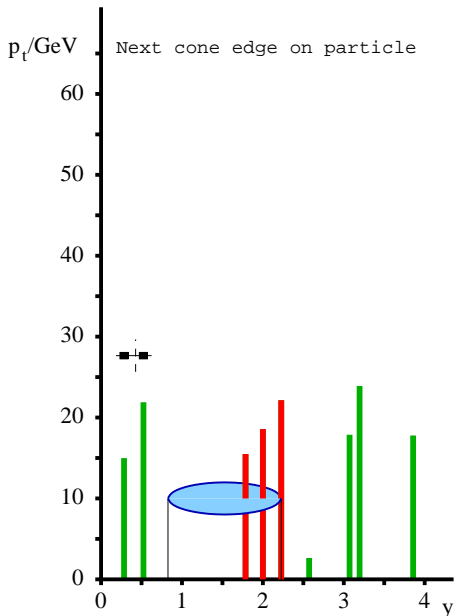
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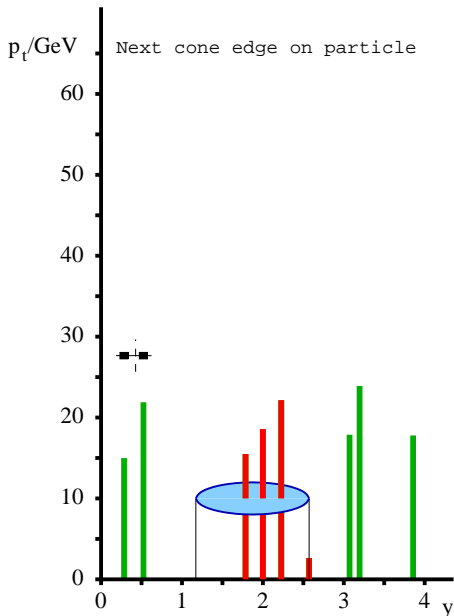
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Aim to identify *all* stable cones, independently of any seeds

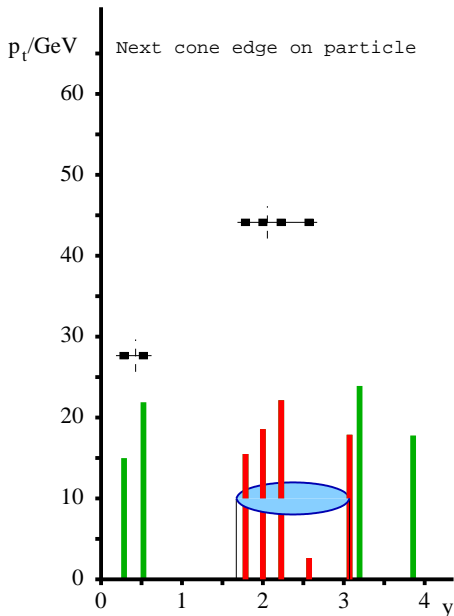
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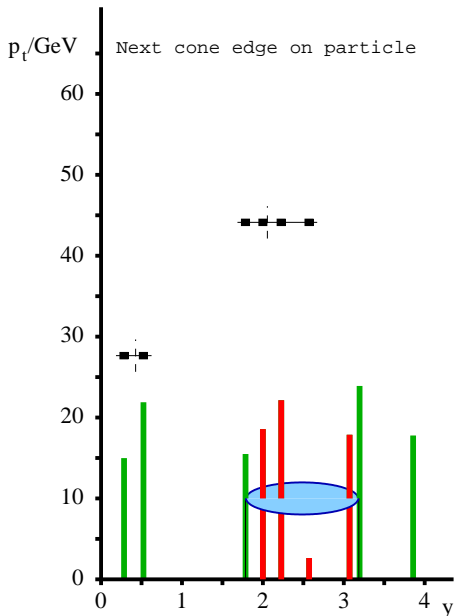
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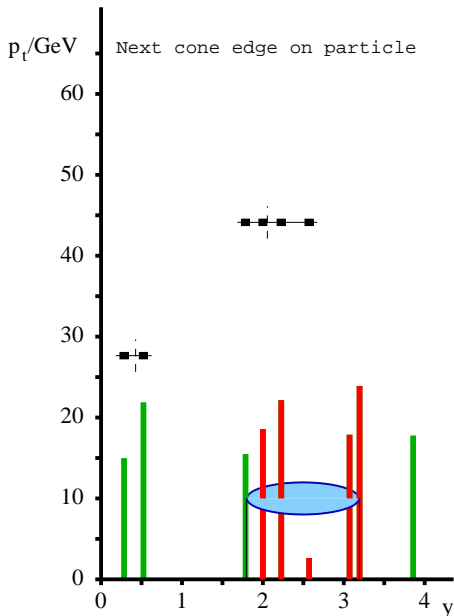
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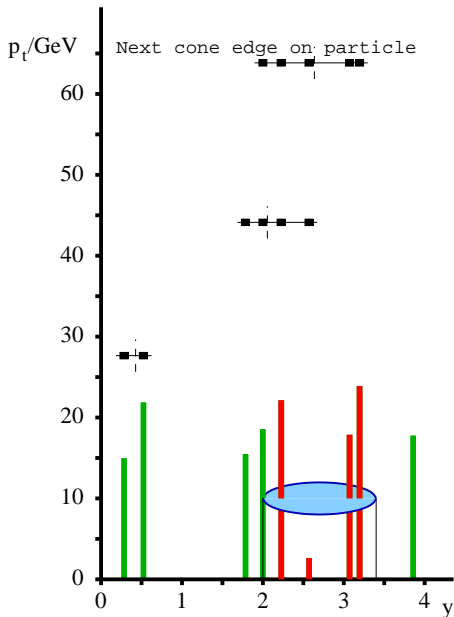
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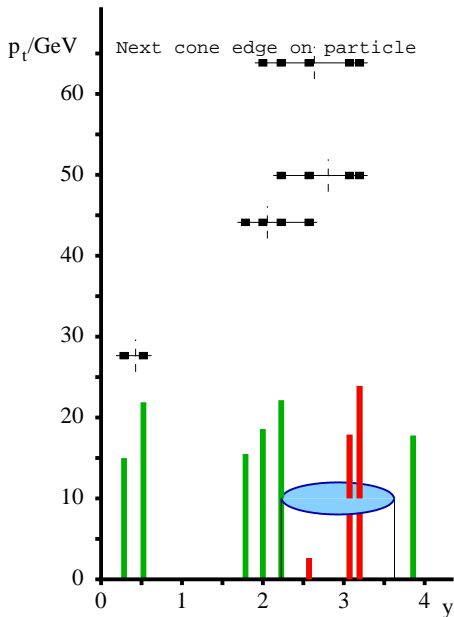
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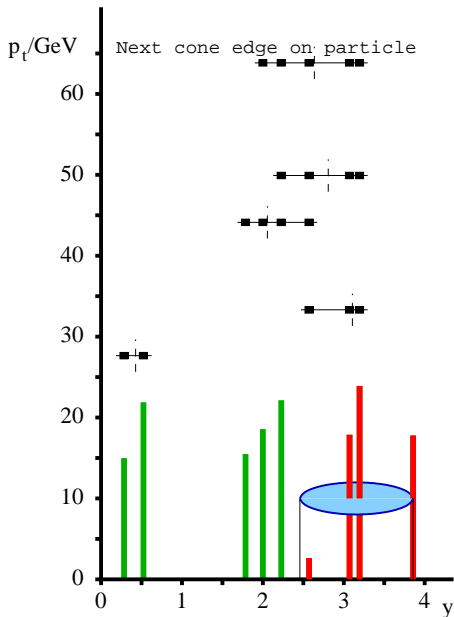
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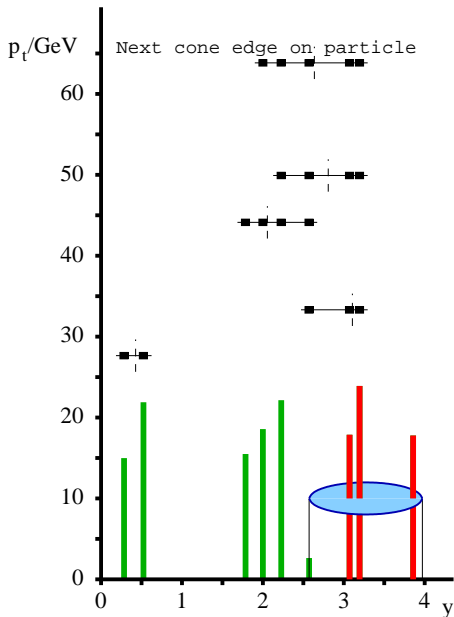
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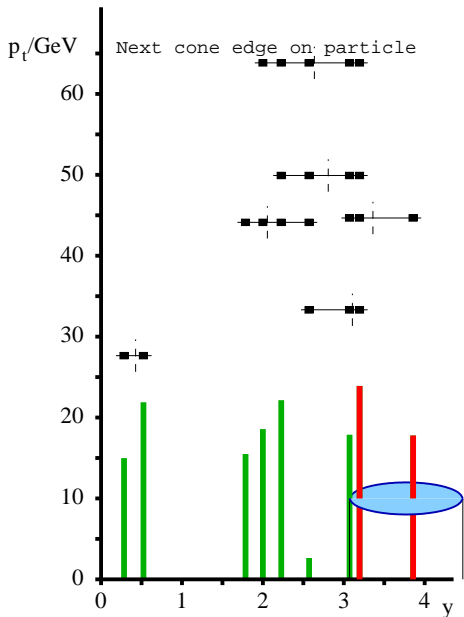
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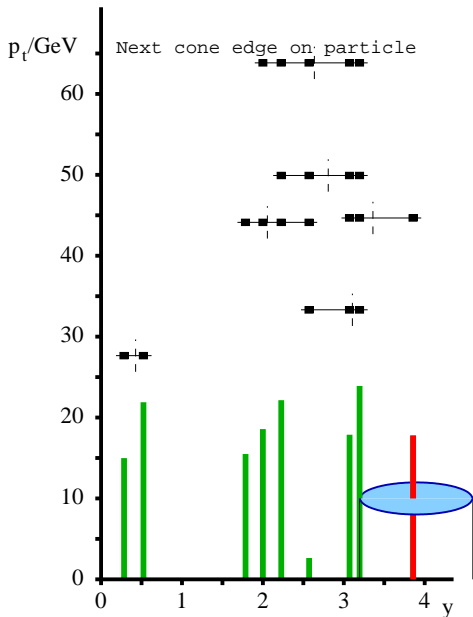
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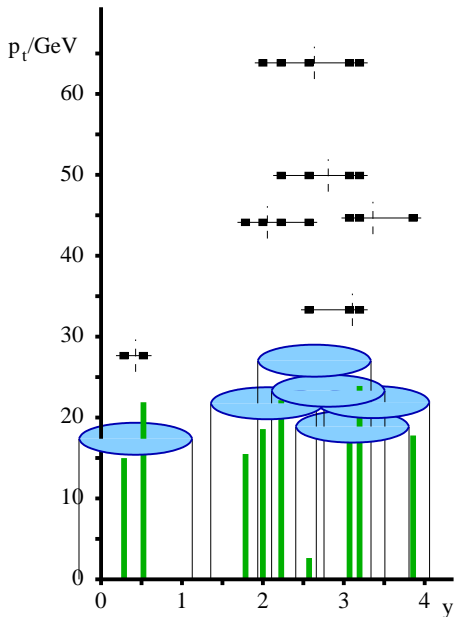
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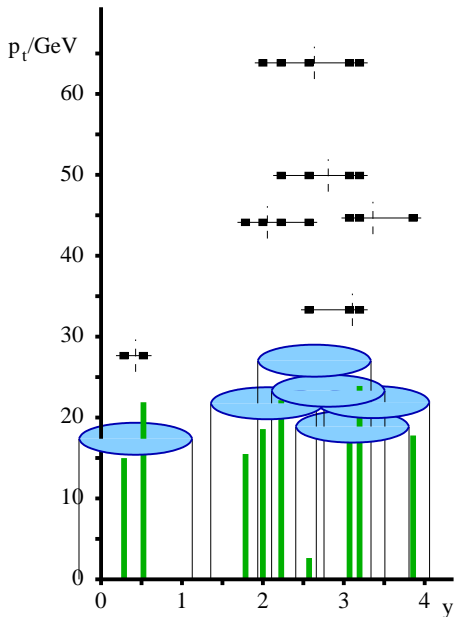
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This gives an IRC safe cone alg.

- ▶ Generate event with $2 < N < 10$ hard particles, find jets
- ▶ Add $1 < N_{soft} < 5$ soft particles, find jets again [repeatedly]
- ▶ If the jets are different, algorithm is IR unsafe.

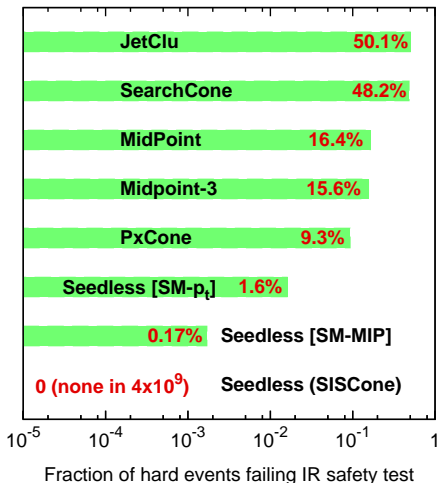
Unsafety level	failure rate
2 hard + 1 soft	~ 50%
3 hard + 1 soft	~ 15%
SISCone	IR safe !

Be careful with split-merge too

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How much does IR safety *really* matter?

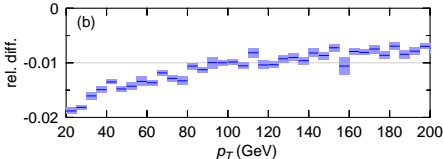
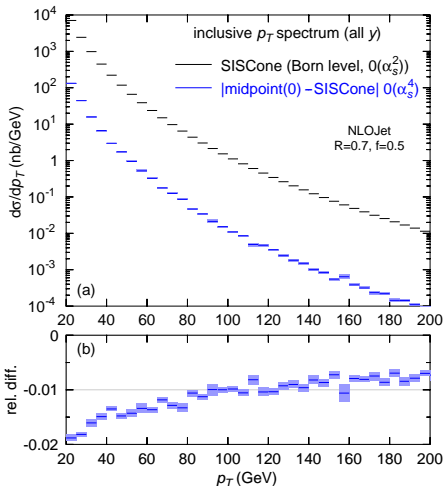
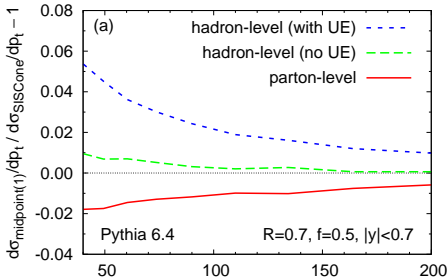
Compare midpoint and SISCone

Result depends on observable:

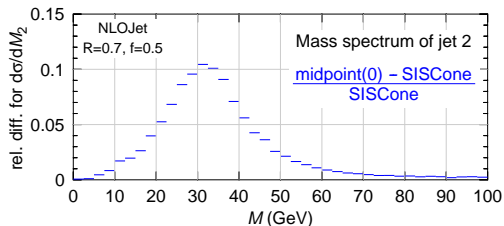
- ▶ inclusive jet spectrum is the least sensitive (affected at NNLO)
- ▶ larger differences (5 – 10%) at hadron level

seedless reduces UE effect

$p\bar{p}$ $\sqrt{s} = 1.96$ TeV



Look at jet masses in multijet events. **NB: Jet masses reconstruct boosted $W/Z/H/top$ in BSM searches**



Select 3-jet events

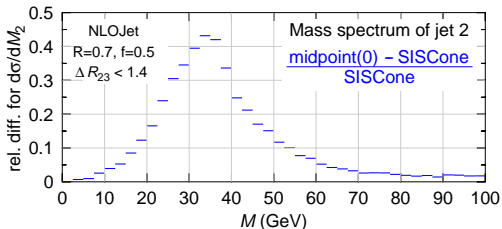
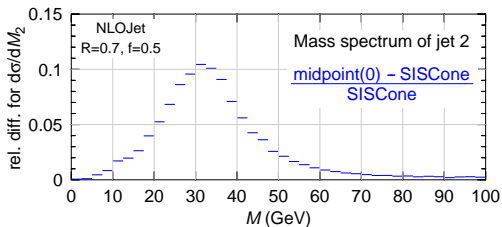
$$p_{t1,2,3} > \{120, 60, 20\} \text{ GeV,}$$

Calculate LO jet-mass spectrum for jet 2, compare midpoint with SISCone.

- ▶ 10% differences by default
- ▶ **40% differences** with extra cut $\Delta R_{2,3} < 1.4$
 e.g. for jets from common decay chain

In complex events, IR safety matters

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In complex events, IR safety matters

- ▶ IR safety often matters less in *inclusive* quantities
- ▶ It matters more in multi-jet cases

- ▶ ATLAS cone, JetClu (IC-SM) are *very bad*
- ▶ CMS cone (IC-PR), Midpoint (IC_{mp}-SM) *moderately bad*

- ▶ An IRC safe cone algorithm exists (SISCone)
- ▶ **Avoid trouble later: use IR-safe algs from the start**

cf. CDF W+jets

What jet definition should I use?

[jet def. \equiv jet alg., R , (f)]

Generalise inclusive-type sequential recombination with

$$d_{ij} = \min(k_{ti}^{2p}, k_{tj}^{2p}) \Delta R_{ij}^2 / R^2 \quad d_{iB} = k_{ti}^{2p}$$

	Alg. name	Comment	time
$p = 1$	k_t CDOSTW '91-93; ES '93	Hierarchical in rel. k_t	$N \ln N$ exp.
$p = 0$	Cambridge/Aachen Dok, Leder, Moretti, Webber '97 Wengler, Wobisch '98	Hierarchical in angle Scan multiple R at once \leftrightarrow QCD angular ordering	$N \ln N$
$p = -1$	anti- k_t Cacciari, GPS, Soyez '08 \sim reverse- k_t Delsart	Hierarchy meaningless, jets like CMS cone (IC-PR)	$N^{3/2}$
SC-SM	SISCone GPS Soyez '07 + Tevatron run II '00	Replaces JetClu, ATLAS MidPoint (xC-SM) cones	$N^2 \ln N$ exp.

*Compromise between having a limited set of algs.
and a good range of complementary properties*

COMMERCIAL BREAK

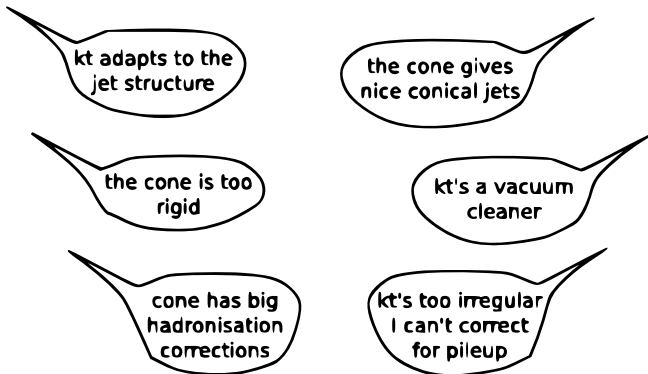
One place to stop for all your jet-finding needs:

FASTJET

`http://www.lpthe.jussieu.fr/~salam/fastjet`
Cacciari, GPS & Soyez '05–07

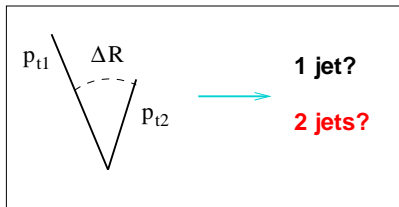
- ▶ Fast, native, computational-geometry methods for k_t , Cam/Aachen
Cacciari & GPS '05-06
- ▶ Plugins for SIScone (plus some other, deprecated cones)
- ▶ Many other features too, e.g. jet areas

Jet discussions: often polarised, driven by unquantified statements

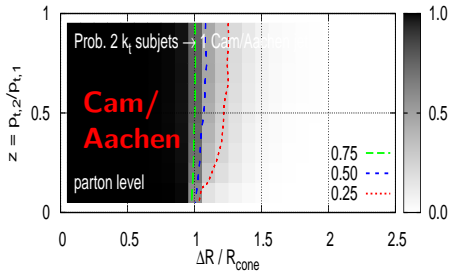
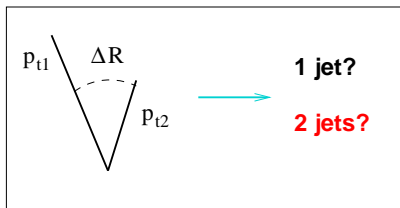


- ▶ Rigorous approach is to quantify similarities & differences
- ▶ Bottom line: grains of truth in the qualitative statements

So want good cone algorithms too [NB: recall, two variants xC-SM & xC-PR]

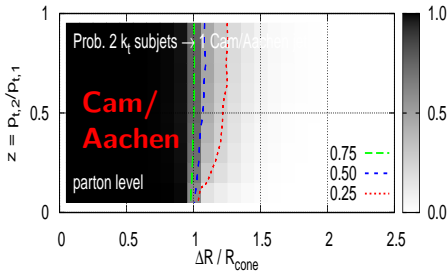
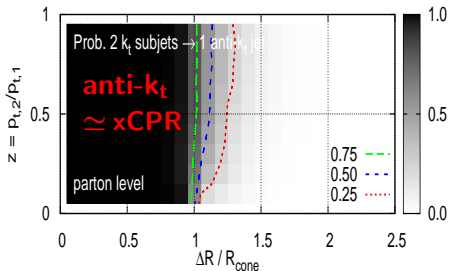
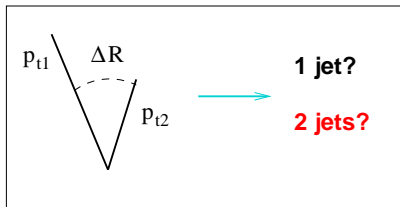


SISCone (xC-SM) reaches further for hard radiation than other algs



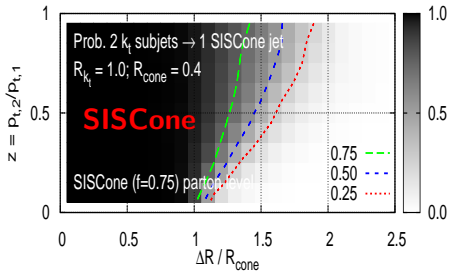
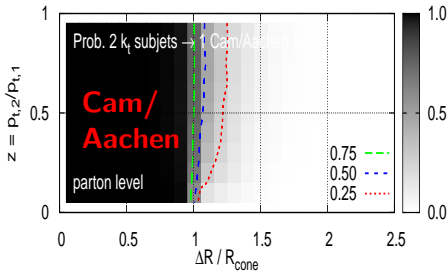
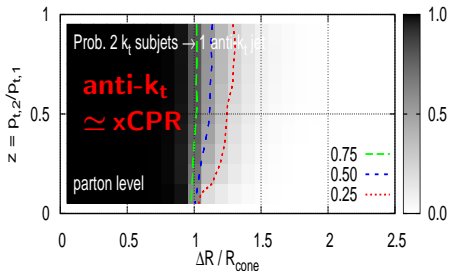
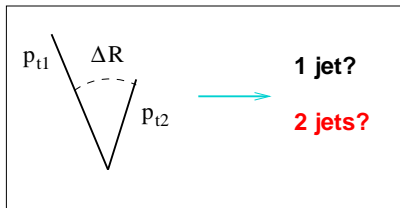
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the *reach* of jet algorithms



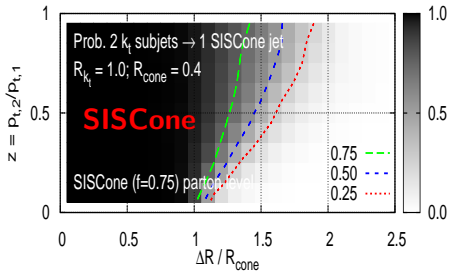
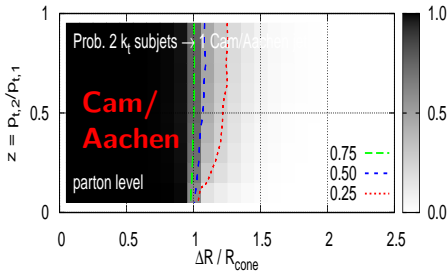
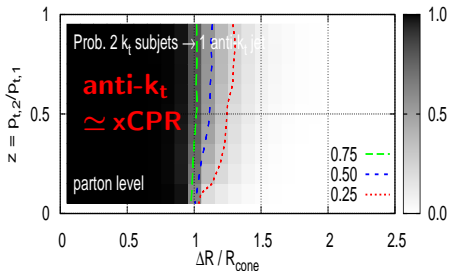
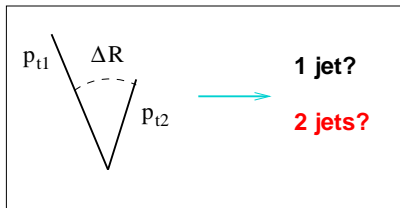
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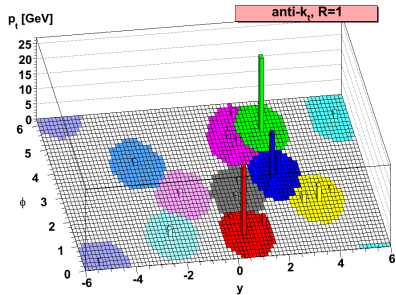
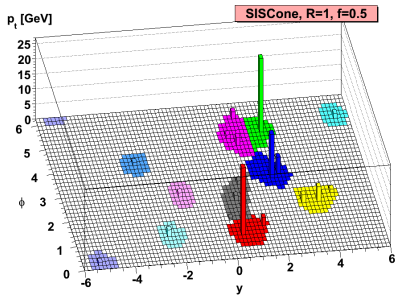
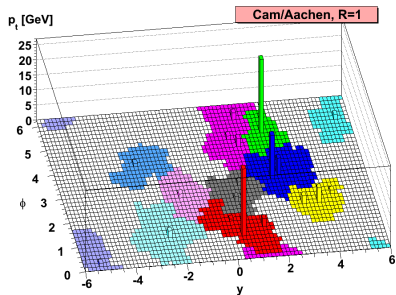
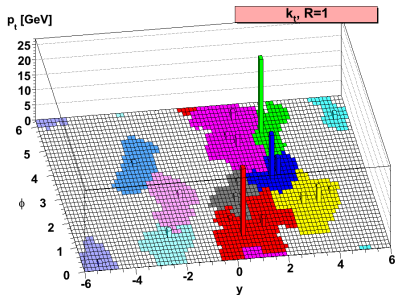


SISCone (xC-SM) reaches further for hard radiation than other algs

the reach of jet algorithms



SISCone (xC-SM) reaches further for hard radiation than other algs



To first approx:
various algs. moderately different;
but **R** can matter a lot more

4-way tension in many measurements:

Prefer small R	prefer large R
resolve many jets (e.g. $t\bar{t}$)	minimize QCD radiation loss
limit UE & pileup	limit hadronisation

Parton $p_t \rightarrow$ jet p_t

Ill-defined: MC “parton”

PT radiation:

$$q : \quad \Delta p_t \simeq \frac{\alpha_s C_F}{\pi} p_t \ln R$$

$$g : \quad \Delta p_t \simeq \frac{\alpha_s C_A}{\pi} p_t \ln R$$

Hadronisation:

$$q : \quad \Delta p_t \simeq \frac{C_F}{R} \cdot 0.4 \text{ GeV}$$

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Underlying event:

$$q, g : \quad \Delta p_t \simeq \frac{R^2}{2} \cdot 2.5\text{--}15 \text{ GeV}$$

crude analytical estimates
cf. Dasgupta, Magnea & GPS '07

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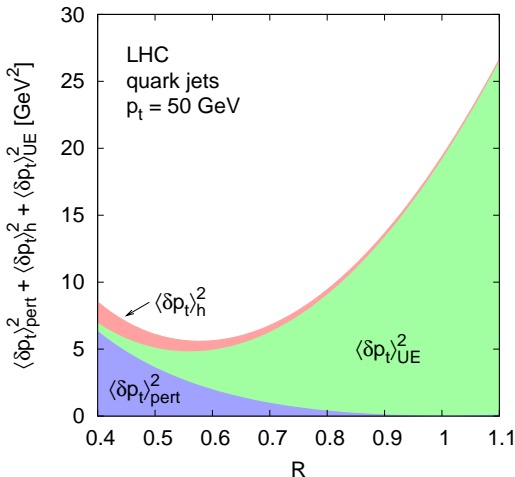
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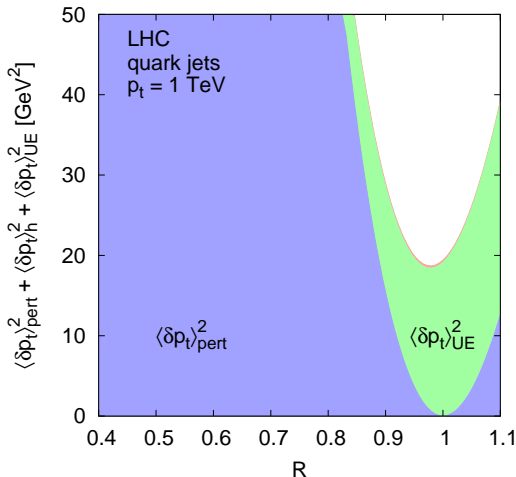
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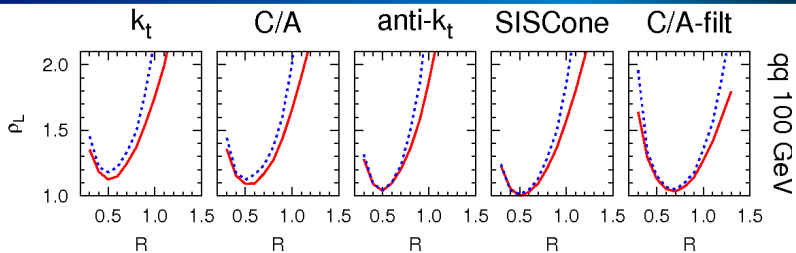
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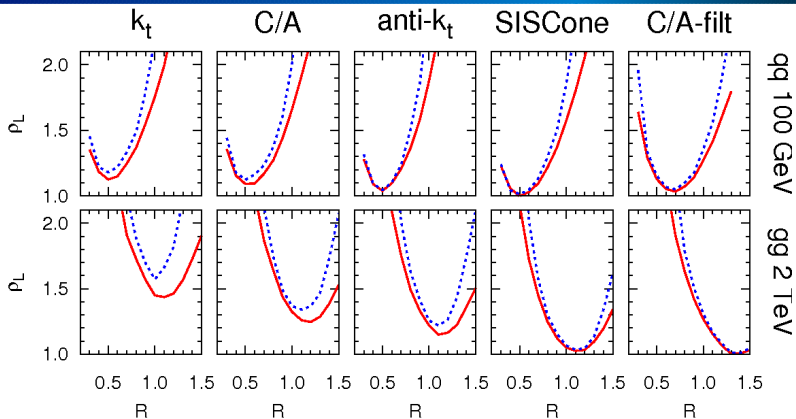
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Relative peak quality (lumi ratios ρ_L), LHC



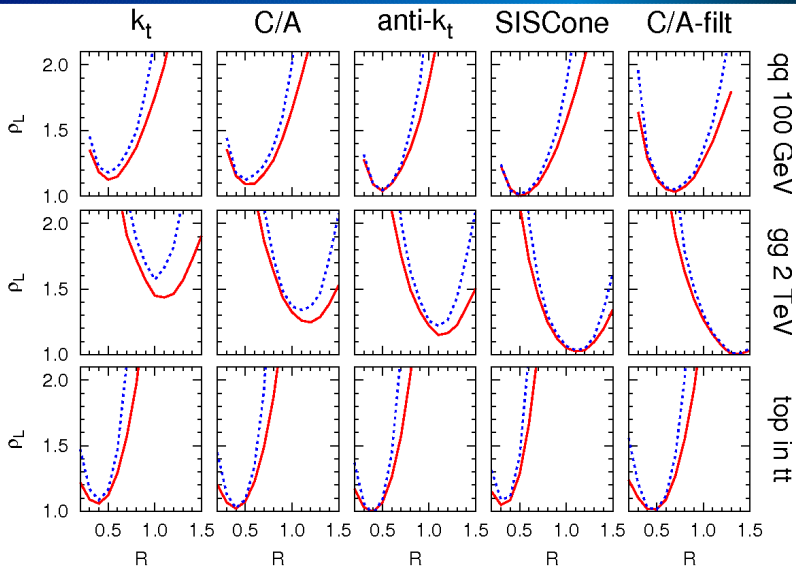
PRELIMINARY

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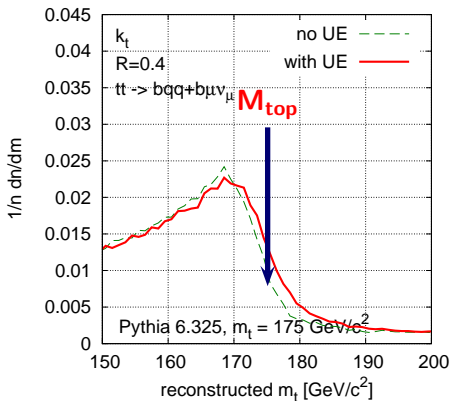


PRELIMINARY

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PRELIMINARY



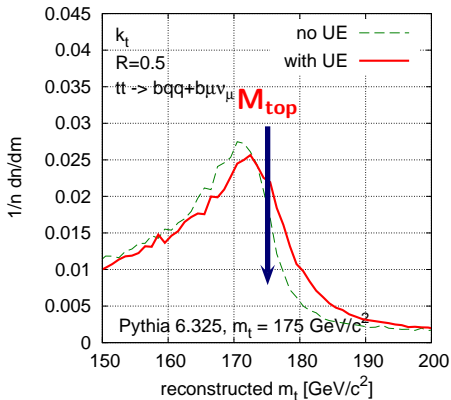
Game: measure top mass to 1 GeV
 example for Tevatron
 $m_t = 175 \text{ GeV}$

► Small R : lose 6 GeV to PT radiation and hadronisation, UE and pileup irrelevant

► Large R : hadronisation and PT radiation leave mass at $\sim 175 \text{ GeV}$, UE adds 2 – 4 GeV.

Is the final top mass (after W jet-energy-scale and Monte Carlo unfolding) independent of R used to measure jets?

Powerful cross-check of systematic effects
 cf. Seymour & Tevlin '06

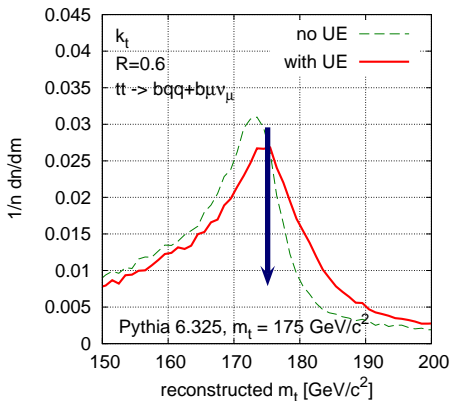


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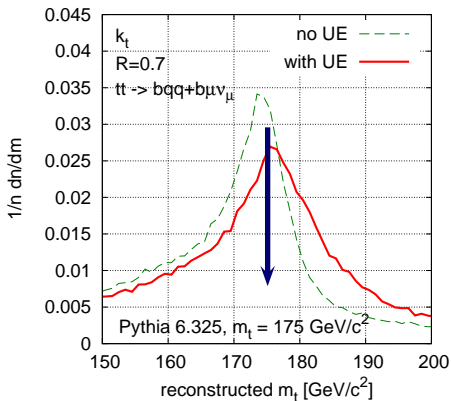


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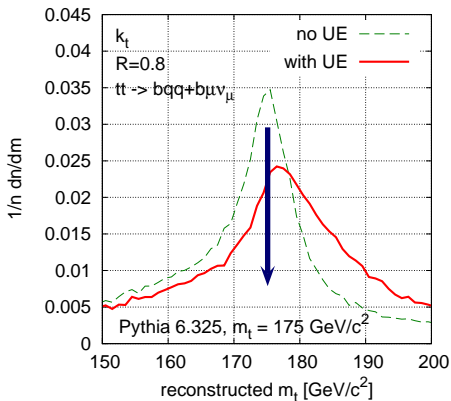


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- ▶ Small R : lose 6 GeV to PT radiation and hadronisation, UE and pileup irrelevant
- ▶ Large R : hadronisation and PT radiation leave mass at $\sim 175 \text{ GeV}$, UE adds 2 – 4 GeV.

Is the final top mass (after W jet-energy-scale and Monte Carlo unfolding) independent of R used to measure jets?

Powerful cross-check of systematic effects
cf. Seymour & Tevlin '06

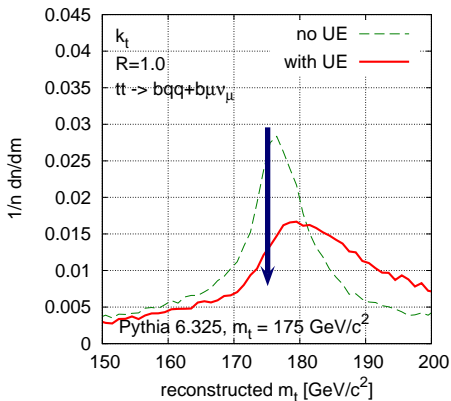


Game: measure top mass to 1 GeV
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Jets without hard partons:

Most jet algorithms give you $\sim 50 - 100$
“jets,” mostly not hard.

provide window on UE and min-bias

Making use of *all* jets

iev 0 (irepeat 24): number of particles = 1428

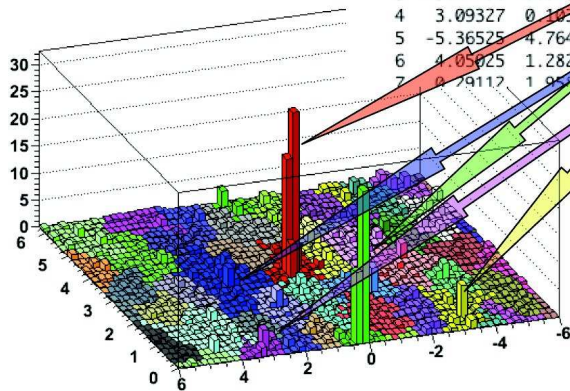
strategy used = NlnN

number of particles = 9051

Total area: 76.0265

Expected area: 76.0265

ijet	eta	phi	Pt	area +- err
0	0.15050	3.24498	69.970	2.625 +- 0.020
1	0.18579	0.13150	59.133	1.896 +- 0.020
2	2.33840	3.23960	31.976	4.749 +- 0.028
3	-3.41796	0.52394	26.585	3.084 +- 0.021
4	3.09327	0.10350	21.072	2.688 +- 0.023
5	-5.36525	4.76491	19.593	2.780 +- 0.012
6	4.05025	1.28279	15.861	3.592 +- 0.028
7	0.29117	1.95745	11.566	2.114 +- 0.018



Approximate linear relation
between Pt and area for
minimum bias jets.

Can be used on an event-by-
event basis to correct the hard
jets

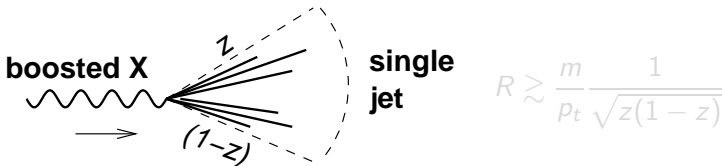
Pushing jets to their limit:

when a W , Z , H or a top \rightarrow a single jet

Not unusual at LHC: $m_W, m_t \ll 14$ TeV

Illustrate LHC challenges with a recently widely discussed class of problems:

Can you identify hadronically decaying EW bosons when they're produced at high p_t ?



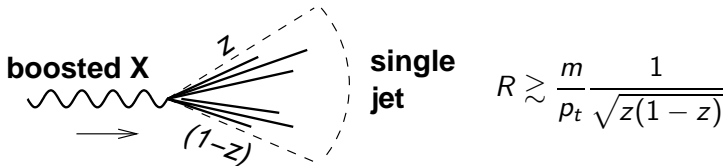
Significant discussion over years: **heavy new things decay to EW states**

- ▶ Seymour '94 [Higgs \rightarrow $WW \rightarrow \nu\ell$ jets]
- ▶ Butterworth, Cox & Forshaw '02 [$WW \rightarrow WW \rightarrow \nu\ell$ jets]
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ETC.

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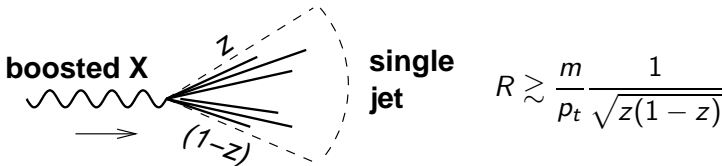
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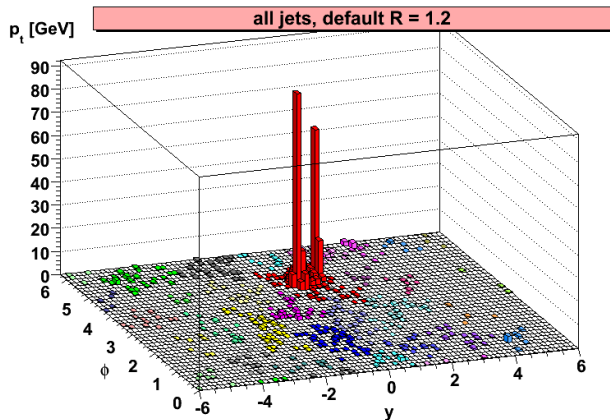
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ETC.

Jets (p. 71)
└ jet \neq a parton
└ 1 jet \gtrsim 2 partons

$pp \rightarrow ZH \rightarrow \nu\bar{\nu}b\bar{b}$, @14 TeV, $m_H = 115$ GeV



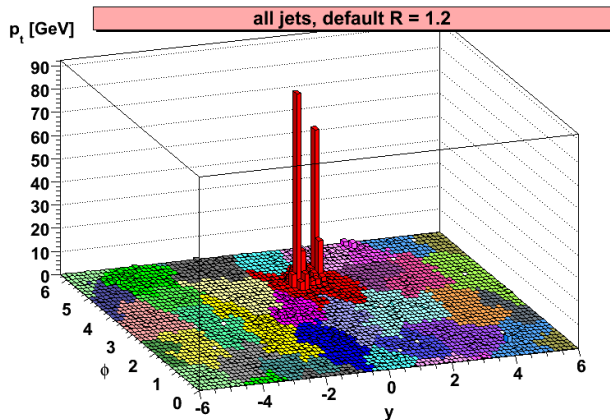
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Butterworth, Davison, Rubin & GPS '08

Possible new (light) Higgs discovery channel

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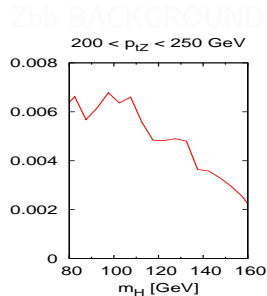
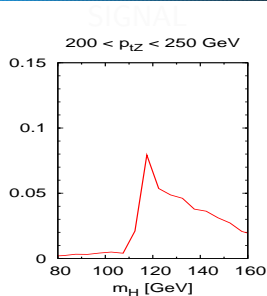
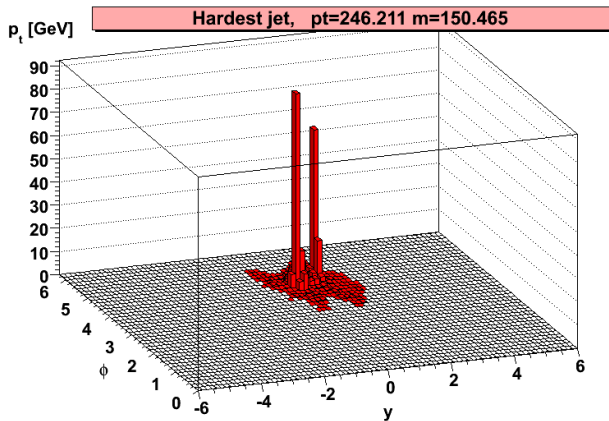
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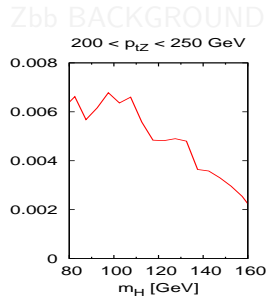
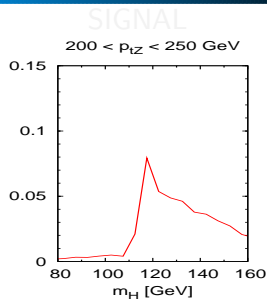
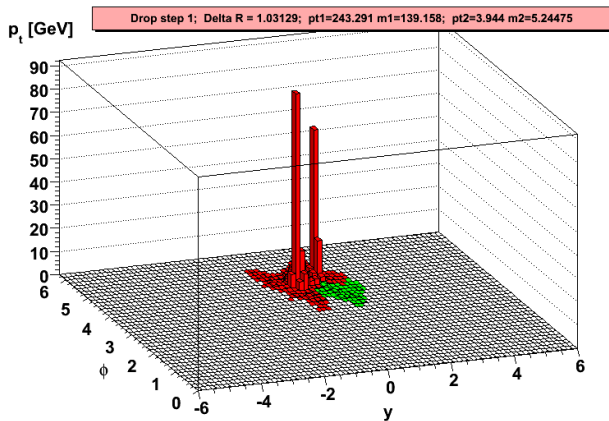
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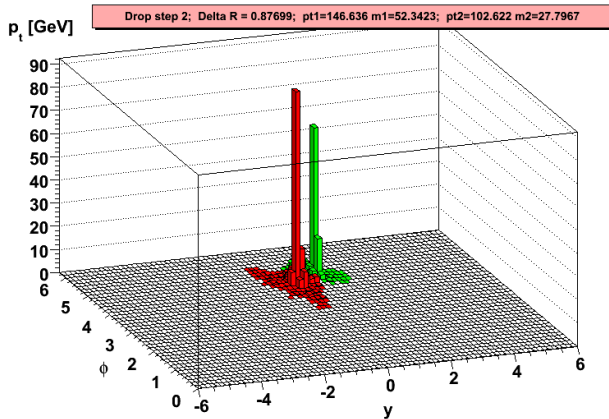
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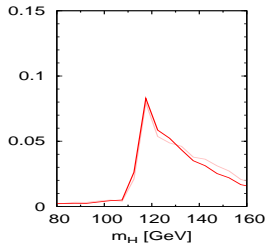
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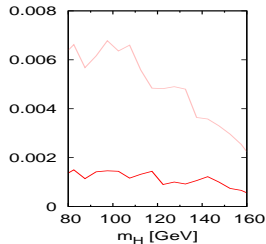
SIGNAL

$200 < p_{tZ} < 250 \text{ GeV}$



Zbb BACKGROUND

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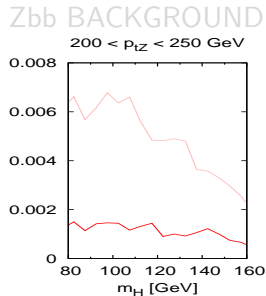
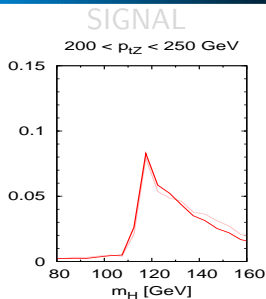
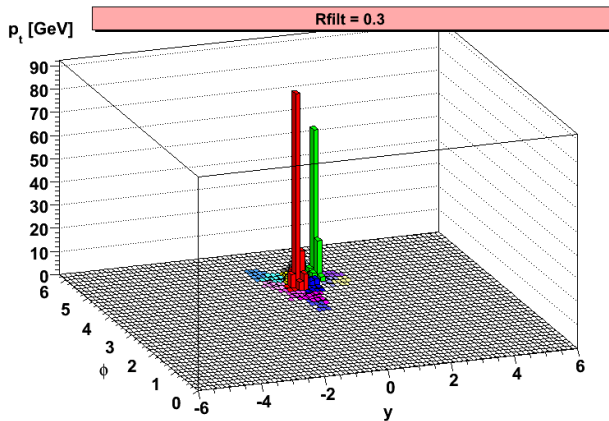
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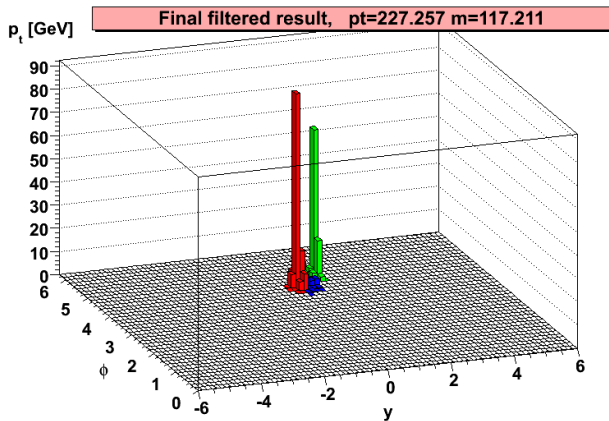
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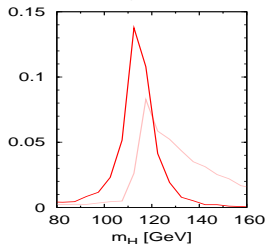
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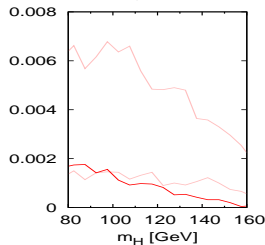
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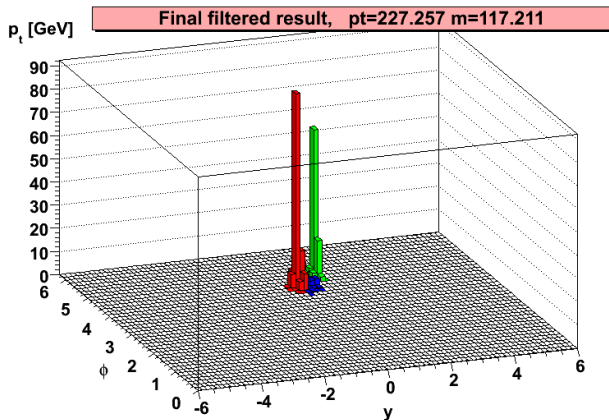
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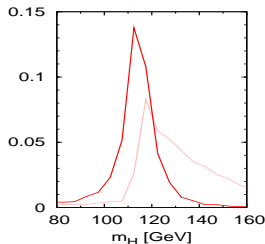
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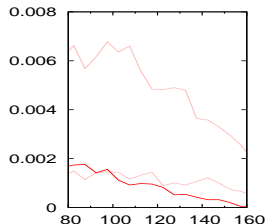
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Much to be learnt still about extracting boosted W/H/Z/top?

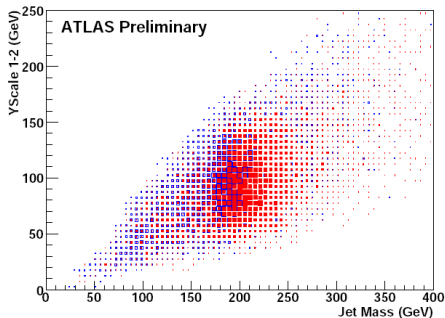
Brooijmans '08 ATL-PHYS-CONF-2008-008, based on k_t algorithm

+ Thaler & Wang '08; Almeida et al. '08 (k_t , jet-shapes)

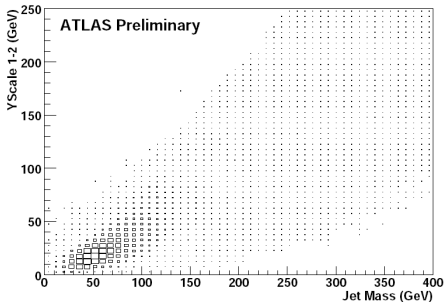
+ Kaplan et al '08 (C/A decomposition)

Use subjet relative transverse-momentum scale ("y-scale") & correlation with jet mass to pick out top quarks from background

top quarks $p_t \sim 1$ TeV



normal jets



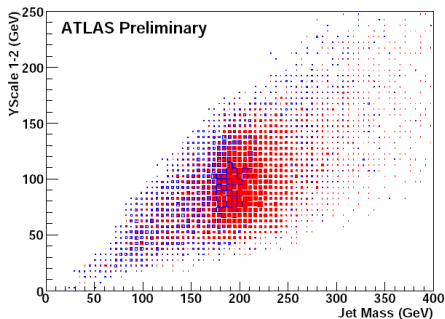
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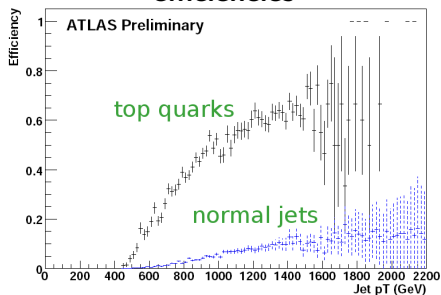
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efficiencies



Conclusions

- ▶ A jet is not a parton: it's (sort of) what you choose it to be.
- ▶ It's easier to think in terms of partons (LO, NLO pQCD) with IR/Collinear safe jet algorithms. And gives sense to pQCD predictions
- ▶ \exists many cones algs. Not equivalent. Many are IR/Coll unsafe.
xC-SM \rightarrow SISCone; xC-PR \rightarrow anti- k_t
- ▶ "The best" jet definition *does not exist*
- ▶ To get the most out of jet-algs.,
 - ▶ Understand the interplay of physical scales high $p_t \rightarrow$ larger R
 - ▶ Try out different combinations of algorithm & R
 - ▶ Check Variations of alg. & R don't change extracted physical quantities
- ▶ Special cases (e.g. boosted W/t/...) benefit from special techniques
e.g. seq. recomb. "jet-decomposition" is a powerful tool