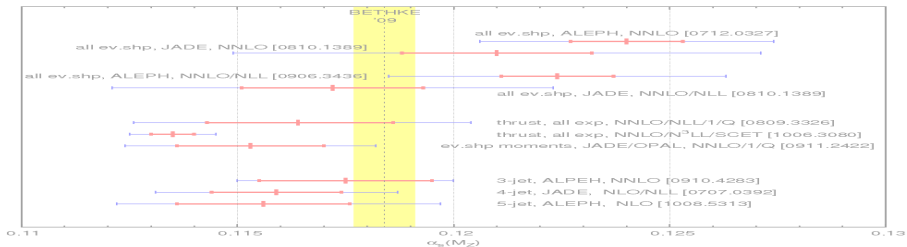


Review of event-shape measurements of α_s

Gavin Salam

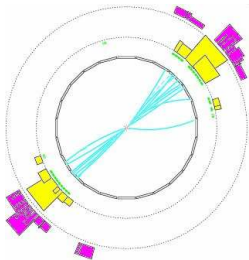
CERN, Princeton & LPTHE/CNRS (Paris)

Workshop on precision measurements of α_s
Munich, Germany, 9–11 February 2011

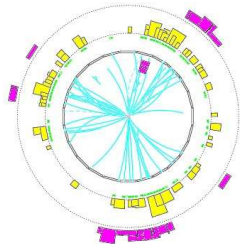


Various proposals to measure shape of events. Most famous example is **Thrust**:

$$T = \max_{\vec{n}_T} \frac{\sum_i |\vec{p}_i \cdot \vec{n}_T|}{\sum_i |\vec{p}_i|},$$



2-jet event: $T \simeq 1$



3-jet event: $T \simeq 2/3$

Fraction of events with $T \simeq 2/3$ is $\propto \alpha_s$

“Jet” masses: $M_i^2 = \left(\sum_{k \in H_i} p_k \right)^2$, $i = 1, 2$;

▶ Heavy-jet mass $\rho \equiv M_H^2 = \frac{\max(M_1^2, M_2^2)}{E_{vis}^2}$

Broadenings: $B_i = \frac{\sum_{k \in H_i} |\vec{p}_k \times \vec{n}_T|}{2 \sum_j |\vec{p}_j|}$, $i = 1, 2$;

▶ Total broadening: $B_T = B_1 + B_2$

▶ Wide-jet broadening: $B_W = \max(B_1, B + 2)$

C-parameter: $C = \frac{3 \sum_{i,j} |\vec{p}_i| |\vec{p}_j| \sin^2 \theta_{ij}}{(\sum_i |\vec{p}_i|)^2}$

Jet rates: most often defined using the k_t algorithm, as a function of jet resolution parameter $y_{cut} \sim k_t^2/Q^2$

NB: n -jet rate with $n > 3$ sensitive to α_s^{n-1} , i.e. enhanced sensitivity to α_s .

4 LEP experiments (ALEPH, DELPHI, L3, OPAL):

- ▶ LEP 1: very high statistics at $Q = M_Z$
- ▶ LEP 2: modest statistics at energies from 133...207 GeV
- ▶ Events with radiated photon giving effective hadronic $Q < M_Z$
Use of these data as “pure QCD with lower Q ” is of arguable legitimacy

SLD @ SLAC:

- ▶ very high statistics at $Q = M_Z$

JADE @ PETRA:

- ▶ good statistics from 14...44 GeV
- ▶ data reanalysed with modern tools and variables

TASSO, AMY, etc.:

- ▶ older data covering 14...55 GeV

Also DIS event shapes from HERA
Won't discuss here.

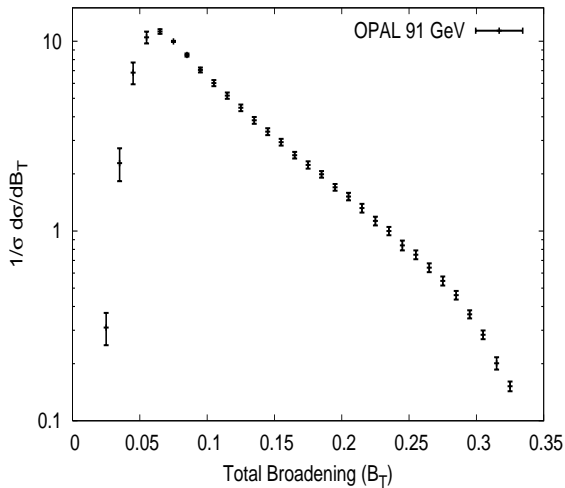
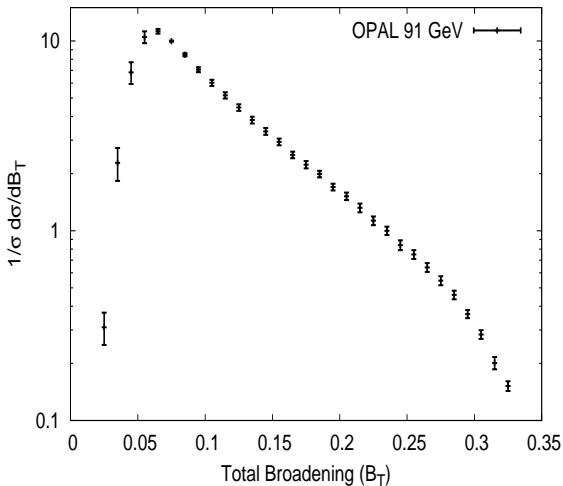


ILLUSTRATION OF
QUALITY OF DATA



Some typical kinds of ingredient

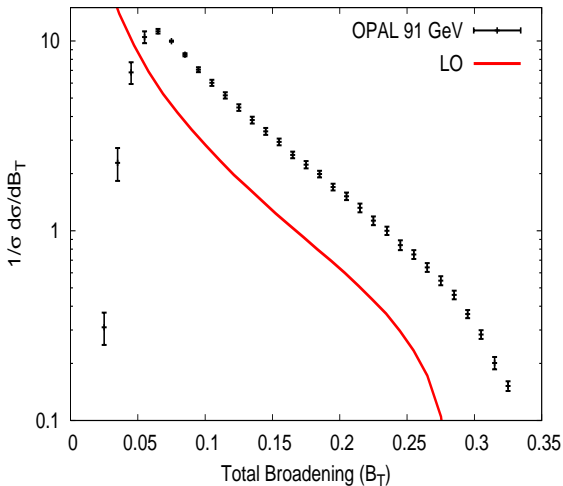
Fixed order calculation:

▶ LO: $\frac{d\sigma}{dB_T} = A_1(B_T)\alpha_s$

▶ NLO: add $A_2(B_T)\alpha_s^2$

Resummation of $\alpha_s^n \ln^{2n} \frac{1}{B_T}$

Hadronisation corrections



Some typical kinds of ingredient

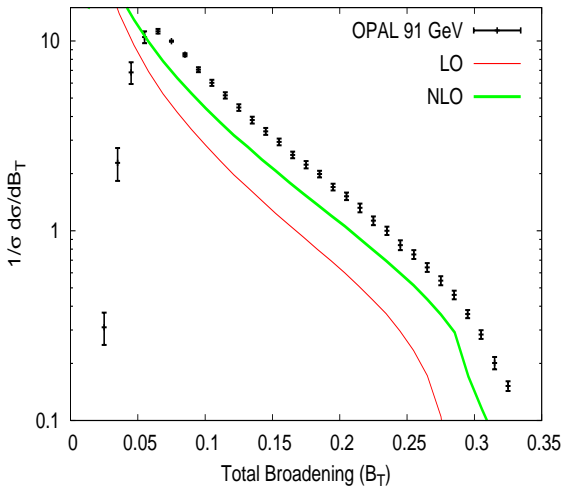
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Some typical kinds of ingredient

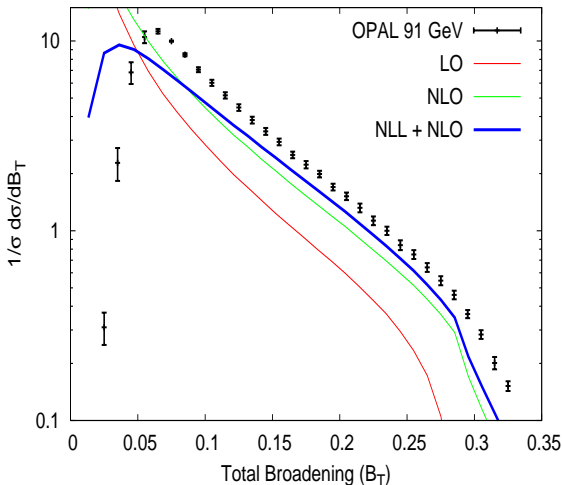
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Some typical kinds of ingredient

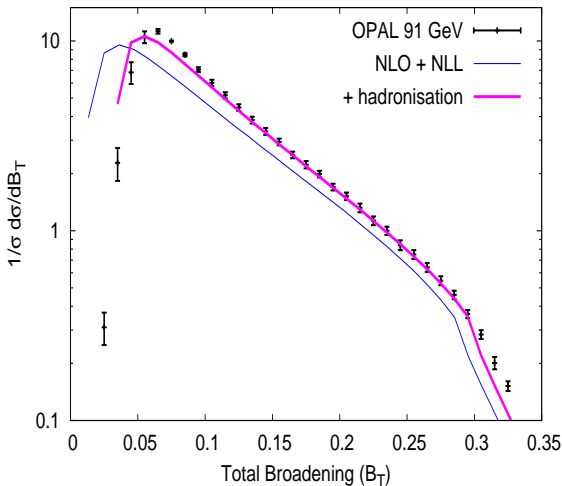
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Hadronisation corrections



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Fixed order calculation:

▶ LO: $\frac{d\sigma}{dB_T} = A_1(B_T)\alpha_s$

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Resummation of $\alpha_s^n \ln^{2n} \frac{1}{B_T}$

Hadronisation corrections

What's happened recently that's relevant for α_s ?

1. NNLO calculations for 3-jet type observables

Gehrmann–De Ridder et al '07; Weinzierl '08

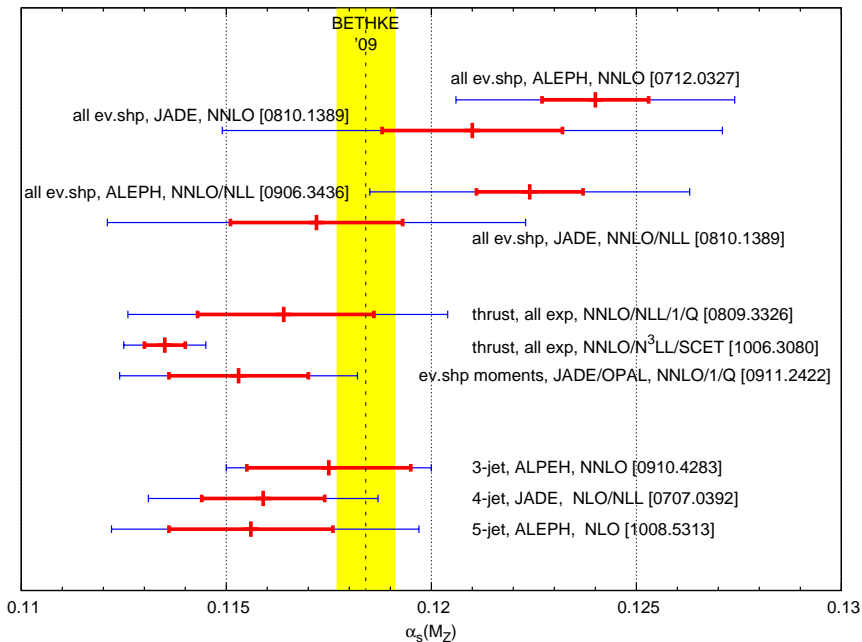
2. More accurate resummation (N^3LL) for thrust and heavy-jet mass in SCET

Becher & Schwartz '08; Chien & Schwartz '10

3. NLO calculation of 5-jet rate

Frederix et al '10

4. Lots of phenomenology



Data choices

- ▶ One experiment (and one energy): simpler combination of systematics
- ▶ Many experiments / many energies: more data, lever-arm in \sqrt{s} , etc.
- ▶ One “excellent” observable? Or many, as cross check of systematics?

Resummation

- ▶ Traditional resummation, available for many observables, but only NLL?
- ▶ Or SCET N³LL, available only for thrust and heavy-jet mass?
- ▶ Or no resummation at all?

Hadronisation

- ▶ Do you correct for it?
- ▶ Do you estimate it from Monte Carlos?
- ▶ Or from analytical models?

Data choices

- ▶ One experiment (and one energy): simpler combination of systematics
- ▶ Many experiments / many energies: more data, lever-arm in \sqrt{s} , etc.
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Re

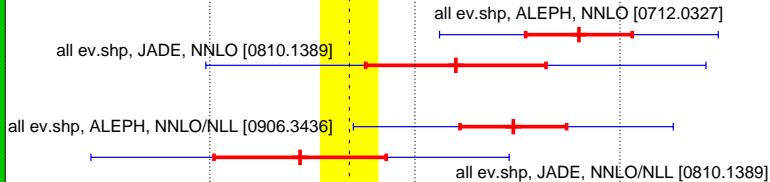
Often the choices are entangled with each other

- ▶ Or SCET N³LL, available only for thrust and heavy-jet mass?
- ▶ Or no resummation at all?

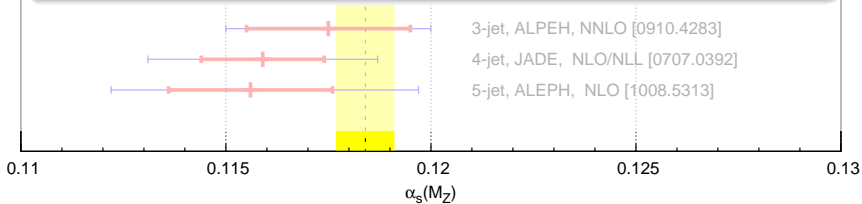
Hadronisation

- ▶ Do you correct for it?
- ▶ Do you estimate it from Monte Carlos?
- ▶ Or from analytical models?

BETHKE
'09

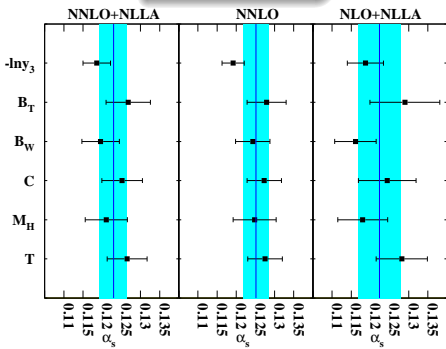


NNLO and NNLO/NLLA with hadronisation from Monte Carlo event generators

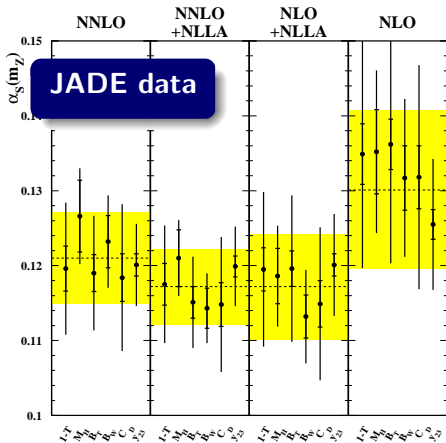


NNLO,	ALEPH	$0.1240 \pm .0008_{\text{stat}} \pm .0010_{\text{exp}} \pm .0011_{\text{had}} \pm .0029_{\text{theo}}$
NNLO	JADE	$0.1210 \pm .0007_{\text{stat}} \pm .0021_{\text{exp}} \pm .0044_{\text{had}} \pm .0036_{\text{theo}}$
NNLO/NLL,	ALEPH	$0.1224 \pm .0009_{\text{stat}} \pm .0009_{\text{exp}} \pm .0012_{\text{had}} \pm .0035_{\text{theo}}$
NNLO/NLL,	JADE	$0.1172 \pm .0006_{\text{stat}} \pm .0020_{\text{exp}} \pm .0035_{\text{had}} \pm .0030_{\text{theo}}$

ALEPH data



Dissertori et al '07, '08



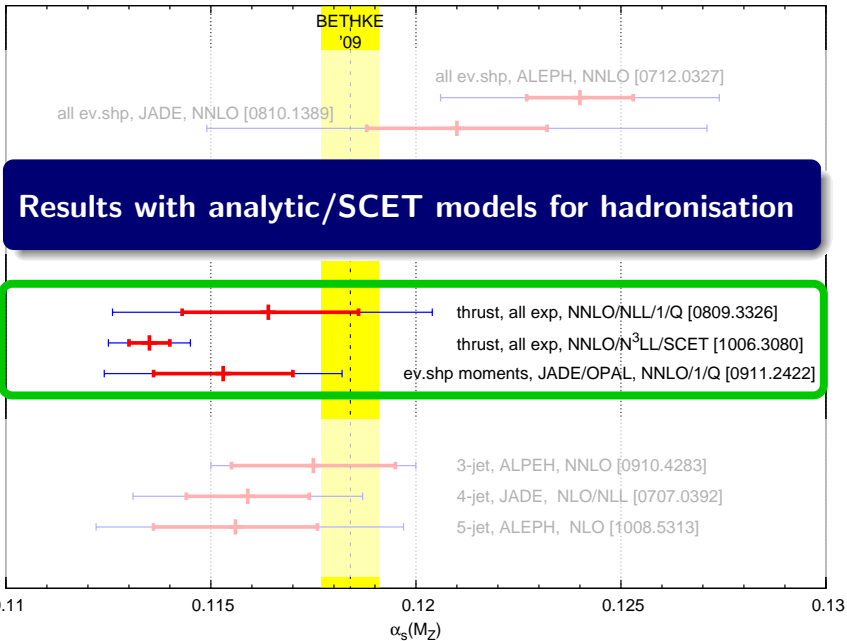
Bethke et al '08

In these results, an $\mathcal{O}(5 - 10\%)$ contribution to α_s comes from **hadronisation corrections**.

Determined by taking Pythia/Herwig/Ariadne MCs and looking at ratio of parton-level event-shape distribution to hadron-level distribution.

But who says MC parton level is in any way related to NNLO/NNLA parton level?

E.g. in MC's \propto radiation $\propto \alpha_s(p_t)$ with a 1 GeV cutoff
In NNLO, radiation $\propto \alpha_s(\mu) + b_0\alpha_s^2(\mu) \ln \mu/p_t + \dots$ with no cutoff



What matters for thrust (etc.) is how hadronisation modifies the soft energy flow at large angles. Parametrise with effective coupling

$$\alpha_0 = \frac{1}{\mu_I} \int_0^{\mu_I} d\mu \alpha_{\text{eff}}(\mu)$$

Argue thrust shifted by amount proportional to α_0

$$\delta T = -\frac{8C_F \mathcal{M}}{\pi^2} \frac{\mu_I}{Q} \left[\underbrace{\alpha_0 - \alpha_s(\mu_R) - \alpha_s^2(\mu_R^2) \left(\ln \frac{\mu_R}{\mu_I} + \dots \right) - \alpha_s^3(\mu_R^2) (\dots)}_{\text{subtract "perturbative" NNLO infrared piece}} \right]$$

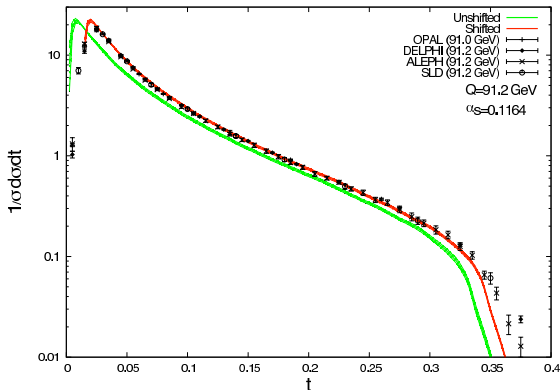
i.e.

$$\frac{d\sigma}{dT}(T) \rightarrow \frac{d\sigma}{dT}(T - \delta T)$$

Dokshitzer & Webber '95, '97 + numerous related contributions
 [Possibly] Valid in 2-jet limit; but seems to work also near 3-jet limit

$$\alpha_s(M_Z) = 0.1164 \pm 0.0022_{\text{exp}} \pm 0.0031_{\text{theo}}$$

Davison & Webber 0809.3326



Results based on all LEP and low-energy data

$$\chi^2/d.o.f. = 1.09$$

Experimental error corresponds to $\delta\chi^2 \simeq 14$ on grounds that fit not perfect.

N^3 LL resummation recently performed in context of Soft Collinear Effective Theory (SCET) for thrust and heavy-jet mass

Becher & Schwartz '08; Chien & Schwartz '10

SCET also argued to provide operator definition of hadronisation and rigorous factorisation (including subtraction terms shown for effective coupling)

Choice of which operators to keep/neglect requires skill?

SCET provides event-shape α_s extraction with by far the smallest errors:

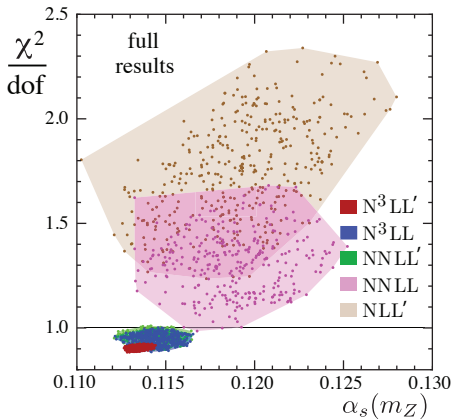
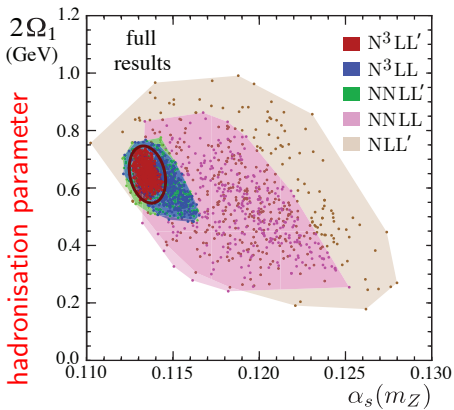
$$\alpha_s(M_Z) = 0.1135 \pm 0.0002_{\text{expt}} \pm 0.0005_{\text{hadr}} \pm 0.0009_{\text{pert}}$$

Abbate et al, arXiv:1006.3080

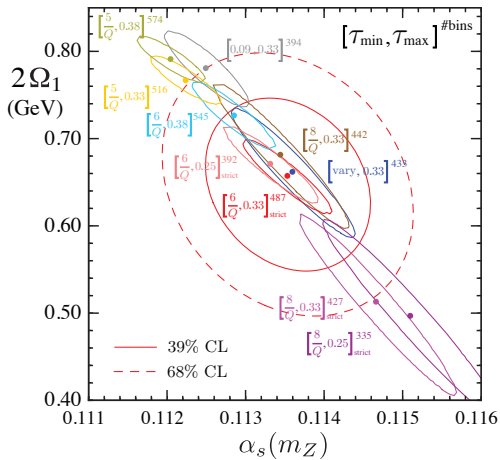
Davison & Webber had larger exp./hadr. err. (similar dataset); related to $\delta\chi^2$?

**Not compatible with world average. Similar to *some* DIS extractions
Is it correct?**

Good converge and fit quality



Dependence on fit range



Suggestion of systematic dependence \gtrsim quoted theory error.

Span of results is 0.003

Particular sensitivity to lower edge

What is basis for specific central choice?

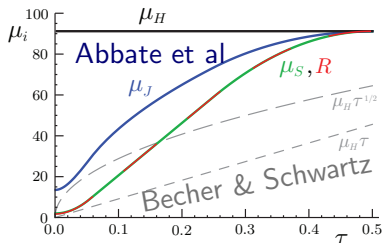
Choice of “profile” function

Experiment	Energy	BS results [20]	our BS profile	default profile
ALEPH	91.2 GeV	0.1168(1)	0.1170	0.1223
ALEPH	133 GeV	0.1183(37)	0.1187	0.1235
ALEPH	161 GeV	0.1263(70)	0.1270	0.1328
• • • •	• • • •	• • • •	• • • •	• • • •
average		0.1172(10)	0.1180	0.1221
global fit (stat)	all Q		0.1188	0.1242
global fit (stat+syst)	all Q		0.1192	0.1245

TABLE VIII: Comparison of the results for $\alpha_s(m_Z)$ quoted by Becher and Schwartz in Ref. [20] with results we obtain from our adapted code [...] and employ their profile functions for the nonsingular, hard, jet and soft scales, with results shown in the column labeled “our BS profile”. In the last column we show results with this same code, but using our default profile functions.

0.005 difference \gg theory error

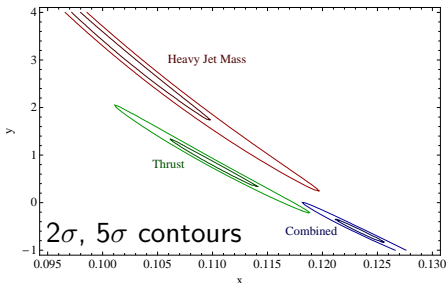
Schwartz agrees his orig. choice not right.
But is Abbate choice sufficiently broad?



Profile function determines

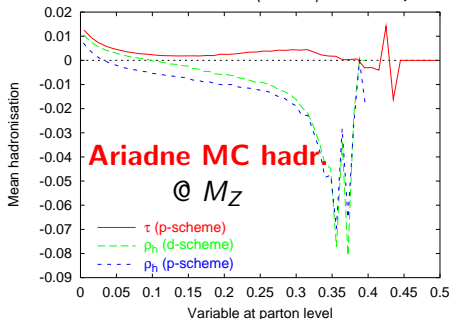
- ▶ Weight of fixed-order v. resummation
- ▶ Scales in resummation [\sim modified logs of CTTW]

hadronisation systematics?



Chien & Schwartz arXiv:1005.1644
 But also Gardi & Rathsman '03

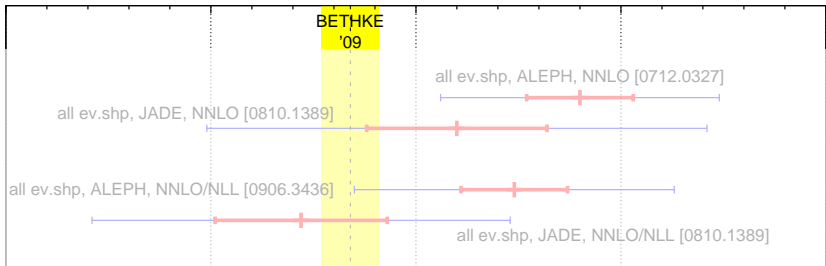
hadronisation $\langle \text{shift} \rangle$ v. τ, ρ



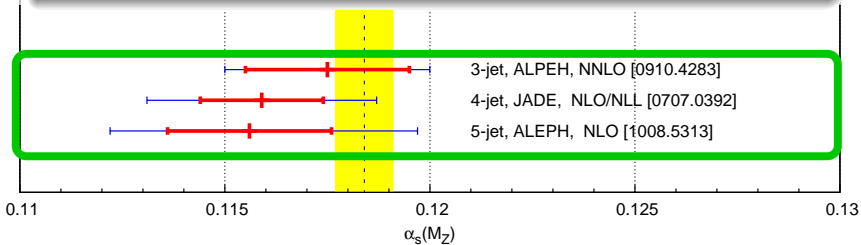
GPS & Wicke '01

- ▶ Jet masses' hadronisation very sensitive to π , K and p/n masses
- ▶ Heavy/wide observables are tricky (esp. for dists, higher moments).
- ▶ Do related systematics propagate also into thrust?

Monte Carlo hadronisation gets these things "right"



Event shapes take all hadrons. **Jet rates** v. y_{cut} use only a subset \rightarrow less sensitive to hadronisation?



	LEP1, hadr. $\sigma_{\text{tot}}^{-1} d\sigma/dy_{45}, R_5$	LEP1, no hadr. $\sigma_{\text{tot}}^{-1} d\sigma/dy_{45}, R_5$
stat.	+0.0002	+0.0002
	-0.0002	-0.0002
syst.	+0.0027	+0.0027
	-0.0029	-0.0029
pert.	+0.0062	+0.0068
	-0.0043	-0.0047
fit range	+0.0014	+0.0005
	-0.0014	-0.0005
hadr.	+0.0012	-
	-0.0012	-
$\alpha_s(M_Z)$	0.1159 ^{+0.0070} -0.0055	0.1163 ^{+0.0073} -0.0055

Hadronisation with Sherpa (multijet-matched) is small effect.

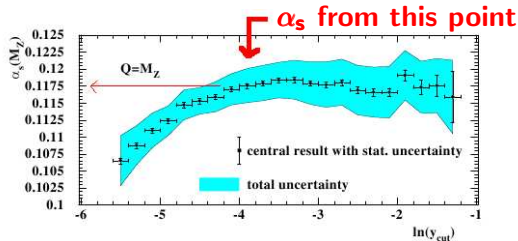
→ do not correct for it
NB: depends on MC

Final result with LEP1 & LEP2 data:

$$\alpha_s(M_Z) = 0.1156^{+0.0041}_{-0.0034}$$

Frederix et al arXiv:1008.5312

Table 2: Values of the strong coupling constant $\alpha_s(M_Z)$ obtained from fits to ALEPH LEP1 data for $\sigma_{\text{tot}}^{-1} d\sigma/dy_{45}$ and R_5 . NLO QCD predictions are used. Hadronization corrections are estimated with SHERPA. Default fit ranges are $3.8 \leq -\ln y_{45} \leq 5.2$, and $4.0 \leq -\ln y_{\text{cut}} \leq 5.6$. See the text for details.



3-jet rate as a function of y_{cut}

Measure α_s from point where

- ▶ determination stable
- ▶ systematics small

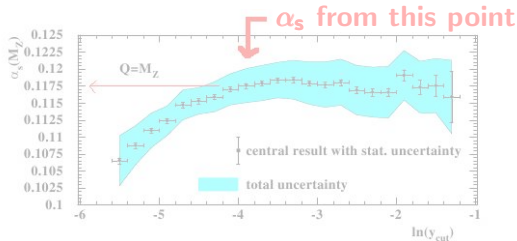
From ALEPH LEP 1 data:

$$\alpha_s(M_Z) = 0.1175 \pm 0.0020_{\text{exp}} \pm 0.0015_{\text{theo}}$$

Dissertori et al., arXiv:0910.4283

- ▶ Aside from SCET analysis, this is measurement with smallest theory error
- ▶ Can exp. error be reduced by combining other LEP experiments' data?
- ▶ Total effect of hadronisation is 0.0015. Adding this (conservatively) as hadronisation error, get $\alpha_s = 0.1175 \pm 0.0029$

Would other jet definitions (e.g. Cambridge) have yet smaller hadronisation?



3-jet rate as a function of y_{cut}

Measure α_s from point where

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Dissertori et al., arXiv:0910.4283

▶ **Interesting combination of modest errors and simplicity of analysis**

hadronisation error, get $\alpha_s = 0.1175 \pm 0.0029$

Would other jet definitions (e.g. Cambridge) have yet smaller hadronisation?

Situation with event shapes is not ideal

Benefits of NNLO somewhat counteracted by large hadronisation effects and limited perturbative convergence

Best “simple” analysis might be from jet rates?

SCET with N³LL has small errors, but disturbing α_s .
How well controlled are the systematics?

EXTRAS

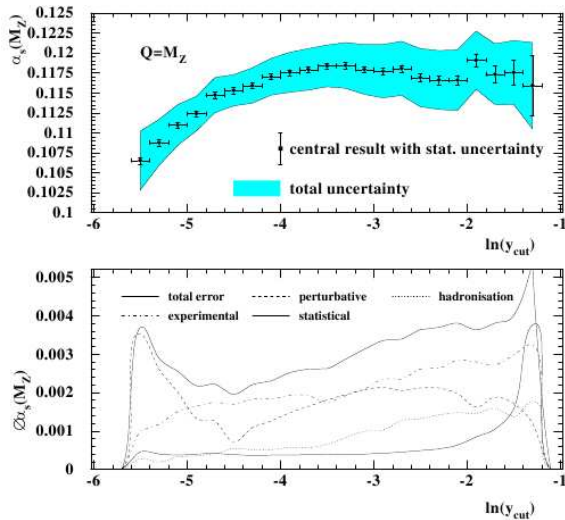


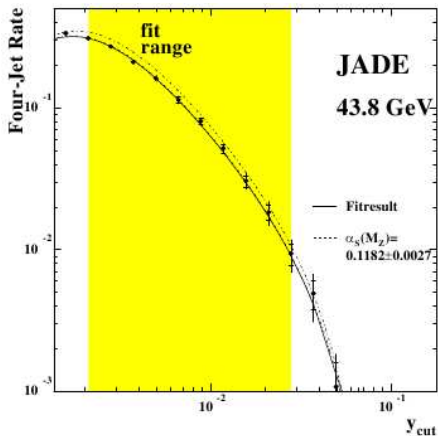
FIG. 1: Determinations of $\alpha_s(M_Z)$ from the three-jet rate, measured by ALEPH at the Z peak, for several values of the jet-resolution parameter y_{cut} . The error bars show the

From LEP 1:

$$\alpha_s(M_Z) = 0.1175 \pm 0.0020_{\text{exp}} \pm 0.0015_{\text{theo}}$$

- ▶ Hadronisation is $\sim 5\%$ correction
- ▶ α_s extracted at $y = 0.02$
- ▶ LEP2 consistent

$$\alpha_s(M_Z) = 0.1159 \pm 0.0004_{\text{stat}} \pm 0.0012_{\text{exp}} \pm 0.0024_{\text{had}} \pm 0.0007_{\text{theo}}$$



- ▶ Uses R_4 (exclusive?)
- ▶ fit ranges chosen by “requiring that hadronization corrections be less than 50% and the detector corrections be less than 50% in the fit range”

$$\alpha_s(M_Z) = 0.1156^{+0.0041}_{-0.0034}$$

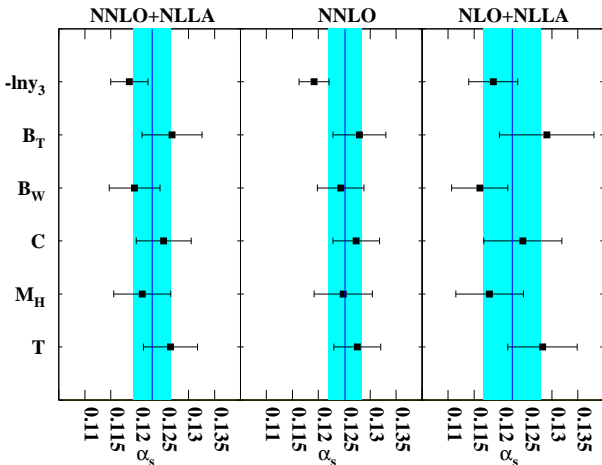
	LEP1, hadr. $\sigma_{\text{tot}}^{-1} d\sigma/dy_{45}, R_5$	LEP1, no hadr. $\sigma_{\text{tot}}^{-1} d\sigma/dy_{45}, R_5$
stat.	+0.0002 -0.0002	+0.0002 -0.0002
syst.	+0.0027 -0.0029	+0.0027 -0.0029
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fit range	+0.0014 -0.0014	+0.0005 -0.0005
hadr.	+0.0012 -0.0012	-
$\alpha_s(M_Z)$	0.1159 ^{+0.0070} -0.0055	0.1163 ^{+0.0073} -0.0055

Table 2: Values of the strong coupling constant $\alpha_s(M_Z)$ obtained from fits to ALEPH LEP1 data for $\sigma_{\text{tot}}^{-1} d\sigma/dy_{45}$ and R_5 . NLO QCD predictions are used. Hadronization corrections are estimated with SHERPA. Default fit ranges are $3.8 \leq -\ln y_{45} \leq 5.2$, and $4.0 \leq -\ln y_{\text{cut}} \leq 5.6$. See the text for details.

	LEP2, no hadr. $\sigma_{\text{tot}}^{-1} d\sigma/dy_{45}$	LEP2, no hadr. R_5	LEP2, no hadr. $\sigma_{\text{tot}}^{-1} d\sigma/dy_{45}, R_5$
stat.	+0.0020 -0.0022	+0.0022 -0.0025	+0.0015 -0.0016
syst.	+0.0008 -0.0009	+0.0012 -0.0012	+0.0008 -0.0008
pert.	+0.0049 -0.0034	+0.0029 -0.0020	+0.0029 -0.0020
fit range	+0.0038 -0.0038	+0.0030 -0.0030	+0.0028 -0.0028
$\alpha_s(M_Z)$	0.1189 ^{+0.0066} -0.0057	0.1120 ^{+0.0050} -0.0047	0.1155 ^{+0.0044} -0.0039

Table 3: Values of the strong coupling constant $\alpha_s(M_Z)$ obtained from fits to ALEPH LEP2 data with $E_{\text{cm}} \geq 183$ GeV for $\sigma_{\text{tot}}^{-1} d\sigma/dy_{45}$ and R_5 . NLO QCD predictions are used. Hadronization corrections are not included. Default fit ranges are $4.8 \leq -\ln y_{45} \leq 6.4$, and $2.1 \leq -\log_{10} y_{\text{cut}} \leq 2.9$. See the text for details.

$$\alpha_s(M_Z) = 0.1224 \pm 0.0009_{\text{stat}} \pm 0.0009_{\text{exp}} \pm 0.0012_{\text{had}} \pm 0.0035_{\text{theo}}$$

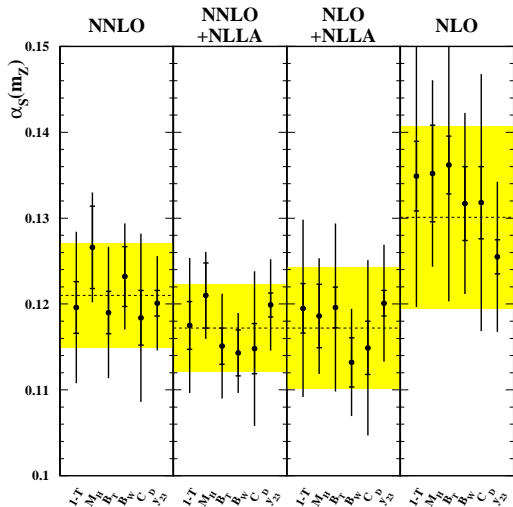


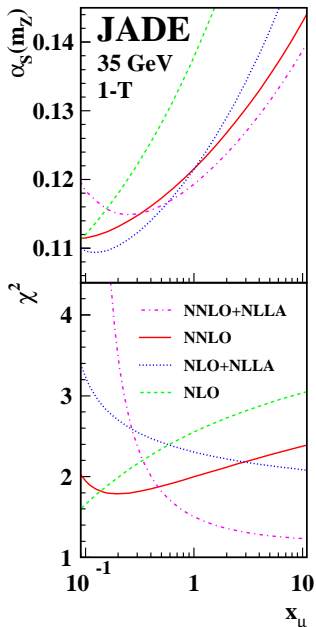
Results using Herwig++ + POWHEG for hadr. are 3% lower

NNLO : $\alpha_s(M_Z) = 0.1210 \pm 0.0007_{\text{stat}} \pm 0.0021_{\text{exp}} \pm 0.0044_{\text{had}} \pm 0.0036_{\text{theo}}$

NNLO+NLL : $\alpha_s(M_Z) = 0.1172 \pm 0.0006_{\text{stat}} \pm 0.0020_{\text{exp}} \pm 0.0035_{\text{had}} \pm 0.0030_{\text{theo}}$

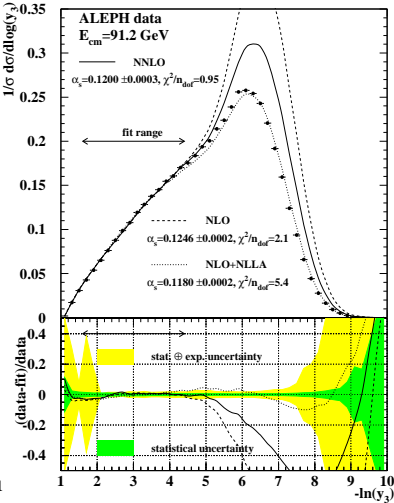
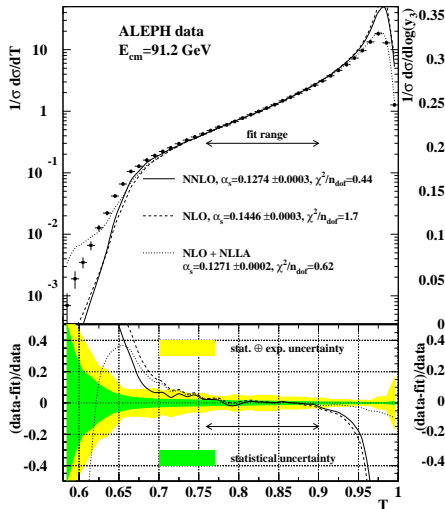
- ▶ Fit range chosen such that LL terms less than unity; same range used for NNLO & NNLO+NLL



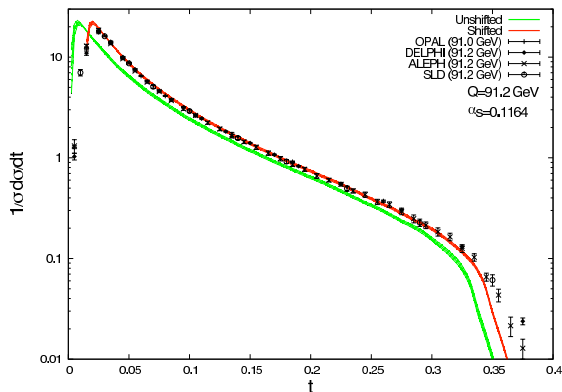


Illustrates clear improvement coming from inclusion of NNLO/NLL information

$$\alpha_s(M_Z) = 0.1240 \pm 0.0008_{\text{stat}} \pm 0.0010_{\text{exp}} \pm 0.0011_{\text{had}} \pm 0.0029_{\text{theo}}$$

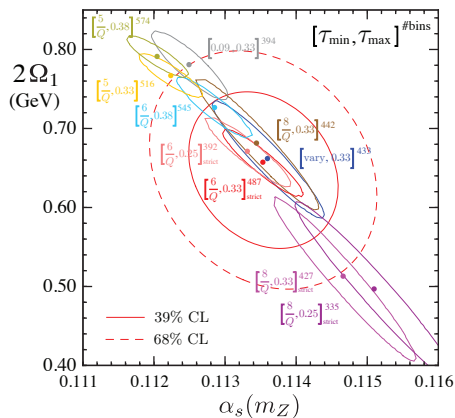


$$\alpha_s(M_Z) = 0.1164 \pm 0.0022_{\text{exp}} \pm 0.00157_{\text{theo}}$$



Note, scales varied in range
 $\sqrt{1/2} < x_\mu < \sqrt{2}$
 With a conventional $\frac{1}{2} < x_\mu < 2$ variation, expect theory error to double roughly?

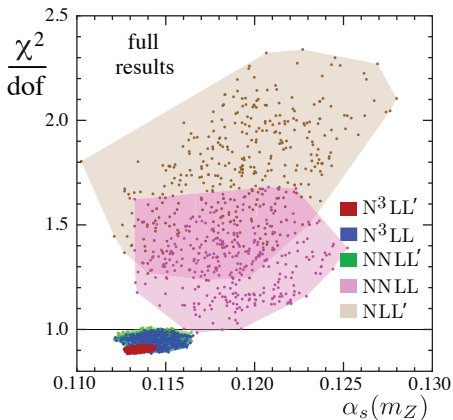
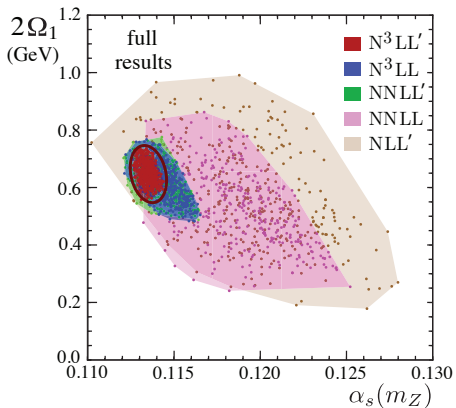
$$\alpha_s(M_Z) = 0.1135 \pm 0.0002_{\text{expt}} \pm 0.0005_{\text{hadr}} \pm 0.0009_{\text{pert}}$$



Experiment	Energy	BS results [20]	our BS profile	default profile
ALEPH	91.2 GeV	0.1168(1)	0.1170	0.1223
ALEPH	133 GeV	0.1183(37)	0.1187	0.1235
ALEPH	161 GeV	0.1263(70)	0.1270	0.1328
• • • • •		• • • • •	• • • • •	• • • • •
average		0.1172(10)	0.1180	0.1221
global fit (stat)	all Q		0.1188	0.1242
global fit (stat+syst)	all Q		0.1192	0.1245

TABLE VIII: Comparison of the results for $\alpha_s(m_Z)$ quoted by Becher and Schwartz in Ref. [20] with results we obtain from our adapted code [...] and employ their profile functions for the nonsingular, hard, jet and soft scales, with results shown in the column labeled “our BS profile”. In the last column we show results with this same code, but using our default profile functions.

profile function matters at $\sim 4\%$



Assuming no hadronisation (ρ): $\alpha_s(M_Z) = 0.1220 \pm 0.0031$

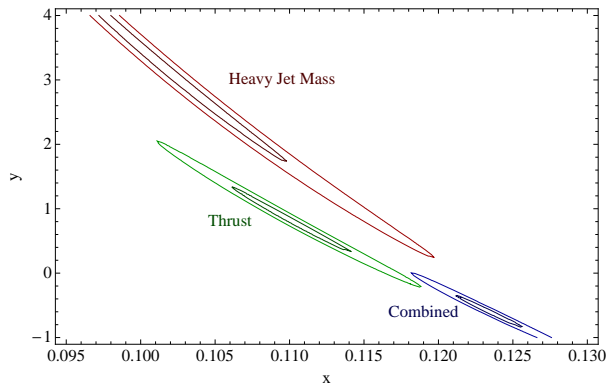
Assuming no hadronisation (T): $\alpha_s(M_Z) = 0.1193 \pm 0.0027$

Combined:

$$\alpha_s(M_Z) = 0.1193 \pm 0.0011_{\text{stat}} \pm 0.0012_{\text{syst}} \pm 0.0017_{\text{had}} \pm 0.0013_{\text{pert}} \pm 0.0005_{\text{soft}}$$

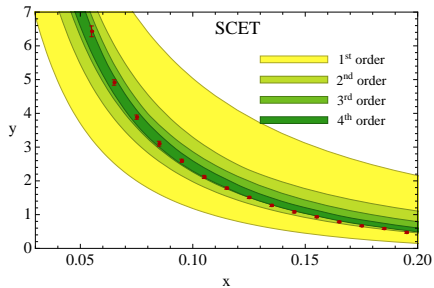
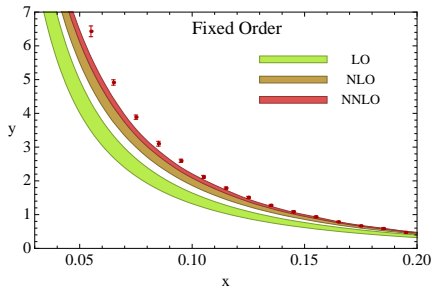
With hadronisation as shift (ρ): $\alpha_s(M_Z) = 0.1017$

With hadronisation as shift (T): $\alpha_s(M_Z) = 0.1101$

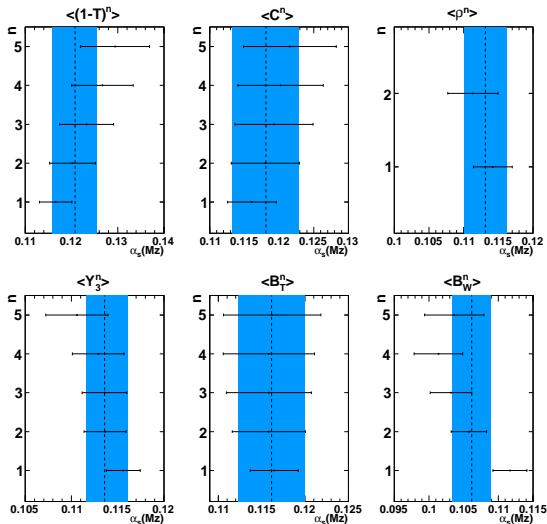


Is fit without hadronisation degrees of freedom reasonable?

Discrepancy between ρ and Thrust in power-correction fits has a long history.



$$\alpha_s(M_Z) = 0.1153 \pm 0.0017_{\text{exp}} \pm 0.0023_{\text{th}}$$



- ▶ PC formulae validity arguable for higher moments
- ▶ hadron-mass effects for the heavy jet mass
- ▶ Why not extract results based on first moments?