

Jets, our window on partons at the LHC

Gavin P. Salam

LPTHE, UPMC Paris 6 & CNRS

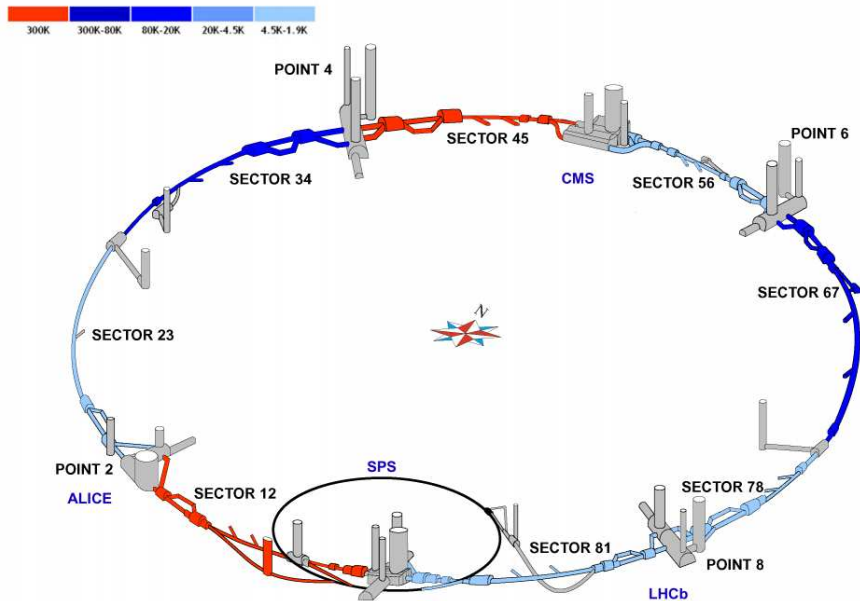
Torino, 23 May 2008

Basics: Cacciari (LPTHE) & Soyez (BNL)

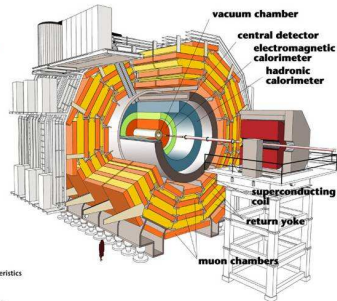
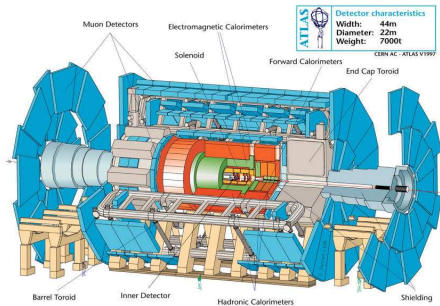
Higgs: Butterworth, Davison (ATLAS UCL) & Rubin (LPTHE)

Other related work: Dasgupta (Manchester), Magnea (Turin), Rojo (LPTHE)





2 general purpose detectors



Compared to current biggest collider (Tevatron)

- ▶ LHC energy will be **7 times higher**
- ▶ Total number of collisions (over 6 years) **50 times higher**

Aims are varied; Higgs discovery top priority

Last undiscovered component of standard model

ϕ has vacuum expectation value v ,

$\phi = v + H \leftrightarrow$ particle masses

$$\mathcal{L} = \dots + (v + H)^2 \bar{q}q + (v + H)^2 W^+ W^- + \dots$$

Excitations H around v are the Higgs \equiv sign of what's going on.

Plus searches for anything NEW in this energy domain

Compared to current biggest collider (Tevatron)

- ▶ LHC energy will be **7 times higher**
- ▶ Total number of collisions (over 6 years) **50 times higher**

Aims are varied; Higgs discovery top priority

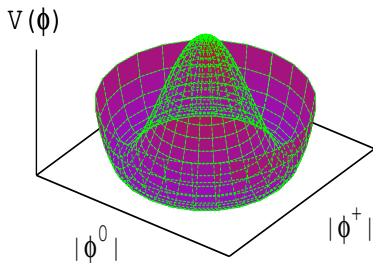
Last undiscovered component of standard model

ϕ has vacuum expectation value v ,

$\phi = v + H \leftrightarrow$ particle masses

$$\mathcal{L} = \dots + (v + H)^2 \bar{q}q + (v + H)^2 W^+ W^- + \dots$$

Excitations H around v are the Higgs \equiv sign of what's going on.



Plus searches for anything NEW in this energy domain

Compared to current biggest collider (Tevatron)

- ▶ LHC energy will be **7 times higher**
- ▶ Total number of collisions (over 6 years) **50 times higher**

Aims are varied; Higgs discovery top priority

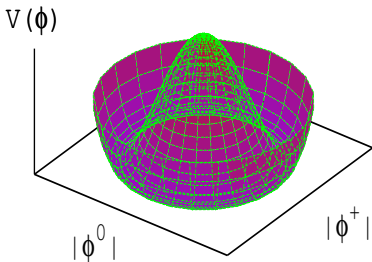
Last undiscovered component of standard model

ϕ has vacuum expectation value v ,

$\phi = v + H \leftrightarrow$ particle masses

$$\mathcal{L} = \dots + (v + H)^2 \bar{q}q + (v + H)^2 W^+ W^- + \dots$$

Excitations H around v are the Higgs \equiv sign of what's going on.



Plus searches for anything NEW in this energy domain

LHC is a parton collider

- ▶ Quarks and gluons are inevitable in initial state
- ▶ and ubiquitous in the final state

Partons — quarks and gluons — are key concepts of QCD.

- ▶ Lagrangian is in terms of quark and gluon fields
- ▶ Perturbative QCD *only* deals with partons

Though we often talk of quarks and gluons, we never see them

- ▶ Not an asymptotic state of the theory — because of confinement
- ▶ But also even in perturbation theory
 - because of collinear divergences (in massless approx.)
- ▶ The closest we can get to handling final-state partons is **jets**

1. Jets Introduction

quark

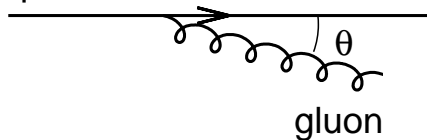
Gluon emission:

$$\int \alpha_s \frac{dE}{E} \frac{d\theta}{\theta} \gg 1$$

At low scales:

$$\alpha_s \rightarrow 1$$

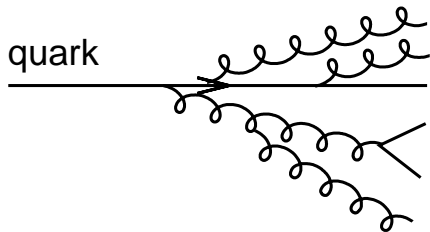
quark

Gluon emission:

$$\int \alpha_s \frac{dE}{E} \frac{d\theta}{\theta} \gg 1$$

At low scales:

$$\alpha_s \rightarrow 1$$

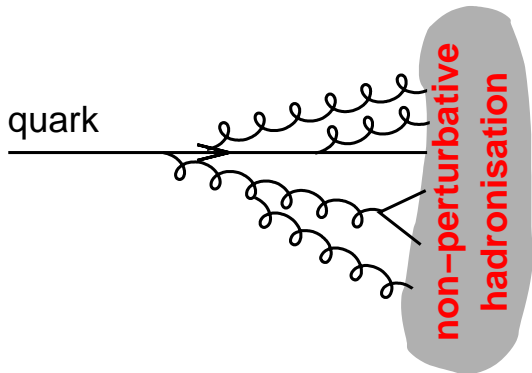


Gluon emission:

$$\int \alpha_s \frac{dE}{E} \frac{d\theta}{\theta} \gg 1$$

At low scales:

$$\alpha_s \rightarrow 1$$

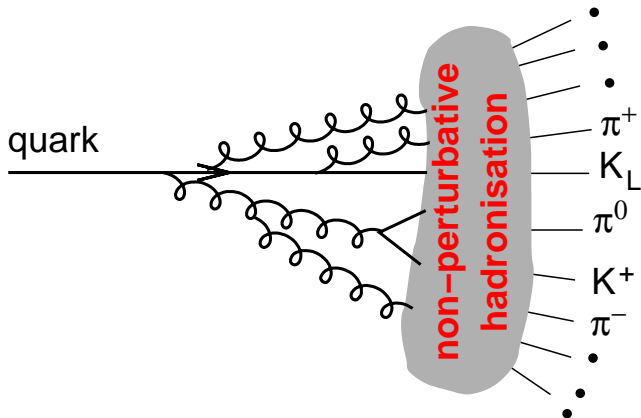


Gluon emission:

$$\int \alpha_s \frac{dE}{E} \frac{d\theta}{\theta} \gg 1$$

At low scales:

$$\alpha_s \rightarrow 1$$

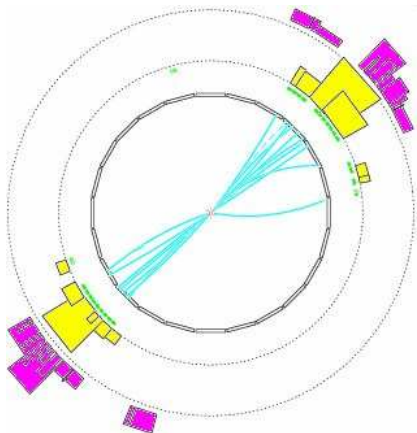


Gluon emission:

$$\int \alpha_s \frac{dE}{E} \frac{d\theta}{\theta} \gg 1$$

At low scales:

$$\alpha_s \rightarrow 1$$



Jets are what we see.

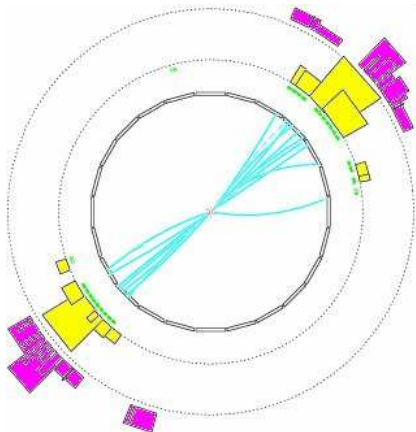
Clearly(?) 2 of them.

2 partons?

$$E_{parton} = M_Z/2?$$

How many jets do you see?

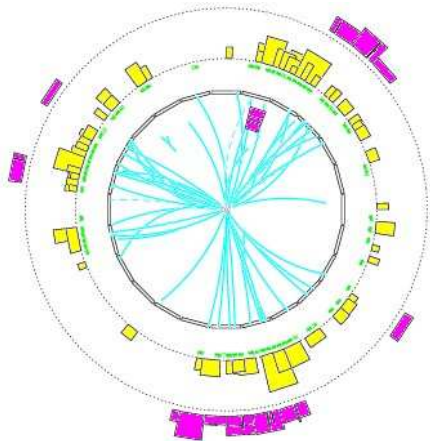
Do you really want to ask yourself
this question for 10^8 events?



Jets are what we see.
Clearly(?) 2 of them.

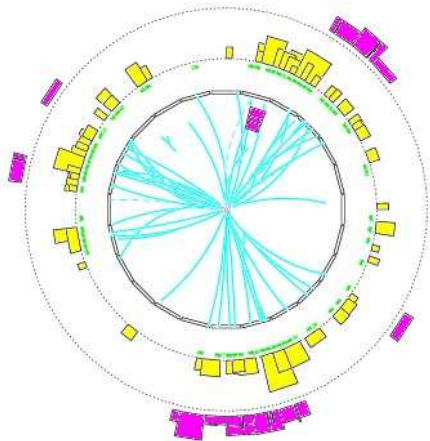
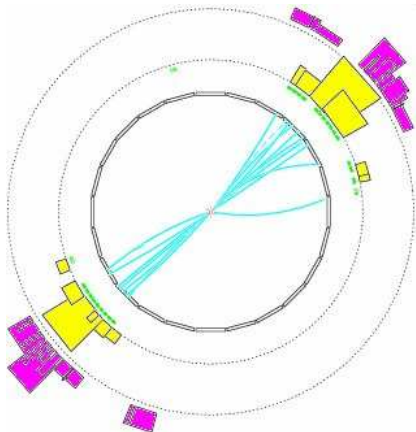
2 partons?

$$E_{parton} = M_Z/2?$$



How many jets do you see?

Do you really want to ask yourself
this question for 10^8 events?



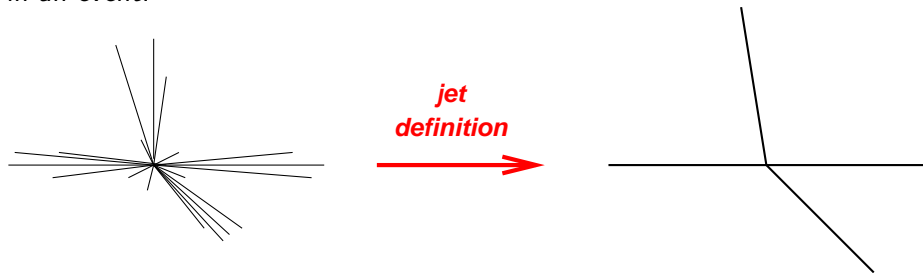
Jets are what we see.
Clearly(?) 2 of them.

2 partons?

$$E_{parton} = M_Z/2?$$

How many jets do you see?
Do you really want to ask yourself
this question for 10^8 events?

A jet definition is a systematic procedure that **projects away the multiparticle dynamics**, so as to leave a simple picture of what happened in an event:

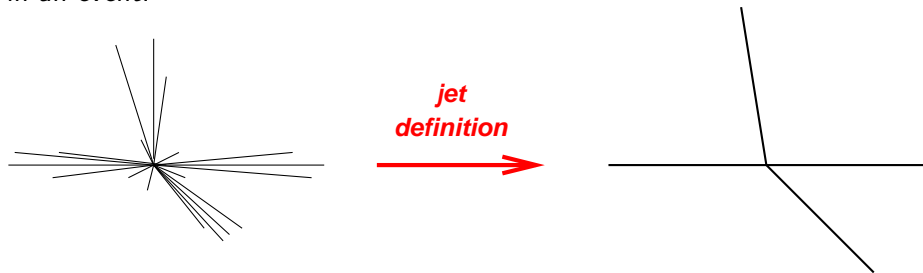


Jets are *as close as we can get to a physical single hard quark or gluon*: with good definitions their properties (multiplicity, energies, [flavour]) are

- ▶ finite at any order of perturbation theory
- ▶ insensitive to the parton \rightarrow hadron transition

NB: finiteness \longleftrightarrow set of jets depends on jet def.

A jet definition is a systematic procedure that **projects away the multiparticle dynamics**, so as to leave a simple picture of what happened in an event:



Jets are *as close as we can get to a physical single hard quark or gluon*: with good definitions their properties (multiplicity, energies, [flavour]) are

- ▶ finite at any order of perturbation theory
- ▶ insensitive to the parton \rightarrow hadron transition

NB: finiteness \longleftrightarrow set of jets depends on jet def.

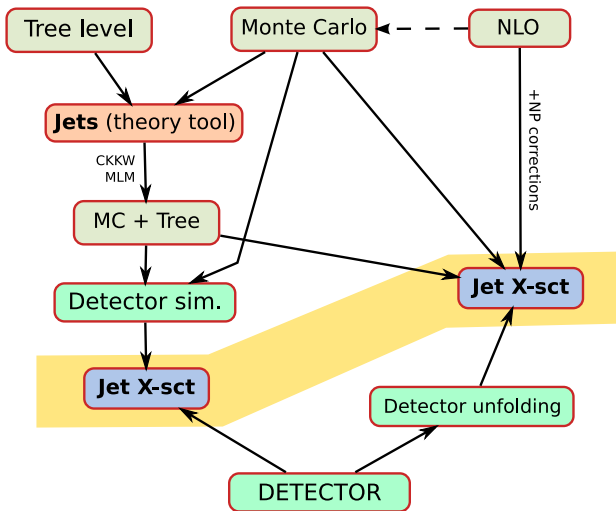
A jet definition is a systematic procedure that **projects away the multiparticle dynamics**, so as to leave a simple picture of what happened in an event:



Jets are *as close as we can get to a physical single hard quark or gluon*: with good definitions their properties (multiplicity, energies, [flavour]) are

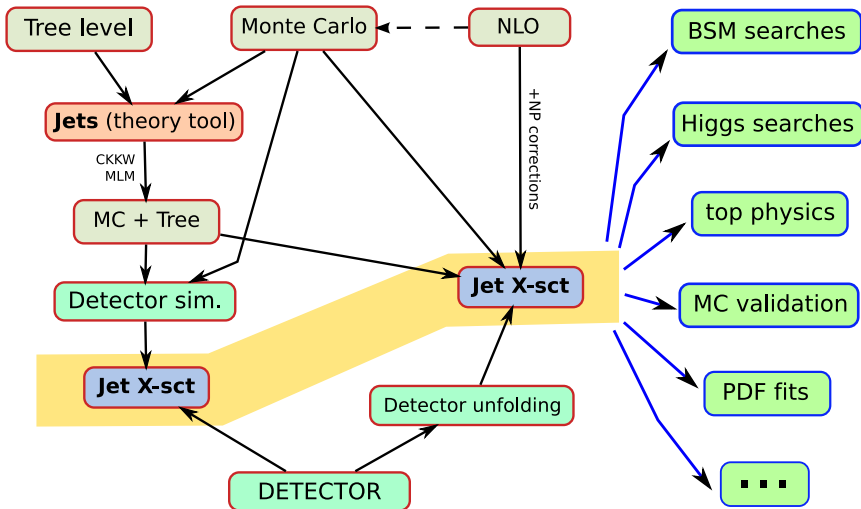
- ▶ finite at any order of perturbation theory
- ▶ insensitive to the parton \rightarrow hadron transition

NB: finiteness \longleftrightarrow set of jets depends on jet def.



Jet (definitions) provide central link between expt., “theory” and theory

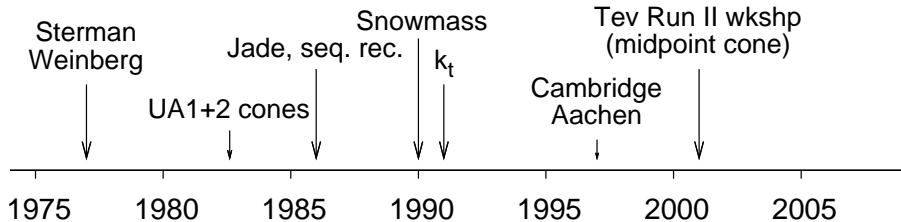
And jets are an input to almost all analyses



Jet (definitions) provide central link between expt., "theory" and theory

And jets are an input to almost all analyses

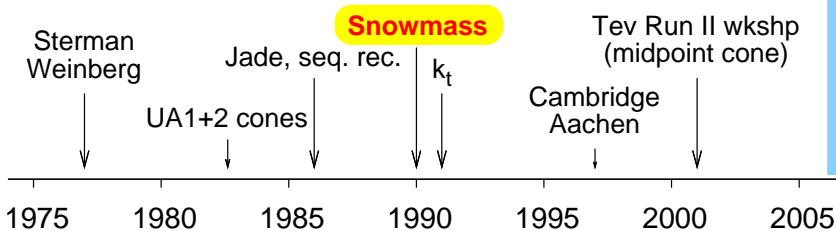
- ▶ Periodic key developments in jet definitions spurred by ever-increasing experimental/theoretical sophistication.
- ▶ Approach of LHC provides motivation for taking a new, fresh, systematic look at jets.
- ▶ This talk: **some of the discoveries along the way**



Definitions shown are those with widest exptl. impact

NB: also ARCLUS, OJF, ...

- ▶ Periodic key developments in jet definitions spurred by ever-increasing experimental/theoretical sophistication.
- ▶ Approach of LHC provides motivation for taking a new, fresh, systematic look at jets.
- ▶ This talk: **some of the discoveries along the way**



Speed, IR safety, Jet Areas
Non-pert. effects, Jet Flavour

Definitions shown are those with widest exptl. impact

NB: also ARCLUS, OJF, ...

Number of particles:

Experiment	N
LEP, HERA	50
Tevatron	100–400
LHC low-lumi	800
LHC high-lumi	4000
LHC PbPb	30000

- ▶ Range & complexity of signatures (jets, $t\bar{t}$, tj , Wj , Hj , $t\bar{t}j$, WWj , Wjj , SUSY, etc.)
- ▶ e.g. ~ 5 million $t\bar{t} \rightarrow 6$ jet events/year
- ▶ Theory investment
 ~ 100 people \times 10 years
 60 – 100 million \$

Physics scales:

Experiment	Physics	Scale
LEP, HERA	Electroweak + Hadronisation	100 GeV 0.5 GeV
Tevatron \rightarrow LHC	+ Underlying event	4 \rightarrow 15 GeV?
LHC	+ BSM + Pileup	1 TeV? 30 – 120 GeV

Number of particles:

Experiment	N
LEP, HERA	50
Tevatron	100–400
LHC low-lumi	800
LHC high-lumi	4000
LHC PbPb	30000

- ▶ Range & complexity of signatures (jets, $t\bar{t}$, tj , Wj , Hj , $t\bar{t}j$, WWj , Wjj , SUSY, etc.)
- ▶ e.g. ~ 5 million $t\bar{t} \rightarrow 6$ jet events/year
- ▶ Theory investment

~ 100 people \times 10 years
60 – 100 million \$

Physics scales:

Experiment	Physics	Scale
LEP, HERA	Electroweak + Hadronisation	100 GeV 0.5 GeV
Tevatron \rightarrow LHC	+ Underlying event	4 \rightarrow 15 GeV?
LHC	+ BSM + Pileup	1 TeV? 30 – 120 GeV

Number of particles:

Experiment	N
LEP, HERA	50
Tevatron	100–400
LHC low-lumi	800
LHC high-lumi	4000
LHC PbPb	30000

- ▶ Range & complexity of signatures (jets, $t\bar{t}$, tj , Wj , Hj , $t\bar{t}j$, WWj , Wjj , SUSY, etc.)
- ▶ e.g. ~ 5 million $t\bar{t} \rightarrow 6$ jet events/year
- ▶ Theory investment

~ 100 people \times 10 years
 60 – 100 million \$

Physics scales:

Experiment	Physics	Scale
LEP, HERA	Electroweak	100 GeV
	+ Hadronisation	0.5 GeV
Tevatron \rightarrow LHC	+ Underlying event	4 \rightarrow 15 GeV?
LHC	+ BSM	1 TeV?
	+ Pileup	30 – 120 GeV

Snowmass Accord (1990):

 FERMILAB-Conf-90/249-E
 [E-741/CDF]

Toward a Standardization of Jet Definitions ·

Several important properties that should be met by a jet definition are [3]:

1. Simple to implement in an experimental analysis;
2. Simple to implement in the theoretical calculation;
3. Defined at any order of perturbation theory;
4. Yields finite cross section at any order of perturbation theory;
5. Yields a cross section that is relatively insensitive to hadronization.

Without these, either the experiment won't use the jet-definition, or the theoretical calculations will be compromised

Long satisfied in e^+e^- and DIS

Satisfied in $\lesssim 10\%$ of jet work at Tevatron

Hardly discussed in LHC TDRs

Snowmass Accord (1990):

 FERMLAB-Conf-90/249-E
 [E-741/CDF]

Toward a Standardization of Jet Definitions ·

Several important properties that should be met by a jet definition are [3]:

1. Simple to implement in an experimental analysis;
2. Simple to implement in the theoretical calculation;
3. Defined at any order of perturbation theory;
4. Yields finite cross section at any order of perturbation theory;
5. Yields a cross section that is relatively insensitive to hadronization.

Without these, either the experiment won't use the jet-definition, or the theoretical calculations will be compromised

Long satisfied in e^+e^- and DIS
 Satisfied in $\lesssim 10\%$ of jet work at Tevatron
 Hardly discussed in LHC TDRs

2. Safe, practical jet-finding

Sequential recombination

k_t , Jade, Cam/Aachen, ...

Bottom-up:

Cluster 'closest' particles repeatedly until few left → jets.

Works because of mapping:

closeness \Leftrightarrow *QCD divergence*

Loved by e^+e^- , ep and theorists

Cone

UA1, JetClu, Midpoint, ...

Top-down:

Find coarse regions of energy flow (cones), and call them jets.

Works because *QCD only modifies energy flow on small scales*

Loved by pp and few(er) theorists

Both had serious issues that got in way of practical use and/or physical validity

Two classes of jet algorithm

Sequential recombination k_t , Jade, Cam/Aachen, ...**Bottom-up:**Cluster 'closest' particles repeatedly until few left \rightarrow jets.

Works because of mapping:

closeness \Leftrightarrow *QCD divergence*Loved by e^+e^- , ep and theorists**Cone**

UA1, JetClu, Midpoint, ...

Top-down:

Find coarse regions of energy flow (cones), and call them jets.

Works because *QCD only modifies energy flow on small scales*Loved by pp and few(er) theorists

Both had serious issues that got in way of practical use and/or physical validity

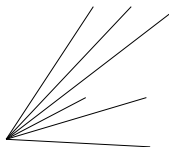
Sequential recombination algorithms

k_t algorithm

Catani, Dokshitzer, Olsson, Seymour, Turnock, Webber '91-'93

Ellis, Soper '93

- ▶ Find smallest of all $d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2$ and $d_{iB} = k_i^2$
- ▶ Recombine i, j (if iB : $i \rightarrow \text{jet}$)
- ▶ Repeat



NB: hadron collider variables

▶ $\Delta R_{ij}^2 = (\phi_i - \phi_j)^2 + (y_i - y_j)^2$

▶ rapidity $y_i = \frac{1}{2} \ln \frac{E_i + p_{zi}}{E_i - p_{zi}}$

▶ ΔR_{ij} is boost invariant angle

R sets jet opening angle

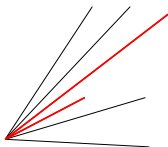
Sequential recombination algorithms

k_t algorithm

Catani, Dokshizter, Olsson, Seymour, Turnock, Webber '91-'93

Ellis, Soper '93

- ▶ Find **smallest of all** $d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2$ and $d_{iB} = k_i^2$
- ▶ Recombine i, j (if iB : $i \rightarrow \text{jet}$)
- ▶ Repeat



NB: hadron collider variables

▶ $\Delta R_{ij}^2 = (\phi_i - \phi_j)^2 + (y_i - y_j)^2$

▶ rapidity $y_i = \frac{1}{2} \ln \frac{E_i + p_{zi}}{E_i - p_{zi}}$

▶ ΔR_{ij} is boost invariant angle

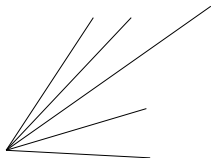
R sets jet opening angle

Sequential recombination algorithms

k_t algorithm

Catani, Dokshizter, Olsson, Seymour, Turnock, Webber '91-'93
 Ellis, Soper '93

- ▶ Find smallest of all $d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2$ and $d_{iB} = k_i^2$
- ▶ **Recombine** i, j (if iB : $i \rightarrow \text{jet}$)
- ▶ Repeat



NB: hadron collider variables

- ▶ $\Delta R_{ij}^2 = (\phi_i - \phi_j)^2 + (y_i - y_j)^2$
- ▶ rapidity $y_i = \frac{1}{2} \ln \frac{E_i + p_{zi}}{E_i - p_{zi}}$
- ▶ ΔR_{ij} is boost invariant angle

R sets jet opening angle

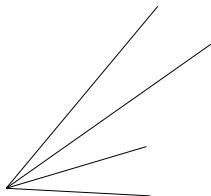
Sequential recombination algorithms

k_t algorithm

Catani, Dokshitzer, Olsson, Seymour, Turnock, Webber '91-'93

Ellis, Soper '93

- ▶ Find smallest of all $d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2$ and $d_{iB} = k_i^2$
- ▶ Recombine i, j (if iB : $i \rightarrow \text{jet}$)
- ▶ Repeat



NB: hadron collider variables

- ▶ $\Delta R_{ij}^2 = (\phi_i - \phi_j)^2 + (y_i - y_j)^2$
- ▶ rapidity $y_i = \frac{1}{2} \ln \frac{E_i + p_{zi}}{E_i - p_{zi}}$
- ▶ ΔR_{ij} is boost invariant angle

R sets jet opening angle

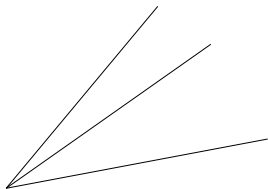
Sequential recombination algorithms

 k_t algorithm

Catani, Dokshitzer, Olsson, Seymour, Turnock, Webber '91-'93

Ellis, Soper '93

- ▶ Find smallest of all $d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2$ and $d_{iB} = k_i^2$
- ▶ Recombine i, j (if iB : $i \rightarrow \text{jet}$)
- ▶ Repeat

NB: hadron collider variables

- ▶ $\Delta R_{ij}^2 = (\phi_i - \phi_j)^2 + (y_i - y_j)^2$
- ▶ rapidity $y_i = \frac{1}{2} \ln \frac{E_i + p_{zi}}{E_i - p_{zi}}$
- ▶ ΔR_{ij} is boost invariant angle

R sets jet opening angle

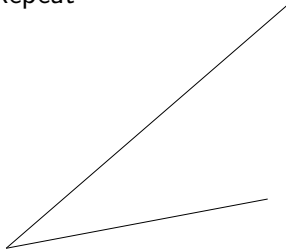
Sequential recombination algorithms

k_t algorithm

Catani, Dokshitzer, Olsson, Seymour, Turnock, Webber '91-'93

Ellis, Soper '93

- ▶ Find smallest of all $d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2$ and $d_{iB} = k_i^2$
- ▶ Recombine i, j (if iB : $i \rightarrow \text{jet}$)
- ▶ Repeat



NB: hadron collider variables

- ▶ $\Delta R_{ij}^2 = (\phi_i - \phi_j)^2 + (y_i - y_j)^2$
- ▶ rapidity $y_i = \frac{1}{2} \ln \frac{E_i + p_{zi}}{E_i - p_{zi}}$
- ▶ ΔR_{ij} is boost invariant angle

R sets jet opening angle

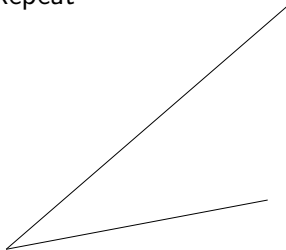
Sequential recombination algorithms

k_t algorithm

Catani, Dokshizter, Olsson, Seymour, Turnock, Webber '91-'93

Ellis, Soper '93

- ▶ Find smallest of all $d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2 / R^2$ and $d_{iB} = k_i^2$
- ▶ Recombine i, j (if iB : $i \rightarrow \text{jet}$)
- ▶ Repeat



NB: hadron collider variables

- ▶ $\Delta R_{ij}^2 = (\phi_i - \phi_j)^2 + (y_i - y_j)^2$
- ▶ rapidity $y_i = \frac{1}{2} \ln \frac{E_i + p_{zi}}{E_i - p_{zi}}$
- ▶ ΔR_{ij} is boost invariant angle

R sets jet opening angle

k_t distance measures

$$d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2, \quad d_{iB} = k_{ti}^2$$

are closely related to structure of divergences for QCD emissions

$$[dk_j] |M_{g \rightarrow g_i g_j}^2(k_j)| \sim \frac{\alpha_s C_A}{2\pi} \frac{dk_{tj}}{\min(k_{ti}, k_{tj})} \frac{d\Delta R_{ij}}{\Delta R_{ij}}, \quad (k_{tj} \ll k_{ti}, \Delta R_{ij} \ll 1)$$

and

$$[dk_i] |M_{Beam \rightarrow Beam + g_i}^2(k_i)| \sim \frac{\alpha_s C_A}{\pi} \frac{dk_{ti}}{k_{ti}} d\eta_i, \quad (k_{ti}^2 \ll \{\hat{s}, \hat{t}, \hat{u}\})$$

k_t algorithm attempts approximate inversion of branching process

k_t distance measures

$$d_{ij} = \min(k_{ti}^2, k_{tj}^2) \Delta R_{ij}^2, \quad d_{iB} = k_{ti}^2$$

are closely related to structure of divergences for QCD emissions

$$[dk_j] |M_{g \rightarrow g_i g_j}^2(k_j)| \sim \frac{\alpha_s C_A}{2\pi} \frac{dk_{tj}}{\min(k_{ti}, k_{tj})} \frac{d\Delta R_{ij}}{\Delta R_{ij}}, \quad (k_{tj} \ll k_{ti}, \Delta R_{ij} \ll 1)$$

and

$$[dk_i] |M_{Beam \rightarrow Beam + g_i}^2(k_i)| \sim \frac{\alpha_s C_A}{\pi} \frac{dk_{ti}}{k_{ti}} d\eta_i, \quad (k_{ti}^2 \ll \{\hat{s}, \hat{t}, \hat{u}\})$$

k_t algorithm attempts approximate inversion of branching process

'Trivial' computational issue:

- ▶ for N particles: $N^2 d_{ij}$ searched through N times = N^3
- ▶ 4000 particles (or calo cells): **1 minute**
NB: often study $10^7 - 10^8$ events (20-200 CPU years)
- ▶ Heavy Ions: 30000 particles: **10 hours/event**

As far as possible physics choices should not be limited by computing.

Even if we're clever about repeating the full search each time, we still have $\mathcal{O}(N^2)$ d_{ij} 's to establish

Fast Hierarchical Clustering and Other Applications of Dynamic Closest Pairs

David Eppstein
UC Irvine

We develop data structures for dynamic closest pair problems with arbitrary distance functions, that do not necessarily come from any geometric structure on the objects. Based on a technique previously used by the author for Euclidean closest pairs, we show how to insert and delete objects from an n -object set, maintaining the closest pair, in $O(n \log^2 n)$ time per update and $O(n)$ space. With quadratic space, we can instead use a quadtree-like structure to achieve an optimal time bound, $O(n)$ per update. We apply these data structures to hierarchical clustering, greedy matching, and TSP heuristics, and discuss other potential applications in machine learning, Gröbner bases, and local improvement algorithms for partition and placement problems. Experiments show our new methods to be faster in practice than previously used heuristics.

Categories and Subject Descriptors: F.2.2 [Analysis of Algorithms]: Nonnumeric Algorithms

General Terms: Closest Pair, Agglomerative Clustering

Additional Key Words and Phrases: TSP, matching, conga line data structure, quadtree, nearest neighbor heuristic

1. INTRODUCTION

Hierarchical clustering has long been a mainstay of statistical analysis, and clustering based methods have attracted attention in other fields: computational biology (reconstruction of evolutionary trees; tree-based multiple sequence alignment), scientific simulation (n -body problems), theoretical computer science (network design and nearest neighbor searching) and of course the web (hierarchical indices such as Yahoo). Many clustering methods have been devised and used in these applications, but less effort has gone into algorithmic speedups of these methods.

In this paper we identify and demonstrate speedups for a key subroutine used in several clustering algorithms, that of maintaining closest pairs in a dynamic set of objects. We also describe several other applications or potential applications of the

k_t alg. is so good it's used throughout science!

NB HEP is not only field to use brute-force...

For general distance measures problem reduces to $\sim N^2$ (factor ~ 20 for $N = 1000$).

Eppstein '99
+ Cardinal '03

Fast Hierarchical Clustering and Other Applications of Dynamic Closest Pairs

David Eppstein
UC Irvine

We develop data structures for dynamic closest pair problems with arbitrary distance functions, that do not necessarily come from any geometric structure on the objects. Based on a technique previously used by the author for Euclidean closest pairs, we show how to insert and delete objects from an n -object set, maintaining the closest pair, in $O(n \log^2 n)$ time per update and $O(n)$ space. With quadratic space, we can instead use a quadtree-like structure to achieve an optimal time bound, $O(n)$ per update. We apply these data structures to hierarchical clustering, greedy matching, and TSP heuristics, and discuss other potential applications in machine learning, Gröbner bases, and local improvement algorithms for partition and placement problems. Experiments show our new methods to be faster in practice than previously used heuristics.

Categories and Subject Descriptors: F.2.2 [Analysis of Algorithms]: Nonnumeric Algorithms

General Terms: Closest Pair, Agglomerative Clustering

Additional Key Words and Phrases: TSP, matching, conga line data structure, quadtree, nearest neighbor heuristic

Of these naive methods, brute force recomputation may be most commonly used, due to its low space requirements and ease of implementation. Three hierarchical clustering codes we examined, Zupan's [Zupan 1982], CLUSTAL W [Thompson et al. 1994], and PHYLIP [Felsenstein 1995] use brute force. (Indeed, they do not even save space by doing so, since they all store the distance matrix.) Pazzani's learning code [Pazzani 1997] also uses brute force (M. Pazzani, personal communication), as does *Mathematica's* Gröbner basis code (D. Lichtblau, personal communication).

k_t alg. is so good it's used throughout science!

NB HEP is not only field to use brute-force. . .

For general distance measures problem reduces to $\sim N^2$ (factor ~ 20 for $N = 1000$).

Eppstein '99
+ Cardinal '03

There are $N(N - 1)/2$ distances d_{ij} — surely we have to calculate them all in order to find smallest?

k_t distance measure is partly *geometrical*:

$$\begin{aligned} \min_{i,j} d_{ij} &\equiv \min_{i,j} (\min\{k_{ti}^2, k_{tj}^2\} \Delta_{ij}^2) \\ &= \min_{i,j} (k_{ti}^2 \Delta_{ij}^2) \\ &= \min_i (k_{ti}^2 \min_j \Delta_{ij}^2) \end{aligned}$$

In words: if i, j form smallest d_{ij} then j is geometrical nearest neighbour (GNN) of i .

k_t distance need only be calculated between GNNs

Each point has 1 GNN \rightarrow need only calculate N d_{ij} 's

There are $N(N - 1)/2$ distances d_{ij} — surely we have to calculate them all in order to find smallest?

k_t distance measure is partly *geometrical*:

$$\begin{aligned} \min_{i,j} d_{ij} &\equiv \min_{i,j} (\min\{k_{ti}^2, k_{tj}^2\} \Delta_{ij}^2) \\ &= \min_{i,j} (k_{ti}^2 \Delta_{ij}^2) \\ &= \min_i (k_{ti}^2 \min_j \Delta_{ij}^2) \end{aligned}$$

In words: if i, j form smallest d_{ij} then j is geometrical nearest neighbour (GNN) of i .

k_t distance need only be calculated between GNNs

Each point has 1 GNN \rightarrow need only calculate N d_{ij} 's

There are $N(N - 1)/2$ distances d_{ij} — surely we have to calculate them all in order to find smallest?

k_t distance measure is partly *geometrical*:

$$\begin{aligned} \min_{i,j} d_{ij} &\equiv \min_{i,j} (\min\{k_{ti}^2, k_{tj}^2\} \Delta_{ij}^2) \\ &= \min_{i,j} (k_{ti}^2 \Delta_{ij}^2) \\ &= \min_i (k_{ti}^2 \min_j \Delta_{ij}^2) \end{aligned}$$

In words: if i, j form smallest d_{ij} then j is geometrical nearest neighbour (GNN) of i .

k_t distance need only be calculated between GNNs

Each point has 1 GNN \rightarrow need only calculate N d_{ij} 's

There are $N(N - 1)/2$ distances d_{ij} — surely we have to calculate them all in order to find smallest?

k_t distance measure is partly *geometrical*:

$$\begin{aligned} \min_{i,j} d_{ij} &\equiv \min_{i,j} (\min\{k_{ti}^2, k_{tj}^2\} \Delta_{ij}^2) \\ &= \min_{i,j} (k_{ti}^2 \Delta_{ij}^2) \\ &= \min_i (k_{ti}^2 \min_j \Delta_{ij}^2) \end{aligned}$$

In words: if i, j form smallest d_{ij} then j is geometrical nearest neighbour (GNN) of i .

k_t distance need only be calculated between GNNs

Each point has 1 GNN \rightarrow need only calculate N d_{ij} 's

There are $N(N - 1)/2$ distances d_{ij} — surely we have to calculate them all in order to find smallest?

k_t distance measure is partly *geometrical*:

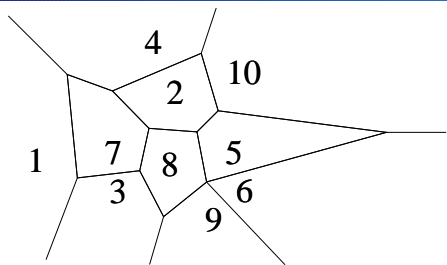
$$\begin{aligned} \min_{i,j} d_{ij} &\equiv \min_{i,j} (\min\{k_{ti}^2, k_{tj}^2\} \Delta_{ij}^2) \\ &= \min_{i,j} (k_{ti}^2 \Delta_{ij}^2) \\ &= \min_i (k_{ti}^2 \min_j \Delta_{ij}^2) \end{aligned}$$

In words: if i, j form smallest d_{ij} then j is geometrical nearest neighbour (GNN) of i .

k_t distance need only be calculated between GNNs

Each point has 1 GNN \rightarrow need only calculate N d_{ij} 's

Finding Geom Nearest Neighbours



Given a set of vertices on plane (1...10) a *Voronoi diagram* partitions plane into cells containing all points closest to each vertex

Dirichlet '1850, Voronoi '1908

A vertex's nearest other vertex is always in an adjacent cell.

E.g. GNN of point 7 will be found among 1,4,2,8,3 (it turns out to be 3)

Construction of Voronoi diagram for N points: $N \ln N$ time Fortune '88

Update of 1 point in Voronoi diagram: $\ln N$ time

Devillers '99 [+ related work by other authors]

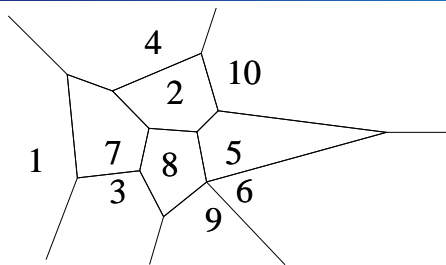
Convenient C++ package available: **CGAL** <http://www.cgal.org>

Assemble with other comp. science methods: **FastJet**

Cacciari & GPS, hep-ph/0512210

<http://www.lpthe.jussieu.fr/~salam/fastjet/>

Finding Geom Nearest Neighbours



Given a set of vertices on plane (1...10) a *Voronoi diagram* partitions plane into cells containing all points closest to each vertex

Dirichlet '1850, Voronoi '1908

A vertex's nearest other vertex is always in an adjacent cell.

E.g. GNN of point 7 will be found among 1,4,2,8,3 (it turns out to be 3)

Construction of Voronoi diagram for N points: $N \ln N$ time Fortune '88

Update of 1 point in Voronoi diagram: $\ln N$ time

Devillers '99 [+ related work by other authors]

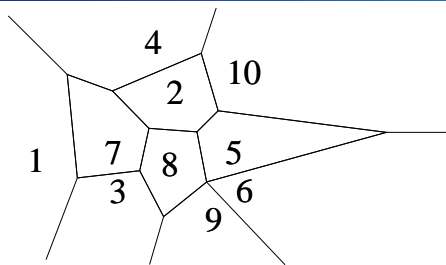
Convenient C++ package available: CGAL <http://www.cgal.org>

Assemble with other comp. science methods: FastJet

Cacciari & GPS, hep-ph/0512210

<http://www.lpthe.jussieu.fr/~salam/fastjet/>

Finding Geom Nearest Neighbours



Given a set of vertices on plane (1...10) a *Voronoi diagram* partitions plane into cells containing all points closest to each vertex

Dirichlet '1850, Voronoi '1908

A vertex's nearest other vertex is always in an adjacent cell.

E.g. GNN of point 7 will be found among 1,4,2,8,3 (it turns out to be 3)

Construction of Voronoi diagram for N points: $N \ln N$ time Fortune '88

Update of 1 point in Voronoi diagram: $\ln N$ time

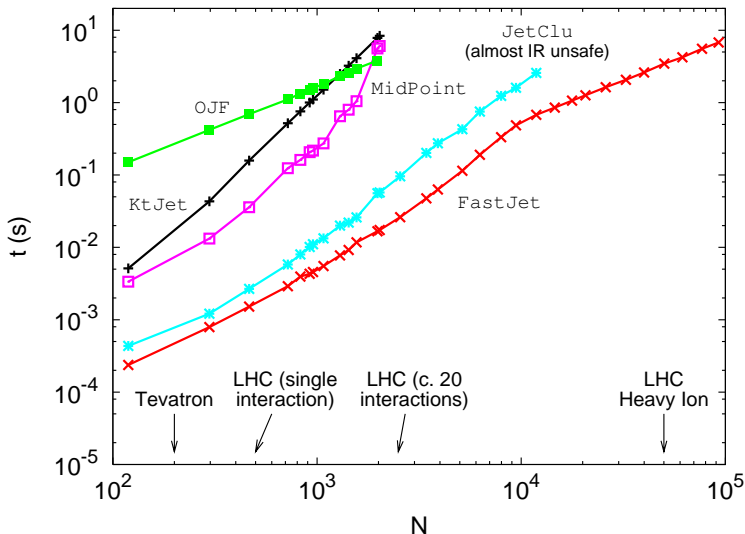
Devillers '99 [+ related work by other authors]

Convenient C++ package available: **CGAL** <http://www.cgal.org>

Assemble with other comp. science methods: **FastJet**

Cacciari & GPS, hep-ph/0512210

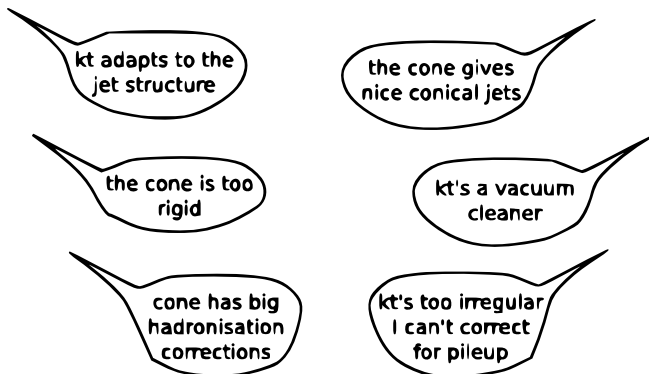
<http://www.lpthe.jussieu.fr/~salam/fastjet/>



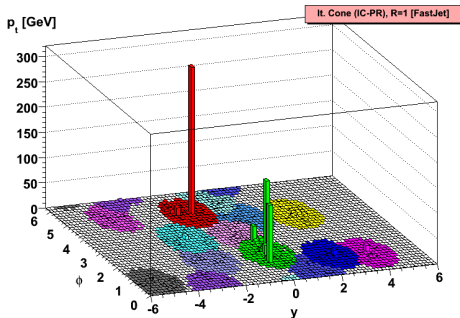
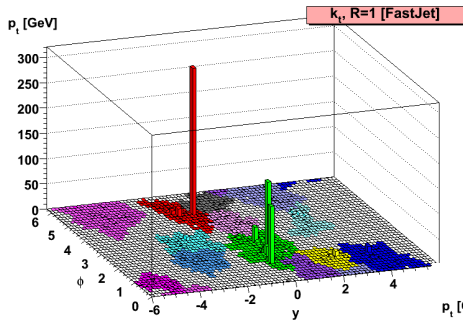
NB: for $N < 10^4$, FastJet switches to a related geometrical N^2 alg.

Conclusion: speed issues for k_t resolved

Jet discussions: often polarised, driven by unquantified statements



- ▶ Rigorous approach is to quantify similarities & differences
Dasgupta, Magnea & GPS '07; Cacciari, GPS & Soyez '08
- ▶ Bottom line: grains of truth in the qualitative statements
So want good cone algorithms too [NB: two varieties, IC-SM & IC-PR]



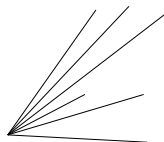
Iterative Cone [with progressive removal]

Procedure:

- ▶ Find one stable cone

By iterating from hardest seed particle

- ▶ Call it a jet; remove its particles from the event; repeat



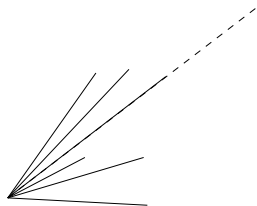
Iterative Cone [with progressive removal]

Procedure:

- ▶ Find one stable cone

By iterating from hardest seed particle

- ▶ Call it a jet; remove its particles from the event; repeat



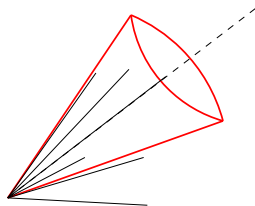
Iterative Cone [with progressive removal]

Procedure:

- ▶ Find one stable cone

By iterating from hardest seed particle

- ▶ Call it a jet; remove its particles from the event; repeat



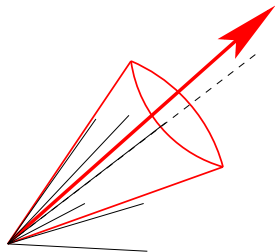
Iterative Cone [with progressive removal]

Procedure:

- ▶ Find one stable cone

By iterating from hardest seed particle

- ▶ Call it a jet; remove its particles from the event; repeat



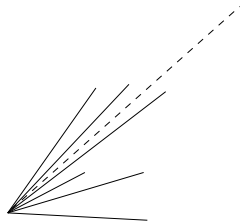
Iterative Cone [with progressive removal]

Procedure:

► Find one stable cone

By iterating from hardest seed particle

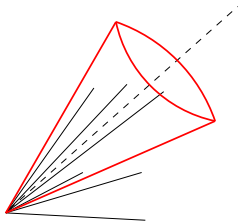
► Call it a jet; remove its particles from the event; repeat



Iterative Cone [with progressive removal]

Procedure:

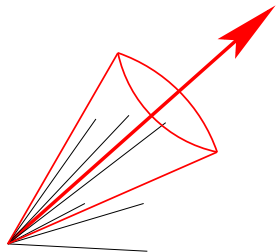
- ▶ Find one stable cone By iterating from hardest seed particle
- ▶ Call it a jet; remove its particles from the event; repeat



Iterative Cone [with progressive removal]

Procedure:

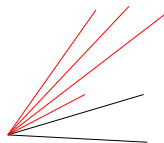
- ▶ Find one stable cone By iterating from hardest seed particle
- ▶ Call it a jet; remove its particles from the event; repeat



Iterative Cone [with progressive removal]

Procedure:

- ▶ Find one stable cone By iterating from hardest seed particle
- ▶ Call it a jet; remove its particles from the event; repeat



Iterative Cone [with progressive removal]

Procedure:

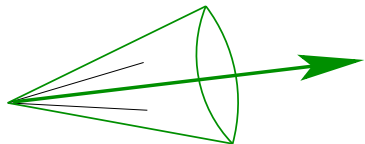
- ▶ Find one stable cone By iterating from hardest seed particle
- ▶ Call it a jet; remove its particles from the event; repeat



Iterative Cone [with progressive removal]

Procedure:

- ▶ Find one stable cone By iterating from hardest seed particle
- ▶ Call it a jet; remove its particles from the event; repeat



Iterative Cone [with progressive removal]

Procedure:

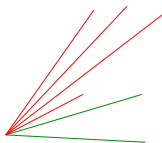
- ▶ Find one stable cone By iterating from hardest seed particle
- ▶ Call it a jet; remove its particles from the event; repeat



Iterative Cone [with progressive removal]

Procedure:

- ▶ Find one stable cone By iterating from hardest seed particle
- ▶ Call it a jet; remove its particles from the event; repeat



Iterative Cone [with progressive removal]

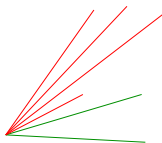
Procedure:

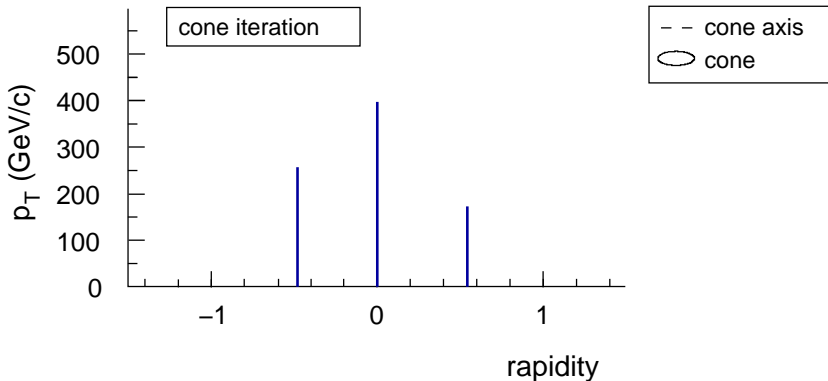
- ▶ Find one stable cone By iterating from hardest seed particle
- ▶ Call it a jet; remove its particles from the event; repeat

Iterative Cone with Progressive Removal (IC-PR)

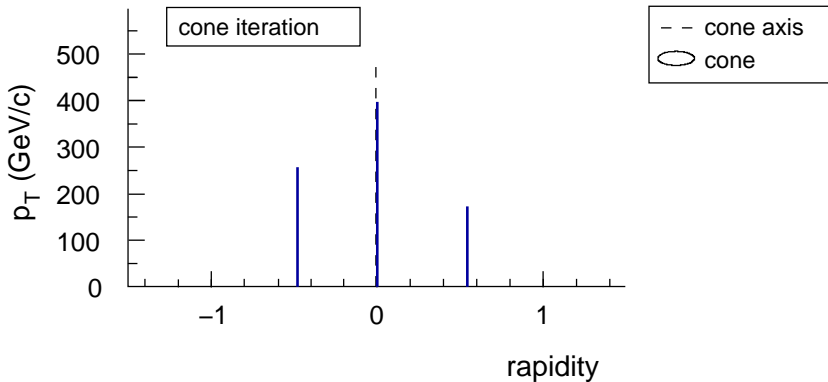
e.g. CMS it. cone, [Pythia Cone, GetJet], ...

- ▶ NB: not same type of algorithm as Atlas Cone, MidPoint, SIScone

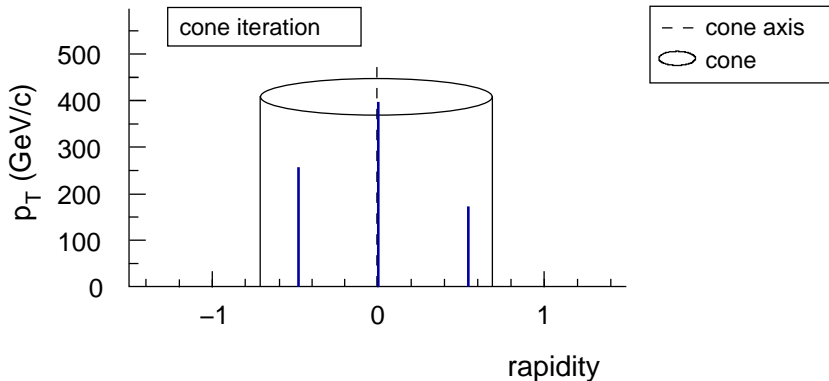




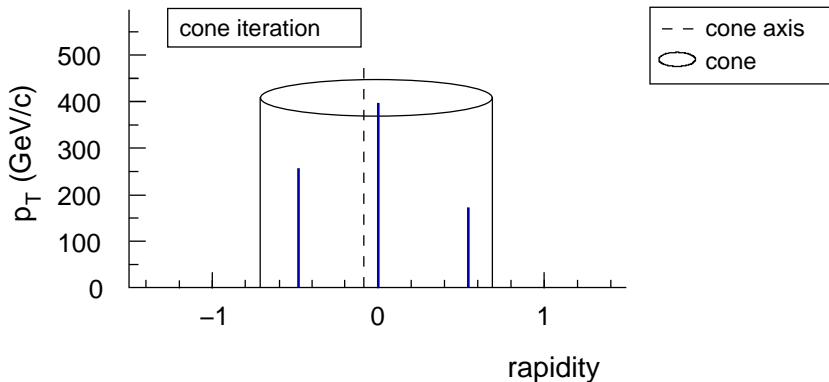
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



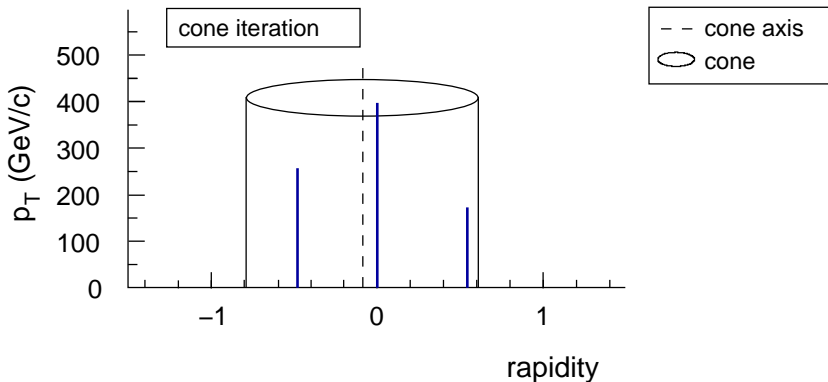
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



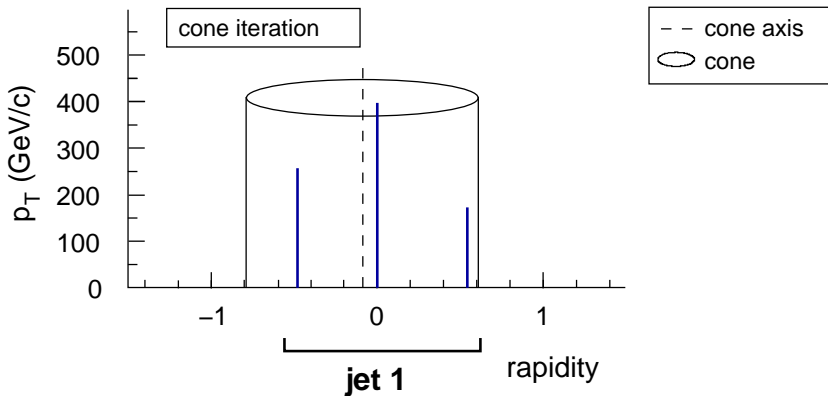
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



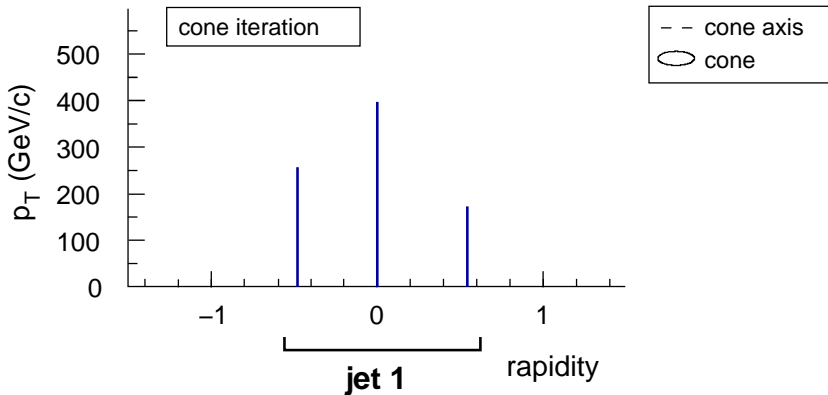
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



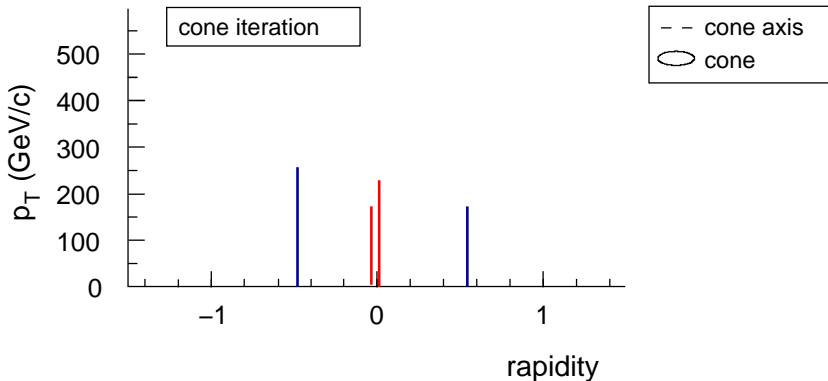
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



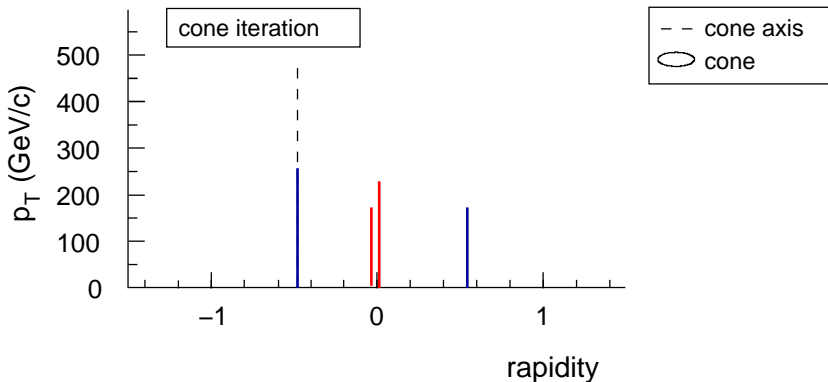
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



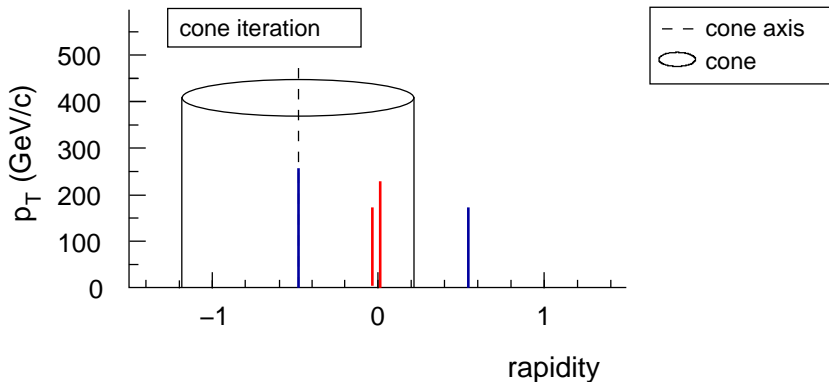
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



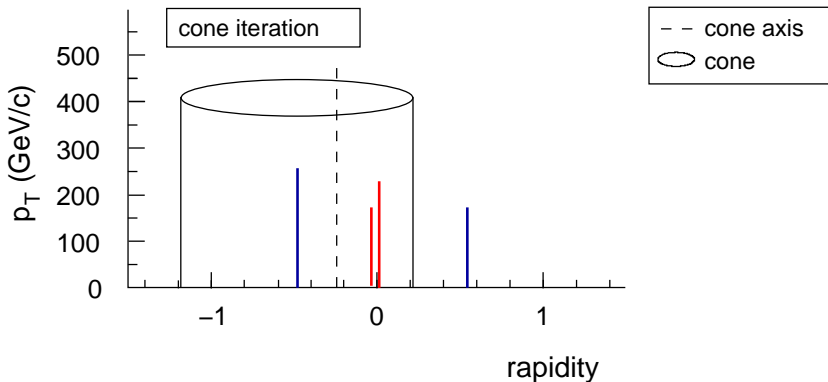
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



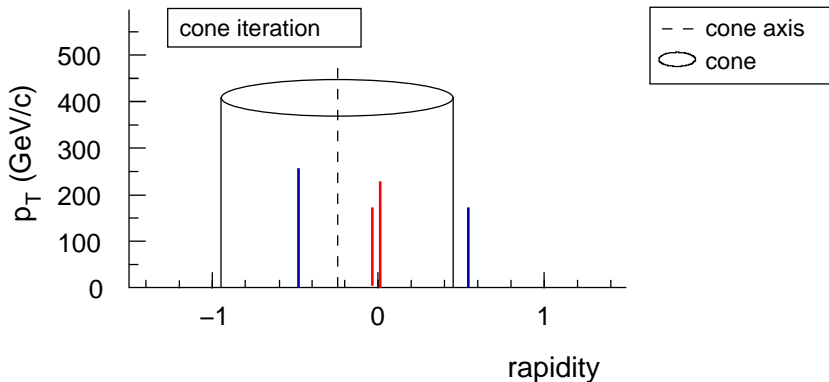
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



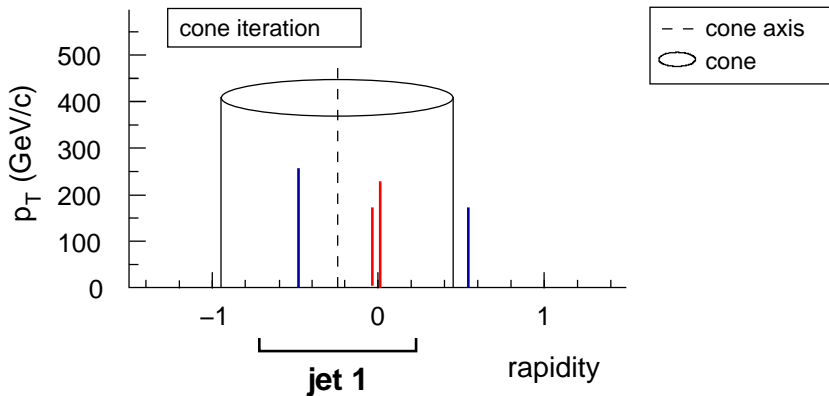
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



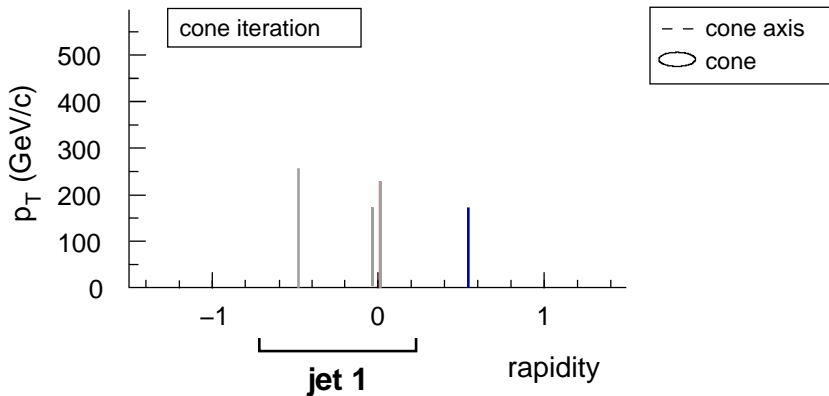
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



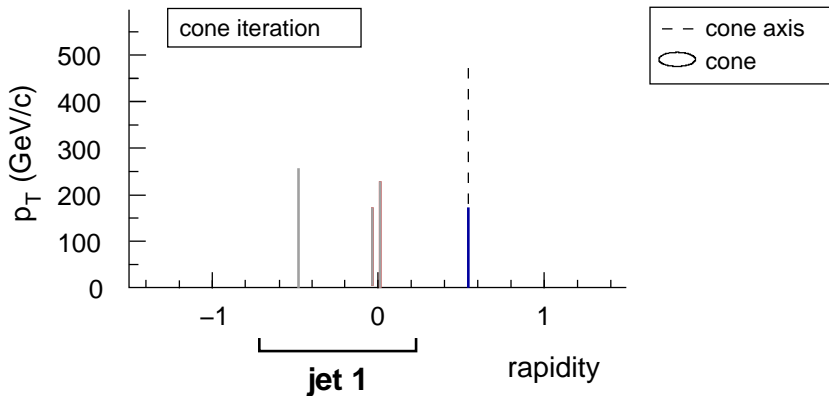
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



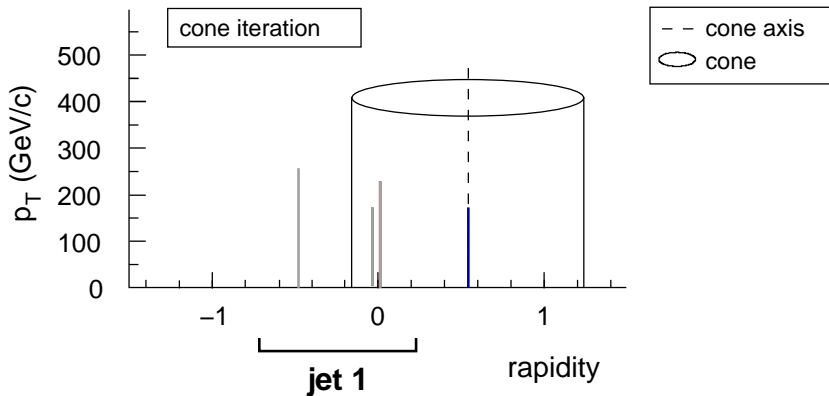
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



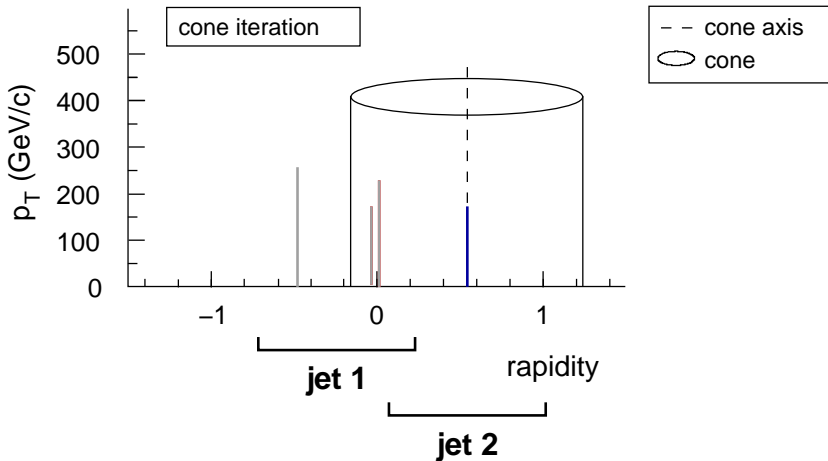
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



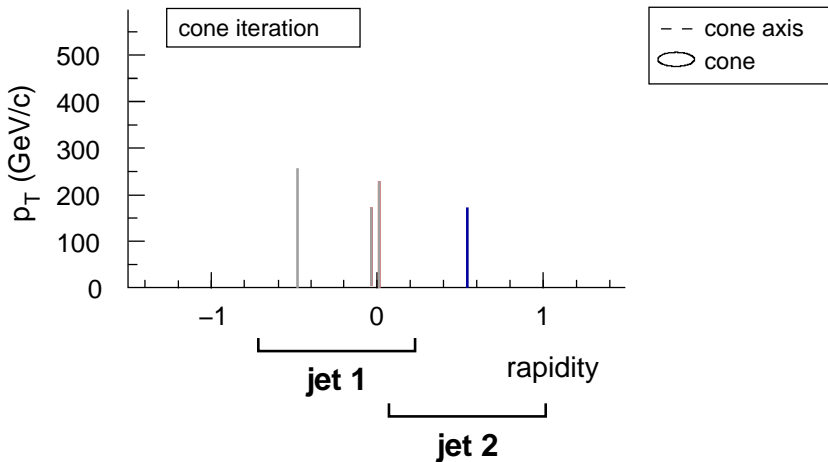
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



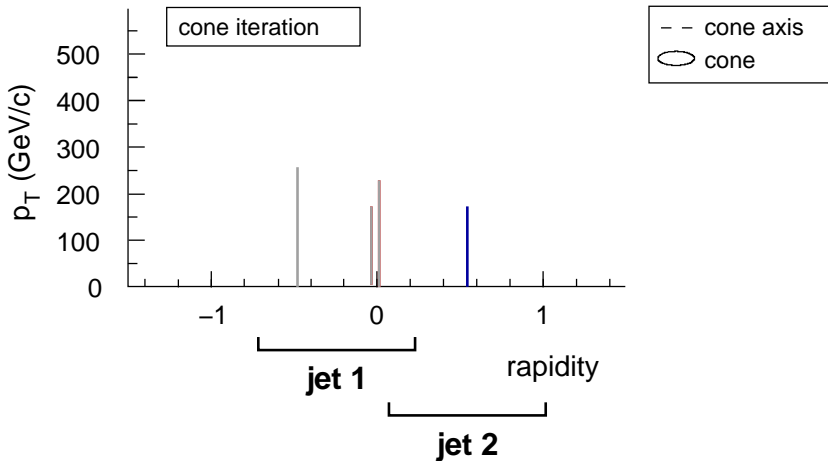
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



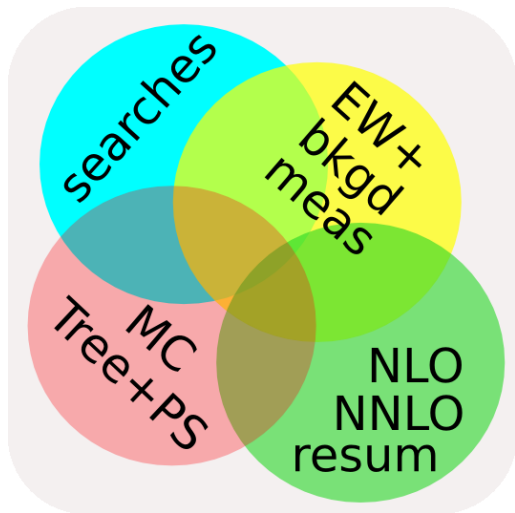
Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞



Collinear splitting can modify the hard jets: ICPR algorithms are collinear unsafe \implies perturbative calculations give ∞

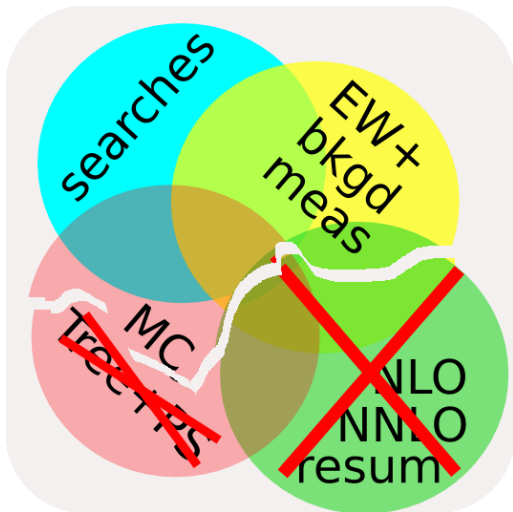


For everything to fit together **all of Snowmass criteria needed.**

Given need to compromise, the IRC safety usually goes first.

This breaks connection between different parts of QCD.

~ 90% of Tevatron and LHC work based on IRC unsafe algs — a pervasive problem.



For everything to fit together **all of Snowmass criteria needed.**

Given need to compromise, the IRC safety usually goes first.

This breaks connection between different parts of QCD.

~ 90% of Tevatron and LHC work based on IRC unsafe algs — a pervasive problem.

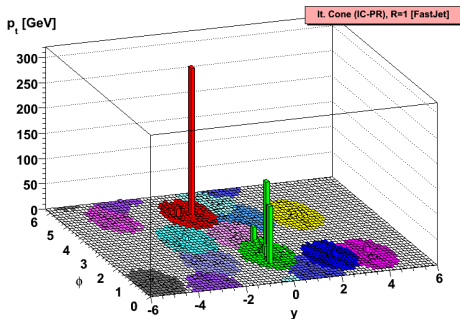
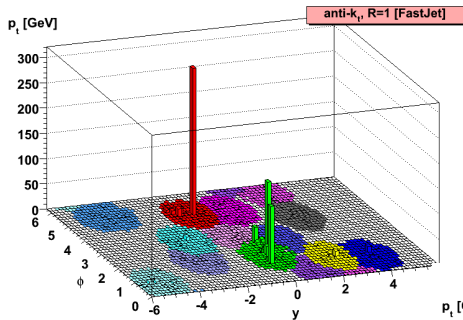
What we want: something that behaves like a cone algorithm (circular jets), but that is IRC safe.

Approach: drop the “cone” in definition, but design an algorithm that still acts like a cone: **anti- k_t**

1. Find smallest of d_{ij} , d_{iB} : $d_{ij} = \min(k_{ti}^{-2}, k_{tj}^{-2}) \Delta R_{ij}^2 / R^2$, $d_{iB} = k_{ti}^{-2}$
2. if ij , recombine them; if iB , call i a jet, and remove from list of particles
3. repeat from step 1 until no particles left.

Cacciari, GPS & Soyez '08

Looks like k_t but momentum in *denominator* causes d_{ij} to involve largest k_t
 → jets grows outward from hard “seeds”.



Complementary set of IR/Collinear safe jet algs \longrightarrow flexibility in studying complex events.

Consider families of jet algs: e.g. sequential recombination with

$$d_{ij} = \min(k_{ti}^{2p}, k_{tj}^{2p}) \Delta R_{ij}^2 / R^2$$

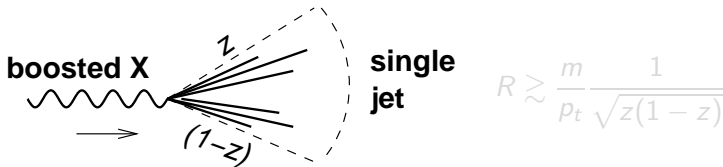
	Alg. name	Comp. Geometry problem	time
$p = 1$	k_t CDOSTW '91-93; ES '93	Dynamic Nearest Neighbour CGAL (Devillers et al)	$N \ln N$ exp.
$p = 0$	Cambridge/Aachen Dok, Leder, Moretti, Webber '97 Wengler, Wobisch '98	Dynamic Closest Pair T. Chan '02	$N \ln N$
$p = -1$	anti- k_t (cone-like) Cacciari, GPS, Soyez, in prep.	Dynamic Nearest Neighbour CGAL (worst case)	$N^{3/2}$
cone	SISCone GPS Soyez '07 + Tevatron run II '00	All circular enclosures previously unconsidered	$N^2 \ln N$ exp.

All accessible in FastJet
FastJet in software of all (4) LHC collaborations

3. An example: boosted Higgs search

Illustrate LHC challenges with a recently widely discussed class of problems:

Can you identify hadronically decaying EW bosons when they're produced at high p_t ?

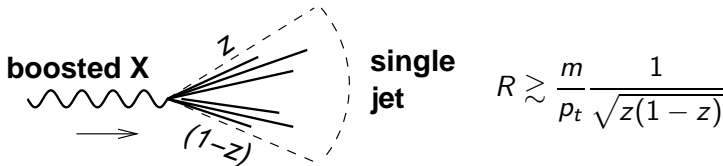


Significant discussion over years: **heavy new things decay to EW states**

- ▶ Seymour '94 [Higgs $\rightarrow WW \rightarrow \nu\ell$ jets]
- ▶ Butterworth, Cox & Forshaw '02 [$WW \rightarrow WW \rightarrow \nu\ell$ jets]
- ▶ Butterworth, Ellis & Raklev '07 [SUSY decay chains $\rightarrow W, H$]
- ▶ Skiba & Tucker-Smith '07 [vector quarks]
- ▶ Contino & Servant '08 [top partners]
- ▶ ...

Illustrate LHC challenges with a recently widely discussed class of problems:

Can you identify hadronically decaying EW bosons when they're produced at high p_t ?

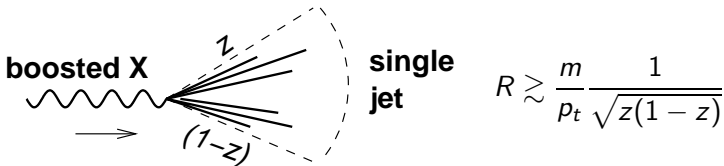


Significant discussion over years: **heavy new things decay to EW states**

- ▶ Seymour '94 [Higgs $\rightarrow WW \rightarrow \nu\ell$ jets]
- ▶ Butterworth, Cox & Forshaw '02 [$WW \rightarrow WW \rightarrow \nu\ell$ jets]
- ▶ Butterworth, Ellis & Raklev '07 [SUSY decay chains $\rightarrow W, H$]
- ▶ Skiba & Tucker-Smith '07 [vector quarks]
- ▶ Contino & Servant '08 [top partners]
- ▶ ...

Illustrate LHC challenges with a recently widely discussed class of problems:

Can you identify hadronically decaying EW bosons when they're produced at high p_t ?



Significant discussion over years: **heavy new things decay to EW states**

- ▶ Seymour '94 [Higgs $\rightarrow WW \rightarrow \nu\ell$ jets]
- ▶ Butterworth, Cox & Forshaw '02 [$WW \rightarrow WW \rightarrow \nu\ell$ jets]
- ▶ Butterworth, Ellis & Raklev '07 [SUSY decay chains $\rightarrow W, H$]
- ▶ Skiba & Tucker-Smith '07 [vector quarks]
- ▶ Contino & Servant '08 [top partners]
- ▶ ...

Most obvious method: look at the **jet mass**, but

- ▶ QCD jets can be massive too → large backgrounds
- ▶ Non-pert mass resolⁿ $\sim \delta M \sim R^4 \Lambda_{UE} \frac{p_t}{M}$ Dasgupta, Magnea & GPS '07

Natural idea: use hierarchical structure of k_t alg to resolve structure

Seymour '93; Butterworth, Cox & Forshaw '02 [Ysplitter]

- ▶ You can cut on d_{ij} (rel. \perp mom.²), correl. with mass helps reject bkgds
- ▶ But not ideal: k_t intrinsic mass resolution often poor

What you really want:

- ▶ Stay with hierarchical-type alg: study two subsets
 - ▶ Dynamically choose R based on p_t & M → best mass resolution
- **Cambridge/Aachen** algorithm

Repeatedly cluster pair of objects closest in angle until all separated by $\geq R$

[Can then undo clustering & look at jet on a range of angular scales]

Most obvious method: look at the **jet mass**, but

- ▶ QCD jets can be massive too → large backgrounds
- ▶ Non-pert mass resolⁿ $\sim \delta M \sim R^4 \Lambda_{UE} \frac{p_t}{M}$ Dasgupta, Magnea & GPS '07

Natural idea: **use hierarchical structure of k_t alg to resolve structure**

Seymour '93; Butterworth, Cox & Forshaw '02 [Ysplitter]

- ▶ You can cut on d_{ij} (rel. \perp mom.²), correl. with mass helps reject bkgds
- ▶ But not ideal: k_t intrinsic mass resolution often poor

What you really want:

- ▶ Stay with hierarchical-type alg: study two subsets
 - ▶ Dynamically choose R based on p_t & $M \rightarrow$ best mass resolution
- **Cambridge/Aachen** algorithm

Repeatedly cluster pair of objects closest in angle until all separated by $\geq R$
 [Can then undo clustering & look at jet on a range of angular scales]

Most obvious method: look at the **jet mass**, but

- ▶ QCD jets can be massive too → large backgrounds
- ▶ Non-pert mass resolⁿ $\sim \delta M \sim R^4 \Lambda_{UE} \frac{p_t}{M}$ Dasgupta, Magnea & GPS '07

Natural idea: **use hierarchical structure of k_t alg to resolve structure**

Seymour '93; Butterworth, Cox & Forshaw '02 [Ysplitter]

- ▶ You can cut on d_{ij} (rel. \perp mom.²), correl. with mass helps reject bkgds
- ▶ But not ideal: k_t intrinsic mass resolution often poor

What you really want:

- ▶ Stay with hierarchical-type alg: study two subsets
 - ▶ Dynamically choose R based on p_t & $M \rightarrow$ best mass resolution
- **Cambridge/Aachen** algorithm

Repeatedly cluster pair of objects closest in angle until all separated by $\geq R$
 [Can then undo clustering & look at jet on a range of angular scales]

E.g.: WH/ZH search channel @ LHC

- ▶ Signal is $W \rightarrow \ell\nu$, $H \rightarrow b\bar{b}$.
- ▶ Backgrounds include $Wb\bar{b}$, $t\bar{t} \rightarrow \ell\nu b\bar{b}jj$, ...

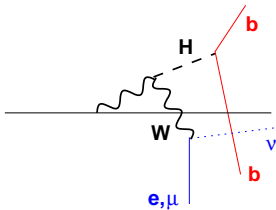
Studied e.g. in ATLAS TDR

Difficulties, e.g.

- ▶ $gg \rightarrow t\bar{t}$ has $\ell\nu b\bar{b}$ with **same intrinsic mass scale**, but much higher partonic luminosity
- ▶ Need exquisite control of bkgd shape

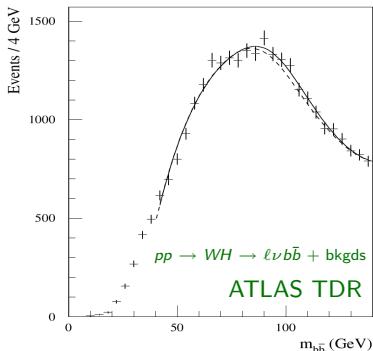
Try a long shot?

- ▶ Go to high p_t ($p_{tH}, p_{tV} > 200$ GeV)
- ▶ Lose 95% of signal, but more efficient?
- ▶ Maybe kill $t\bar{t}$ & gain clarity?



- ▶ Signal is $W \rightarrow \ell\nu$, $H \rightarrow b\bar{b}$.
- ▶ Backgrounds include $Wb\bar{b}$, $t\bar{t} \rightarrow \ell\nu b\bar{b}jj$, ...

Studied e.g. in ATLAS TDR

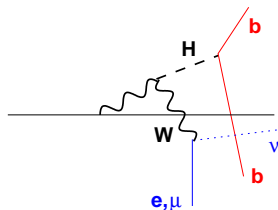


Difficulties, e.g.

- ▶ $gg \rightarrow t\bar{t}$ has $\ell\nu b\bar{b}$ with **same intrinsic mass scale**, but much higher partonic luminosity
- ▶ Need exquisite control of bkgd shape

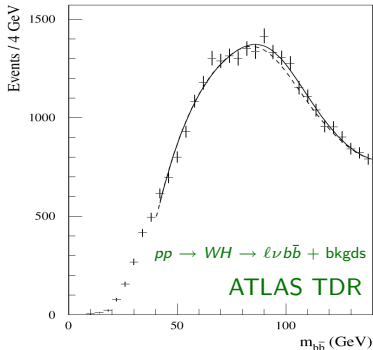
Try a long shot?

- ▶ Go to high p_t ($p_{tH}, p_{tV} > 200$ GeV)
- ▶ Lose 95% of signal, but more efficient?
- ▶ Maybe kill $t\bar{t}$ & gain clarity?



- ▶ Signal is $W \rightarrow \ell\nu$, $H \rightarrow b\bar{b}$.
- ▶ Backgrounds include $Wb\bar{b}$, $t\bar{t} \rightarrow \ell\nu b\bar{b}jj$, ...

Studied e.g. in ATLAS TDR

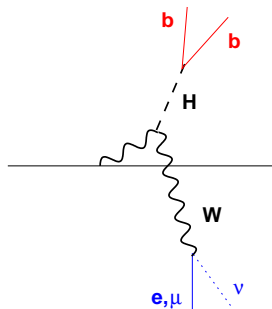


Difficulties, e.g.

- ▶ $gg \rightarrow t\bar{t}$ has $\ell\nu b\bar{b}$ with **same intrinsic mass scale**, but much higher partonic luminosity
- ▶ Need exquisite control of bkgd shape

Try a long shot?

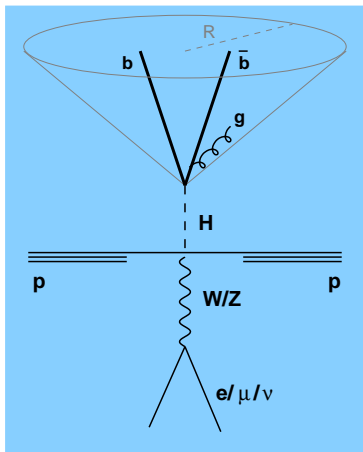
- ▶ Go to high p_t ($p_{tH}, p_{tV} > 200$ GeV)
- ▶ Lose 95% of signal, but more efficient?
- ▶ Maybe kill $t\bar{t}$ & gain clarity?



Searching for high- p_t HW/HZ?

High- p_t light Higgs decays to $b\bar{b}$ inside a single jet. Can this be seen?

Butterworth, Davison, Rubin & GPS '08



Cluster with Cambridge/Aachen

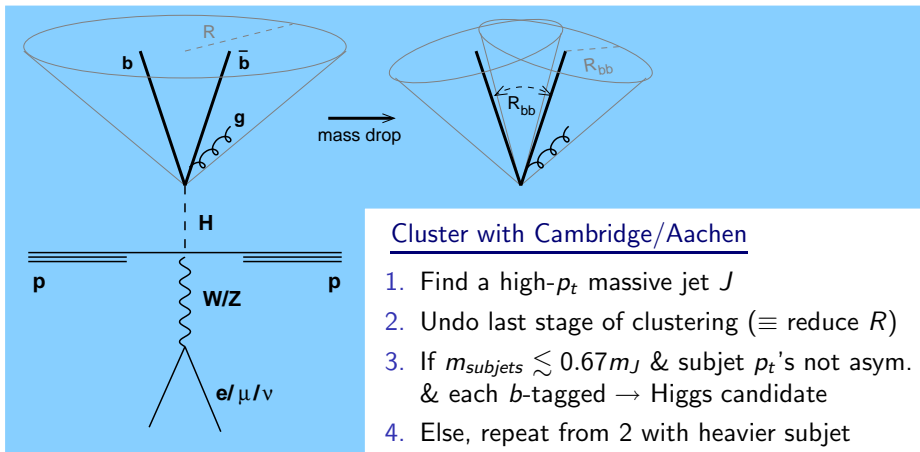
1. Find a high- p_t massive jet J
2. Undo last stage of clustering (\equiv reduce R)
3. If $m_{\text{subjects}} \lesssim 0.67m_J$ & subjet p_t 's not asym. & each b -tagged \rightarrow Higgs candidate
4. Else, repeat from 2 with heavier subjet

Then on the Higgs-candidate: *filter* away UE/pileup by reducing $R \rightarrow R_{\text{filt}}$, take *three hardest subjects* (keep LO gluon radⁿ) + require b -tags on two hardest.

Searching for high- p_t HW/HZ?

High- p_t light Higgs decays to $b\bar{b}$ inside a single jet. Can this be seen?

Butterworth, Davison, Rubin & GPS '08

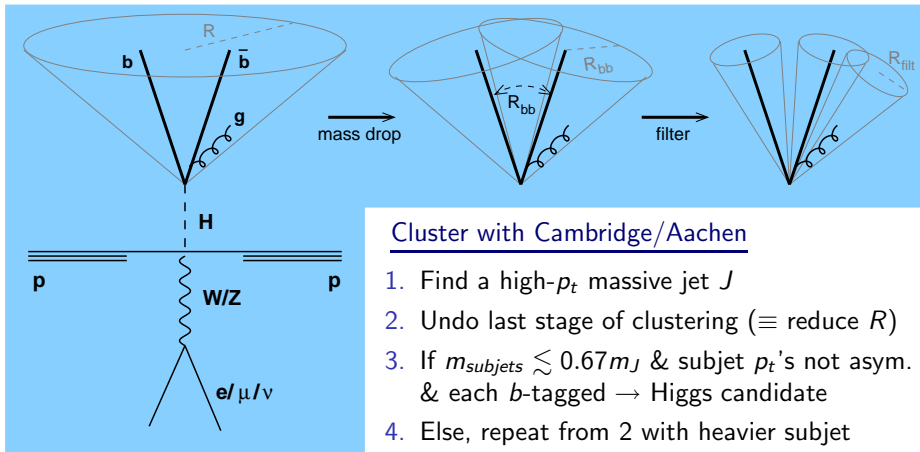


Then on the Higgs-candidate: *filter* away UE/pileup by reducing $R \rightarrow R_{filt}$, take *three hardest subjects* (keep LO gluon radⁿ) + require b -tags on two hardest.

Searching for high- p_t HW/HZ?

High- p_t light Higgs decays to $b\bar{b}$ inside a single jet. Can this be seen?

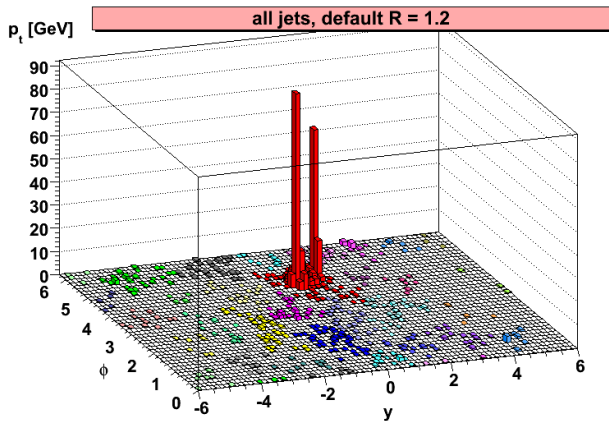
Butterworth, Davison, Rubin & GPS '08



Then on the Higgs-candidate: *filter* away UE/pileup by reducing $R \rightarrow R_{filt}$, take *three hardest subjets* (keep LO gluon radⁿ) + require b -tags on two hardest.

$$pp \rightarrow ZH \rightarrow \nu\bar{\nu}b\bar{b}, @14 \text{ TeV}, m_H = 115 \text{ GeV}$$

SIGNAL

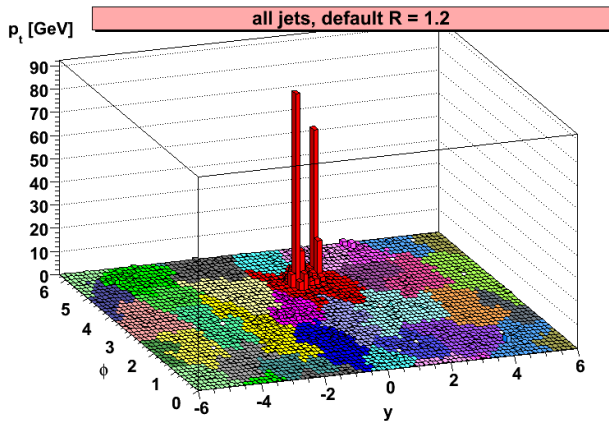


Zbb BACKGROUND

arbitrary norm.

$$pp \rightarrow ZH \rightarrow \nu\bar{\nu}b\bar{b}, @14 \text{ TeV}, m_H = 115 \text{ GeV}$$

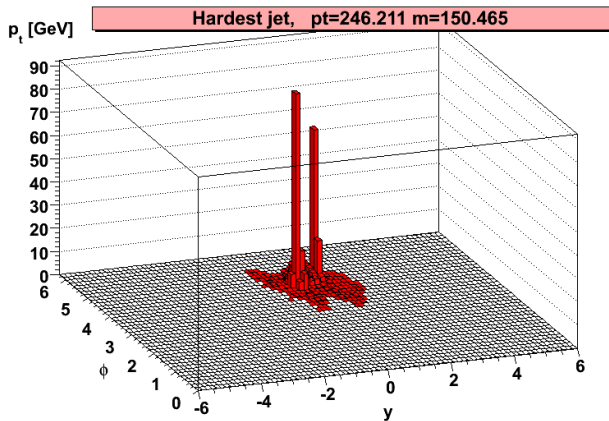
SIGNAL



Zbb BACKGROUND

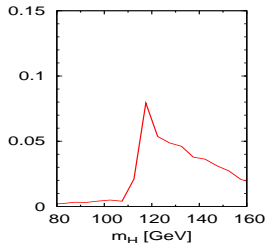
arbitrary norm.

$$pp \rightarrow ZH \rightarrow \nu\bar{\nu}b\bar{b}, @14\text{ TeV}, m_H = 115\text{ GeV}$$



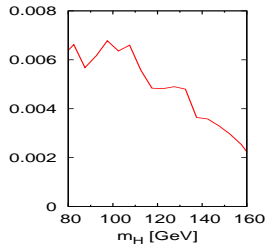
SIGNAL

$200 < p_{tZ} < 250$ GeV



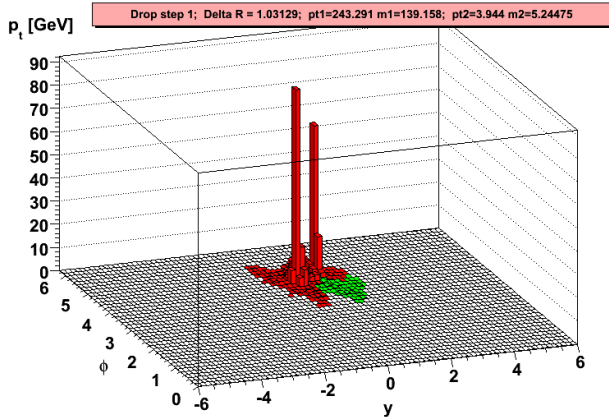
Zbb BACKGROUND

$200 < p_{tZ} < 250$ GeV



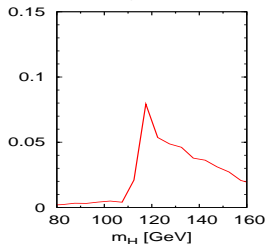
arbitrary norm.

$$pp \rightarrow ZH \rightarrow \nu\bar{\nu}b\bar{b}, @14\text{ TeV}, m_H = 115\text{ GeV}$$



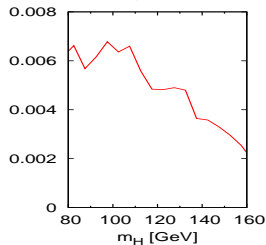
SIGNAL

$200 < p_{tZ} < 250\text{ GeV}$



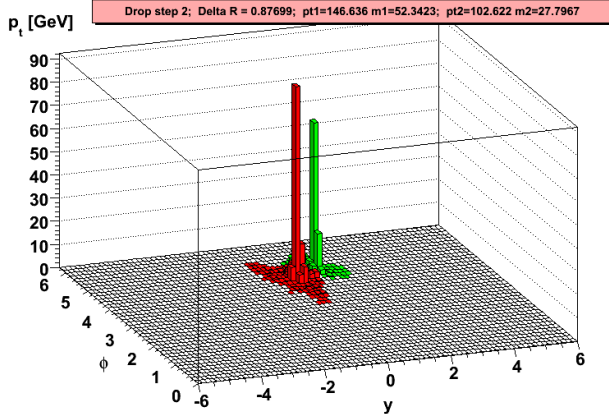
Zbb BACKGROUND

$200 < p_{tZ} < 250\text{ GeV}$



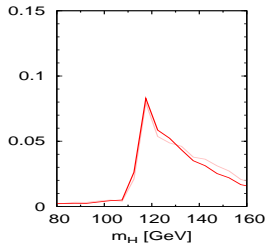
arbitrary norm.

$$pp \rightarrow ZH \rightarrow \nu\bar{\nu}b\bar{b}, @14\text{ TeV}, m_H = 115\text{ GeV}$$



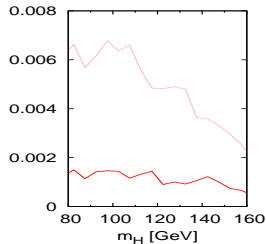
SIGNAL

$200 < p_{tZ} < 250\text{ GeV}$



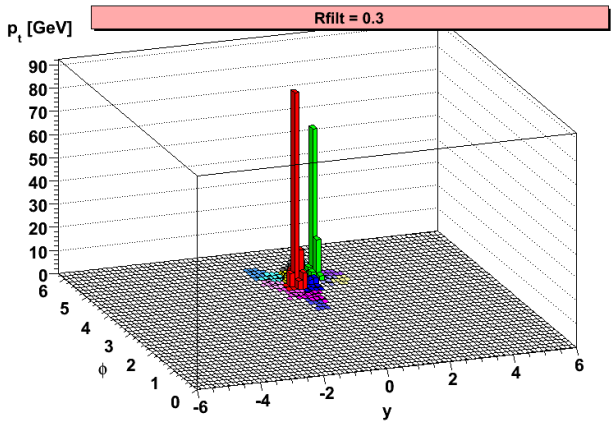
Zbb BACKGROUND

$200 < p_{tZ} < 250\text{ GeV}$



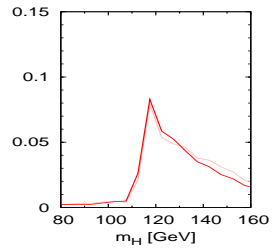
arbitrary norm.

$$pp \rightarrow ZH \rightarrow \nu\bar{\nu}b\bar{b}, @14\text{ TeV}, m_H = 115\text{ GeV}$$



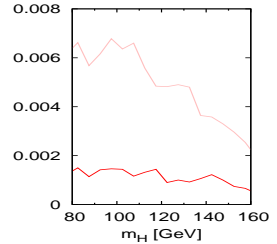
SIGNAL

$200 < p_{tZ} < 250$ GeV



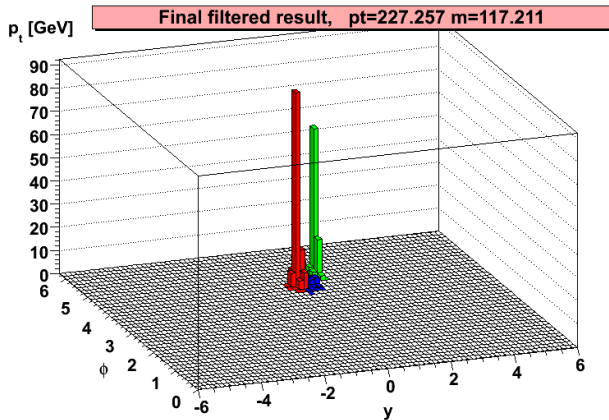
Zbb BACKGROUND

$200 < p_{tZ} < 250$ GeV



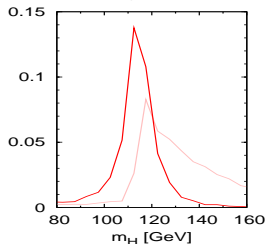
arbitrary norm.

$$pp \rightarrow ZH \rightarrow \nu\bar{\nu}b\bar{b}, @14 \text{ TeV}, m_H = 115 \text{ GeV}$$



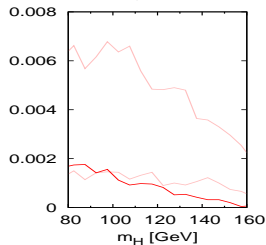
SIGNAL

$200 < p_{tZ} < 250 \text{ GeV}$



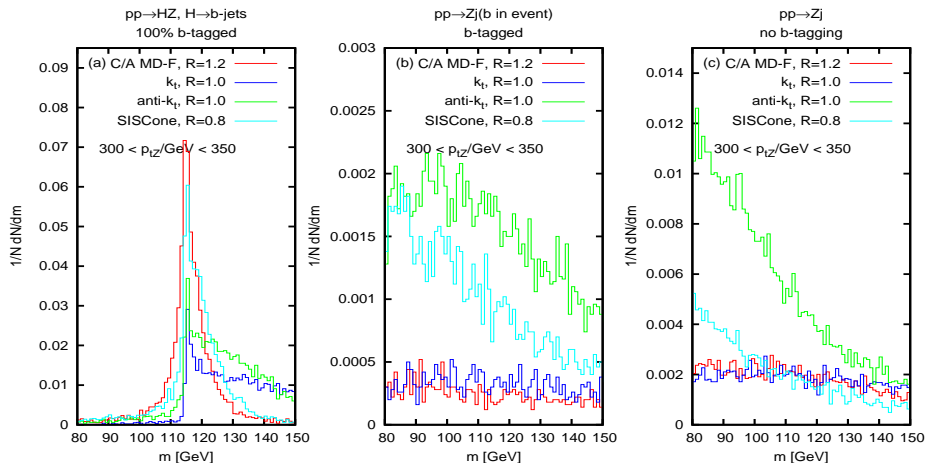
Zbb BACKGROUND

$200 < p_{tZ} < 250 \text{ GeV}$



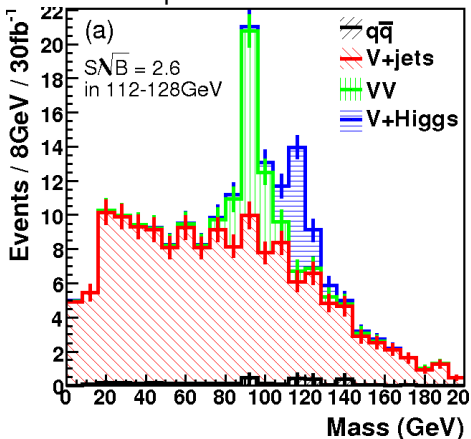
arbitrary norm.

Check mass spectra in HZ channel, $H \rightarrow b\bar{b}$, $Z \rightarrow \ell^+\ell^-$



Cambridge/Aachen (C/A) with mass-drop and filtering (MD/F) works best

Leptonic channel



Common cuts

- ▶ $p_{tV}, p_{tH} > 200$ GeV
- ▶ $|\eta_H| < 2.5$
- ▶ $[p_{t,\ell} > 30$ GeV, $|\eta_\ell| < 2.5]$
- ▶ No extra ℓ , b 's with $|\eta| < 2.5$
- ▶ Real/fake b -tag rates: 0.7/0.01
- ▶ S/\sqrt{B} from 18 GeV window

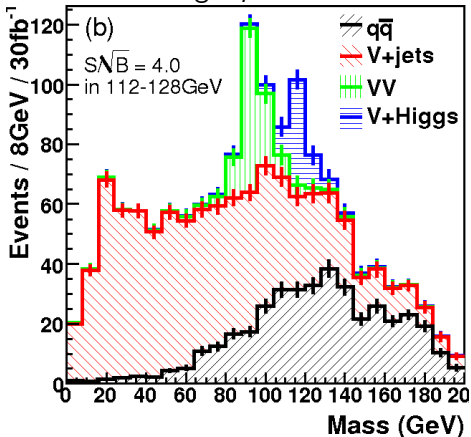
Leptonic channel

$$Z \rightarrow \mu^+\mu^-, e^+e^-$$

- ▶ $80 < m_{\ell^+\ell^-} < 100$ GeV

At 5.9σ for 30 fb^{-1} for $m_H = 115$ GeV this looks like a possible new channel for light Higgs discovery. **Deserves serious exp. study!**

Missing E_T channel



Common cuts

- ▶ $p_{tV}, p_{tH} > 200$ GeV
- ▶ $|\eta_H| < 2.5$
- ▶ $[p_{t,\ell} > 30$ GeV, $|\eta_\ell| < 2.5]$
- ▶ No extra ℓ , b 's with $|\eta| < 2.5$
- ▶ Real/fake b -tag rates: 0.7/0.01
- ▶ S/\sqrt{B} from 18 GeV window

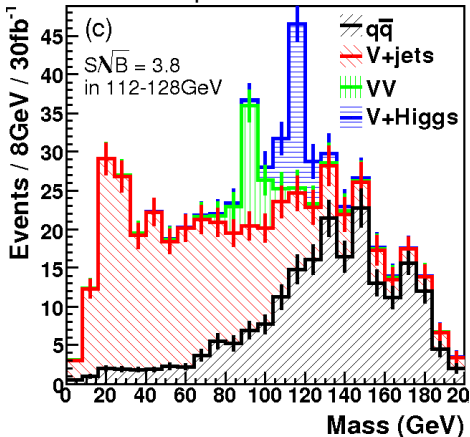
Missing- E_t channel

$$Z \rightarrow \nu\bar{\nu}, W \rightarrow \nu[\ell]$$

- ▶ $\cancel{E}_T > 200$ GeV

At 5.9σ for 30 fb^{-1} for $m_H = 115$ GeV this looks like a possible new channel for light Higgs discovery. **Deserves serious exp. study!**

Semi-leptonic channel



Common cuts

- ▶ $p_{tV}, p_{tH} > 200$ GeV
- ▶ $|\eta_H| < 2.5$
- ▶ $[p_{t,\ell} > 30$ GeV, $|\eta_\ell| < 2.5]$
- ▶ No extra ℓ , b 's with $|\eta| < 2.5$
- ▶ Real/fake b -tag rates: 0.7/0.01
- ▶ S/\sqrt{B} from 18 GeV window

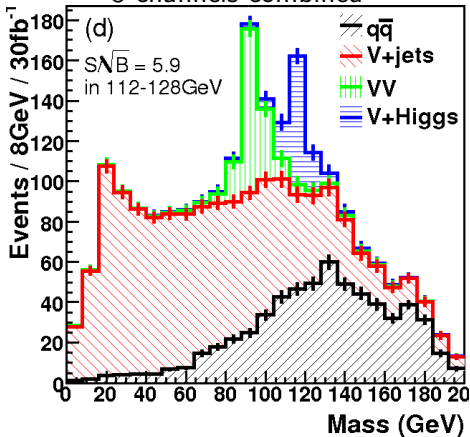
Semi-leptonic channel

$W \rightarrow \nu\ell$

- ▶ $\cancel{E}_T > 30$ GeV (& consistent W .)
- ▶ no extra jets $|\eta| < 3, p_t > 30$

At 5.9σ for 30 fb^{-1} for $m_H = 115$ GeV this looks like a possible new channel for light Higgs discovery. **Deserves serious exp. study!**

3 channels combined



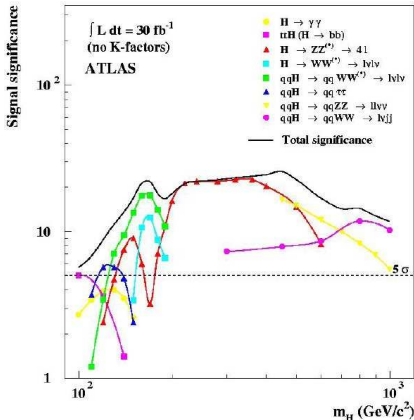
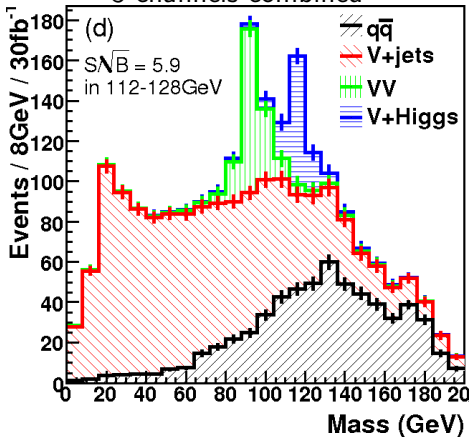
Common cuts

- ▶ $p_{tV}, p_{tH} > 200$ GeV
- ▶ $|\eta_H| < 2.5$
- ▶ $[p_{t,\ell} > 30$ GeV, $|\eta_\ell| < 2.5]$
- ▶ No extra ℓ, b 's with $|\eta| < 2.5$
- ▶ Real/fake b -tag rates: 0.7/0.01
- ▶ S/\sqrt{B} from 18 GeV window

3 channels combined

At 5.9σ for 30 fb^{-1} for $m_H = 115$ GeV this looks like a possible new channel for light Higgs discovery. **Deserves serious exp. study!**

3 channels combined



At 5.9σ for 30 fb^{-1} for $m_H = 115$ GeV this looks like a possible new channel for light Higgs discovery. **Deserves serious exp. study!**

Closing

- ▶ Jets are the closest we can get to seeing and giving meaning to partons
- ▶ Play a pivotal role in experimental analyses, comparisons to QCD calculations
- ▶ Significant progress in past 2 years towards making them *consistent* (IR/Collinear safe) and *practical* Link with computational geometry
All tools are made public:
<http://www.lpthe.jussieu.fr/~salam/fastjet/>
- ▶ The physics of how jets behave in a hadron-collider environment is a rich subject — much to be understood, and potential for significant impact in how jets are used at LHC E.g. Boosted higgs search

EXTRA SLIDES

Real life does not have infinities, but pert. infinity leaves a real-life trace

$$\alpha_s^2 + \alpha_s^3 + \alpha_s^4 \times \infty \rightarrow \alpha_s^2 + \alpha_s^3 + \alpha_s^4 \times \ln p_t/\Lambda \rightarrow \alpha_s^2 + \underbrace{\alpha_s^3 + \alpha_s^3}_{\text{BOTH WASTED}}$$

Among consequences of IR unsafety:

	<i>Last meaningful order</i>			Known at
	JetClu, ATLAS cone [IC-SM]	MidPoint [IC _{mp} -SM]	CMS it. cone [IC-PR]	
Inclusive jets	LO	NLO	NLO	NLO (→ NNLO)
W/Z + 1 jet	LO	NLO	NLO	NLO
3 jets	none	LO	LO	NLO [nlojet++]
W/Z + 2 jets	none	LO	LO	NLO [MCFM]
m_{jet} in $2j + X$	none	none	none	LO

NB: \$30 – 50M investment in NLO

Multi-jet contexts much more sensitive: **ubiquitous at LHC**

And LHC will rely on QCD for background double-checks
extraction of cross sections, extraction of parameters

Real life does not have infinities, but pert. infinity leaves a real-life trace

$$\alpha_s^2 + \alpha_s^3 + \alpha_s^4 \times \infty \rightarrow \alpha_s^2 + \alpha_s^3 + \alpha_s^4 \times \ln p_t/\Lambda \rightarrow \alpha_s^2 + \underbrace{\alpha_s^3 + \alpha_s^3}_{\text{BOTH WASTED}}$$

Among consequences of IR unsafety:

	<i>Last meaningful order</i>			Known at
	JetClu, ATLAS cone [IC-SM]	MidPoint [IC _{mp} -SM]	CMS it. cone [IC-PR]	
Inclusive jets	LO	NLO	NLO	NLO (→ NNLO)
W/Z + 1 jet	LO	NLO	NLO	NLO
3 jets	none	LO	LO	NLO [nlojet++]
W/Z + 2 jets	none	LO	LO	NLO [MCFM]
m _{jet} in 2j + X	none	none	none	LO

NB: \$30 – 50M investment in NLO

Multi-jet contexts much more sensitive: **ubiquitous at LHC**

And LHC will rely on QCD for background double-checks
extraction of cross sections, extraction of parameters

Real life does not have infinities, but pert. infinity leaves a real-life trace

$$\alpha_s^2 + \alpha_s^3 + \alpha_s^4 \times \infty \rightarrow \alpha_s^2 + \alpha_s^3 + \alpha_s^4 \times \ln p_t/\Lambda \rightarrow \alpha_s^2 + \underbrace{\alpha_s^3 + \alpha_s^3}_{\text{BOTH WASTED}}$$

Among consequences of IR unsafety:

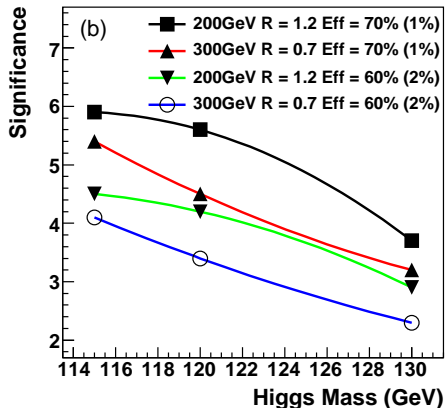
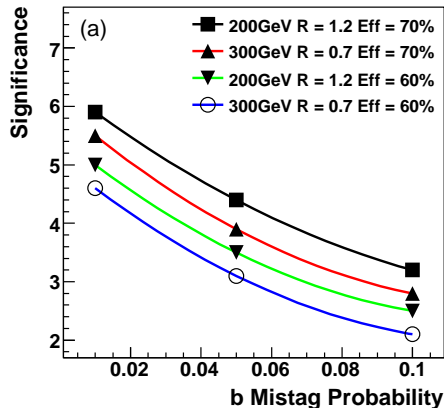
	<i>Last meaningful order</i>			Known at
	JetClu, ATLAS cone [IC-SM]	MidPoint [IC _{mp} -SM]	CMS it. cone [IC-PR]	
Inclusive jets	LO	NLO	NLO	NLO (→ NNLO)
W/Z + 1 jet	LO	NLO	NLO	NLO
3 jets	none	LO	LO	NLO [nlojet++]
W/Z + 2 jets	none	LO	LO	NLO [MCFM]
m _{jet} in 2j + X	none	none	none	LO

NB: \$30 – 50M investment in NLO

Multi-jet contexts much more sensitive: **ubiquitous at LHC**

And LHC will rely on QCD for background double-checks
extraction of cross sections, extraction of parameters

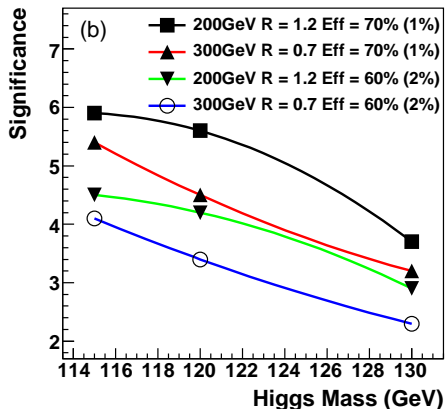
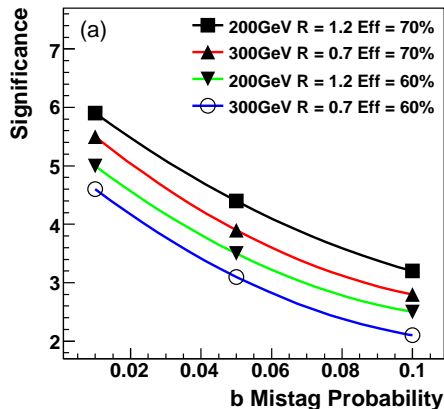
Impact of b -tagging, Higgs mass



Most scenarios above 3σ ; still much work to be done, notably on verification of experimental resolution.

Regardless of final outcome, illustrates value of choosing appropriate “jet-methods,” and of potential for progress with new ideas.

Impact of b -tagging, Higgs mass



Most scenarios above 3σ ; still much work to be done, notably on verification of experimental resolution.

Regardless of final outcome, illustrates value of choosing appropriate “jet-methods,” and of potential for progress with new ideas.