#### Giant K factors

#### Gavin Salam

CERN, Princeton & LPTHE/CNRS (Paris)

Work performed with Mathieu Rubin and Sebastian Sapeta, arXiv:1006.2144

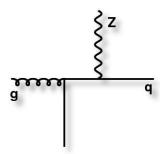
Workshop on Heavy Particles at the LHC Pauli Center, Zürich, 5–7 January 2011

Many searches for New Physics (eg. SUSY) rely on Leading Order predictions for backgrounds.

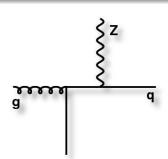
eg. Z+4jet background to gluino pair production with NLO technology rapidly becoming mature for such cases

LO often considered good to within a factor of 2 NLO to within 10-20%

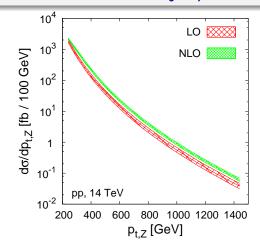
This talk is about cases where such "rules of thumb" fail (spectacularly)



Use MCFM to examine various properties of such events at LO and NLO.



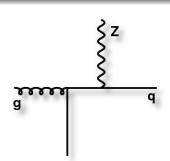
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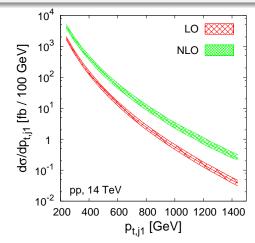
 $p_t$  spectrum of **Z boson** gets K-factor of 1.5

Fairly standard kind of occurrence

## The Z+jet process

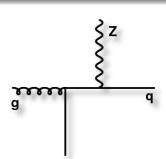


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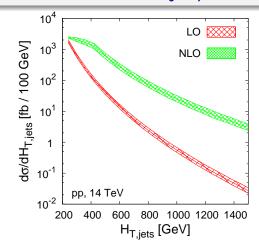


#### $p_t$ spectrum of **leading jet** gets K-factor of 5–10

related issues in Butterworth, Davison, Rubin & GPS '08 Bauer & Lange '09; Denner, Dittmaier, Kasprzik & Much '09

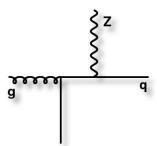


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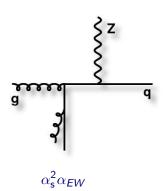
$$H_{T,jets} \equiv \sum_{i \in iets} p_{t,i}$$
 gets K-factor of up to 100

Such things are not supposed to happen with  $\alpha_s = 0.1!$ 



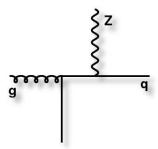
 $\alpha_{\mathsf{s}}\alpha_{\mathsf{EW}}$ 

#### **Next-to-Leading Order**



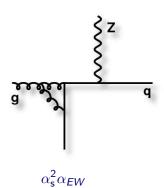
LHC probes scales well above EW scale,  $\sqrt{s} \gg M_Z$ . EW bosons are **light**. New log-enhanced topologies appear.

 $H_{T,jets}$  is extreme, because at LO  $H_{T,jets} \simeq p_{t,jet 1}$ ; NLO:  $H_{T,jets} \simeq 2p_{t,jet}$ 



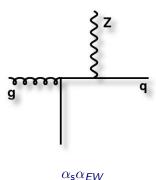
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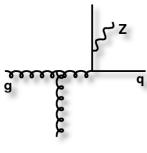


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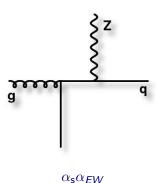
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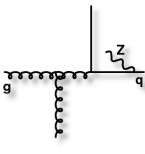
$$\alpha_s^2 \alpha_{EW} \ln^2 \frac{p_t}{M_Z}$$

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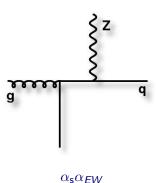
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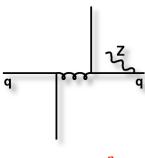
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$$\alpha_{\rm s}^2 \alpha_{EW} \ln^2 \frac{p_t}{M_7}$$

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Though we calculate Z+jet@NLO, giant K-factors really dominated by ("LO") Z+2-parton piece of Z+jet@NLO.

We know LO calculations aren't reliable.

We really want to combine Z+jet@NLO with Z+2-jet@NLO

without double counting

without having to do full Z+jet@NNLO calculation

# First try the following:

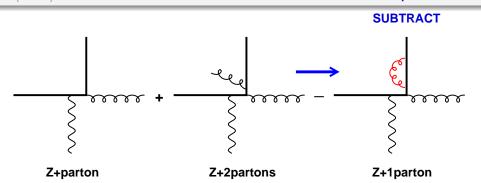
Take the "leading" process 
$$[Z + jet @ LO]$$

and add in process with one extra jet.

[i.e. include Z + 2 jets @ LO]

**approximate** the 1-loop Z+jet term, by requiring cancellation of all divergences

#### cartoon of the LoopSim idea



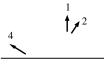
softest particle of Z+2 is "looped" = removed from event (kinematics reshuffled)

- ▶ For every Z + 2 parton (2  $\rightarrow$  3) event, figure out what what 2  $\rightarrow$  2 event it would really have come from "Loop" the softest parton [Don't actually explicitly calculate any loop diagrams: simulate the loops]
- Subtract that 2 → 2 event

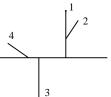
Unlike MLM, no cutoffs on  $2 \rightarrow 3$  events

## cartoon of LoopSim "looping" procedure

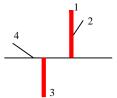
(a) Input event



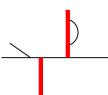
(b) Attributed emission seq.



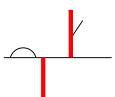
(c) Born particle ID



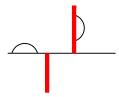
(d) Output 1-loop event



(e) 2nd output 1-loop event



(f) Output 2-loop event



- ▶ Use jet algorithm to assign a branching structure to event à la CKKW
- ► The particles that are softest are the ones that will be "looped"

## Define operators:

 $U_{\ell}(\text{event E}) \equiv \text{all simulated } \ell\text{-loop events from E}$ 

$$U_{orall}(\mathsf{event}) \equiv \sum_{\ell=0} U_\ell(\mathsf{event})$$

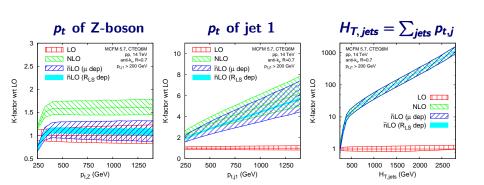
"U" stands for unitarisation (cancellation of all divergences) sum of all diagrams (essentially) adds up to zero

To combine Z+j with Z+2j take

$$Z+j@\bar{n}LO \equiv Z+j@LO + U_{\forall}(Z+2j@LO)$$

we use " $\bar{n}$ LO" to emphasize that this is a crude approximation to an actual NLO calculation — the exact loops are missing NB:  $U_{\forall}$  here includes  $\ell=0,1$ 

## $\bar{n}$ LO results (*K*-factors, normalised to LO)



When the K-factors are large,  $\bar{n}LO$  agrees well with NLO

MLM matching also does a similar job

# Differences between and LoopSim and MLM/CKKW matching:

- Does not rely on shower (✓: simple; ✗: not easily integrated with shower MCs)
- 2. Does not need arbitrary separation of Z+1/Z+2/etc. samples with (hard-to-choose) momentum cutoff
  - 3. Can easily be extended beyond LO matching

## $\bar{n}$ NLO: merging Z+j and Z+2j, both @NLO

Just replace simulated loops with exact loops Apply LoopSim to exact 1-loop to get (e.g.) simulated 2-loop terms

$$E_{n,\ell} \equiv$$
 event with  $n$  partons and  $\ell$  exact loops  $U_{orall,\ell} \equiv$  operator to apply when  $\ell$  exact loops known  $U_{orall,1}(E_{n,0}) = U_{orall}(E_{n,0}) - U_{orall}(U_1(E_{n,0}))$   $U_{orall,1}(E_{n,1}) = U_{orall}(E_{n,1})$ 

$$\mathsf{Z} + \mathsf{j}@\mathsf{\bar{n}}\mathsf{NLO} = \mathsf{Z} + \mathsf{j}@\mathsf{NLO} + U_{\forall,1}(\mathsf{Z} + 2\mathsf{j}@\mathsf{NLO}_{\mathsf{only}})$$

Extension to NLO, NNLO, multi-leg, etc. is almost trivial in LoopSim

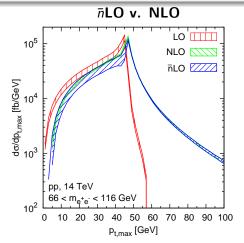
Not the case in methods that merge with parton showers too

## Testing NLO Merging, in 3 processes

- 1. Z@NLO with Z+j@NLO
- 2. Z+j@NLO with Z+2j@NLO
  - 3. 2j@NLO with 3j@NLO



# Validation: Drell-Yan lepton $p_t$ , $\bar{n}$ NLO v. NNLO

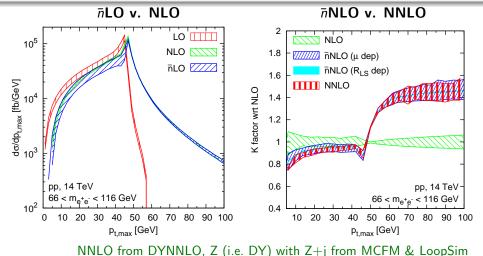


Z (i.e. DY) with  $Z{+}j$  from MCFM & LoopSim

For  $p_{tell} \gtrsim \frac{1}{2}M_Z + \Gamma_Z$  (giant K-factor!) it had to work For  $p_{t,\ell} \lesssim \frac{1}{2}M_Z + \Gamma_Z$  it's remarkable that it still works



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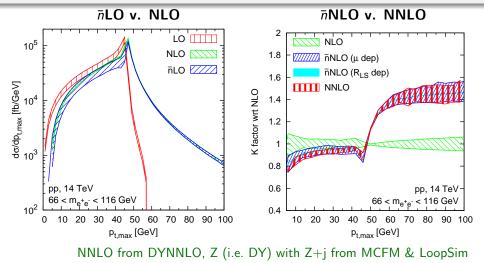


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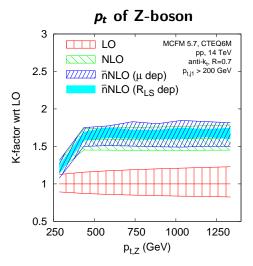
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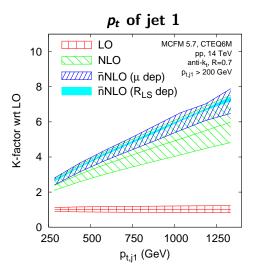


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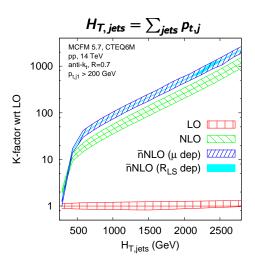


- ▶ p<sub>tZ</sub> distribution didn't have giant K-factors.
- ▶ *n*NLO brings no benefit

To get improvement you would need exact 2-loop terms



- ▶  $p_{tj}$  distribution seems to converge at  $\bar{n}$ NLO
- ▶ scale uncertainties reduced by ~ factor 2



- ► Signficiant further enhancement for *H*<sub>T.iets</sub>
- $ightharpoonup \bar{n}$ NLO brings clear message:

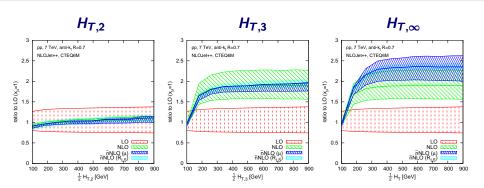
 $H_{T,jets}$  is not a good observable!

## $H_T$ (effective mass) type observables are widely used in searches

- ▶  $H_T$  has a steeply falling distribution (like  $p_{tj}$ ,  $p_{tZ}$ )
- ► At each order (NLO, NNLO), an extra (soft) jet contributes to the H<sub>T</sub> sum
  e.g. from ISR
- ▶ That shifts  $H_T$  up, which translates to a substantial increase in the cross section

We can test this hypothesis for plain jet events, using a truncated sum,

$$H_{T,n} = \sum_{i=1}^{n} p_{t, \text{jet } i}$$



A clear message: for a process with n objects at lowest order, use  $H_{T,n}$ 

Do you know what gets used in your experiment's searches?

Many writers of LHC SUSY proceedings didn't...

## Be aware that giant K-factors exist

Always look one order beyond the leading order, for example with  ${\sf MLM}/{\sf CKKW}\ {\sf matching}$ 

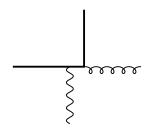
# New tool to get good predictions in such cases: LoopSim

Basically an "operator" to generate approximations to unknown loops Combine Z+j@NLO, Z+2j@NLO to get " $\bar{n}$ NLO" Z+jet It sometimes works even beyond "giant-K-factor" regions

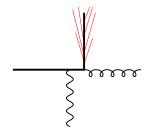
## Watch out for $H_T$

Even for simple processes, it converges very poorly unless you define it carefully (limit number of objects in sum)

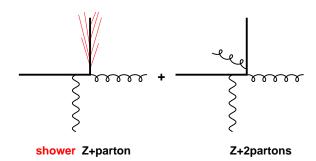
# **EXTRAS**

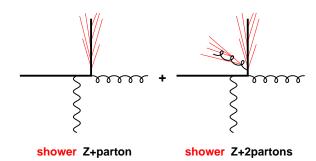


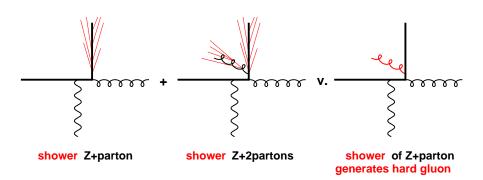
Z+parton

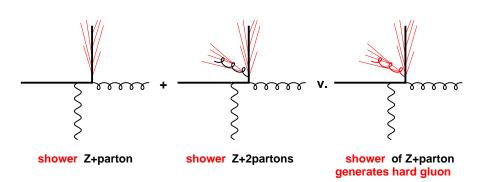


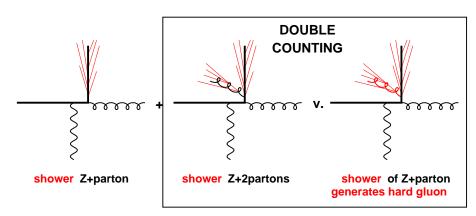
shower Z+parton





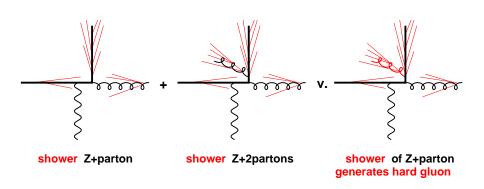




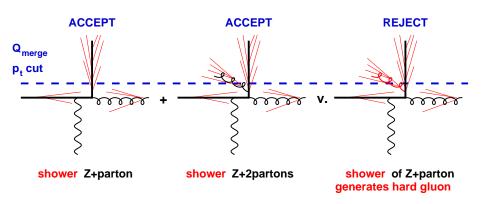


 $\mathsf{Z}$  + parton implicitly includes part of  $\mathsf{Z}$  + 2 partons It's just that the 2nd parton isn't always explicitly "visible"

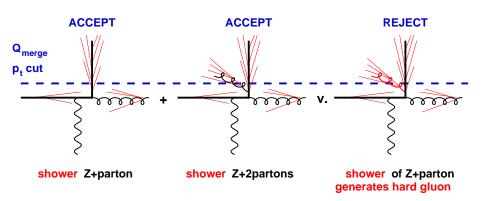
#### cartoon of MLM merging of Z+j and Z+2j



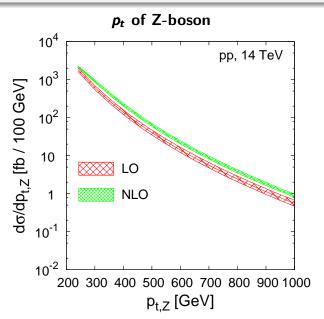
- ► MLM merging relies on parton shower to help figure out what fraction of Z + parton is really Z + 2 partons.
- ▶ Our aim is to do that without the parton shower

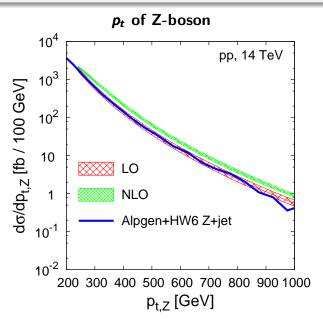


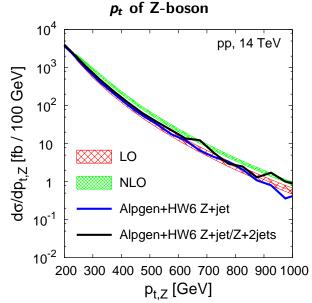
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All predictions similar and stable

