Towards Jetography

Gavin Salam

LPTHE, CNRS and UPMC (Univ. Paris 6)

NIKHEF Theory Group Seminar
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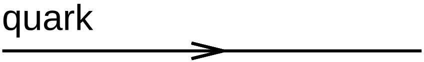
Based on work with
Jon Butterworth, Matteo Cacciari, Mrinal Dasgupta, Adam Davison,
Lorenzo Magnea, Juan Rojo, Mathieu Rubin & Gregory Soyez

Parton fragmentation



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$$\int \alpha_s \frac{dE}{E}\frac{d\theta}{\theta} \gg 1$$

At low scales:

$$\alpha_s \to 1$$

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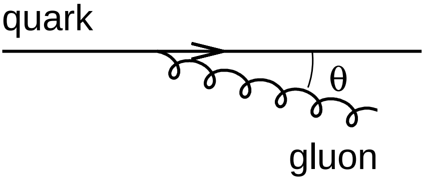

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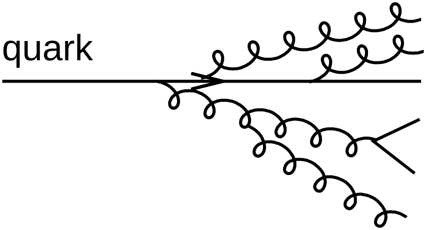

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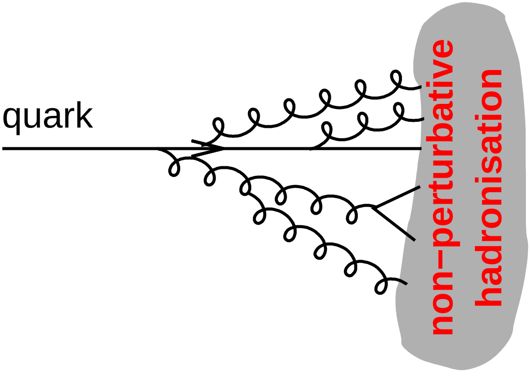

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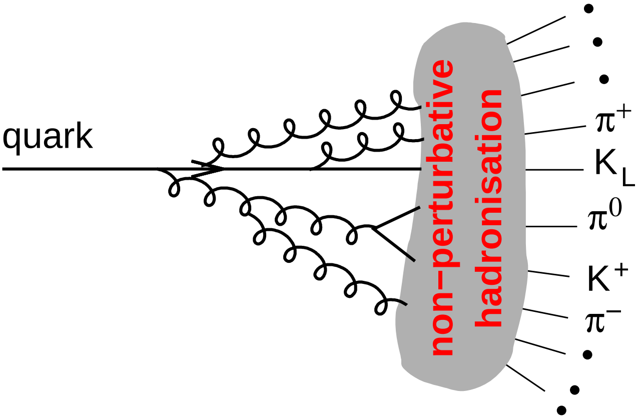

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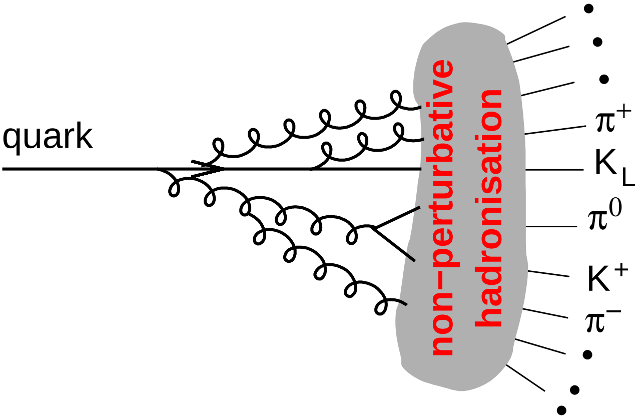

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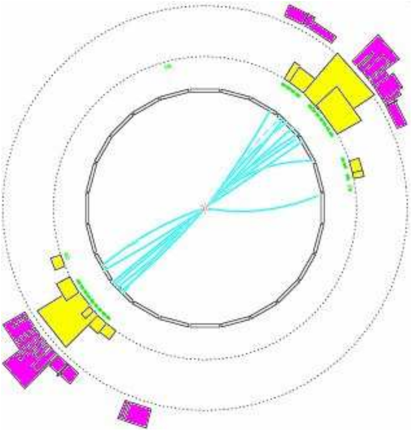
This is a jet

Seeing v. defining jets



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Clearly(?) 2 jets here

How many jets do you see?
Do you really want to ask yourself this question for 10^9 events?

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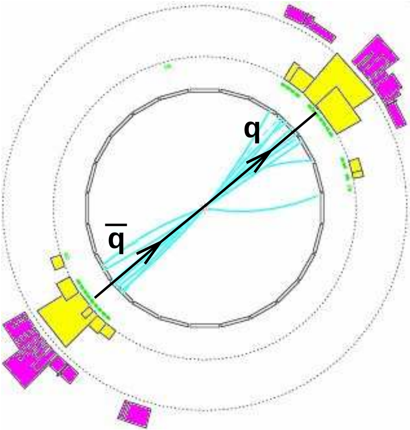

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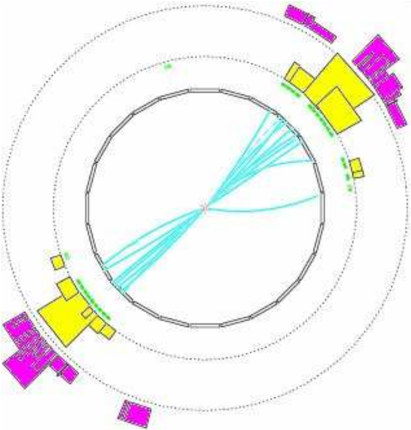

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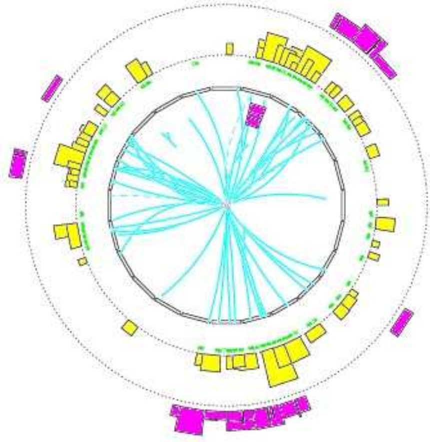

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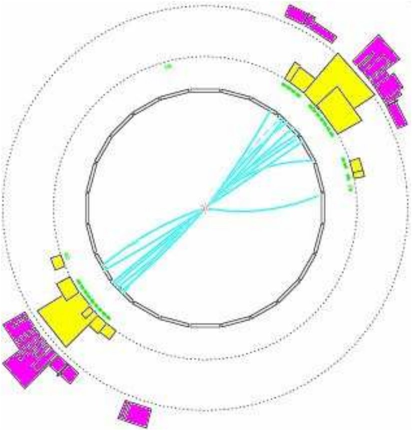

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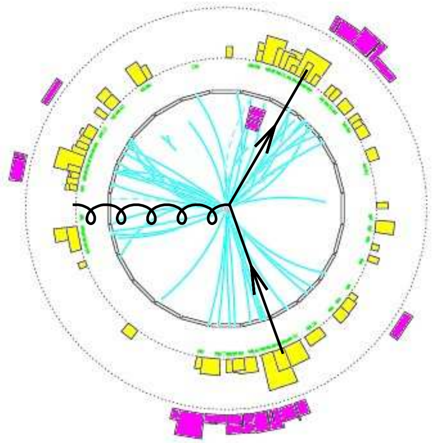

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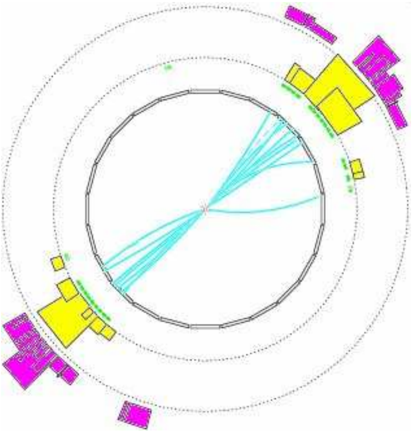

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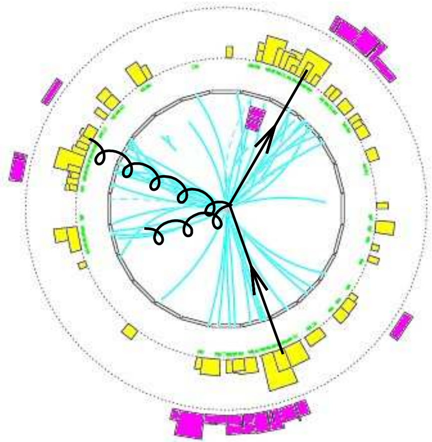

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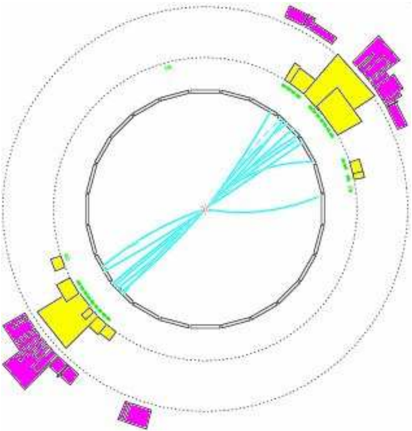

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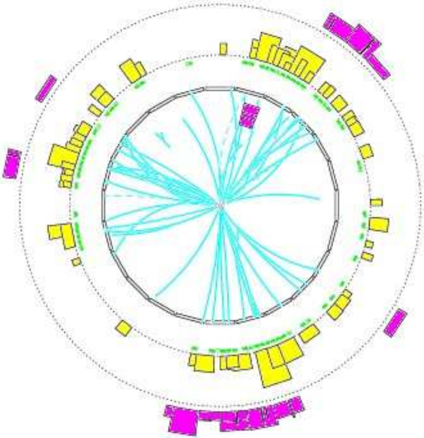

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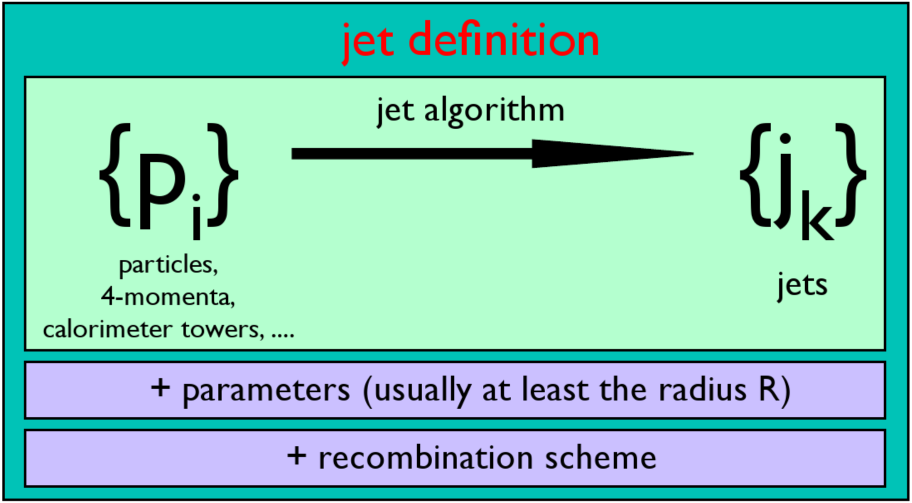
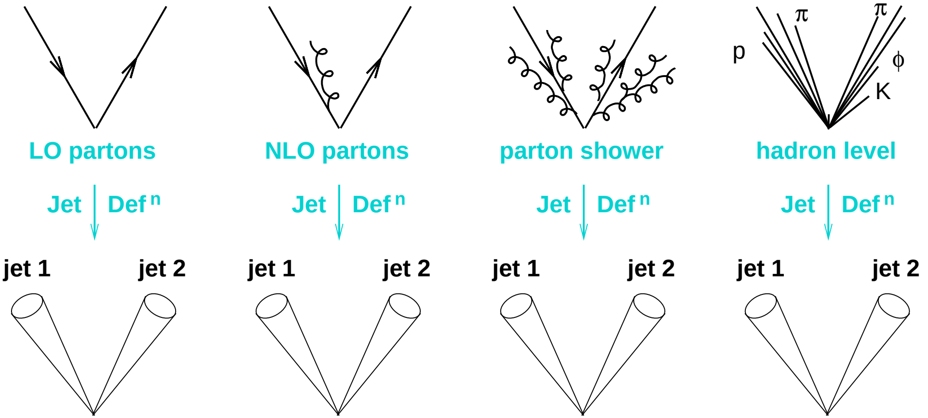
Les Houches 2007 proceedings, arXiv:0803.0678



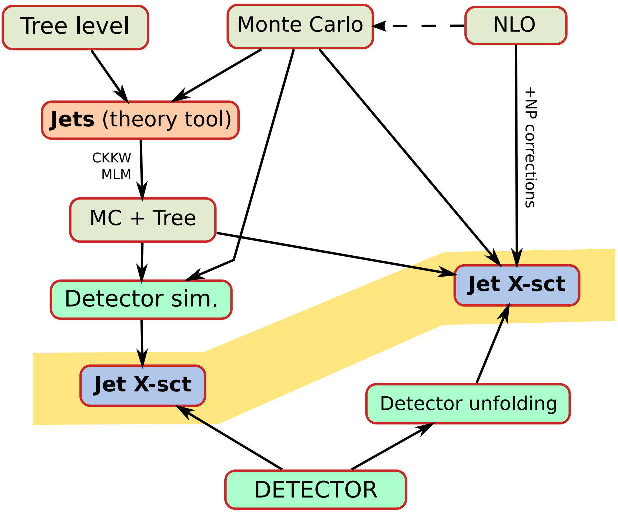
Reminder: running a jet definition gives a well defined physical observable, which we can measure and, hopefully, calculate

Jets as projections



Projection to jets should be resilient to QCD effects

QCD jets flowchart



Jet (definitions) provide central link between expt., "theory" and theory

And jets are an input to almost all analyses

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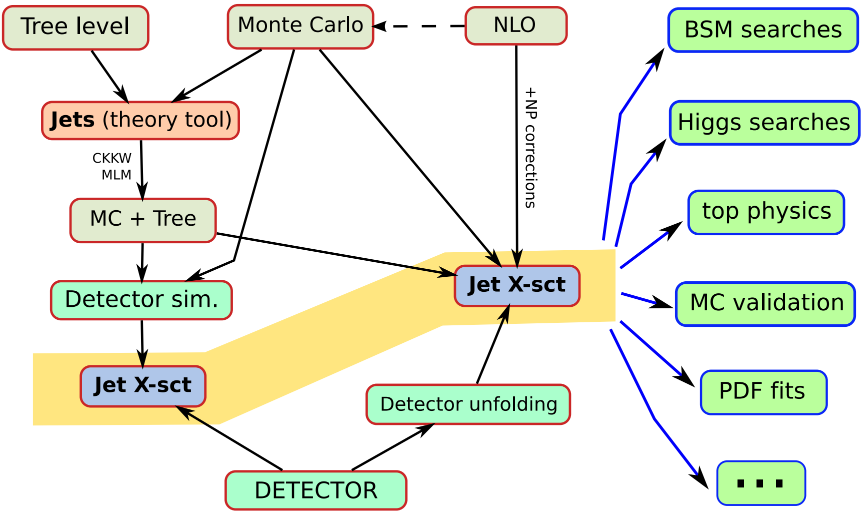

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What jet algorithms are out there?

sequential recombination (k_t)
& cone type

Two classes of jet algorithm

Sequential recombination	**Cone**
k_t, Jade, Cam/Aachen, ...	UA1, JetClu, Midpoint, ...
Bottom-up: Cluster 'closest' particles repeatedly until few left $\rightarrow$ jets.	Top-down: Find coarse regions of energy flow (cones), and call them jets.
Works because of mapping: *closeness* $\Leftrightarrow$ *QCD divergence*	Works because *QCD only modifies energy flow on small scales*
Loved by e^+e^-, ep and theorists	Loved by pp and few(er) theorists

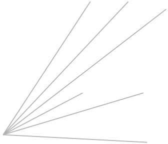
Sequential recombination algorithms

k_t algorithm — Catani, Dokshizter, Olsson, Seymour, Turnock, Webber '91–'93; Ellis, Soper '93

- Find smallest of all $d_{ij} = \min(k_{ti}^2, k_{tj}^2)\Delta R_{ij}^2/R^2$ and $d_{iB} = k_i^2$
- Recombine i, j
- Repeat

Bottom-up jets:

Sequential recombination



NB: hadron collider variables

- $\Delta R_{ij}^2 = (\phi_i - \phi_j)^2 + (y_i - y_j)^2$
- rapidity $y_i = \frac{1}{2} \ln \frac{E_i + p_{zi}}{E_i - p_{zi}}$
- ΔR_{ij} is boost invariant angle

R sets jet opening angle

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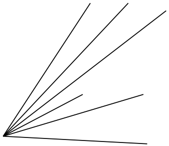

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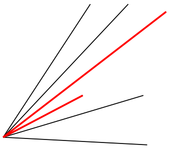

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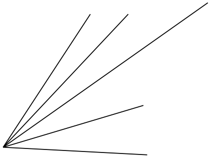

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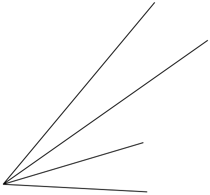

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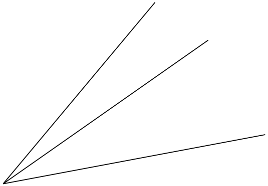

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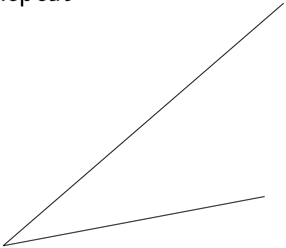

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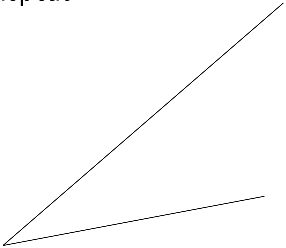

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Why k_t?

k_t distance measures

$$d_{ij} = \min(k_{ti}^2, k_{tj}^2)\Delta R_{ij}^2, \qquad d_{iB} = k_{ti}^2$$

are closely related to structure of divergences for QCD emissions

$$[dk_j]|M^2_{g\to g_i g_j}(k_j)| \sim \frac{\alpha_s C_A}{2\pi}\frac{dk_{tj}}{\min(k_{ti}, k_{tj})}\frac{d\Delta R_{ij}}{\Delta R_{ij}}, \qquad (k_{tj} \ll k_{ti}, \;\; \Delta R_{ij} \ll 1)$$

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k_t algorithm attempts approximate inversion of branching process

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Cones with Split Merge (SM)

Tevatron & ATLAS cone algs have two main steps:

- Find some/all stable cones

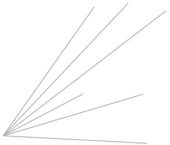
  ≡ cone pointing in same direction as the momentum of its contents

- Resolve cases of overlapping stable cones

  By running a 'split–merge' procedure

Top-down jets:

cone algorithms



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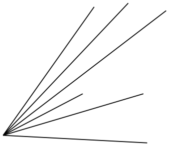

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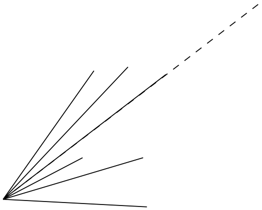

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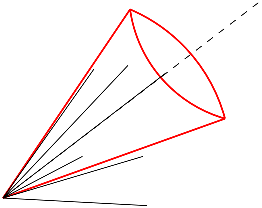

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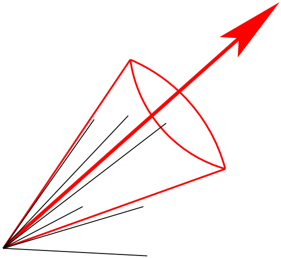

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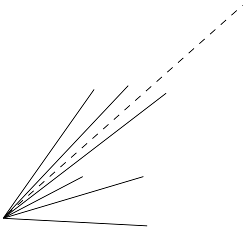

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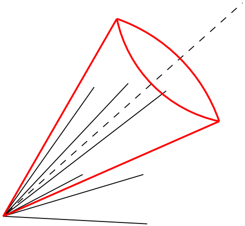

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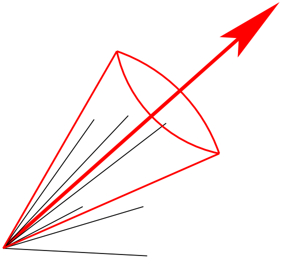

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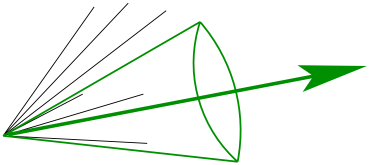

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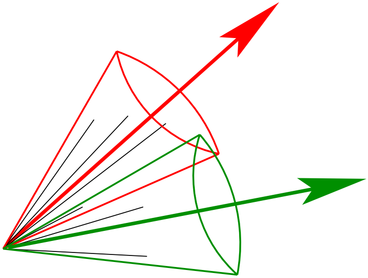

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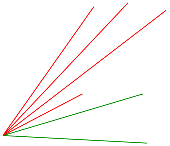

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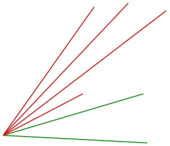

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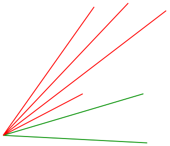

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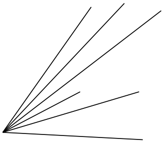

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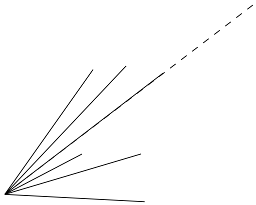

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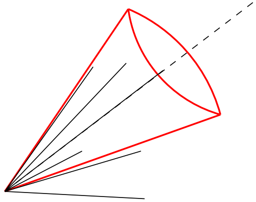

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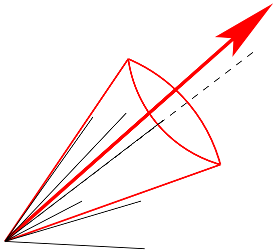

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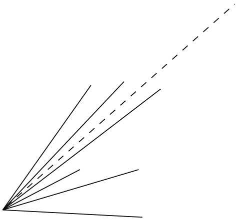

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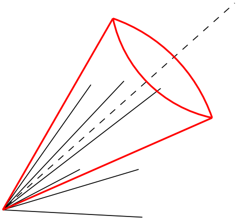

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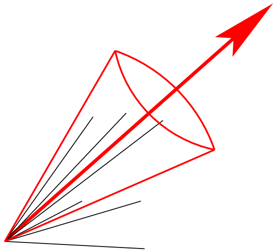

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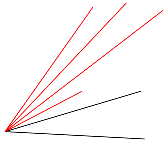

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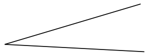

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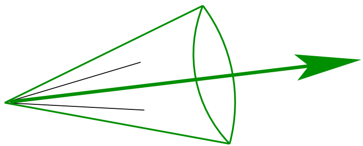

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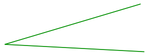

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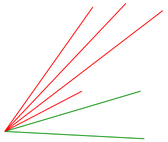

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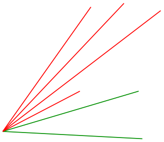

Iterative Cone with Progressive Removal (IC-PR)

e.g. CMS it. cone, [Pythia Cone, GetJet], ...

- NB: not same type of algorithm as Atlas Cone, MidPoint, SISCone

Readying jet "technology" for the LHC era

[a.k.a. satisfying Snowmass]

Snowmass accords

Snowmass Accord (1990):

FERMILAB-Conf-90/249-E
[E-741/CDF]

Toward a Standardization of Jet Definitions *

Several important properties that should be met by a jet definition are [3]:

1. Simple to implement in an experimental analysis;
2. Simple to implement in the theoretical calculation;
3. Defined at any order of perturbation theory;
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Property 1 $\Leftrightarrow$ speed. (+other aspects)

- LHC events may have up to $N = 4000$ particles (at high-lumi)
- Sequential recombination algs. (k_t) slow, $\sim N^3 \to 60s$ for $N = 4000$

$\mathbf{k_t}$ not practical for $\mathcal{O}\left(10^9\right)$ events

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Property 4 ≡ Infrared and Collinear (IRC) Safety. It helps ensure:

- Soft (low-energy) emissions & collinear splittings don't change jets
- Each order of perturbation theory is smaller than previous (at high p_t)

Wasn't satisfied by the cone algorithms

Computing and k_t

Snowmass issue #1

The k_t algorithm and its speed

<u>'Trivial' computational issue:</u>

- for N parti… searched through N times $\to N^3$
- 4000 parti…
  …2000 CPU years)
- Heavy Ions: 30000 particles: 10 hours/event

As far as possible physics choices should not be limited by computing.

Even if we're clever about repeating the full search each time, we still have $\mathcal{O}\left(N^2\right)$ d_{ij}'s to establish

Computing and k_t

'Trivial' computational issue:

- for N particles: N^2 d_{ij} searched through N times $= N^3$
- 4000 particles (or calo cells): 1 minute

NB: often study $10^7 - 10^9$ events (20-2000 CPU years)

- Heavy Ions: 30000 particles: 10 hours/event

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k_t and geometry

There are $N(N-1)/2$ distances d_{ij} — surely we have to calculate them all in order to find smallest?

k_t distance measure is partly *geometrical:*

$$\min_{i,j} d_{ij} \equiv \min_{i,j}\left(\min\{k_{ti}^2, k_{tj}^2\}\Delta R_{ij}^2\right)$$

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In words: for each i look only at the k_t distance to its 2D geometrical nearest neighbour (GNN).

k_t distance need only be calculated between GNNs

Each point has 1 GNN → need only calculate N d_{ij}'s

Cacciari & GPS, '05

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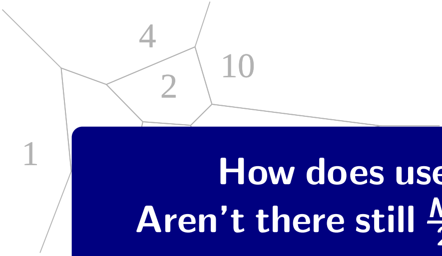

Given a set of vertices on plane (1...10) a *Voronoi diagram* partitions plane into cells containing all points closest to each vertex

...i '1908

...x is al-

...it is 3)

How does use of GNN help?
Aren't there still $\frac{N^2}{2}$ ΔR_{ij}^2 to check...?

Geometrical nearest neighbour finding is a classic problem in the field of Computational Geometry

Const... ...une '88

Updat...

Devillers '99 [+ related work by other authors]

Convenient C++ package available: CGAL, `http://www.cgal.org`

with help of CGAL, k_t clustering can be done in N ln N time
Coded in the FastJet package (v1), Cacciari & GPS '06

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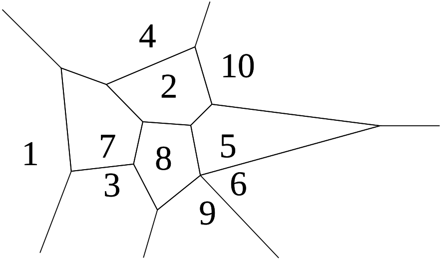

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Dirichlet '1850, Voronoi '1908

A vertex's nearest other vertex is always in an adjacent cell.

E.g. GNN of point 7 must be among 1,4,2,8,3 (it is 3)

Construction of Voronoi diagram for N points: $N \ln N$ time — Fortune '88

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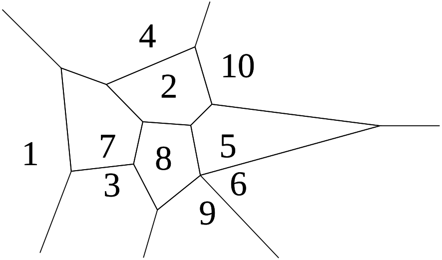

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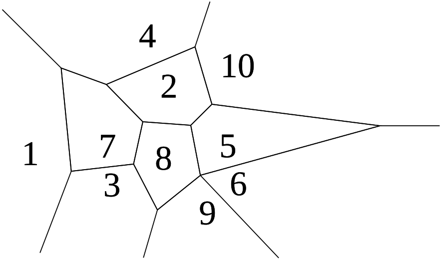

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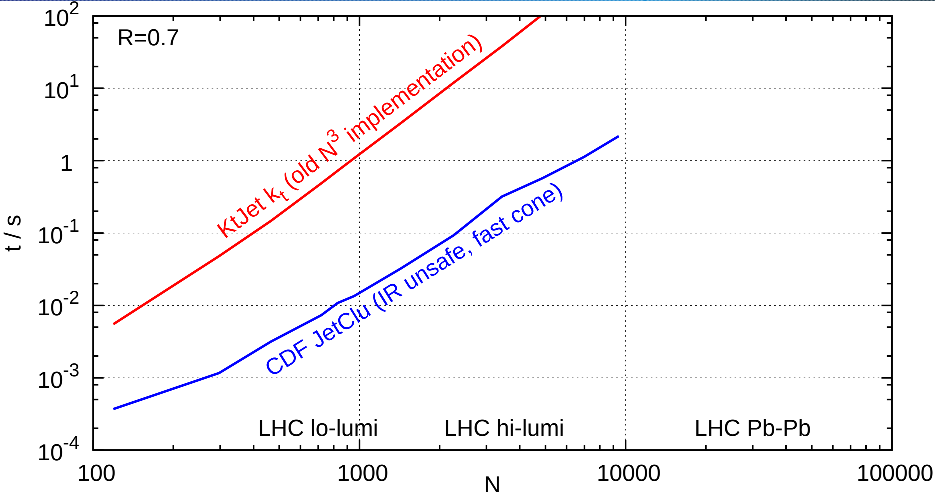

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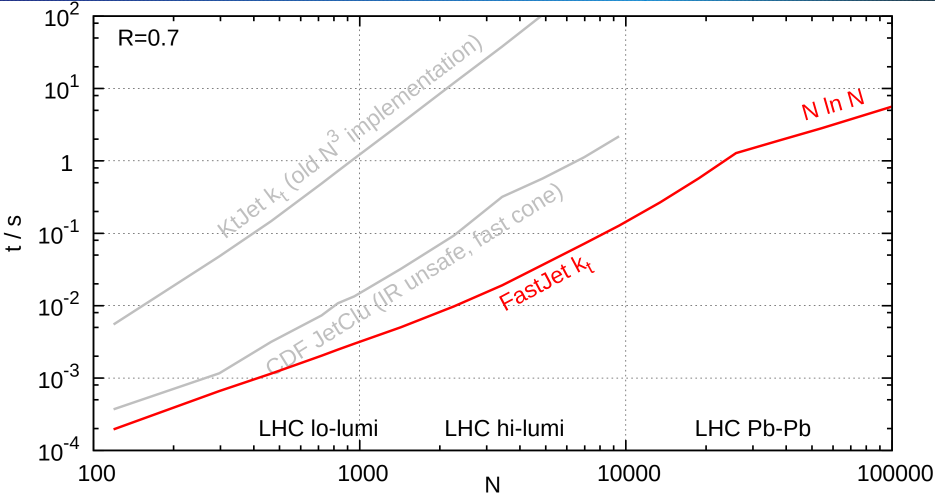
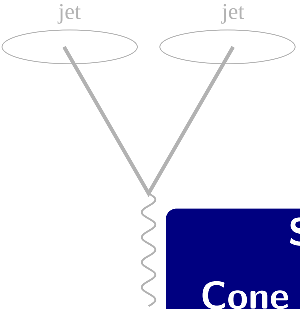


Factorisation of momentum & geometry

→ **2–3 orders of magnitude gain in speed!**

Speed competitive with fast cone algorithms

JetClu (& Atlas Cone) in Wjj @ NLO



Snowmass issue #4

Cone algorithms and IR safety

	$\alpha_s^2\alpha_{EW}$	$\alpha_s^3\alpha_{EW}$	$\alpha_s^3\alpha_{EW}$
1-jet			$+\infty$
2-jet	$\mathcal{O}(1)$	$-\infty$	0

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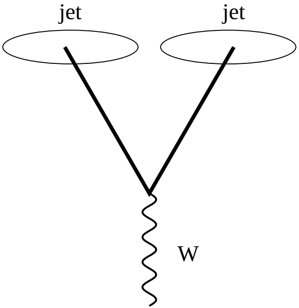

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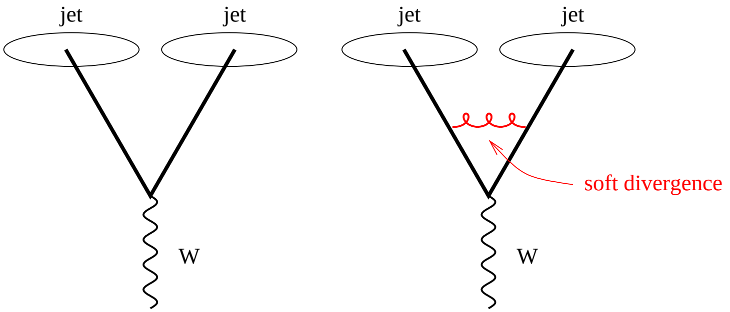

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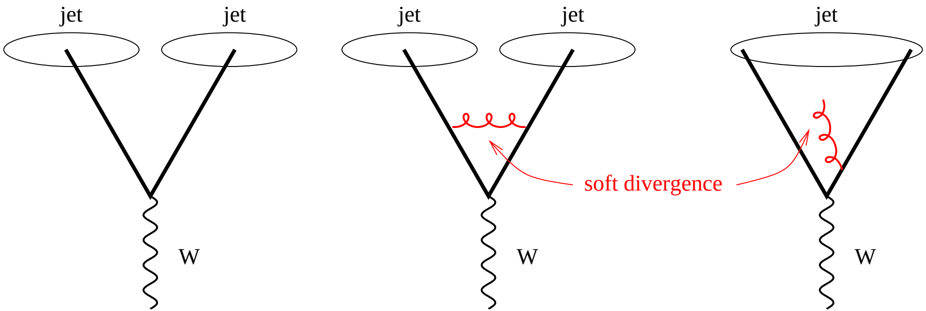

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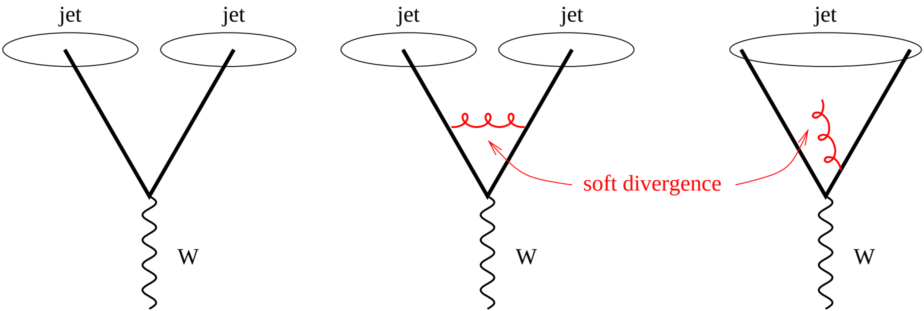

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IRC safety & real-life

Real life does not have infinities, but pert. infinity leaves a real-life trace

$$\alpha_s^2 + \alpha_s^3 + \alpha_s^4 \times \infty \to \alpha_s^2 + \alpha_s^3 + \alpha_s^4 \times \ln p_t/\Lambda \to \alpha_s^2 + \underbrace{\alpha_s^3 + \alpha_s^3}_{\text{BOTH WASTED}}$$

Among consequences of IR unsafety:

	Last meaningful order			
	JetClu, ATLAS cone [IC-SM]	MidPoint [IC$_{mp}$-SM]	CMS it. cone [IC-PR]	Known at
Inclusive jets	LO	NLO	NLO	NLO ($\to$ NNLO)
W/Z + 1 jet	LO	NLO	NLO	NLO
3 jets	**none**	LO	LO	NLO [nlojet++]
W/Z + 2 jets	**none**	LO	LO	NLO [MCFM]
m_{jet} in $2j + X$	**none**	**none**	**none**	LO

NB: 50,000,000\$/£/CHF/€ investment in NLO

Multi-jet contexts much more sensitive: **ubiquitous at LHC**

And LHC will rely on QCD for background double-checks

extraction of cross sections, extraction of parameters

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	Last meaningful order			
	JetClu, ATLAS cone [IC-SM]	MidPoint [IC$_{mp}$-SM]	CMS it. cone [IC-PR]	Known at
Inclusive jets	LO	NLO	NLO	NLO ($\rightarrow$ NNLO)
W/Z + 1 jet	LO	NLO	NLO	NLO
3 jets	**none**	LO	LO	NLO [nlojet++]
W/Z + 2 jets	**none**	LO	LO	NLO [MCFM]
m_{jet} in $2j + X$	**none**	**none**	**none**	LO

NB: 50,000,000\$/£/CHF/€ investment in NLO

Multi-jet contexts much more sensitive: **ubiquitous at LHC**

And LHC will rely on QCD for background double-checks

extraction of cross sections, extraction of parameters

Does lack of IRC safety matter?

I do searches, not QCD. Why should I care about IRC safety?

- Are you looking for a mass-peak? ➥ you needn't care much
- Are you looking for an excess over bkgd? ➥ you need control samples, validated against QCD

$\underbrace{\text{W+1,2,3 jets}}_{\text{NLO v. data}} \longleftrightarrow \underbrace{\text{W+n jets}}_{\text{LO, LO+MC v. data}} \longleftrightarrow \underbrace{\text{new-physics search}}_{\text{LO+MC v. data}}$

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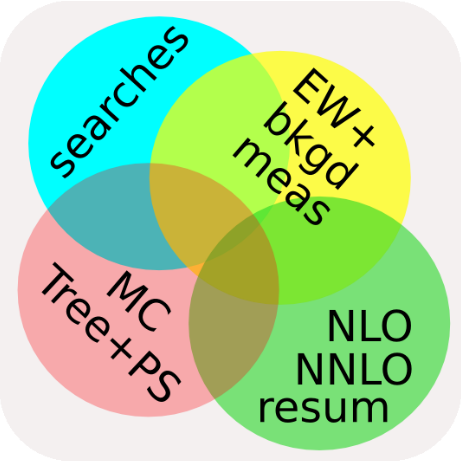

W+1,2,3 jets	⟷	W+n jets	⟷	new-physics search
NLO v. data		LO, LO+MC v. data		LO+MC v. data
IR safe alg.		IR safe alg.		IR safe alg.

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W+1,2,3 jets	$\longleftrightarrow$	W+n jets	$\not\longleftrightarrow$	new-physics search
NLO v. data		LO, LO+MC v. data		LO+MC v. data
IR safe alg.		IR safe alg.		IR **unsafe** alg.

Two directions

How do we solve cone IR safety problems?

Fix stable-cone finding

↓

SISCone

GPS & Soyez '07

Same family as Tev. Run II alg

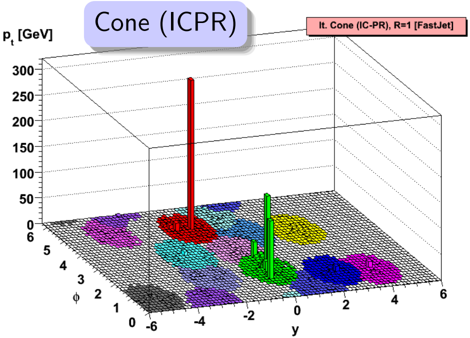
Invent "cone-like" alg.

↓

anti-kt

Cacciari, GPS & Soyez '08

Essential characteristic of cones?



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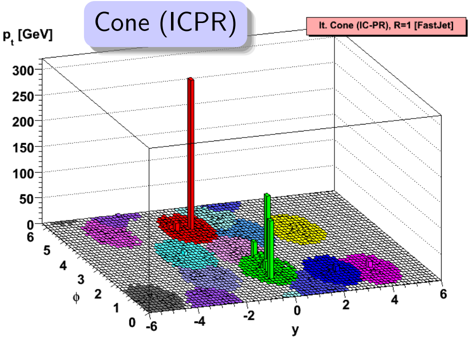
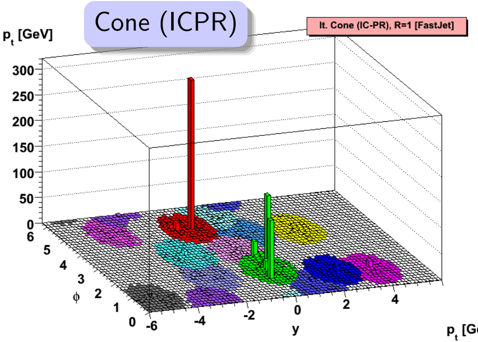


(Some) cone algorithms give **circular** jets in $y - \phi$ plane

Much appreciated by experiments e.g. for acceptance corrections

Essential characteristic of cones?



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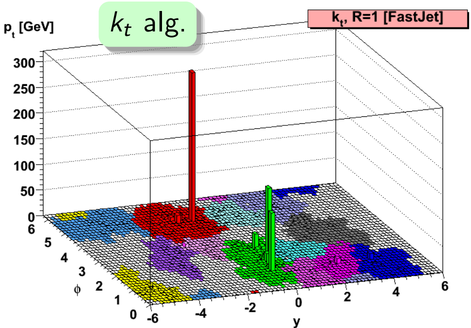
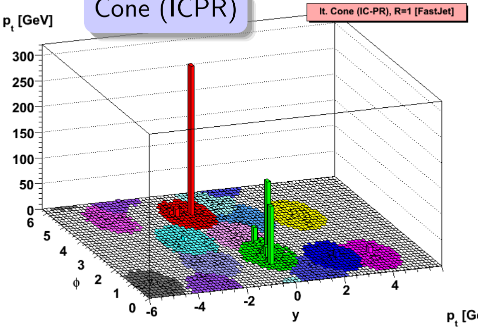


Essential characteristic of cones?

Cone (ICPR)



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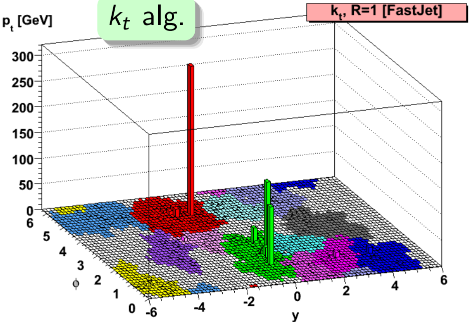
k_t alg.



k_t jets are **irregular**

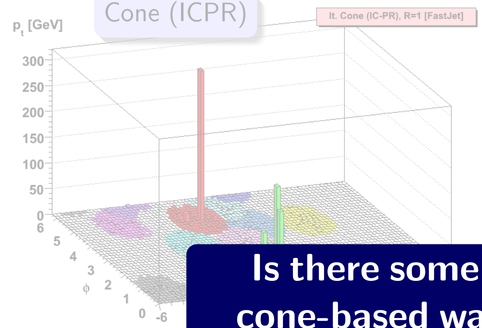
Because soft junk clusters together first:

$$d_{ij} = \min(k_{ti}^2, k_{tj}^2)\Delta R_{ij}^2$$

Regularly held against k_t

Essential characteristic of cones?

Cone (ICPR)



(Some) cone algorithms give **circular** jets in $y - \phi$ plane

Much appreciated by experiments e.g. for acceptance corrections

Is there some other, non cone-based way of getting circular jets?

k_t jets are

Because soft junk clusters together first:

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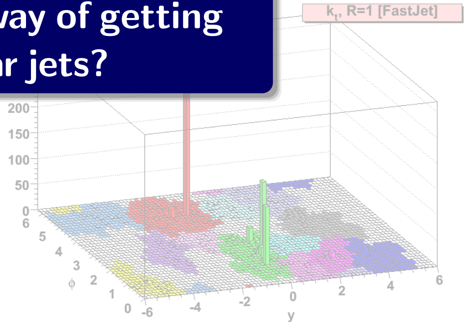
Regularly held against k_t



Adapting seq. rec. to give circular jets

Soft stuff clusters with nearest neighbour

$$k_t\text{: } d_{ij} = \min(k_{ti}^2, k_{tj}^2)\Delta R_{ij}^2 \longrightarrow \textbf{anti-k}_\mathbf{t}\text{: } d_{ij} = \frac{\Delta R_{ij}^2}{\max(k_{ti}^2, k_{tj}^2)}$$

Hard stuff clusters with nearest neighbour

Privilege collinear divergence over soft divergence

Adapting seq. rec. to give circular jets

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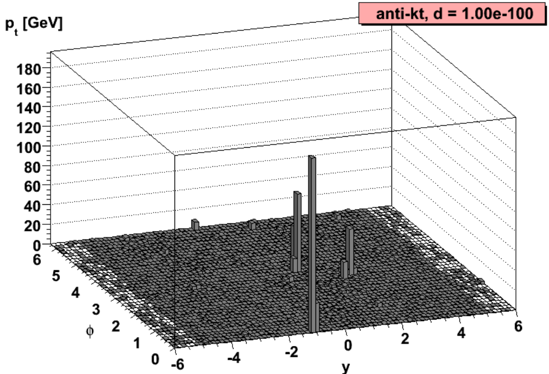
Hard stuff clusters with nearest neighbour

divergence over soft divergence



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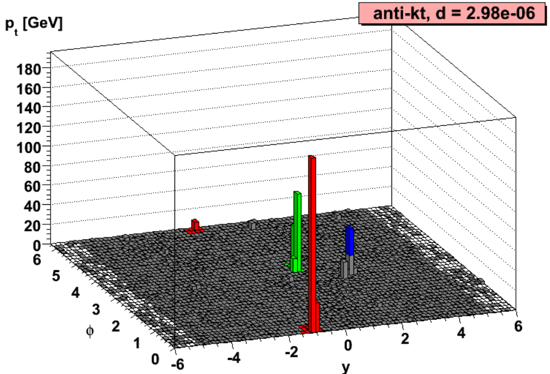

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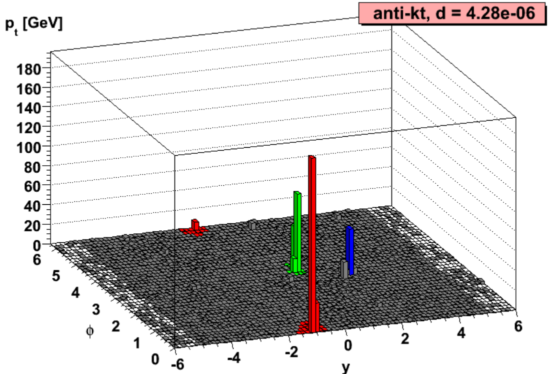

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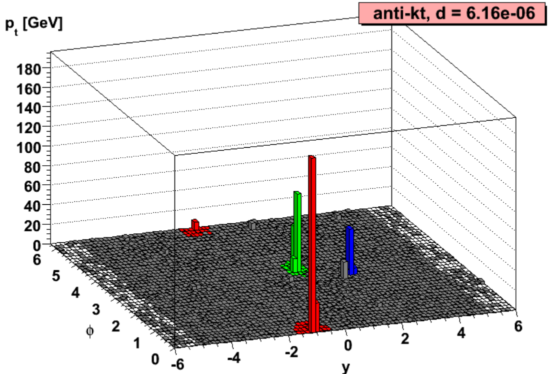

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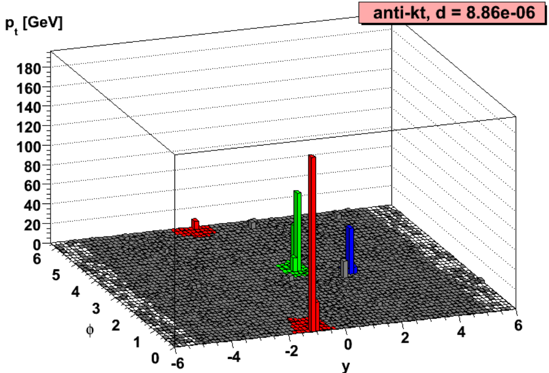

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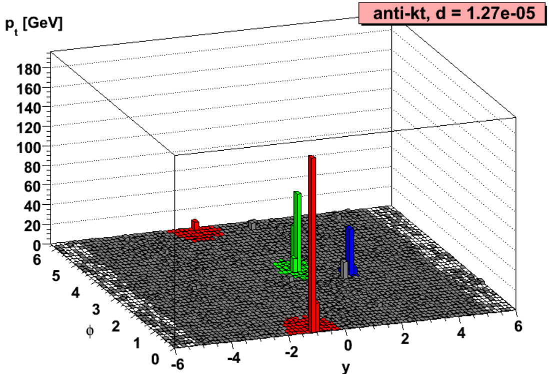

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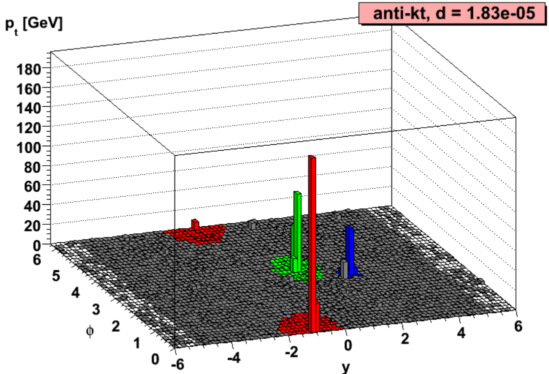

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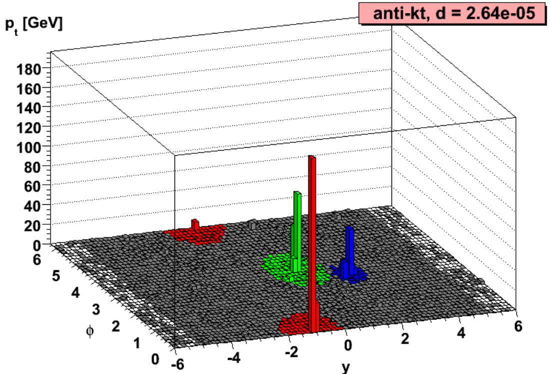

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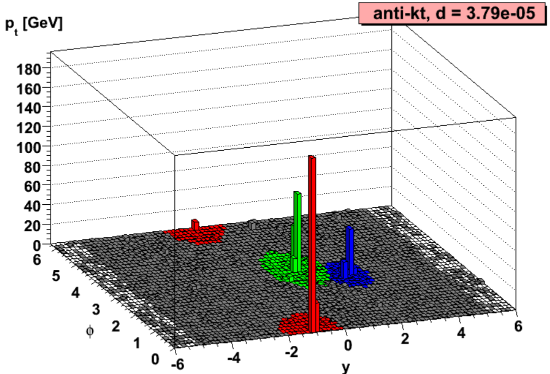

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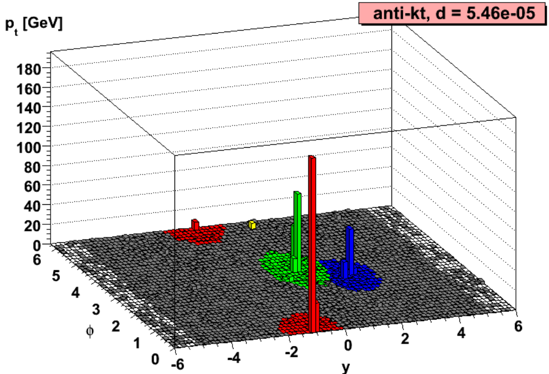

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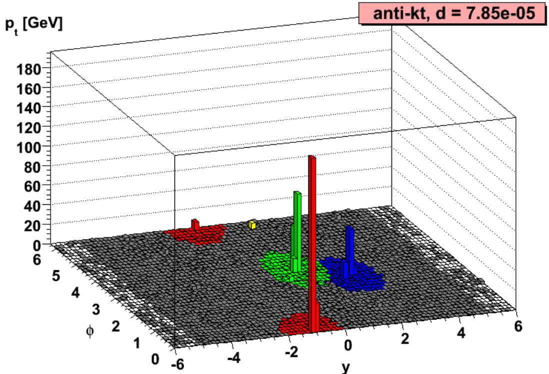

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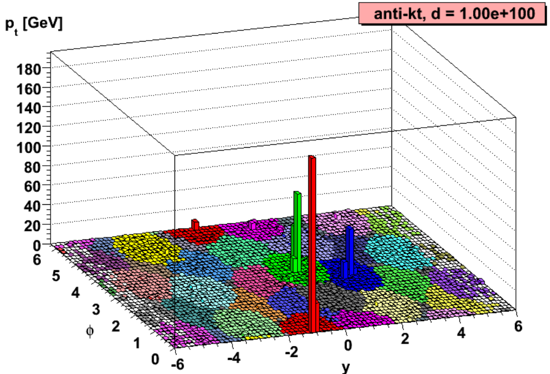

anti-k_t gives cone-like jets without using stable cones

A full set of IRC-safe jet algorithms

Generalise inclusive-type sequential recombination with

$$d_{ij} = \min(k_{ti}^{2\mathbf{p}}, k_{tj}^{2\mathbf{p}})\Delta R_{ij}^2/R^2 \qquad d_{iB} = k_{ti}^{2\mathbf{p}}$$

	Alg. name	Comment	time
$p = 1$	k_t CDOSTW '91-93; ES '93	Hierarchical in rel. k_t	$N \ln N$ exp.
$p = 0$	Cambridge/Aachen Dok, Leder, Moretti, Webber '97 Wengler, Wobisch '98	Hierarchical in angle Scan multiple R at once $\leftrightarrow$ QCD angular ordering	$N \ln N$
$p = -1$	anti-k_t Cacciari, GPS, Soyez '08 $\sim$ reverse-k_t Delsart	Hierarchy meaningless, jets like CMS cone (IC-PR)	$N^{3/2}$
SC-SM	SISCone GPS Soyez '07 + Tevatron run II '00	Replaces JetClu, ATLAS MidPoint (xC-SM) cones	$N^2 \ln N$ exp.

All these algorithms coded in (efficient) C++ at
http://fastjet.fr/ (Cacciari, GPS & Soyez '05-08)

Evolution since 2005

Algorithm	Type	IRC status	Evolution
exclusive k_t	$SR_{p=1}$	OK	$N^3 \to N \ln N$
inclusive k_t	$SR_{p=1}$	OK	$N^3 \to N \ln N$
Cambridge/Aachen	$SR_{p=0}$	OK	$N^3 \to N \ln N$
Run II Seedless cone	SC-SM	OK	$\to$ SISCone
CDF JetClu	IC_r-SM	IR_{2+1}	[$\to$ SISCone]
CDF MidPoint cone	IC_{mp}-SM	IR_{3+1}	$\to$ SISCone
CDF MidPoint searchcone	$IC_{se,mp}$-SM	IR_{2+1}	[$\to$ SISCone]
D0 Run II cone	IC_{mp}-SM	IR_{3+1}	$\to$ SISCone [with p_t cut?]
ATLAS Cone	IC-SM	IR_{2+1}	$\to$ SISCone
PxCone	IC_{mp}-SD	IR_{3+1}	[little used]
CMS Iterative Cone	IC-PR	$Coll_{3+1}$	$\to$ anti-k_t
PyCell/CellJet (from Pythia)	FC-PR	$Coll_{3+1}$	$\to$ anti-k_t
GetJet (from ISAJET)	FC-PR	$Coll_{3+1}$	$\to$ anti-k_t

SR = seq.rec.; IC = it.cone; FC = fixed cone;

SM = split–merge; SD = split–drop; PR = progressive removal

Snowmass is solved

But it was a problem from the 1990s

What are the problems we *should* be trying to solve for LHC?

Which jet definition(s) for LHC?

Choice of algorithm (k_t, SISCone, ...)

Choice of parameters (R, ...)

Can we address this question scientifically?

Jetography

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Choice of algorithm (k_t, SISCone, ...)

Choice of parameters (R, ...)

Can we address this question scientifically?

Jetography

Jet def^n differences

Jet definitions differ mainly in:

$\underbrace{\text{Jet definitions}}_{\text{alg} + R}$

1. How close two particles must be to end up in same jet
[discussed in the '90s, e.g. Ellis & Soper]

2. How much perturbative radiation is lost from a jet
[indirectly discussed in the '90s (analytic NLO for inclusive jets)]

3. How much non-perturbative contamination (hadronisation, UE, pileup) a jet receives
[partially discussed in '90s — Korchemsky & Sterman '95, Seymour '97]

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Jet p_t v. parton p_t: perturbatively?

The question's dangerous: a "parton" is an ambiguous concept

Three limits can help you:

- Threshold limit — e.g. de Florian & Vogelsang '07
- Parton from color-neutral object decay (Z')
- Small-R (radius) limit for jet

One simple result

$$\frac{\langle p_{t,jet} - p_{t,parton} \rangle}{p_t} = \frac{\alpha_s}{\pi} \ln R \times \begin{cases} 1.01 C_F & \textit{quarks} \\ 0.94 C_A + 0.07 n_f & \textit{gluons} \end{cases} + \mathcal{O}(\alpha_s)$$

only $\mathcal{O}(\alpha_s)$ depends on algorithm & process
cf. Dasgupta, Magnea & GPS '07

Jet p_t v. parton p_t: hadronisation?

Hadronisation: the "parton-shower" $\rightarrow$ hadrons transition

Method:

- "infrared finite α_s" — à la Dokshitzer & Webber '95
- **prediction** based on e^+e^- event shape data
- could have been deduced from old work — Korchemsky & Sterman '95, Seymour '97

Main result

$$\langle p_{t,jet} - p_{t,parton-shower} \rangle \simeq -\frac{0.4\text{ GeV}}{R} \times \begin{cases} C_F & quarks \\ C_A & gluons \end{cases}$$

cf. Dasgupta, Magnea & GPS '07
coefficient holds for anti-k_t; see Mrinal's talk for k_t alg.

Underlying Event (UE)

"Naive" prediction (UE $\simeq$ colour dipole between *pp*):

$$\Delta p_t \simeq 0.4\ \text{GeV} \times \frac{R^2}{2} \times \begin{cases} C_F & q\bar{q}\ \text{dipole} \\ C_A & \text{gluon dipole} \end{cases}$$

DWT Pythia tune or ATLAS Jimmy tune tell you:

$$\Delta p_t \simeq \mathbf{10 - 15\ GeV} \times \frac{R^2}{2}$$

This big coefficient motivates special effort to understand interplay between jet algorithm and UE: "jet areas"

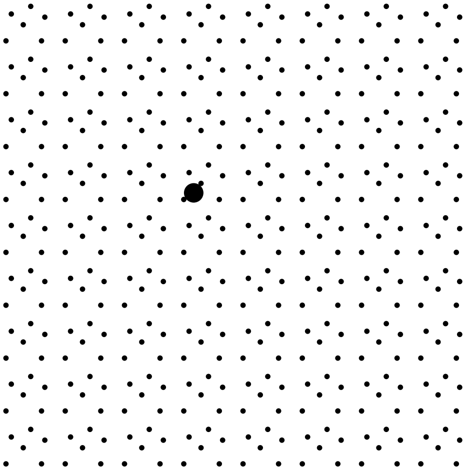
How does coefficient depend on algorithm?

How does it depend on jet p_t? How does it fluctuate?

cf. Cacciari, GPS & Soyez '08

E.g. SISCone jet area

1. One hard particle, many soft



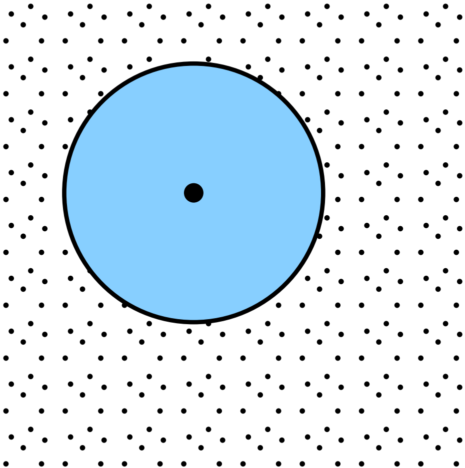
SISCone, any R, $f \gtrsim 0.391$

Jet area =
Measure of jet's susceptibility to uniform soft radiation

Depends on details of an algorithm's clustering dynamics.

E.g. SISCone jet area

2. One hard stable cone, area $= \pi R^2$



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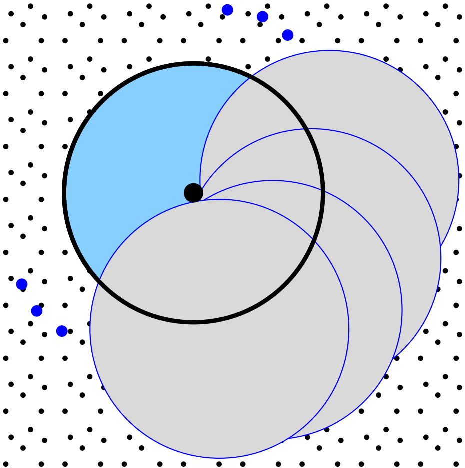
E.g. SISCone jet area

3. Overlapping "soft" stable cones



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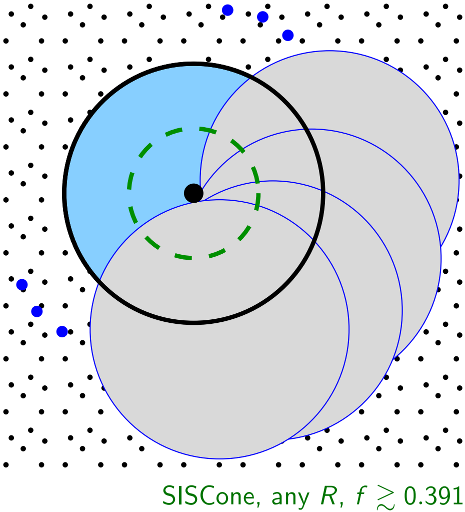
E.g. SISCone jet area

4. "Split" the overlapping parts



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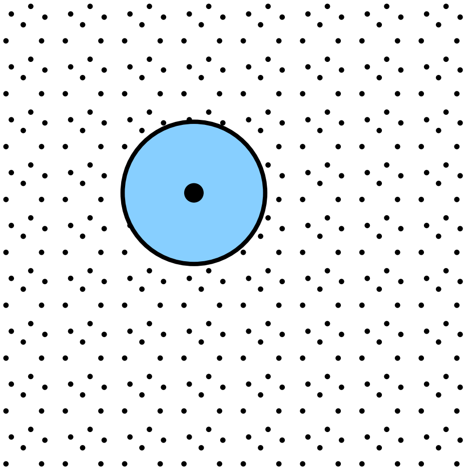
Measure of jet's susceptibility to uniform soft radiation

Depends on details of an algorithm's clustering dynamics.

E.g. SISCone jet area

5. Final hard jet (reduced area)



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Jet area =
Measure of jet's susceptibility to uniform soft radiation

Depends on details of an algorithm's clustering dynamics.

SISCone's area (1 hard particle)
$$= \frac{\mathbf{1}}{\mathbf{4}}\,\pi R^2$$

Jet algorithm properties: summary

	k_t	Cam/Aachen	anti-k_t	SISCone
reach	R	R	R	$(1+\frac{p_{t2}}{p_{t2}})R$
$\Delta p_{t,PT} \simeq \frac{\alpha_s C_i}{\pi} \times$	$\ln R$	$\ln R$	$\ln R$	$\ln 1.35R$
$\Delta p_{t,hadr} \simeq -\frac{0.4\text{ GeV}\, C_i}{R} \times$	0.7	?	1	?
area $= \pi R^2 \times$	0.81 ± 0.28	0.81 ± 0.26	1	0.25
$+\pi R^2 \frac{C_i}{\pi b_0} \ln \frac{\alpha_s(Q_0)}{\alpha_s(Rp_t)} \times$	0.52 ± 0.41	0.08 ± 0.19	0	0.12 ± 0.07

In words:

- k_t: area fluctuates a lot, depends on p_t (bad for UE)
- Cam/Aachen: area fluctuates somewhat, depends less on p_t
- anti-k_t: area is constant (circular jets)
- SISCone: reaches far for hard radiation (good for resolution, bad for multijets), area is smaller (good for UE)

Jet momentum significantly affected by R

So what R should we choose?

Examine this in context of reconstruction of dijet resonance

What R is best for an isolated jet?

E.g. to reconstruct $m_X \sim (p_{tq} + p_{t\bar{q}})$

PT radiation:

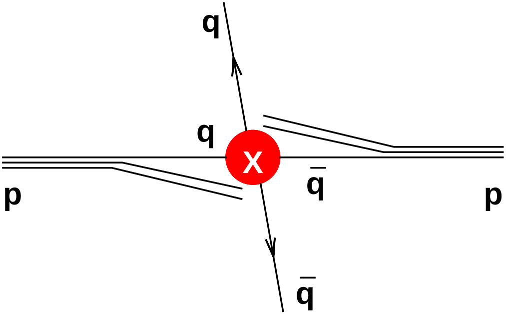
$$q: \quad \langle \Delta p_t \rangle \simeq \frac{\alpha_s C_F}{\pi} p_t \ln R$$

Hadronisation:

$$q: \quad \langle \Delta p_t \rangle \simeq -\frac{C_F}{R} \cdot 0.4 \text{ GeV}$$

Underlying event:

$$q, g: \quad \langle \Delta p_t \rangle \simeq \frac{R^2}{2} \cdot 2.5\text{–}15 \text{ GeV}$$



Minimise fluctuations in p_t

Use crude approximation:

$$\langle \Delta p_t^2 \rangle \simeq \langle \Delta p_t \rangle^2$$

in small-R limit (?!)
cf. Dasgupta, Magnea & GPS '07

What R is best for an isolated jet?

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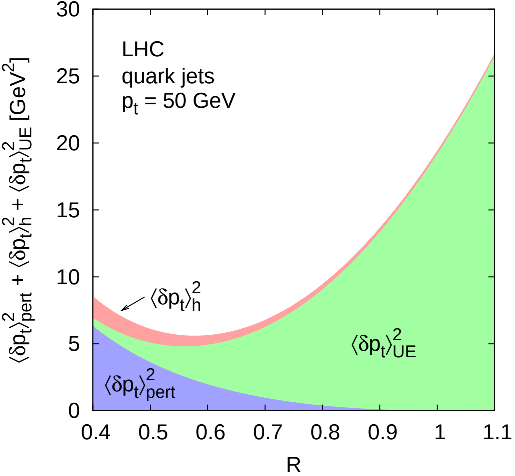
50 GeV quark jet



in small-R limit (?!)

cf. Dasgupta, Magnea & GPS '07

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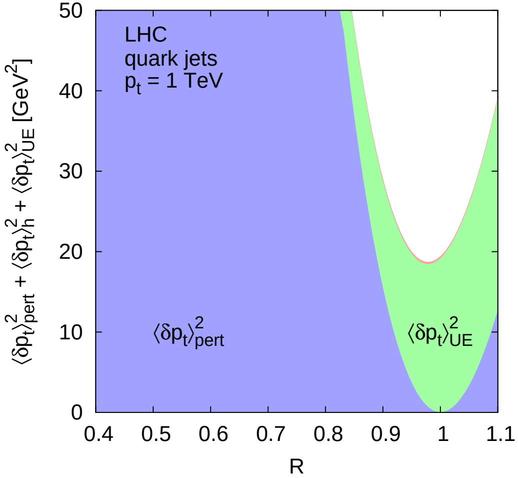
1 TeV quark jet



in small-R limit (?!)

cf. Dasgupta, Magnea & GPS '07

What R is best for an isolated jet?

PT radiation:

$$q: \quad \langle \Delta p_t \rangle \simeq \frac{\alpha_s C_F}{\pi} p_t \ln R$$

Hadr

$q:$

At low p_t, small R limits relative impact of UE

At high p_t, perturbative effects dominate over non-perturbative $\rightarrow R_{best} \sim 1$.

Underlying event:

$$q, g: \quad \langle \Delta p_t \rangle \simeq \frac{R^2}{2} \cdot 2.5\text{—}15 \text{ GeV}$$

Minimise fluctuations in p_t

Use crude approximation:

$$\langle \Delta p_t^2 \rangle \simeq \langle \Delta p_t \rangle^2$$

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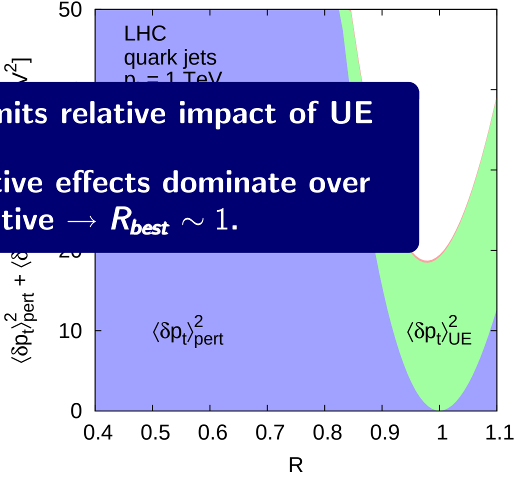
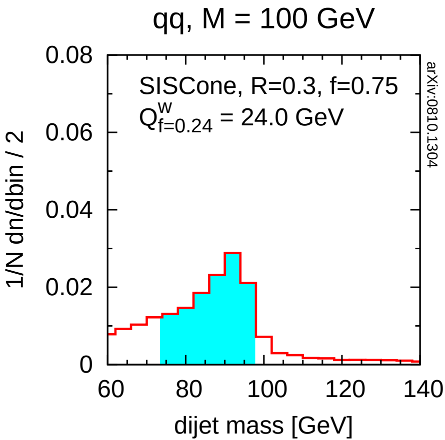


in small-R limit (?!)

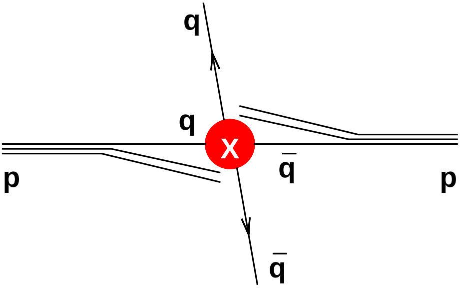
cf. Dasgupta, Magnea & GPS '07

Dijet mass: scan over *R* [Pythia 6.4]

R = 0.3



Resonance X → dijets



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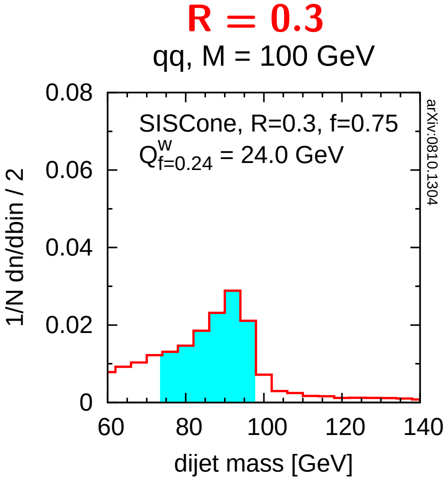
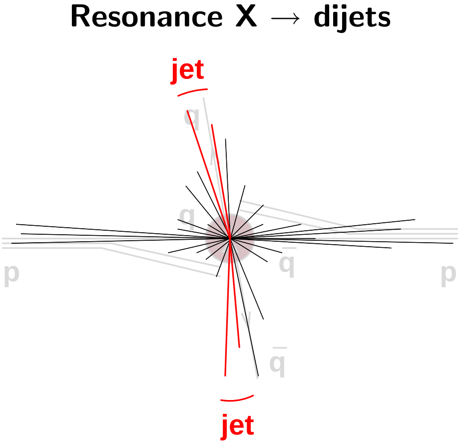
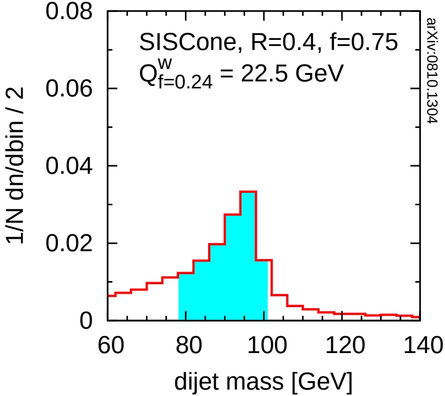




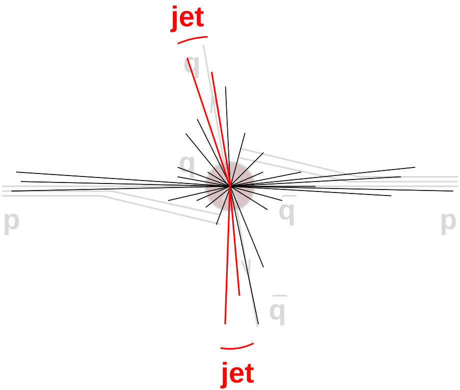
Dijet mass: scan over *R* [Pythia 6.4]

R = 0.4

qq, M = 100 GeV

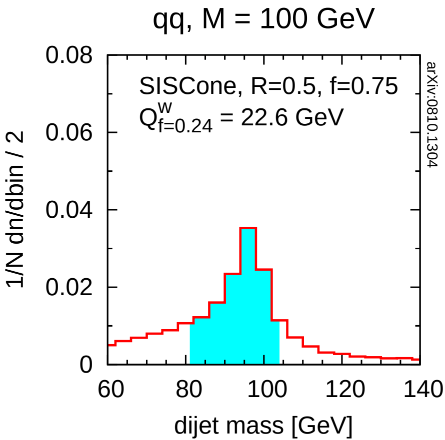


Resonance X → dijets

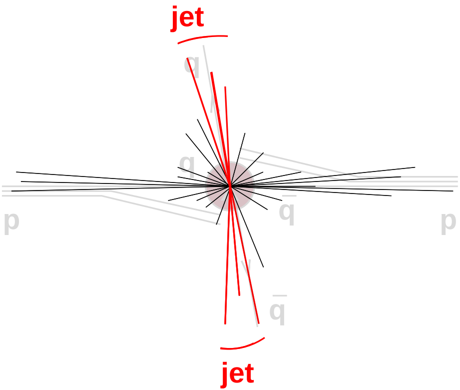


Dijet mass: scan over *R* [Pythia 6.4]

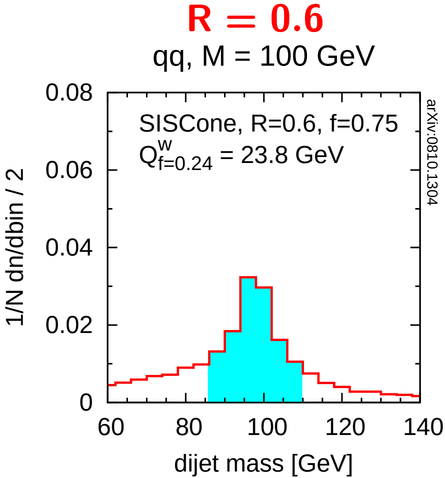
R = 0.5

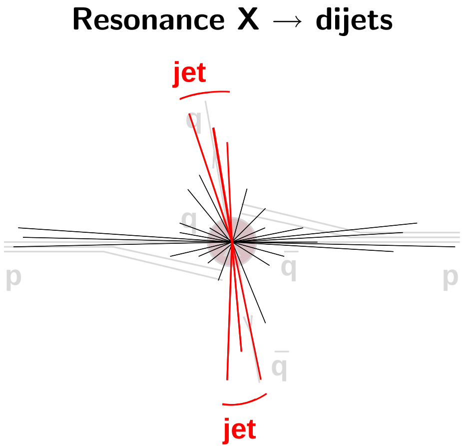


Resonance X → dijets

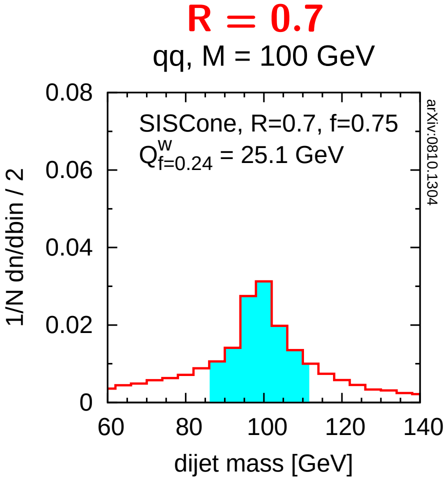


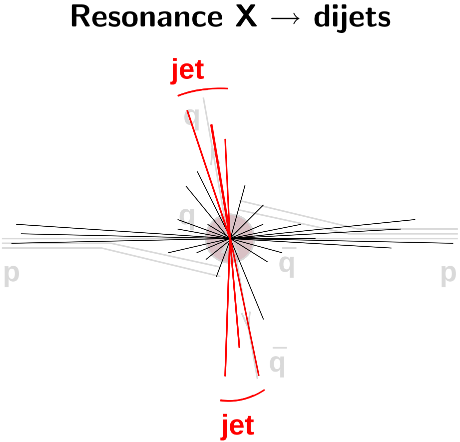
Dijet mass: scan over *R* [Pythia 6.4]





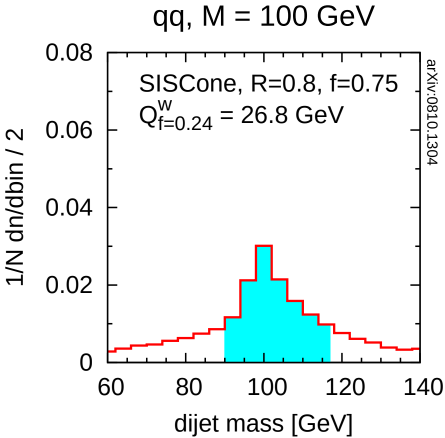
Dijet mass: scan over *R* [Pythia 6.4]



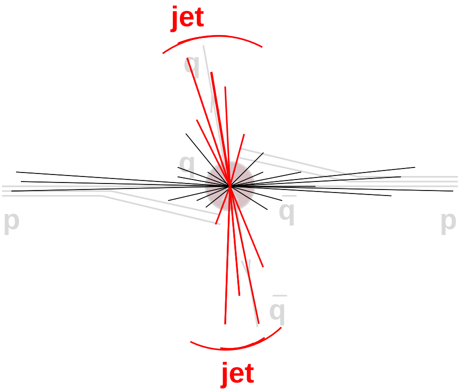


Dijet mass: scan over *R* [Pythia 6.4]

R = 0.8



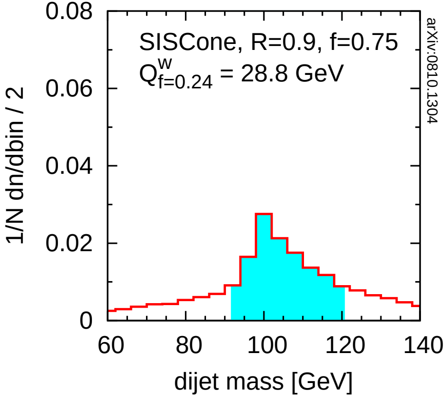
Resonance X → dijets



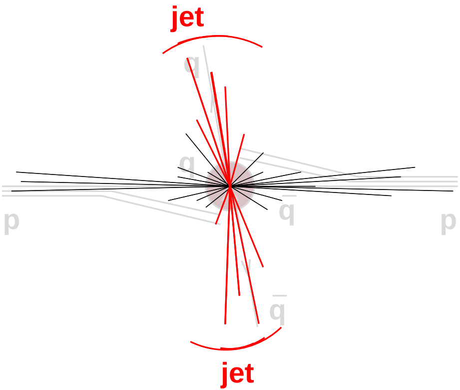
Dijet mass: scan over *R* [Pythia 6.4]

R = 0.9

qq, M = 100 GeV



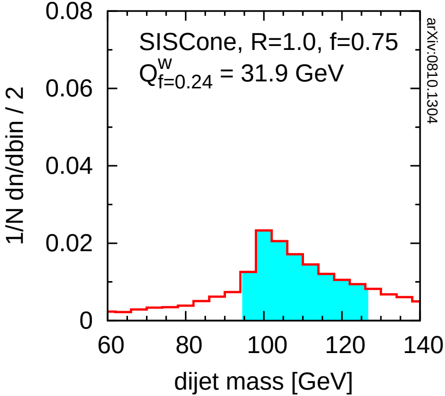
Resonance X → dijets



Dijet mass: scan over *R* [Pythia 6.4]

R = 1.0

qq, M = 100 GeV



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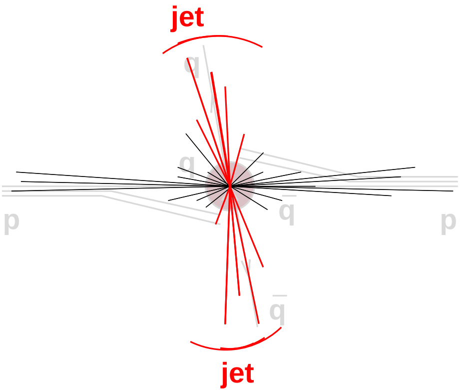
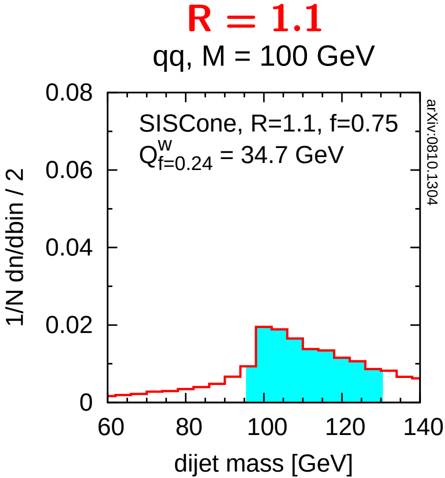
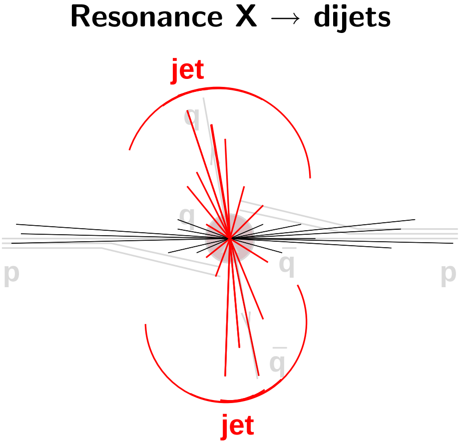


Dijet mass: scan over *R* [Pythia 6.4]

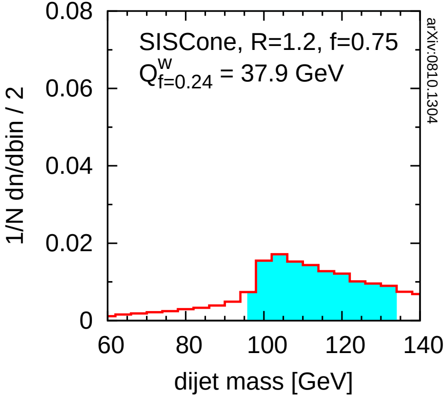




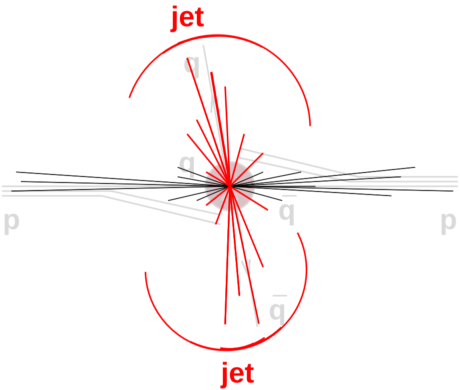
Dijet mass: scan over *R* [Pythia 6.4]

R = 1.2

qq, M = 100 GeV



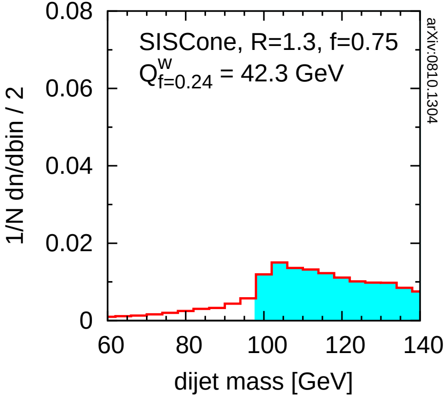
Resonance X → dijets



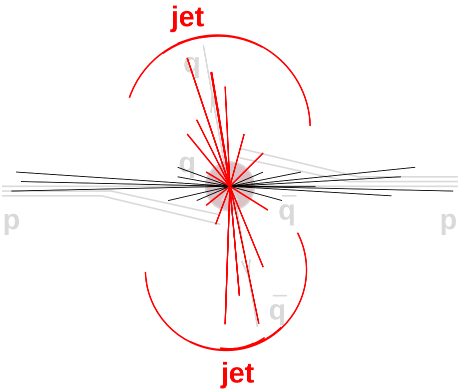
Dijet mass: scan over *R* [Pythia 6.4]

R = 1.3

qq, M = 100 GeV

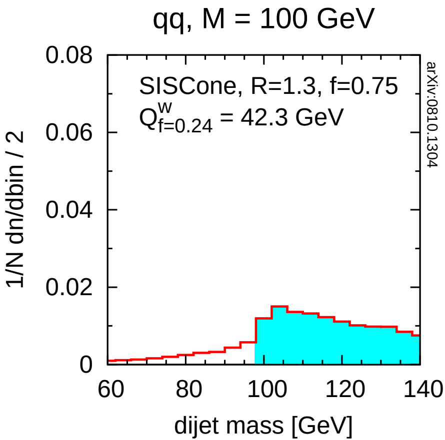


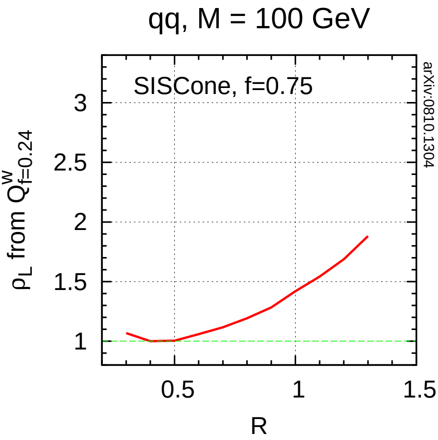
Resonance X → dijets



Dijet mass: scan over *R* [Pythia 6.4]

R = 1.3



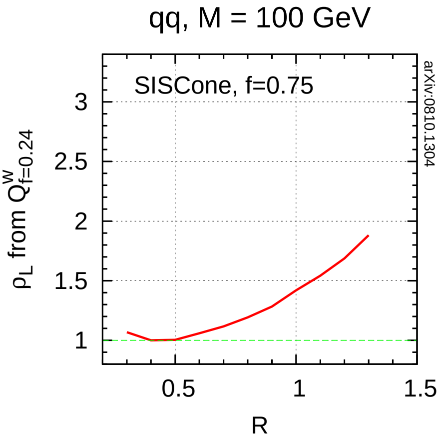


After scanning, summarise "quality" v. *R*. Minimum ≡ BEST

picture not so different from crude analytical estimate

Scan through $q\bar{q}$ mass values

$\mathbf{m_{qq} = 100\ GeV}$



Best R is at minimum of curve

- Best R depends strongly on mass of system
- Increases with mass, just like crude analytical prediction

  NB: current analytics too crude

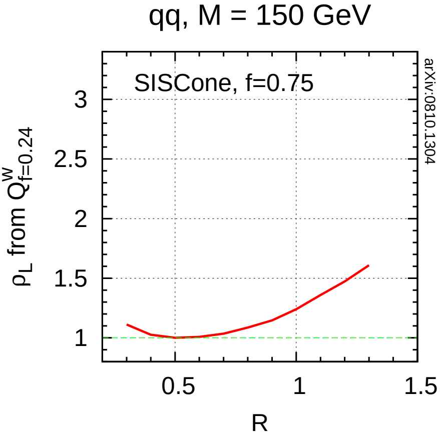
BUT: so far, LHC's plans involve running with fixed smallish R values

e.g. CMS arXiv:0807.4961

NB: 100,000 plots for various jet algorithms, narrow qq and gg resonances from `http://quality.fastjet.fr` Cacciari, Rojo, GPS & Soyez '08

Scan through $q\bar{q}$ mass values

$\mathbf{m_{qq} = 150\ GeV}$



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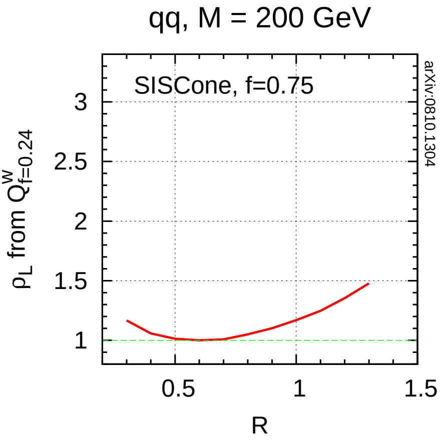
Scan through $q\bar{q}$ mass values

$m_{qq} = 200$ GeV



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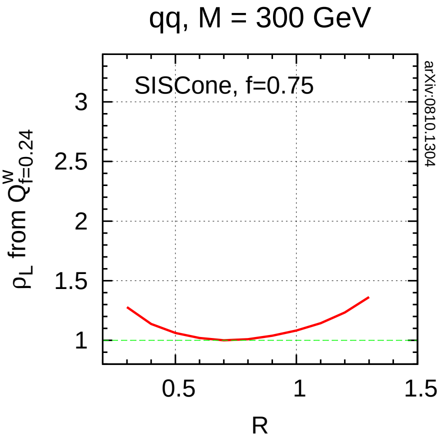
Scan through $q\bar{q}$ mass values

$\mathbf{m_{qq} = 300\ GeV}$



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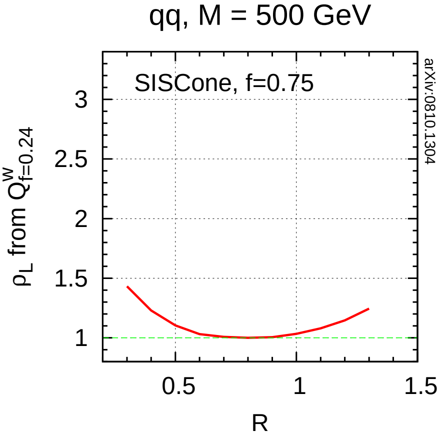
Scan through $q\bar{q}$ mass values

$\mathbf{m_{qq} = 500\ GeV}$



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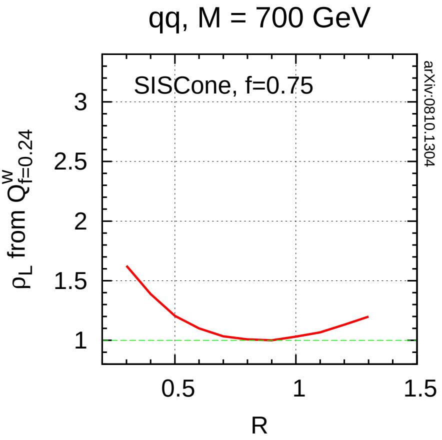
e.g. CMS arXiv:0807.4961

NB: 100,000 plots for various jet algorithms, narrow qq and gg resonances from `http://quality.fastjet.fr` Cacciari, Rojo, GPS & Soyez '08

Scan through $q\bar{q}$ mass values

$\mathbf{m_{qq} = 700\ GeV}$



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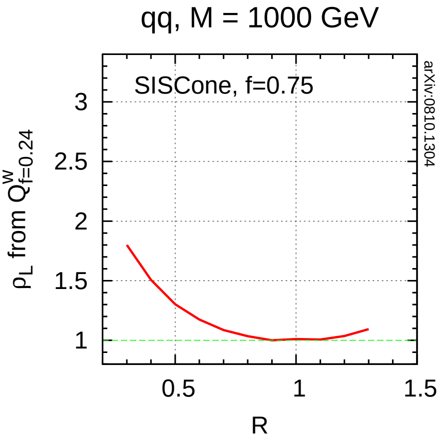
e.g. CMS arXiv:0807.4961

NB: 100,000 plots for various jet algorithms, narrow qq and gg resonances from `http://quality.fastjet.fr` Cacciari, Rojo, GPS & Soyez '08

Scan through $q\bar{q}$ mass values

$\mathbf{m_{qq} = 1000\ GeV}$



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BUT: so far, LHC's plans involve running with fixed smallish R values

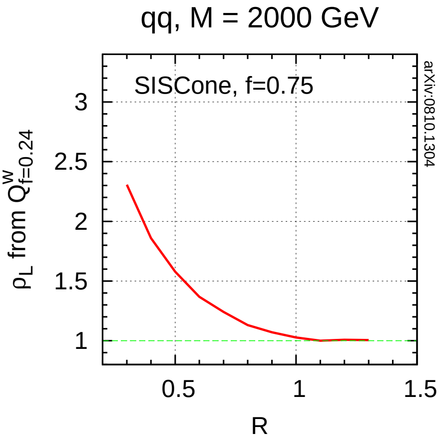
e.g. CMS arXiv:0807.4961

NB: 100,000 plots for various jet algorithms, narrow qq and gg resonances from http://quality.fastjet.fr

Cacciari, Rojo, GPS & Soyez '08

Scan through $q\bar{q}$ mass values

$m_{qq} = 2000$ GeV



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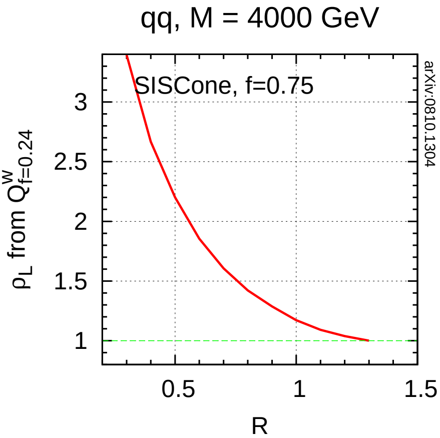
Scan through $q\bar{q}$ mass values

$m_{qq} = 4000$ GeV



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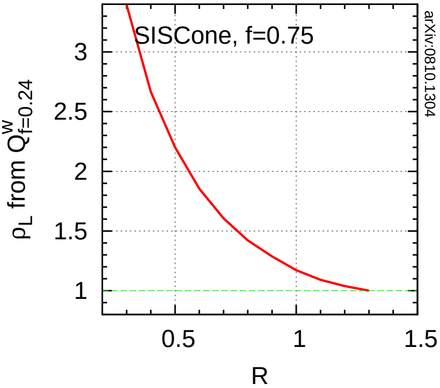

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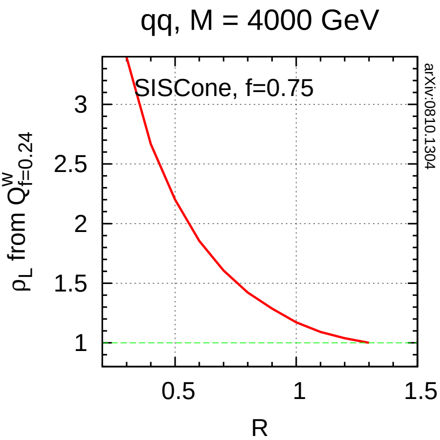

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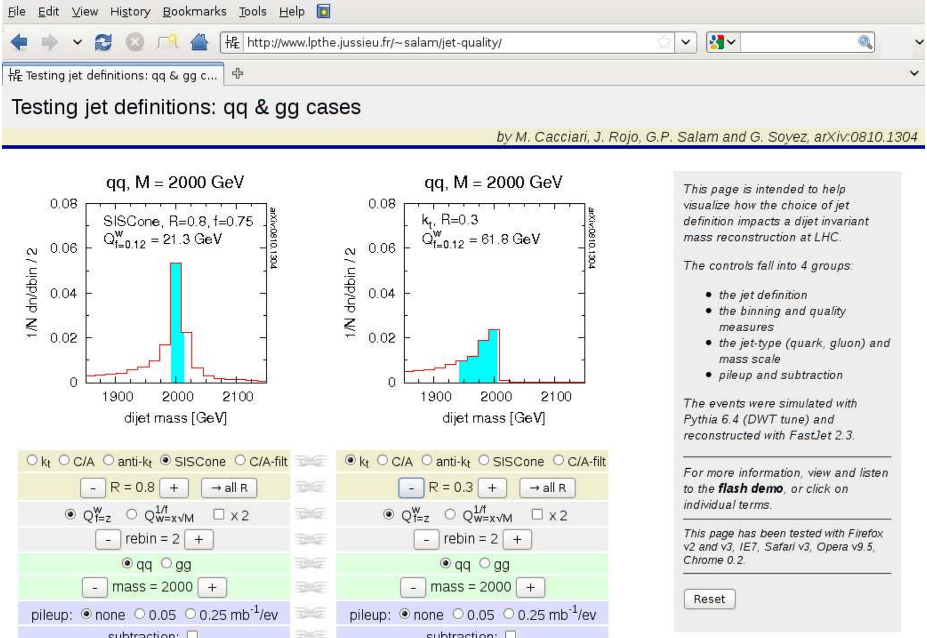
http://quality.fastjet.fr/



How about task of resolving separate jets from separate partons?

Illustrate in context of boosted $H \longrightarrow b\bar{b}$ reconstruction

E.g.: WH/ZH search channel @ LHC

- Signal is $W \to \ell\nu$, $H \to b\bar{b}$.
- Backgrounds include $Wb\bar{b}$, $t\bar{t} \to \ell\nu b\bar{b}jj$, ...

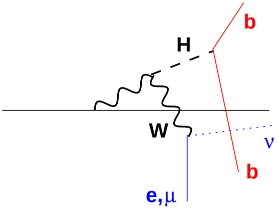
Studied e.g. in ATLAS TDR

Difficulties, e.g.

- $gg \to t\bar{t}$ has $\ell\nu b\bar{b}$ with **same intrinsic mass scale**, but much higher partonic luminosity
- Need exquisite control of bkgd shape

Try a long shot?

- Go to high p_t ($p_{tH}, p_{tV} > 200$ GeV)
- Lose 95% of signal, but more efficient?
- Maybe kill $t\bar{t}$ & gain clarity?



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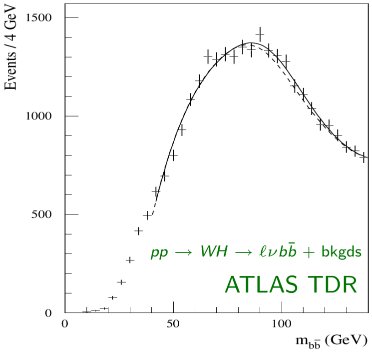

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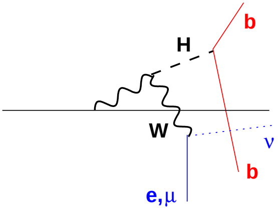

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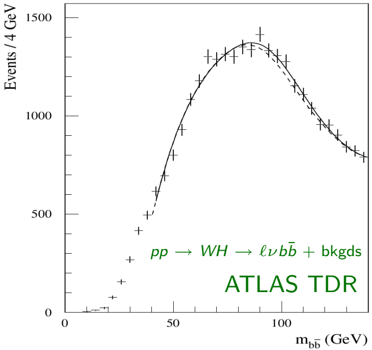
- Backgrounds include $Wb\bar{b}$, $t\bar{t} \to \ell\nu b\bar{b}jj$, ...

Studied e.g. in ATLAS TDR

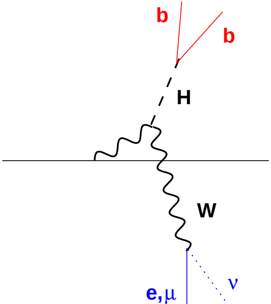


Difficulties, e.g.

- $gg \to t\bar{t}$ has $\ell\nu b\bar{b}$ with **same intrinsic mass scale**, but much higher partonic luminosity
- Need exquisite control of bkgd shape

Try a long shot?

- Go to high p_t ($p_{tH}, p_{tV} > 200$ GeV)
- Lose 95% of signal, but more efficient?
- Maybe kill $t\bar{t}$ & gain clarity?



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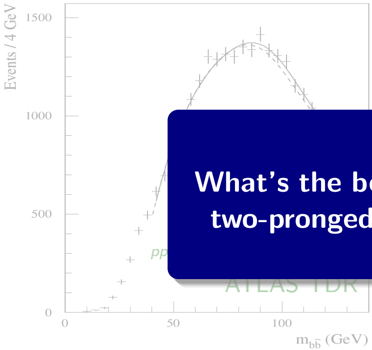

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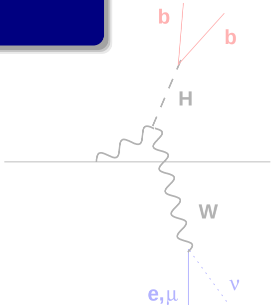
- kgd shape



Question:

What's the best strategy to identify the two-pronged structure of the boosted Higgs decay?

Try a long shot?

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Past methods

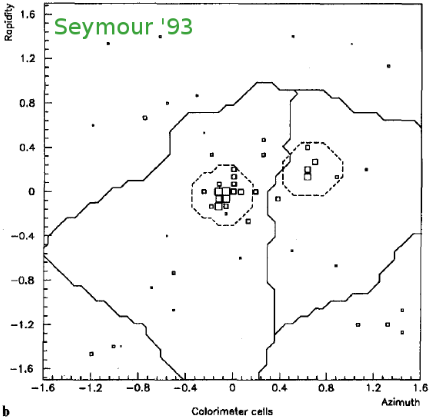
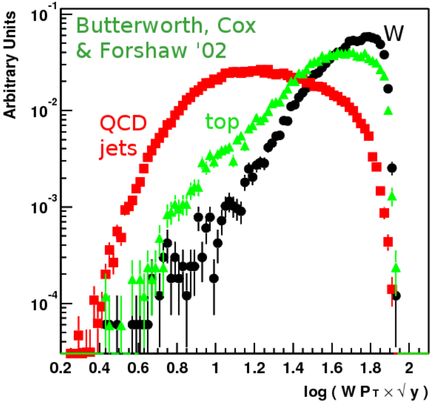


Fig. 2. A hadronic W decay, as seen at calorimeter level, **a** without, and **b** with, particles from the underlying event. Box sizes are logarithmic in the cell energy, lines show the borders of the sub-jets for infinitely soft emission according to the cluster (solid) and cone (dashed) algorithms

Use k_t jet-algorithm's hierarchy to split the jets



Use k_t alg.'s distance measure (rel. trans. mom.) to cut out QCD bkgd:

$$d_{ij}^{k_t} = \min(p_{ti}^2, p_{tj}^2)\Delta R_{ij}^2$$

Y-splitter

only partially correlated with mass

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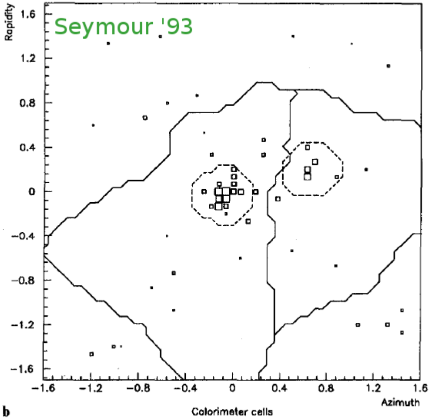

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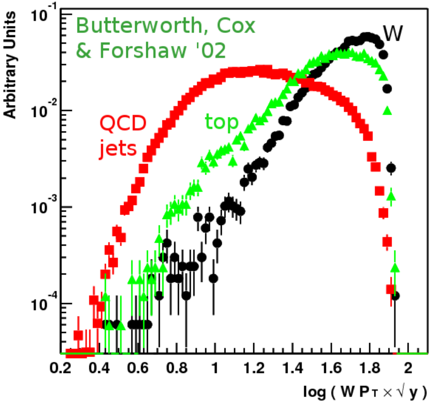

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Our tool

The Cambridge/Aachen jet alg.

Dokshitzer et al '97
Wengler & Wobisch '98

Work out $\Delta R_{ij}^2 = \Delta y_{ij}^2 + \Delta\phi_{ij}^2$ between all pairs of objects i, j;
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Repeat until all objects separated by $\Delta R_{ij} > R$. [in FastJet]

Gives "hierarchical" view of the event; work through it backwards to analyse jet

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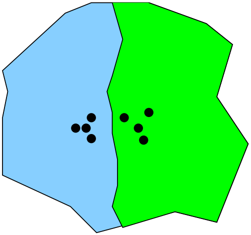
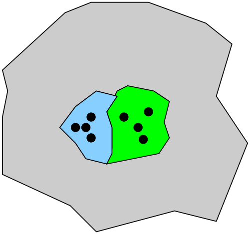
k_t algorithm



Cam/Aachen algorithm

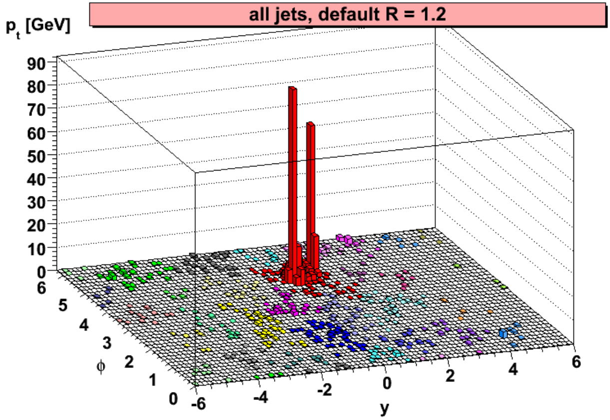


Allows you to "dial" the correct R to keep perturbative radiation, but throw out UE

$pp \longrightarrow ZH \longrightarrow \nu\bar{\nu}b\bar{b}$, @14 TeV, $m_H = 115$ GeV

SIGNAL

Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



Zbb BACKGROUND

Cluster event, C/A, R=1.2

arbitrary norm.

$pp \longrightarrow ZH \longrightarrow \nu\bar{\nu}b\bar{b}$, @14 TeV, $m_H = 115$ GeV

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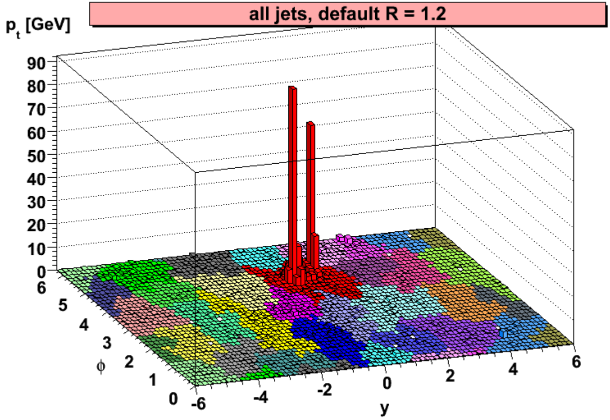

Zbb BACKGROUND

Fill it in, → show jets more clearly

arbitrary norm.

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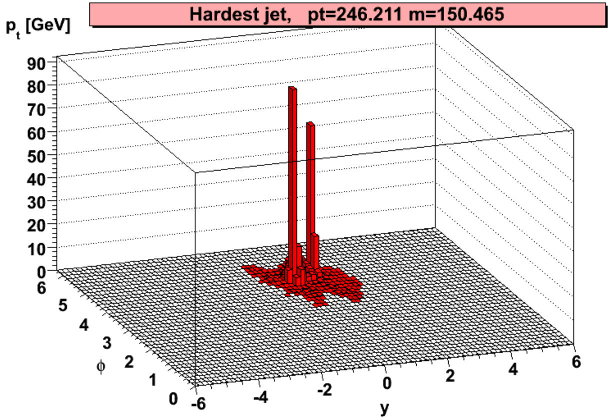
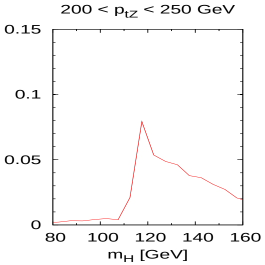
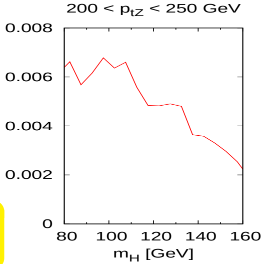


SIGNAL



Zbb BACKGROUND



Consider hardest jet, $m = 150$ GeV

arbitrary norm.

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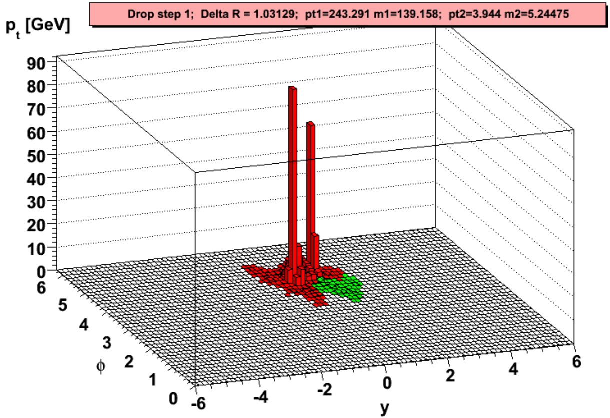
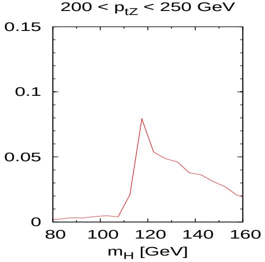


SIGNAL



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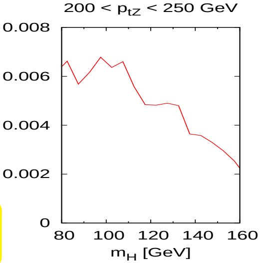
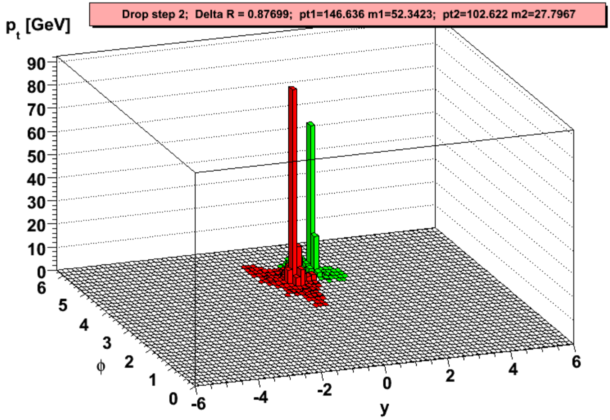


split: $m = 150$ GeV, $\frac{\max(m_1, m2)}{m} = 0.92 \longrightarrow$ repeat

arbitrary norm.

$pp \to ZH \to \nu\bar{\nu}b\bar{b}$, @14 TeV, $m_H = 115$ GeV

Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



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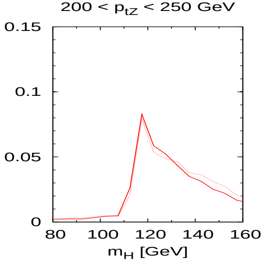
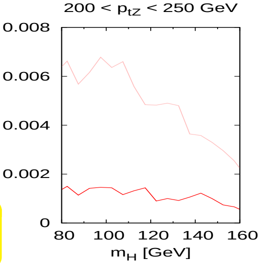


Zbb BACKGROUND

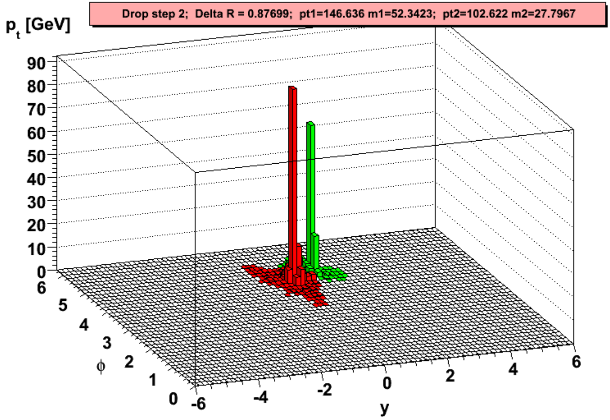


split: $m = 139$ GeV, $\frac{\max(m_1, m2)}{m} = 0.37 \to$ mass drop

arbitrary norm.

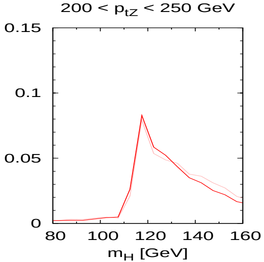
$pp \to ZH \to \nu\bar{\nu}b\bar{b}$, @14 TeV, $m_H = 115$ GeV

Herwig 6.510 + Jimmy 4.31 + FastJet 2.3

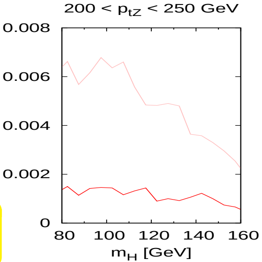


check: $y_{12} \simeq \frac{p_{t2}}{p_{t1}} \simeq 0.7 \to$ OK + 2 b-tags (anti-QCD)

SIGNAL



Zbb BACKGROUND



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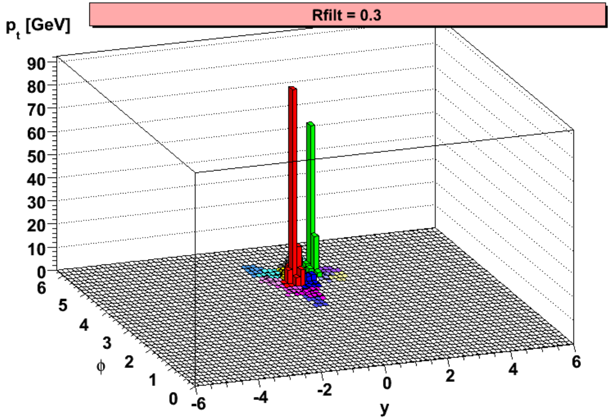
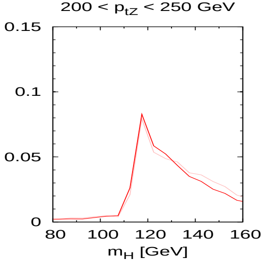


$R_{filt} = 0.3$

SIGNAL



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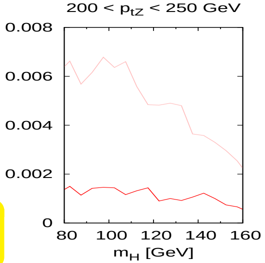

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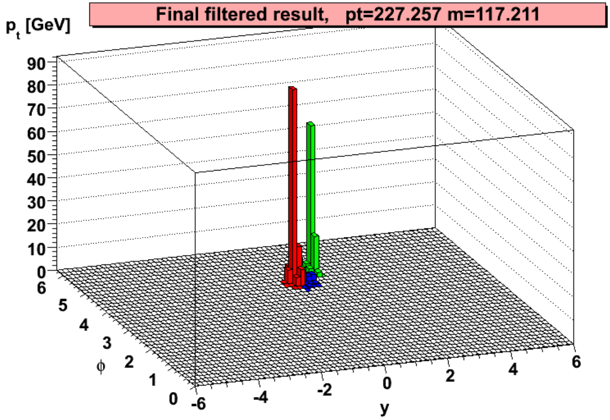
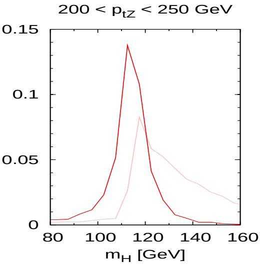


$R_{filt} = 0.3$: take 3 hardest, **m = 117 GeV**

SIGNAL



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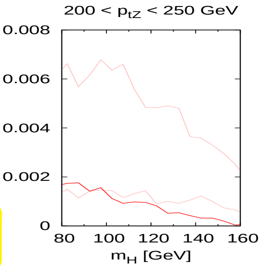


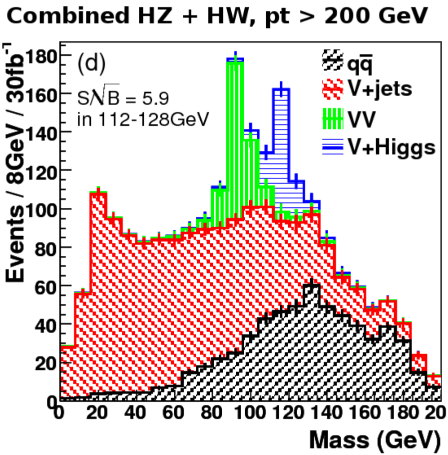
arbitrary norm.

Jet-alg comparison

Cross section for signal and the Z+jets background in the leptonic Z channel for $200 < p_{TZ}/\,\mathrm{GeV} < 600$ and $110 < m_J/\,\mathrm{GeV} < 125$, with perfect b-tagging; shown for our jet definition (C/A MD-F), and other standard ones close to their optimal R values.

Jet definition	σ_S/fb	σ_B/fb	$S/\sqrt{B \cdot \mathrm{fb}}$
C/A, $R = 1.2$, MD-F	0.57	0.51	0.80
k_t, $R = 1.0$, y_{cut}	0.19	0.74	0.22
SISCone, $R = 0.8$	0.49	1.33	0.42
anti-k_t, $R = 0.8$	0.22	1.06	0.21

combine HZ and HW, $p_t > 200$ GeV



- Take $Z \to \ell^+\ell^-$, $Z \to \nu\bar{\nu}$, $W \to \ell\nu$ $\qquad \ell = e, \mu$
- $p_{tV}, p_{tH} > 200$ GeV
- $|\eta_V|, |\eta_H| < 2.5$
- Assume real/fake b-tag rates of 0.7/0.01.
- Some extra cuts in HW channels to reject $t\bar{t}$.
- Assume $m_H = 115$ GeV.

At $\sim 5\sigma$ for 30 fb^{-1} this looks like a competitive channel for light Higgs discovery. **Deserves serious exp. study!**

Conclusions

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- There are no longer any valid excuses for using jet algorithms that are incompatible with the Snowmass criteria.

  LHC experiments are adopting the new tools

  Individual analyses need to follow suit

- It's time to move forwards with the question of how best to use jets in searches

- Examples here show two things:
  - Good jet-finding brings significant gains
  - There's room for serious QCD theory input into optimising jet use

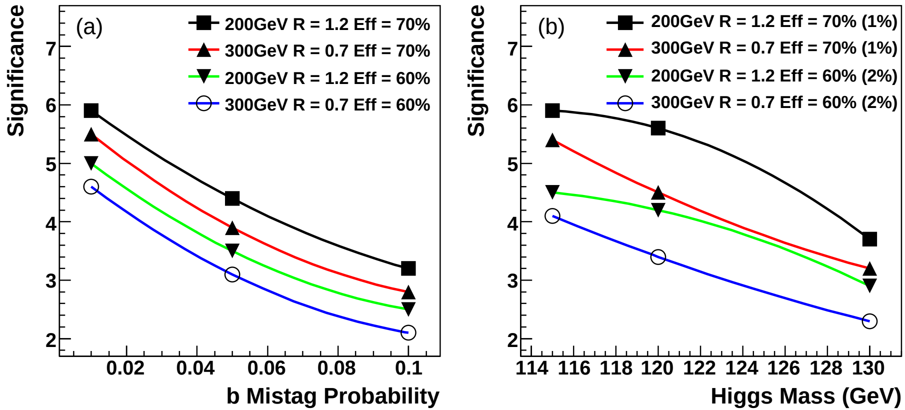
  Not the *only* way of doing things

  But brings more insight than trial & error MC

This opens the road towards *Jetography*, QCD-based autofocus for jets

EXTRAS

Impact of *b*-tagging, Higgs mass



Most scenarios above 3σ

For it to be a significant discovery channel requires decent *b*-tagging, lowish mass Higgs [and good experimental resolution]

In nearly all cases, looks feasible for extracting *WH*, *ZH* couplings

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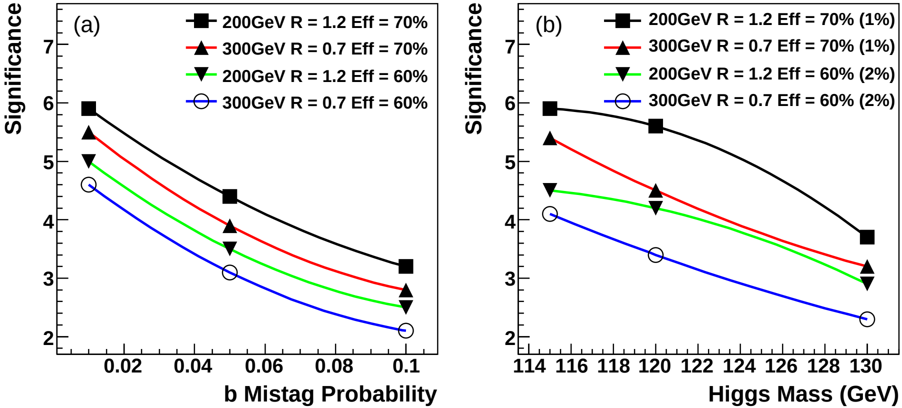

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