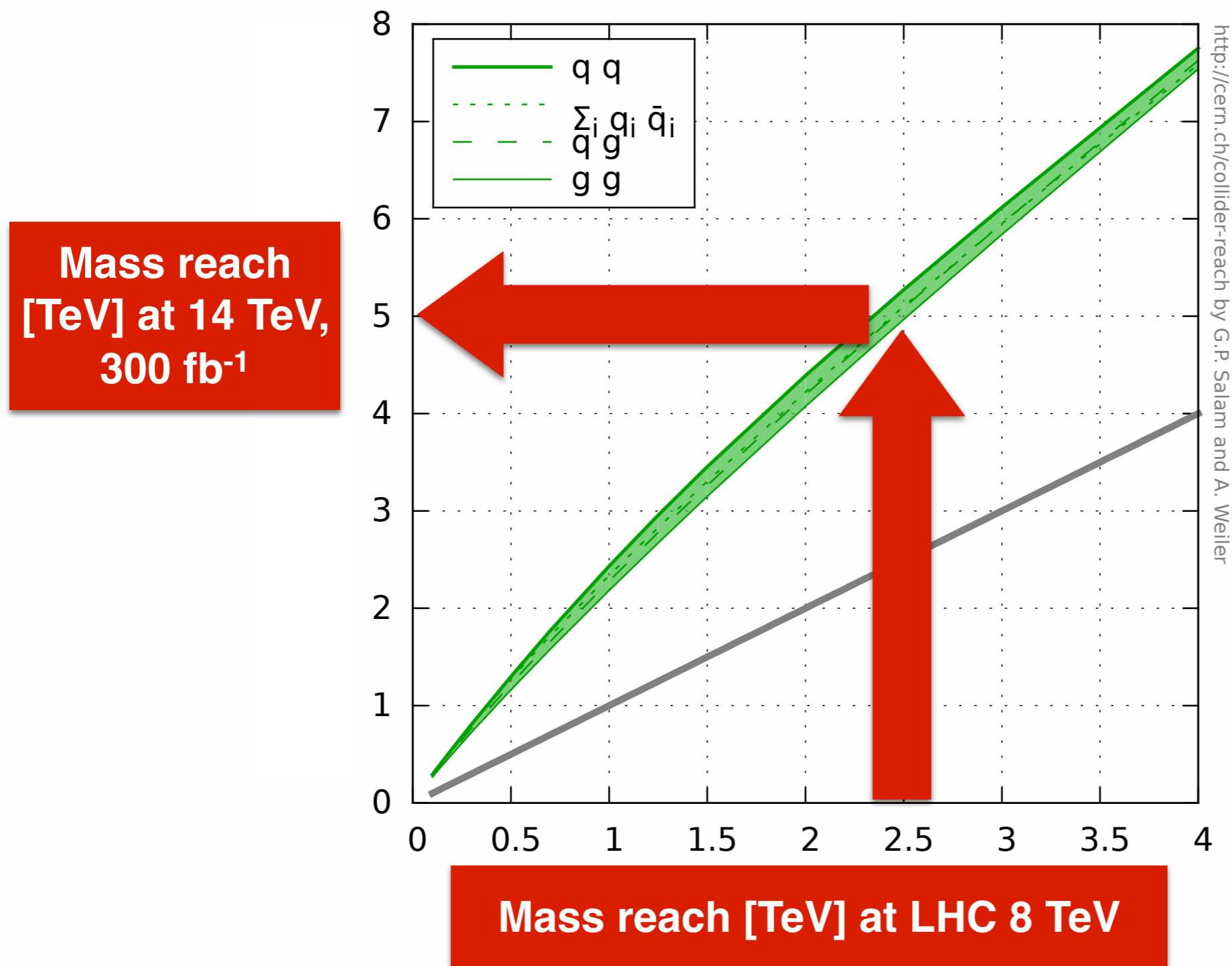


A theorist's perspective on boosted boson tagging

Gavin Salam (CERN-TH)

Boosted Boson Workshop,
CERN, March 25, 2014

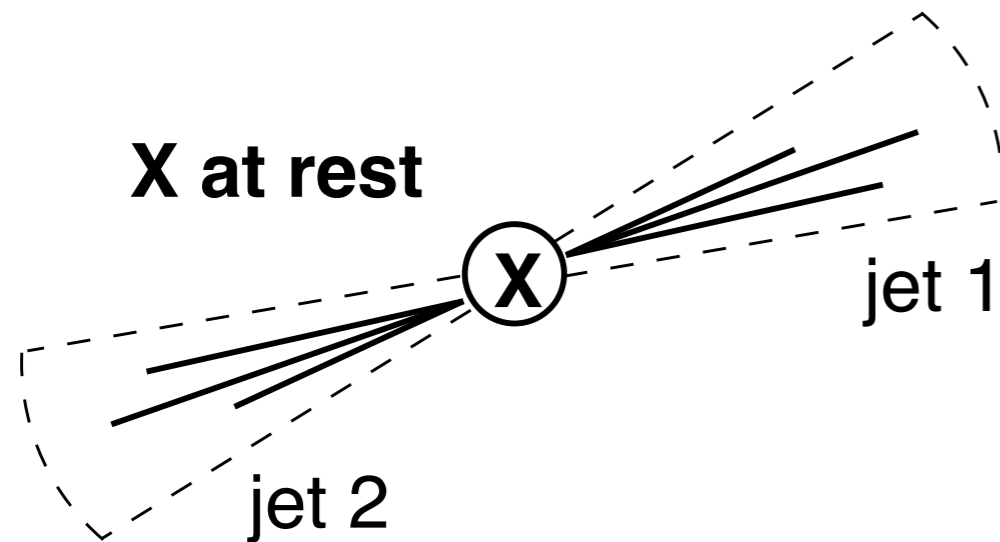
14 TeV, after a few years, will roughly double the mass reach of LHC searches



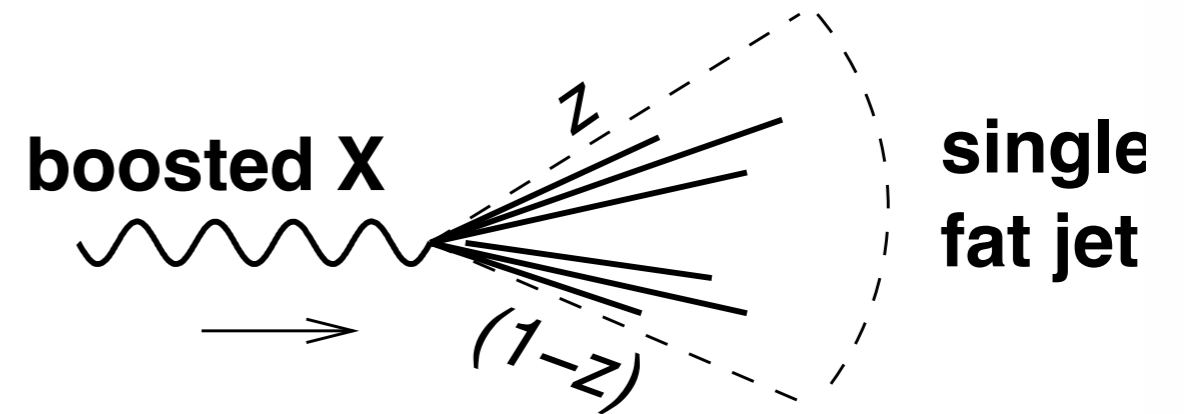
Boosted hadronic decays

($X = W, Z, H, \text{top, new particle}$)

Normal analyses: two quarks from $X \rightarrow q\bar{q}$ reconstructed as two jets

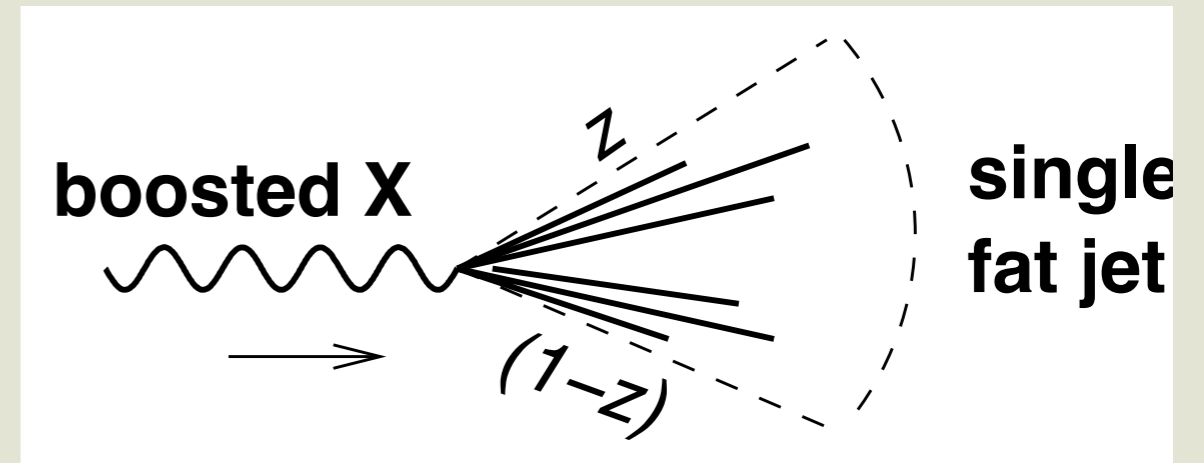


High- p_t regime: EW object X is boosted, decay is collimated, $q\bar{q}$ both in same jet



Happens for $p_t \gtrsim 2m/R$
 $p_t > 400 \text{ GeV}$ for $m = m_W, R = 0.4$

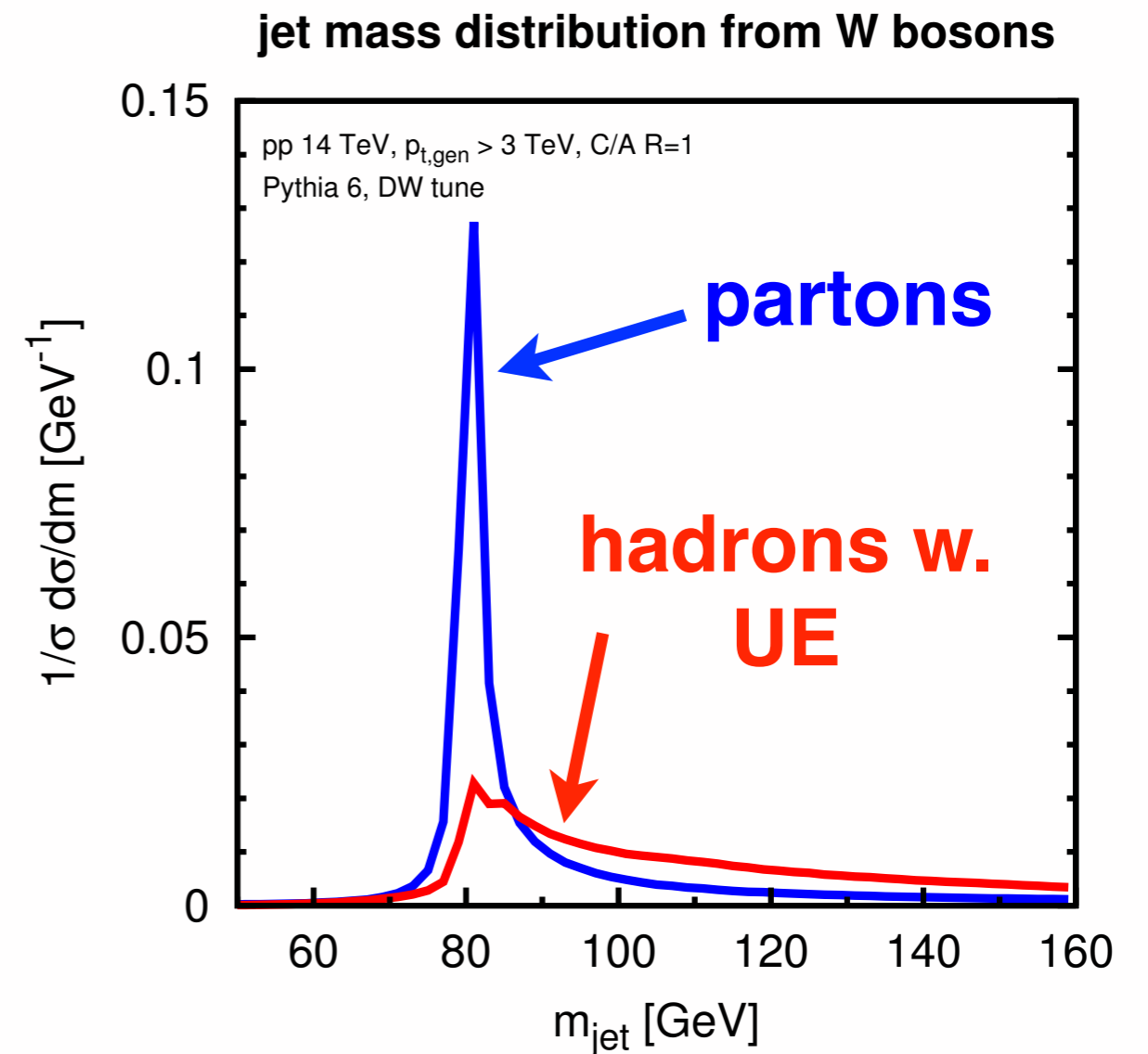
Most obvious way of detecting a boosted decay is through the mass of the jet



But jet mass is **poor** in practice:

e.g., narrow W resonance highly smeared by QCD radiation

(mainly underlying event/pileup)

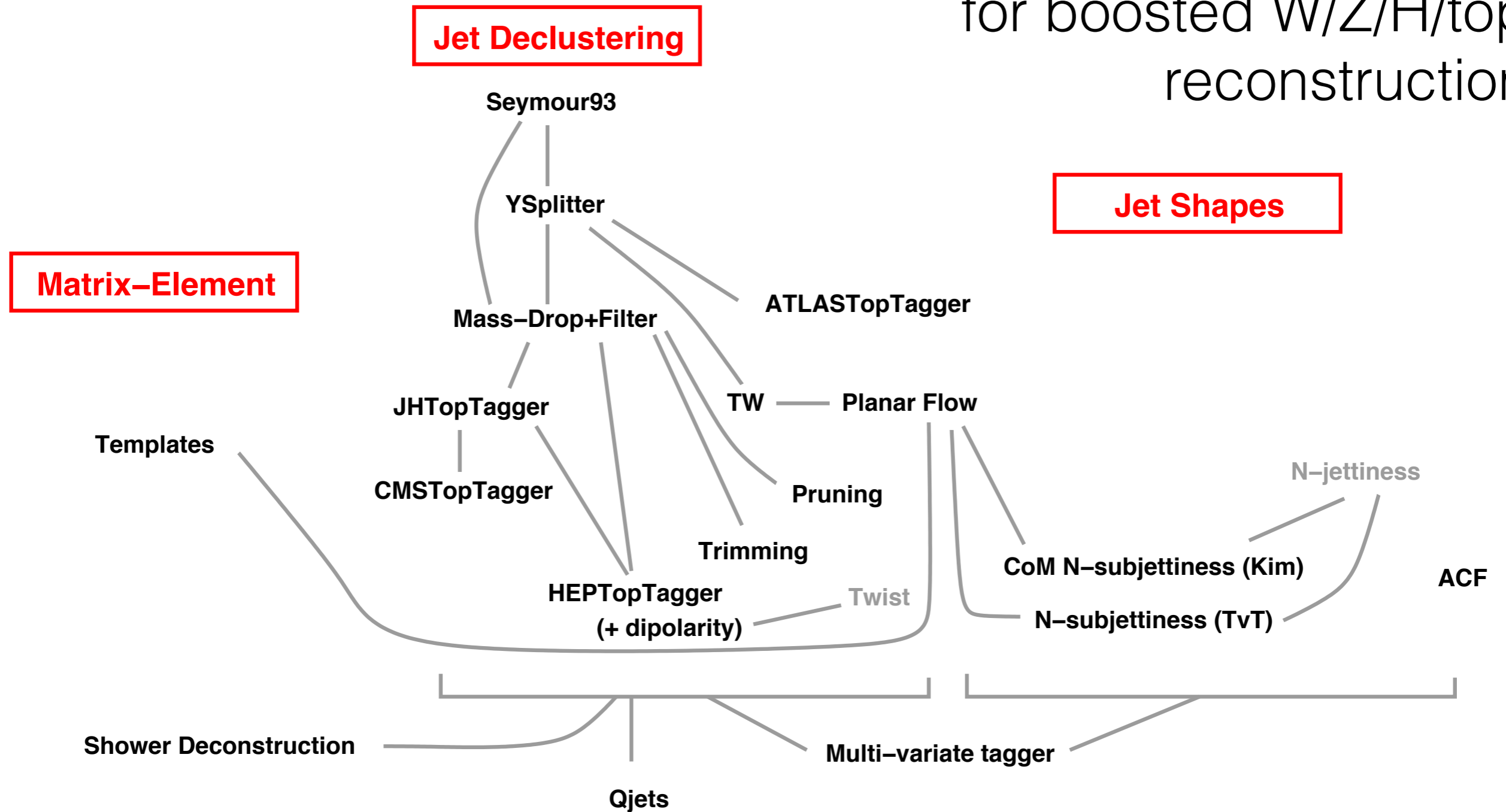


As a field we've devised O(10-20) powerful methods to tag jet substructure.

Many of the methods have been tried out in searches and work; these kinds of methods will be crucial for searches in the years to come.

Very active research field

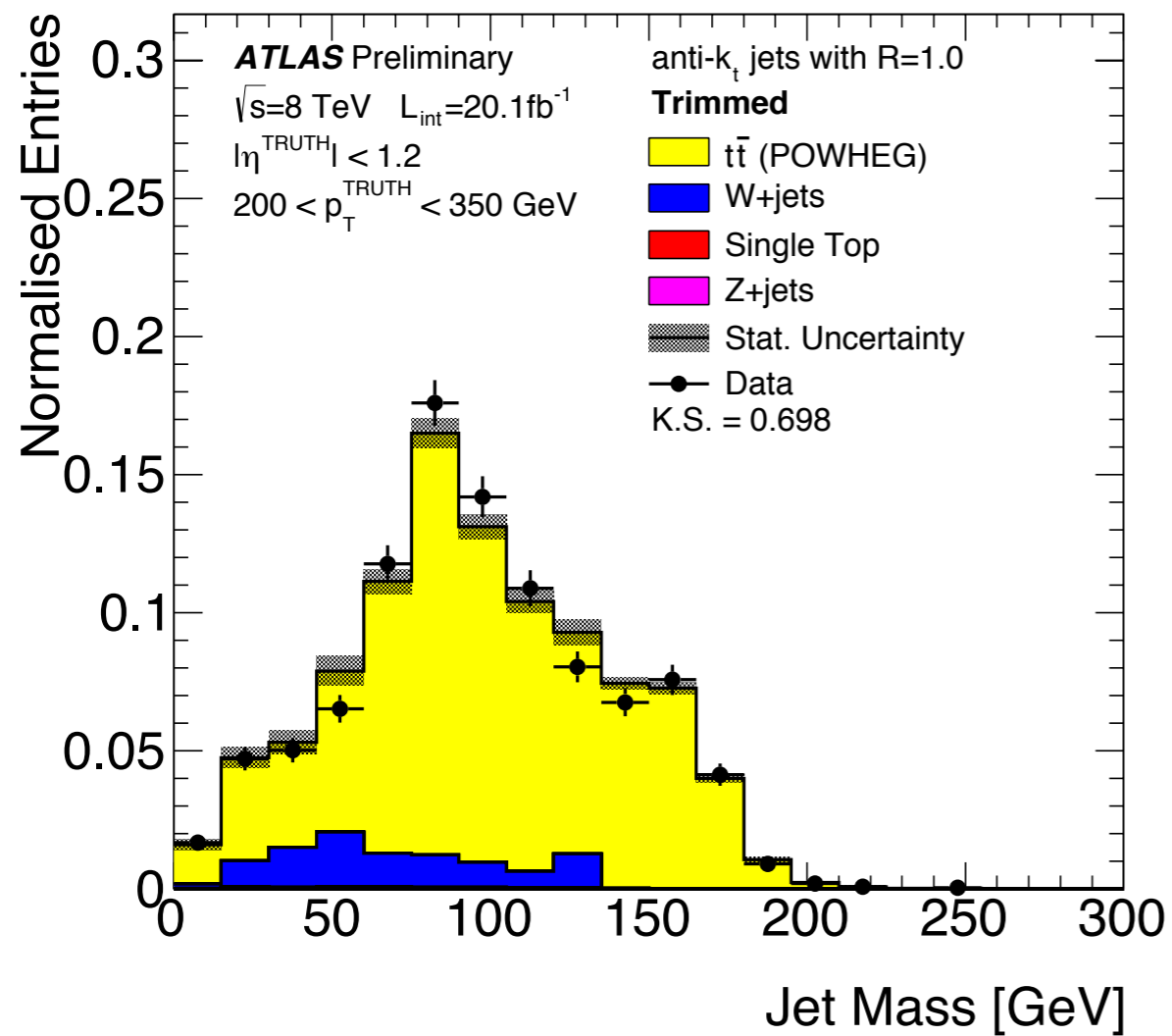
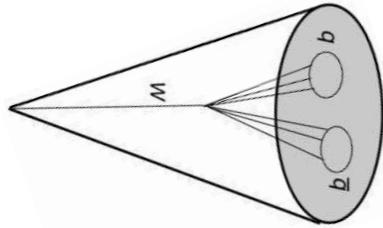
Some of the tools developed for boosted W/Z/H/top reconstruction



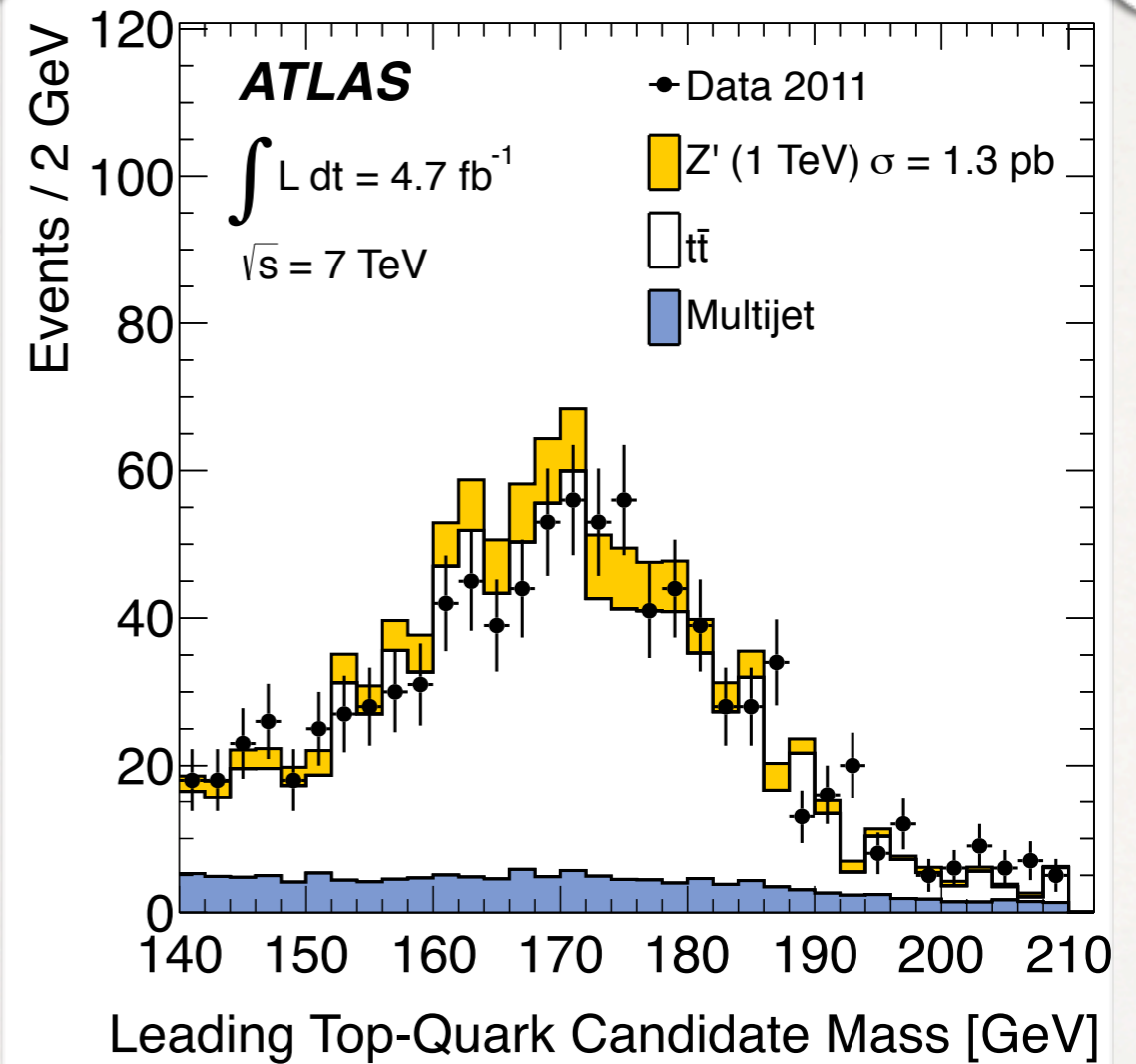
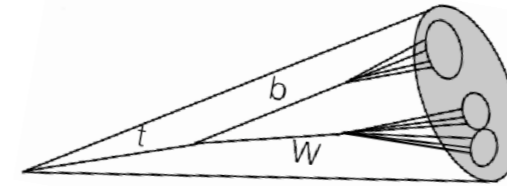
apologies for omitted taggers, arguable links, etc.

Seeing W's and tops in a single jet

W's in a single jet



tops in a single jet



As a field we've devised $O(10-20)$ powerful methods to tag jet substructure.

Many of the methods have been tried out in searches and work; these kinds of methods will be crucial for searches in the years to come.

But from outside, the many methods make the field look pretty confusing.

And from inside, I get the impression we don't always know *why* or *how* the methods work – which is bad if we're looking for robustness.

**Is it time to get back to basics?
I.e. to start understanding our tools?**

What was the original motivation?

Normal $R=0.4/0.5$ jet finding fails to find one jet per prong of a boosted $[W/Z/H/top/NP]$ hadronic decay.

We need to make sure that this doesn't prevent us from using EW-scale particles in TeV scale searches.

Question #1:

To what extent are the things we do with “normal” jets (and leptons) mirrored in the things we're doing with “fat” jets?

What have we found out in the meantime?

There's a huge number of things you can do with jet substructure.

Many of the things appear to improve mass resolution, background rejection, etc. [at least in MC simulation]

Question #2:

How should we balance improvements v. “complexity” of method?

What are we comfortable with?

Resolved Analysis

Find one jet/prong

Cut on jet p_t , Δy , ...

What are we comfortable with?

Resolved Analysis

Find one jet/prong

Cut on jet p_t , Δy , ...



Fat-jet Analysis

Find subjets



Cut on subjet z , ΔR , ...

[MDT/Prune/Trim/Filt/XYZTopTagger/Template ...]

What are we comfortable with?

Resolved Analysis

Find one jet/prong

Cut on jet p_t , Δy , ...

Isolation cut for
colourless leptons, γ

Cut on radiation in jet
for q/g discrimination



Fat-jet Analysis

Find subjets

Cut on subjet z , ΔR , ...

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Cut on subjet z , ΔR , ...

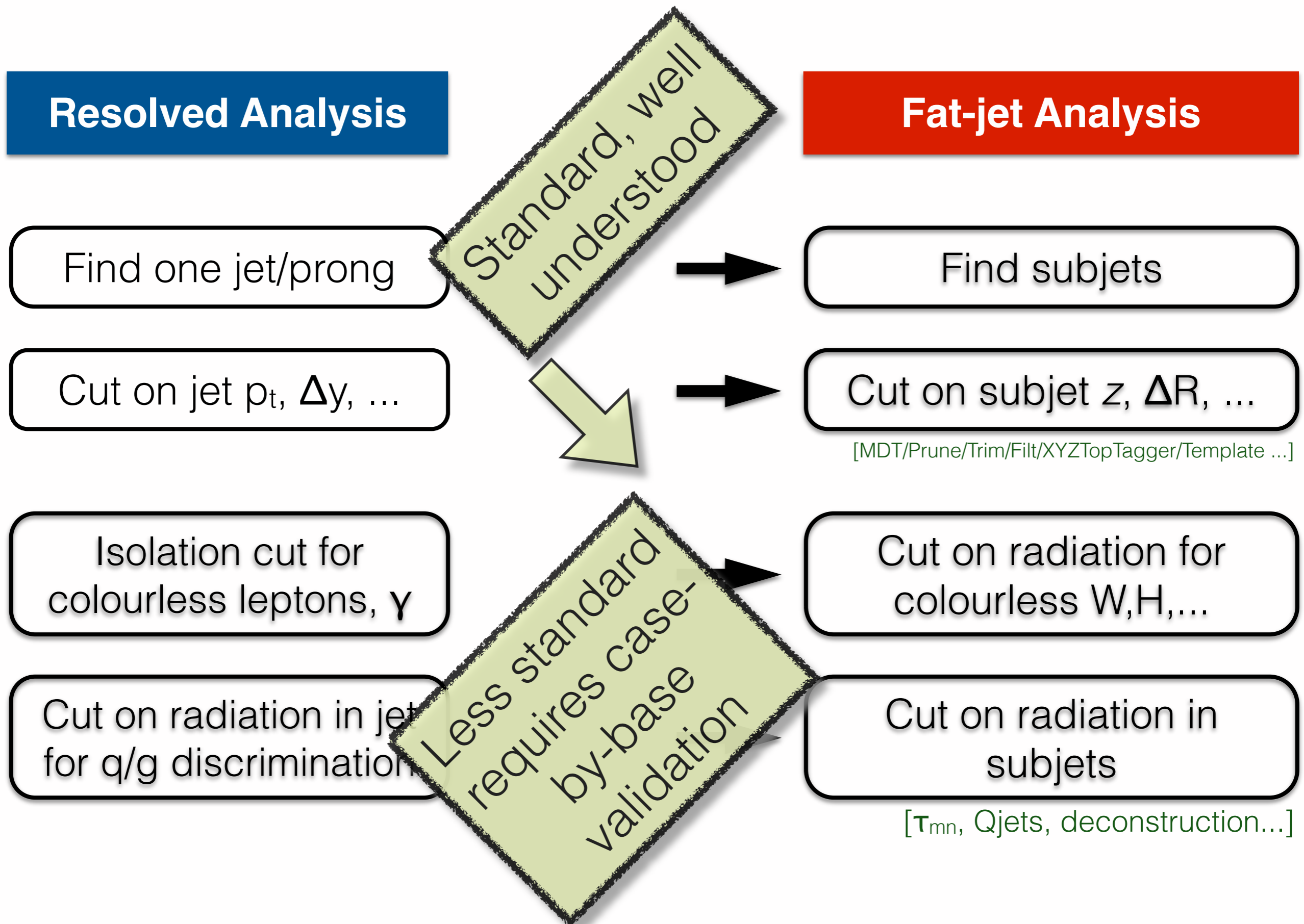
[MDT/Prune/Trim/Filt/XYZTopTagger/Template ...]

Cut on radiation for
colourless W,H,...

Cut on radiation in
subjets

[τ_{mn} , Qjets, deconstruction...]

What are we comfortable with?



analytical understanding: why?

Better Insight

Can guide taggers' use
in experimental
analyses

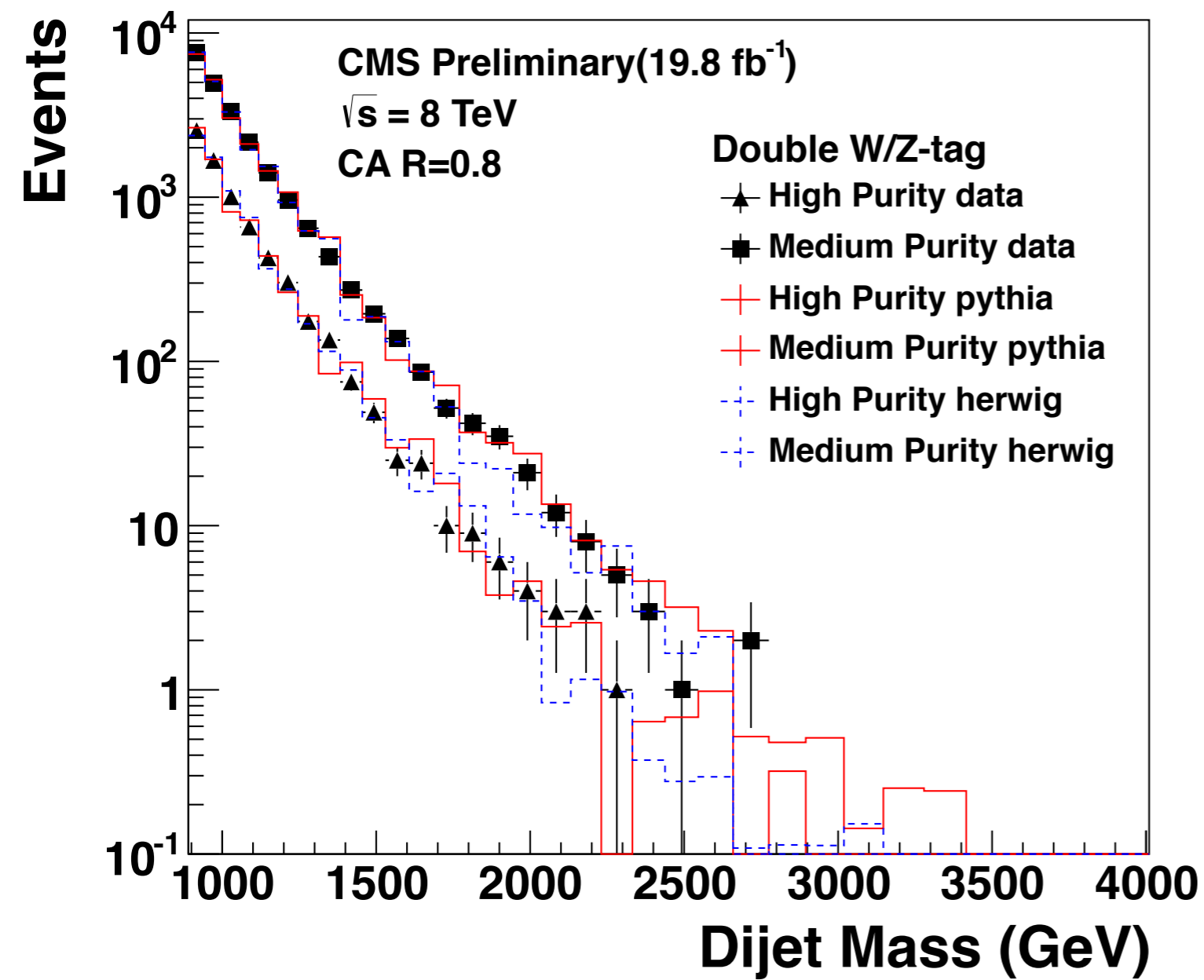
It may help us design
better taggers

Robustness

You know what you
predict, what you don't

Unlike MC, you have
powerful handles for
cross-checks & accuracy
estimates

There is a "right" answer



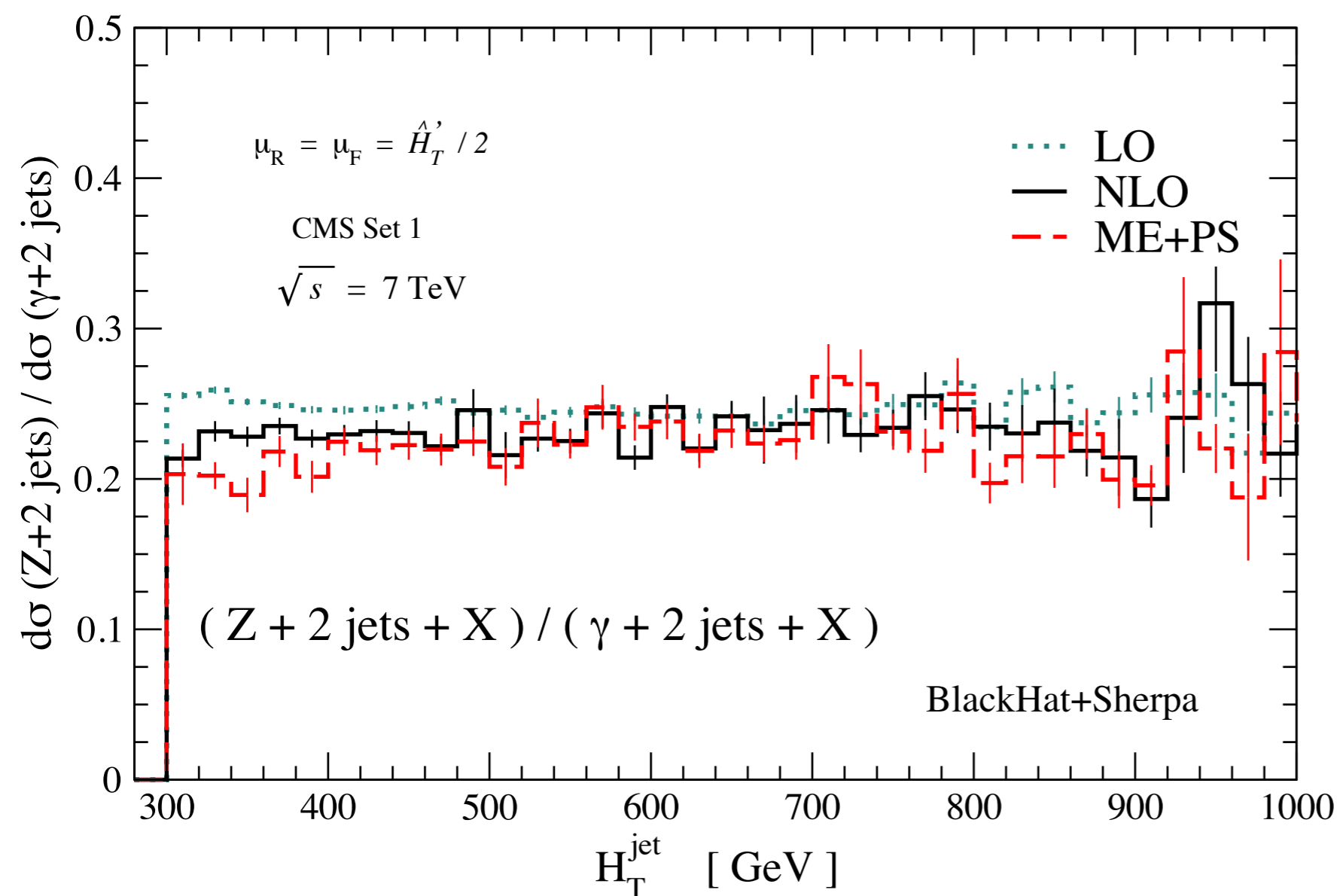
Search for resonances
 in doubly-tagged dijet
 events.

Tagging = pruning +
 tau21 cut

Note different Herwig++
 and Pythia6 shapes

Understanding your taggers means you know what tools you can safely use with them

For robustness, you can then choose taggers whose distributions can be predicted in many ways



Just like
MET($Z \rightarrow \nu\nu$) in
multijets is reliably
estimated from
 γ +jets because
multiple types of
calculations of the
ratio agree

The key variables

For phenomenology

Jet mass: m

*[as compared to W/Z/H
or top mass]*

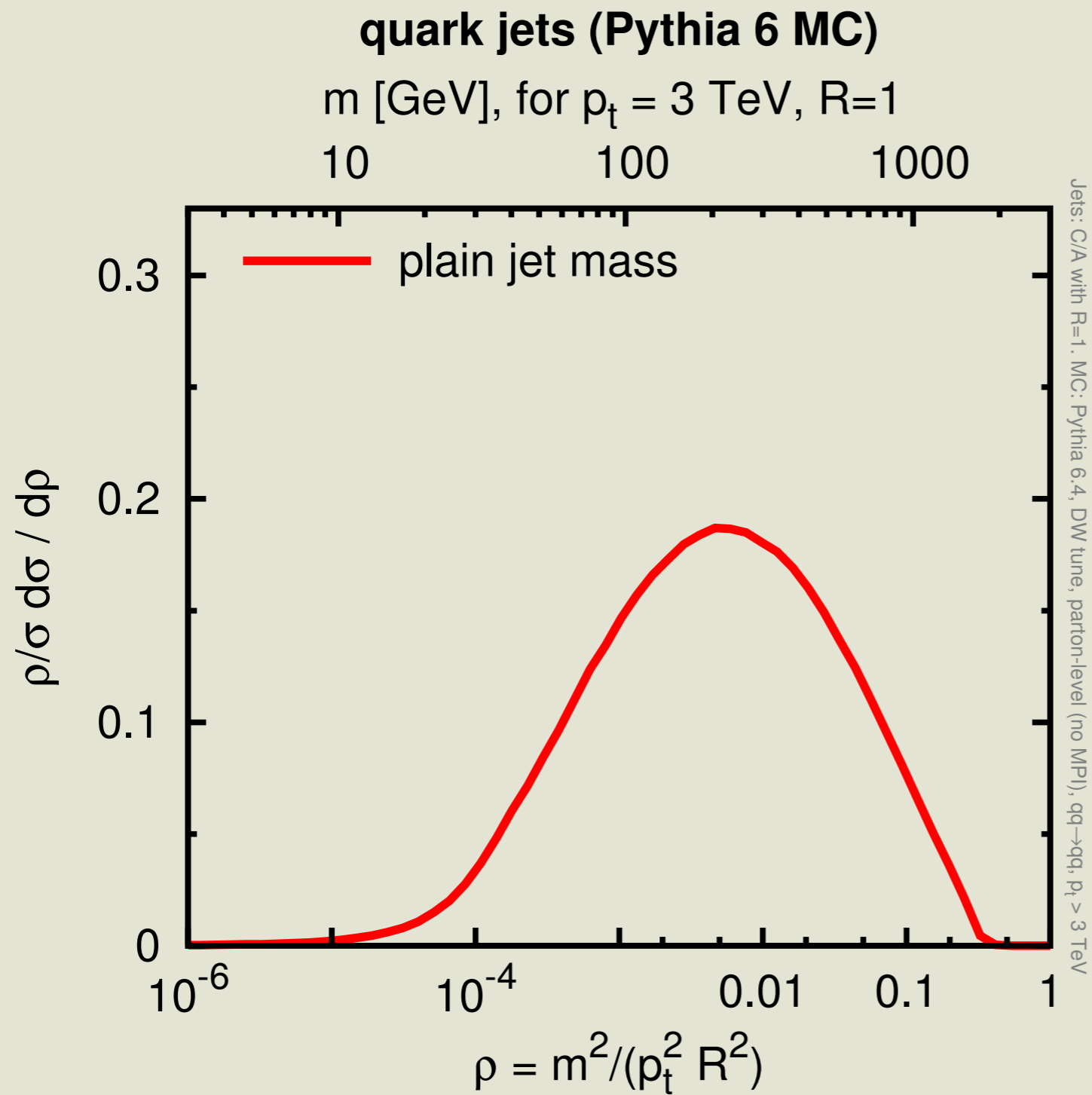
For QCD calculations

$$\rho = \frac{m^2}{p_t^2 R^2}$$

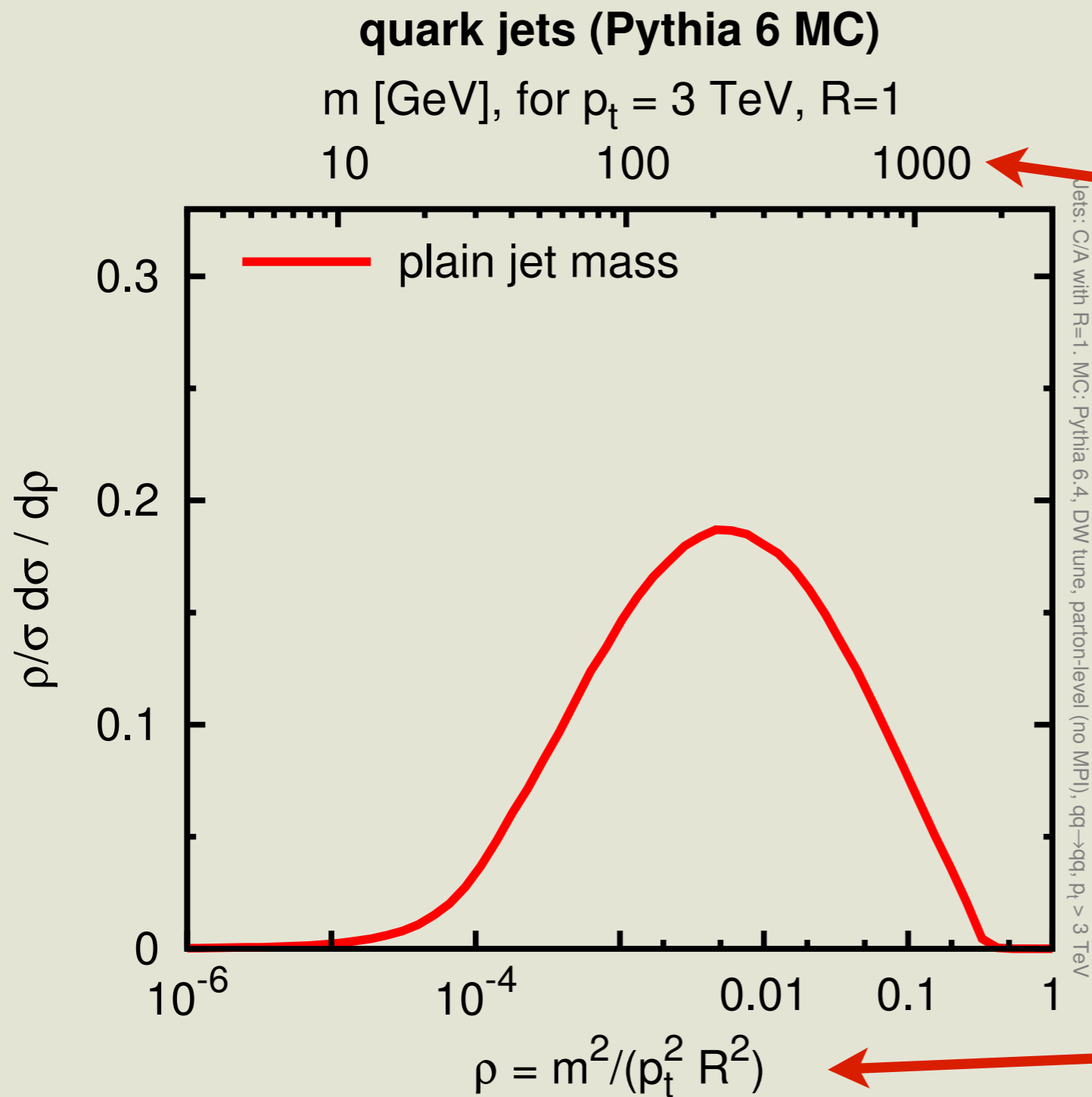
*[R is jet opening angle
– or radius]*

Because ρ is invariant under
boosts along jet direction

Start with “plain” jet mass



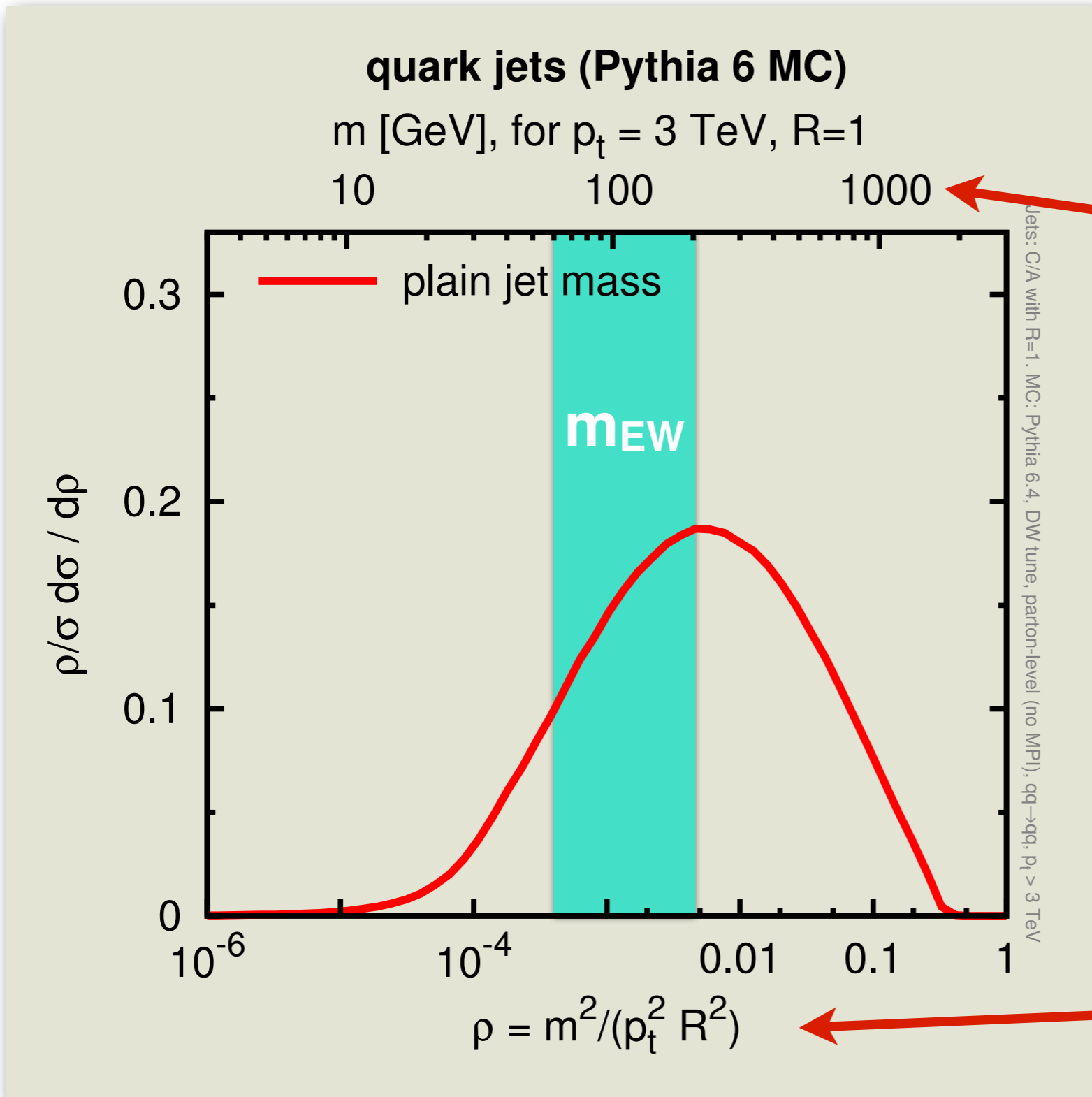
Start with “plain” jet mass



Physical mass for
3 TeV, R=1 jets

$\rho \sim$ Rescaled mass²
(i.e. the QCD variable)

Start with “plain” jet mass



Physical mass for
3 TeV, R=1 jets

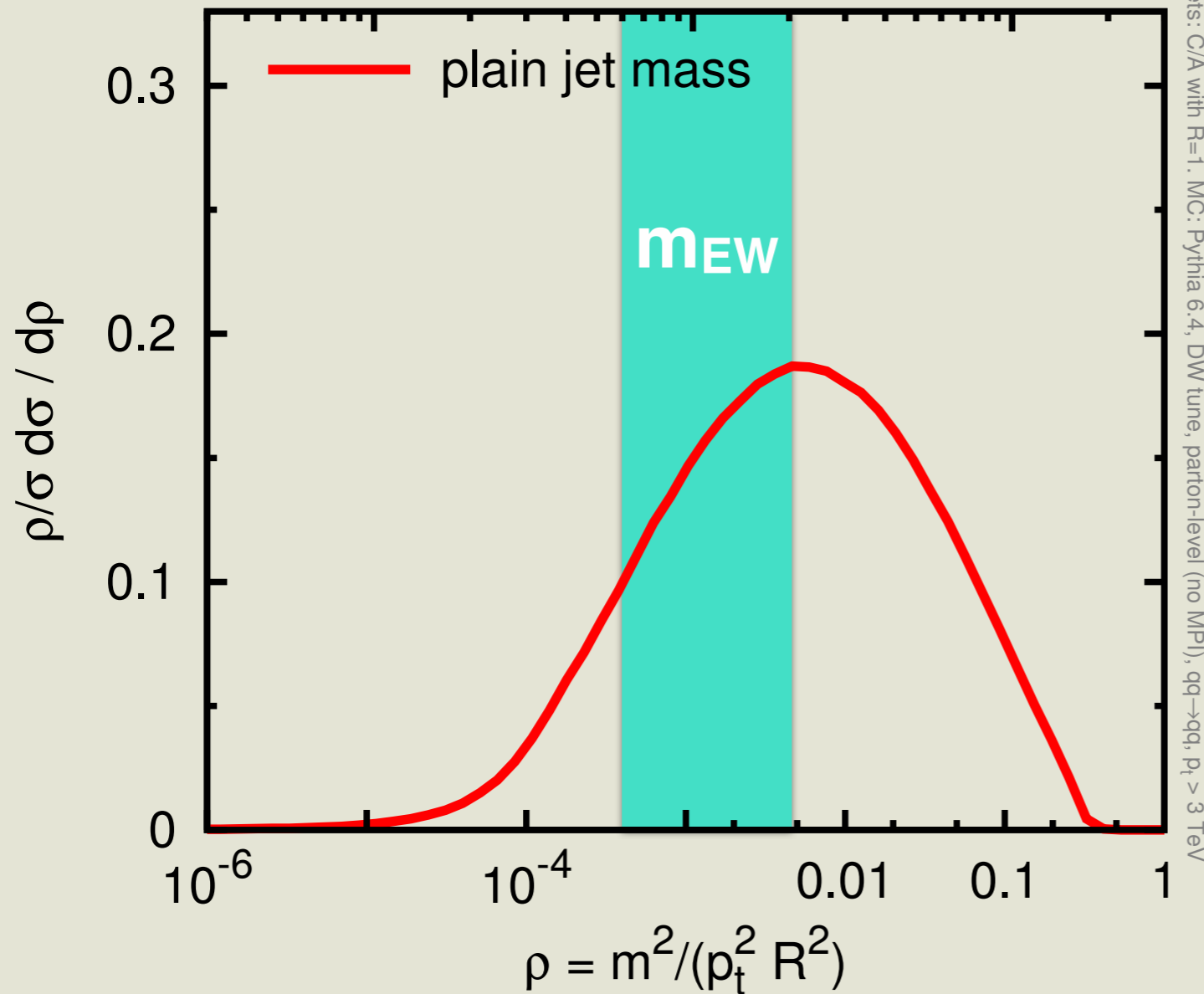
$\rho \sim$ Rescaled mass²
(i.e. the QCD variable)

Start with “plain” jet mass

quark jets (Pythia 6 MC)

m [GeV], for $p_t = 3$ TeV, $R=1$

10 100 1000



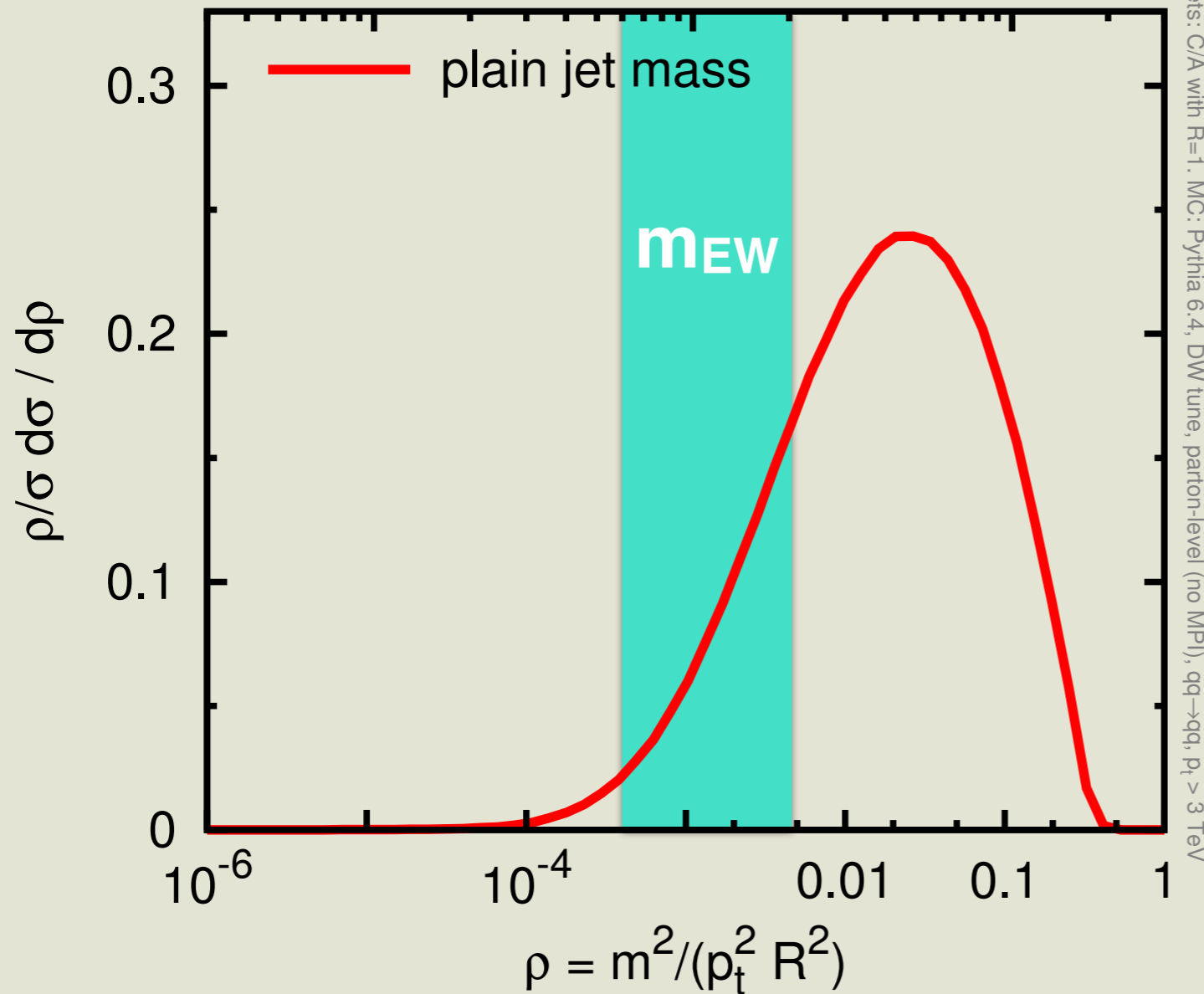
Quark and gluon jets are different, so we treat them separately

Start with “plain” jet mass

gluon jets (Pythia 6 MC)

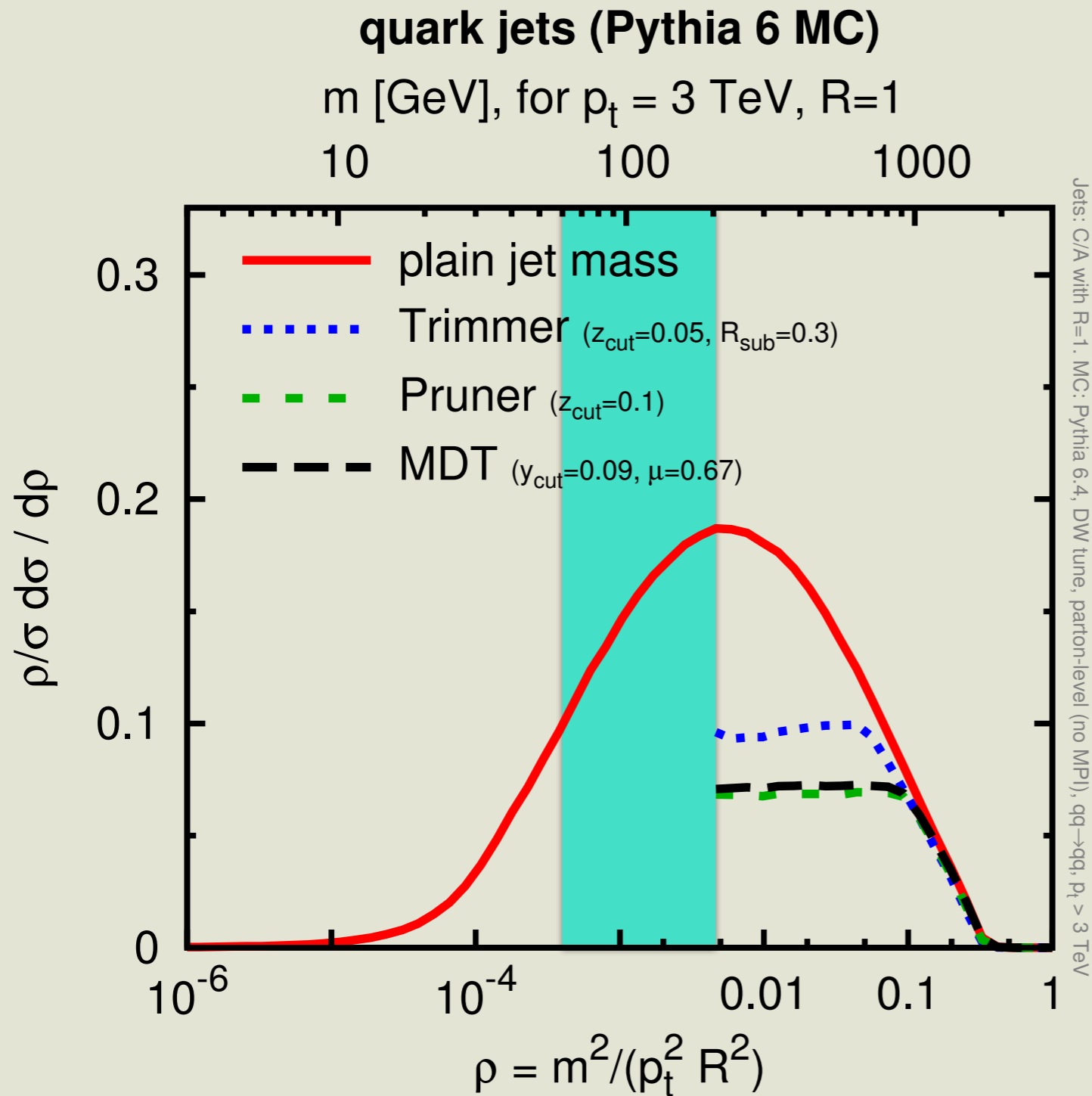
m [GeV], for $p_t = 3$ TeV, $R=1$

10 100 1000



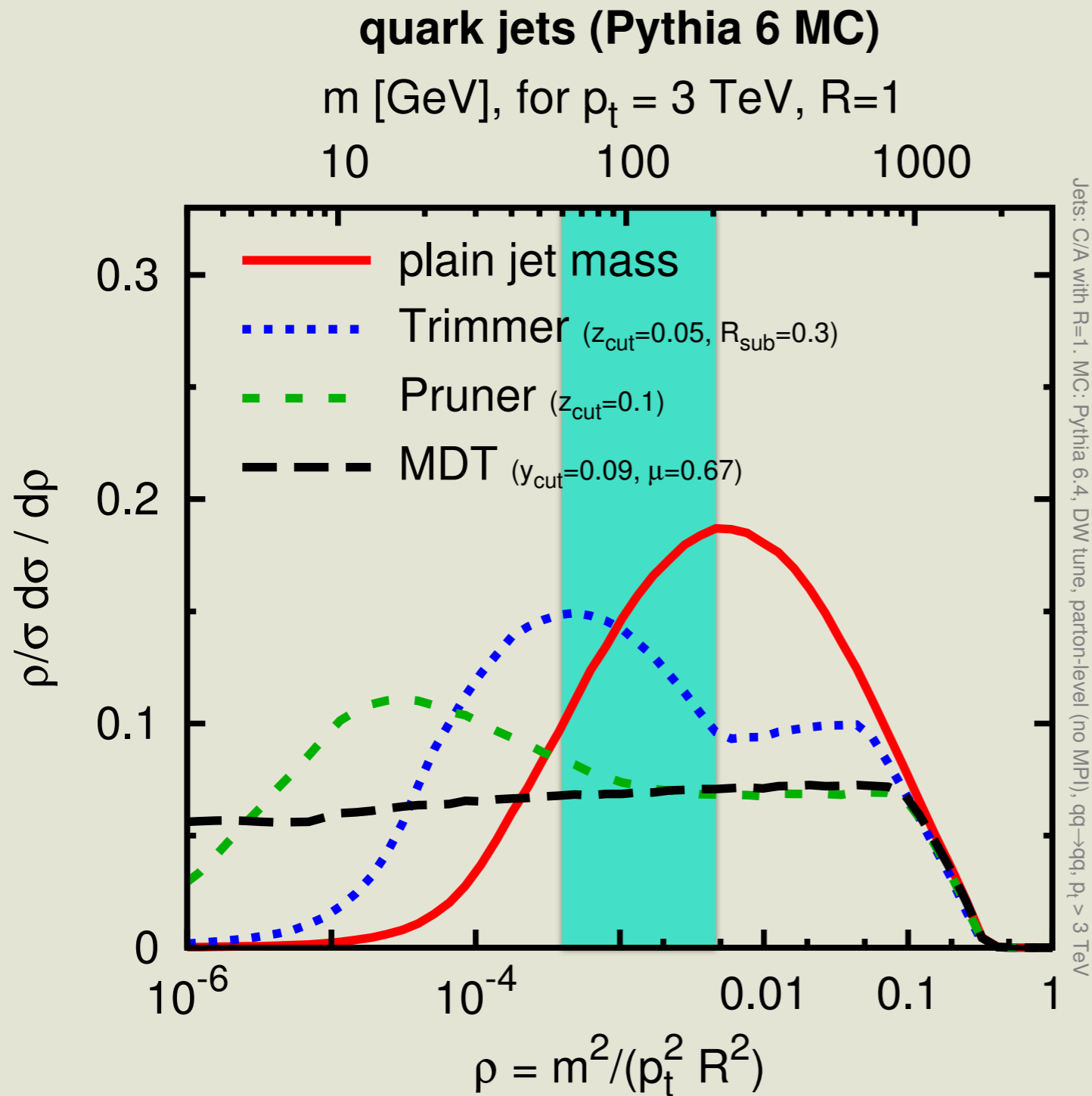
Quark and **gluon** jets are different, so we treat them separately

Now examine “taggers/groomers”



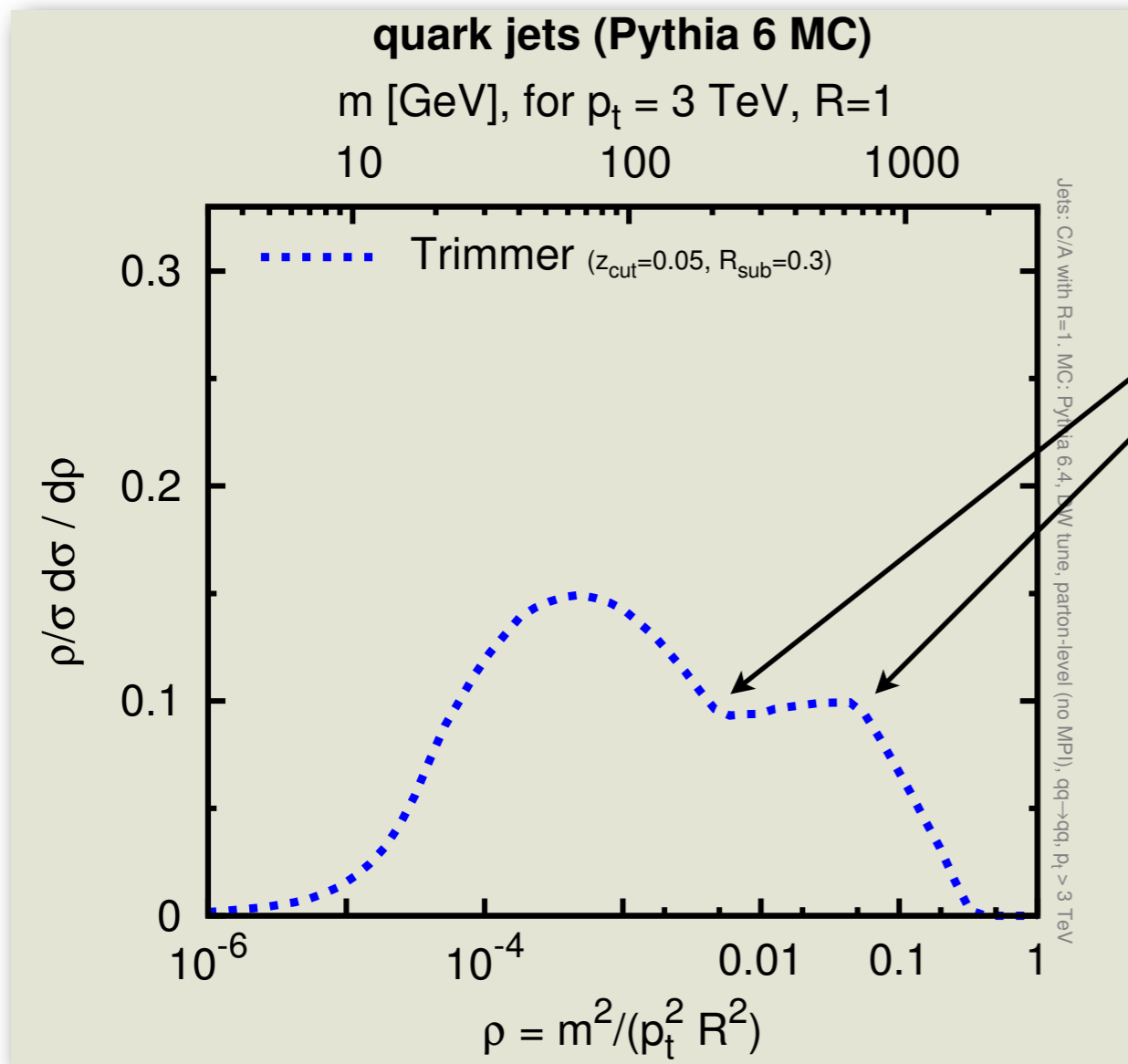
Different taggers
can be
quite similar

Now examine “taggers/groomers”



But only for a limited range of masses

What do we want to find out?



Where exactly are the kinks?
How do their locations depend
on $z_{\text{cut}}, R_{\text{sub}}$?

Kinks are especially
dangerous for data-
driver backgrounds

What physics is relevant in the
different regions?

Because then you have
an idea of how well you
control it

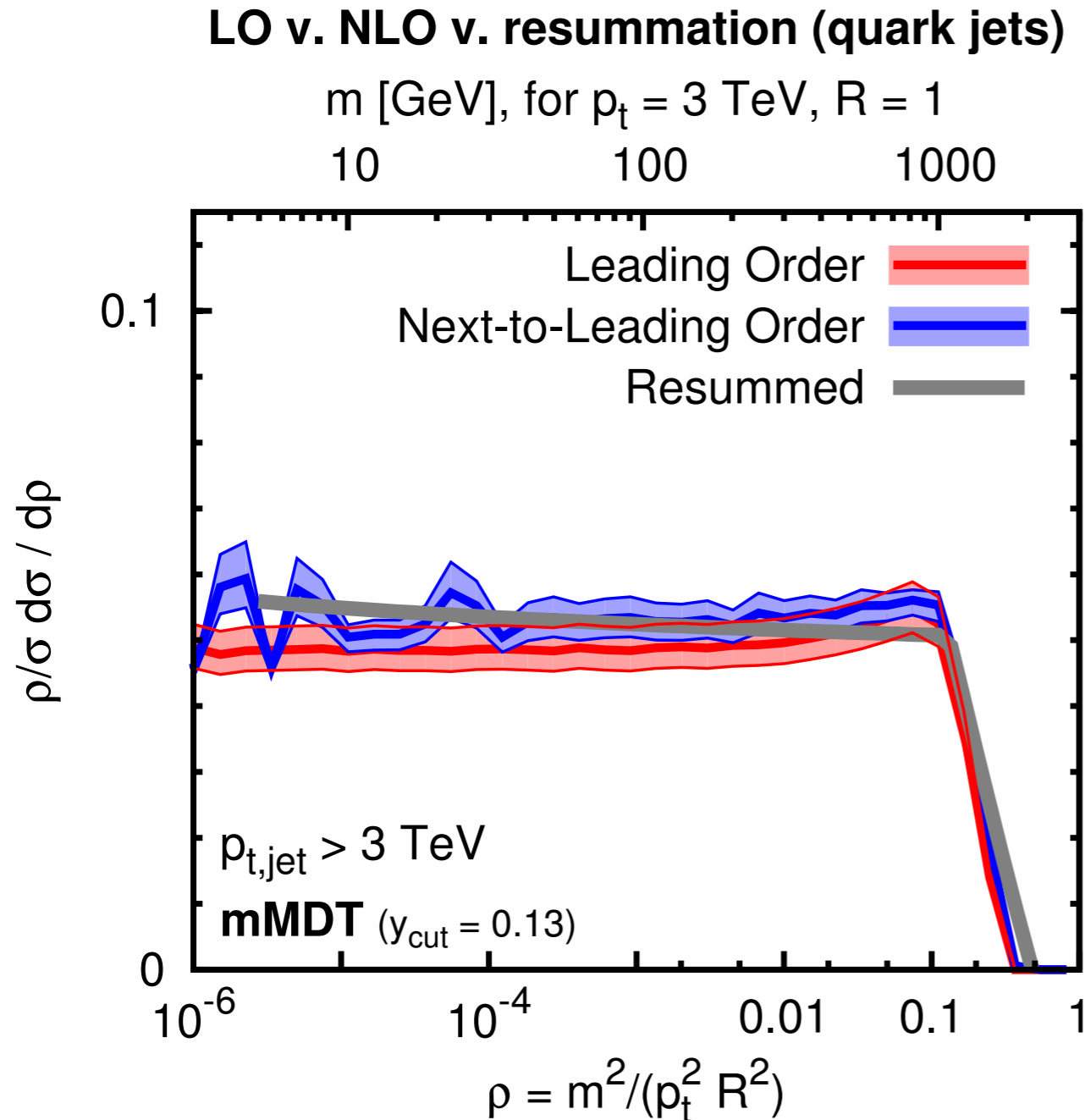
And maybe you can
make better taggers

Does tagging/grooming make precise QCD impossible?

Matt Schwartz @ Boost 2012



mMDT resummation v. fixed order



NLO from NLOJet++

Because we only have single logs, fixed-order is valid over a broader than usual range of scales

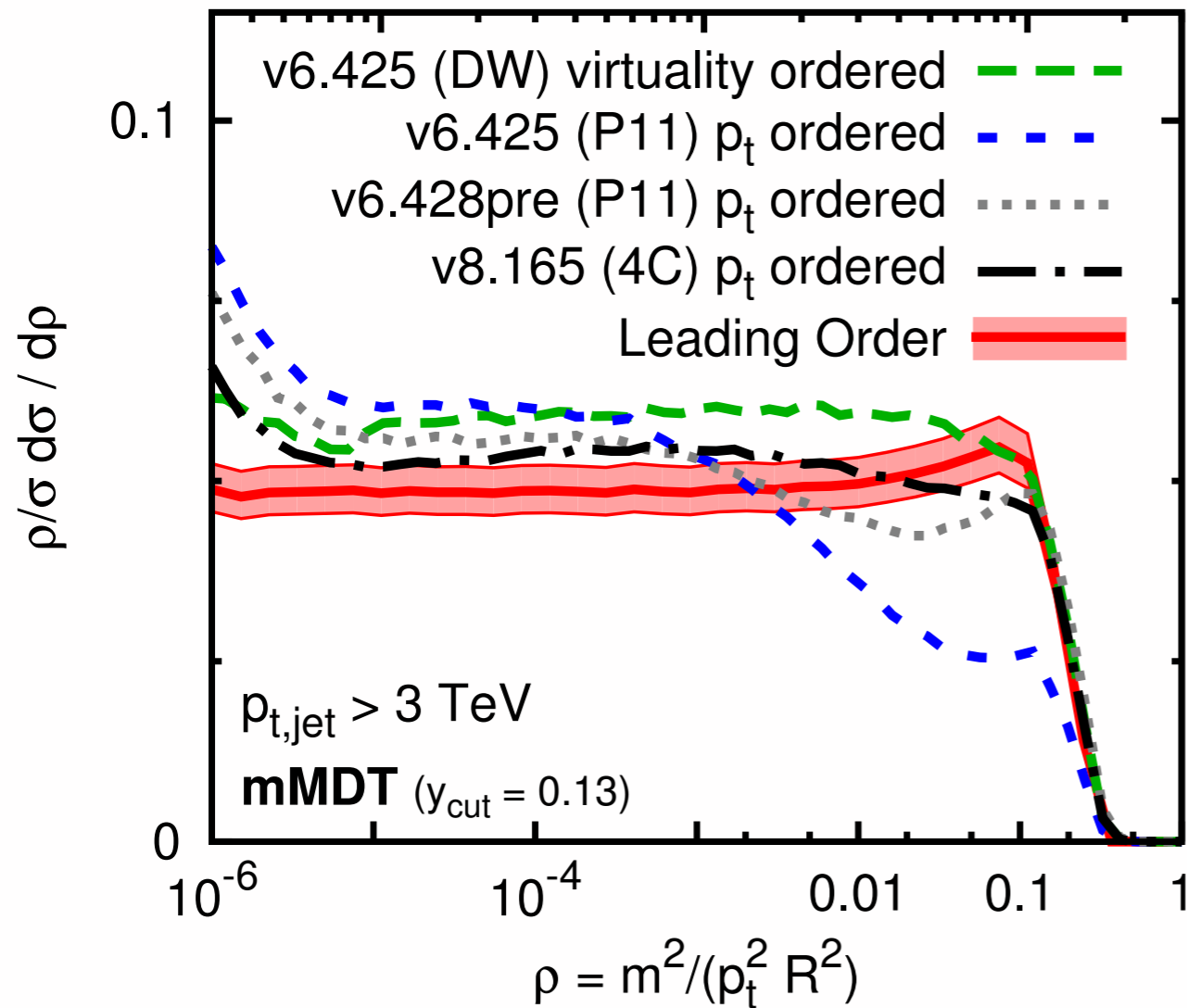
(helped by fortuitous cancellation between running coupling and single-log Sudakov)

mMDT: comparing many showers

LO v. Pythia showers (quark jets)

m [GeV], for $p_t = 3$ TeV, $R = 1$

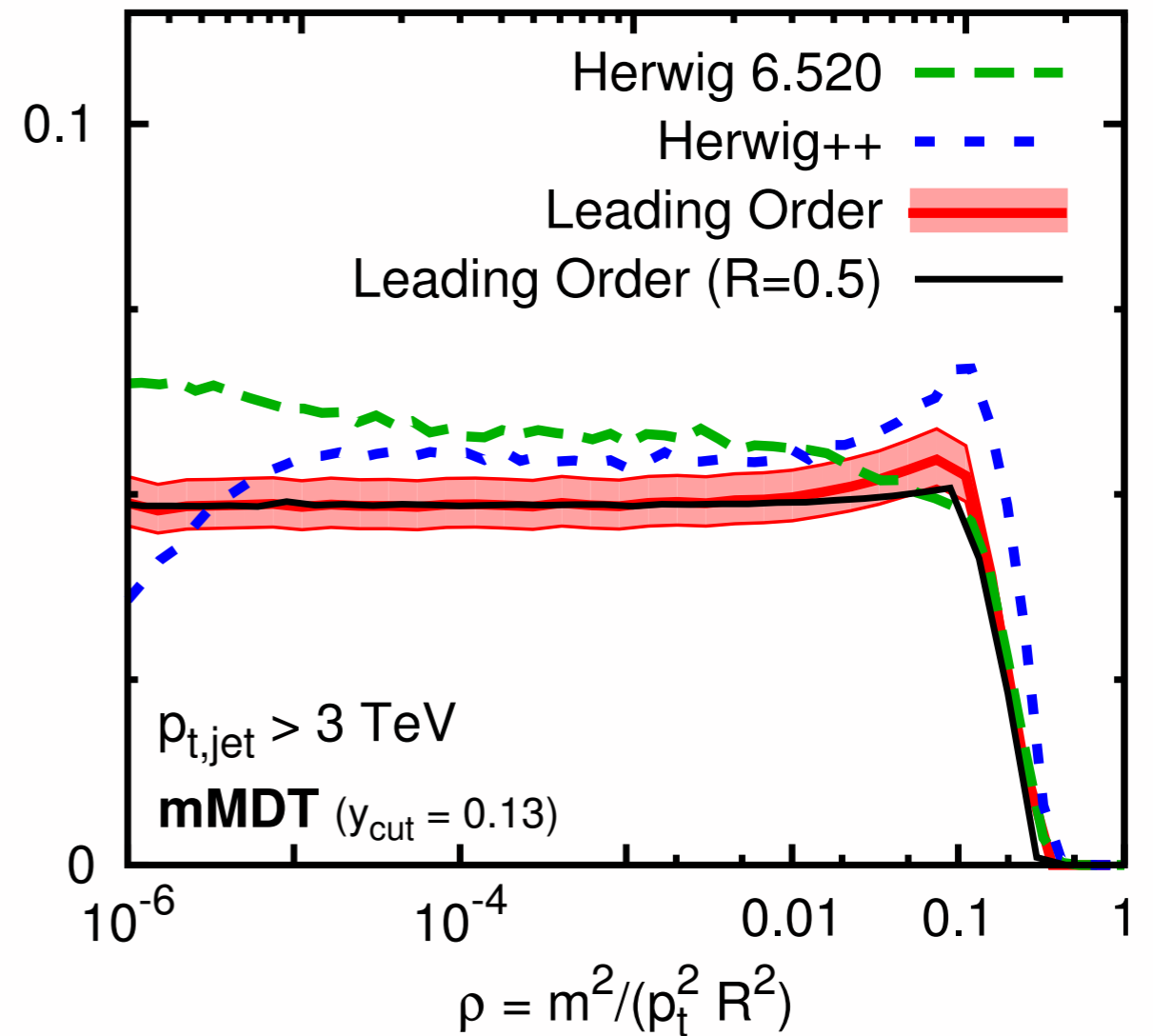
10 100 1000



LO v. Herwig showers (quark jets)

m [GeV], for $p_t = 3$ TeV, $R = 1$

10 100 1000

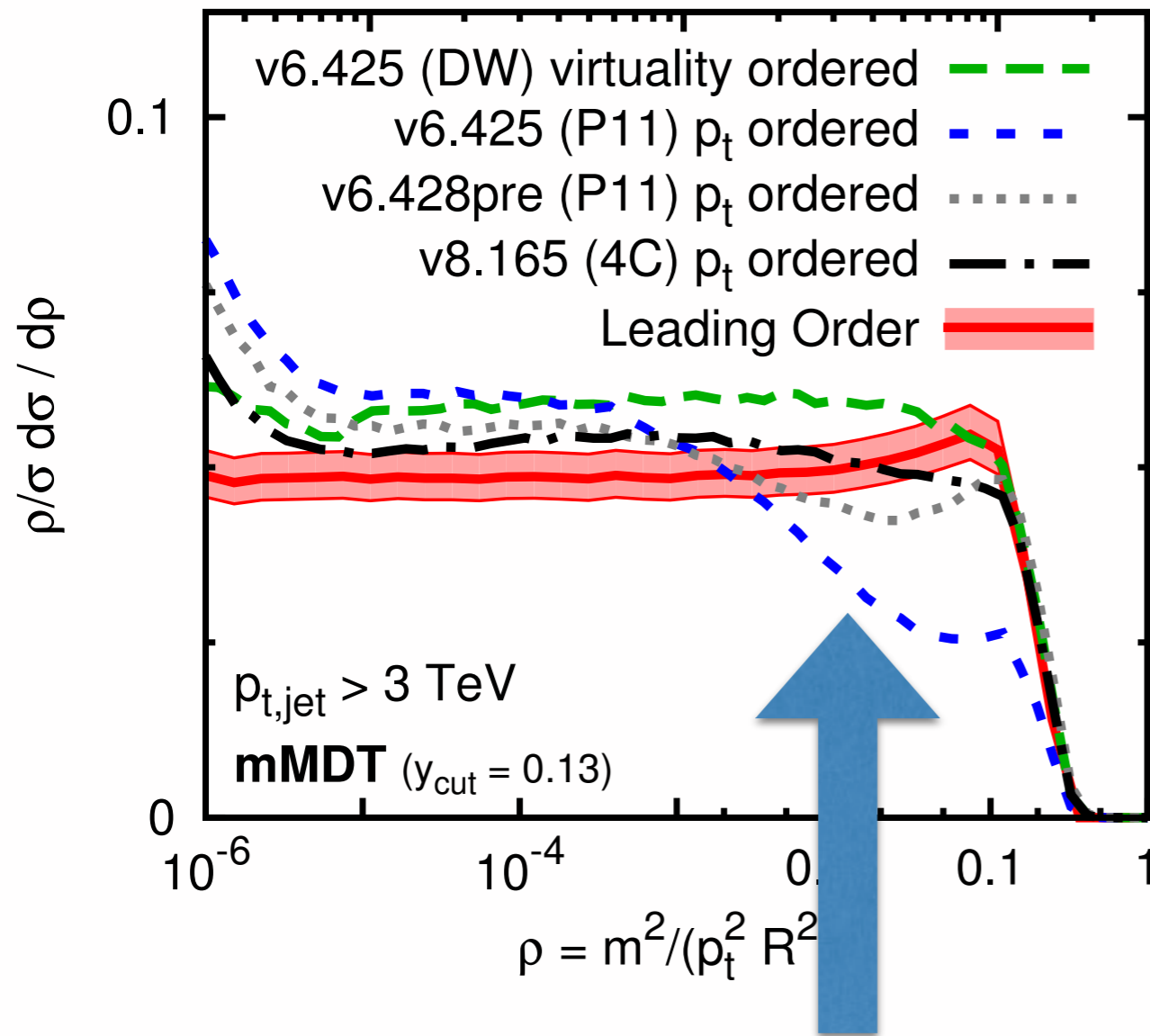


mMDT: comparing many showers

LO v. Pythia showers (quark jets)

m [GeV], for $p_t = 3$ TeV, $R = 1$

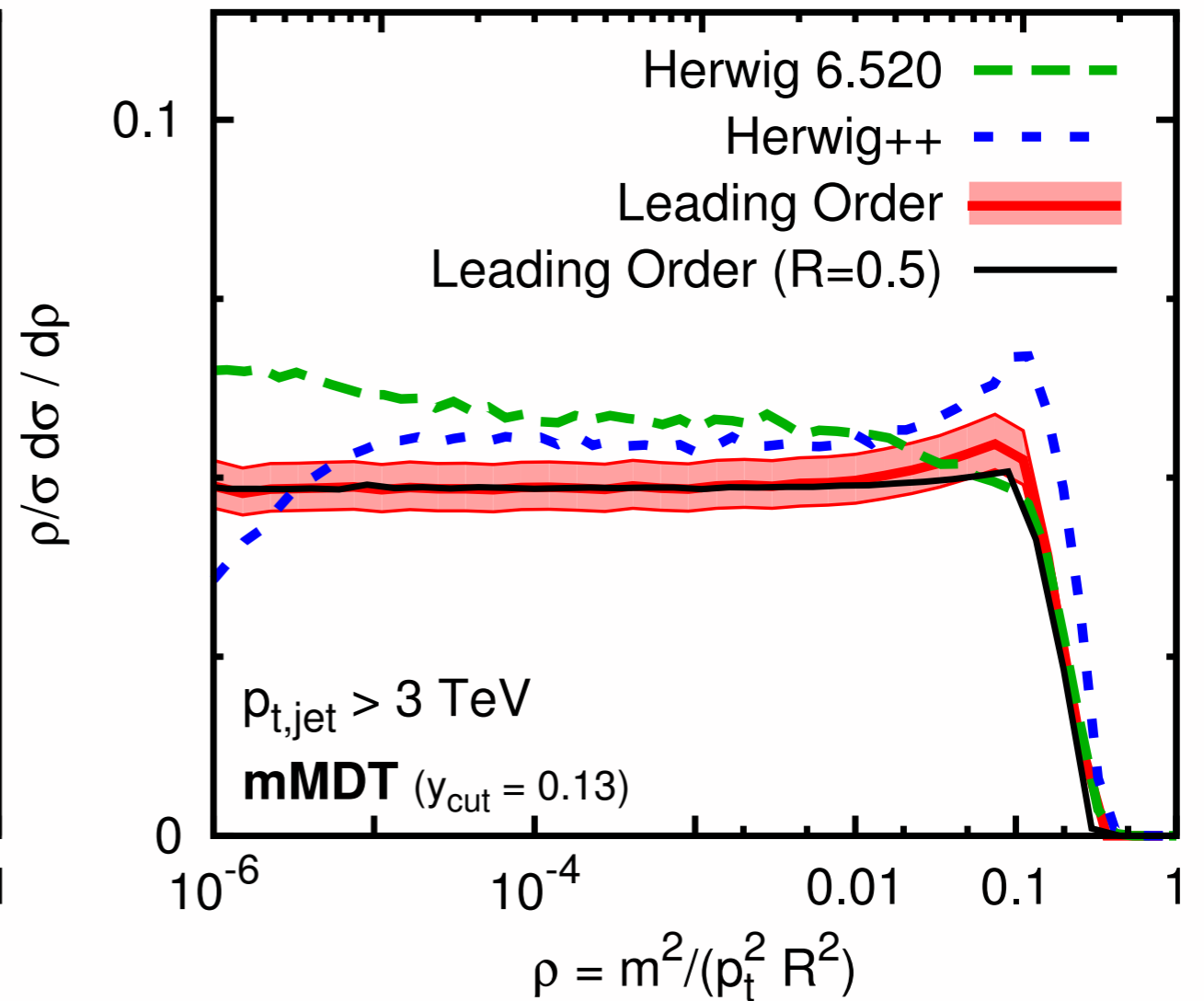
10 100 1000



LO v. Herwig showers (quark jets)

m [GeV], for $p_t = 3$ TeV, $R = 1$

10 100 1000



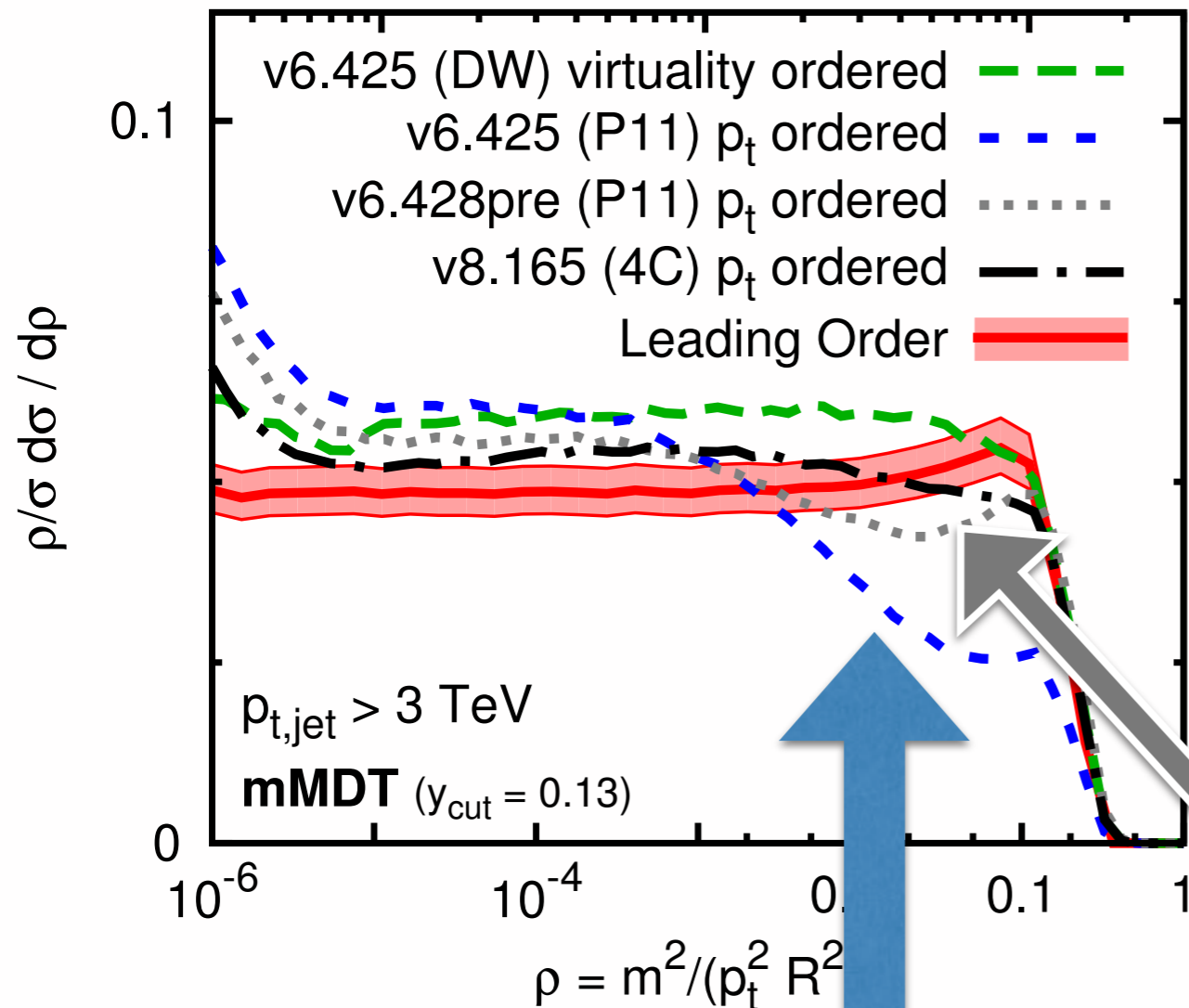
Issue found in Pythia 6 p_t -ordered shower → promptly identified and fixed by Pythia authors!

mMDT: comparing many showers

LO v. Pythia showers (quark jets)

m [GeV], for $p_t = 3$ TeV, $R = 1$

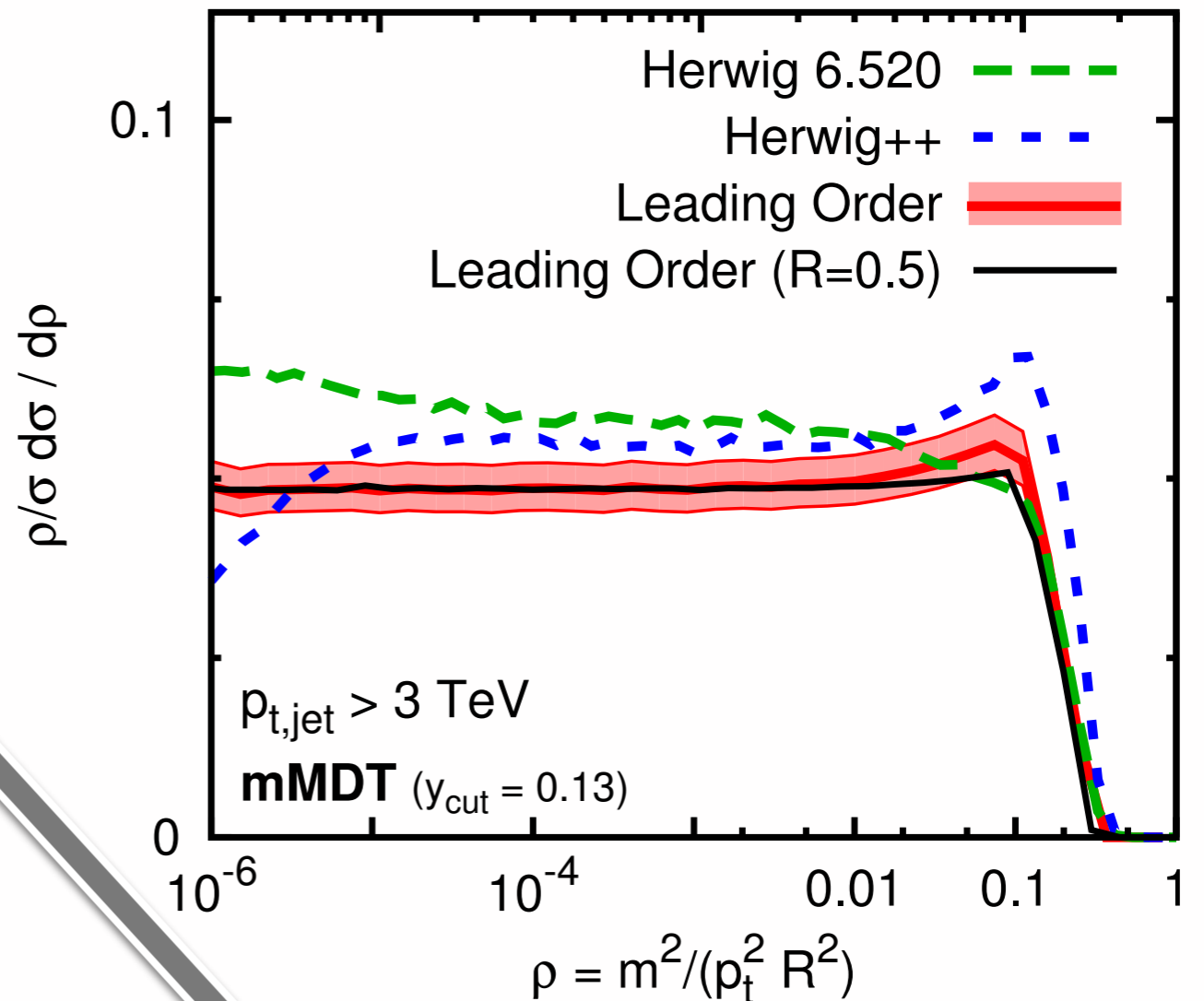
10 100 1000



LO v. Herwig showers (quark jets)

m [GeV], for $p_t = 3$ TeV, $R = 1$

10 100 1000



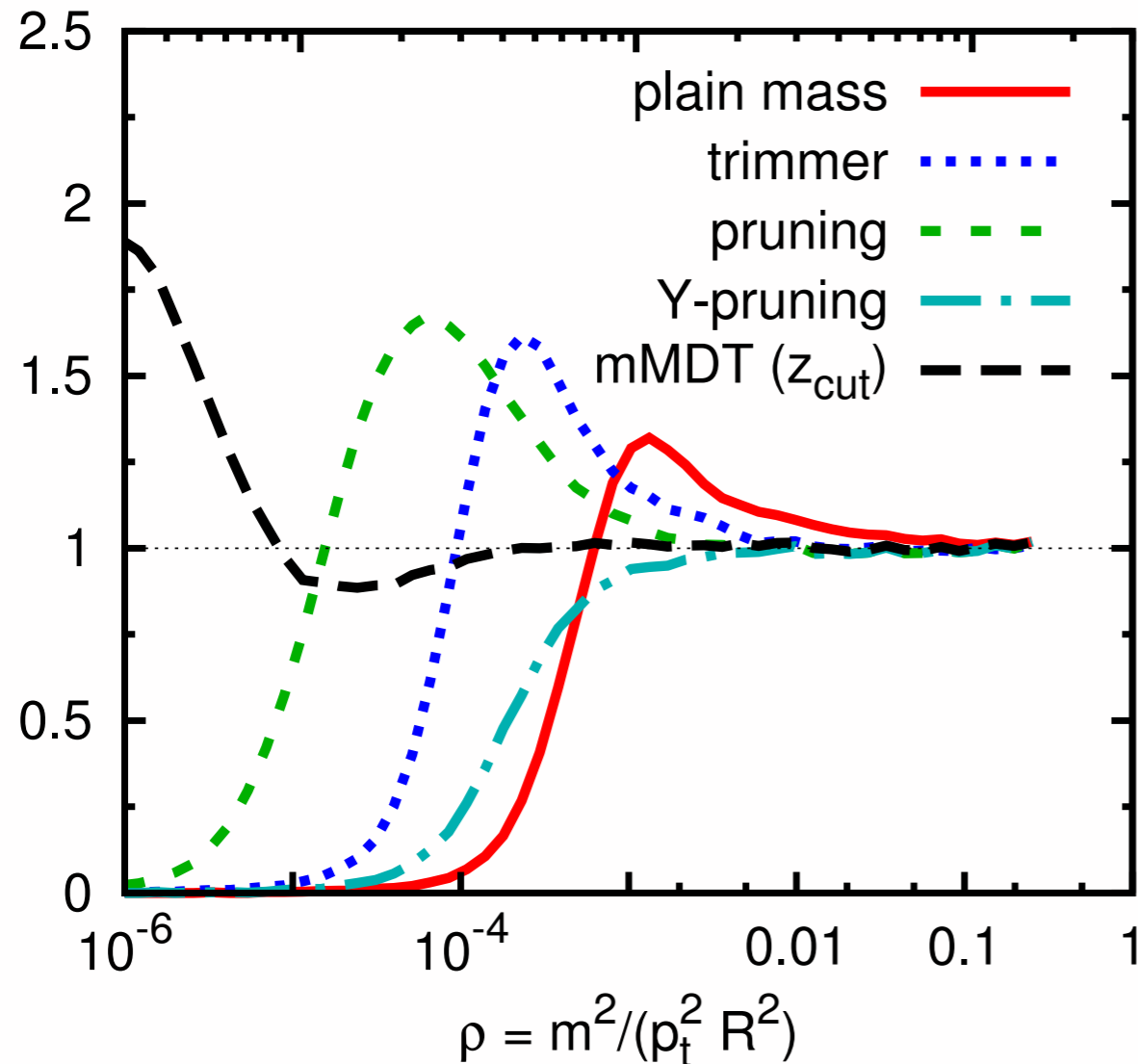
Issue found in Pythia 6 p_t -ordered shower → promptly identified and fixed by Pythia authors!

Hadronisation effects

hadronisation summary (quark jets)

m [GeV], for $p_t = 3$ TeV, $R = 1$

10 100 1000



Nearly all taggers have
large hadronisation
effects:

15 – 60%

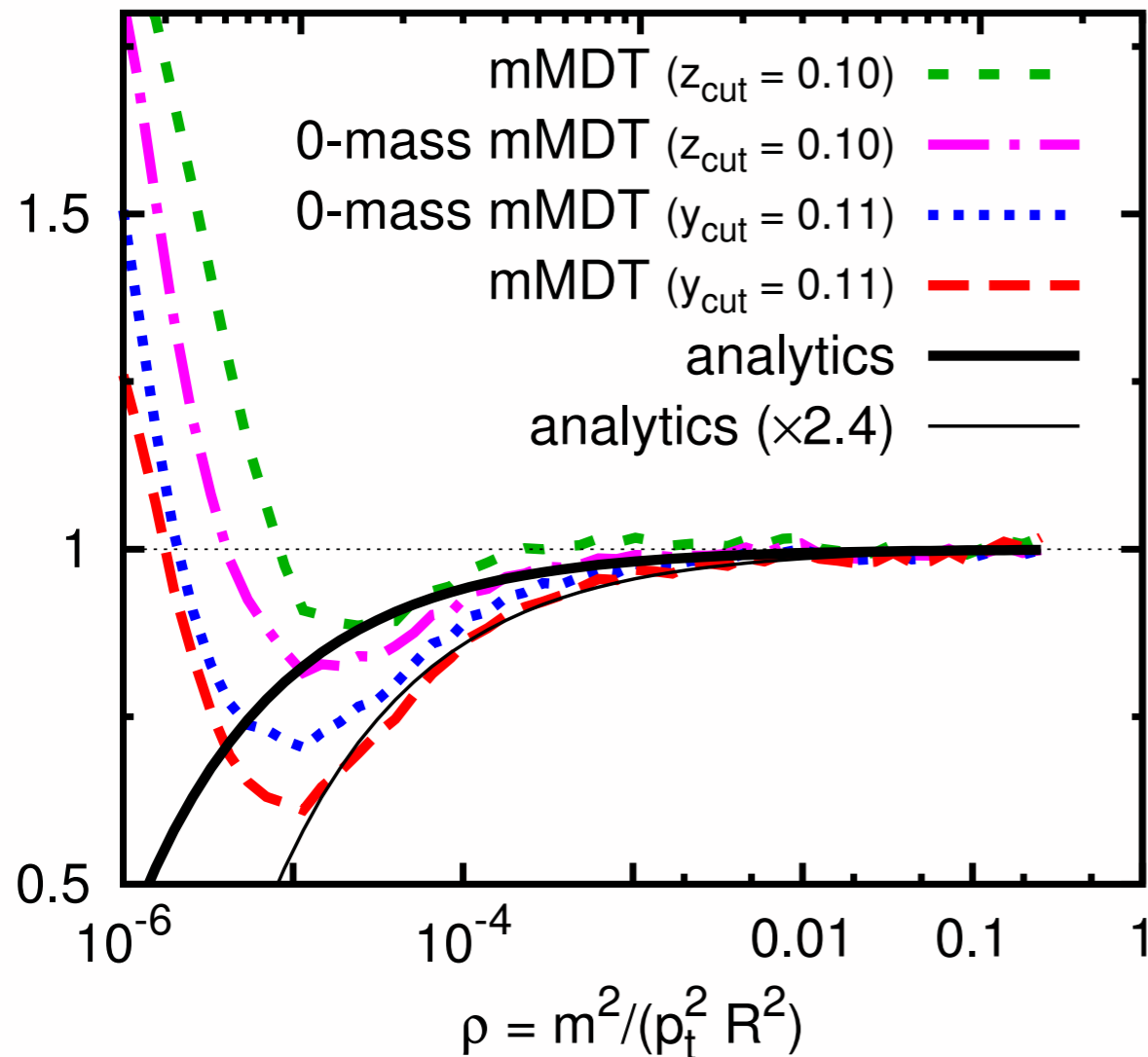
for $m = 30 - 100$ GeV

Hadronisation effects

hadronisation v. analytics (quark jets)

m [GeV], for $p_t = 3$ TeV, $R = 1$

10 100 1000



Exception is (m)MDT.

In some cases
just few % effect.

m -dependence of
hadronisation even
understood analytically!

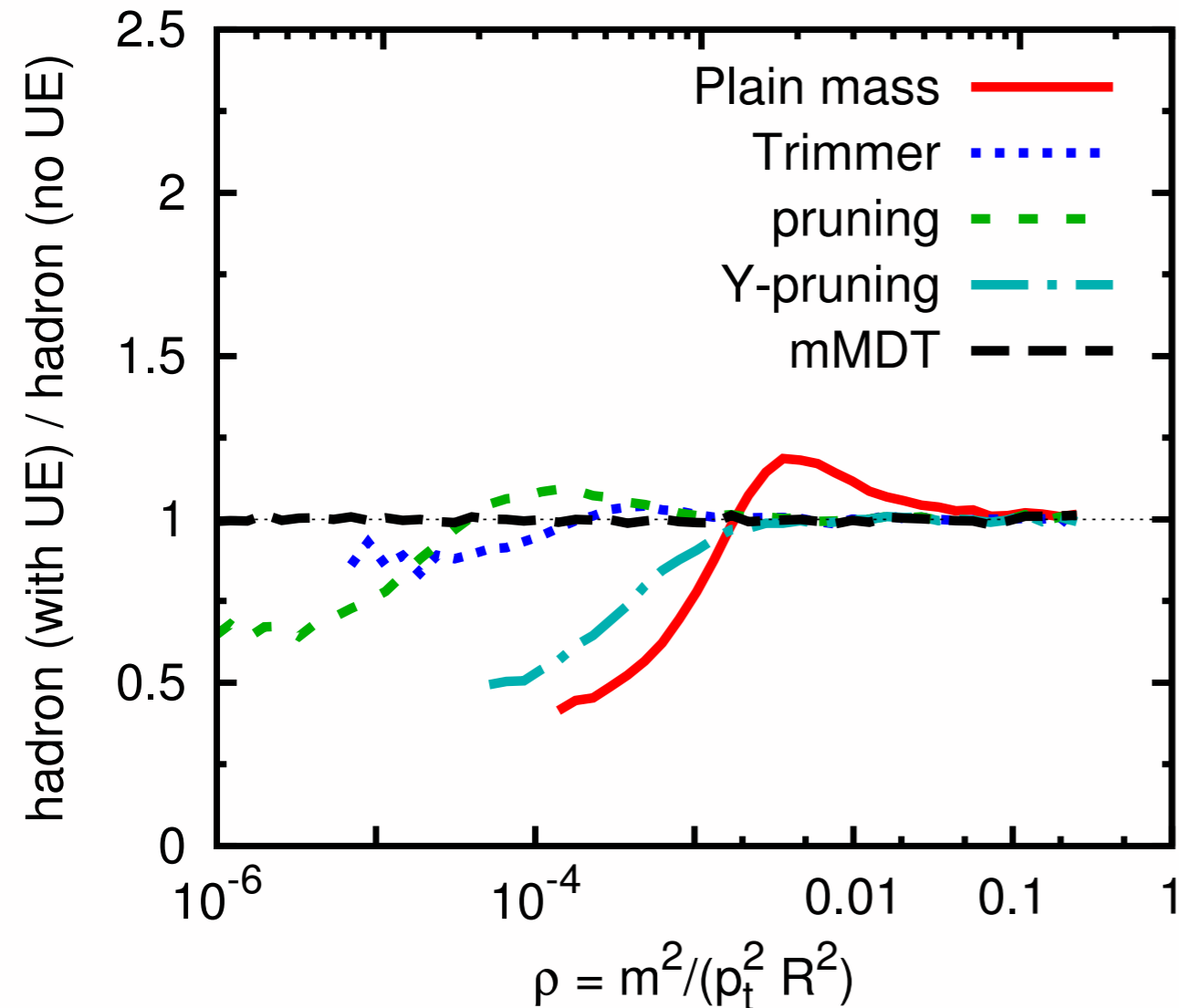
$$\frac{d\sigma}{dm}^{\text{hadron}} \simeq \frac{d\sigma}{dm}^{\text{parton}} \left(1 - c \frac{\Lambda}{m} \right)$$

Underlying Event (UE)

UE summary (quark jets)

m [GeV], for $p_t = 3$ TeV, $R = 1$

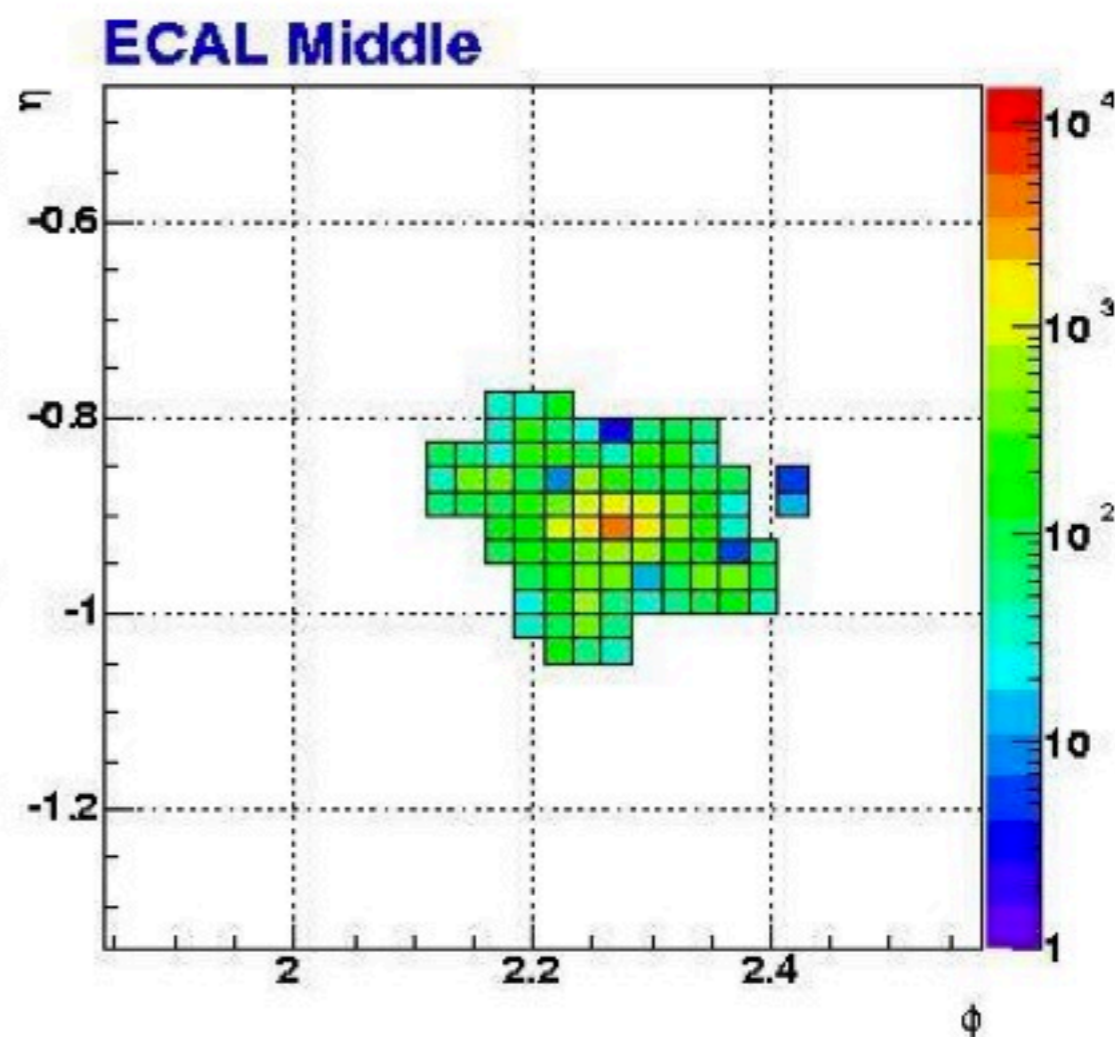
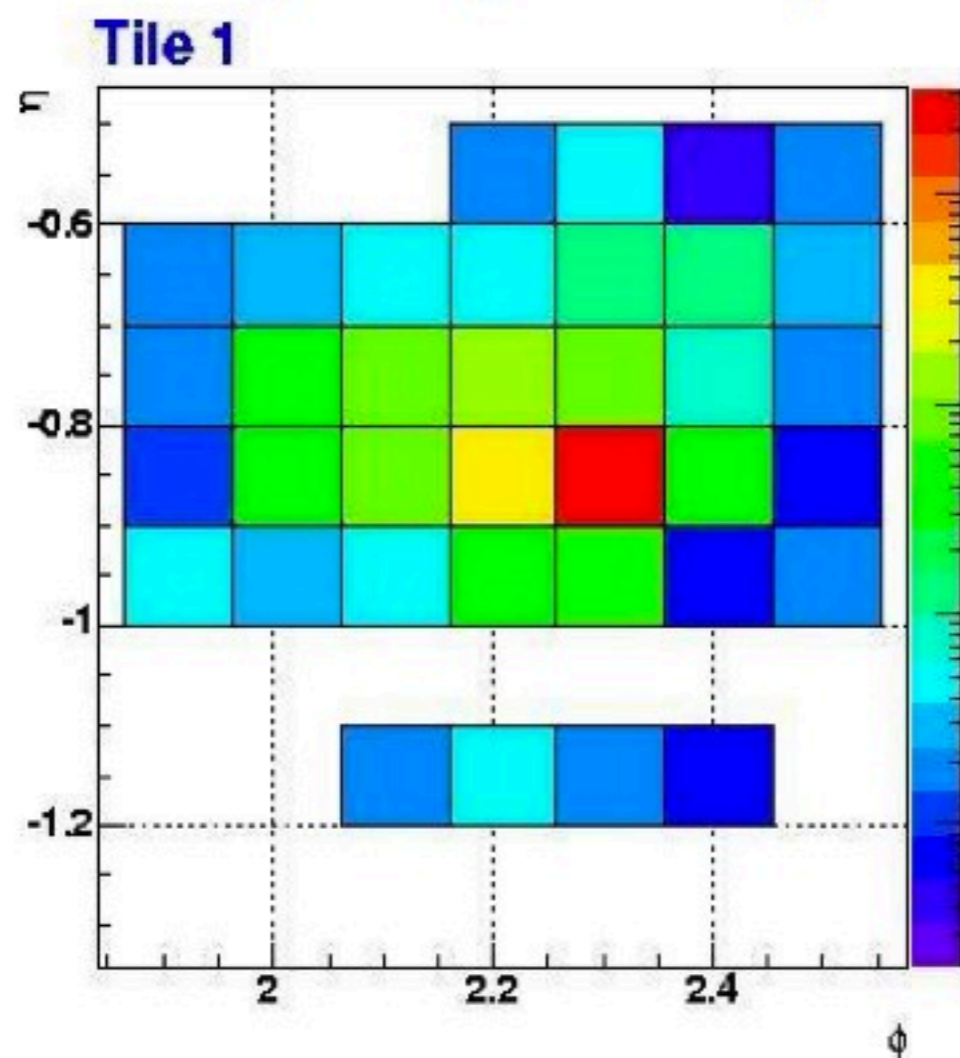
10 100 1000



Underlying event impact
much reduced relative to
jet mass

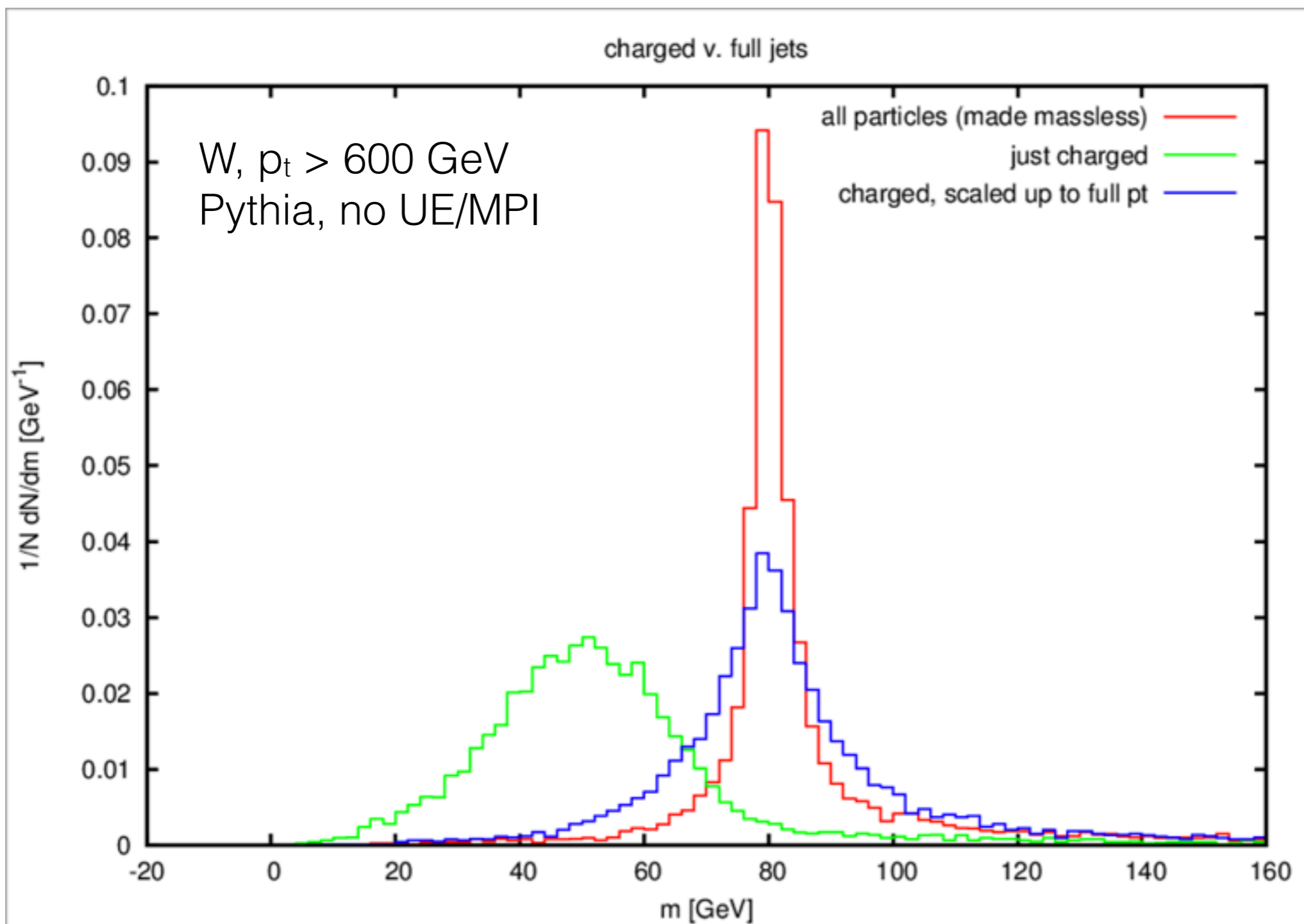
Almost zero for mMDT
(this depends on jet p_t)

A question from theorists to experimenters:
How well can you work around
detector granularity at high p_t ?



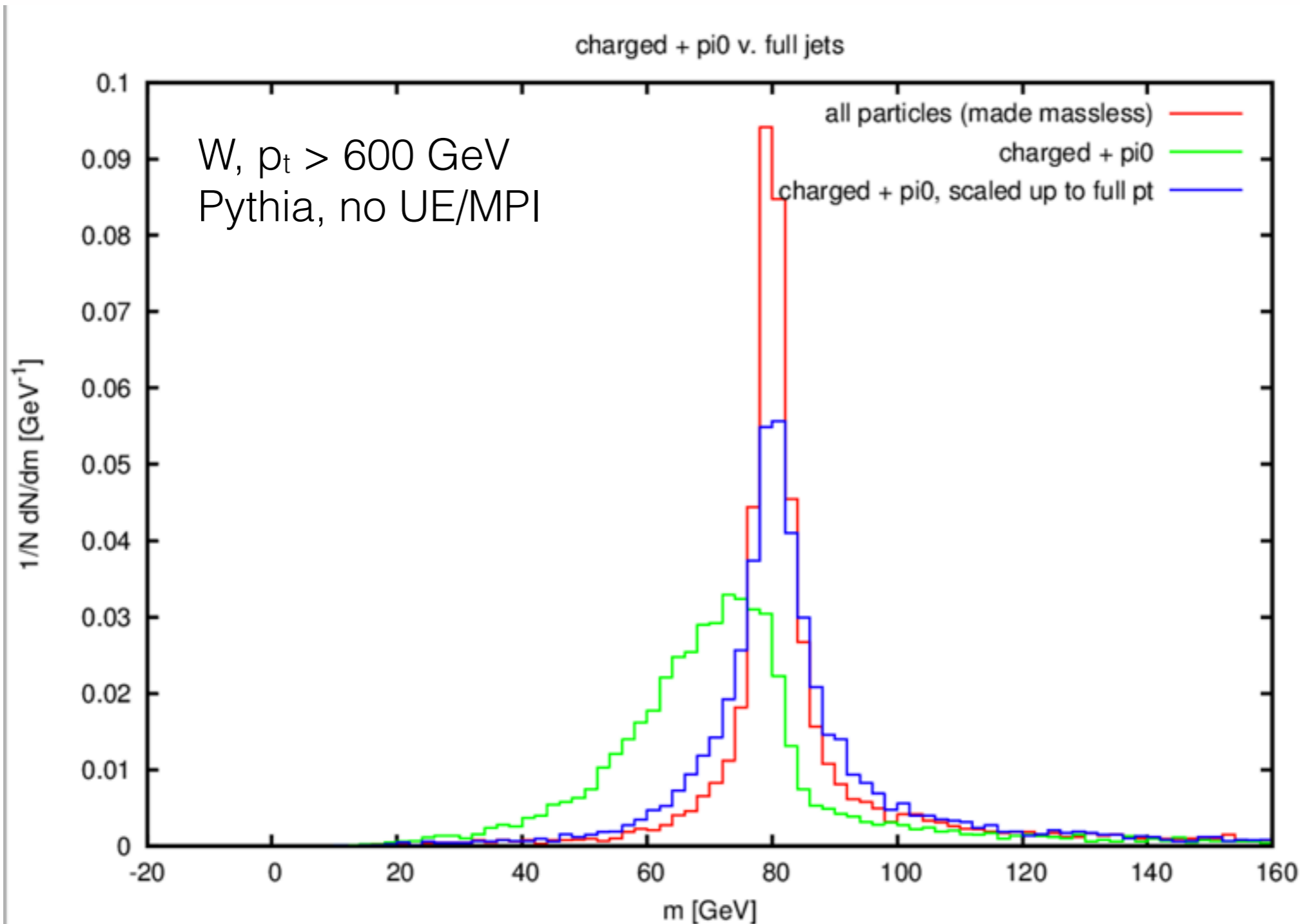
Detector resolution? Use Topocluster/PFlow directly.
Or charged tracks scaled to total jet energy?

cf. work by Tweedie et al



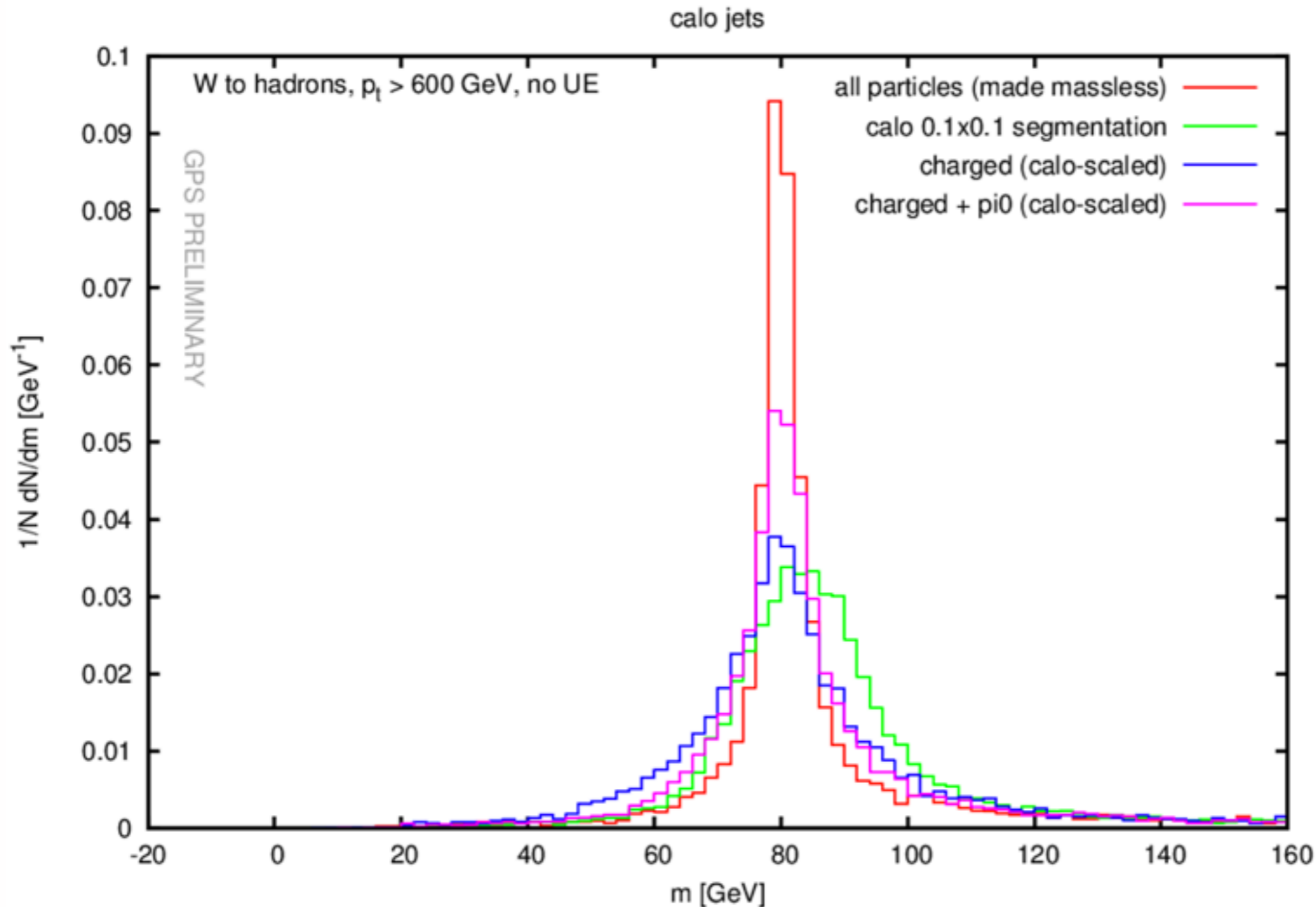
Detector resolution? Use Topocluster/PFlow directly.
Or charged tracks + π^0 scaled to total jet energy?

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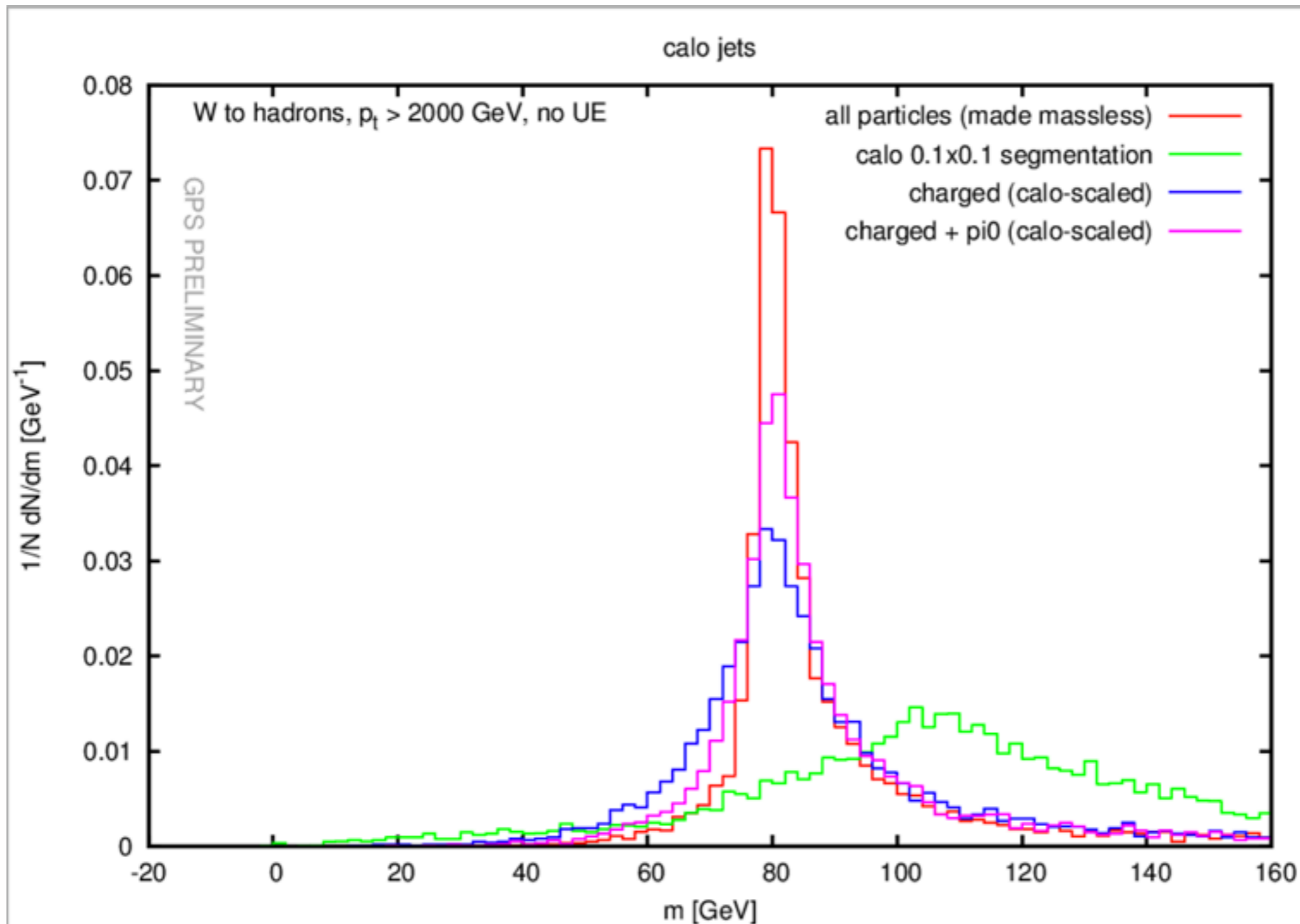
cf. work by Tweedie et al



illustrative
“Calorimeter”
only has
granularity, no
energy
fluctuations

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