

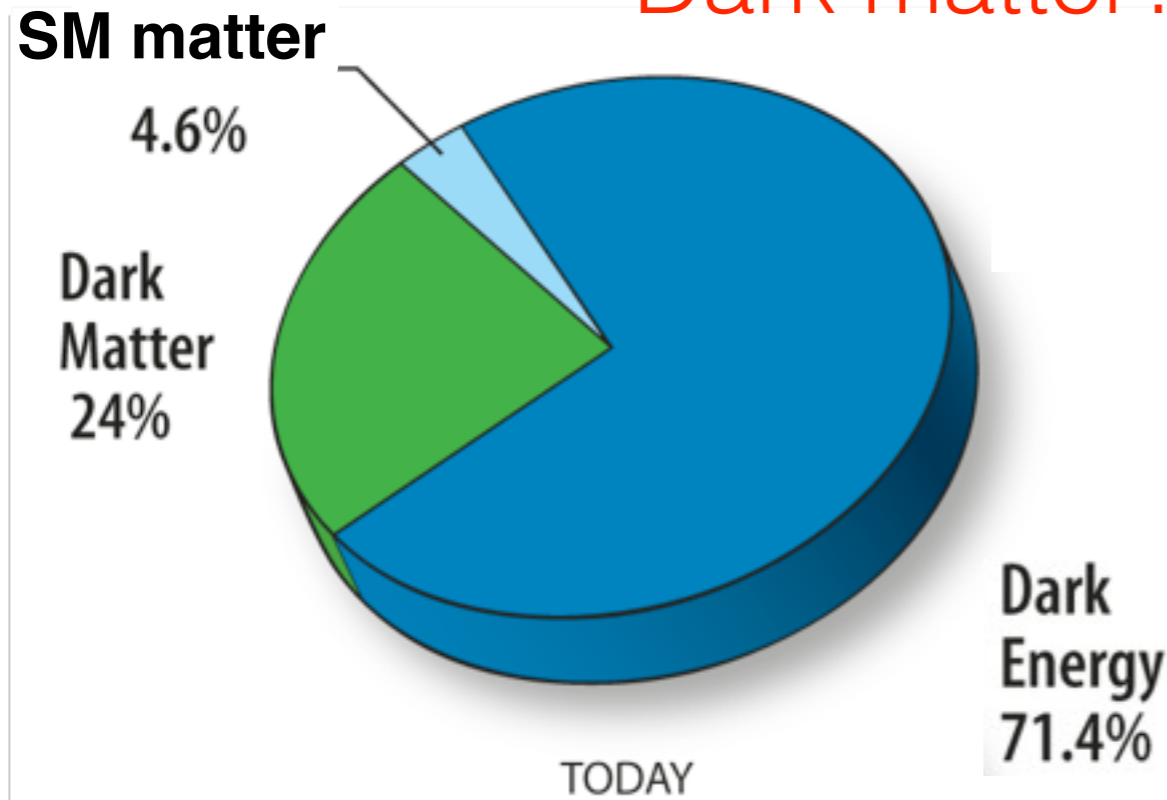
# Theory Perspectives on the HL-LHC

Andi Weiler (CERN & DESY)  
Gavin P. Salam (CERN)

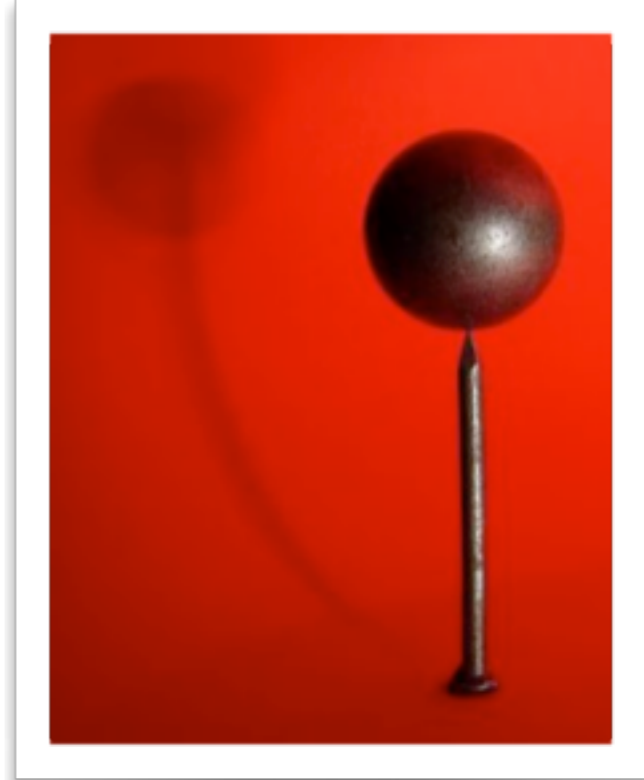
ECFA High Luminosity LHC Experiments Workshop  
Aix-les-Bains, France, 1–3 October 2013

# The SM is incomplete: big questions

## Dark matter?

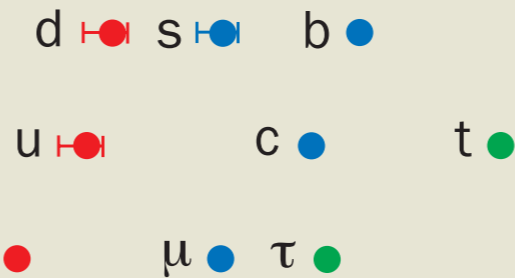


## Fine-tuning?



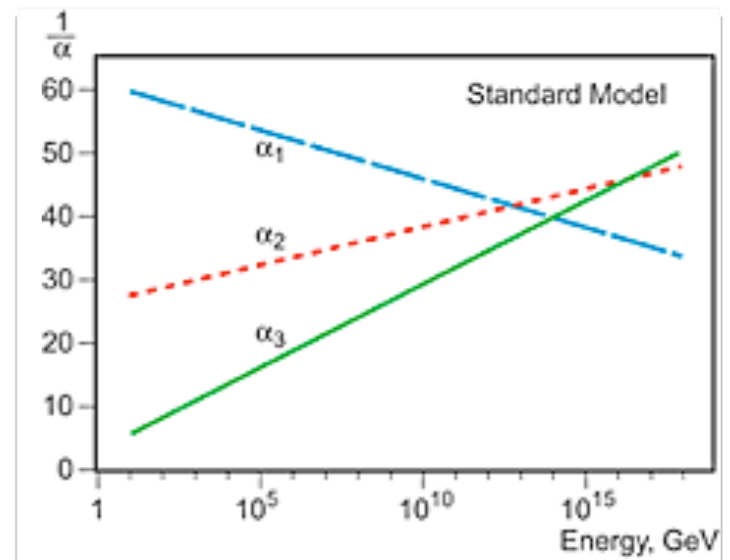
## Origin of SM matter and flavor?

$$Y_U \approx \begin{pmatrix} 10^{-5} & -0.002 & 0.007 + 0.004i \\ 10^{-6} & 0.007 & -0.04 + 0.0008i \\ 10^{-8} + 10^{-7}i & 0.0003 & 0.92 \end{pmatrix}$$



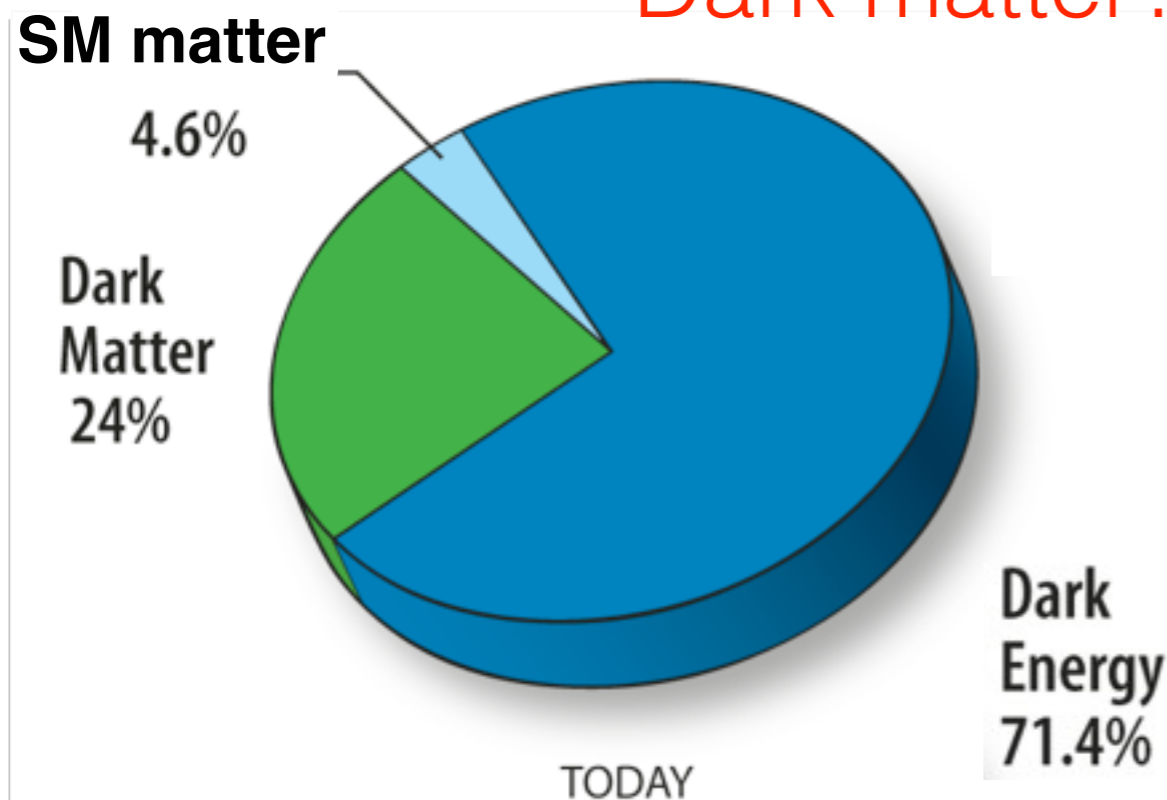
μeV    meV    eV    keV    MeV    GeV    TeV

## Unity of forces?

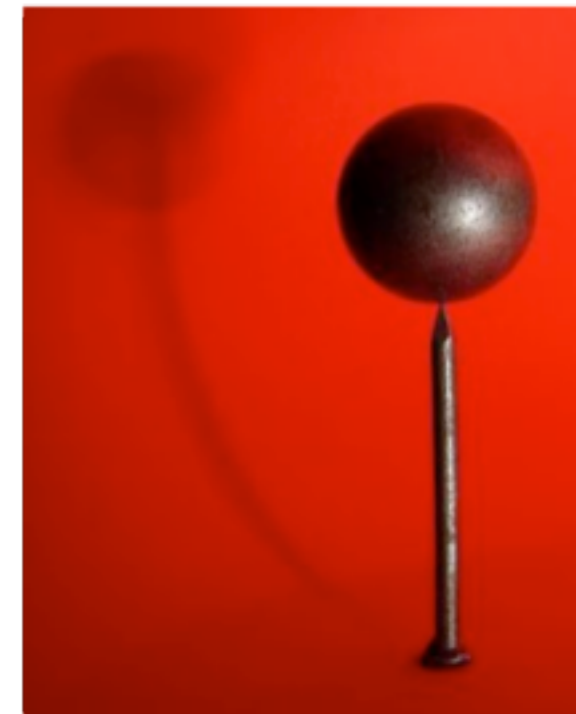


# The SM is incomplete: big questions

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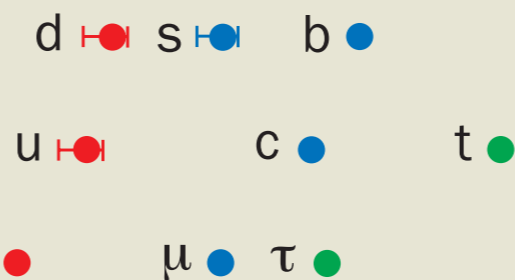


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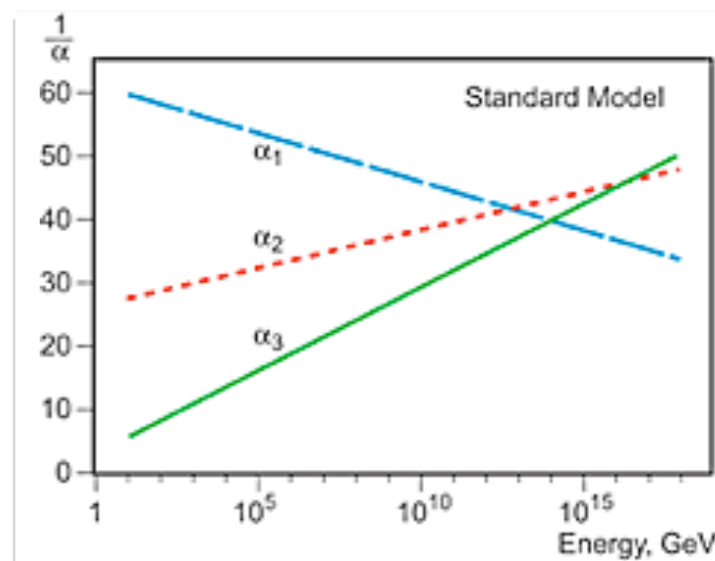


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## Unity of forces?



# Motivation for new physics at the TeV scale

The SM alone suffers from destabilizing quantum fluctuations sensitive to the highest scales in nature (e.g. GUT, Planck scale).

Given a fundamental theory at a high scale  $\Lambda$ , a light Higgs requires fine-tuning in the fundamental theory of order  $(m_H/\Lambda)^2$ .

Fine-tuning?



**Naturalness:** absence of special conspiracies between phenomena occurring at very different length scales.

A strong motivation for new physics at TeV colliders. Most interesting theories offer solutions to open problems of the SM.

What are the theory motivations for HL-LHC?  
I.e. what is  $300 \rightarrow 3000 \text{ fb}^{-1}$  good for?

This is not a report from a working group (so far there is no TH WG), but rather our survey of some of the main points, including consultations with colleagues.

1. Higgs sector in SM & beyond
2. Direct searches for New Physics
3. SM sector, including theory prospects  
(e.g. for use in Higgs physics)

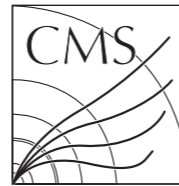
NB: given breadth of topics, references will be illustrative, not exhaustive

# 1. Higgs sector in SM & beyond

# What does 14 TeV @3000fb<sup>-1</sup> bring?

**ATLAS Preliminary (Simulation)**

$\sqrt{s} = 14$  TeV:  $\int L dt = 300 \text{ fb}^{-1}$ ;  $\int L dt = 3000 \text{ fb}^{-1}$   
 $\int L dt = 300 \text{ fb}^{-1}$  extrapolated from 7+8 TeV



arXiv:1307.7135

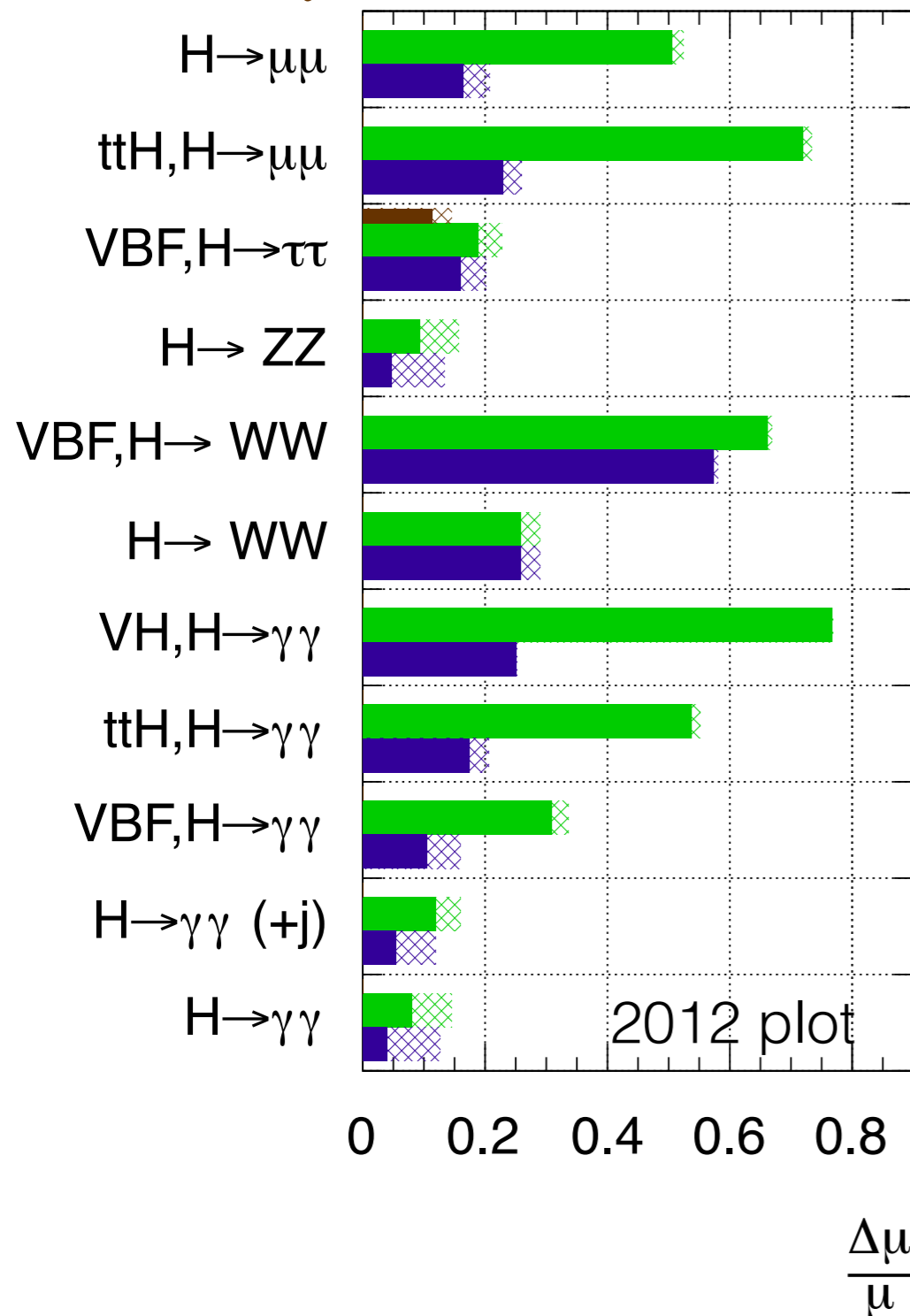


Table 2: Precision on the measurements of the signal strength for a SM-like Higgs boson. These values are obtained at  $\sqrt{s} = 14$  TeV using an integrated dataset of 300 and 3000  $\text{fb}^{-1}$ . Numbers in brackets are % uncertainties on the measurements estimated under [Scenario2, Scenario1], as described in the text. For the direct search for invisible Higgs decays the 95% CL on the branching fraction is given.

L (fb <sup>-1</sup> )	H → γγ	H → WW	H → ZZ	H → bb
300	[6, 12]	[6, 11]	[7, 11]	[11, 14]
3000	[4, 8]	[4, 7]	[4, 7]	[5, 7]

H → ττ	H → Zγ	H → inv.
[8, 14]	[62, 62]	[17, 28]
[5, 8]	[20, 24]	[6, 17]

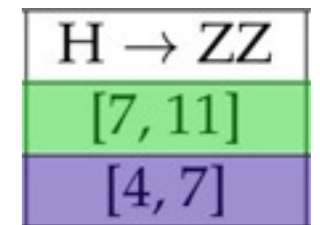
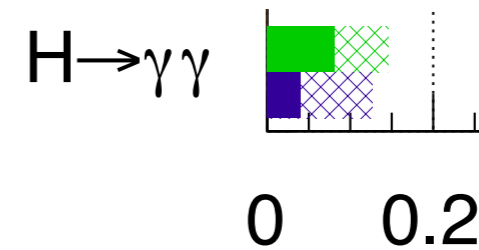


# What does 14 TeV @3000fb<sup>-1</sup> bring?

- Increased precision on existing channels
- Observation of rare decays  $Z\gamma$ ,  $\mu^+\mu^-$ ,  $cc$
- Double Higgs production
- Longitudinal Vector-Boson Scattering

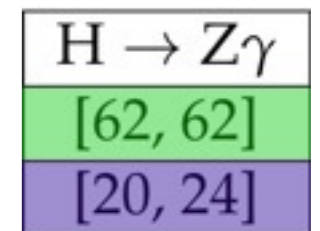
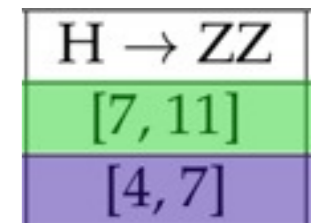
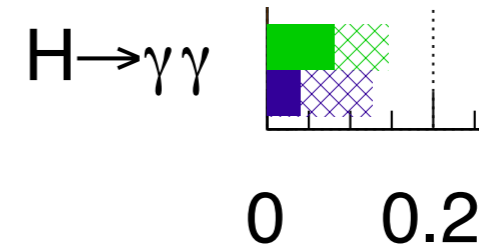
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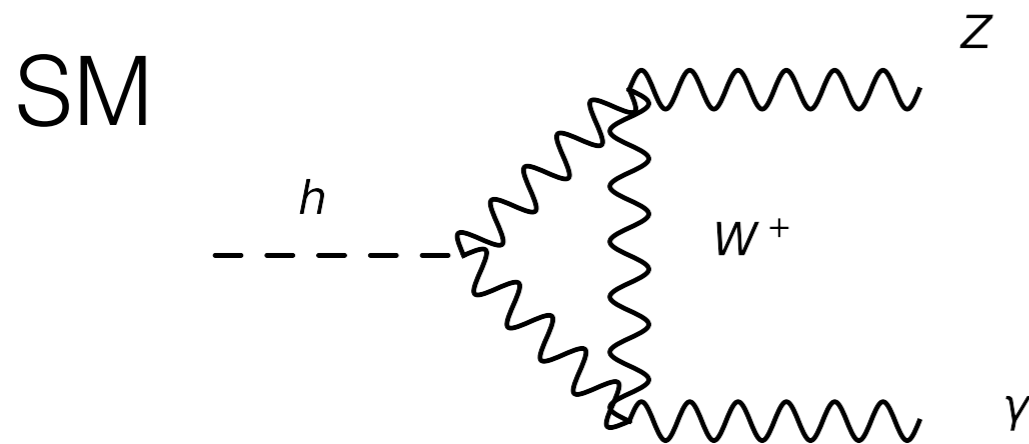
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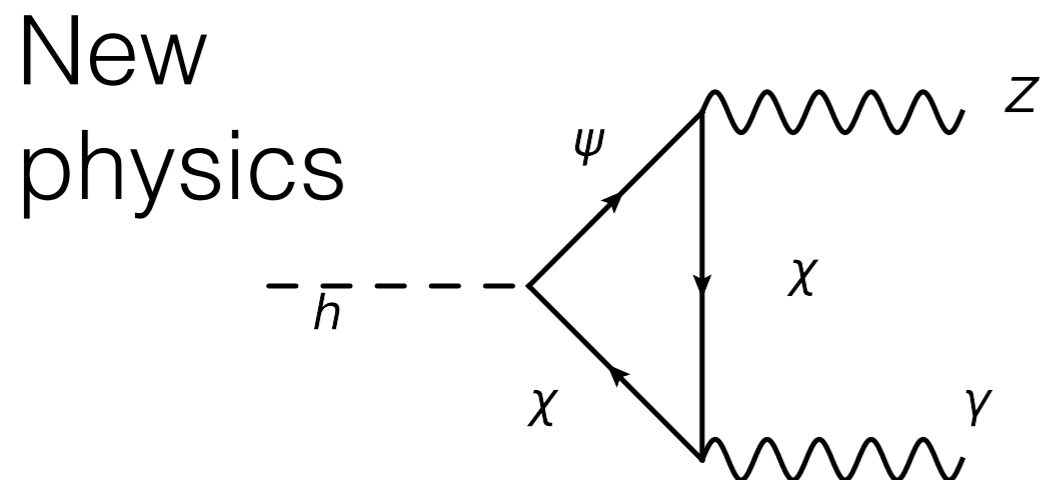
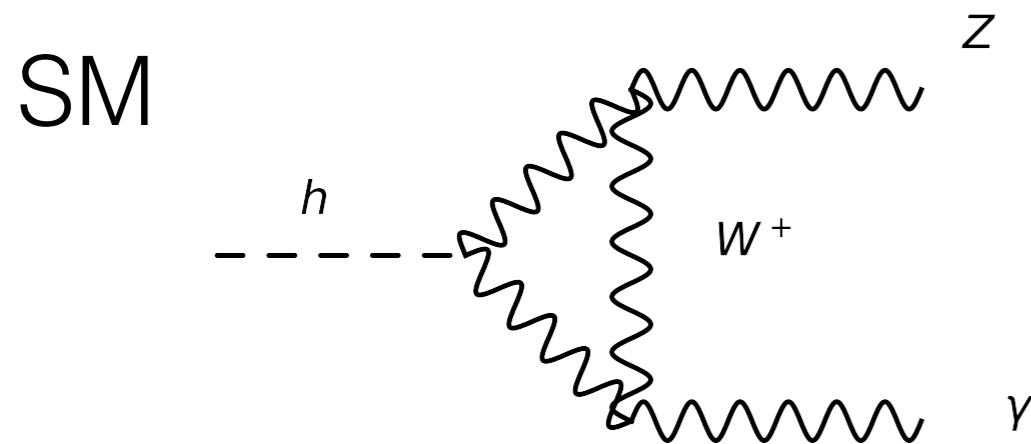
# Why $\gamma Z$ is important

- $\gamma Z$  like  $\gamma\gamma$  and  $gg$  loop induced, but sensitive to effects invisible in  $\gamma\gamma$  and  $gg$  (because of chiral couplings)
- In composite Higgs: Not protected by Goldstone symmetry, large  $\gamma Z$  while  $\gamma\gamma$  and  $gg$  small



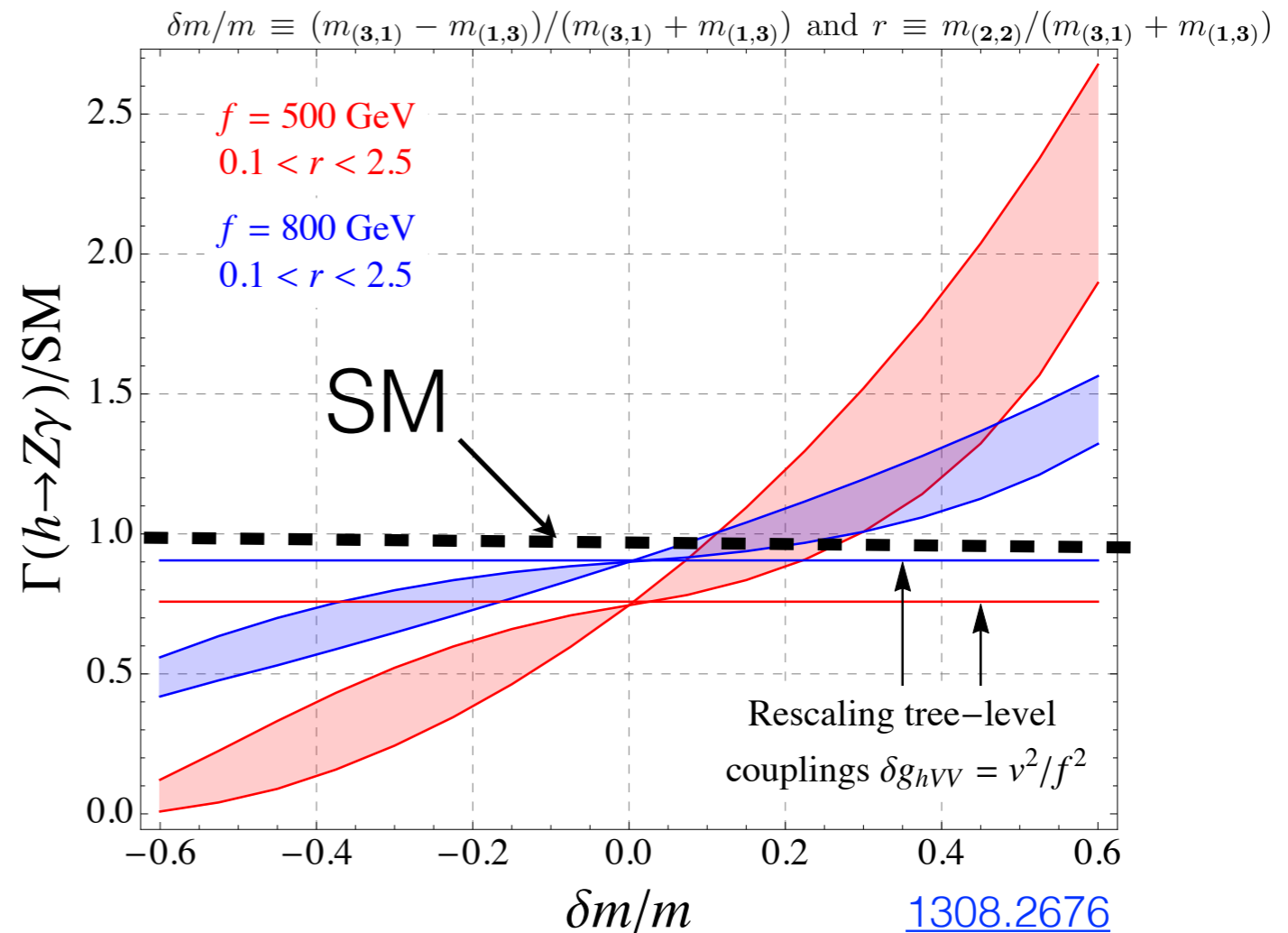
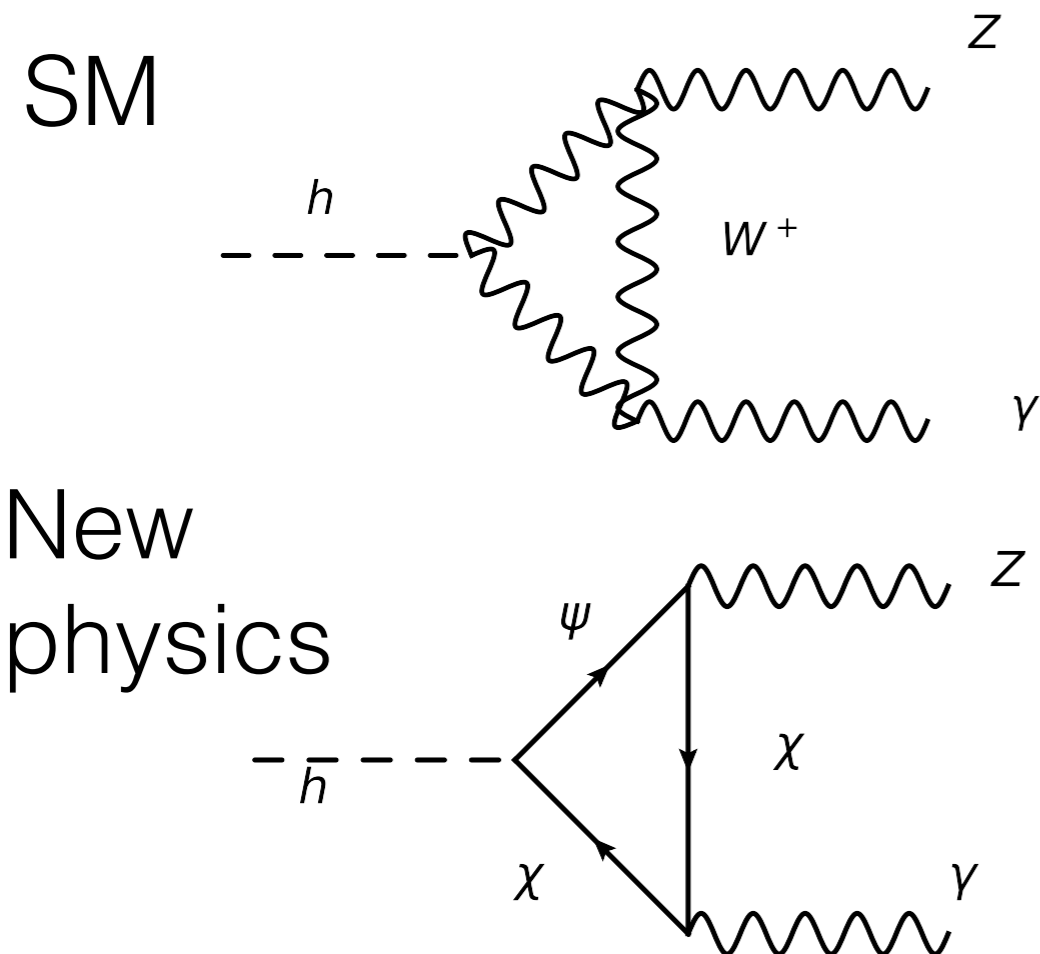
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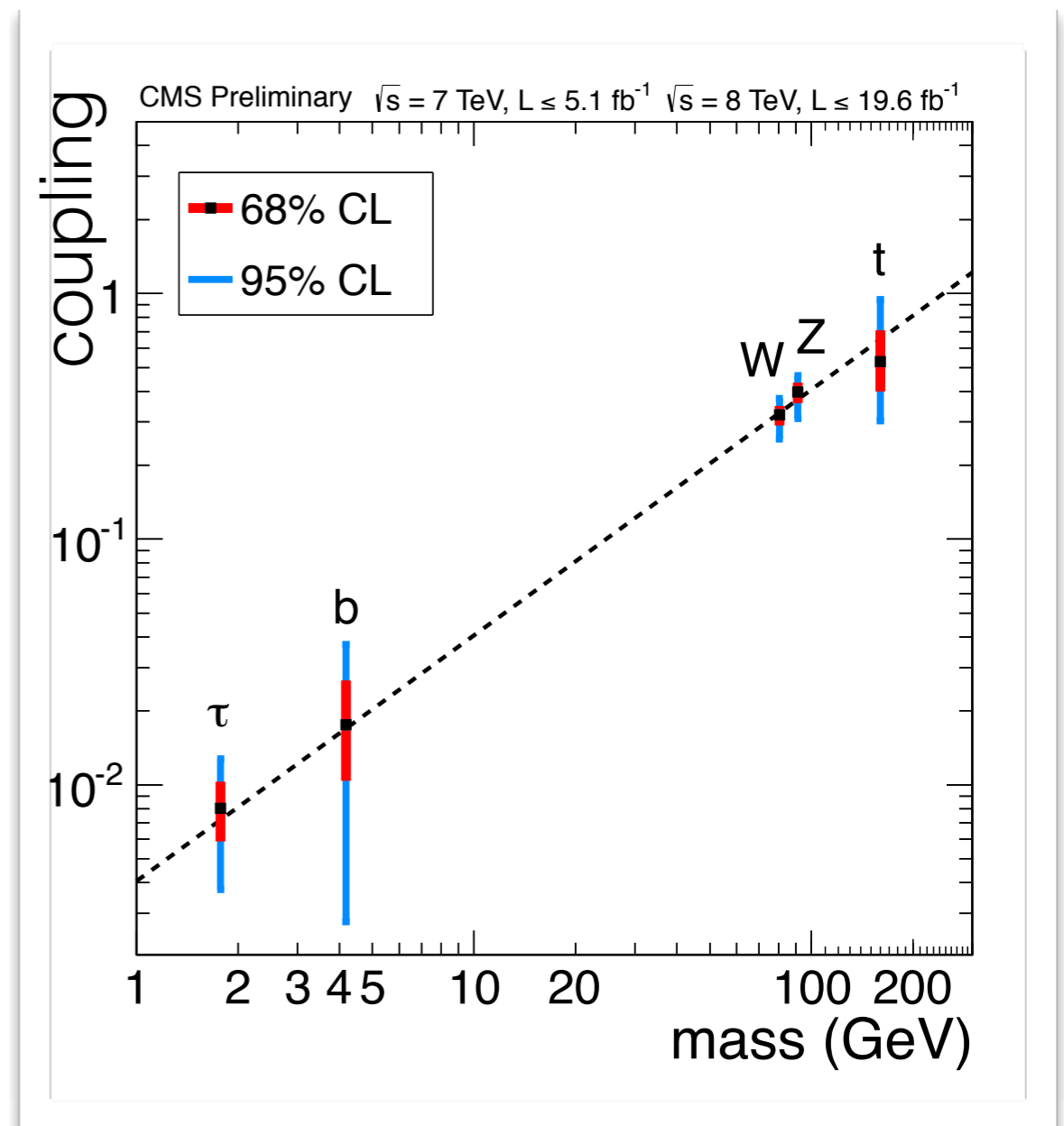
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- By LHC14@300, we'll have probed all 3rd generation fermion couplings to  $O(10\text{--}20\%)$
- $H \rightarrow \mu^+\mu^-$  gives us access to 2nd lepton generation, i.e. is the mass-generation mechanism same for all generations, for quarks and leptons?

mass  $\propto$  coupling to Higgs ?

$$Br(H \rightarrow \mu^+\mu^-)_{SM} = 2.2 \cdot 10^{-4}$$

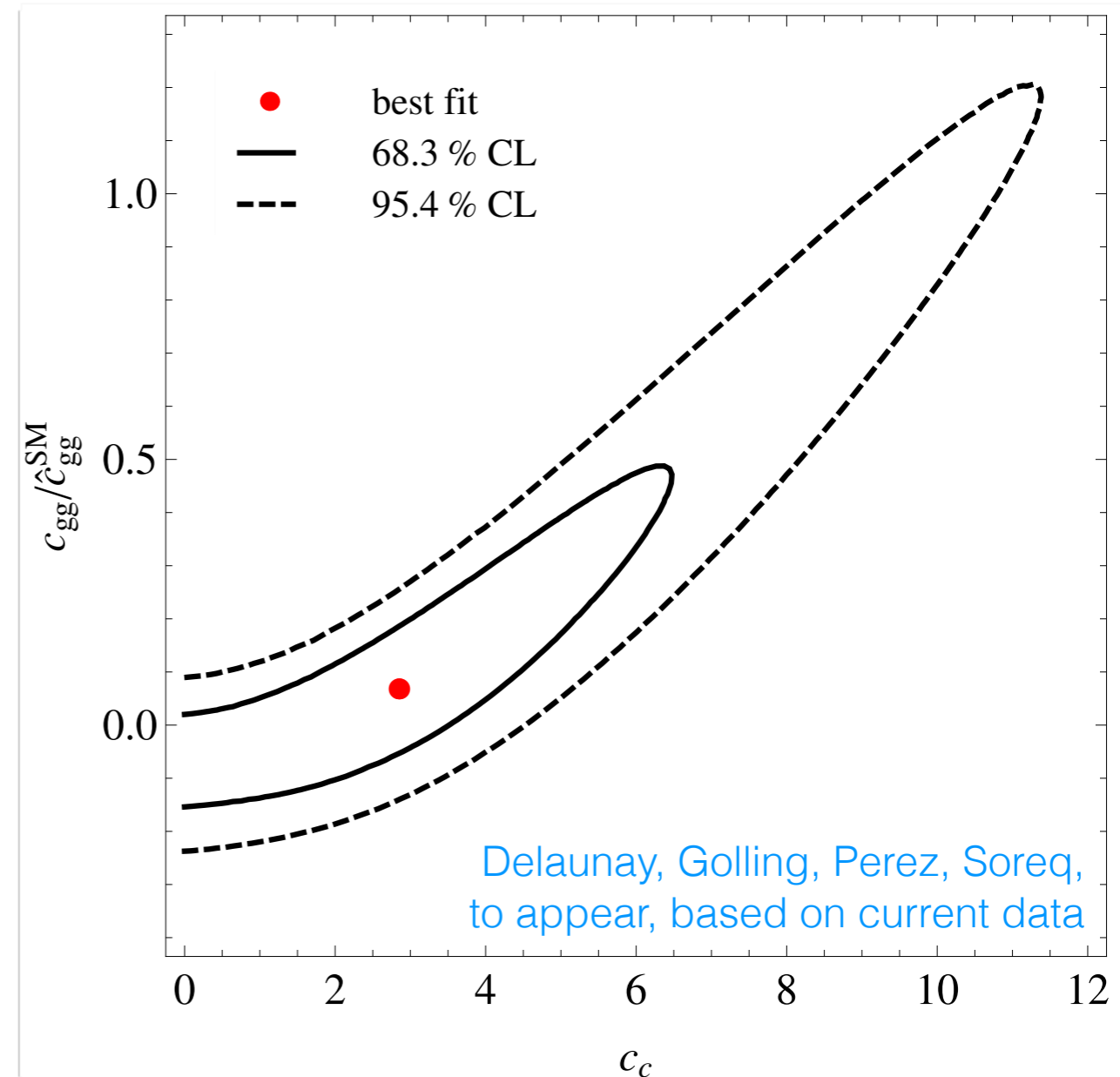


- Hcc coupling can still be 4-8 x SM
- In composite Higgs

$$c_c \simeq 1 + \mathcal{O}\left(\frac{v^2}{f^2}\right) + \mathcal{O}\left(\epsilon_c^2 \frac{g_\psi^2 v^2}{m_\psi^2}\right)$$

large for composite charm and light charm partners

$$\mathcal{L} = c_c h \frac{m_c}{v} \bar{c}c + \dots$$





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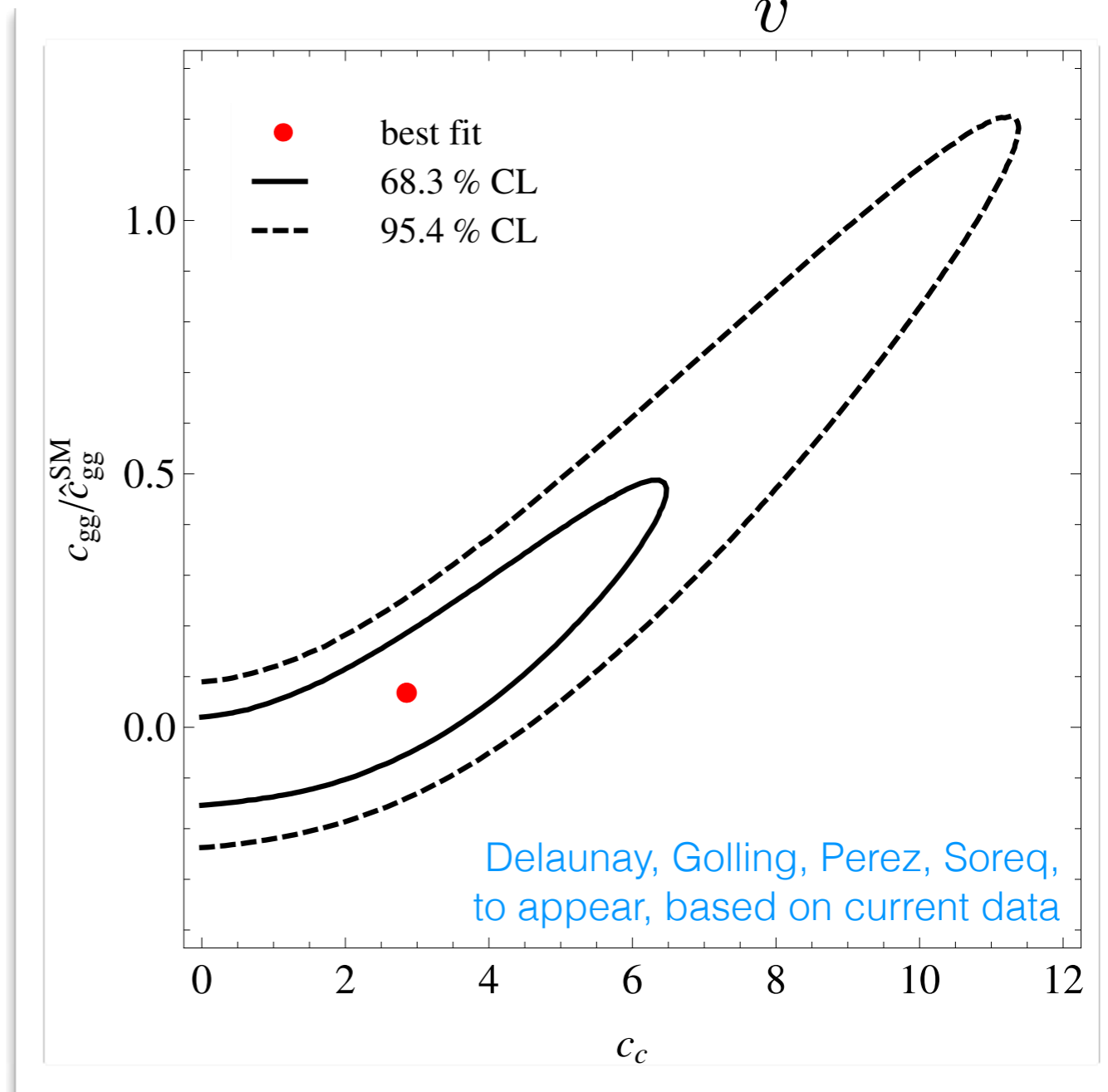
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### Measuring it?

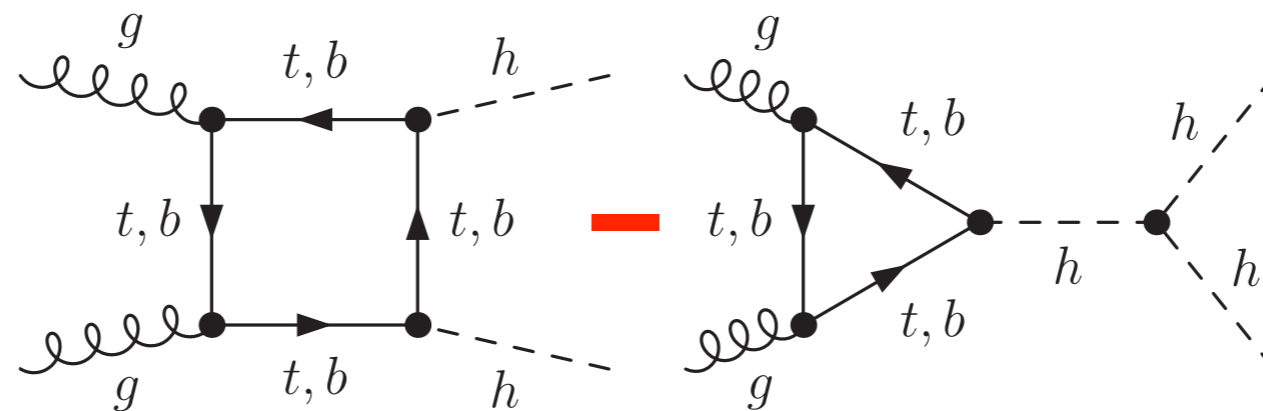
Like H → bb, but with charm tagging?

Or via H → J/ψ γ ? [1306.5770](#)



# Double Higgs production beyond the SM

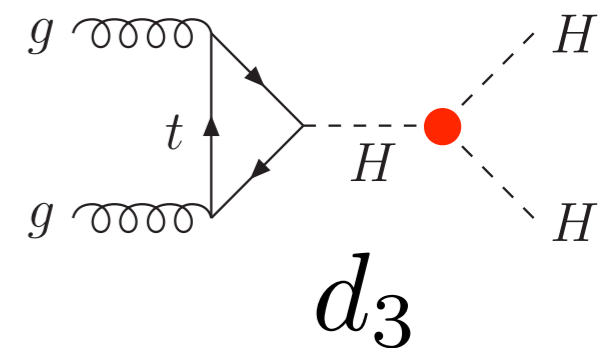
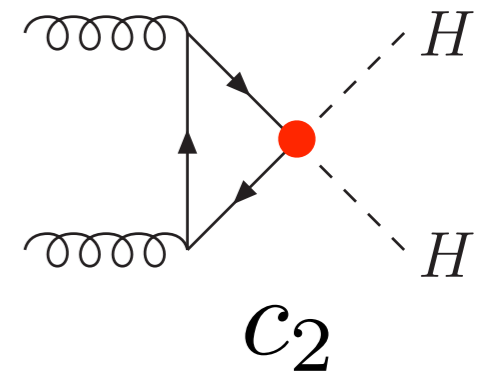
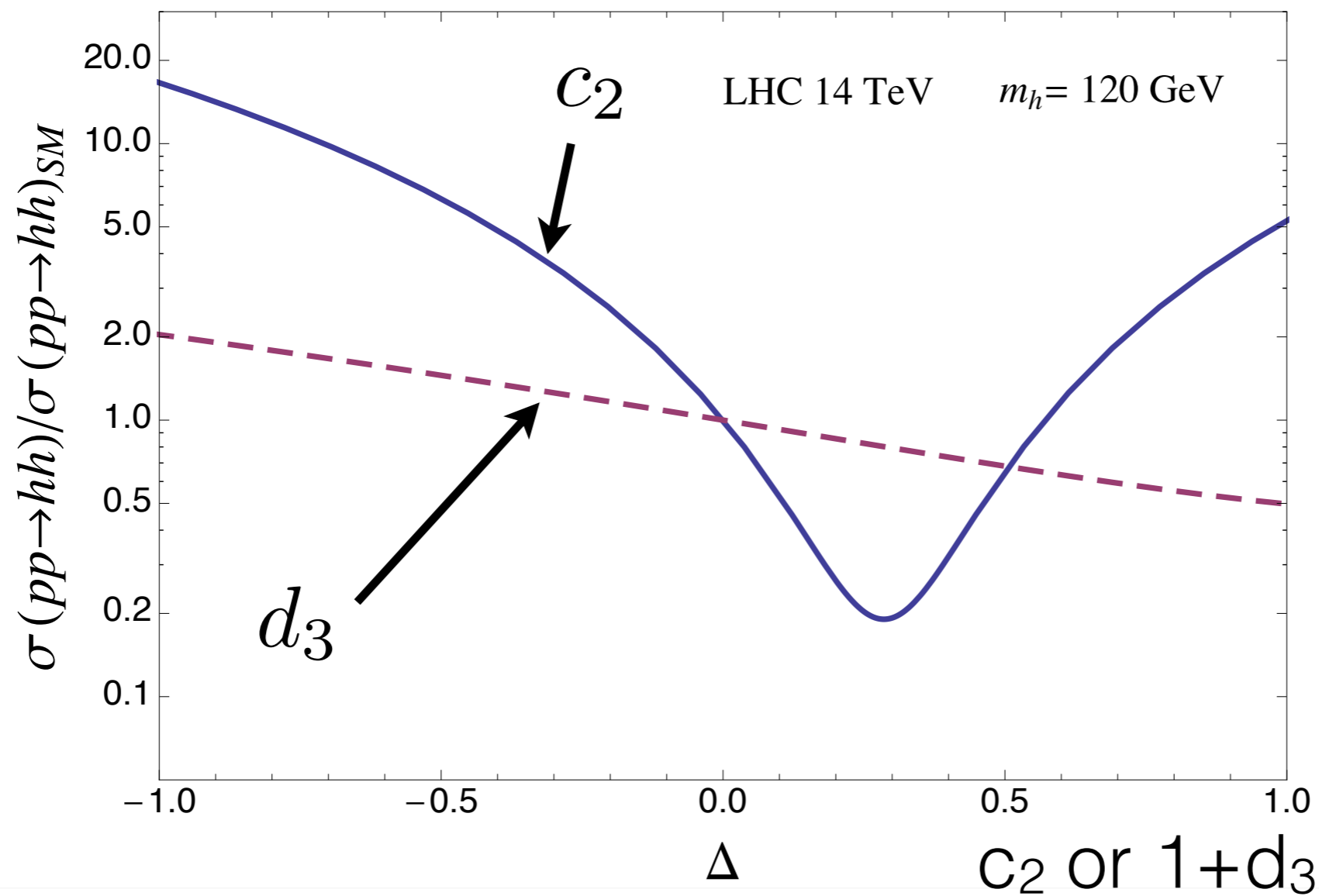
- In SM, access to self-coupling — test of structure of Higgs potential
- Composite Higgs, MSSM at low  $\tan\beta$ , dilaton drastically affect HH rate
- Not an easy measurement, because of box–vertex destructive interference:  
cross section of  $40\pm 3$  fb [de Florian & Mazzitelli 1309.6594]



# Double Higgs production beyond the SM

$$\mathcal{L} = -\frac{1}{2}m_h^2 h^2 - d_3 \frac{1}{6} \left( \frac{3m_h^2}{v} \right) h^3 - m_t \bar{q}_L t_R \left( 1 + c_t \frac{h}{v} + c_2 \frac{h^2}{v^2} \right)$$

SM :  $d_3 = 1, c_2 = 0$



[1205.5444](https://arxiv.org/abs/1205.5444)

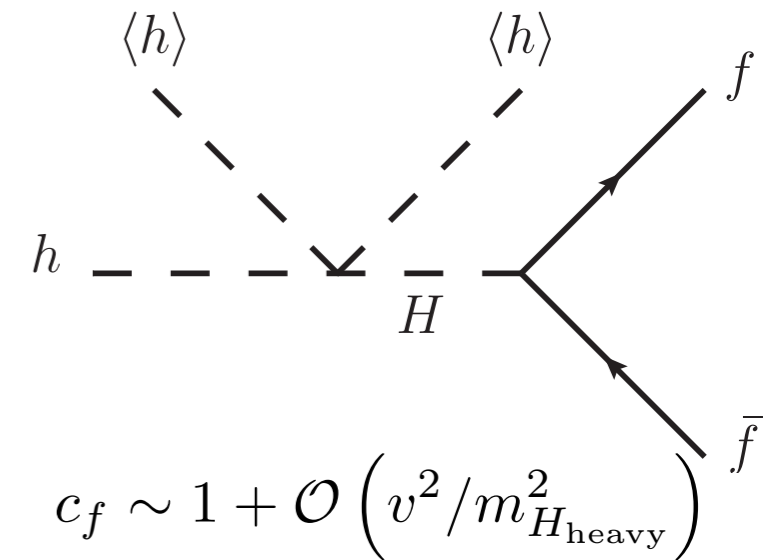
# Why higher precision is important

Supersymmetry example:

MSSM tree-level **Higgs mass**  $\leq M_Z$ , lift via

- 1) large quantum fluctuation from the **stop**
- 2) new interactions, e.g add **singlet**: NMSSM

Unravel with precise  $ffH$  coupling measurement



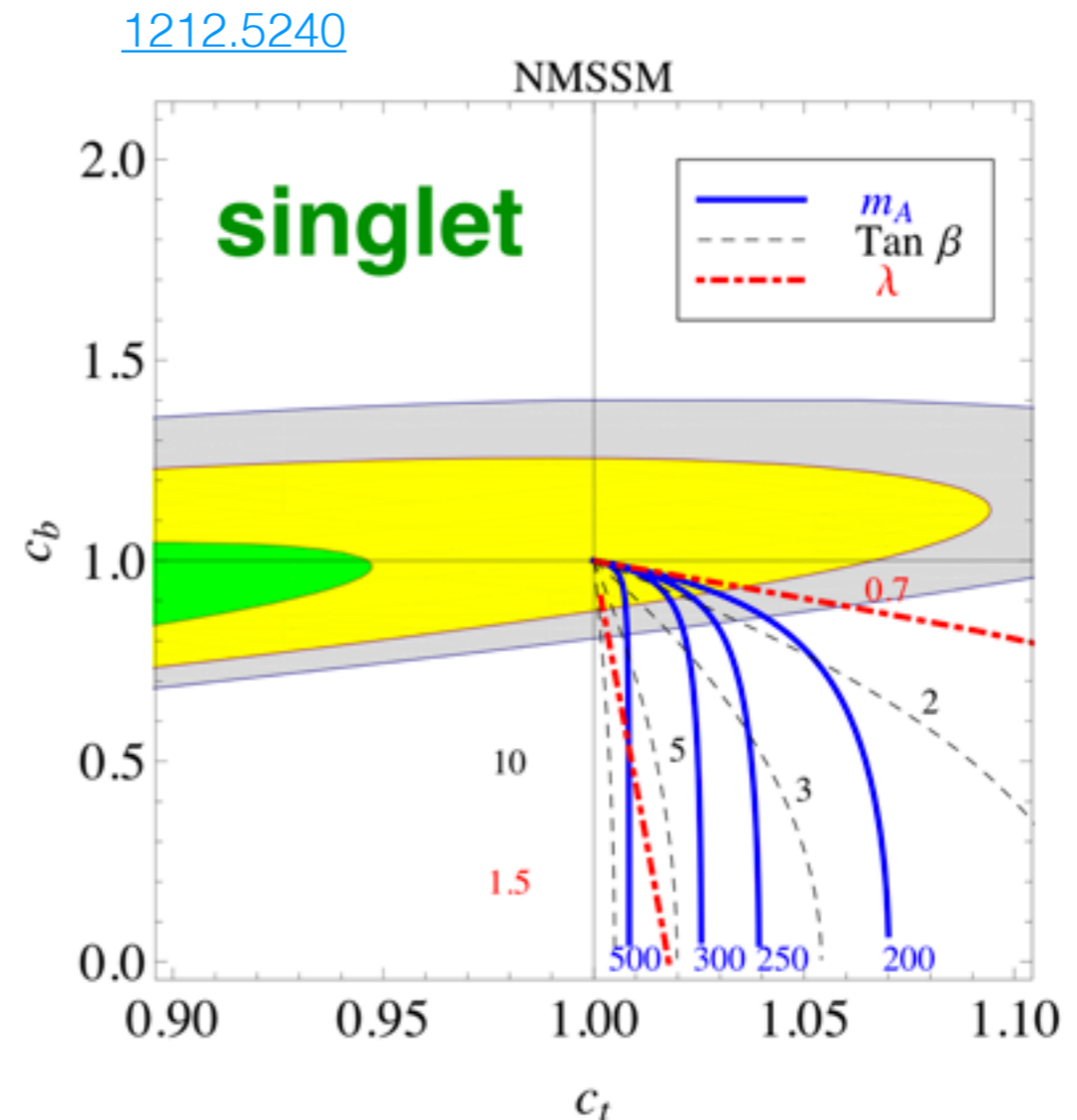
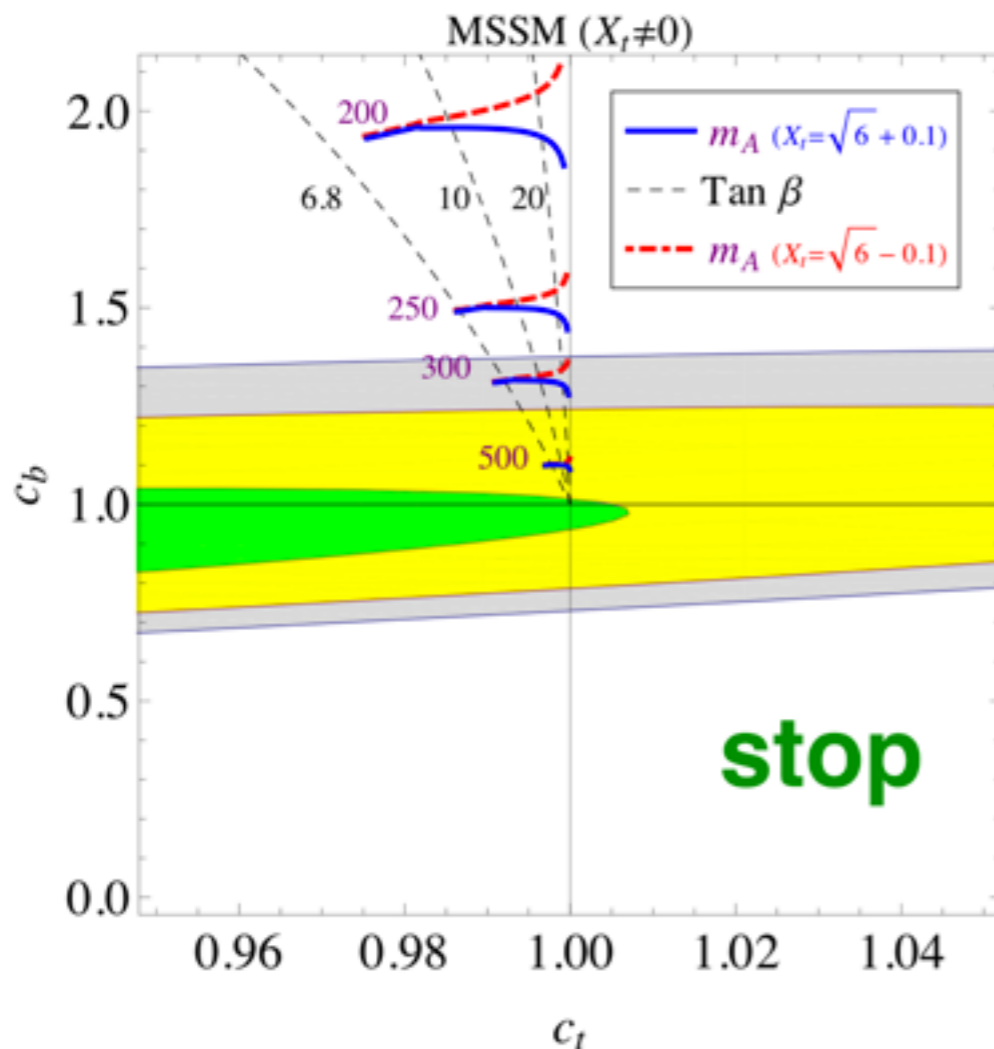
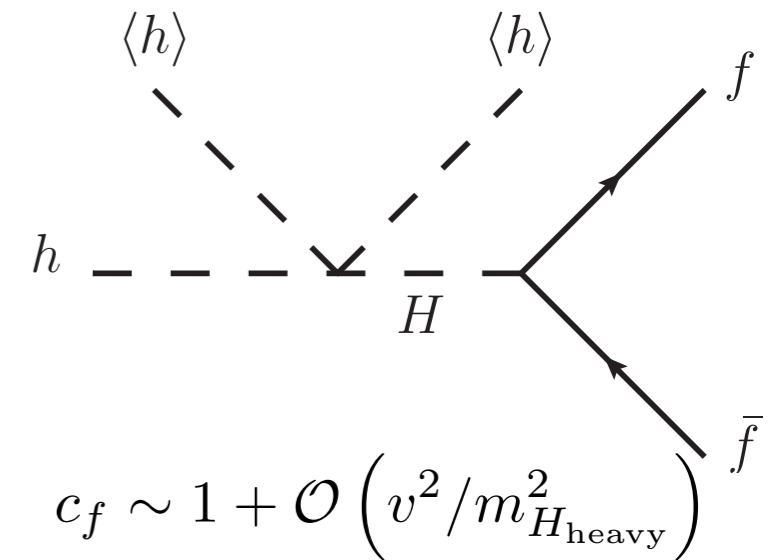
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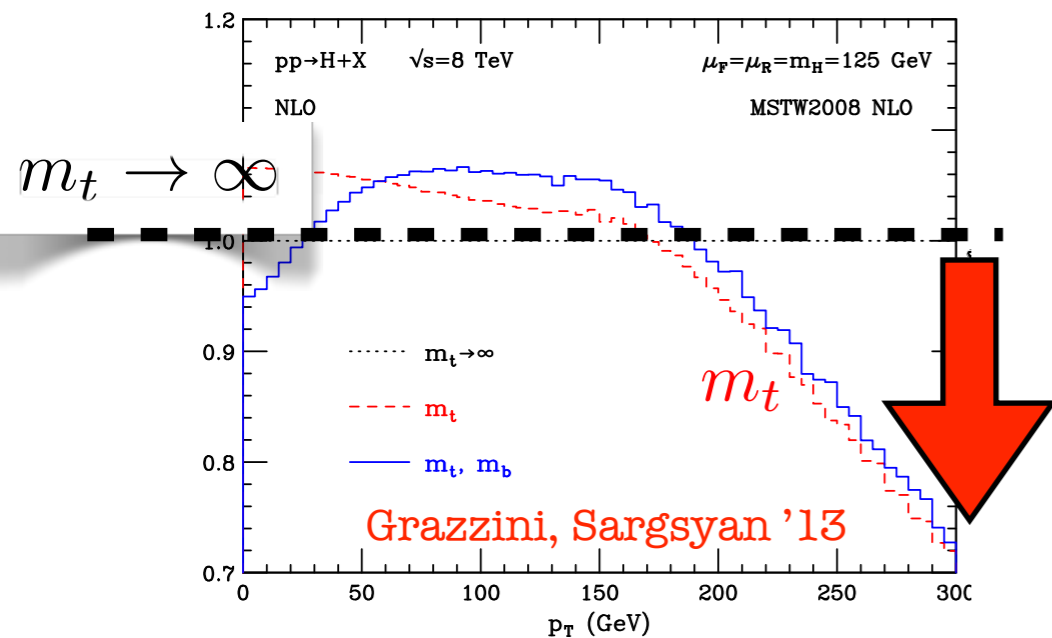


# Why high $p_T$ Higgs is important

$$\text{SM} + \mathcal{L} = \frac{\alpha_s c_g}{12\pi} |H|^2 G_{\mu\nu}^a{}^2 + \frac{\alpha c_\gamma}{2\pi} |H|^2 F_{\mu\nu}^2 + y_t c_t \bar{q}_L \tilde{H} t_R |H|^2$$

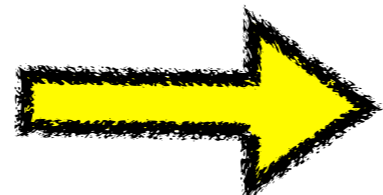
$$\frac{\sigma(gg \rightarrow h)}{\text{SM}} = (1 + (c_g - c_t)v^2)^2$$

Degeneracy 'short-distance' vs 'long-distance'



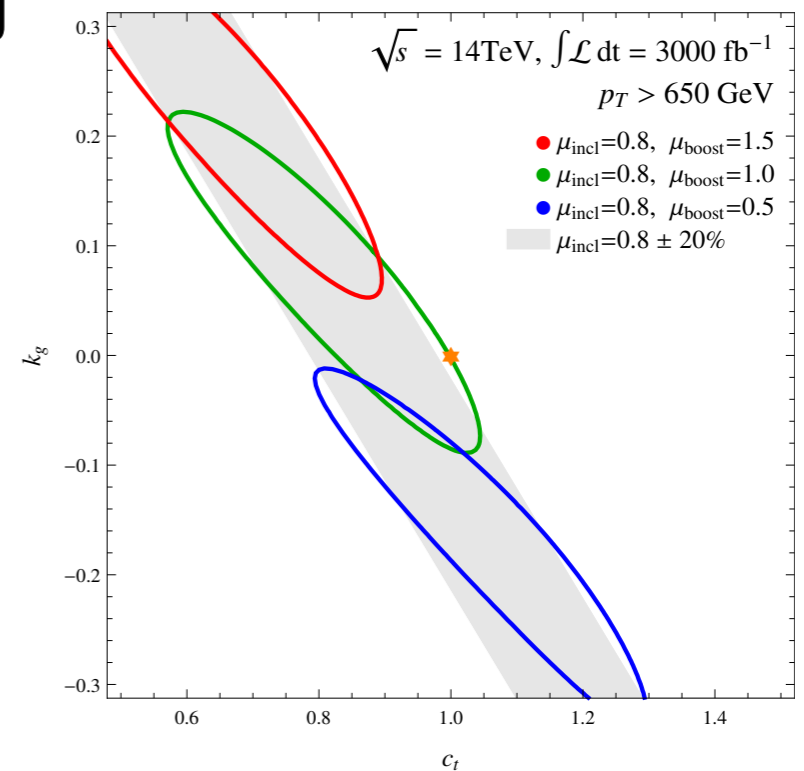
Baur, Glover '90,  
Langenegger et. al '06,

higgs -  $p_T$



high  $p_T$  tail resolves  
loop dynamics,  
breaks degeneracy

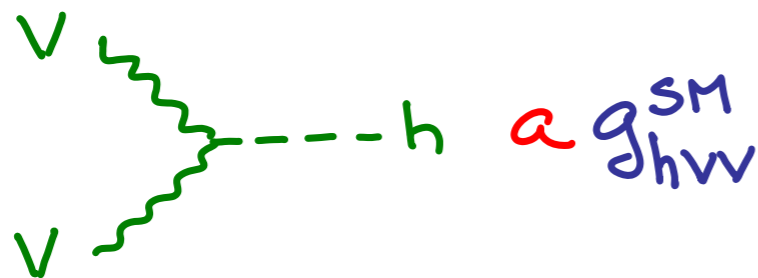
Grojean, Salvioni, Schlafer, Weiler



**C<sub>t</sub>**

- Any deviation in  $WWH$  coupling necessarily implies problems with tree-level unitarity in  $WW$  scattering (H alone doesn't completely unitarize it).
- A given deviation in the couplings  $\rightarrow$  upper bound at which strong interaction effects must appear.  
Signatures: growth of  $WW$  scattering amplitude or resonances
- E.g. for 20% deviations, scale is  $\sim 5$  TeV

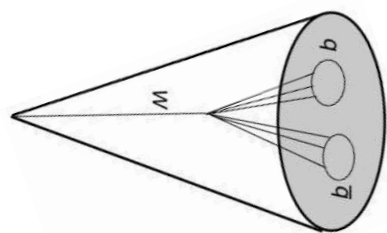
see e.g. [1005.4269](#)



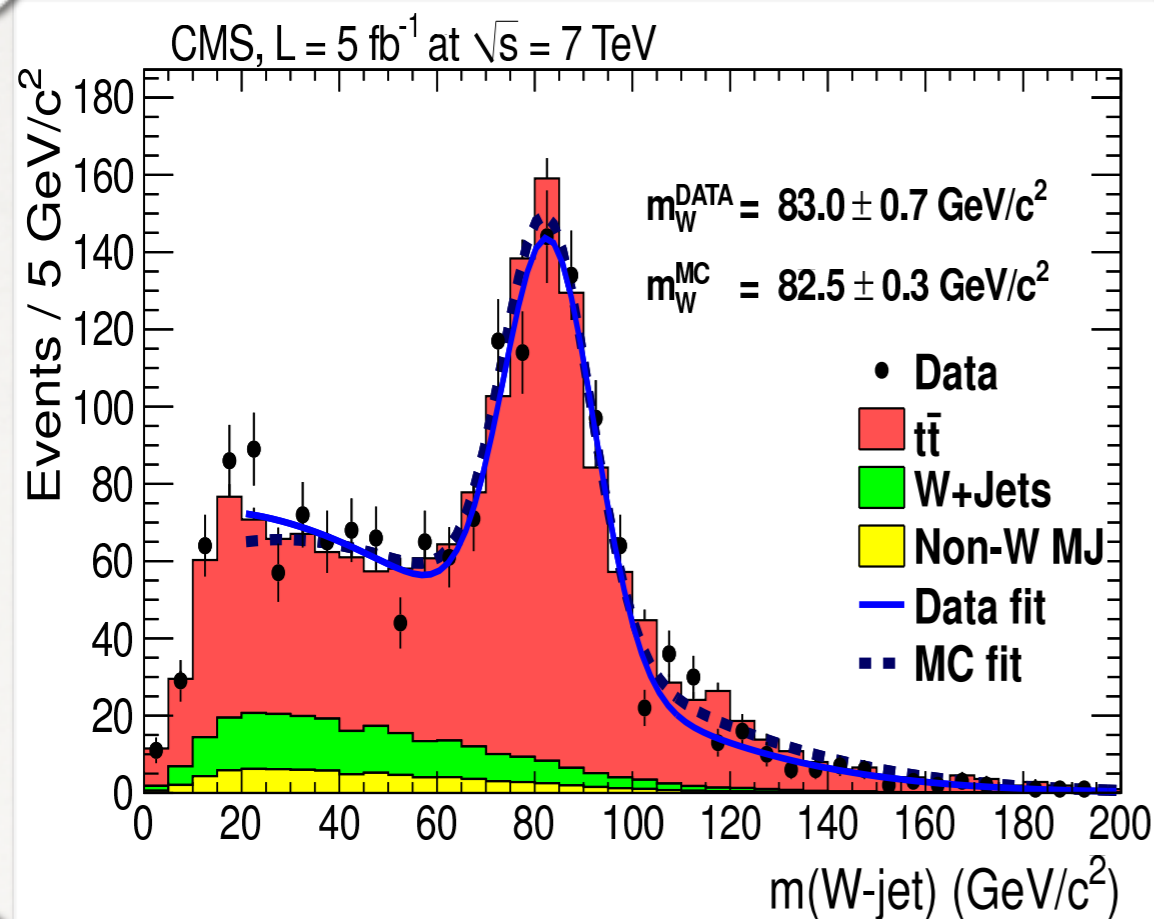
$$\Lambda \approx \frac{4\pi v}{\sqrt{1 - a^2}}$$

# Vector Boson Scattering and jet substructure

## W's in a single jet



At high  $p_t$ , W decays are collimated: for hadronic decays, both “jets” clustered into a single jet. Such “W”-jets can be **tagged**.



- Jet substructure can play an important role in VBS studies [larger hadronic BR, full  $M_{WV}$  reco]
- Crucial in many BSM studies, with W's, Z's, H's, tops coming from high-mass new-physics objects
- Important to **keep substructure in mind in performance studies**



## 2. Direct searches for New Physics

# What reach does $3\text{ab}^{-1}$ @ 14 TeV bring?

Consider a given search today (e.g.  $19\text{fb}^{-1}$ @8TeV), sensitive to a mass scale  **$M_{\text{low}}$** , e.g.  $M_Z$ , or  $2m_{\text{squark}}$ .

How can we estimate the corresponding “reach”,  $M_{\text{new}}$ , of a search at 14 TeV, with 300 or 3000  $\text{fb}^{-1}$ ?

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Very basic estimate: solve following equation for  $M_{\text{high}}$

$$\frac{N_{\text{signal-events}}(M_{\text{high}}^2, 14 \text{ TeV}, \text{Lumi})}{N_{\text{signal-events}}(M_{\text{low}}^2, 8 \text{ TeV}, 19\text{fb}^{-1})} = 1$$

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Even simpler: instead of ratio of # of events, **use ratio of “partonic luminosities”** (e.g.  $q\bar{q}$  lumi,  $gg$  lumi)

[including correct dimensional factors]

NB valid if bkgd & signal driven by same partonic channel

# Luminosity method: check 2

1208.1447  
ATLAS-CONF-2013-024

gg  
stop limits [expected] ( $L_{sp} = 0 \text{ geV}$ )  
7tev, 4.7 ifb 500 gev  
8tev, 20.5 ifb 650 gev ----> 675 GeV

qqbar ATLAS EXOT-2011-06  
ATLAS-CONF-2012-129  
ATLAS-CONF-2013-017

sequential z-prime [expected]  
7tev, 1.1 ifb 1800 gev  
8tev, 6 ifb, 2500 gev ----> 2450 GeV  
8 tev, 20 ifb 2800 gee ----> 2790 GeV

EXOT-2011-07  
ATLAS-CONF-2012-088  
ATLAS-CONF-2012-148

qg  
excited quark  $q^*$  [expected] (NB, sig  $\neq$  bgd scaling )  
7 tev, 1 ifb 2900 gev  
8 tev, 5.8 ifb 3500 gev ----> 3700 GeV  
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Lumi  
method

qg

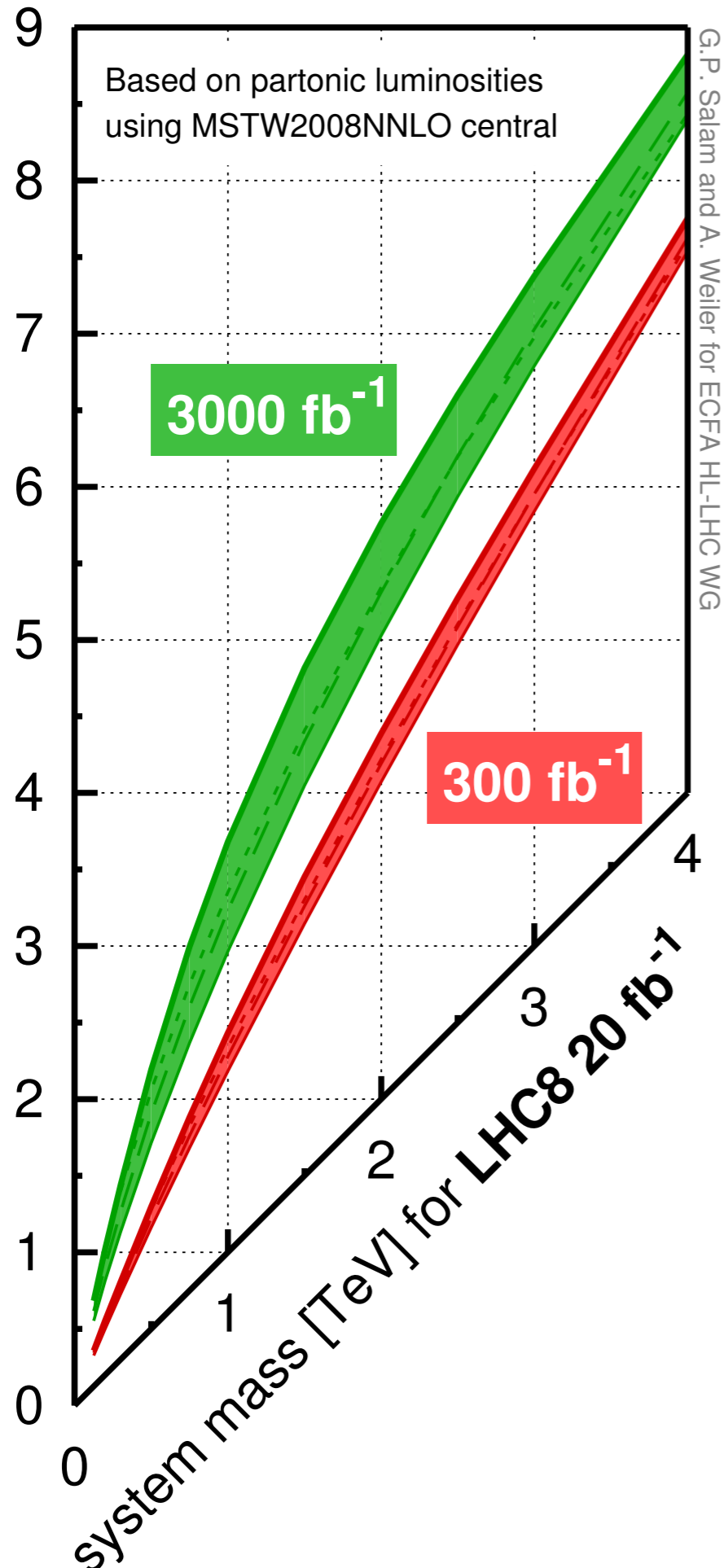
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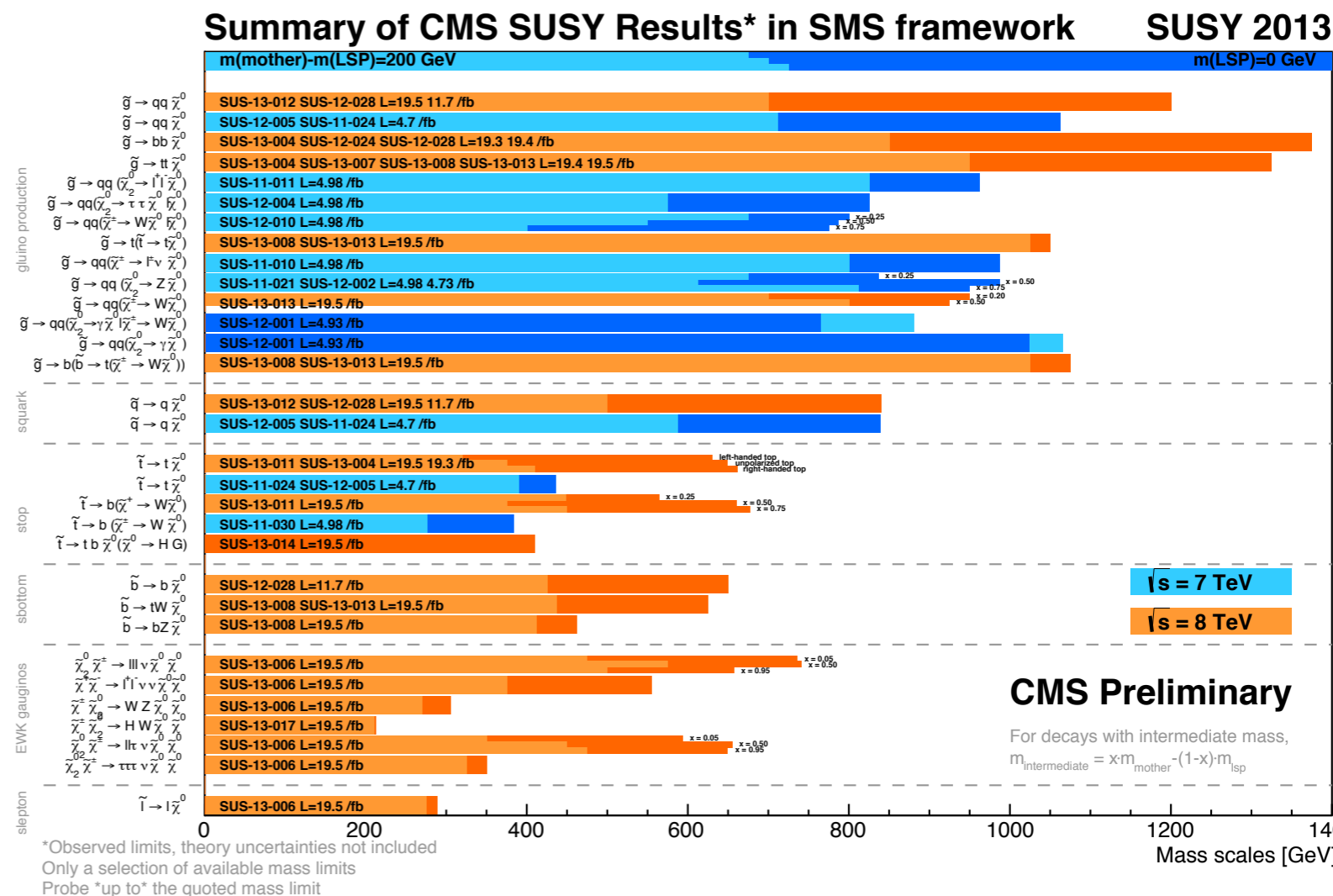
(NB, sig ≠ bgd scaling )  
----> 3700 GeV  
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system mass [TeV] for LHC14



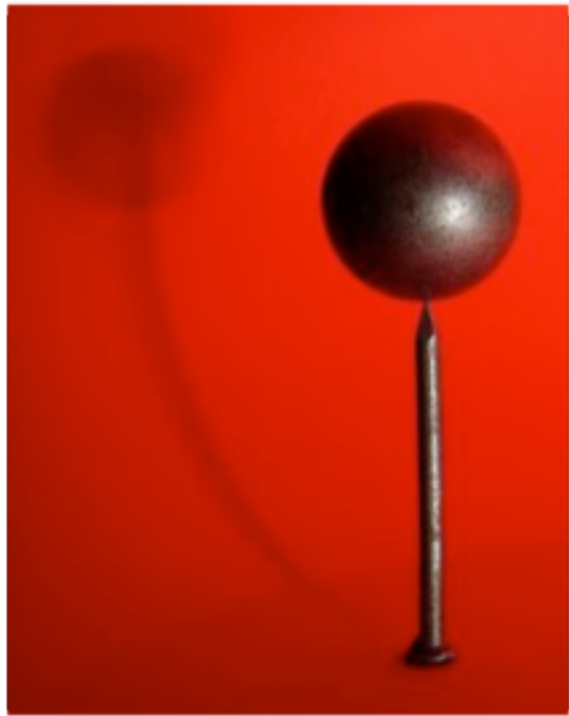
- $\Sigma\Sigma$
- -  $\Sigma g$
- · -  $\Sigma_i q_i \bar{q}_i$
- $gg$

Take existing searches and figure out reach at 14 TeV, for different lumis

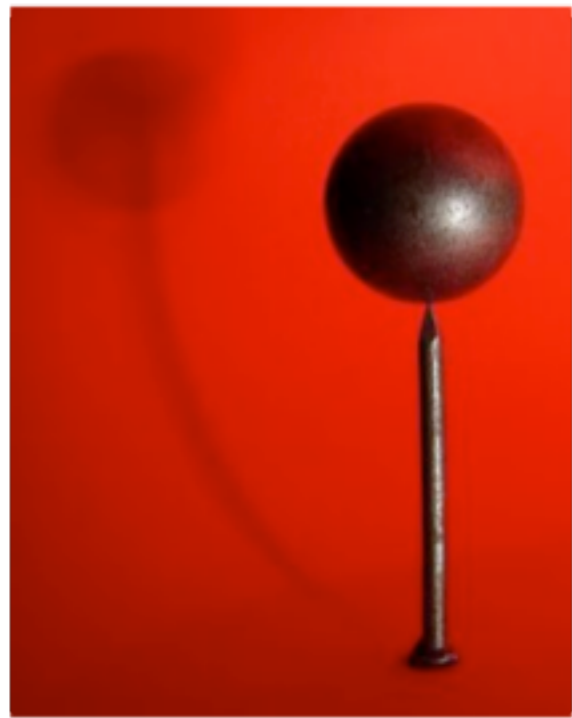




## Fine-tuning in the Higgs sector



## Fine-tuning in the Higgs sector



Light Higgs

stabilizing  
quantum corrections

?

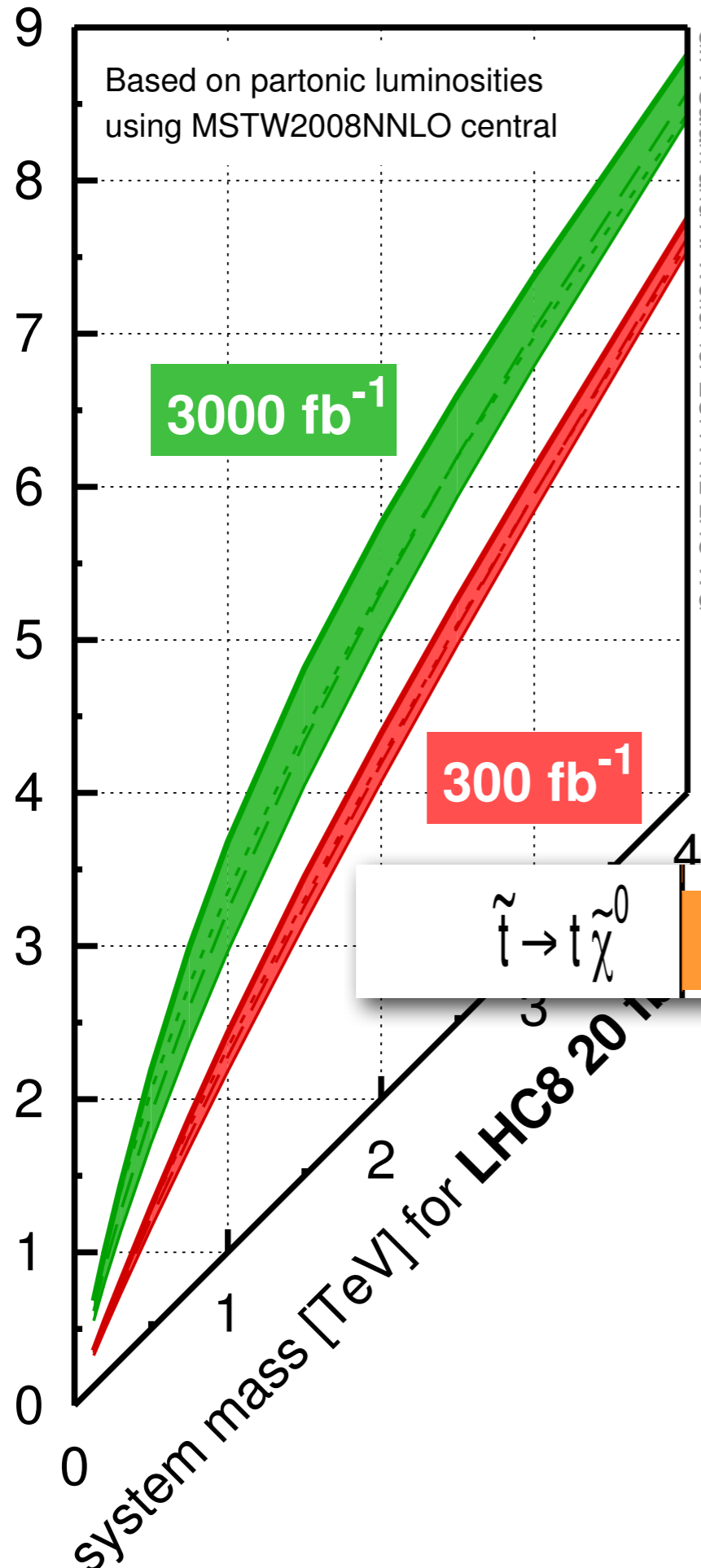
light stops<sub>1,2</sub>, sbottom<sub>L</sub>,  
higgsinos, gluinos, ...

supersymmetry

light top partners  
( $Q=5/3, 2/3, 1/3$ )

composite Higgs

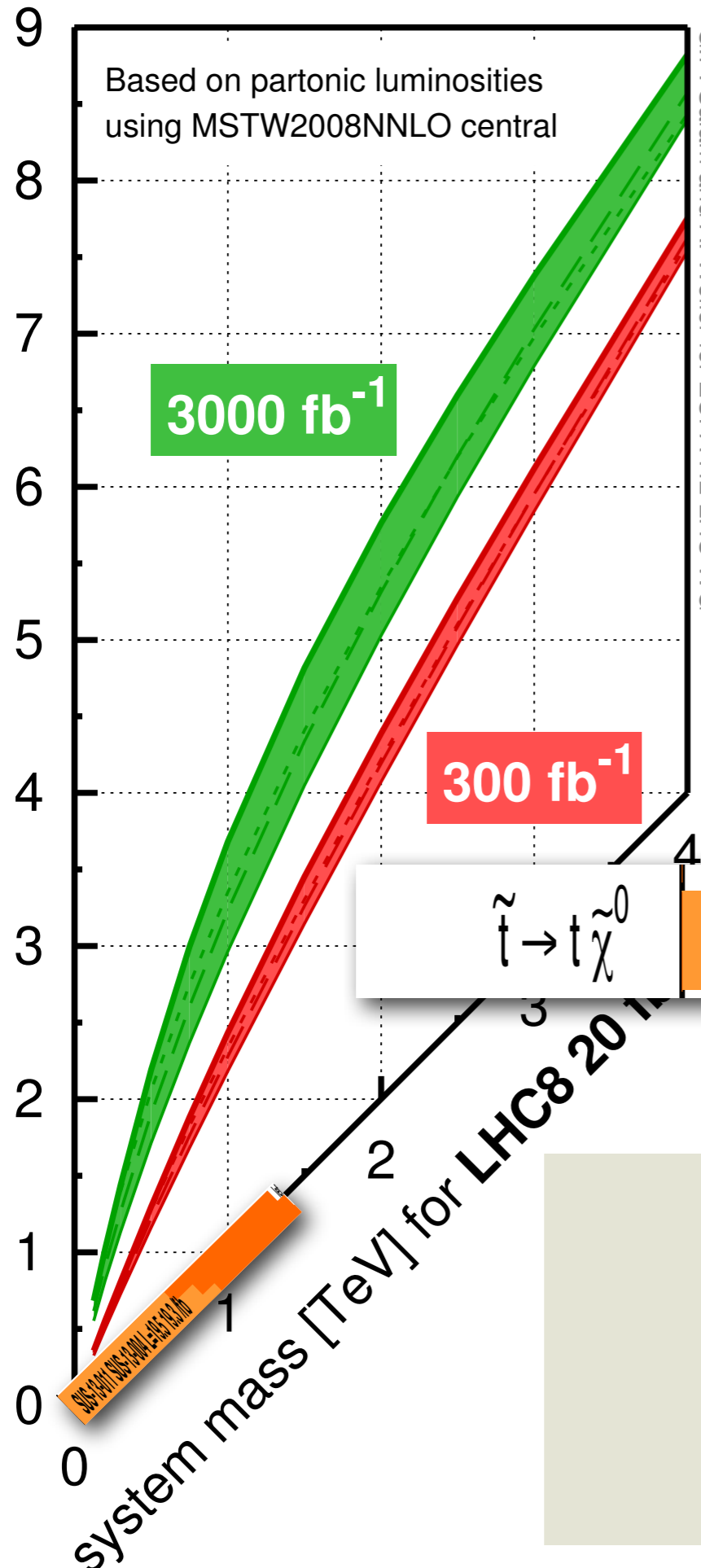
system mass [TeV] for LHC14



- ΣΣ
- - Σg
- · - · Σ<sub>i</sub> q<sub>i</sub> q̄<sub>i</sub>
- gg

SUS-13-011 SUS-13-004 L=19.5 19.3 /fb m<sub>t̃</sub> ≳ 650 GeV

system mass [TeV] for LHC14



- $\Sigma\Sigma$
- -  $\Sigma g$
- ⋯  $\Sigma_i q_i \bar{q}_i$
- · -  $gg$

3000 fb<sup>-1</sup>

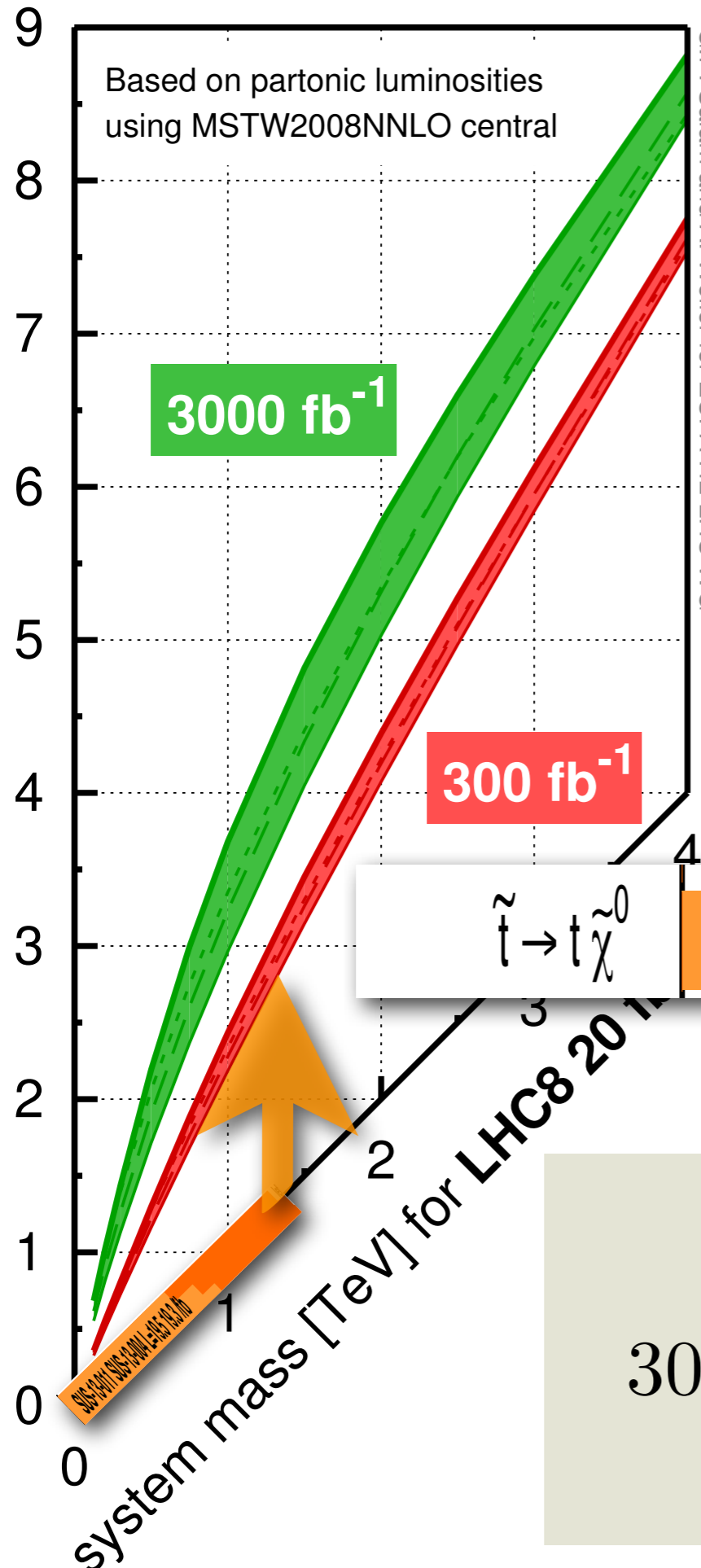
300 fb<sup>-1</sup>

$\tilde{t}^2 \rightarrow t \tilde{\chi}^0$

SUS-13-011 SUS-13-004 L=19.5 19.3 /fb  $m_{\tilde{t}} \gtrsim 650 \text{ GeV}$

$m_{\tilde{t}} \gtrsim 650 \text{ GeV} \equiv m_{\tilde{t}\tilde{t}} \gtrsim 1.3 \text{ TeV}$

system mass [TeV] for LHC14

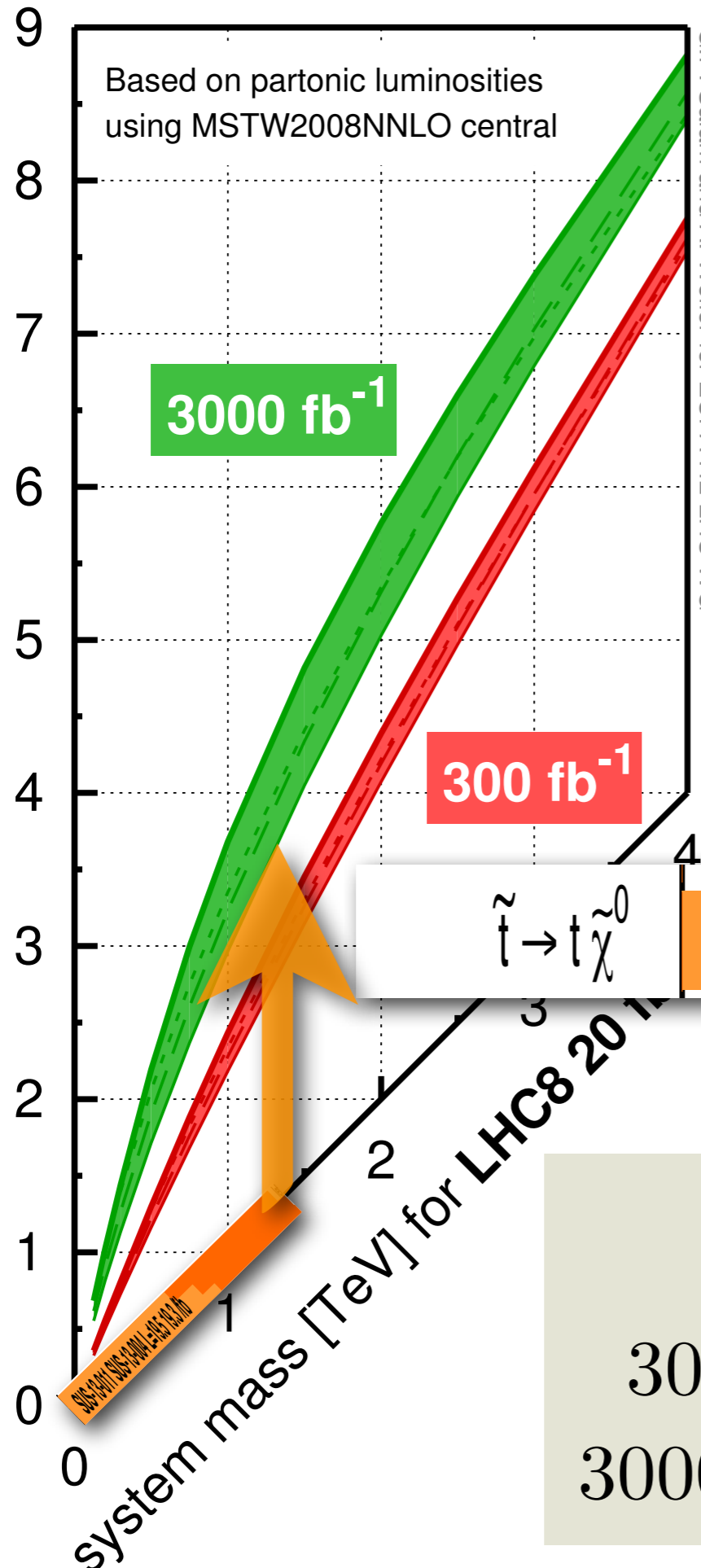


- $\Sigma\Sigma$
- -  $\Sigma g$
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 $300 \text{ fb}^{-1} @ 14 \text{ TeV}: 2.8 \text{ TeV} (m_{\tilde{t}} > 1.4 \text{ TeV})$

system mass [TeV] for LHC14



G.P. Salam and A. Weiler for ECFA HL-LHC WG

- $\Sigma\Sigma$
- -  $\Sigma g$
- ...  $\Sigma_i q_i \bar{q}_i$
- $gg$

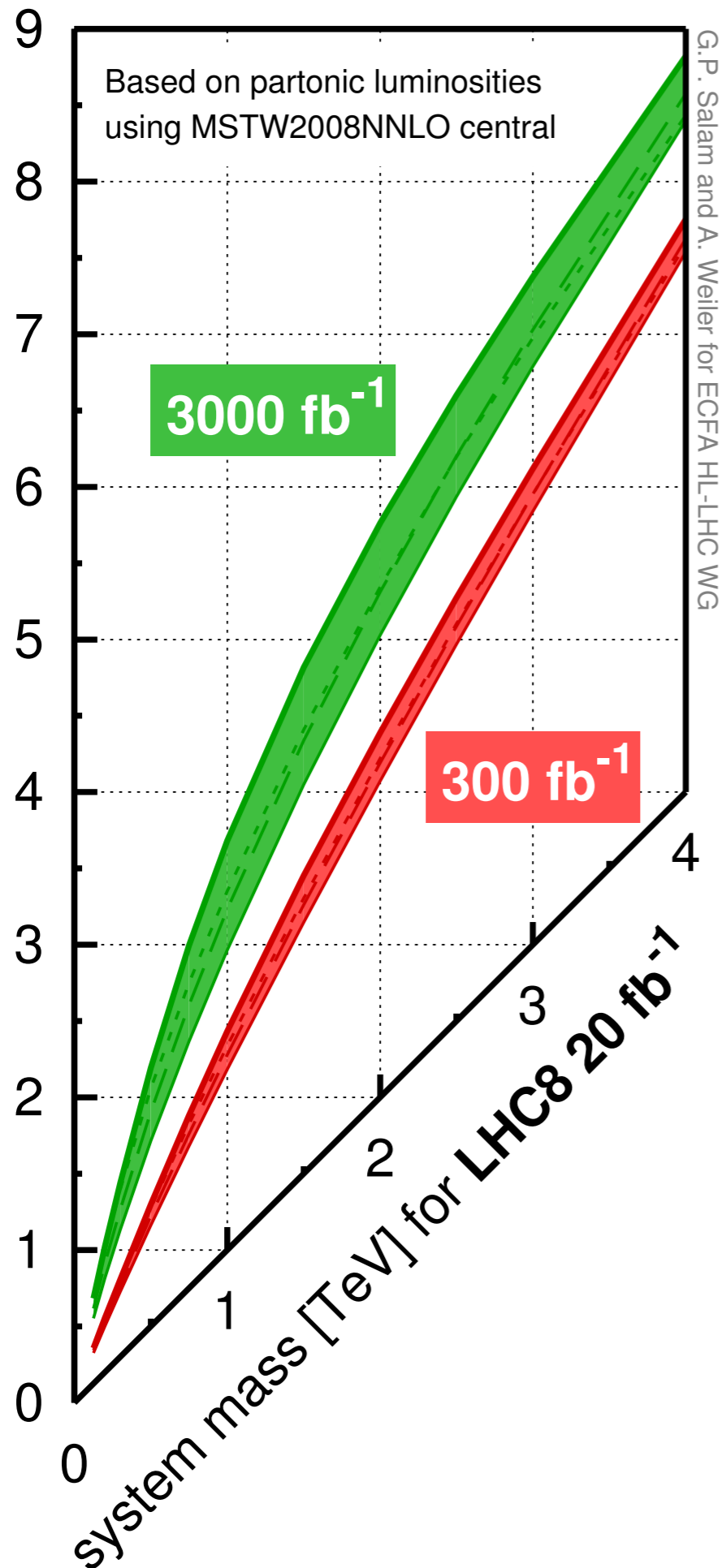
**SUS-13-011 SUS-13-004 L=19.5 19.3 /fb**  **$m_{\tilde{t}} \gtrsim 650 \text{ GeV}$**

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300 fb<sup>-1</sup> @ 14 TeV: 2.8 TeV ( $m_{\tilde{t}} > 1.4 \text{ TeV}$ )

3000 fb<sup>-1</sup> @ 14 TeV: 3.6 TeV ( $m_{\tilde{t}} > 1.8 \text{ TeV}$ )

system mass [TeV] for LHC14



Increase in c.o.m. energy to 14 TeV brings substantial extra reach

**x10 in lumi also brings extra TeV**  
in reach

That can be quite significant (e.g. 35-50%) at the lower end of the range, e.g. when today's limits are ~ 1 TeV.

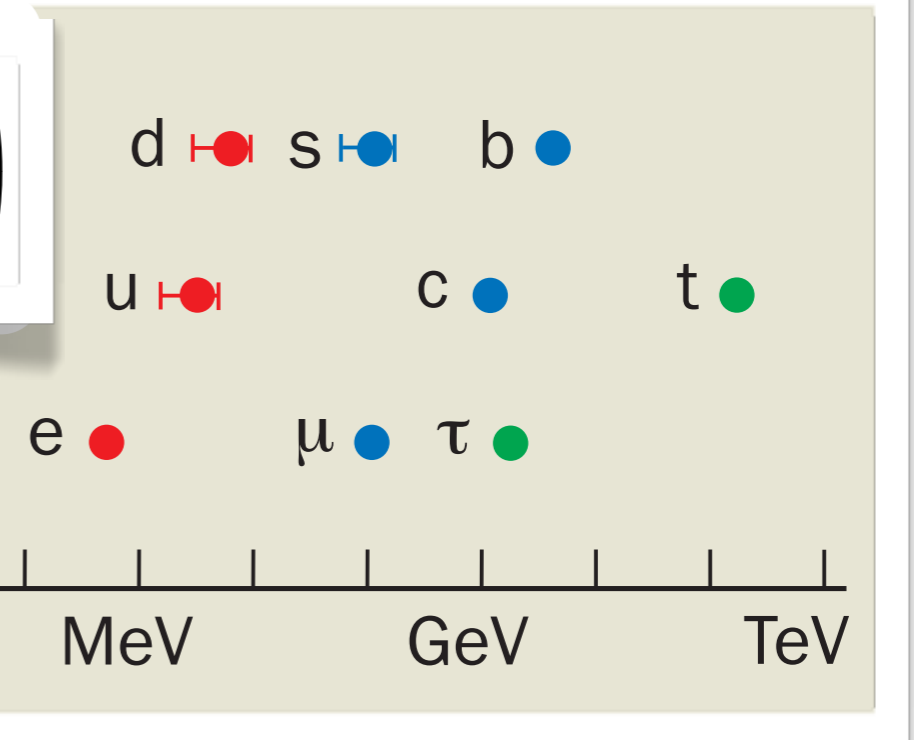
These numbers are to be taken as **indicative of what might be possible**

—  
real life may be tougher,  
future analyses may be cleverer

## Origin of SM matter and flavour?

e.g. up/charm/top mass matrix:

$$Y_U \approx \begin{pmatrix} 10^{-5} & -0.002 & 0.007 + 0.004i \\ 10^{-6} & 0.007 & -0.04 + 0.0008i \\ 10^{-8} + 10^{-7}i & 0.0003 & 0.92 \end{pmatrix}$$



Compare to other couplings of the SM:

$$g_s \sim 1, \quad g \sim 0.6, \quad g' \sim 0.3, \quad \lambda_{\text{Higgs}} \sim 1/8$$



# FCNCs are sensitive SM flavour structure

- The SM flavour sector has passed all tests at B-factories and recently at LHCb, ATLAS and CMS
- Less likely to find  $O(1)$  deviations from the SM, need theoretically clean observables.
- High potential flavour measurements, some examples

$$R = \frac{\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)}{\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)} \quad q_0^2 A_{\text{FB}}(K^{*0} \mu^+ \mu^-) \quad \text{CKM } \gamma$$

$$\phi_s(B_s^0 \rightarrow J/\psi \phi)$$

- Top FCNC rate sensitivity can be improved by  $\sim 4x$

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$$R = \frac{\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)}{\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)}$$

Clean prediction [1208.0934](#)

$$\delta R/R = \pm 0.06 \pm 2\sigma_{f_{s/d}^r}$$

$$q_0^2 A_{\text{FB}}(K^{*0} \mu^+ \mu^-)$$

clean theory, BSM sensitivity

$$\phi_s(B_s^0 \rightarrow J/\psi \phi)$$

CKM  $\gamma$

up to

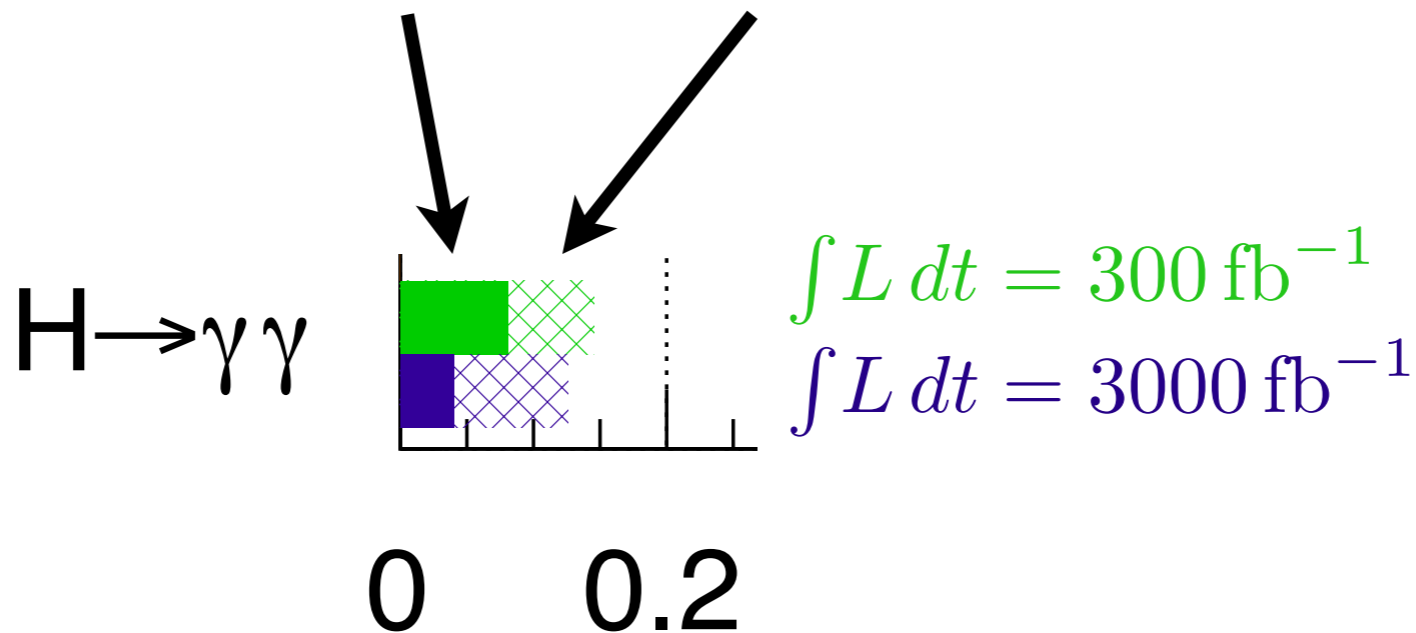
1<sup>o</sup> precision measurement

- Top FCNC rate sensitivity can be improved by  $\sim 4x$

### 3. SM sector, including theory prospects (e.g. for use in Higgs physics)

Assuming  
no theory  
uncertainty

Assuming  
today's theory  
(and PDF) uncertainty



**Many uncertainties cancel in ratios, e.g.  $\gamma\gamma / ZZ$ .**

**But what are prospects for doing  
better directly on cross sections?**

# Sources of uncertainty?

- **Missing higher orders** (“scale”) in QCD and EW. **~ 10%**
- Extra QCD uncertainties in the presence of **cuts** (e.g. jet vetoes). **~ 10%**
- **PDF uncertainties** (within/between groups), **~ 7%**
- **Fundamental constants** ( $\alpha_s$ ,  $m_b$ , etc.) **~ few %**

$m_H$ (GeV)	Cross Section (pb)	+error %	- error %	+scale %	-scale %	+(PDF+ $\alpha_s$ ) %	-(PDF+ $\alpha_s$ ) %
125	49.85	19.6	-14.6	12.2	-8.4	7.4	-6.2

HXSWG  $gg \rightarrow H$

# Sources of uncertainty?

- **Missing higher orders** (“scale”) in QCD and EW. **~ 10%**  
[Need higher-order inclusive X-sections]
- Extra QCD uncertainties in the presence of **cuts** (e.g. jet vetoes). **~ 10%**  
[Need higher-order differential X-sections, fixed-order, resummed & MC]
- **PDF uncertainties** (within/between groups), **~ 7%**  
[Need better data & better theory to interpret it]
- **Fundamental constants** ( $\alpha_s$ ,  $m_b$ , etc.) **~ few %**

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125	49.85	19.6	-14.6	12.2	-8.4	7.4	-6.2

HXSWG  $gg \rightarrow H$

## gg → H

Since 11 years: NNLO (+ threshold)

1208.3130, 1211.6559

In 1-2 years? NNNLO (+ threshold)

1302.4379, 1306.2223

Since 10 years: NLO + parton shower

Past month: NNLO + parton shower

1309.0017

**gg → H**

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**pp → H+jet** (matters for searches with jet vetoes, jet selections, etc.)

Since 10–14 years: NLO, NNLL H  $p_t$  resummation

Past year: NNLL jet resummation

1206.4998, 1307.0025, 1307.1808

NNLO (gluonic part)

1302.6216



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1206.4998, 1307.0025, 1307.1808

NNLO (gluonic part)

1302.6216

NNLO for other  $2 \rightarrow 2$  is **here now (e.g.  $\gamma\gamma$ ,  $t\bar{t}$ )** or soon

# Projected uncertainties for $gg \rightarrow H$

Buehler & Lazopoulos, arXiv:1306.2223

## Expected $N^3LO$ (scale) uncertainties depend on size of $N^3LO$ corrections (parametrised by $K$ )

Order	Cross section [pb]	$\sigma/\sigma_{NNLO}$	$\sigma/\sigma_{LO}$
LO	10.31 $^{+26.9\%}_{-16.6\%}$	0.51	1.00
NLO	17.41 $^{+20.8\%}_{-12.7\%}$	0.86	1.69
NNLO	20.27 $^{+8.3\%}_{-7.1\%}$	1.00	1.97
$N^3LO$ ( $K=0$ )	18.53 $^{+1.2\%}_{-7.9\%}$	0.91	1.80
$N^3LO$ ( $K=5$ )	19.23 $^{+0.3\%}_{-5.1\%}$	0.95	1.87
$N^3LO$ ( $K=10$ )	19.92 $^{+0.0\%}_{-2.6\%}$	0.98	1.93
$N^3LO$ ( $K=15$ )	20.62 $^{+0.4\%}_{-2.2\%}$	1.02	2.00
$N^3LO$ ( $K=20$ )	21.31 $^{+2.0\%}_{-3.1\%}$	1.05	2.07
$N^3LO$ ( $K=30$ )	22.70 $^{+6.0\%}_{-4.9\%}$	1.12	2.20
$N^3LO$ ( $K=40$ )	24.09 $^{+9.6\%}_{-6.5\%}$	1.19	2.34

**Today**

**Near future:**

Depending on size of  $N^3LO$  correction ( $K$ ), final uncertainty could be anywhere from 2% to 8%



# Projected uncertainties for $gg \rightarrow H$

Buehler & Lazopoulos, arXiv:1306.2223

## Expected N<sup>3</sup>LO (scale) uncertainties depend on size of N<sup>3</sup>LO corrections (parametrised by K)

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LO	10.31 <sup>+26.9%</sup> <sub>-16.6%</sub>	0.51	1.00
NLO	17.41 <sup>+20.8%</sup> <sub>-12.7%</sub>	0.86	1.69
NNLO	20.27 <sup>+8.3%</sup> <sub>-7.1%</sub>	1.00	1.97
N <sup>3</sup> LO (K=0)	18.53 <sup>+1.2%</sup> <sub>-7.9%</sub>	0.91	1.80
N <sup>3</sup> LO (K=5)	19.23 <sup>+0.3%</sup> <sub>-5.1%</sub>	0.95	1.87
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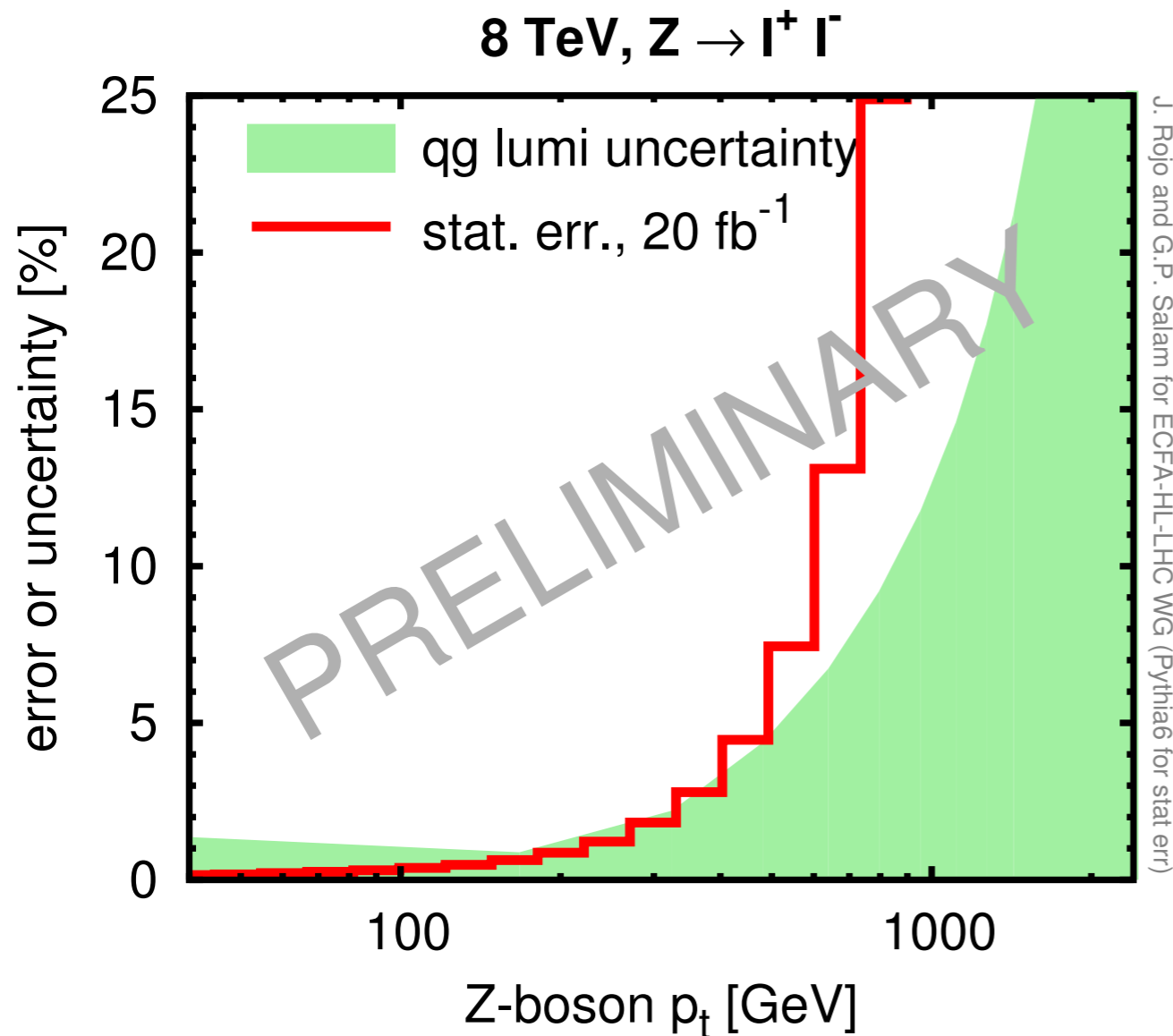
Today

Near future:

Roughly what comes out of N<sup>3</sup>LO estimate from Ball et al, 1303.3590



# e.g. of HL-LHC precision SM measurement: $Z p_t$ spectrum



[Thanks to J. Rojo for partonic lumi uncertainties]

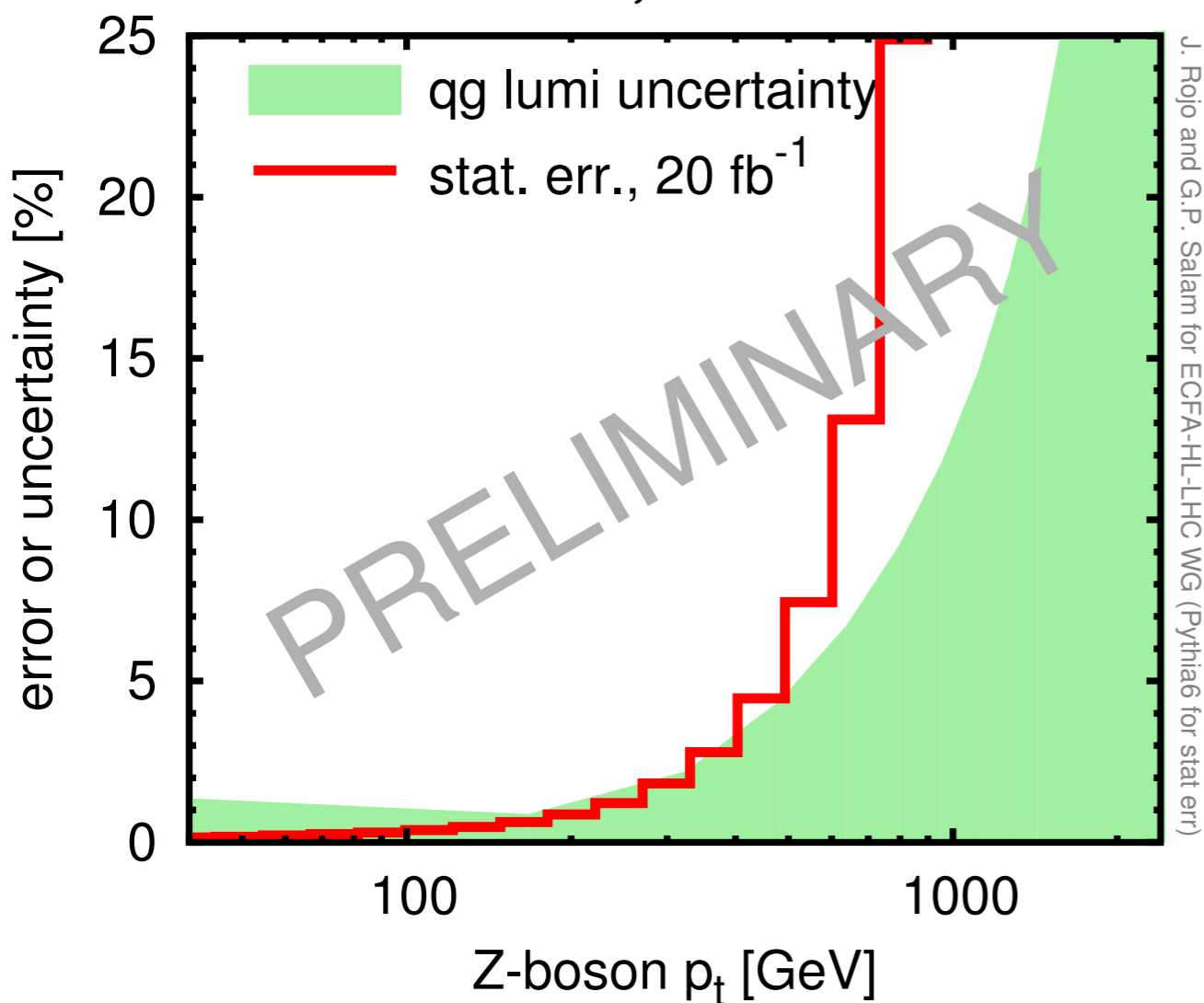
Emerging realisation that the  $Z p_t$  spectrum is a potentially very precise handle on PDFs  
[quark  $\times$  glue  $\times \alpha_s$ ]

Today, will mainly be a vital confirmation(?) of existing knowledge.

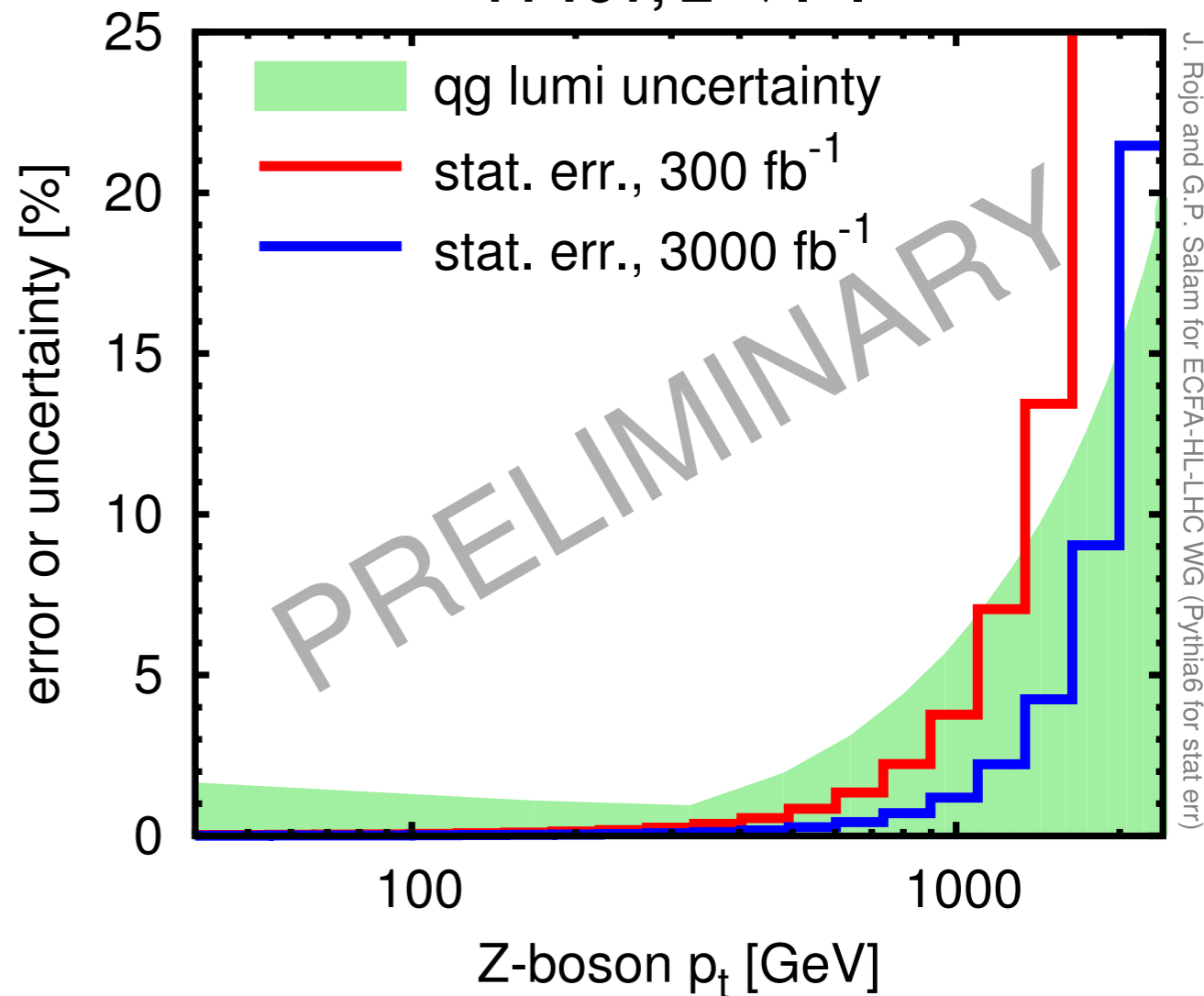
$t\bar{t}$  is also a powerful handle, cf. [1303.7215](#)

# e.g. of HL-LHC precision SM measurement: $Z p_t$ spectrum

8 TeV,  $Z \rightarrow l^+ l^-$



14 TeV,  $Z \rightarrow l^+ l^-$



For  $p_t \sim 1$  TeV, HL-LHC could bring **5x gain in precision!**  
[but only if theory prediction is good enough — today only NLO]

The thermal dimension to the exploration of the standard model  
HL-LHC brings  $\times 10-100$  in data:  $10\text{nb}^{-1} \sim 0.4\text{fb}^{-1}$  pp equivalent

## Same processes as pp , but different motivations & regimes

e.g.  $Z/\gamma+\text{jet}$  at high  $p_T$ , to study balance between jet and the Z.

Rare “probes”: open heavy flavour (jets, elliptic flow)

Quarkonium dissociation: low-pt charmonium elliptic flow, multi-differential  $\Upsilon$  production.

Low-mass di-leptons: thermal radiation  $\gamma$  ( $\rightarrow e^+e^-$ ) to map temperature during system evolution; modification of  $\rho$  spectral function to probe chiral symmetry restoration

See talk by A. Dainese for details

# CONCLUSIONS

## HL-LHC:

- Opens new channels for Higgs studies and significantly improved precision
- Gives substantial extra reach in some new-physics searches, especially for signals with small cross sections
- Will need theory improvements to get full benefits and will offer precision SM measurement opportunities

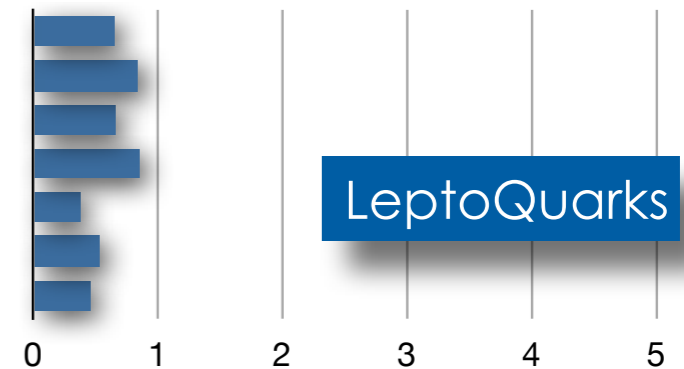
**Many of the details in the coming talks!**



# BACKUP SLIDES

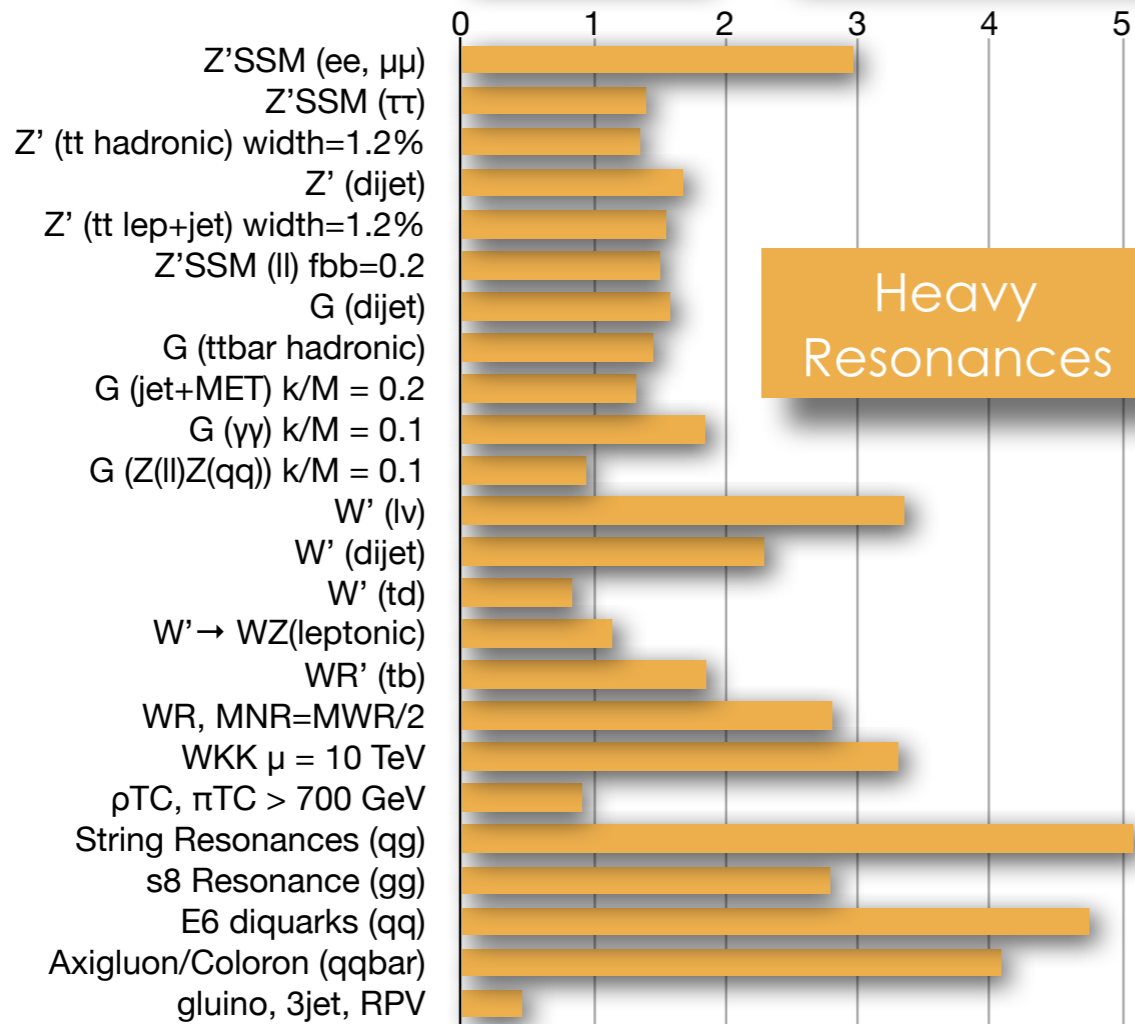
# CMS EXOTICA 95% CL EXCLUSION LIMITS (TeV)

LQ1,  $\beta=0.5$   
 LQ1,  $\beta=1.0$   
 LQ2,  $\beta=0.5$   
 LQ2,  $\beta=1.0$   
 LQ3 (bv),  $Q=\pm 1/3, \beta=0.0$   
 LQ3 (b $\tau$ ),  $Q=\pm 2/3$  or  $\pm 4/3, \beta=1.0$   
 stop (b $\tau$ )



$q^*$  (qg), dijet  
 $q^*$  (qW)  
 $q^*$  (qZ)  
 $q^*$ , dijet pair  
 $q^*$ , boosted Z  
 $e^*$ ,  $\Lambda = 2$  TeV  
 $\mu^*$ ,  $\Lambda = 2$  TeV

Compositeness



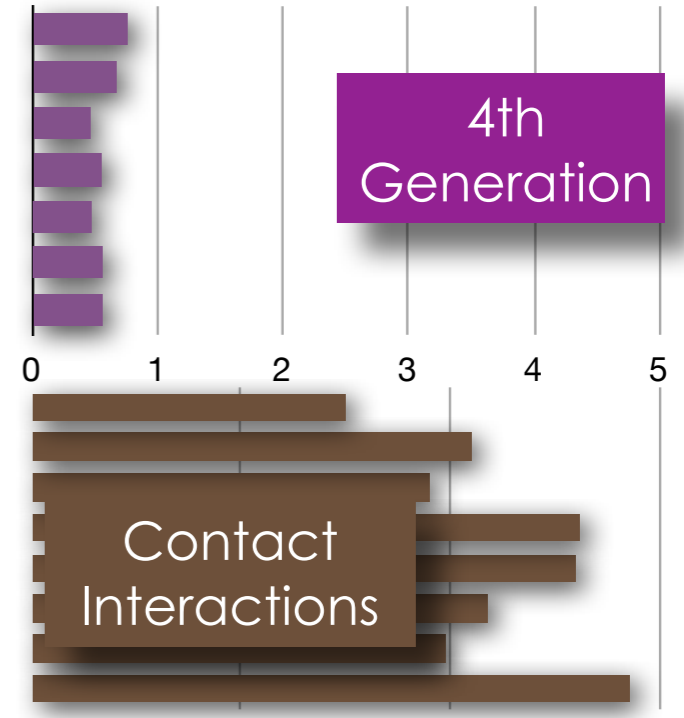
Heavy Resonances

$b' \rightarrow tW$ , (3l, 2l) + b-jet  
 $q', b'/t'$  degenerate,  $V_{tb}=1$   
 $b' \rightarrow tW$ , l+jets  
 $B' \rightarrow bZ$  (100%)  
 $T' \rightarrow tZ$  (100%)  
 $t' \rightarrow bW$  (100%), l+jets  
 $t' \rightarrow bW$  (100%), l+l

4th Generation

C.I.  $\Lambda$ , X analysis,  $\Lambda+$  LL/RR  
 C.I.  $\Lambda$ , X analysis,  $\Lambda-$  LL/RR  
 C.I.,  $\mu\mu$ , destructive LLIM  
 C.I.,  $\mu\mu$ , constructive LLIM  
 C.I., single e (HnCM)  
 C.I., single  $\mu$  (HnCM)  
 C.I., incl. jet, destructive  
 C.I., incl. jet, constructive

Contact Interactions

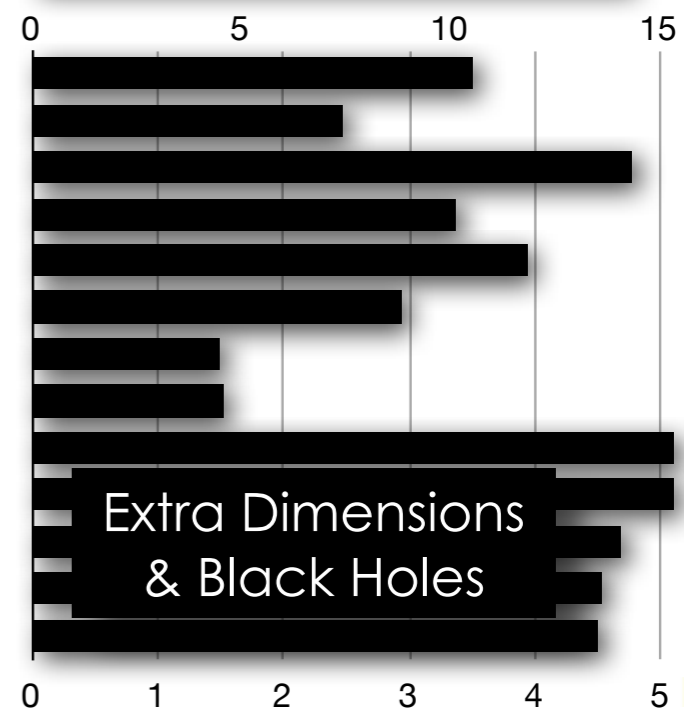


gluino, Stopped Gluino  
 stop, HSCP  
 stop, Stopped Gluino  
 stau, HSCP, GMSB  
 hyper-K, hyper- $\rho=1.2$  TeV  
 neutralino,  $c\tau < 50$ cm

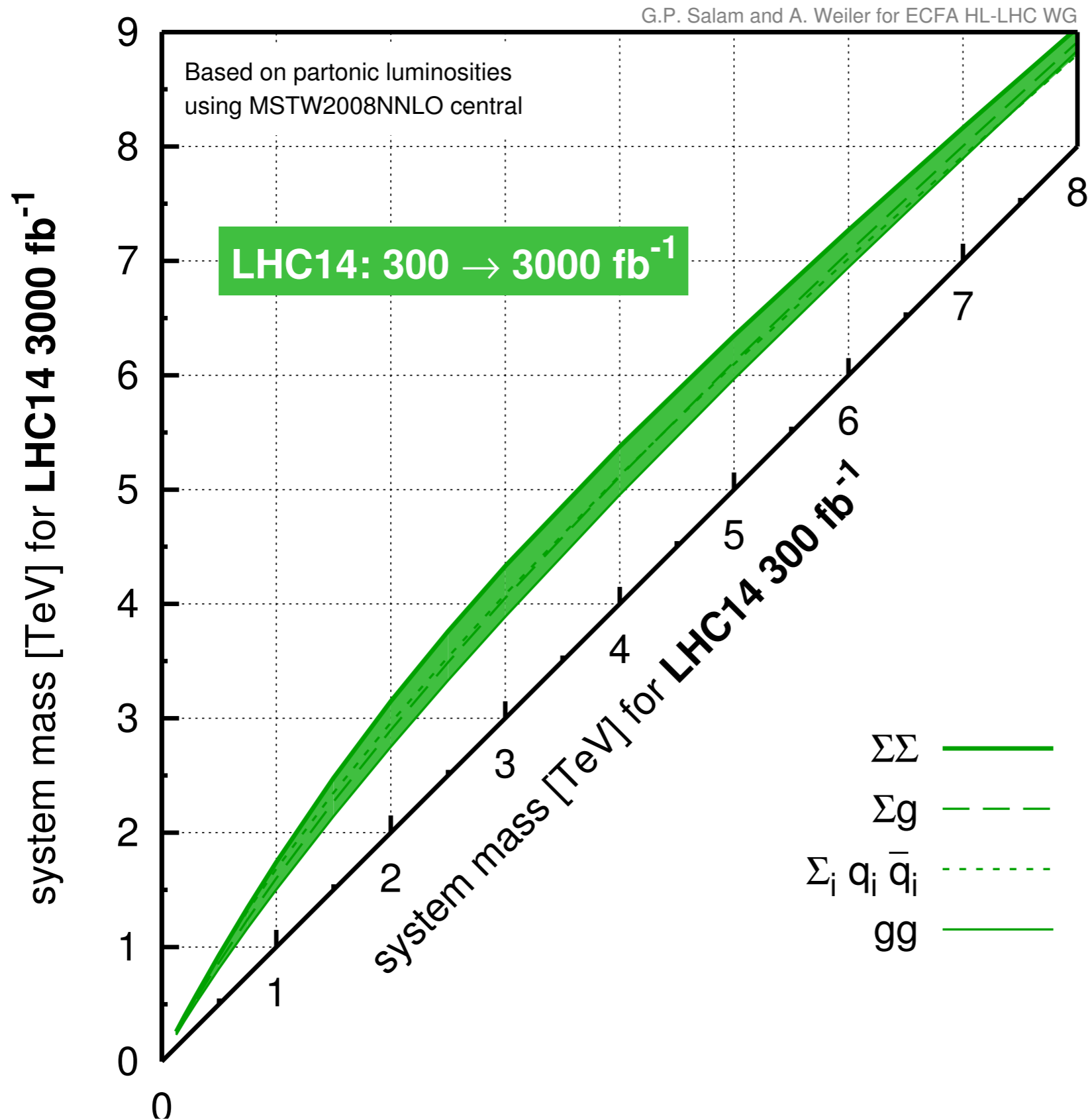
Long Lived

$M_s, \gamma\gamma, HLZ, nED = 3$   
 $M_s, \gamma\gamma, HLZ, nED = 6$   
 $M_s, ll, HLZ, nED = 3$   
 $M_s, ll, HLZ, nED = 6$   
 MD, monojet,  $nED = 3$   
 MD, monojet,  $nED = 6$   
 MD, mono- $\gamma, nED = 3$   
 MD, mono- $\gamma, nED = 6$   
 MBH, rotating, MD=3TeV,  $nED = 2$   
 MBH, non-rot, MD=3TeV,  $nED = 2$   
 MBH, boil. remn., MD=3TeV,  $nED = 2$   
 MBH, stable remn., MD=3TeV,  $nED = 2$   
 MBH, Quantum BH, MD=3TeV,  $nED = 2$

Extra Dimensions & Black Holes



Sh. Rastou



# Expected sensitivities of key heavy-flavour observables

		LHC era			HL-LHC era	
		2010–12	2015–17	2019–21	2024–26	2028–30+
$\frac{\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-)}{\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-)}$	CMS	108%	71%	47%	...	21%
	LHCb	220%	110%	60%	40%	28%
$\phi_s(B_s^0 \rightarrow J/\psi \phi)$	ATLAS	0.11	0.05–0.07	0.04–0.05	...	0.020
	LHCb	0.05	0.025	0.013	0.009	0.006
$\phi_s(B_s^0 \rightarrow \phi \phi)$	LHCb	0.18	0.12	0.04	0.026	0.017
$\gamma$	LHCb	7°	4°	1.7°	1.1°	0.7°
	Belle2	—	11°	2°	1.5°	—
$A_\Gamma(D^0 \rightarrow K^+ K^-)$	LHCb	$3.4 \times 10^{-4}$	$2.2 \times 10^{-4}$	$0.9 \times 10^{-4}$	$0.5 \times 10^{-4}$	$0.3 \times 10^{-4}$
	Belle2	—	$18 \times 10^{-4}$	$4\text{--}6 \times 10^{-4}$	$3\text{--}5 \times 10^{-4}$	—
$q_0^2 A_{\text{FB}}(K^{*0} \mu^+ \mu^-)$	LHCb	11%	5%	3.1%	2.1%	1.4%
	Belle2	—	16%	2.2%	1.6%	—
$t \rightarrow qZ$	ATLAS CMS					
$t \rightarrow q\gamma$	ATLAS CMS					

See talk by M.-H. Schune for more numbers & details

E.g. Ahrens et al (1008.3162), who say uncertainty can be made small,  $O(3\%)$  in their framework even without  $N^3LO$ ,

Baglio & Djouadi (1012.0530), David & Passarino (1307.1843), who argue it may be larger than widely accepted so far.

## Bottom line?

We'll know more in a year or two, once  $N^3LO$  appears.

On 15 year timescale, TH community will learn yet more and hopefully move towards greater consensus on the uncertainties.

But ratios of observables (e.g.  $ZZ^* / \gamma\gamma$ ) not affected.