

Perspectives on SM, Higgs and beyond at LHC

Gavin P. Salam (CERN)

*Theoretical Physics Colloquium
University of Oxford
17 February 2017*

**A typical introduction to a particle physics colloquium often starts
with “big unanswered questions”**

Nature of dark matter (& dark energy)

Fine-tuning (e.g. supersymmetry and similar)

Flavour-asymmetry of the universe

[...]

and less about the standard model (SM)...

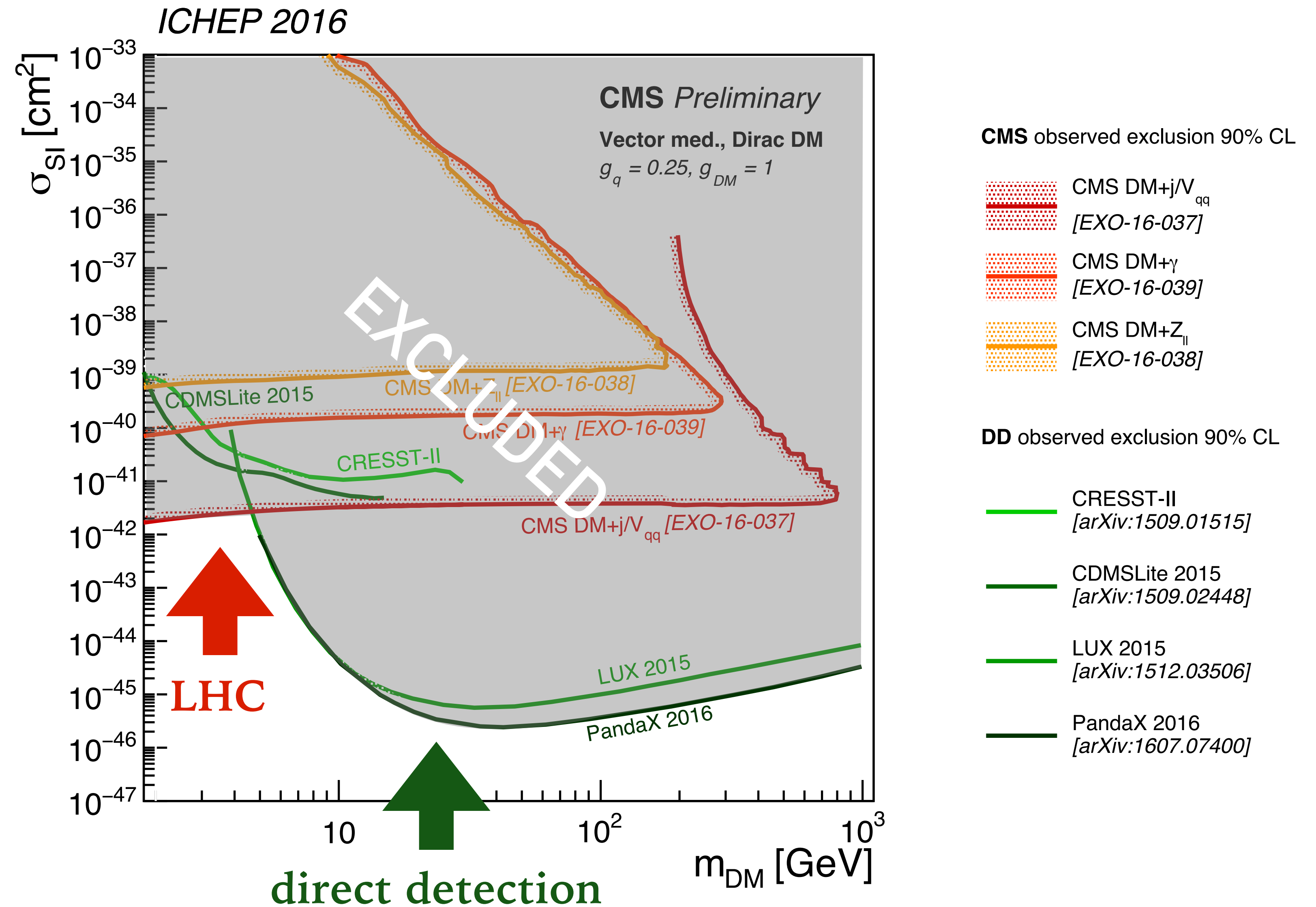
	mass →	$\approx 2.3 \text{ MeV}/c^2$	$\approx 1.275 \text{ GeV}/c^2$	$\approx 173.07 \text{ GeV}/c^2$	0	$\approx 125 \text{ GeV}/c^2$
QUARKS	charge →	$2/3$	$2/3$	$2/3$	0	0
	spin →	$1/2$	$1/2$	$1/2$	1	0
		u up	c charm	t top	g gluon	H Higgs boson
		$\approx 4.8 \text{ MeV}/c^2$	$\approx 95 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	
		$-1/3$	$-1/3$	$-1/3$	0	
		$1/2$	$1/2$	$1/2$	1	
		d down	s strange	b bottom	γ photon	
LEPTONS		$0.511 \text{ MeV}/c^2$	$105.7 \text{ MeV}/c^2$	$1.777 \text{ GeV}/c^2$	$91.2 \text{ GeV}/c^2$	
		-1	-1	-1	0	
		$1/2$	$1/2$	$1/2$	1	
		e electron	μ muon	τ tau	Z Z boson	
	$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 15.5 \text{ MeV}/c^2$	$80.4 \text{ GeV}/c^2$		
	0	0	0	± 1		
	$1/2$	$1/2$	$1/2$	1		
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson		
					GAUGE BOSONS	

since experiments have already found all its particles...

Looking beyond the SM: searches for dark matter at LHC & elsewhere

Classic dark-matter candidate: a weakly-interacting massive particle (WIMP, e.g. from supersymmetry).

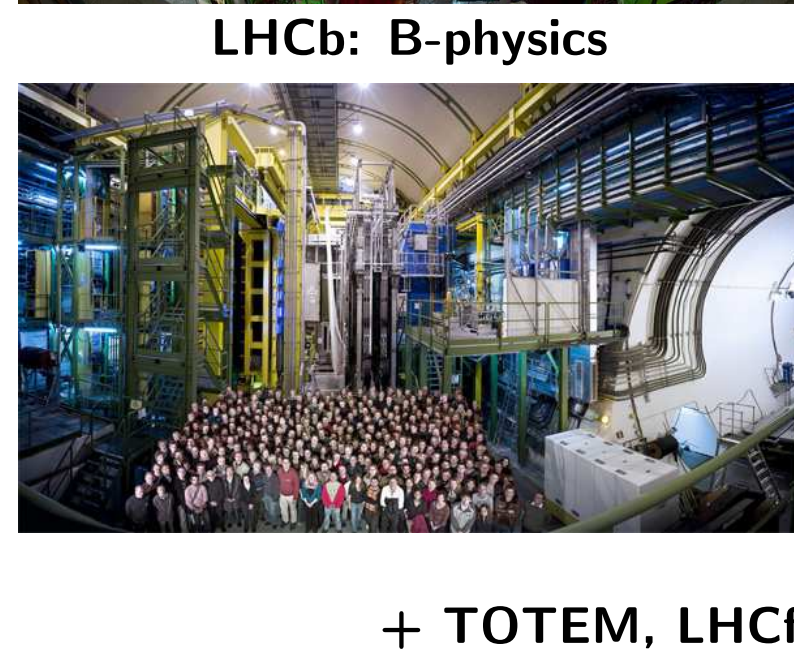
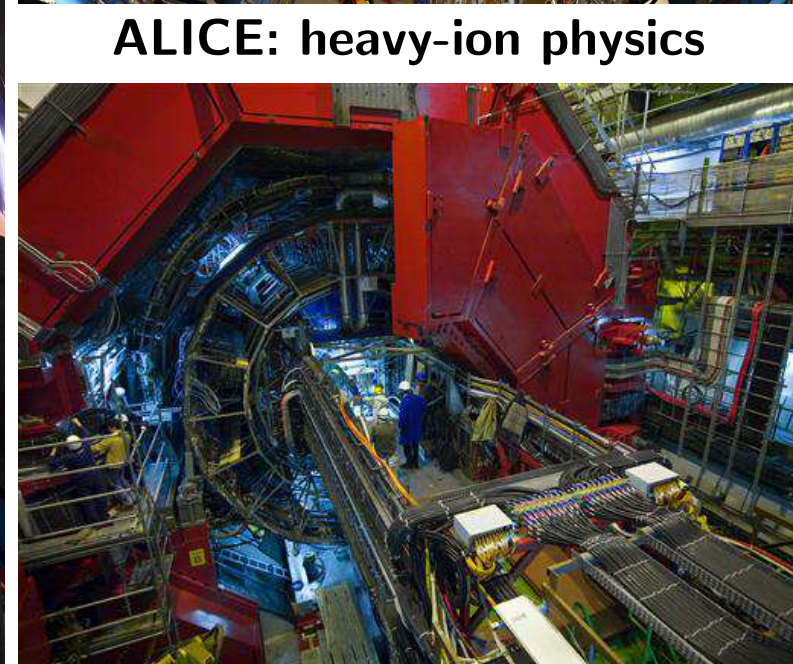
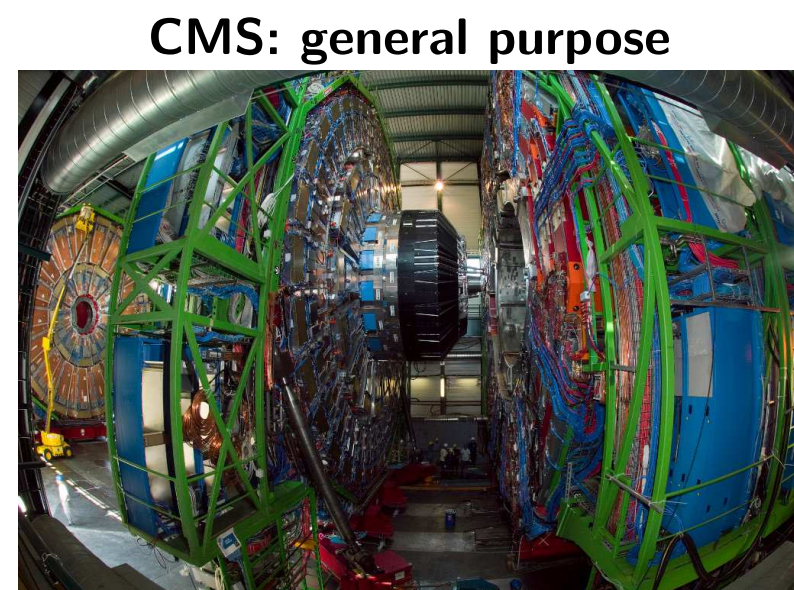
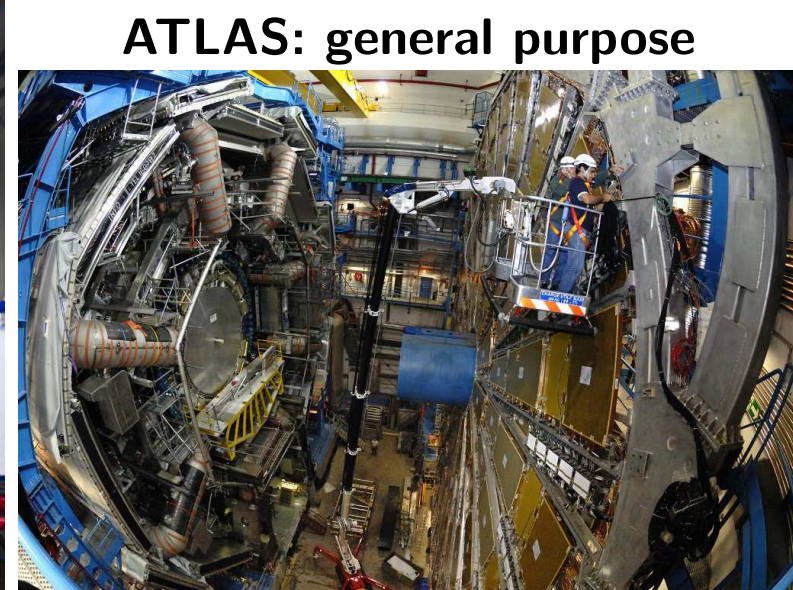
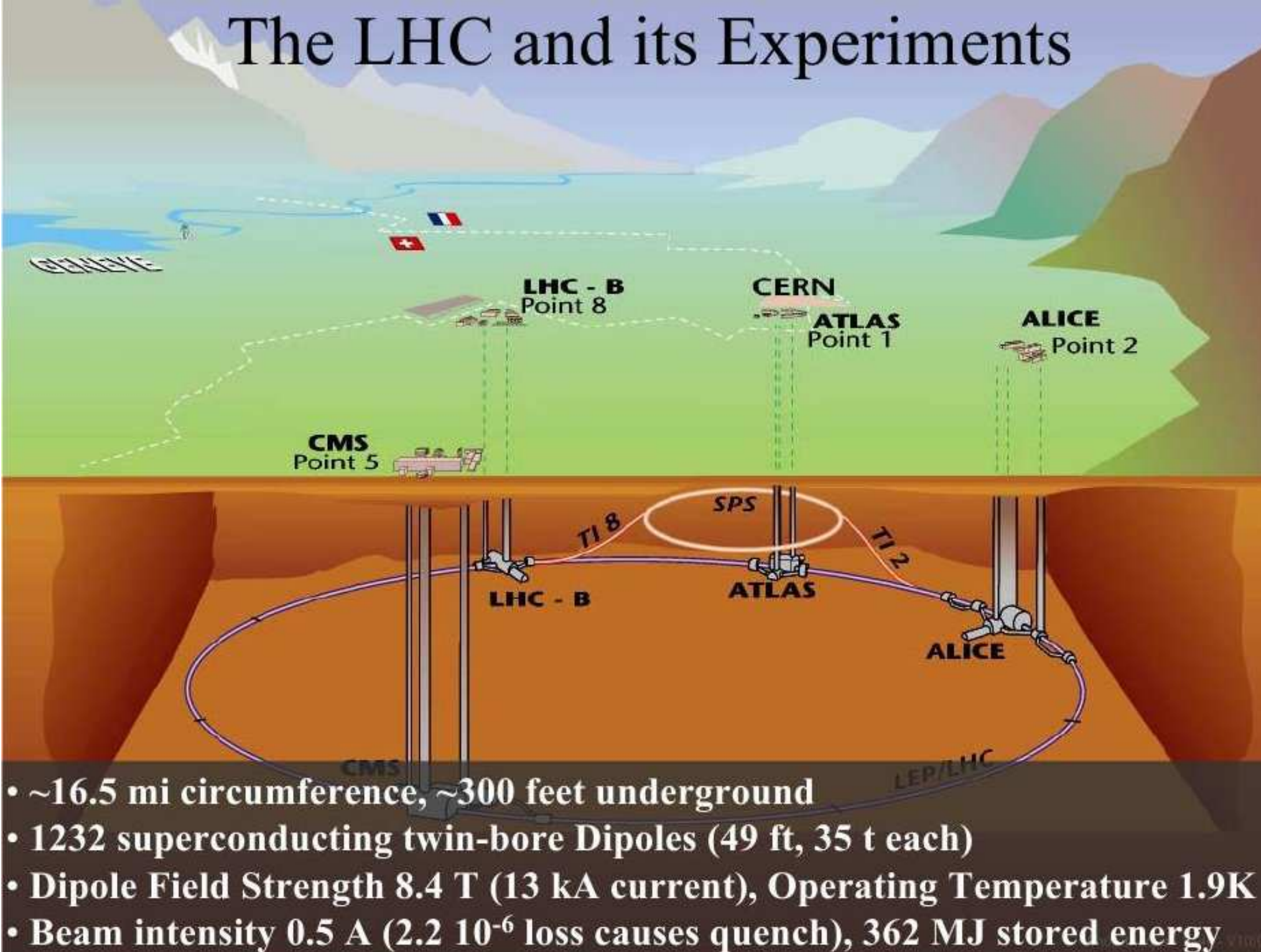
Masses \sim GeV upwards
(search interpretations strongly model dependent)



Searching for answers to the
“big unanswered questions” is vitally important,
(even if there’s no way of knowing if it will pay off)

But we also shouldn’t forget the importance of
“big answerable questions”
and the issue of how we go about answering them

The LHC and its Experiments



+ TOTEM, LHCf

perspective in context of LHC

Today

- 20 fb⁻¹ @ 8 TeV
- 13 fb⁻¹ @ 13 TeV (analysed)

Future

- 2018: 100 fb⁻¹ @ 13 TeV
- 2023: 300 fb⁻¹ @ 1? TeV
- 2035: 3000 fb⁻¹ @ 14 TeV

1 fb⁻¹ = 10¹⁴ collisions

Current analyses based on < 1% of the ultimate dataset

STANDARD MODEL — KNOWABLE UNKNOWNNS

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi \\ & + \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi)\end{aligned}$$

This is what you get when you buy one of those famous CERN T-shirts

This equation neatly sums up our current understanding of fundamental particles and forces.

STANDARD MODEL — KNOWABLE UNKNOWNNS

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi \\ & + \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi) \end{aligned}$$

This is what you get when you buy one of those famous CERN T-shirts

“understanding” = knowledge ?

“understanding” = assumption ?

This equation neatly sums up our **current understanding** of fundamental particles and forces.

NOTATION

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \not{D} \psi \\ & + \sum_i Y_{ij} \psi_i \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi)\end{aligned}$$

A_μ : gauge field

ψ : fermion field

ϕ : Higgs field

$$= \phi_0(\text{VEV}) + H(\text{Higgs})$$

$$D_\mu = \partial_\mu + ieA_\mu \text{ etc.}$$

$$F_{\mu\nu} \sim [D_\mu, D_\nu]$$

e.g. $\bar{\psi} D \psi \rightarrow \psi A_\mu \psi \rightarrow$ fermion-fermion-gauge vertex

i.e. terms of \mathcal{L} map to particle interactions

GAUGE-MATTER PART

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi \\ & + \psi_i Y_{ij} \psi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi)\end{aligned}$$

e.g. $qq\gamma$, qqZ , qqg , $e\nu W$ interactions
— well established in ep , e^+e^- , pp
collisions, etc.

≡ KNOWLEDGE

(also being studied at LHC — e.g.
jets, $DY/Z/W$, V +jets, $t\bar{t}$, etc.)

This equation neatly sums up our current understanding of fundamental particles and forces.

GAUGE-MATTER PART

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi \\ & + \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi)\end{aligned}$$

e.g. $qq\gamma$, qqZ , qqg , $e\nu W$ interactions
— well established in ep , e^+e^- , pp
collisions, etc.

≡ KNOWLEDGE

(also being studied at LHC — e.g.
jets, $DY/Z/W$, V +jets, $t\bar{t}$, etc.)

Do we “know” everything about this part?
E.g. direct emission of photon from top
quarks is, today, at edge of observability.
But it’s so much like any other gauge-matter
interaction that we almost take it for
granted

This equation neatly sums up our
current understanding of fundamental
particles and forces.

PURE GAUGE

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} + |D_\mu \phi|^2 - V(\phi)$$

e.g. ZWW, 3-gluon interactions — well established at LEP (e.g. $e^+e^- \rightarrow W^+W^-$)

≡ KNOWLEDGE

& also being studied at LHC

This equation neatly sums up our current understanding of fundamental particles and forces.

PURE GAUGE

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} + |D_\mu \phi|^2 - V(\phi)$$

e.g. ZWW, 3-gluon interactions — well established at LEP (e.g. $e^+e^- \rightarrow W^+W^-$)

≡ KNOWLEDGE

& also being studied at LHC

**We've seen gauge sectors work over and over again
→ gives us the illusion that the SM is established**

This equation neatly sums up our current understanding of fundamental particles and forces.

HIGGS BOSON

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi \\ & + \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi) \end{aligned}$$

LEP precision made it compelling,
LHC discovered it

≡ KNOWLEDGE

it behaves in every way like a scalar

≡ KNOWLEDGE

is it fundamental/pointlike?

to find out need

~ high- p_T /offshell Higgses

→ data barely sensitive...

This equation neatly sums up our current understanding of fundamental particles and forces.

HIGGS BOSON

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} \not{D} \psi \\ & + \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} \\ & + |D_\mu \phi|^2 - V(\phi) \end{aligned}$$

LEP precision made it compelling,
LHC discovered it

≡ KNOWLEDGE

it behaves in every way like a scalar

≡ KNOWLEDGE

is it fundamental/pointlike?

to find out need

~ high- p_T /offshell Higgses

→ data barely sensitive...

**Novelty? If fundamental, very
(the only fundamental scalar we know of)**

This equation neatly sums up our current understanding of fundamental particles and forces.

GAUGE-HIGGS INTERACTIONS

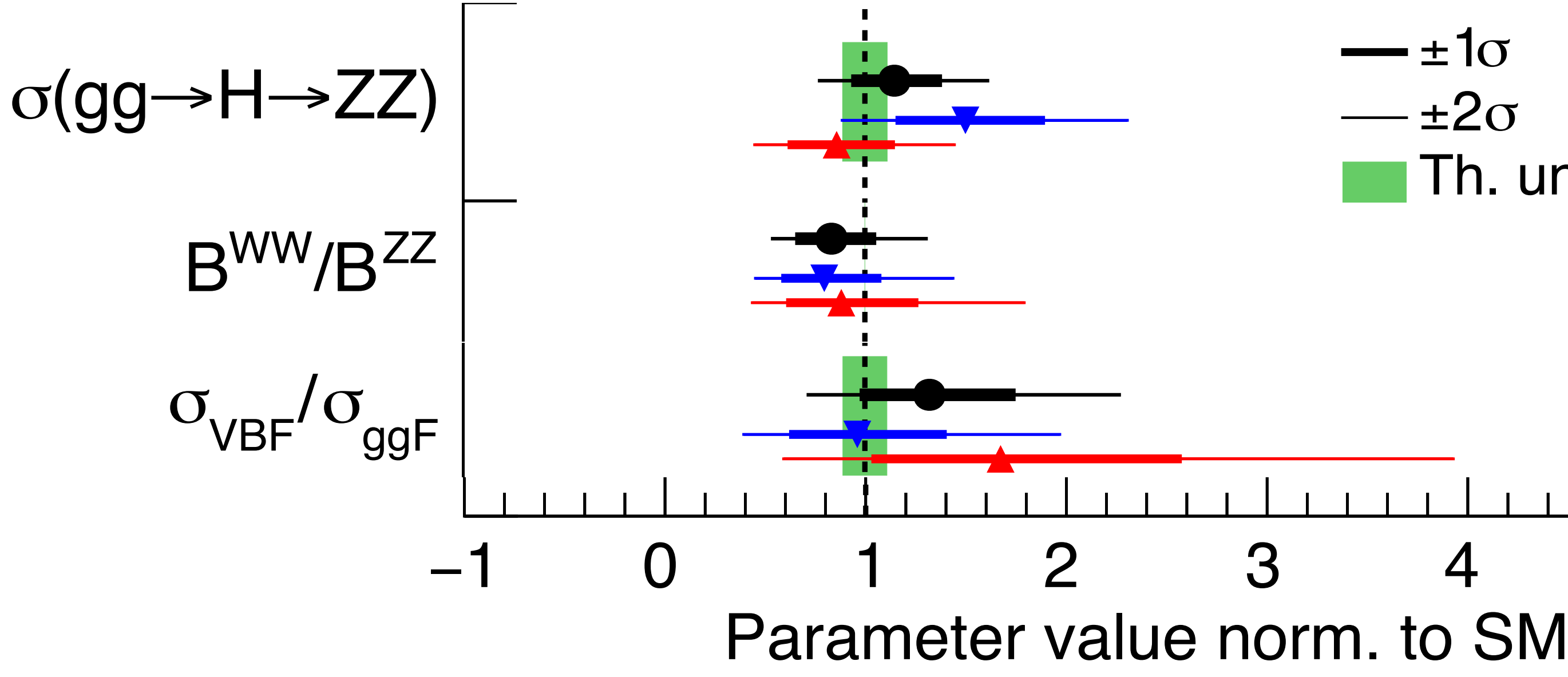
$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + \sum_i \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} + |D_\mu \phi|^2 - V(\phi)$$

Definitely non-zero.

$H \rightarrow ZZ, H \rightarrow WW, \text{VBF}$

(would require "conspiracy" of couplings in order to be substantially different)

≡ PROBABLY TRUE



This equation neatly sums up our current understanding of fundamental particles and forces.

GAUGE-HIGGS INTERACTIONS

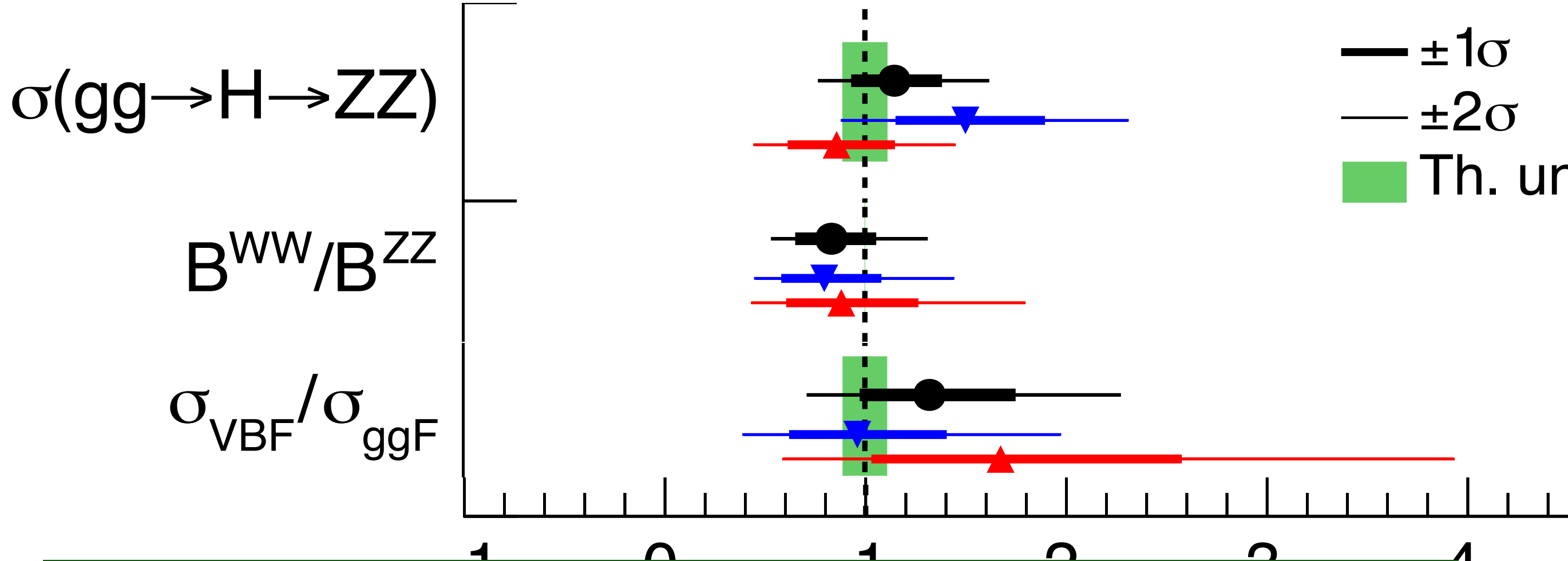
$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} + |D_\mu \phi|^2 - V(\phi)$$

Definitely non-zero.

$H \rightarrow ZZ, H \rightarrow WW, \text{VBF}$

(would require "conspiracy" of couplings in order to be substantially different)

≡ PROBABLY TRUE



Novelty? Covariant derivative D is widespread, but first time we see it with a scalar

This equation neatly sums up our current understanding of fundamental particles and forces.

YUKAWA COUPLINGS

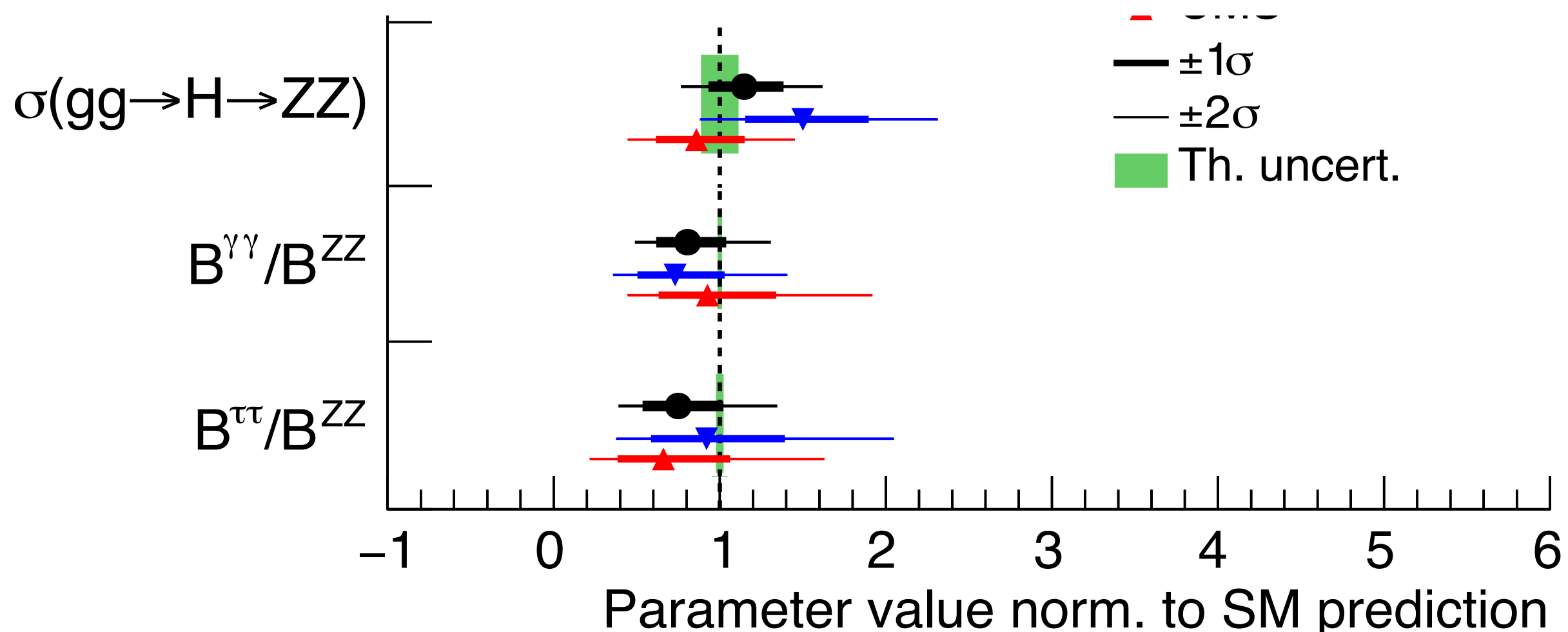
$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + \psi_i Y_{ij} \psi_j \phi + \text{h.c.} + |D_\mu \phi|^2 - V(\phi)$$

top? $gg \rightarrow H, H \rightarrow \gamma\gamma \equiv$ **INDIRECT**

bottom? H branching ratios \equiv **INDIRECT**

tau? \sim observed $\equiv \sim$ **KNOWLEDGE**

1st & 2nd gen? \equiv **IGNORANCE**



This equation neatly sums up our current understanding of fundamental particles and forces.

YUKAWA COUPLINGS

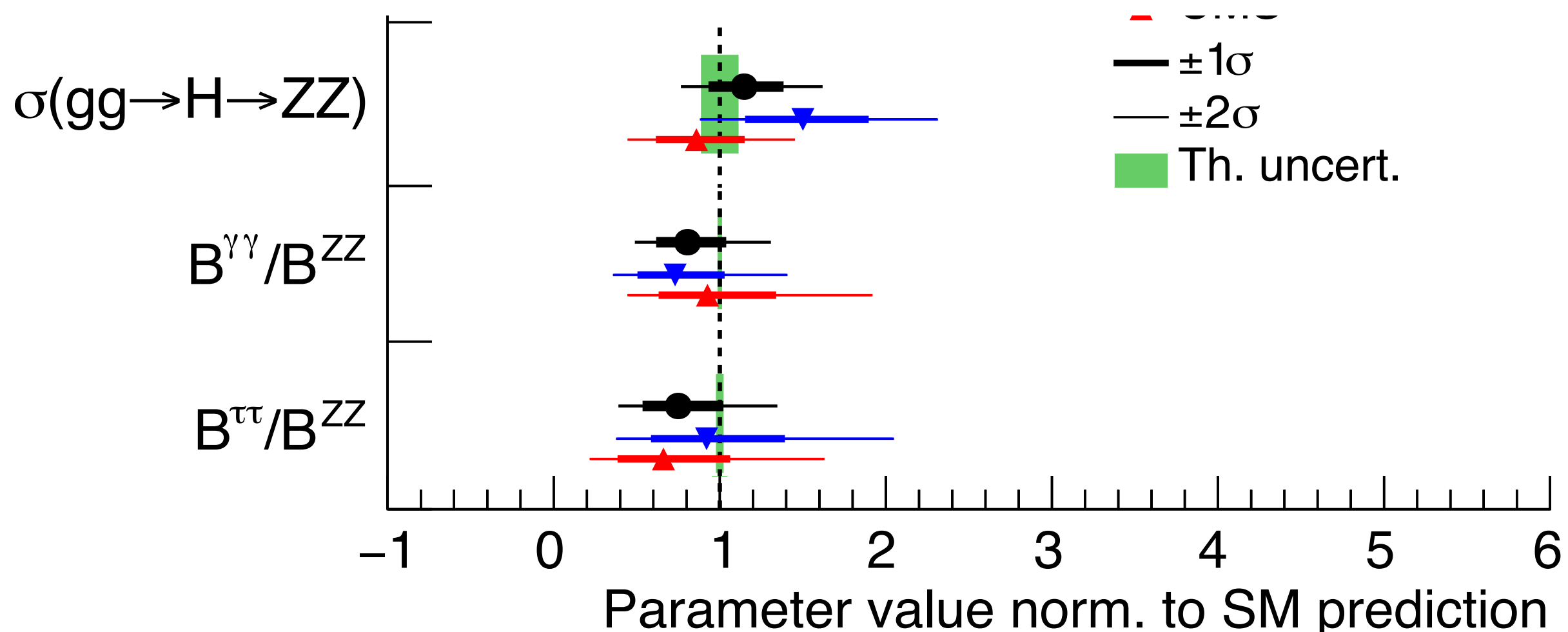
$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i \bar{\psi} \not{D} \psi + \sum_i \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} + |D_\mu \phi|^2 - V(\phi)$$

top? $gg \rightarrow H, H \rightarrow \gamma\gamma \equiv$ **INDIRECT**

bottom? H branching ratios \equiv **INDIRECT**

tau? \sim observed $\equiv \sim$ **KNOWLEDGE**

1st & 2nd gen? \equiv **IGNORANCE**



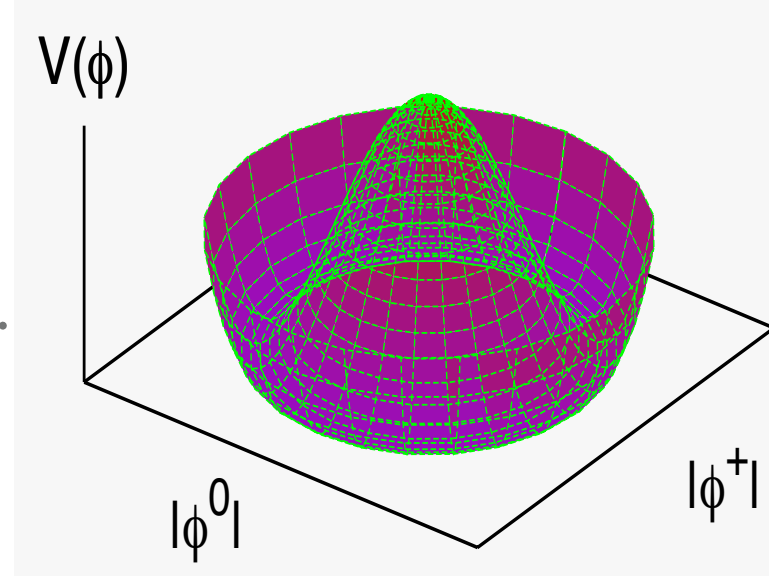
This equation neatly sums up our current understanding of fundamental particles and forces.

Novelty? We've never seen anything like it
 → **mystery of 5 orders of magnitude in mass between electron & top, CKM**

$$\begin{aligned}
\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\
& + i \bar{\psi} \not{D} \psi \\
& + \bar{\psi}_i Y_{ij} \psi_j \phi + h.c. \\
& + |D_\mu \phi|^2 - \boxed{V(\phi)}
\end{aligned}$$

This equation neatly sums up our current understanding of fundamental particles and forces.

HIGGS POTENTIAL



Vacuum Expectation Value (VEV)?

≡ KNOWLEDGE

2nd derivative ($\sim m_H$)?

[not a prediction of the theory
& any realistic theory must have a
minimum & 2nd derivative]

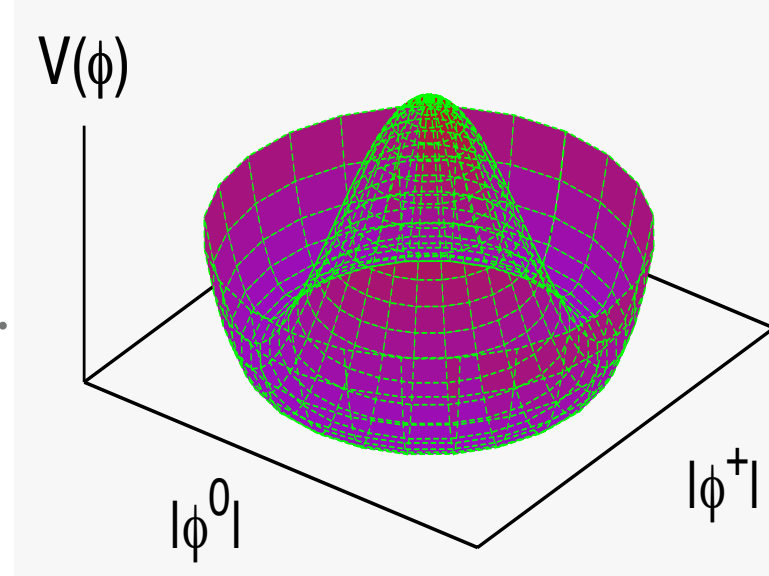
≡ KNOWLEDGE

$\phi^2 + \phi^4$? ≡ ASSUMPTION

$$\begin{aligned}
\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\
& + i \bar{\psi} \not{D} \psi \\
& + \chi_i Y_{ij} \chi_j \phi + h.c. \\
& + |D_\mu \phi|^2 - \boxed{V(\phi)}
\end{aligned}$$

This equation neatly sums up our current understanding of fundamental particles and forces.

HIGGS POTENTIAL



Vacuum Expectation Value (VEV)?

≡ KNOWLEDGE

2nd derivative ($\sim m_H$)?

[not a prediction of the theory
& any realistic theory must have a
minimum & 2nd derivative]

≡ KNOWLEDGE

$\phi^2 + \phi^4$? ≡ ASSUMPTION

**Novelty? Theorists' toy model,
never seen in nature (as fundamental);
Connects with stability of universe**

OVERALL TODAY

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + \bar{\psi}_i Y_{ij} \psi_j \phi + \text{h.c.} + |D_\mu \phi|^2 - V(\phi)$$




There remains a lot to establish
in the Higgs sector

This equation neatly sums up our current understanding of fundamental particles and forces.

OVERALL TODAY

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + \psi_i y_{ij} \psi_j \phi + \text{h.c.} + |D_\mu \phi|^2 - V(\phi)$$



t	b	τ
c	s	μ
u	s	e



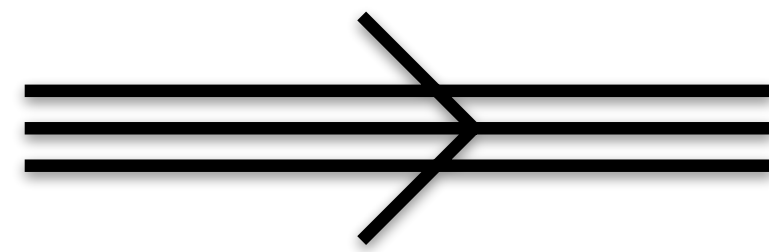
ASSUMPTION

KNOWLEDGE

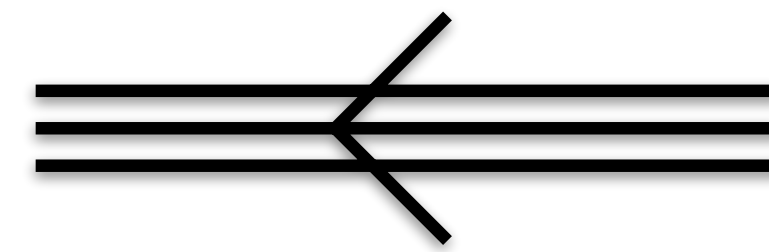
There remains a lot to establish in the Higgs sector

This equation neatly sums up our current understanding of fundamental particles and forces.

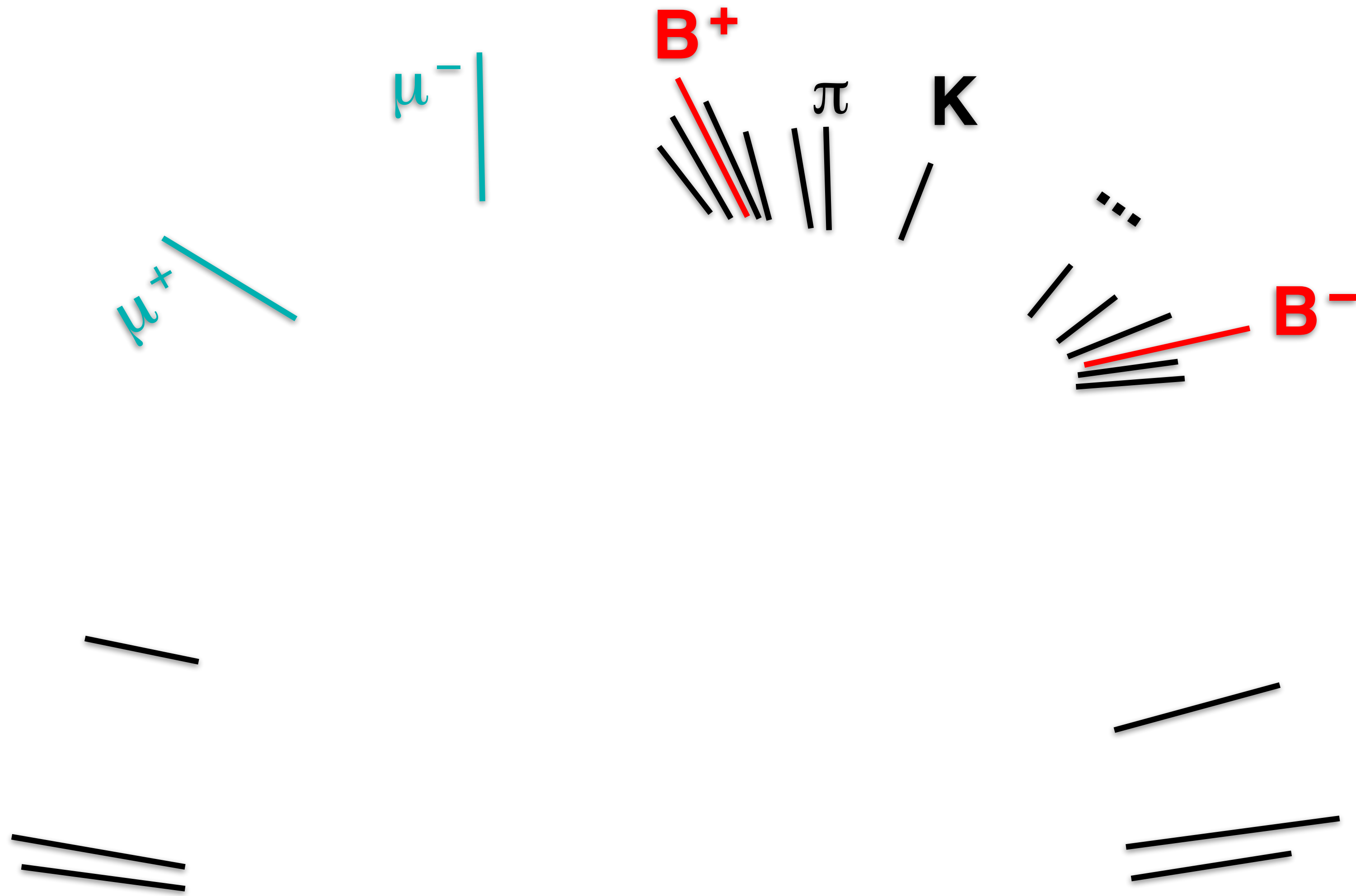
HOW?



proton

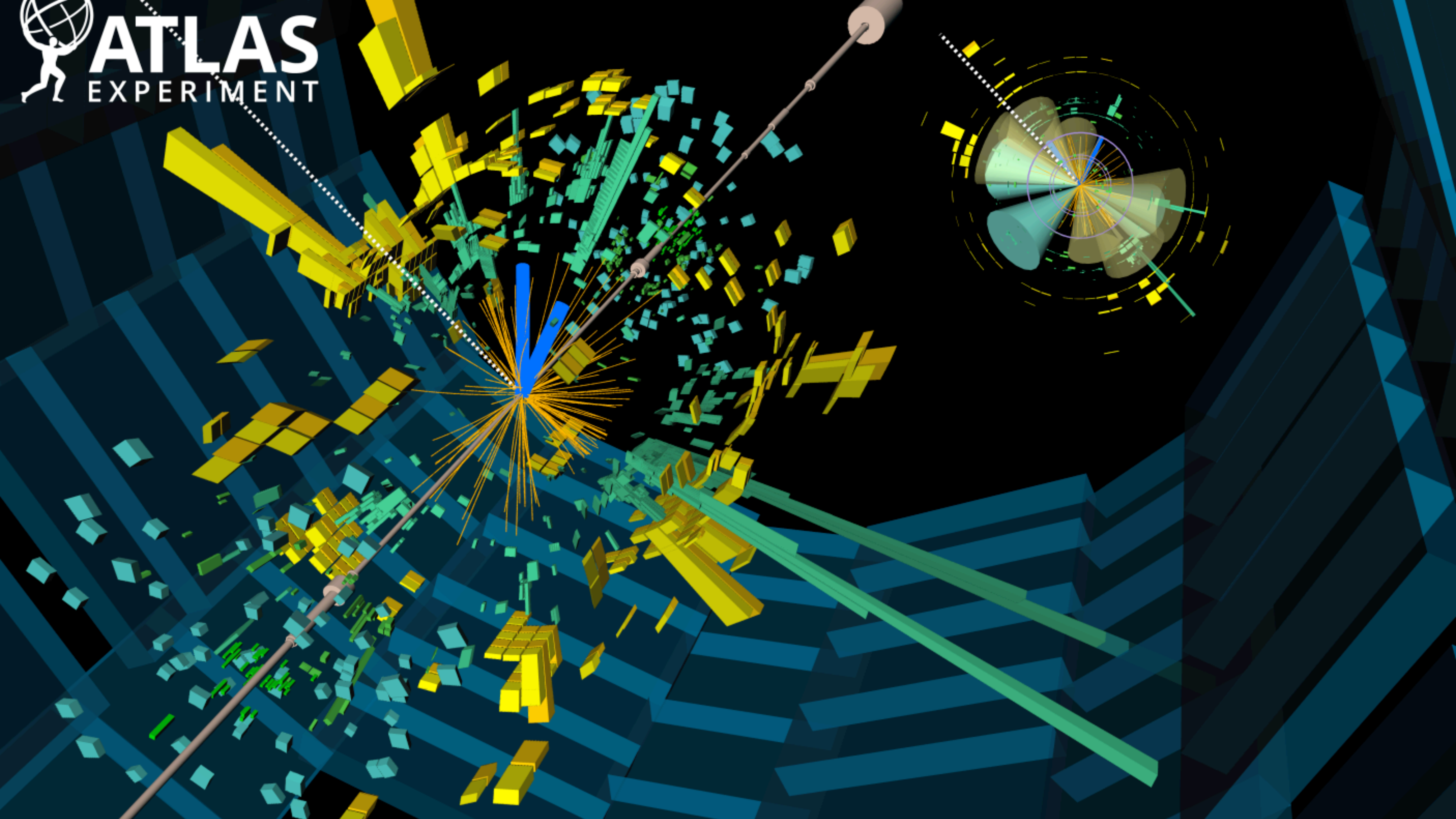


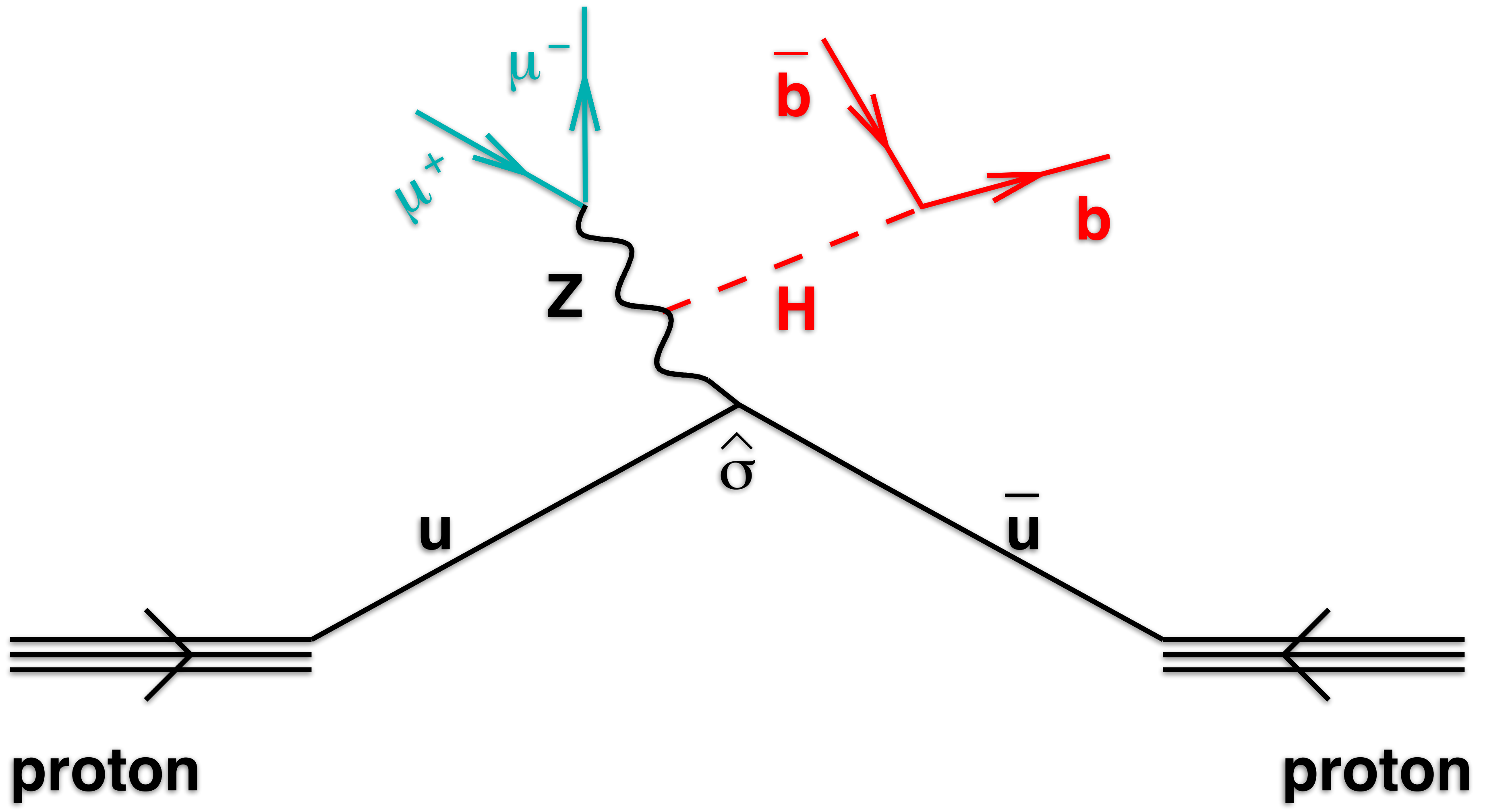
proton

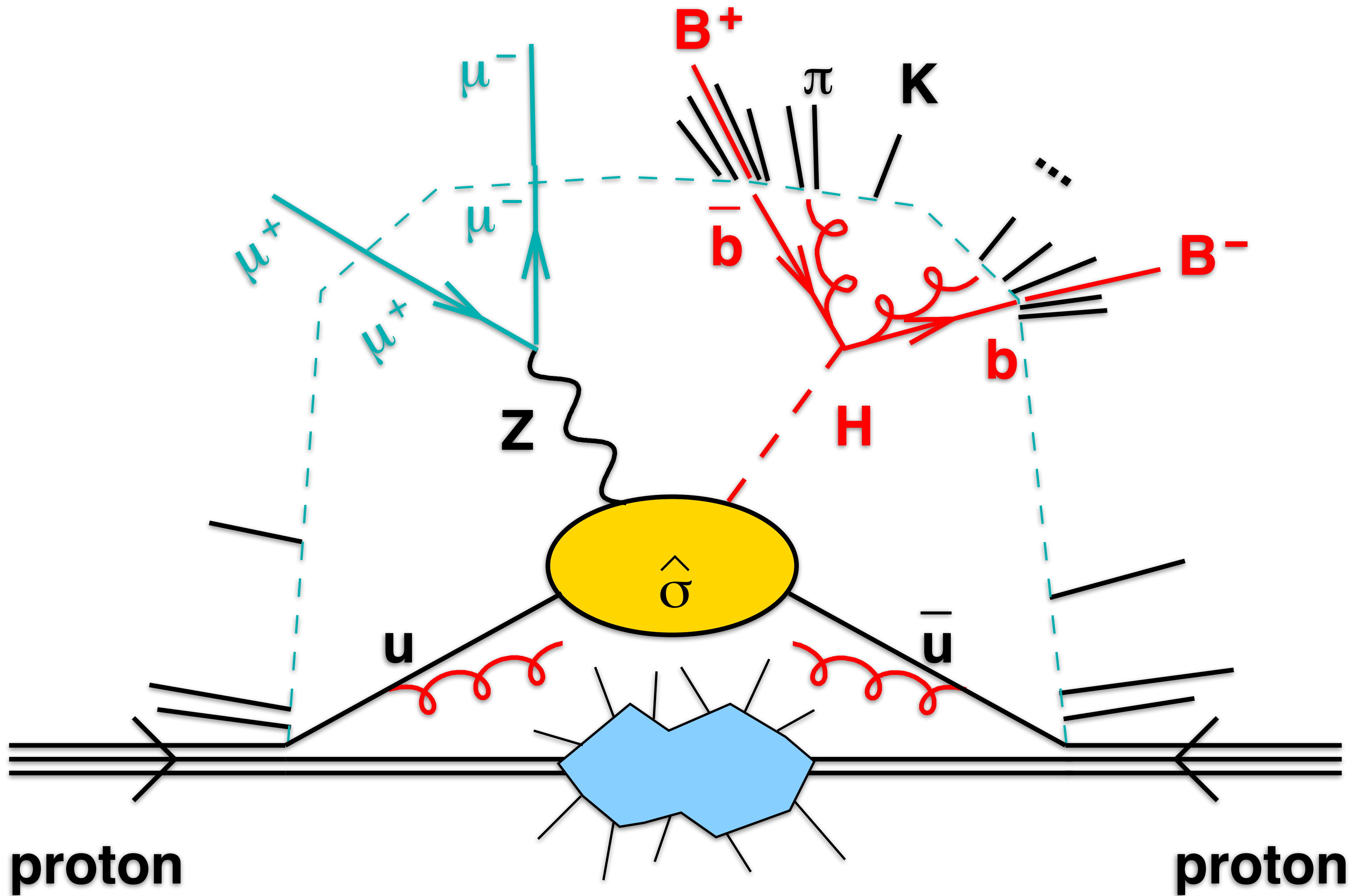




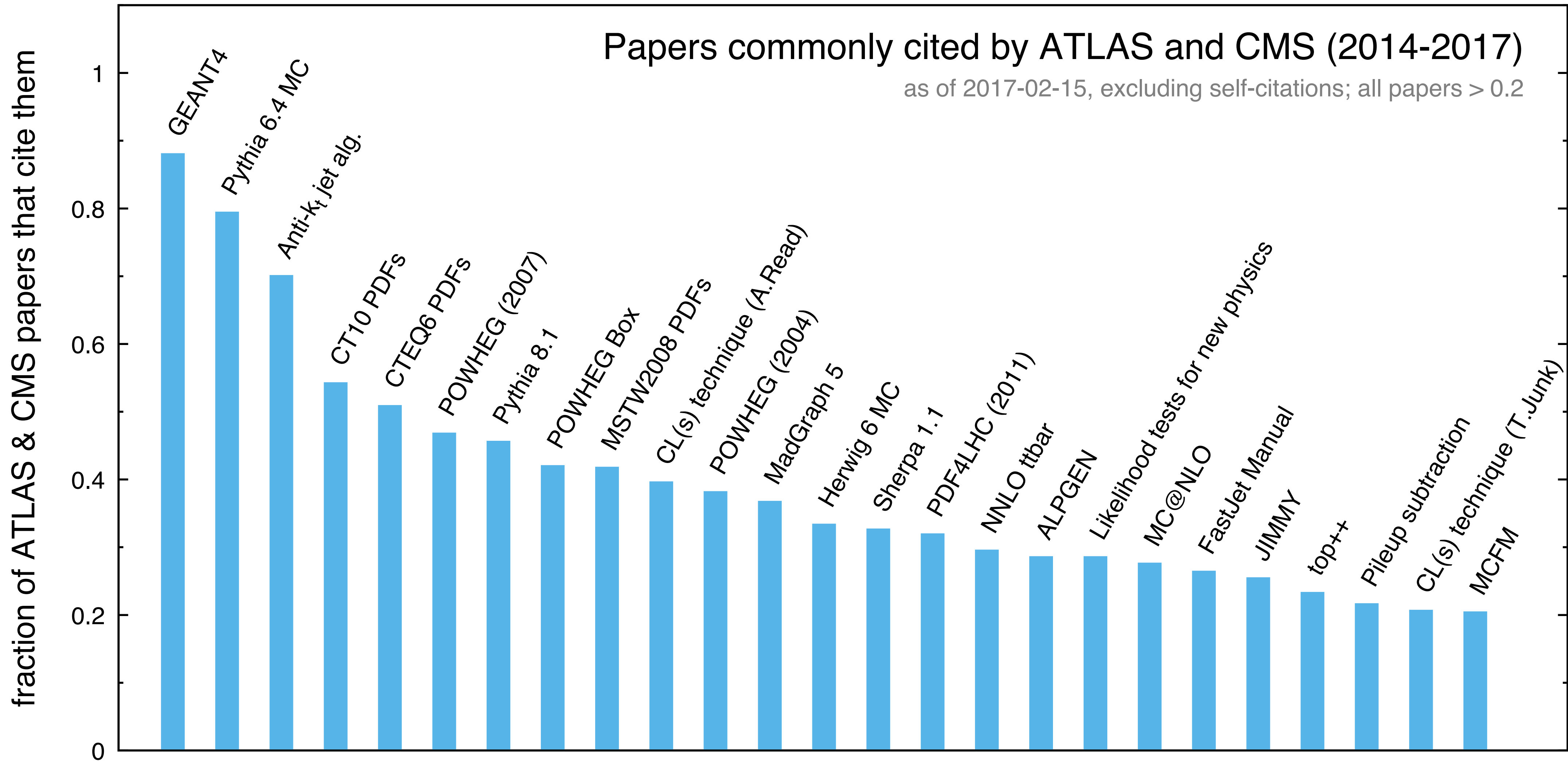
ATLAS
EXPERIMENT





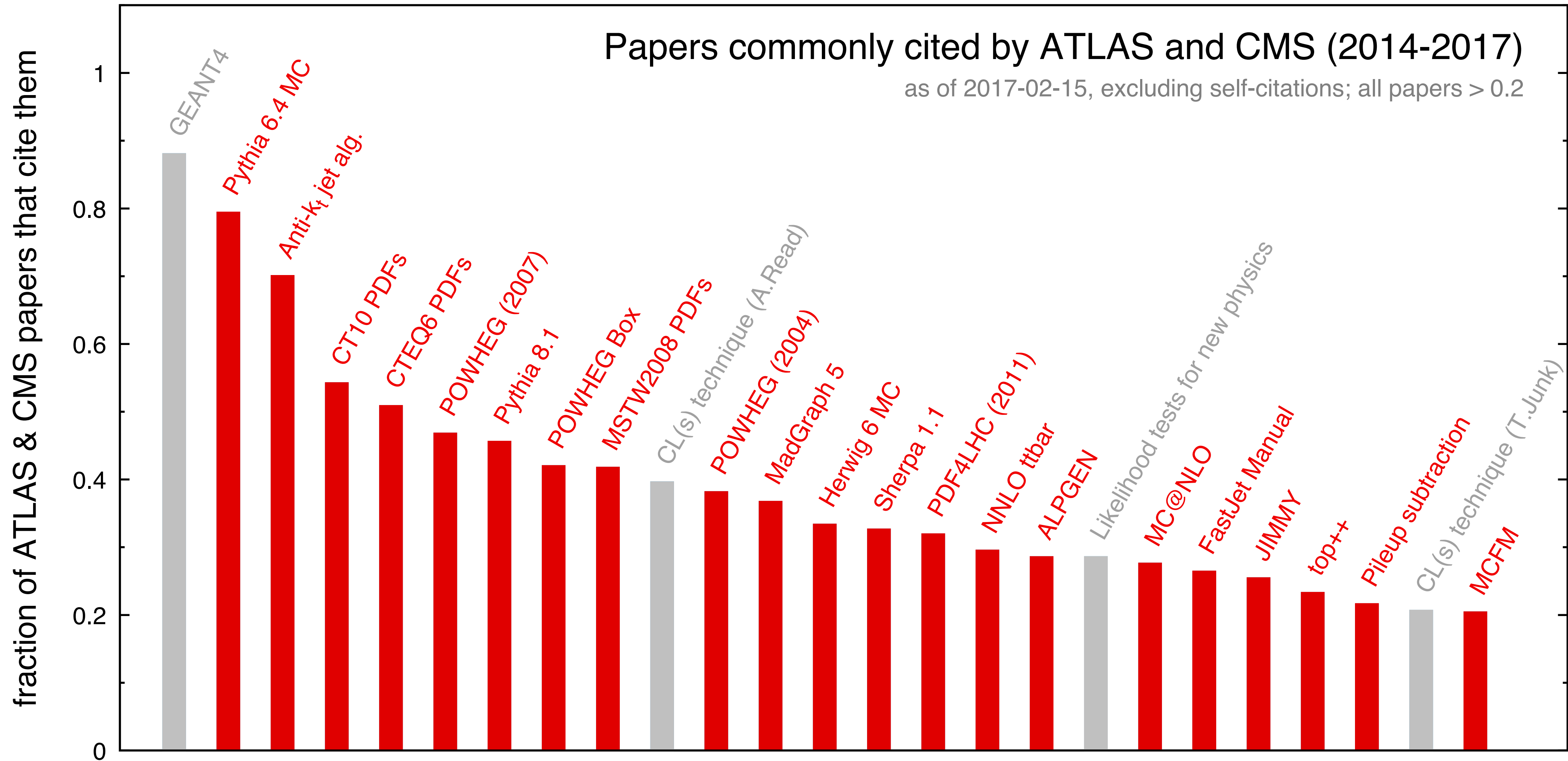


What do ATLAS & CMS use to make sense of this?



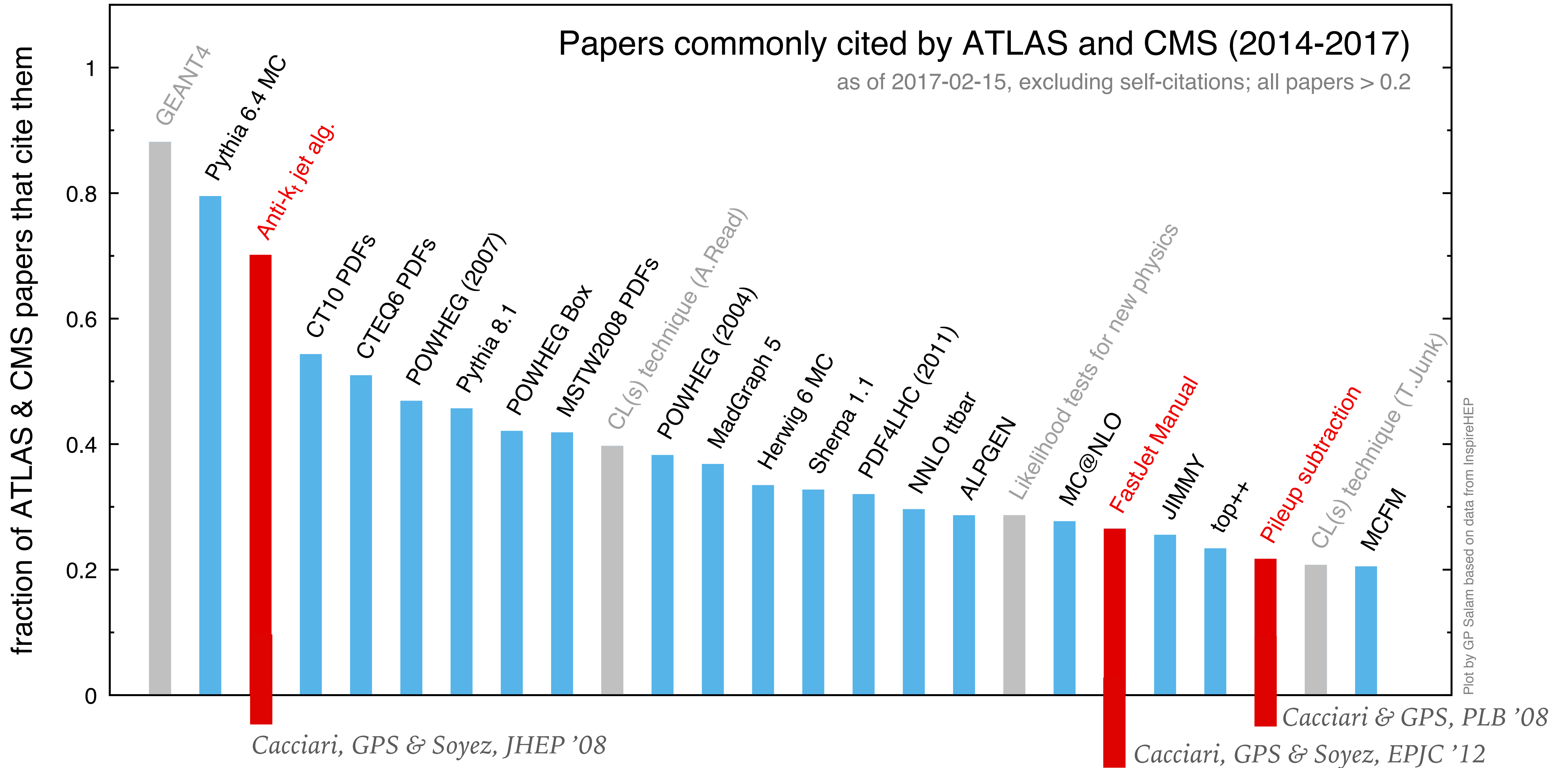
Plot by GP Salam based on data from InspireHEP

In very large part: theoretical work

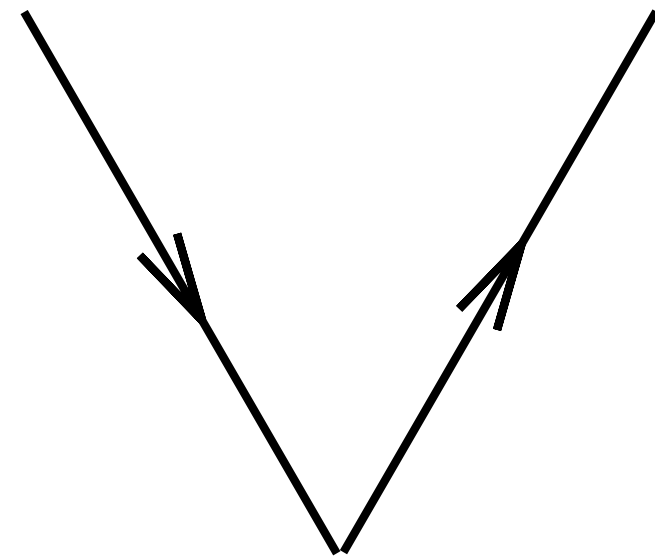


Plot by GP Salam based on data from InspireHEP

In very large part: theoretical work

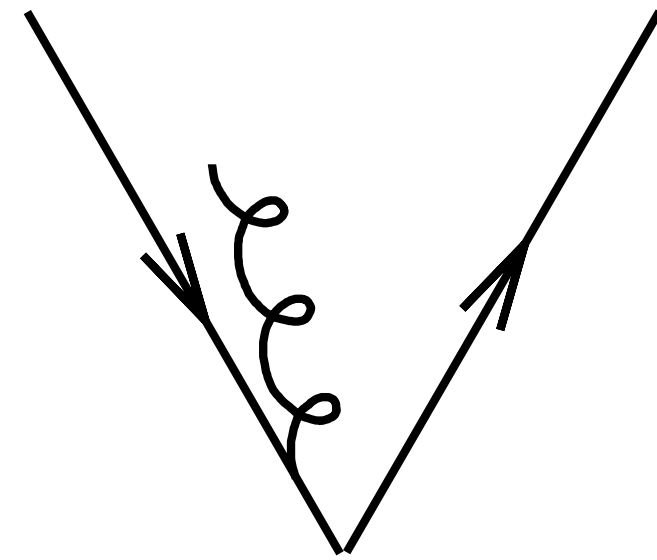
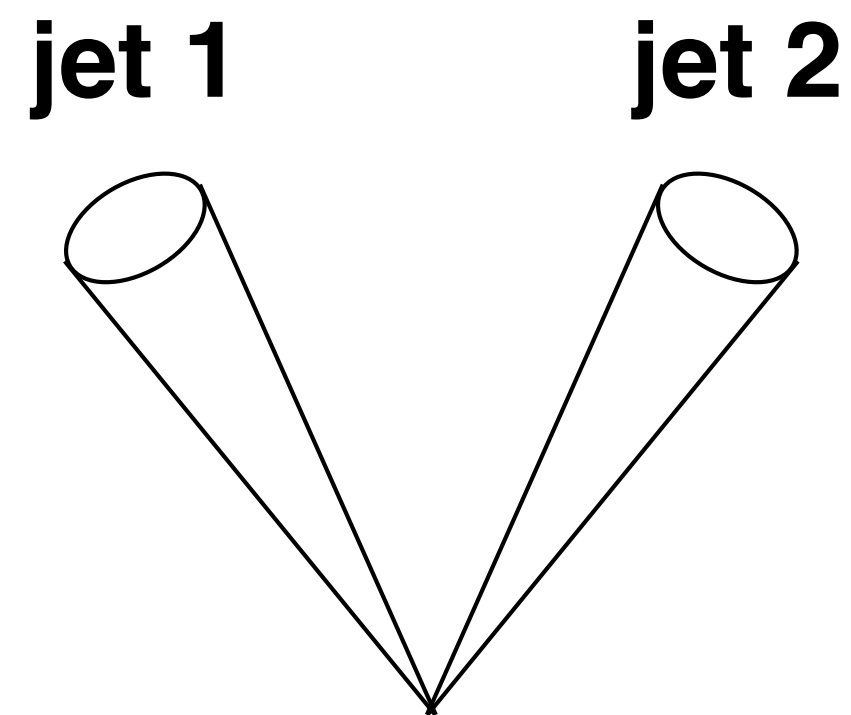


Jet definitions



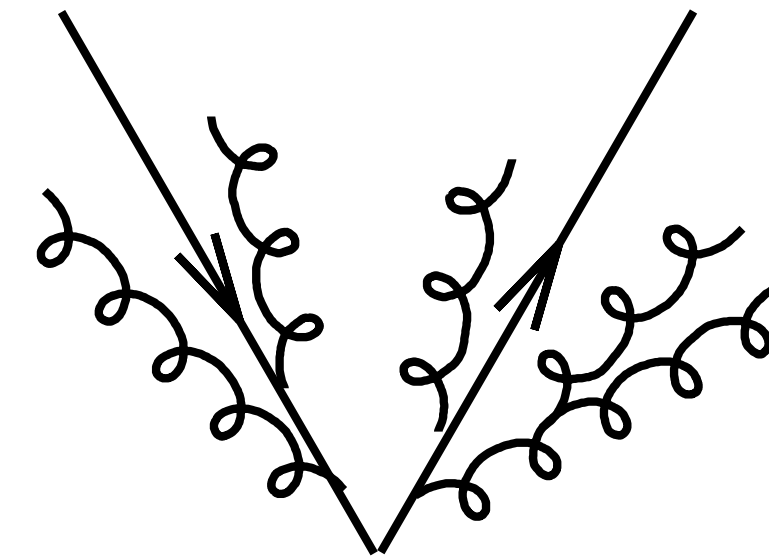
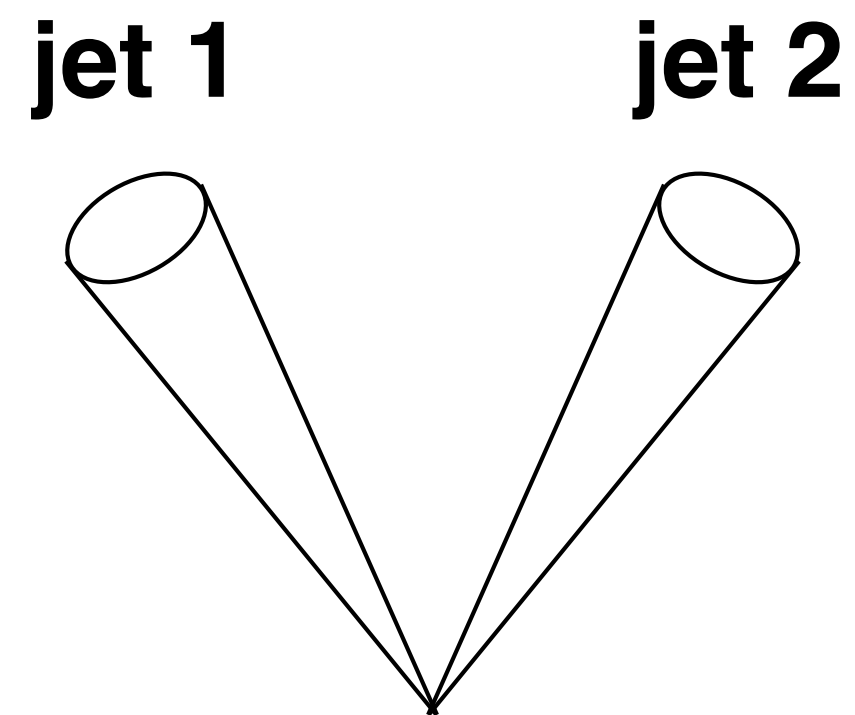
LO partons

Jet ↓ Defⁿ



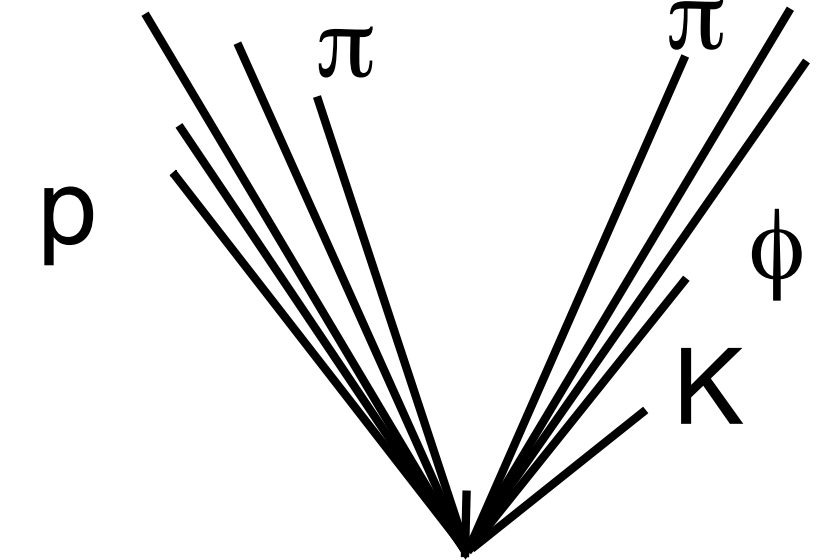
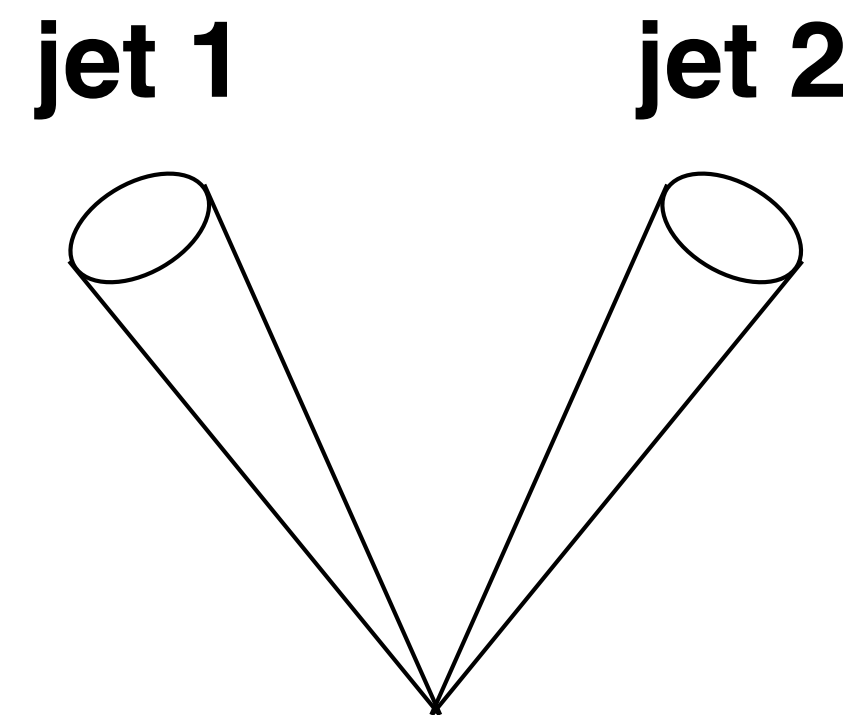
NLO partons

Jet ↓ Defⁿ



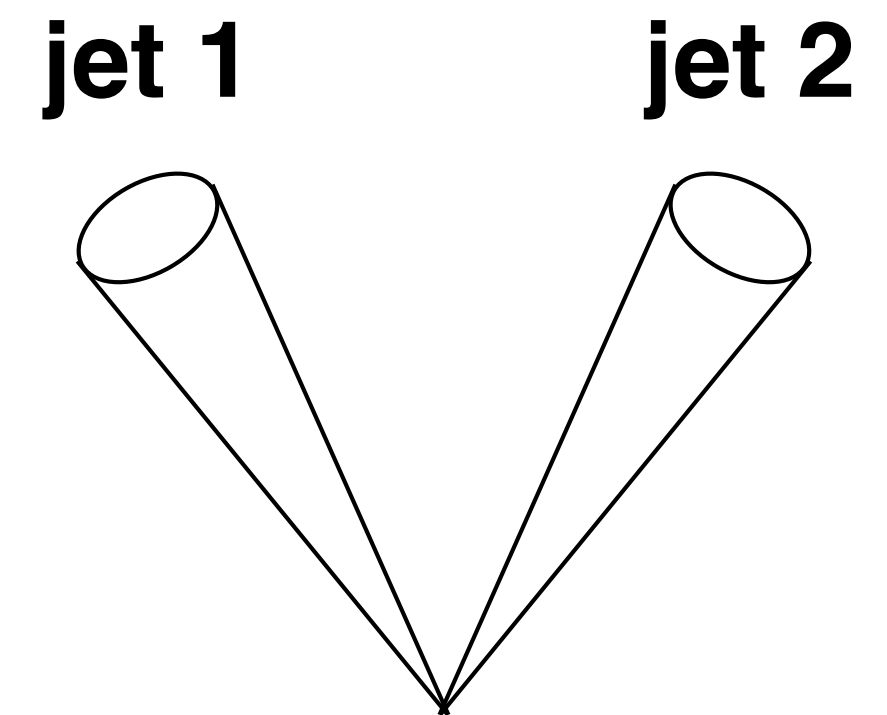
parton shower

Jet ↓ Defⁿ



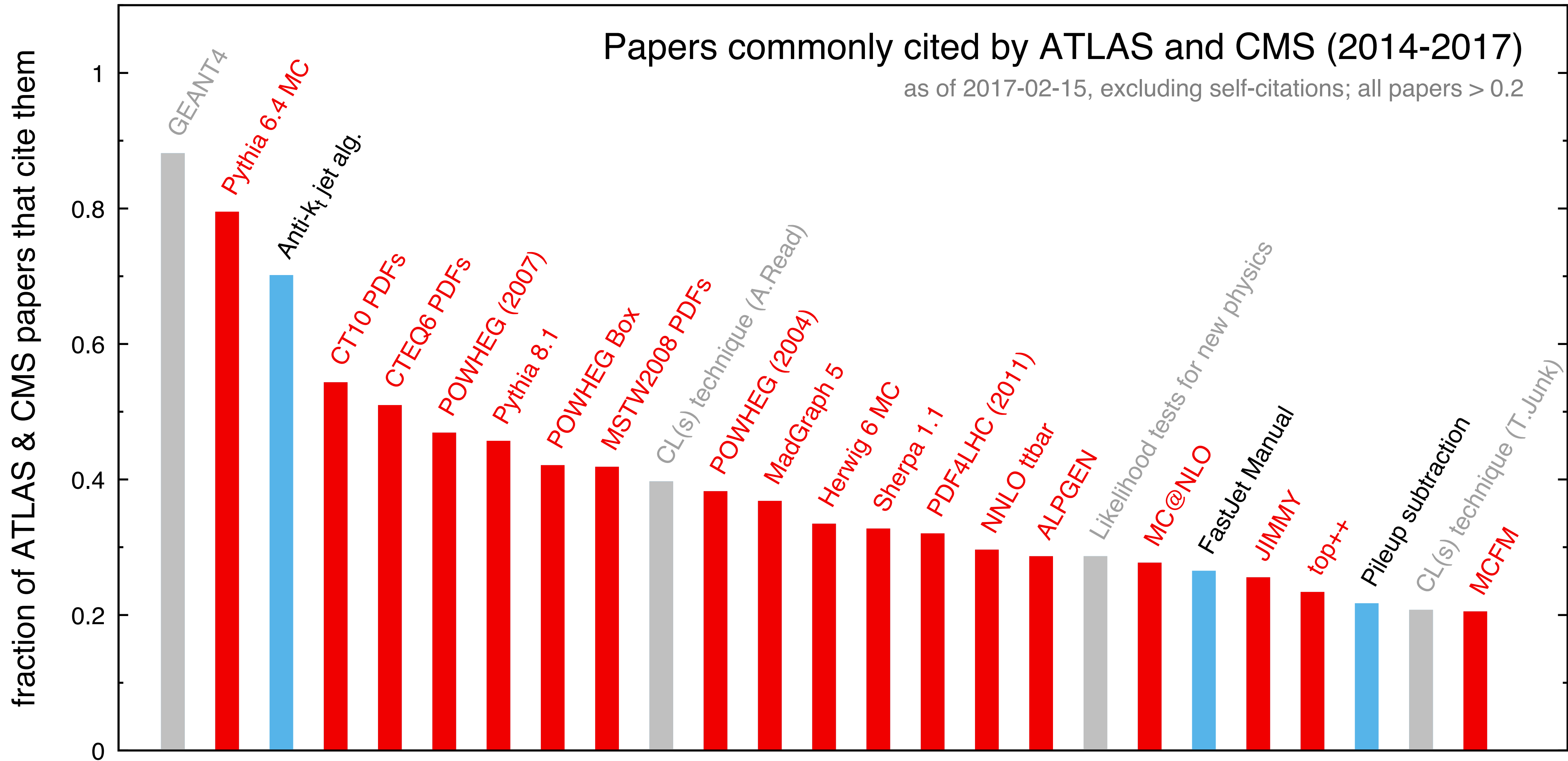
hadron level

Jet ↓ Defⁿ



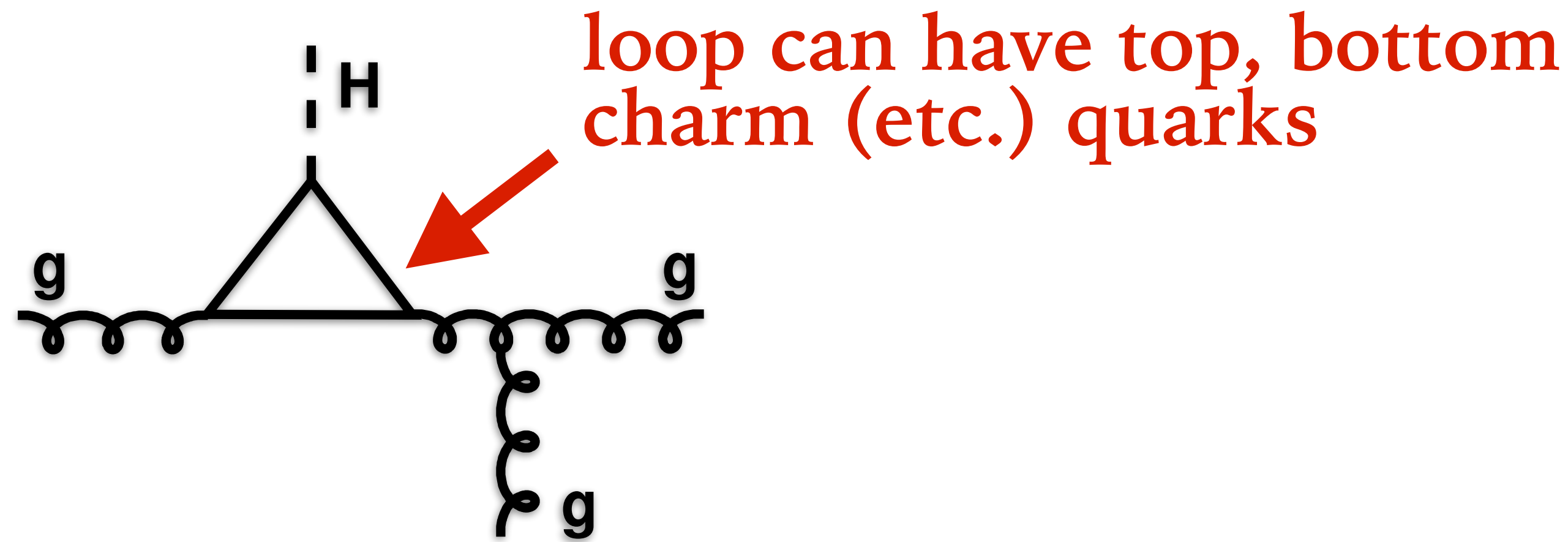
projection to jets gives simplified view of the essence of an event

Most of the rest: predicting what happens in hadron collisions



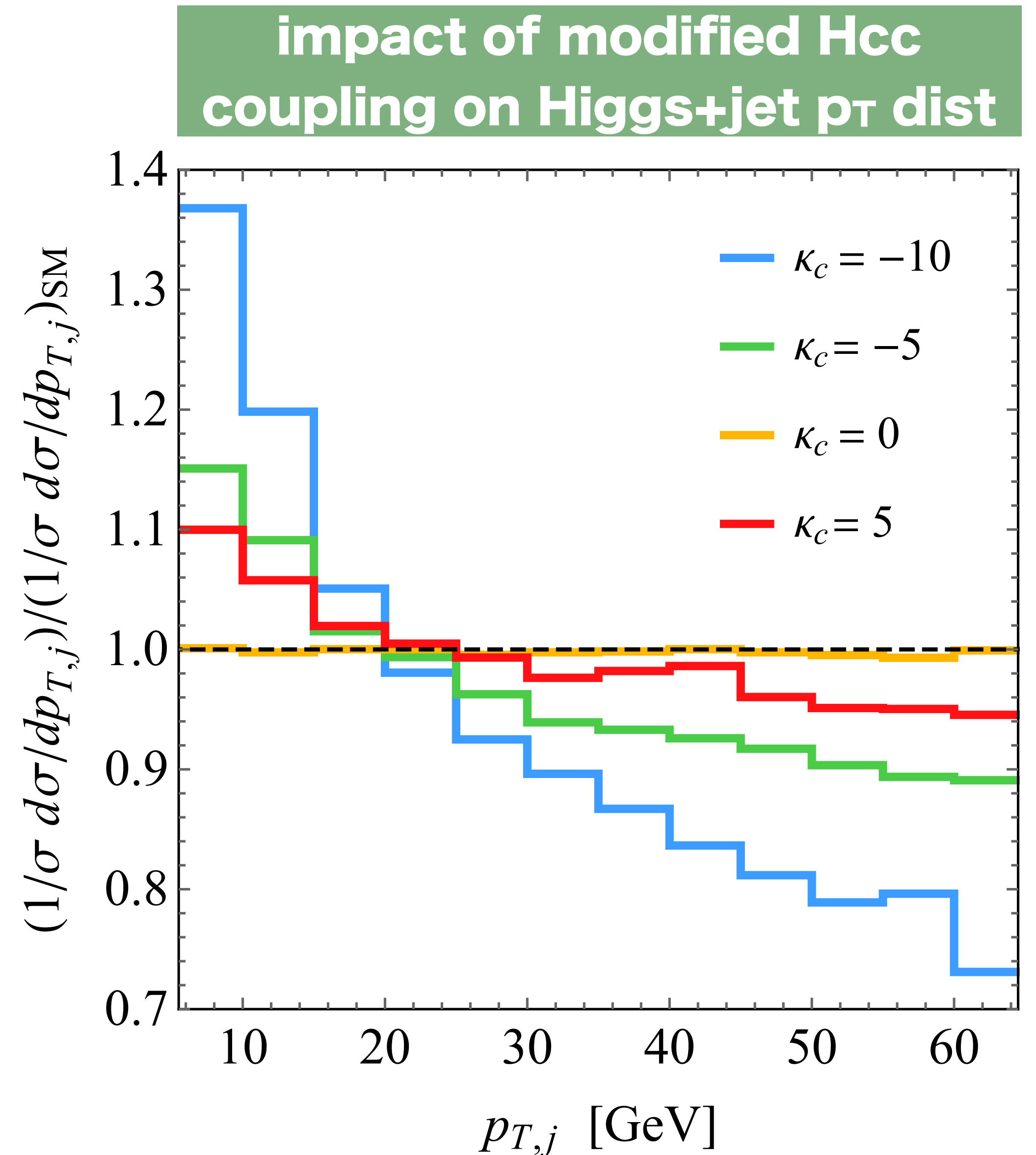
Plot by GP Salam based on data from InspireHEP

Key question today: **precision** of predictions. Here's an example why

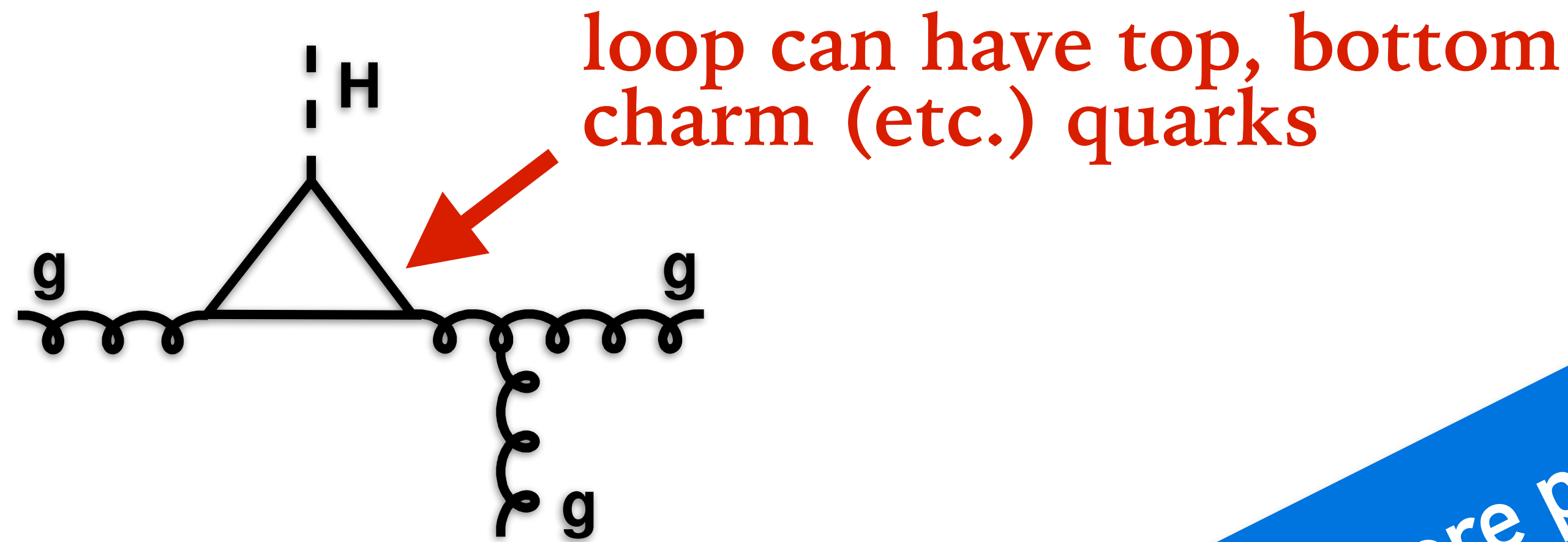


- quark mass affects momentum distribution of Higgs
- full distribution affected by relative contributions of top, bottom and charm
- sensitivity to Hcc Yukawa coupling

*Bishara, Haisch, Monni & Re, arXiv:1606.09253
(cf. also Soreq, Zhu, Zupan, arXiv:1606.09621)*



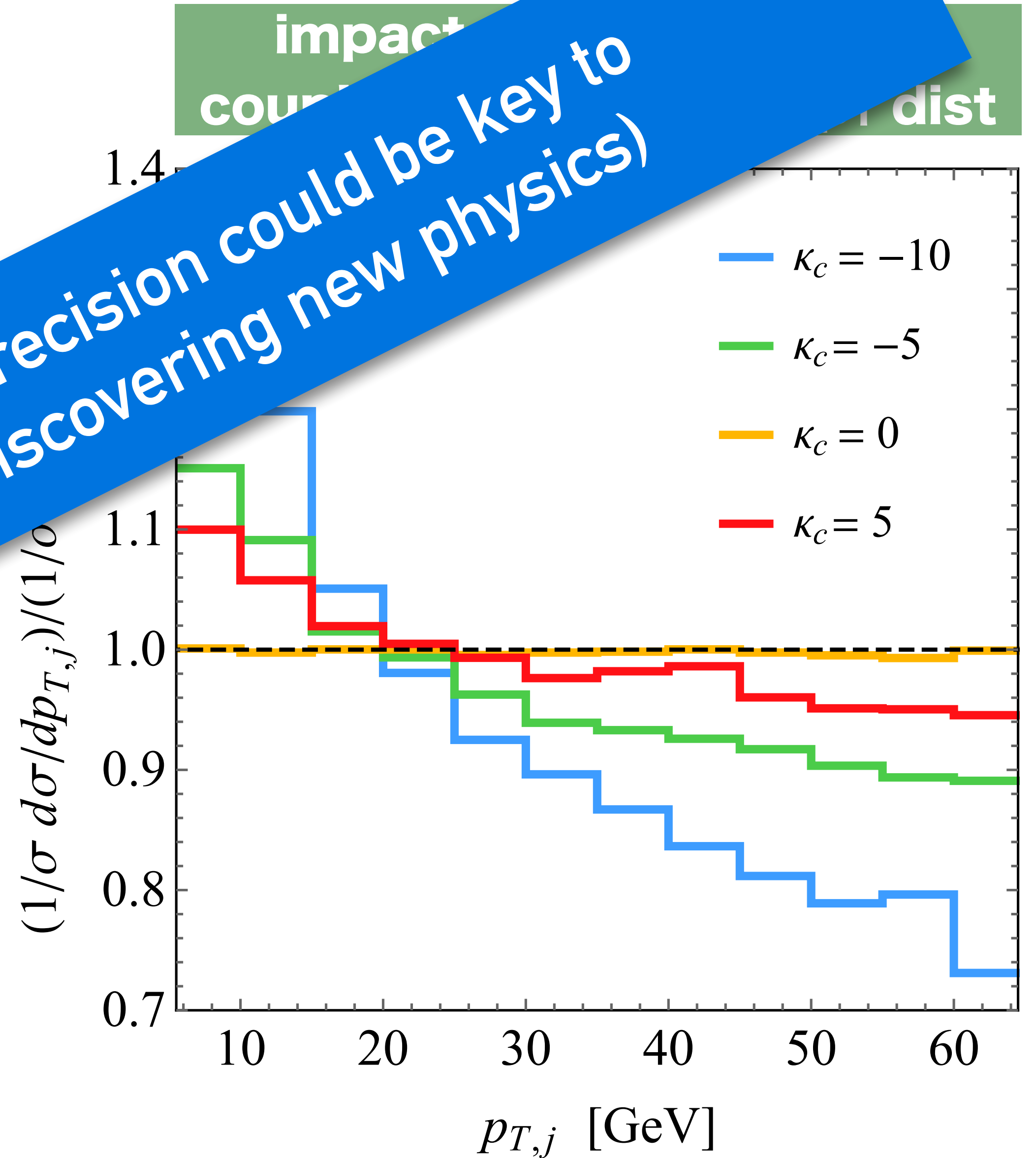
Key question today: **precision** of predictions. Here's an example why



- quark mass affects momentum of Higgs
- full distribution of Higgs decay products receive contributions from top and charm quarks
- Yukawa coupling

Biswas, Monni & Re, *arXiv:1606.09253*
 (cf. *arXiv:1606.09621*)

One of countless examples where precision could be key to establishing Standard Model (or discovering new physics)



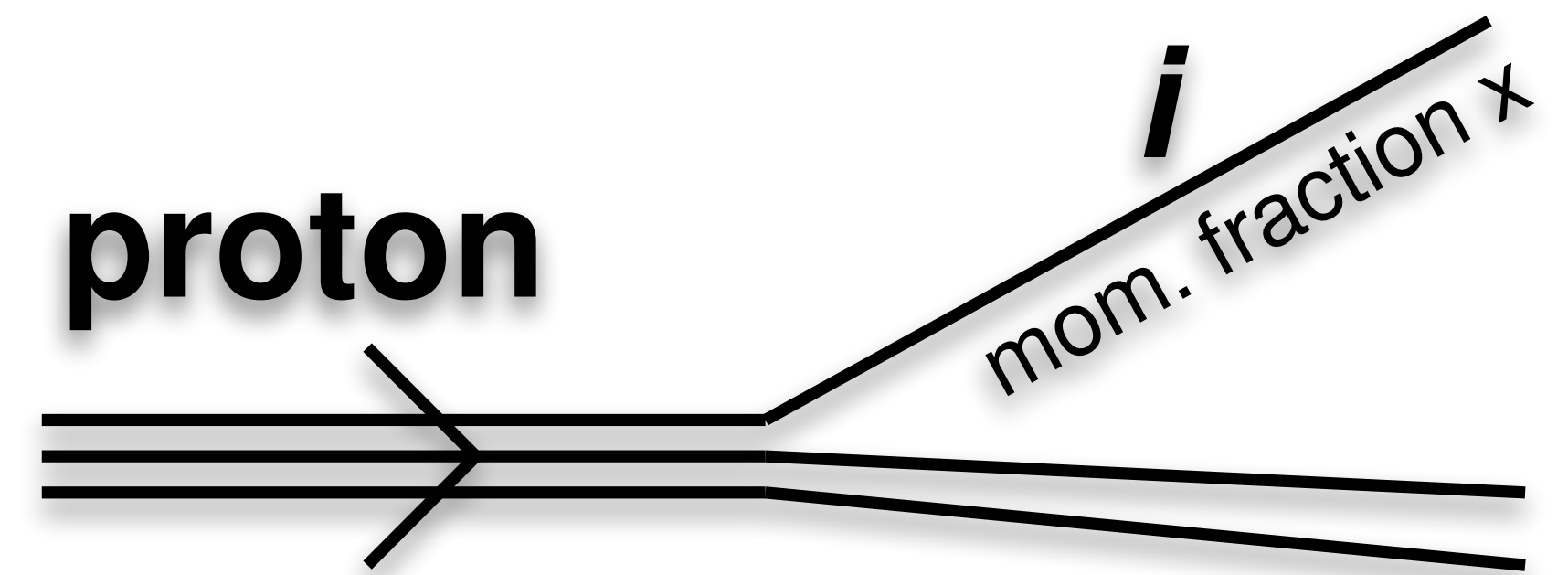
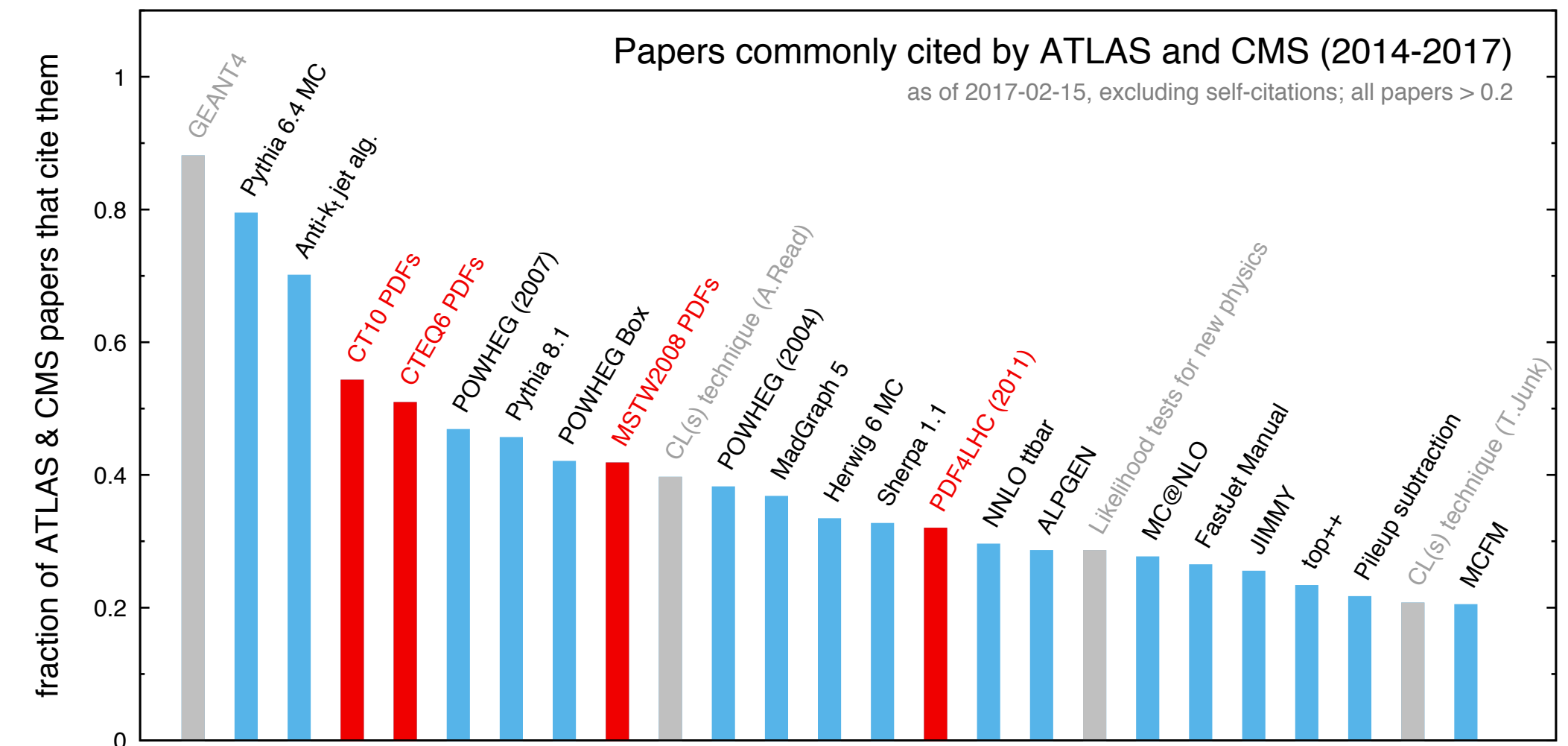
parton distribution functions (PDFs)

Protons are composite objects (uud + gluons + ...)

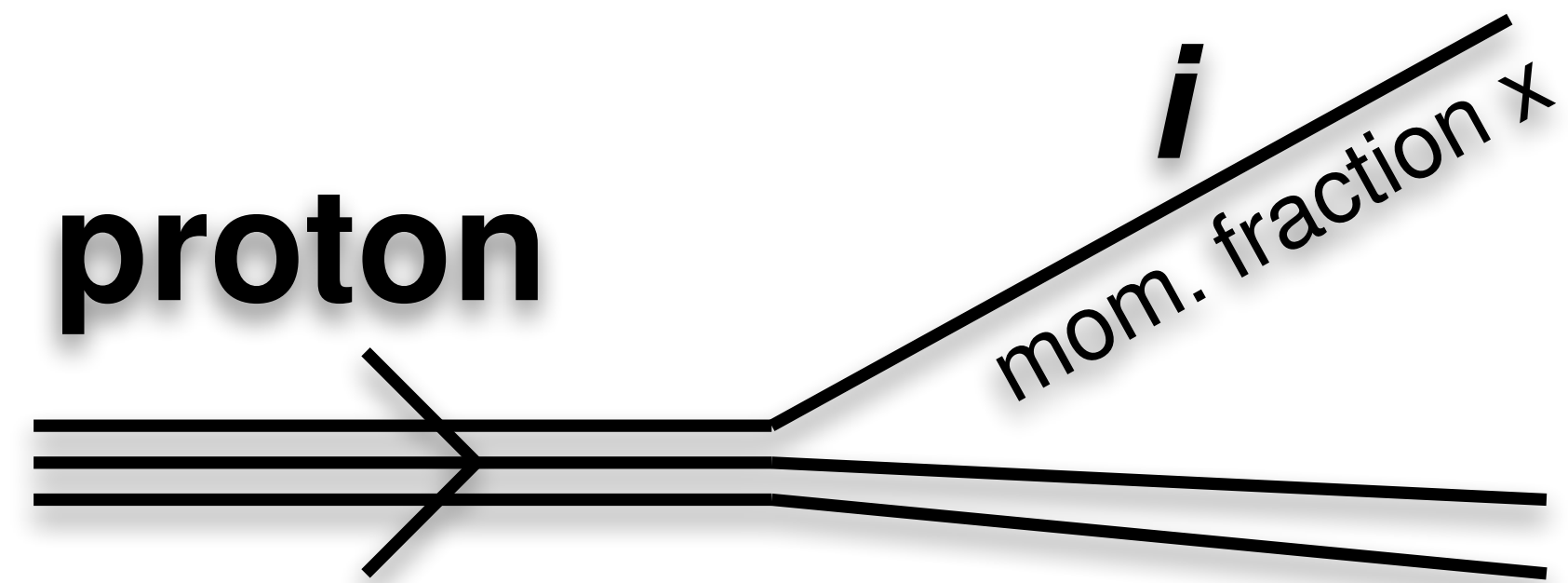
Quantitative LHC physics requires knowledge of PDFs:

$$f_{i/p}(x, \mu^2)$$

- number of partons of flavour i [=u, d, g, ...]
- inside a fast-moving proton p
- carrying a fraction x of the proton's momentum
- when viewed with resolution momentum scale μ [$\sim 1/\text{wavelength of probe}$]



parton distribution functions (PDFs)

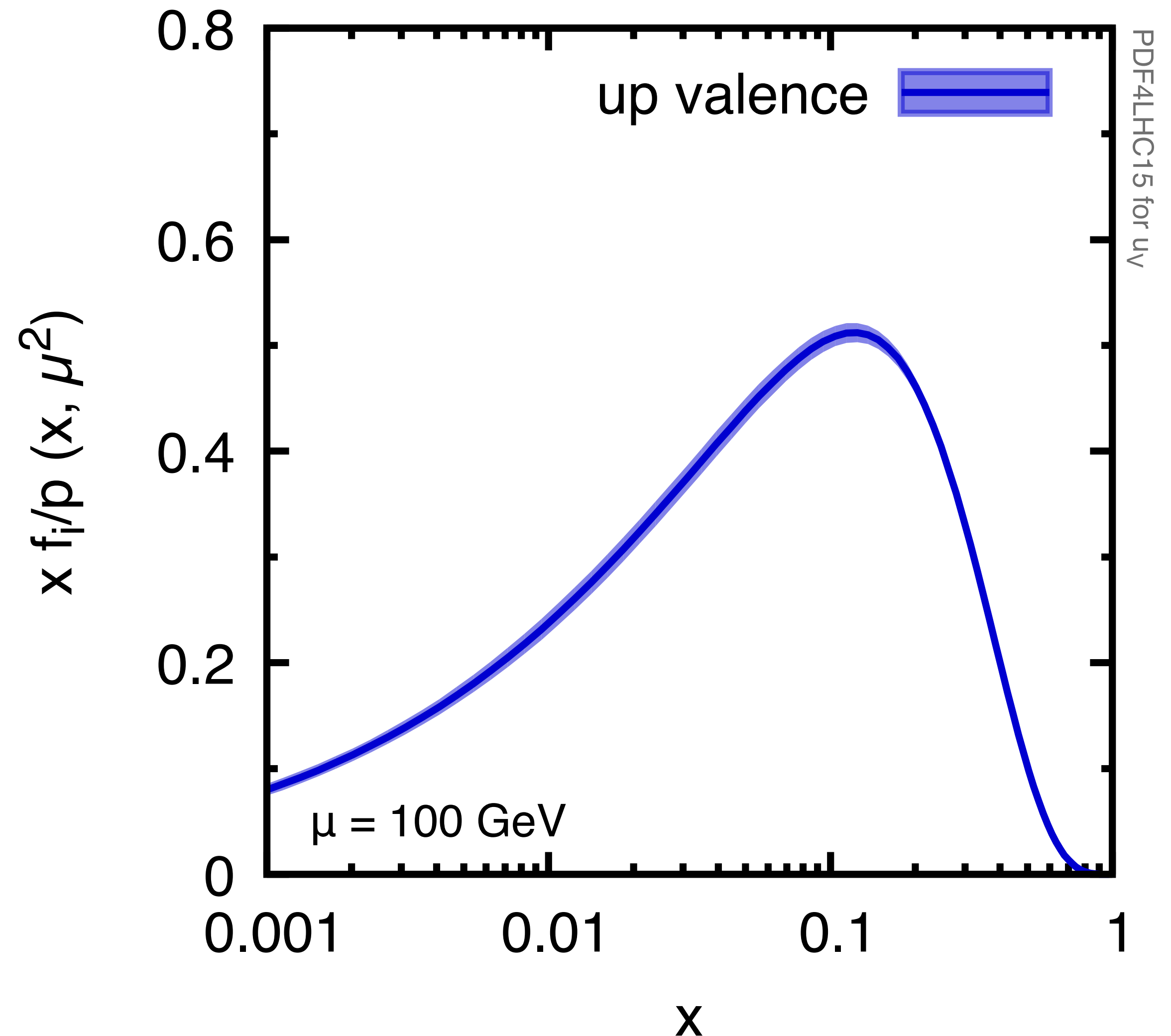


LHC physics
needs PDFs in region
 $\sim 10^{-3} - 0.5$

Typically known with good
precision $\sim 1-3\%$

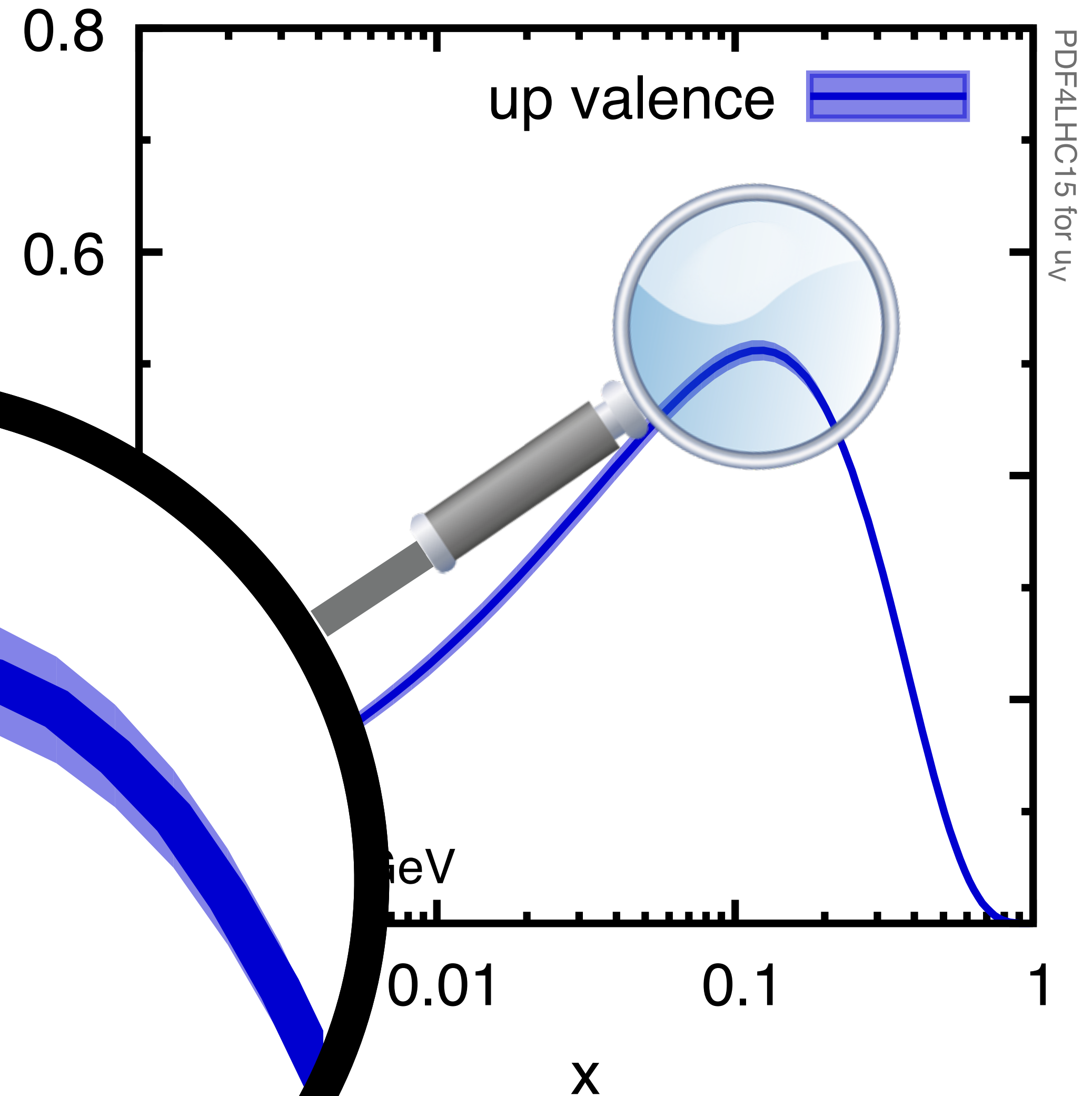
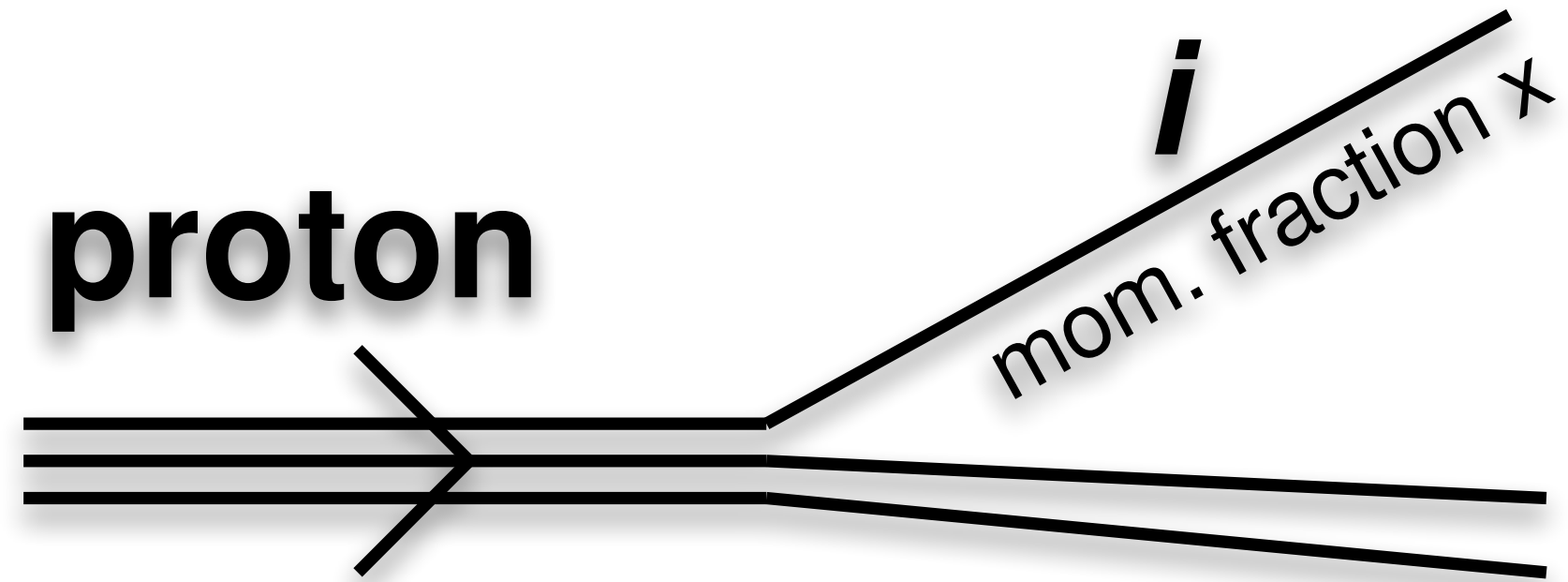
E.g. NNPDF, MMHT, CT & PDF4LHC working group (+ also HERAPDF, ABM, ...)

$$f_{i/p}(x, \mu^2)$$



parton distribution functions (PDFs)

$$f_{i/p}(x, \mu^2)$$



LHC physics
needs PDFs in re
 $\sim 10^{-3} - 0.$

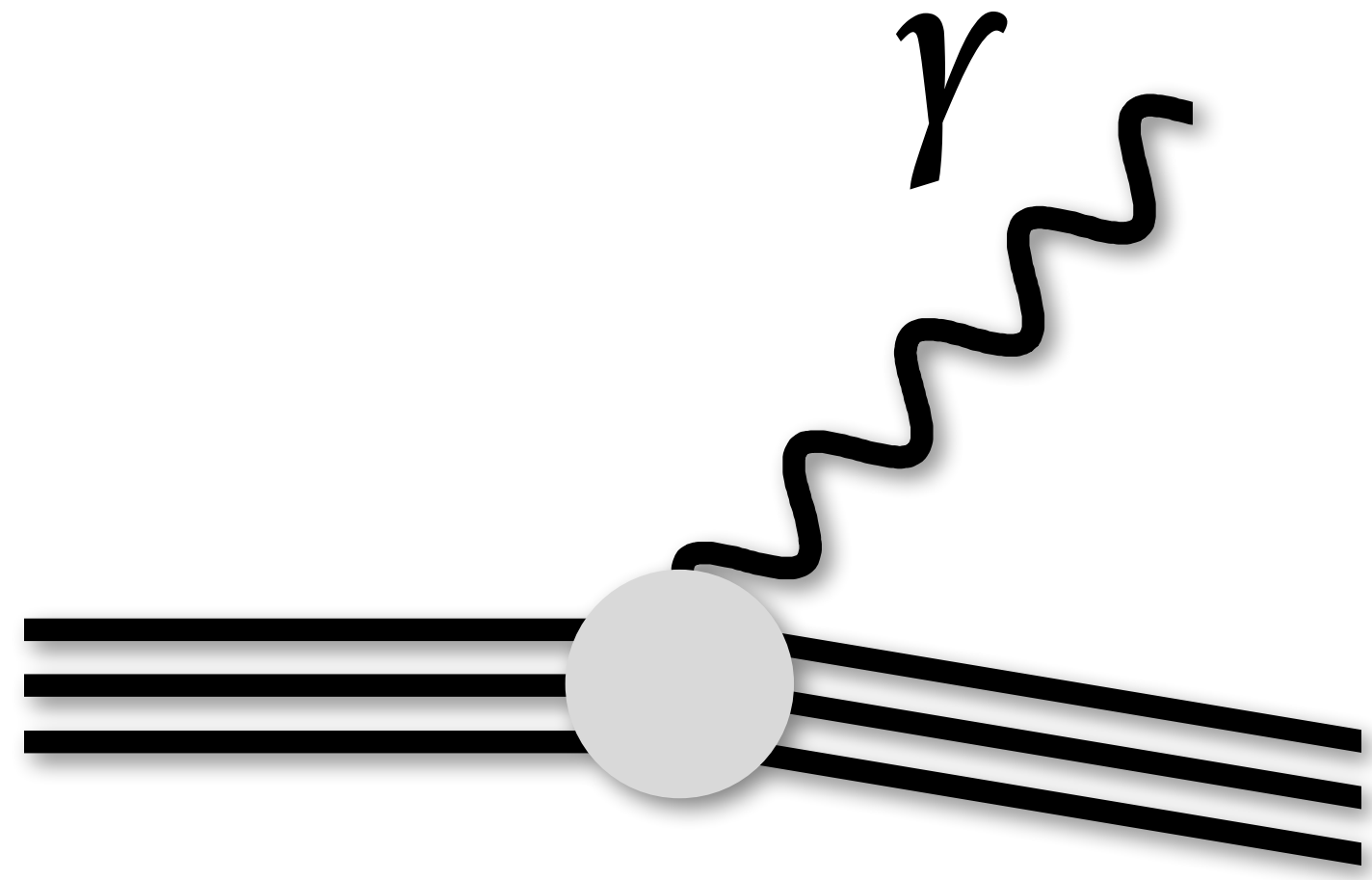
Typically known w

precision $\sim 1-$

E.g. NNPDF, MMHT, CT & PDF4LHC

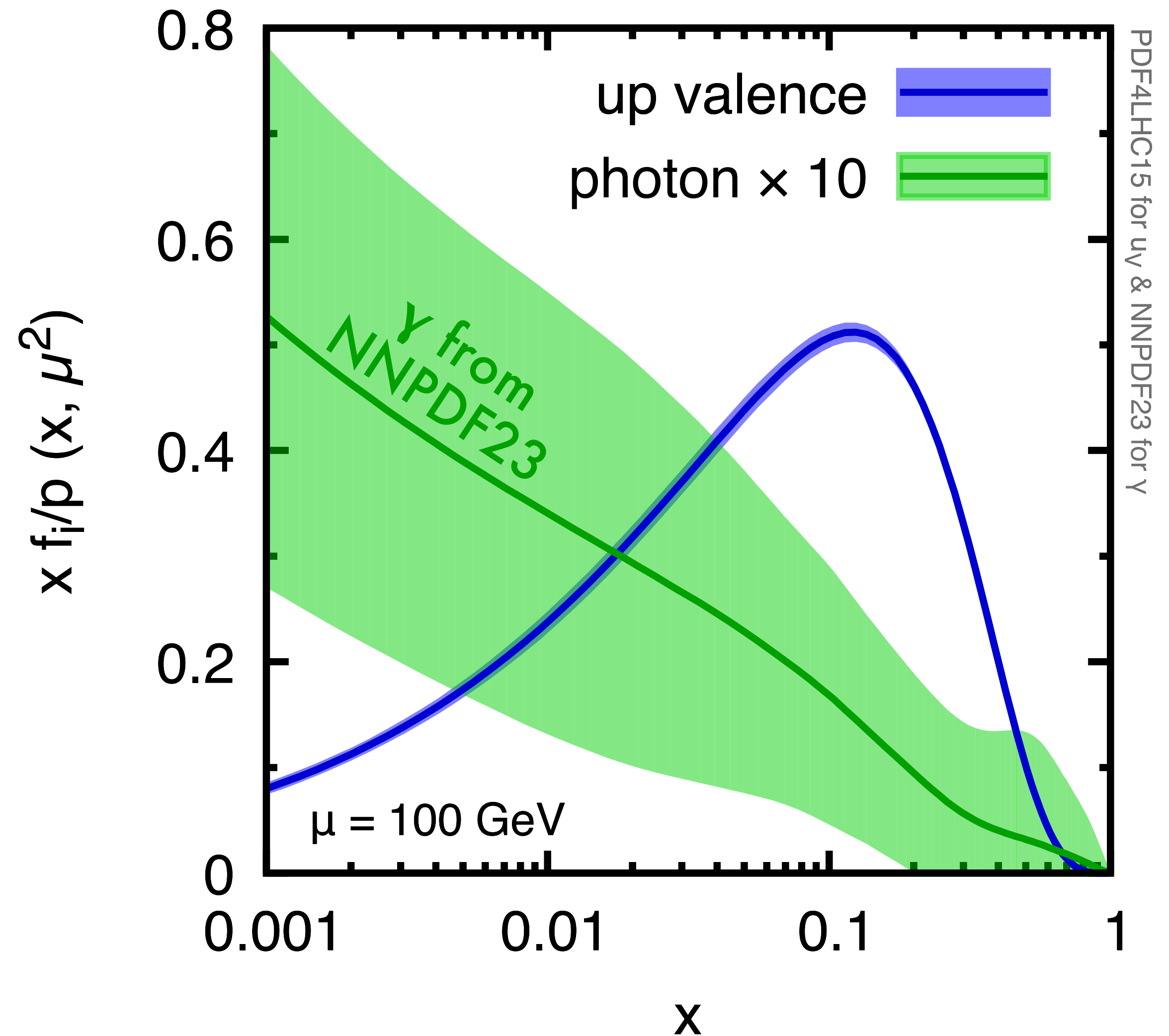
(ABM, ...)

parton distribution functions (PDFs)



One exception:
the photon distribution
inside the proton
(had up to 100% uncertainty)

$$f_{i/p}(x, \mu^2)$$



the **photon** distribution inside the proton

or “how much light accompanies a fast-moving proton?”

based on Manohar, Nason, GPS & Zanderighi

PRL '16 (Editors' Suggestion)

+ work in progress

photon induced contribution to HW production

$pp \rightarrow H W^+ (\rightarrow l^+ \nu) + X$ at 13 TeV

non-photon induced contributions

91.2 ± 1.8 fb

photon-induced contriibs (NNPDF23)

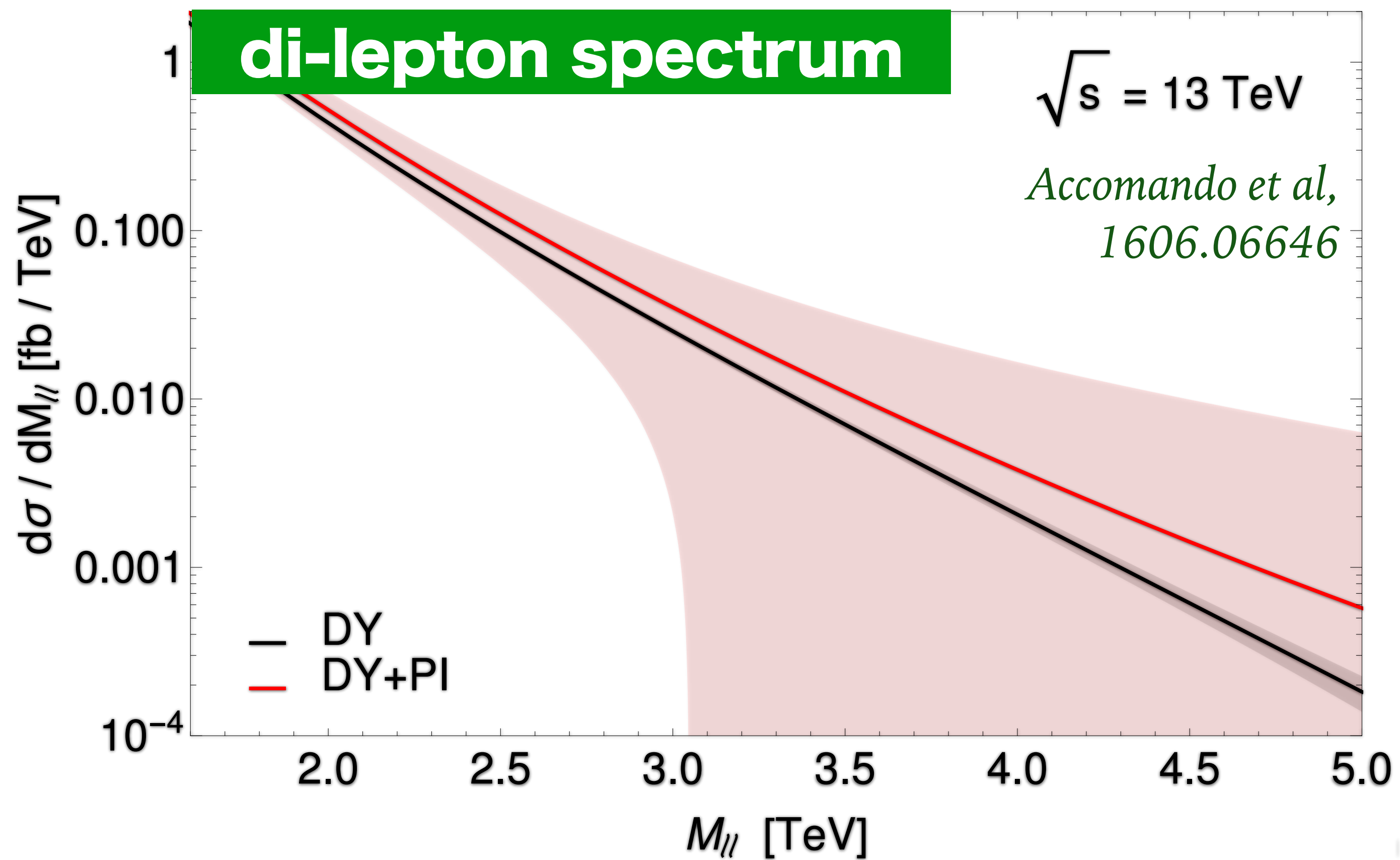
$6.0^{+4.4}_{-2.9}$ fb

**photon
contribution
brings the
largest overall
uncertainty**

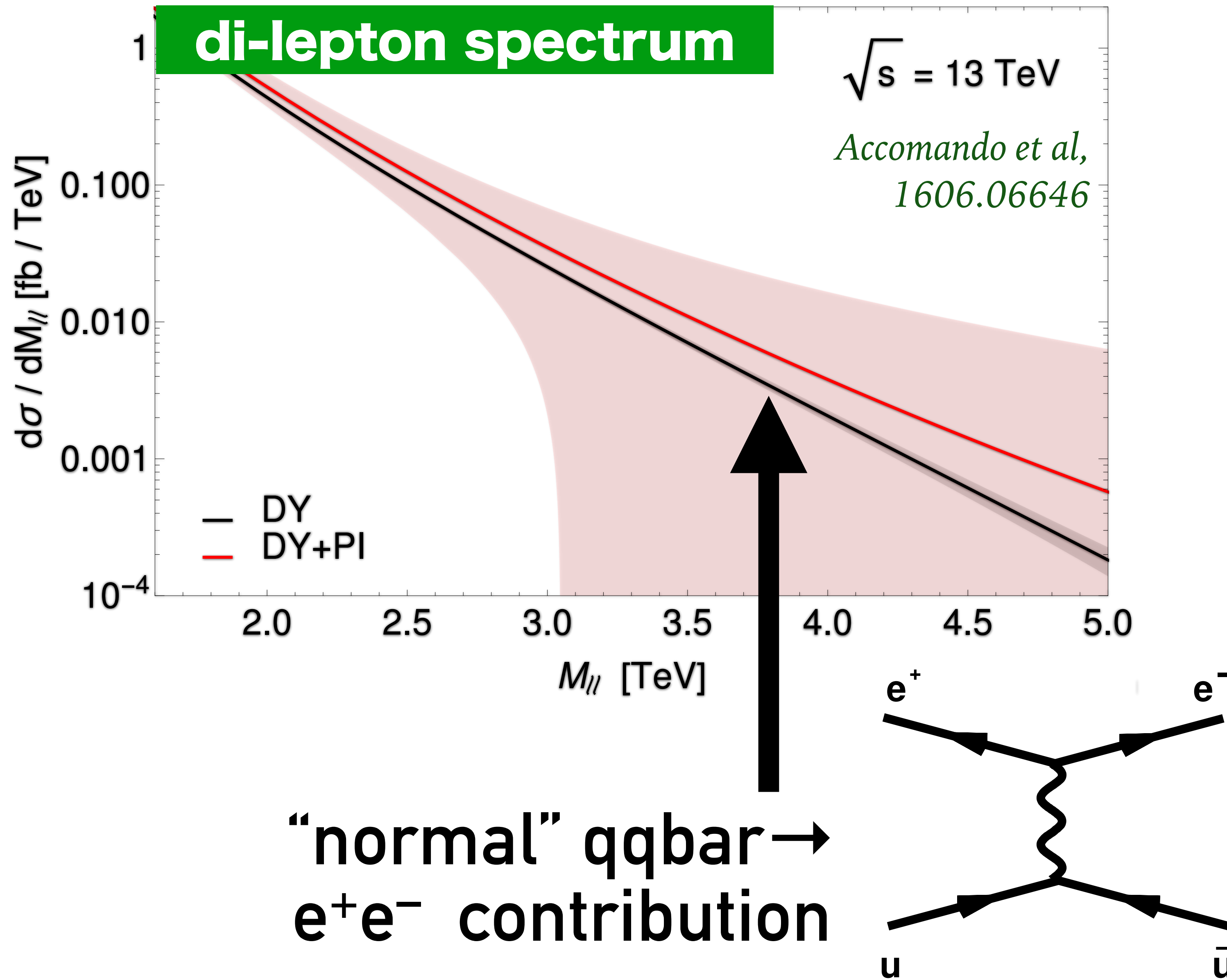


*non-photon numbers from LHCHSWG (YR4)
including PDF uncertainties*

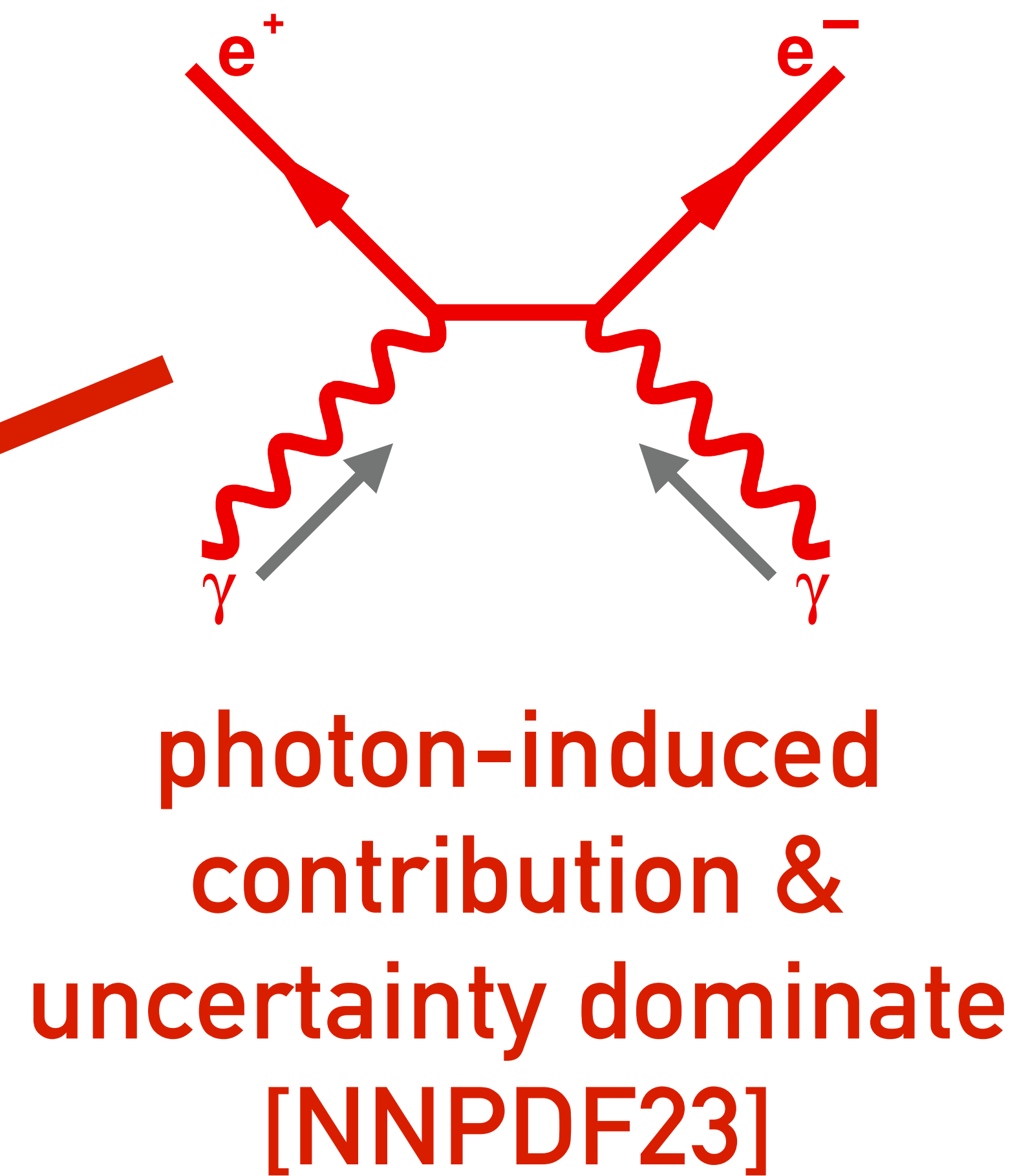
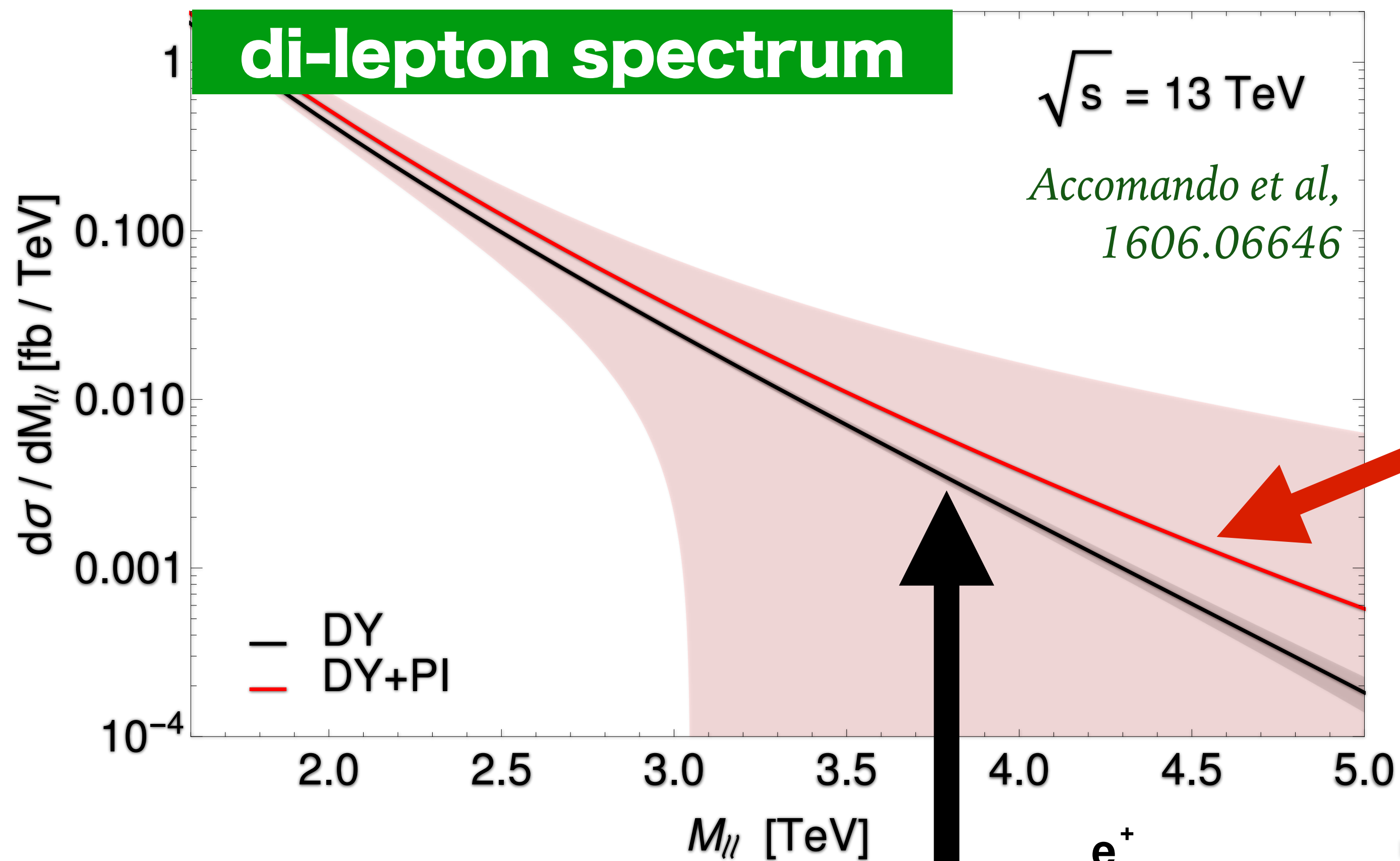
it matters in new-physics searches



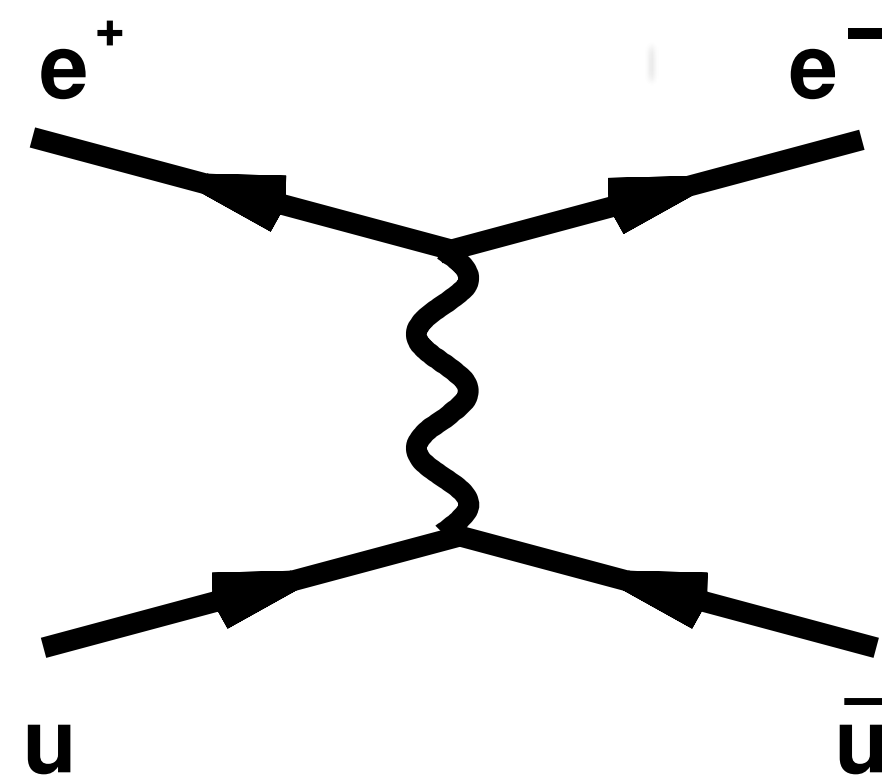
it matters in new-physics searches



it matters in new-physics searches



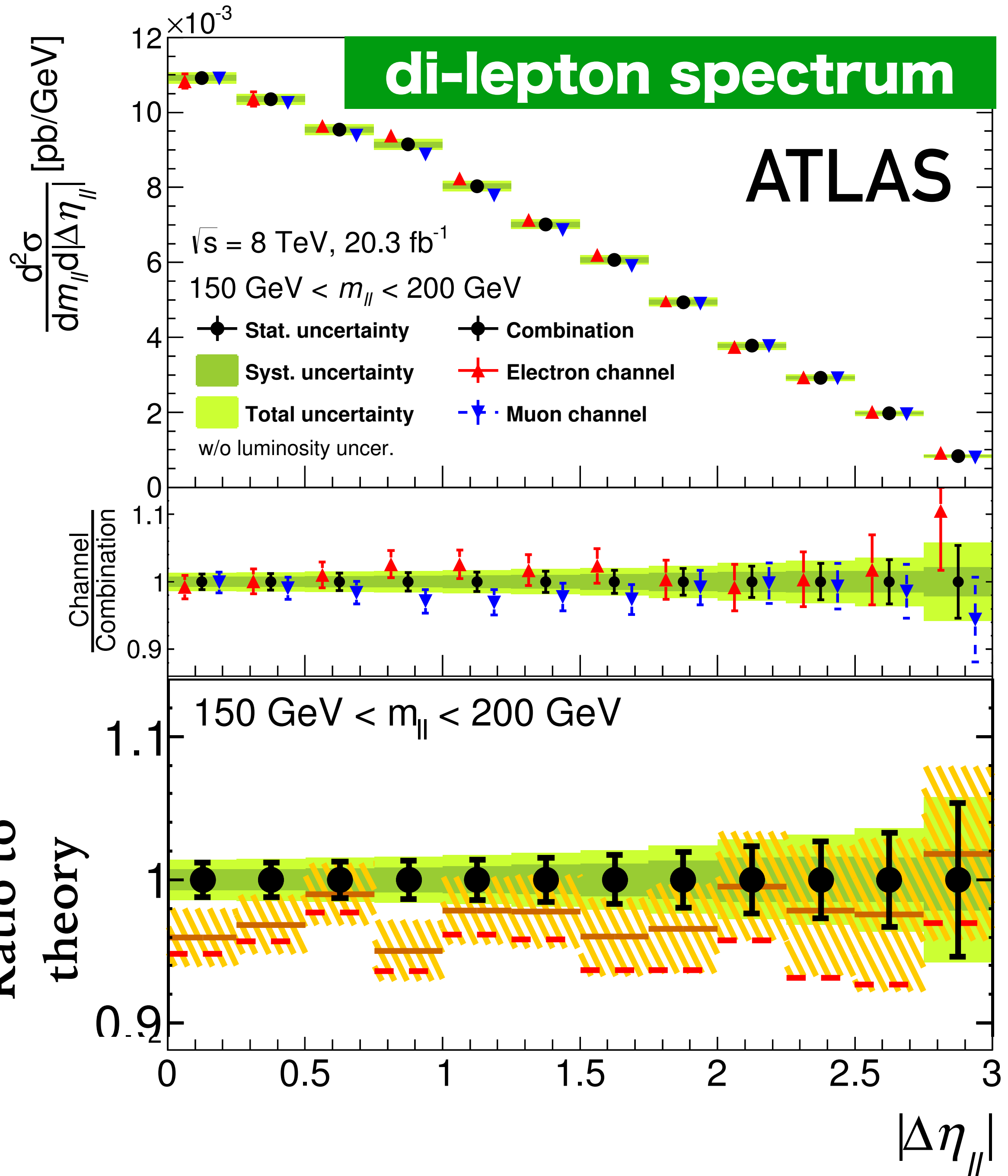
“normal” $q\bar{q} \rightarrow e^+e^-$ contribution



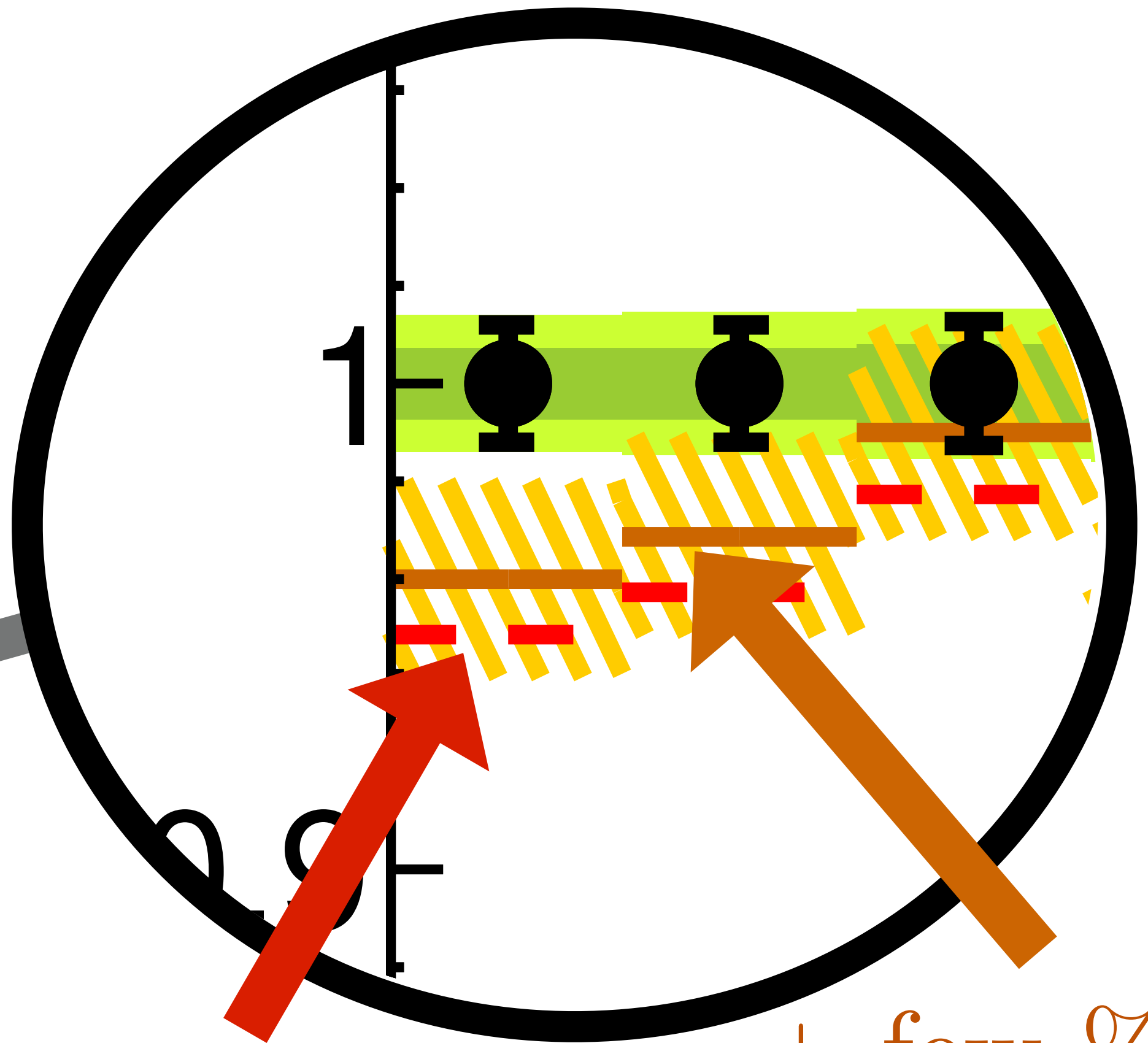
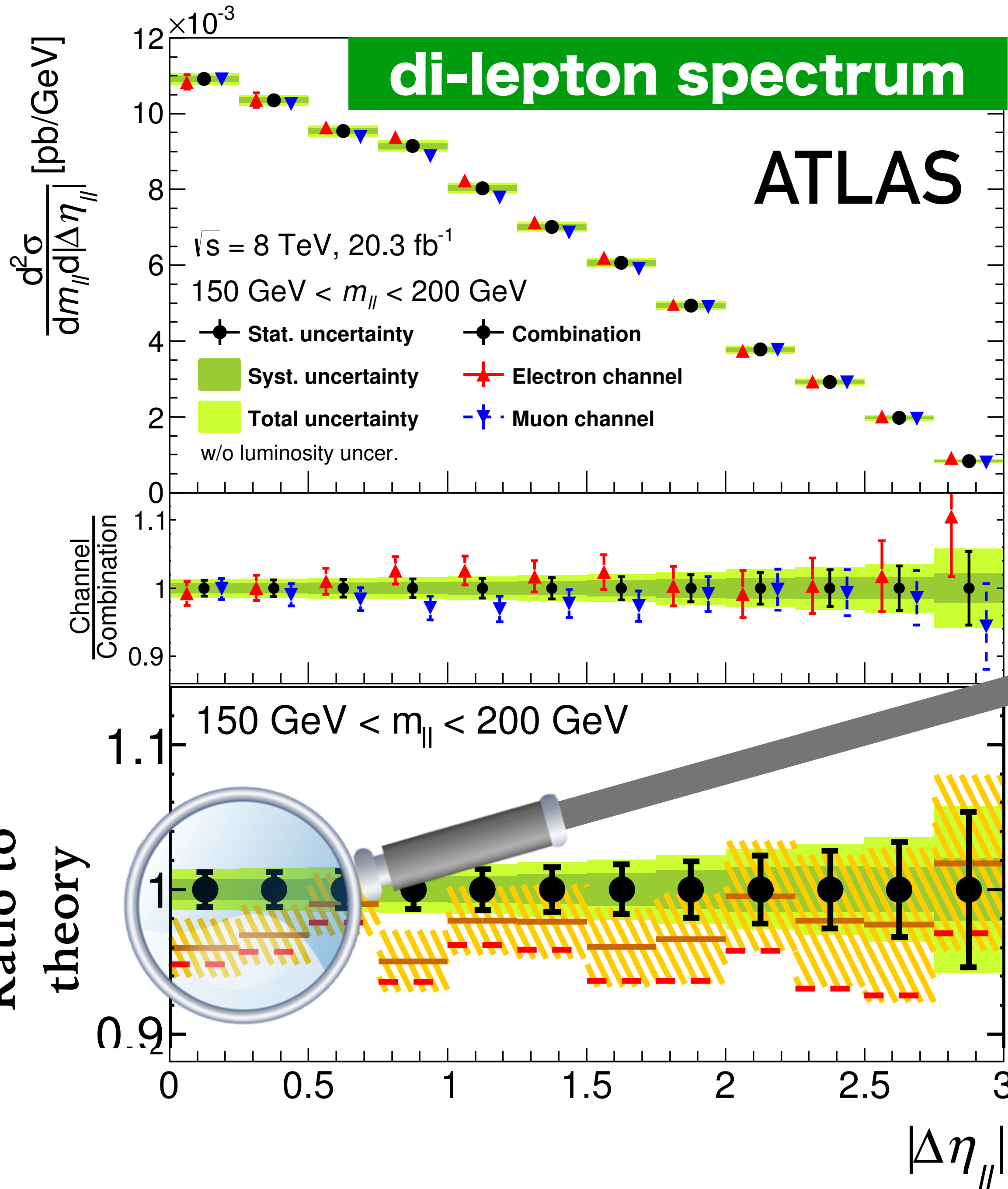
model-independent γ PDF fit (c. 2013)

di-lepton spectrum

ATLAS



model-independent γ PDF fit (c. 2013)



95-99% from $q\bar{q} \rightarrow e^+e^-$

+ few % from $\gamma\gamma \rightarrow e^+e^-$

**is there another
way of doing this?**

photon distribution from fast-moving charged particle

Point-like particle, e.g. electrons

► Fermi, Z. Phys. 1924 ; von Weizsäcker, Z. Phys 1924; Williams, Phys.Rev. 1934

$$f_{\gamma/e}(x, \mu^2) = \frac{\alpha}{2\pi} \left[\frac{1 + (1-x)^2}{x} \log \left(\frac{1-x}{x^2} \frac{\mu^2}{m_e^2} \right) - 2 \frac{1-x - x^2 \frac{m_e^2}{\mu^2}}{x} \right]$$

photon distribution from fast-moving charged particle

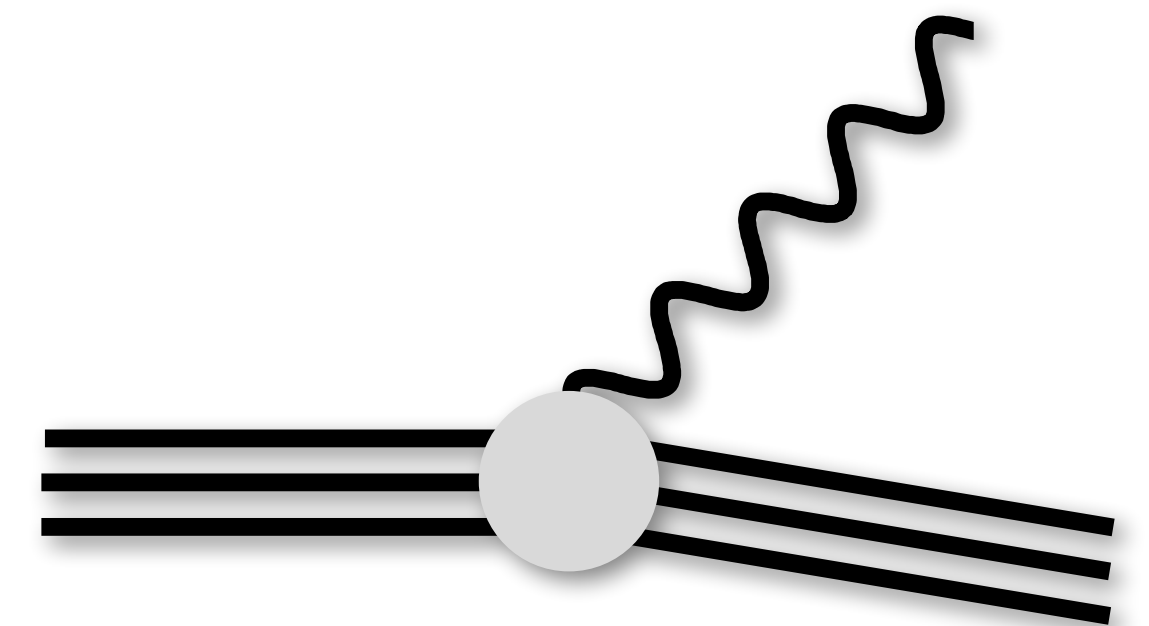
Point-like particle, e.g. electrons

- Fermi, Z. Phys. 1924 ; von Weizsäcker, Z. Phys 1924; Williams, Phys.Rev. 1934

$$f_{\gamma/e}(x, \mu^2) = \frac{\alpha}{2\pi} \left[\frac{1 + (1-x)^2}{x} \log \left(\frac{1-x}{x^2} \frac{\mu^2}{m_e^2} \right) - 2 \frac{1-x - x^2 \frac{m_e^2}{\mu^2}}{x} \right]$$

But protons are not point-like...

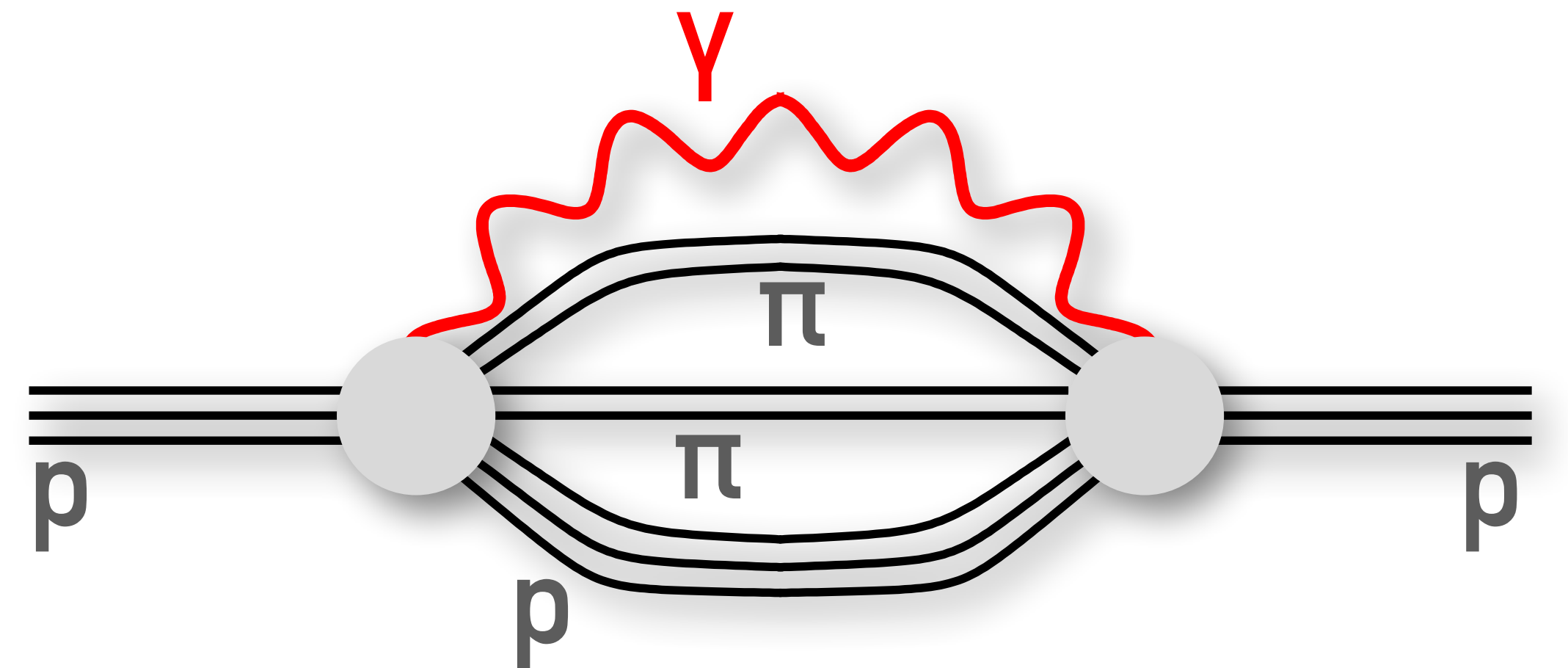
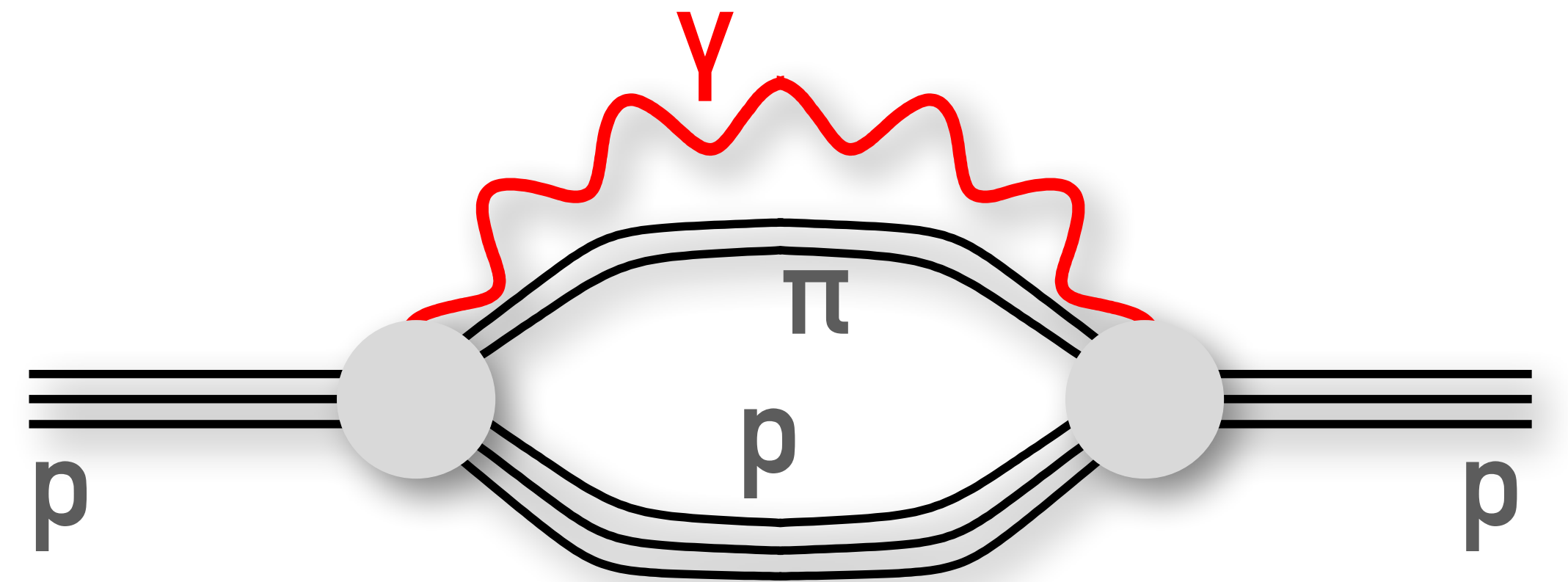
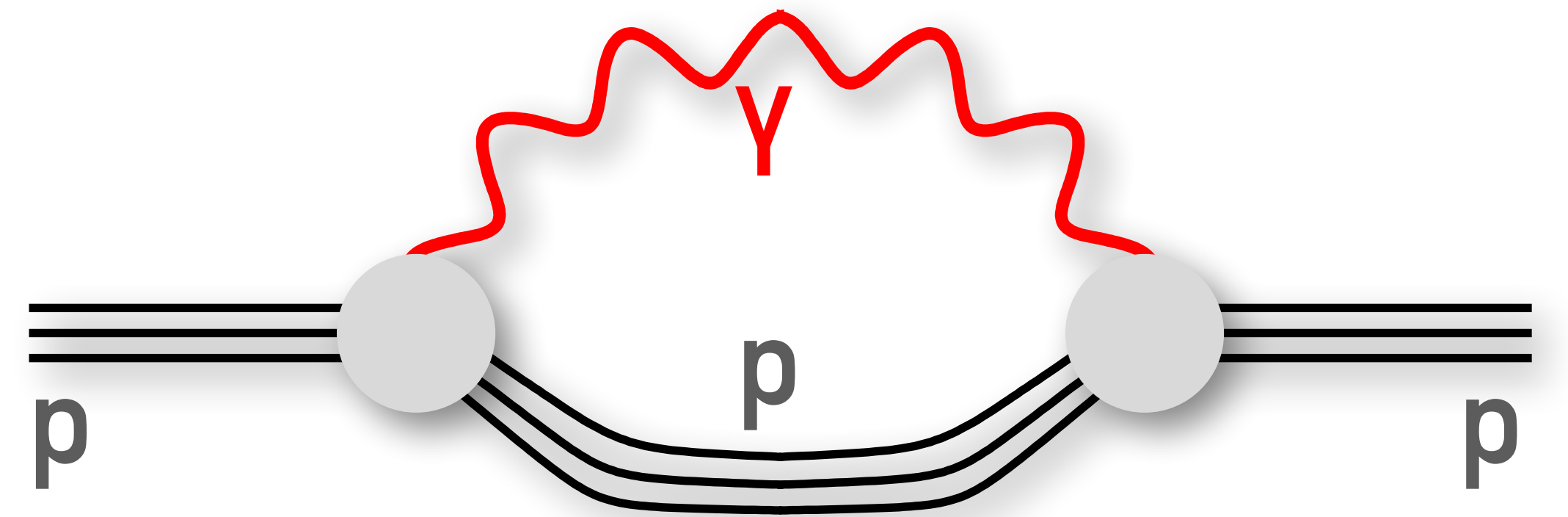
- Budnev, Ginzburg, Meledin & Serbo, Phys.Rept. 1974
→ an answer for the case where the proton remains intact after photon emission



given in terms of “proton form factors” (measurable from elastic ep scattering)

“number of photons” inside a proton?

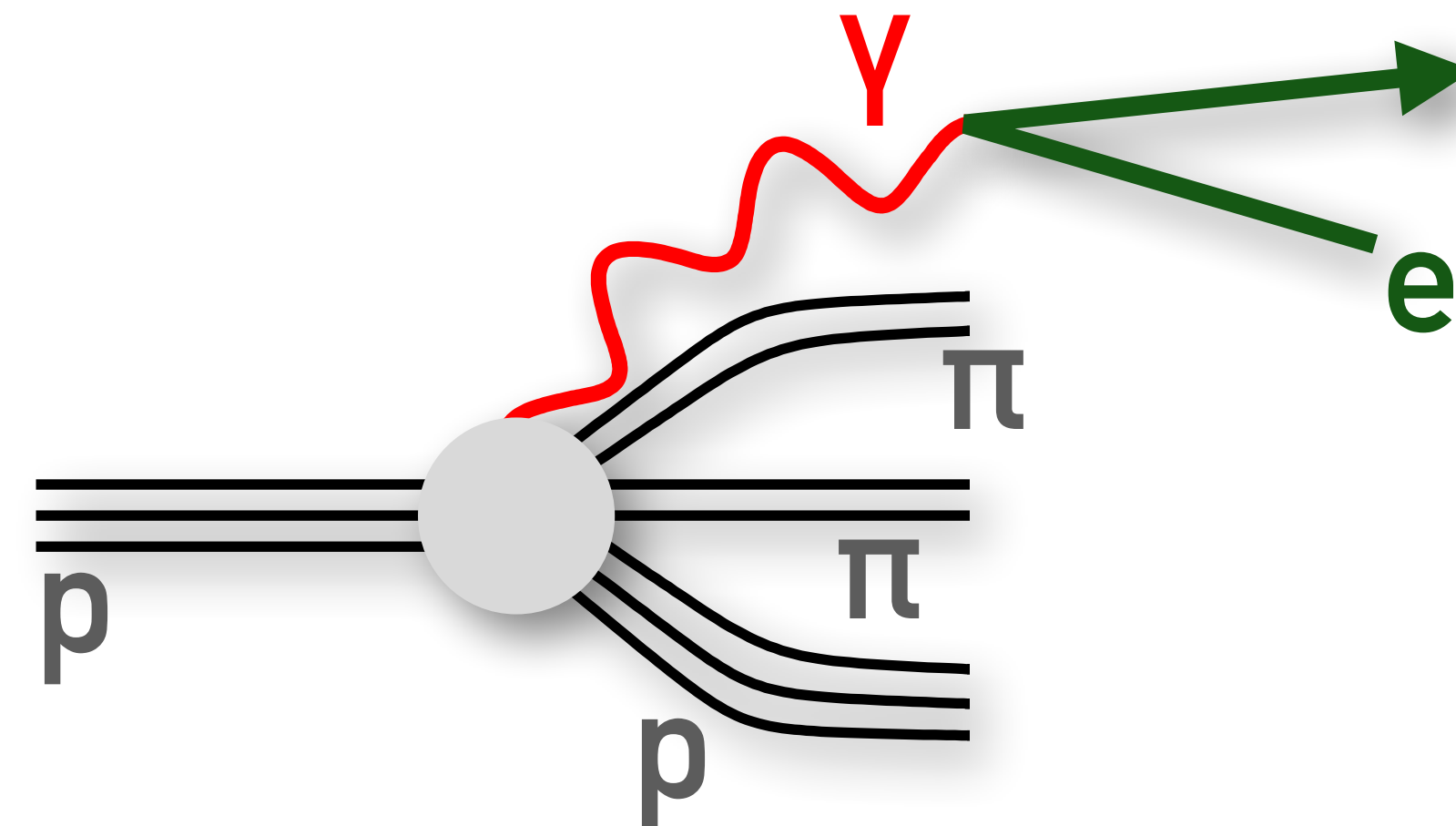
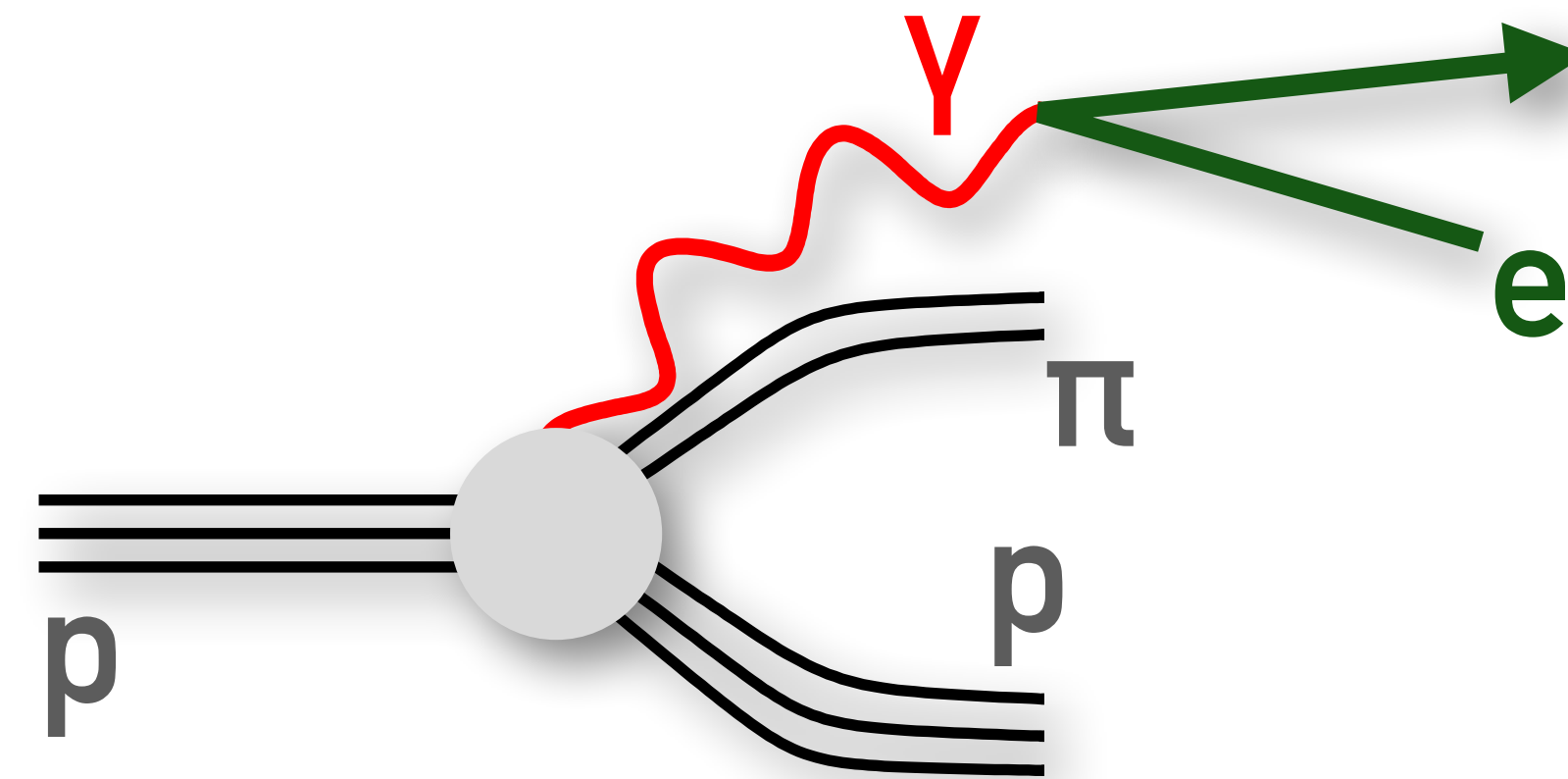
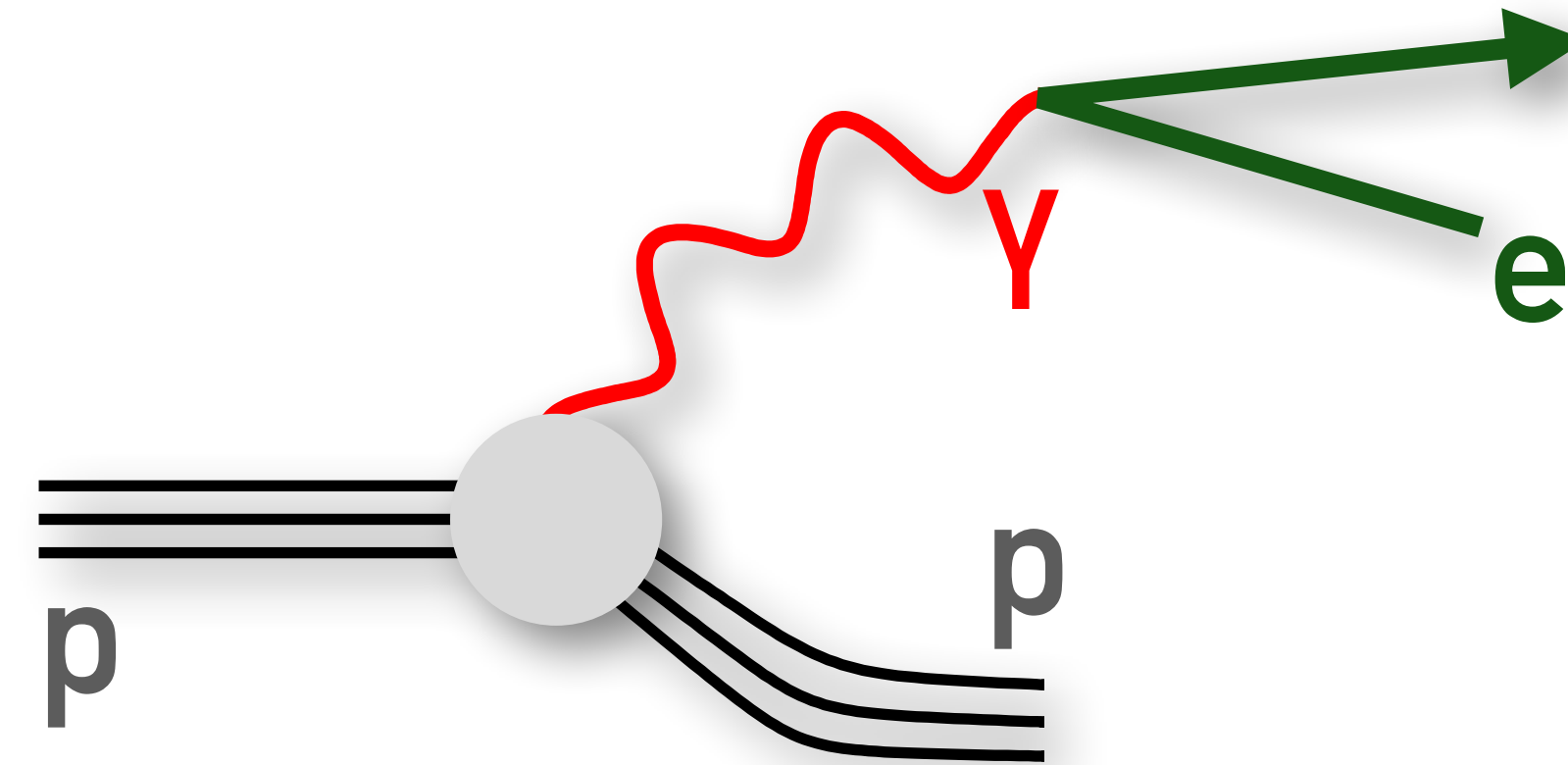
Proton constantly fluctuates in & out of different Fock states, some of which have a photon.



“number of photons” inside a proton?

Proton constantly fluctuates in & out of different Fock states, some of which have a photon.

If you absorb the γ , proton breaks up.



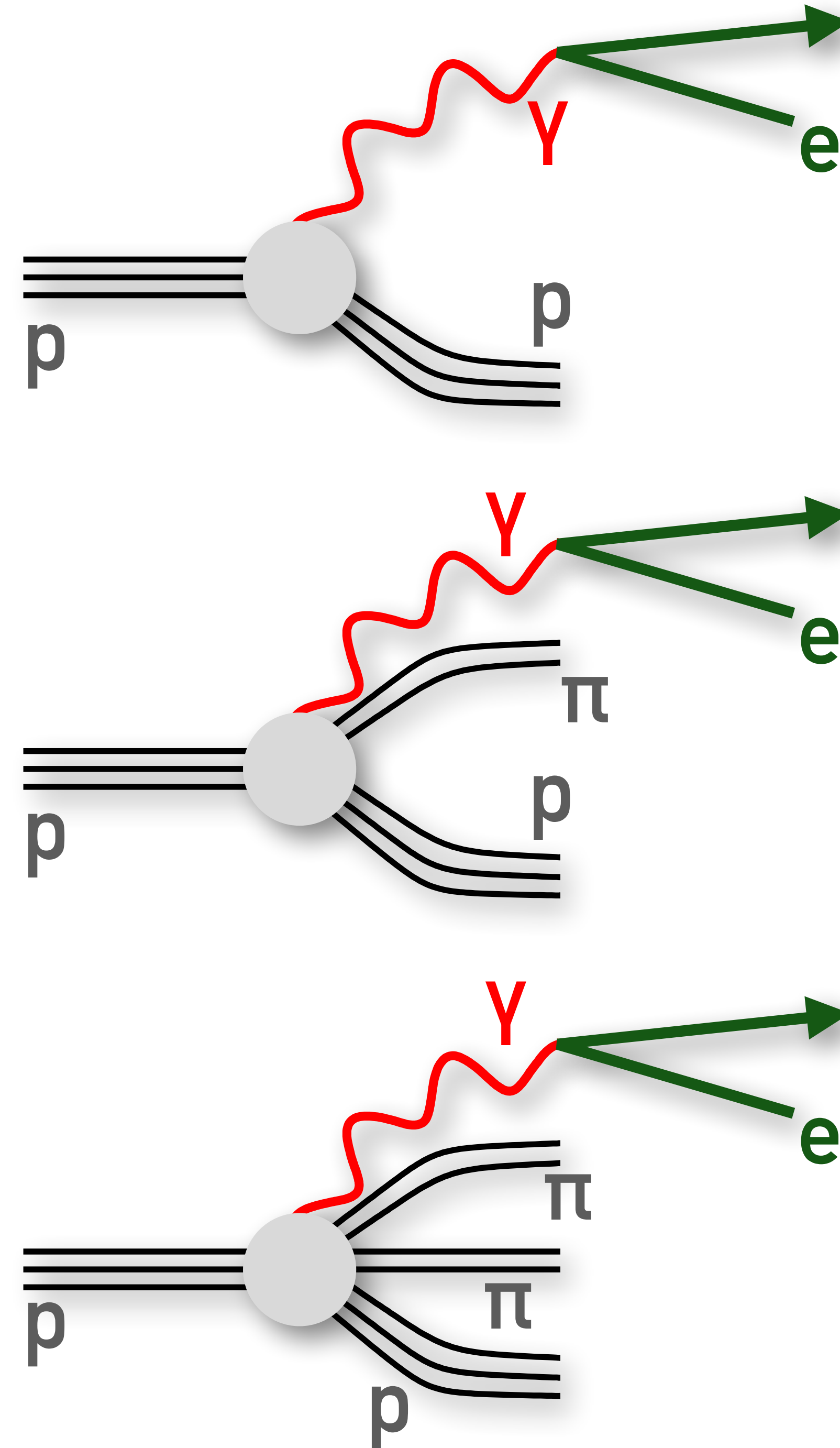
“number of photons” inside a proton?

Proton constantly fluctuates in & out of different Fock states, some of which have a photon.

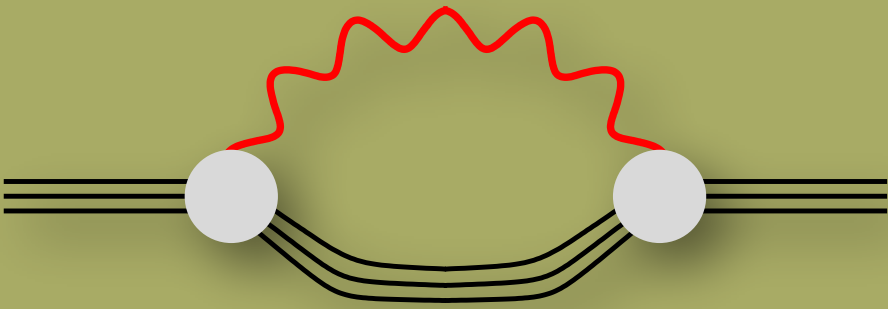
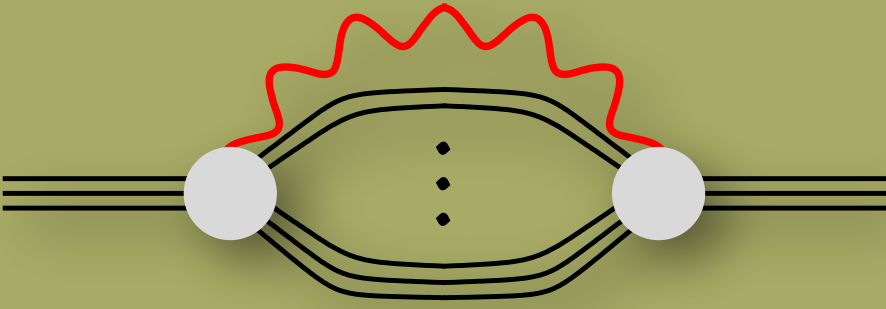
If you absorb the γ , proton breaks up.

Understanding this from first principles is a strong-coupling non-perturbative problem (beyond ability of lattice QCD)

Main approach in widely used γ determinations: **models.**

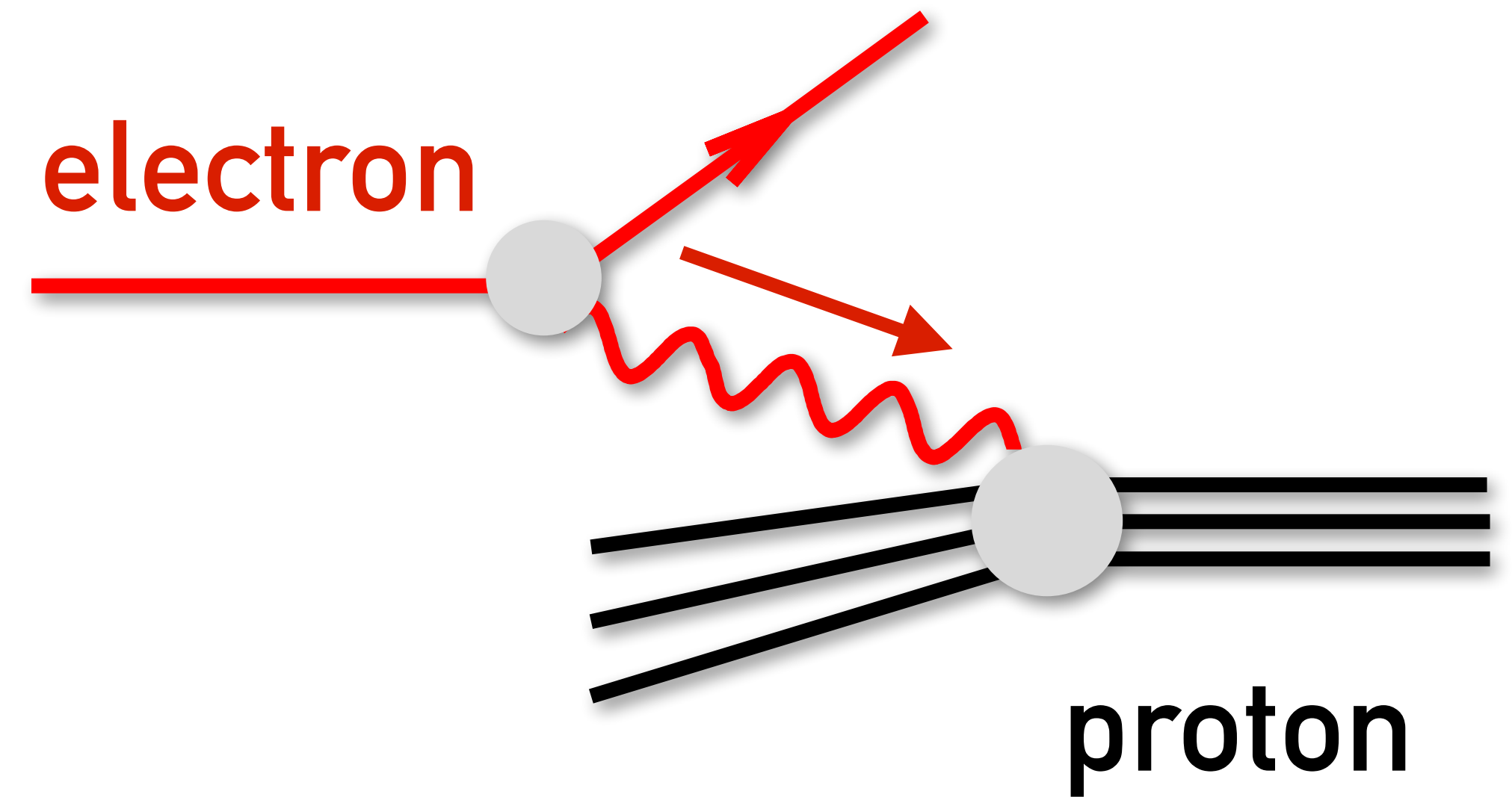


Widely discussed photon-PDF estimates

	 elastic	 inelastic	public computer-readable form?
Gluck Pisano Reya 2002	dipole	model	✗
MRST2004qed	✗	model	✓
CT14qed_inc	dipole	model (data-constrained)	✓
Martin Ryskin 2014	dipole (only electric part)	model	✗
Harland-Lang, Khoze Ryskin 2016	dipole	model	✗
NNPDF23qed	no separation; fit to data		✓

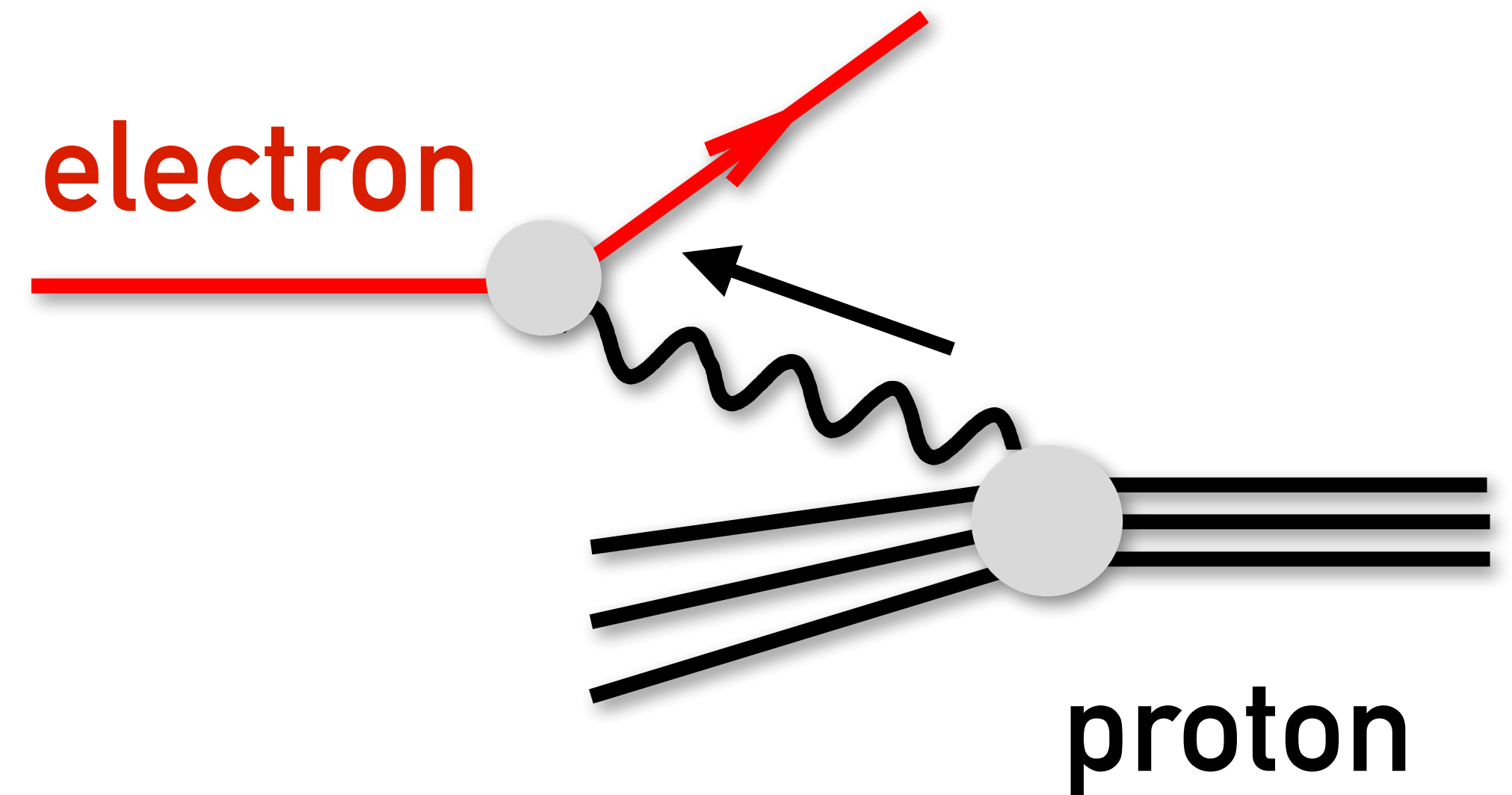
electron–proton scattering

- Experiments have been going on for decades
- Usually seen as photons from electron probing proton structure



electron–proton scattering

- ▶ Experiments have been going on for decades
- ▶ Usually seen as photons from electron probing proton structure
- ▶ But **can be viewed as electron probing proton's photonic field**
- ▶ Everything about electron–proton interaction encoded in two “structure functions” $F_2(x, Q^2)$ & $F_L(x, Q^2)$

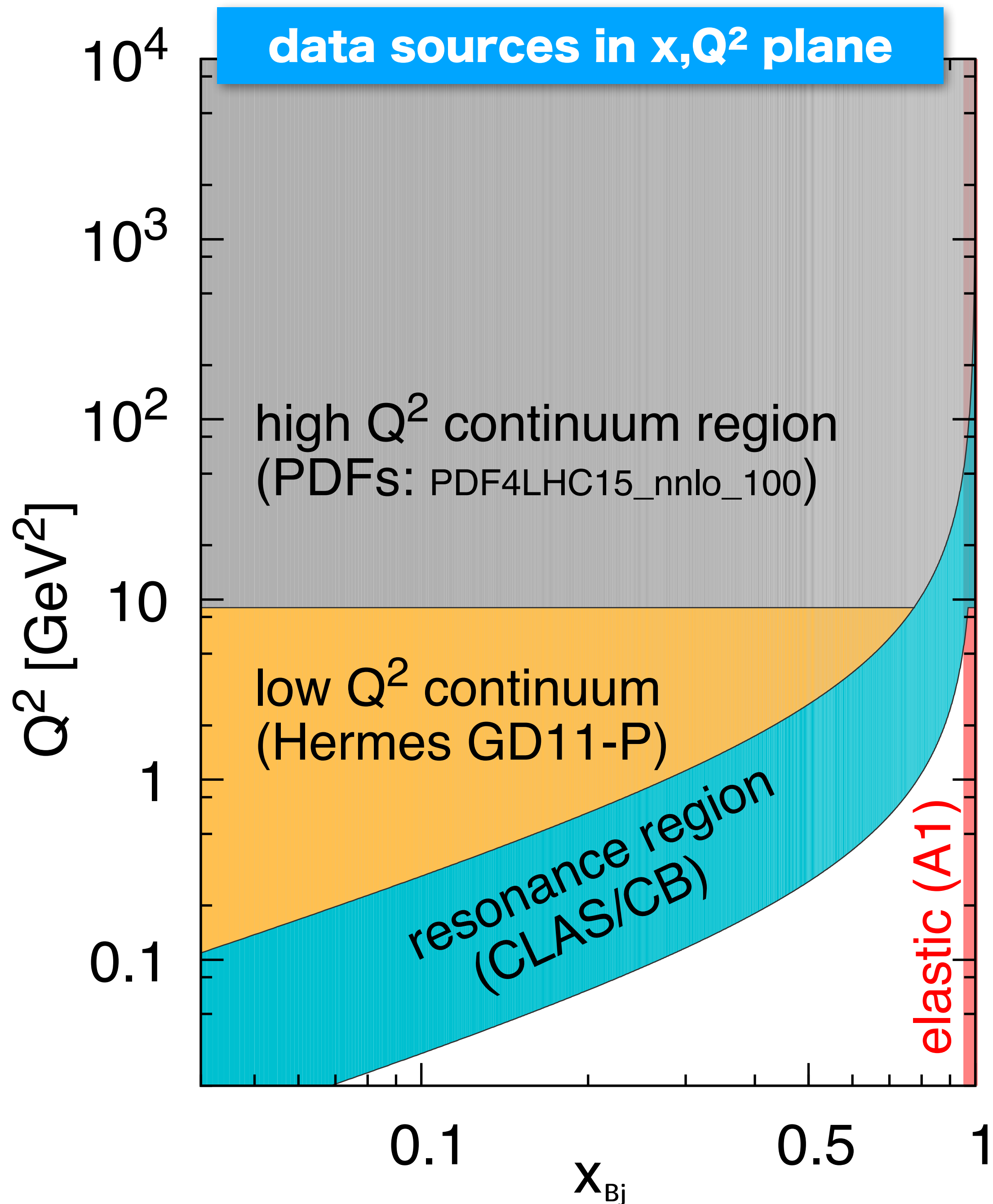


$$\frac{d\sigma}{dx dQ^2} = \frac{4\pi\alpha^2}{xQ^4} \left(\left(1 - y + \frac{y^2}{2} \left(1 + 2x^2 \frac{m_p^2}{Q^2} \right) \right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right)$$

Photon PDF in terms of F_2 and F_L — the **LUXqed** approach

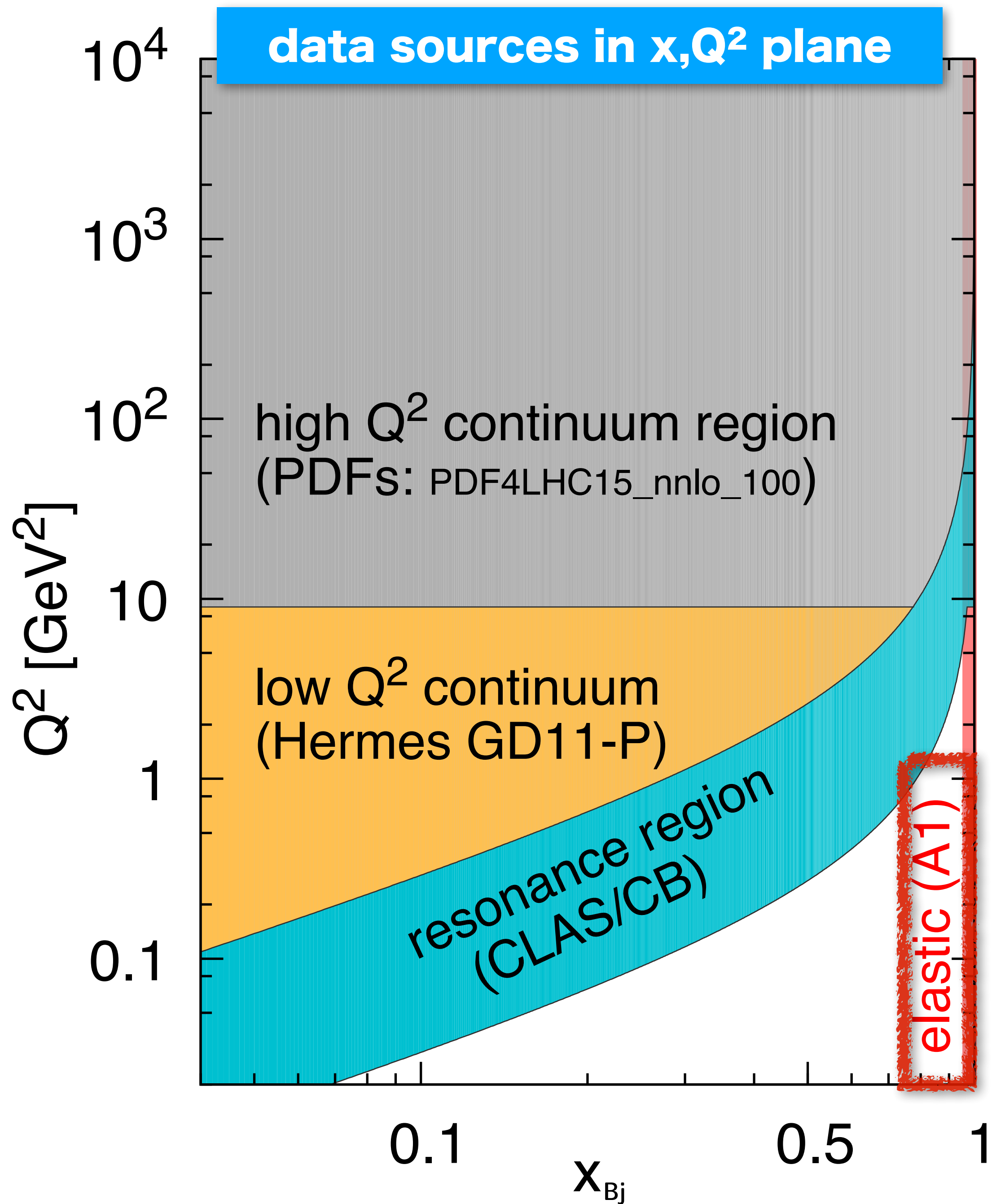
$$\begin{aligned}
 x f_{\gamma/p}(x, \mu^2) = & \frac{1}{2\pi\alpha(\mu^2)} \int_x^1 \frac{dz}{z} \left\{ \int_{\frac{x^2 m_p^2}{1-z}}^{\frac{\mu^2}{1-z}} \frac{dQ^2}{Q^2} \alpha^2(Q^2) \right. \\
 & \left[\left(z p_{\gamma q}(z) + \frac{2x^2 m_p^2}{Q^2} \right) F_2(x/z, Q^2) - z^2 F_L\left(\frac{x}{z}, Q^2\right) \right] \\
 & \left. - \alpha^2(\mu^2) z^2 F_2\left(\frac{x}{z}, \mu^2\right) \right\}
 \end{aligned}$$

It subsequently emerged that two “forgotten” papers, Anlauf et. al, CPC70(1992)97 Mukherjee & Pisano, hep-ph/0306275, had the correct integrand (but not the limits)

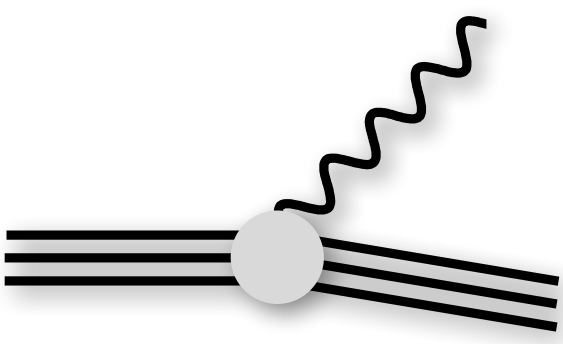


DATA

- x, Q^2 plane naturally breaks up into regions with different physical behaviours and data sources
- We don't use F_2 and F_L data directly, but rather various fits to data



ELASTIC COMPONENT

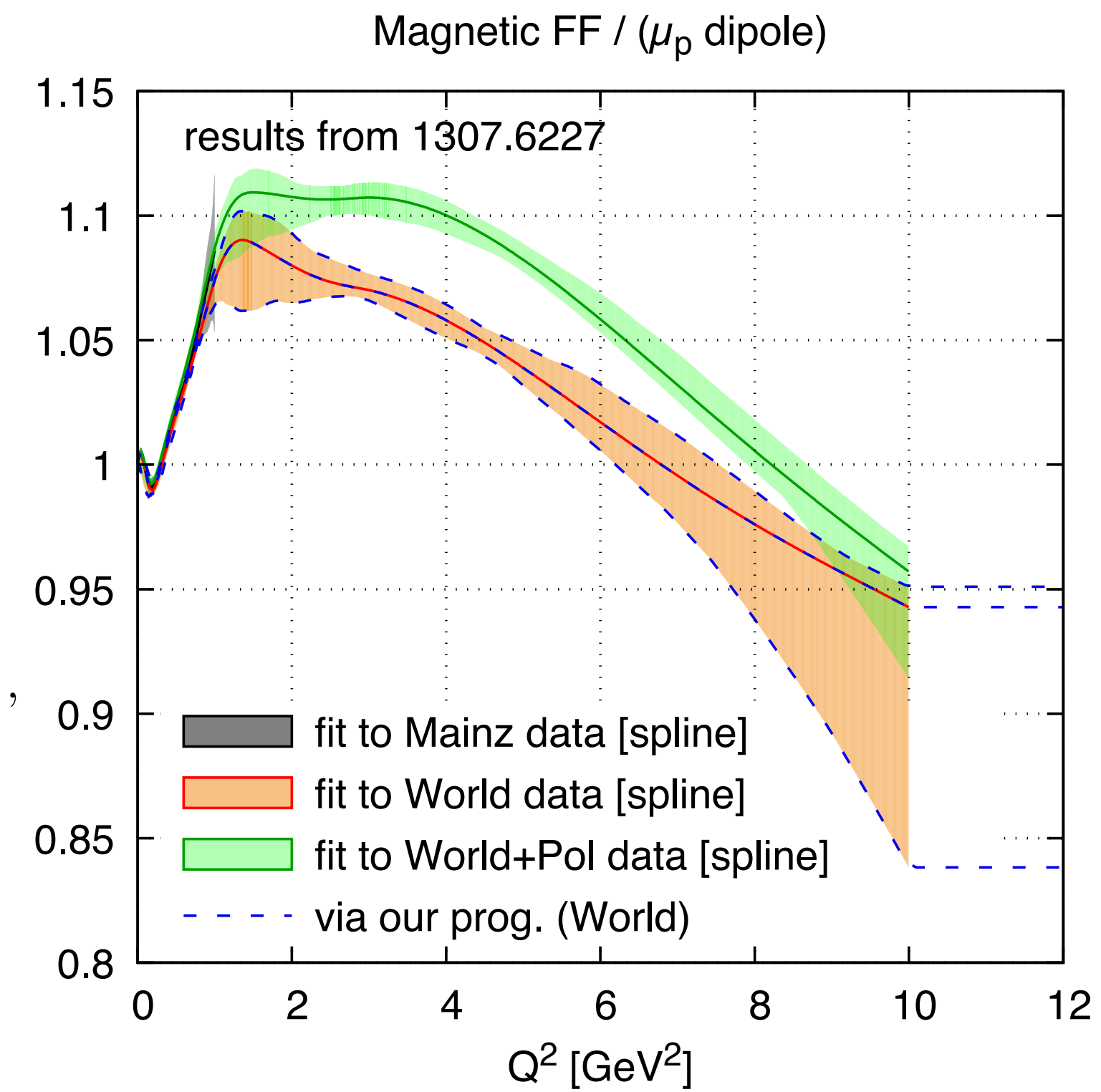


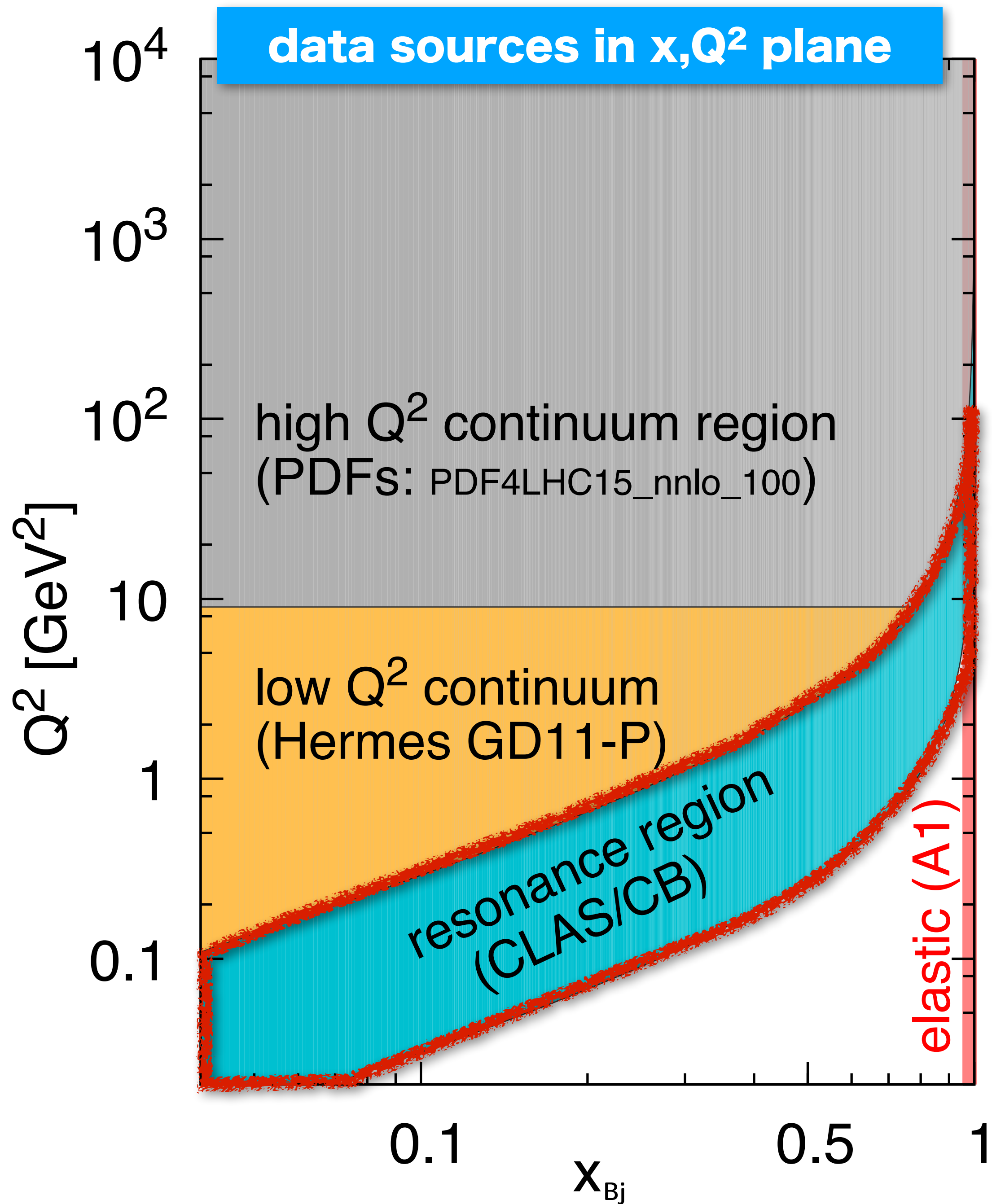
- Elastic component of $F_{2/L}$ lives at $x = 1$
- Express in terms of Sachs Form factors

$$\tau = Q^2 / (4m_p^2)$$

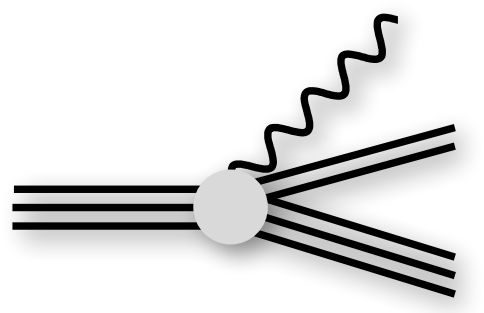
$$F_2^{\text{el}} = \frac{[G_E(Q^2)]^2 + [G_M(Q^2)]^2 \tau}{1 + \tau} \delta(1 - x),$$

$$F_L^{\text{el}} = \frac{[G_E(Q^2)]^2}{\tau} \delta(1 - x),$$

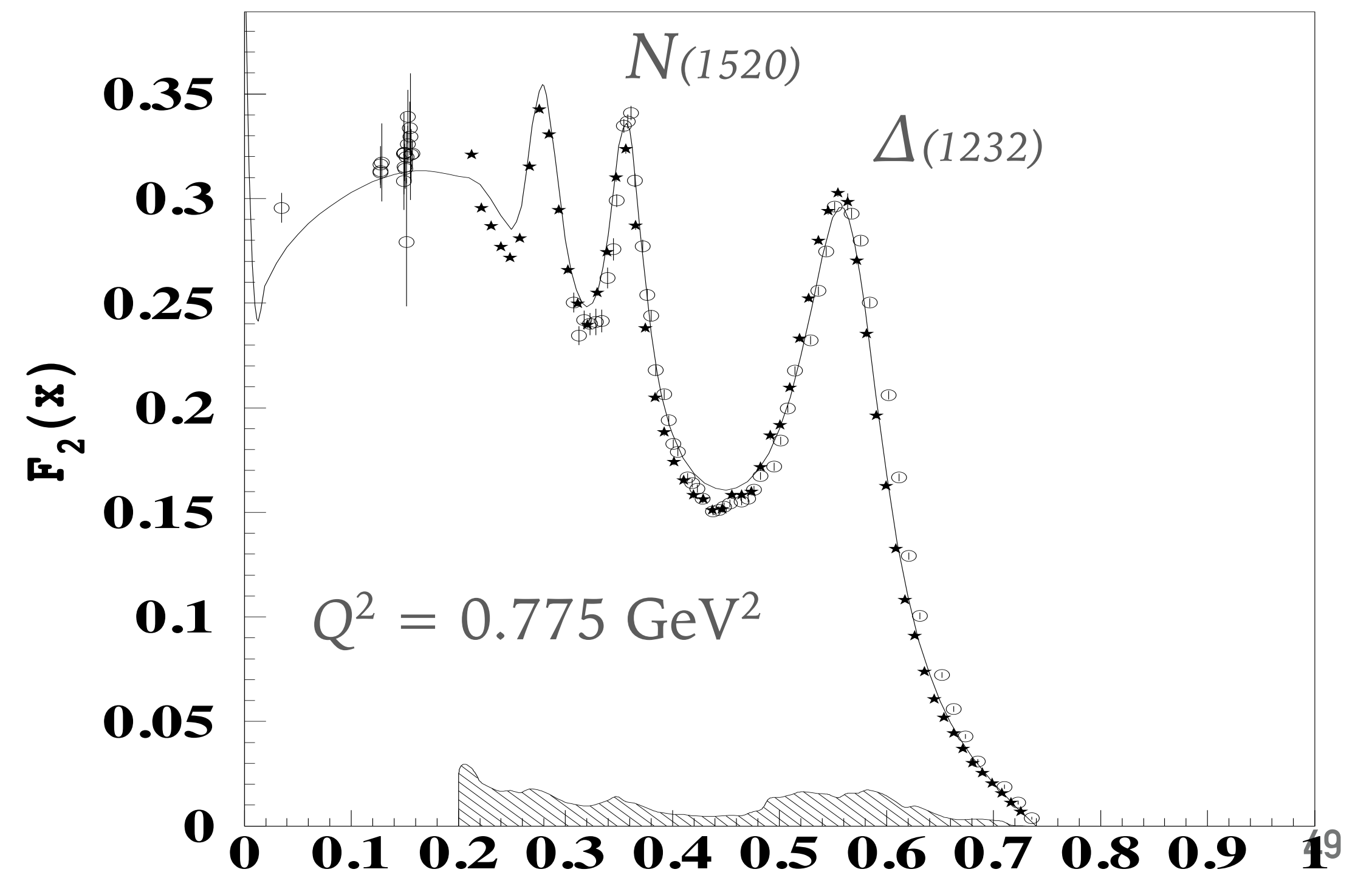


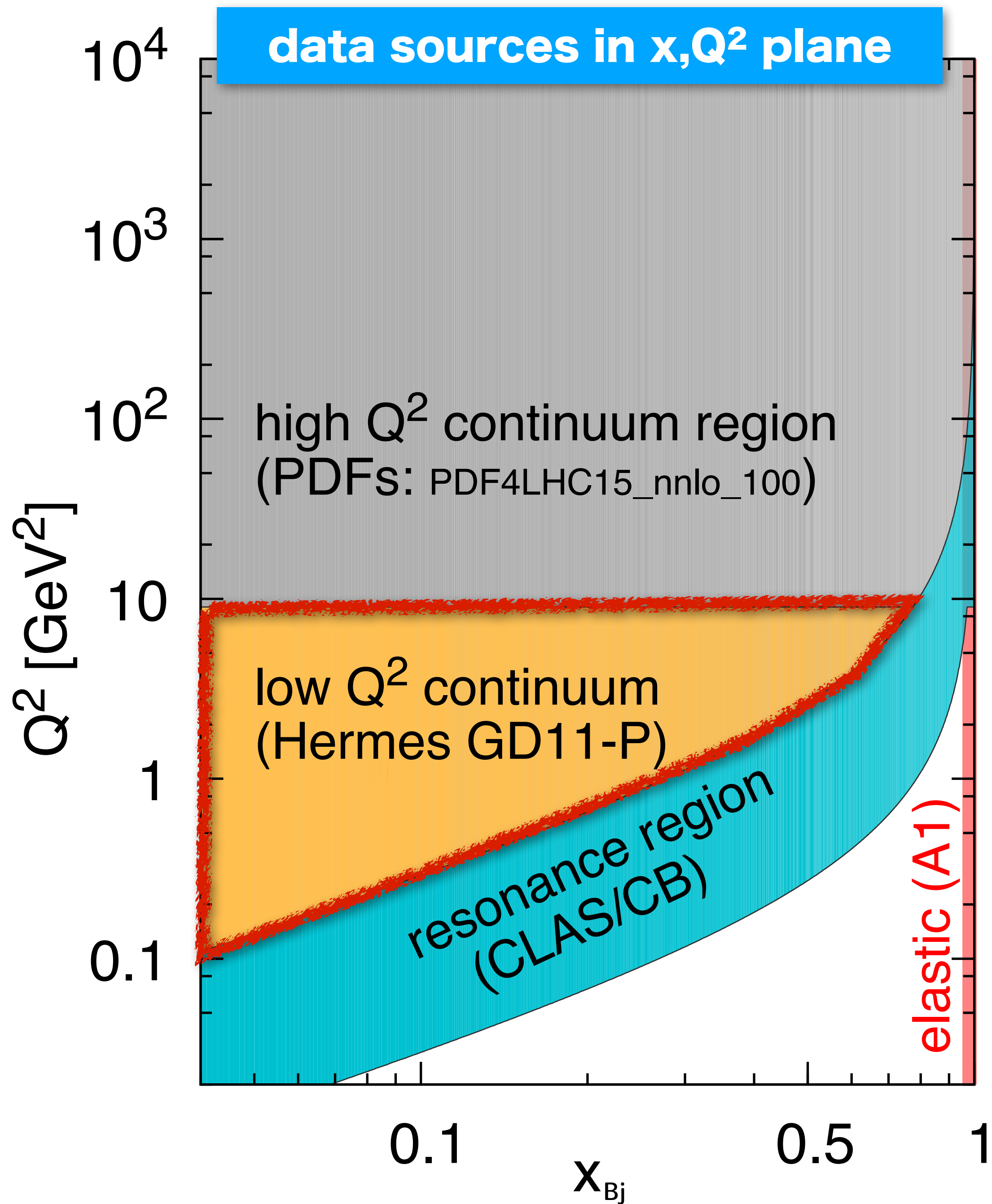


RESONANCE COMPONENT

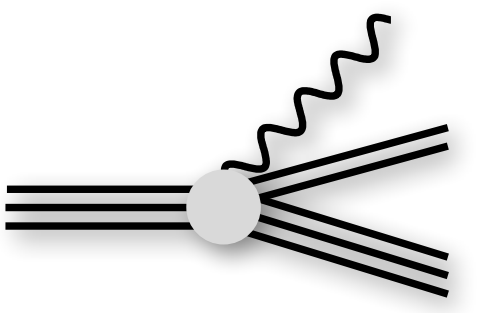


- proton gets excited, e.g. to $\Delta \rightarrow p\pi$ and higher resonances
- relevant for $(m_p + m_\pi)^2 < W^2 < 3.5 \text{ GeV}^2$

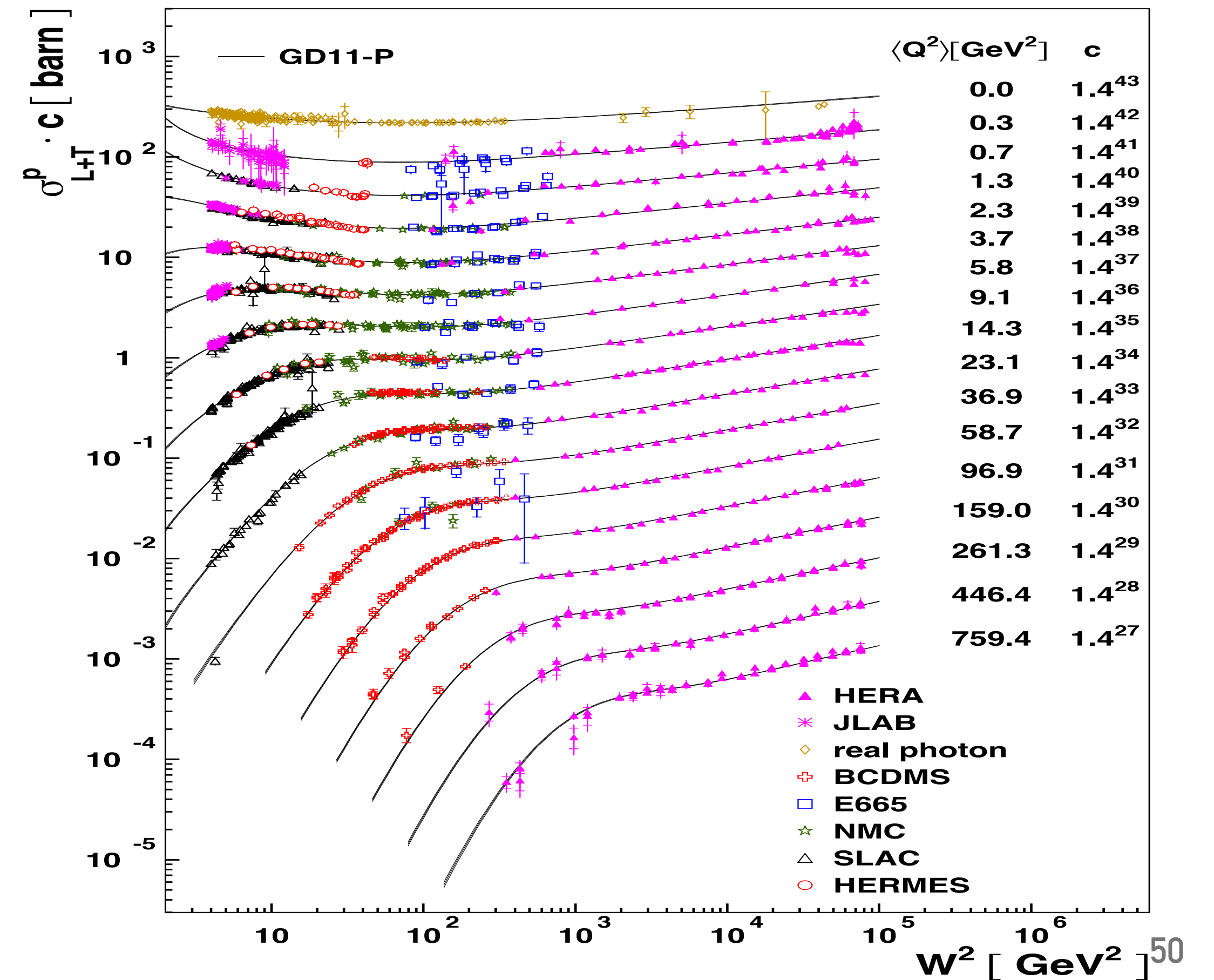


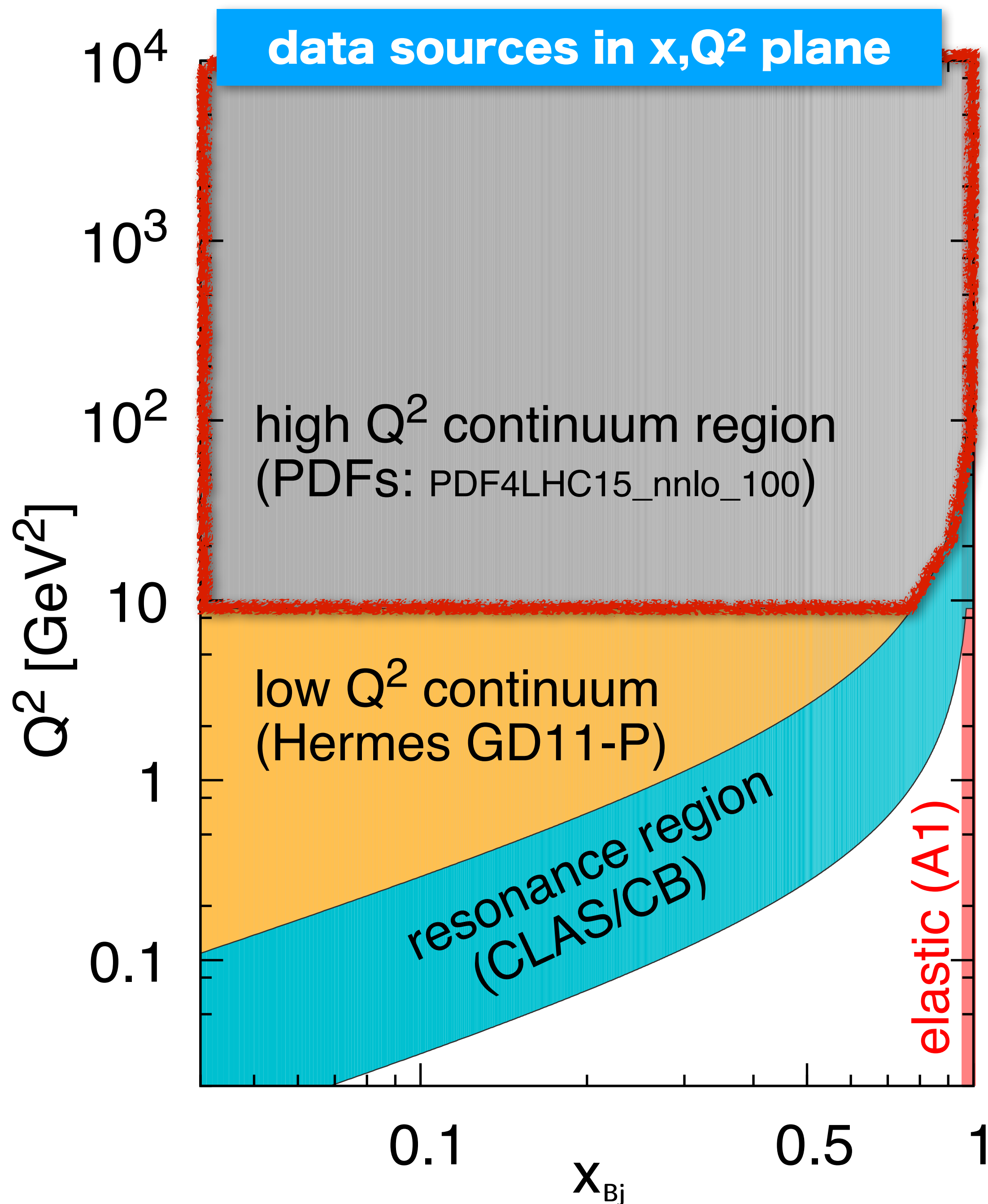


CONTINUUM COMPONENT

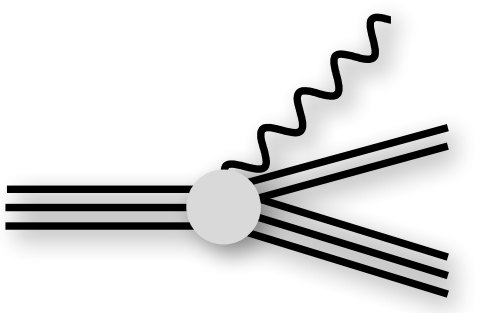


- Much data
- For $Q^2 \rightarrow 0$, σ_{yp} indep. of Q^2 at fixed W^2





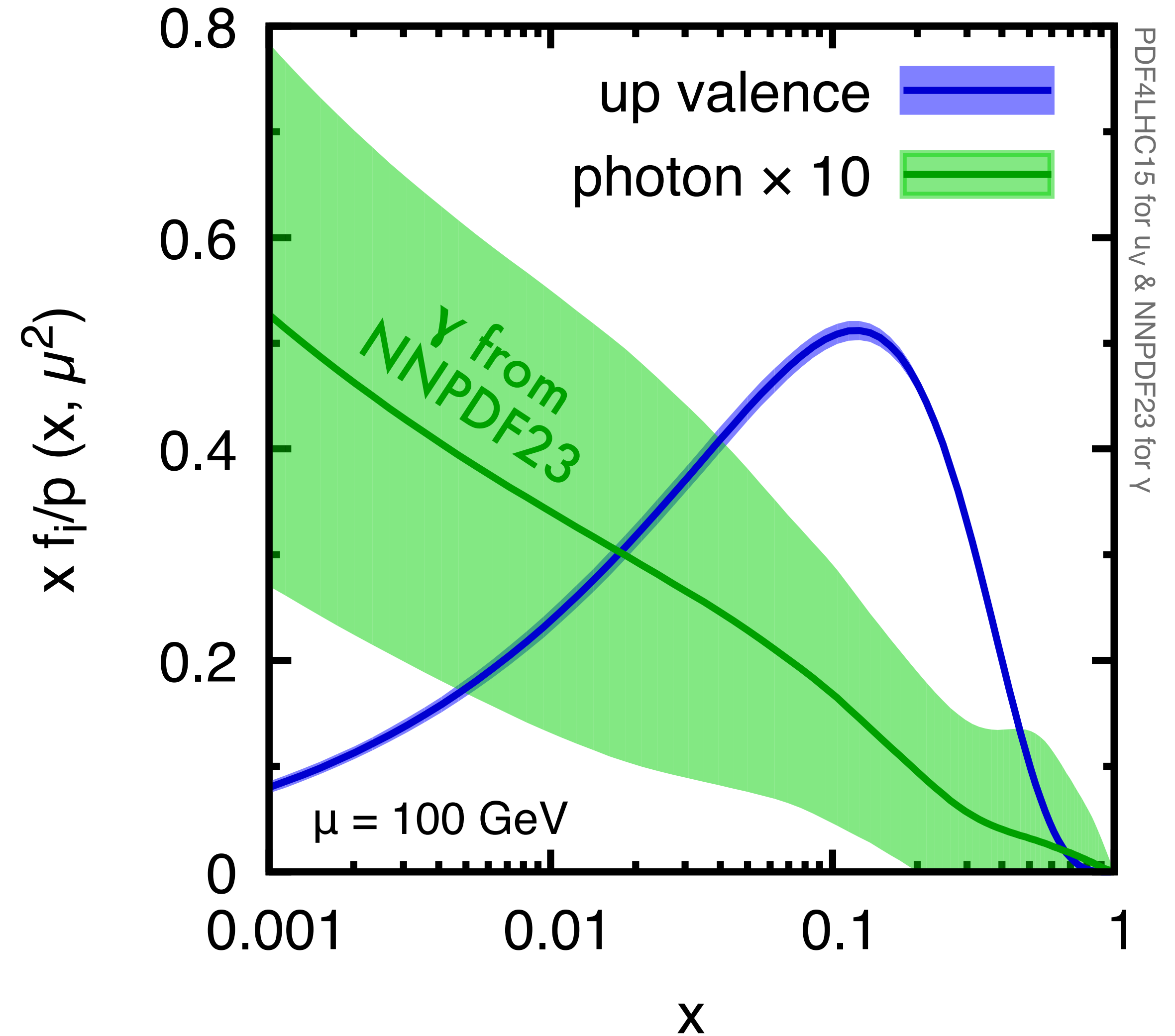
CONTINUUM COMPONENT



- ▶ Less direct data for F_2 and F_L at high Q^2
- ▶ But we can reliably use PDFs and coefficient functions (up to NNLO) to calculate them
- ▶ Our default choice is PDF4LHC15_nnlo_100 (and zero-mass variable flavour-number scheme)

photon PDF results

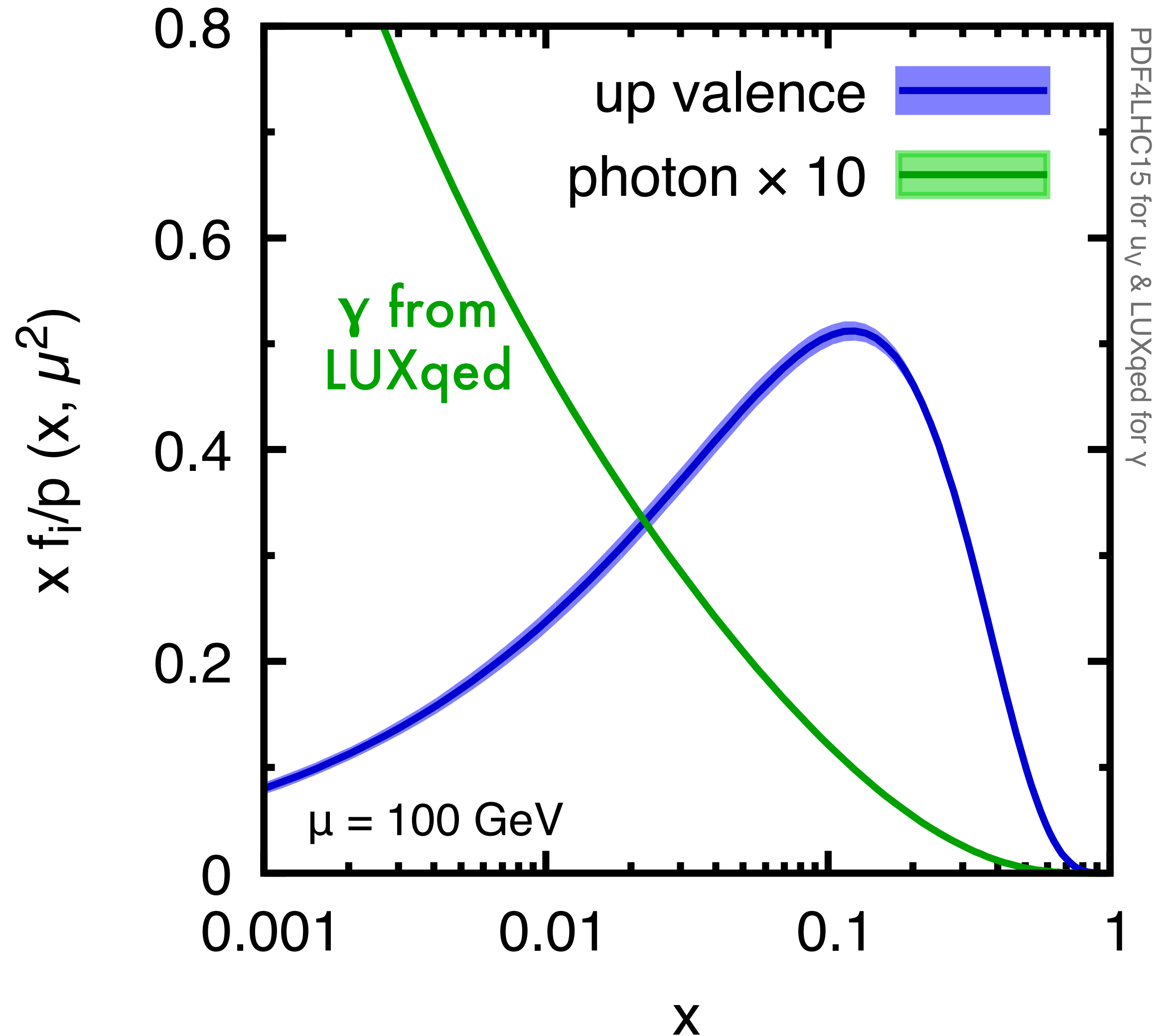
- Model-independent uncertainty (NNPDF) was 50–100%



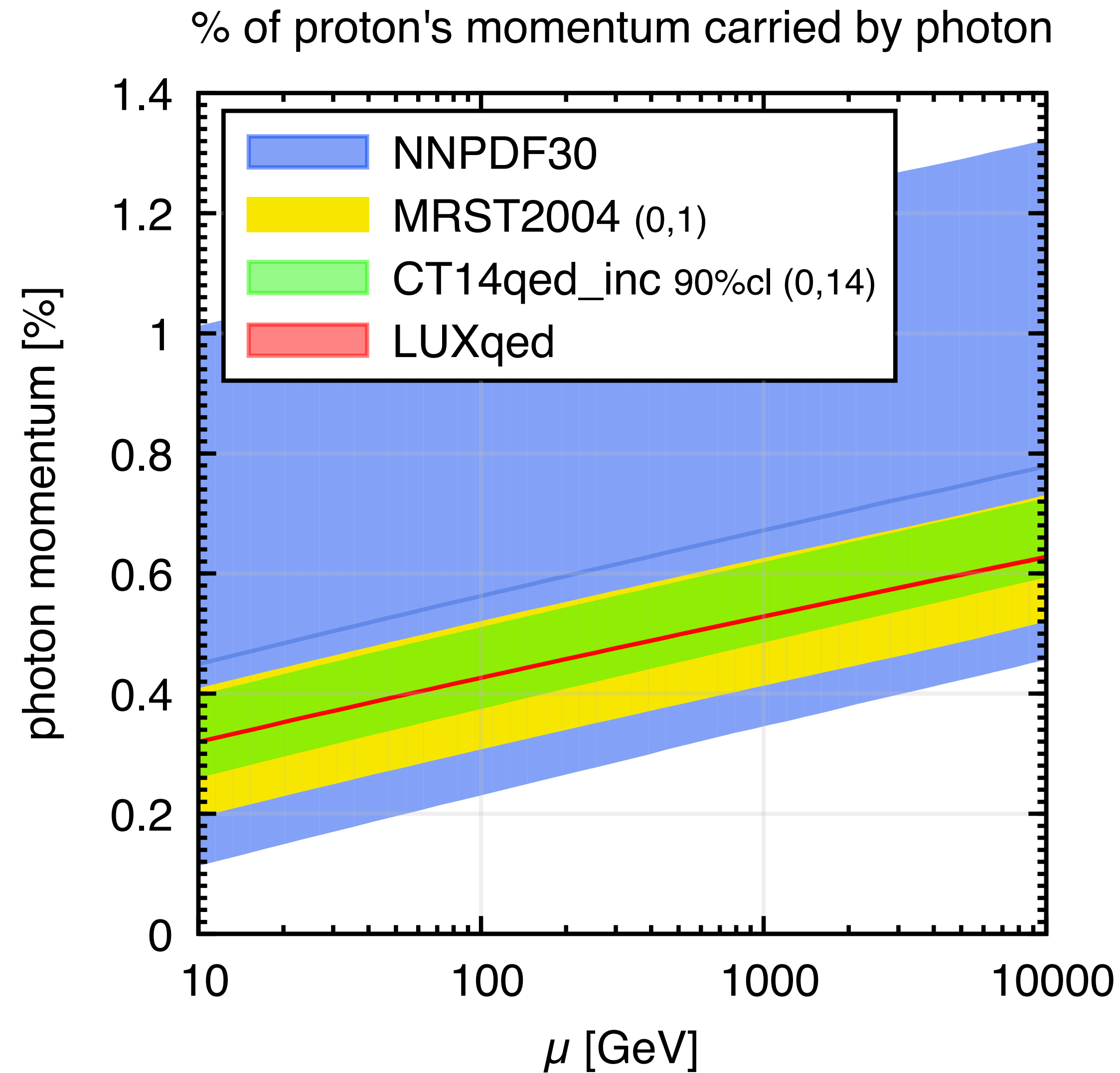
photon PDF results

- Model-independent uncertainty (NNPDF) was 50–100%
- Goes down to $O(1\%)$ with LUXqed determination

pp \rightarrow H W ⁺ (\rightarrow l ⁺ v) + X at 13 TeV	
non-photon induced contributions	91.2 ± 1.8 fb
photon-induced contriibs (NNPDF23)	$6.0^{+4.4}_{-2.9}$ fb
photon-induced contriibs (LUXqed)	4.4 ± 0.1 fb



How much light is there in the proton? [Momentum fraction]



momentum ($\mu = 100$ GeV)	
gluon	$46.8 \pm 0.4\%$
up valence	$18.2 \pm 0.3\%$
down valence	$7.5 \pm 0.2\%$
light sea quarks	$20.7 \pm 0.4\%$
charm	$4.0 \pm 0.1\%$
bottom	$2.5 \pm 0.1\%$
photon	$0.426 \pm 0.003\%$

LUXqed_plus_PDF4LHC15_nnlo_100

(1+107 members, symmhessian, errors handled by LHAPDF out of the box,

valid for $\mu > 10$ GeV)

CONCLUSIONS

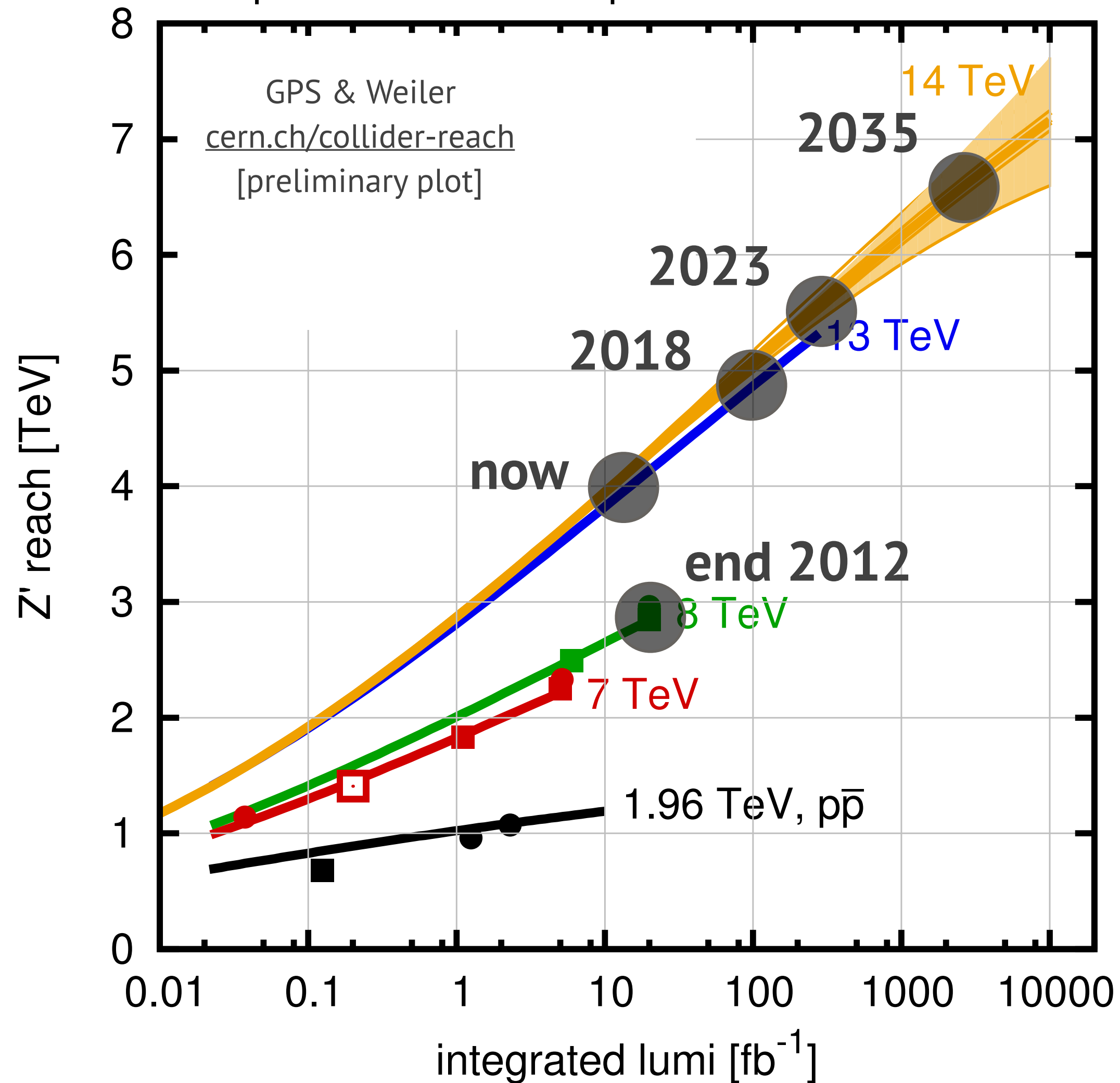
Summary

- The LHC has a rich programme in the years ahead to **establish** large parts of the **Higgs sector** of the Standard Model
- Extracting **qualitatively new information**, depends on **quantitative precision** in understanding proton–proton collisions
- New ways of thinking about precision LHC physics can bring big payoffs (and solve long-standing basic physics problems, e.g. **γ PDF**)
- Search for New Physics (BSM) continues, aided by the progress

EXTRA SLIDES

Discovery potential: (now → HL-LHC) > (run I → now)

Z' exclusion reach v. lumi



Today

- 20 fb⁻¹ @ 8 TeV
- 13 fb⁻¹ @ 13 TeV (results)

Future

- 2018: 100 fb⁻¹ @ 13 TeV
- 2023: 300 fb⁻¹ @ 1? TeV
- 2035: 3000 fb⁻¹ @ 14 TeV

1 fb⁻¹ = 10¹⁴ collisions

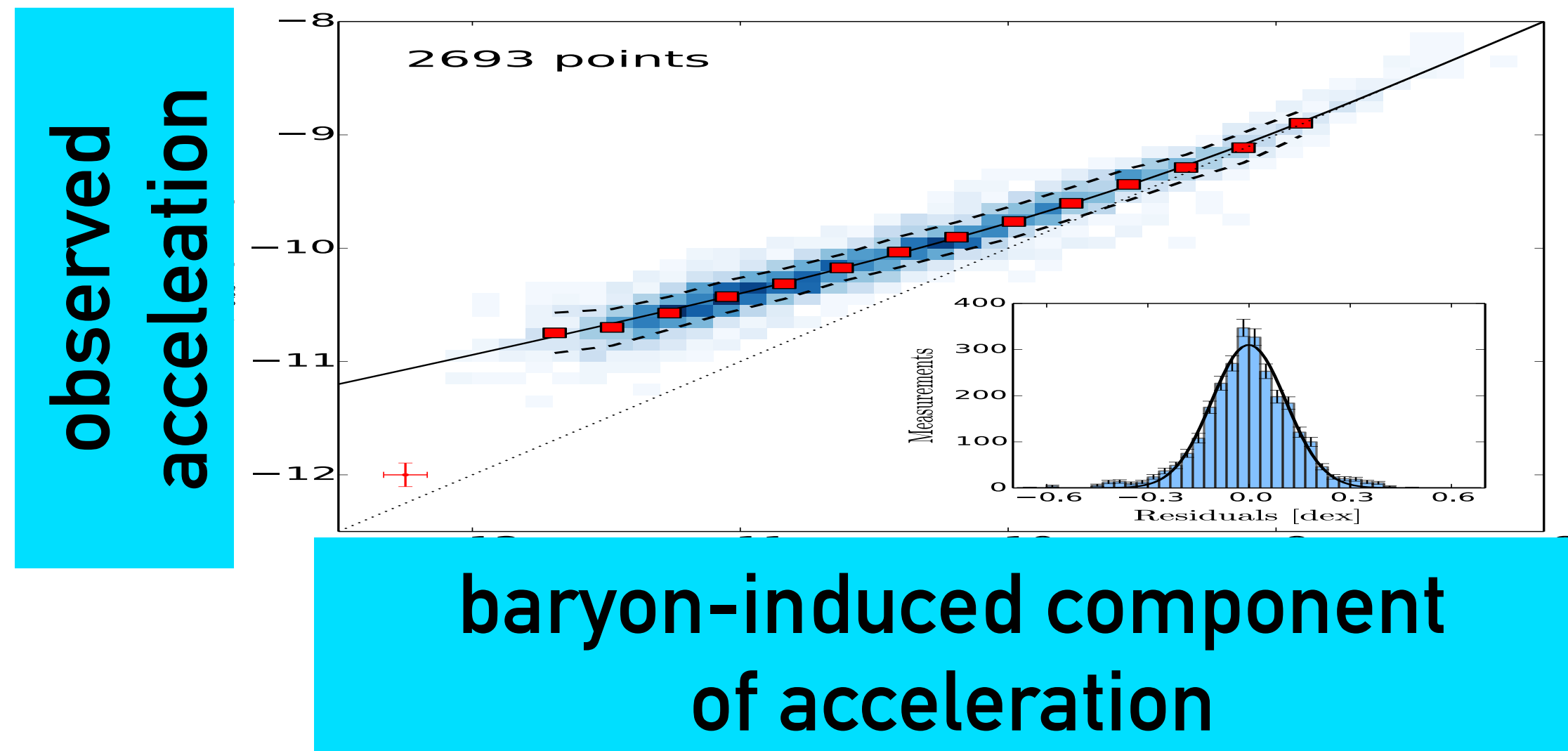
But not clear that dark matter is “standard” WIMP-like

[...] **Standard cosmological**
[...] simulations with] **dark**
matter halos [...] **do not**
naturally lead to realistic
galaxies [44, 46].

Complicated [...] “feedback”
must be invoked [...]

Whether such processes can
satisfactorily explain the
radial acceleration relation
and its small scatter remains
to be demonstrated [47, 48].

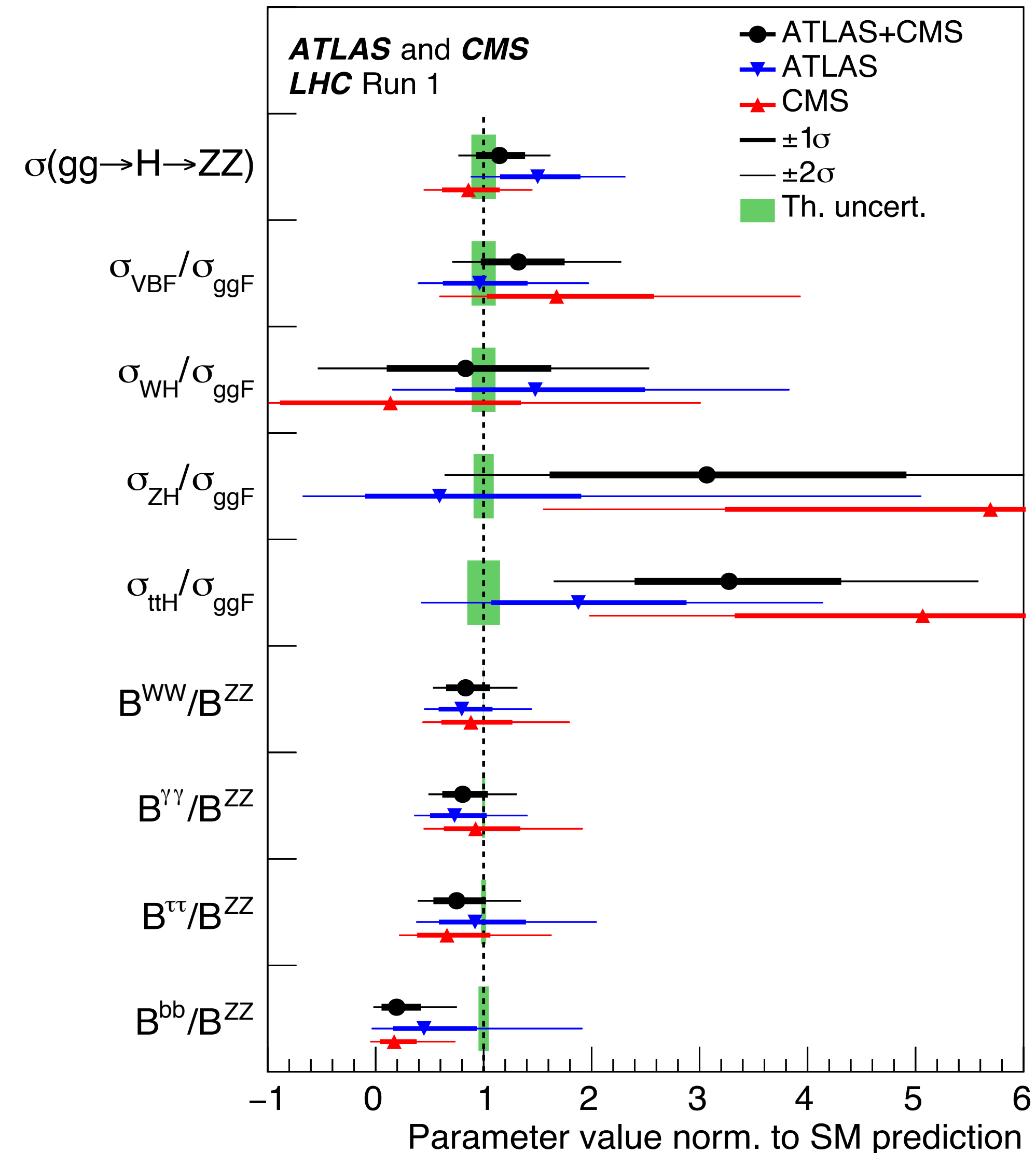
PRL117, 201101 (2016)



HIGGS

now & future

What do we know today? Broad picture looks standard-model like



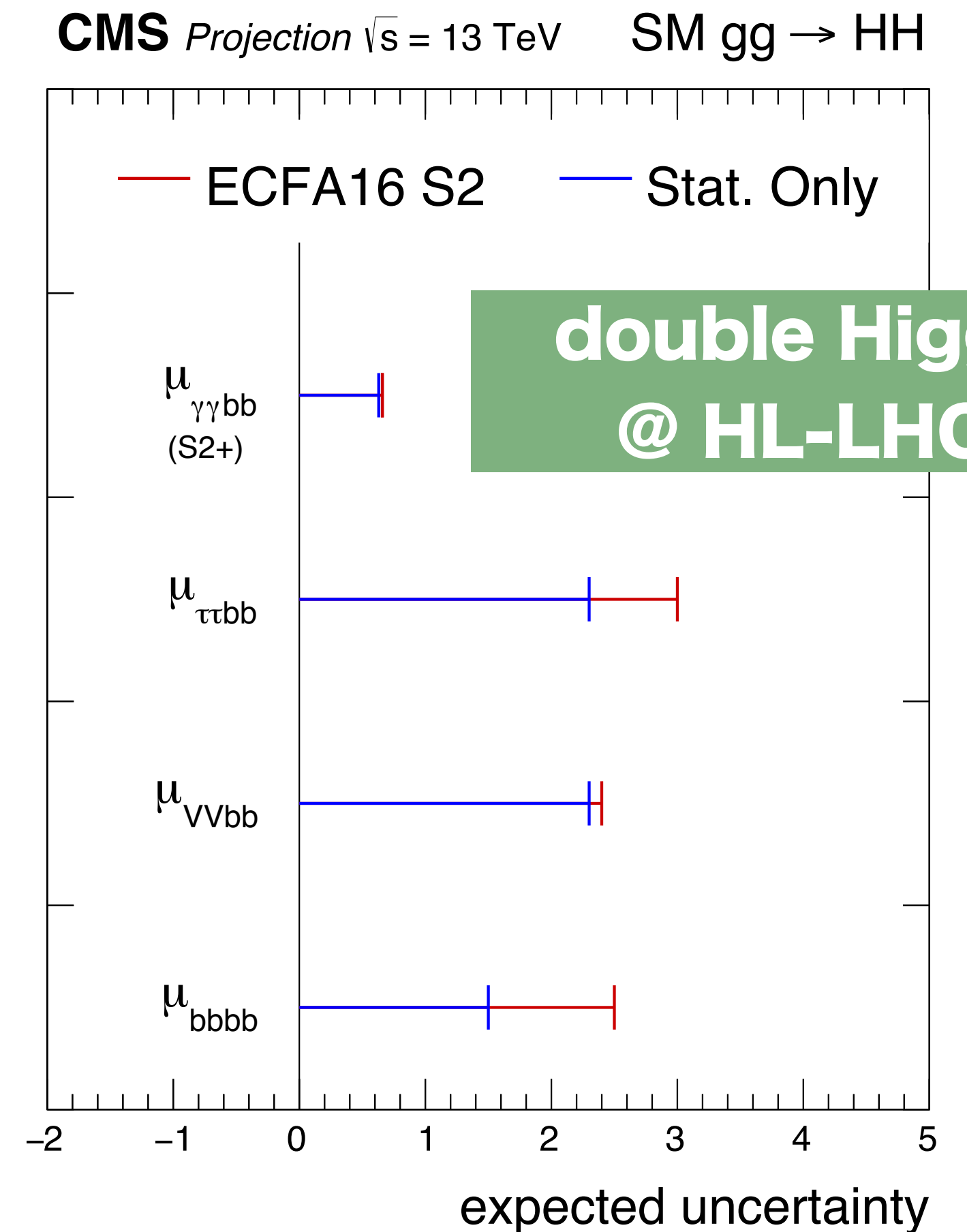
Coupling to electroweak and 3rd generation looks standard

- ▶ we see expected rate of decays to ZZ and WW (and some evidence of VBF/VH)
- ▶ observation consistent with $\sigma(\text{gluon fusion})$ means top-coupling is probably standard
- ▶ fact that all cross sections look right also means b-coupling is probably standard (because it dominates in denominator of branching ratios)
- ▶ reasonable evidence that coupling to tau is standard (direct observation)

To see the data, as is, with very non-standard (t,b, τ ,W,Z) couplings would require some degree of conspiracy.

WHAT WILL THE LHC BRING?

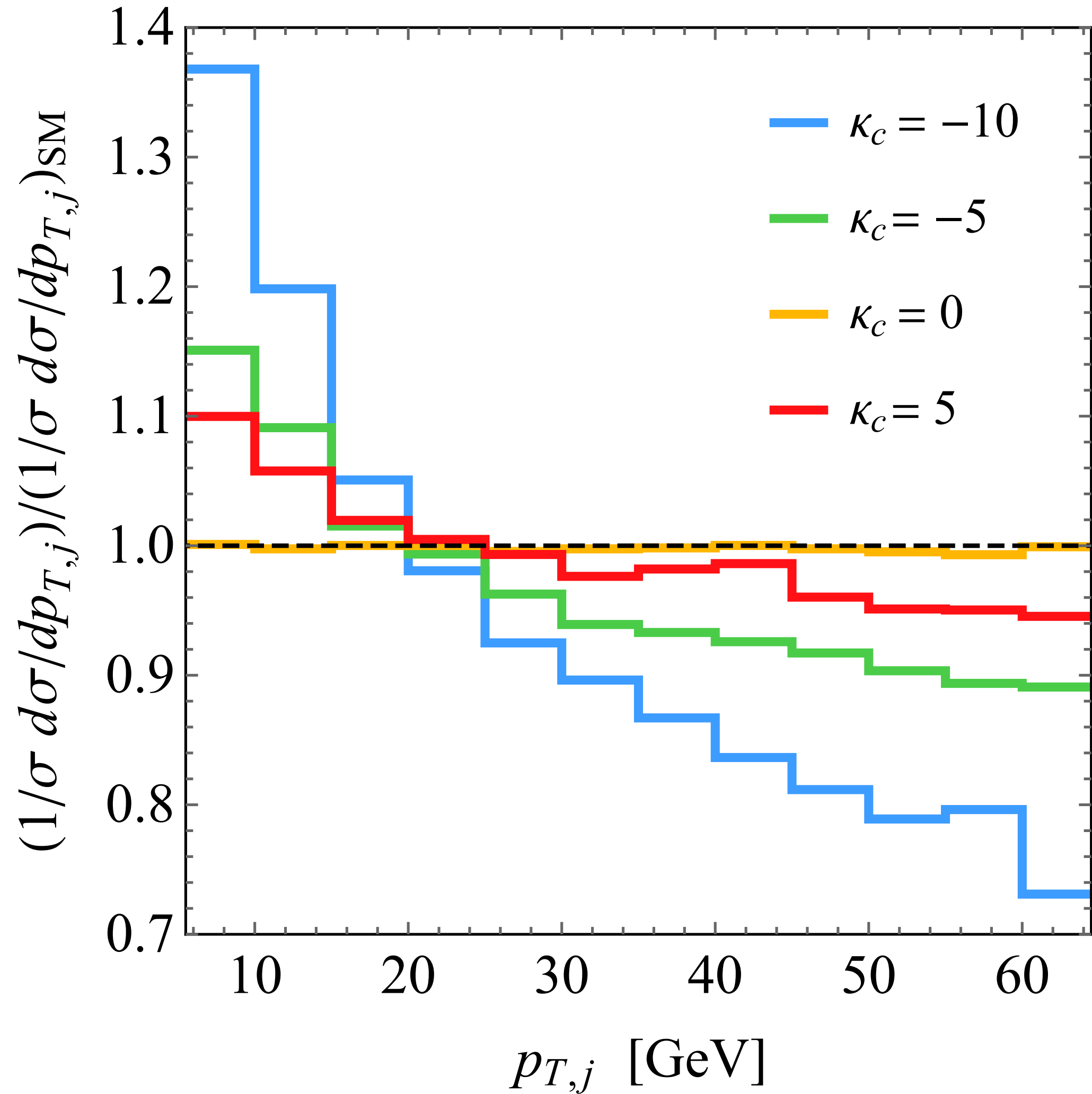
- Run 2: observation of $H \rightarrow bb$ (Yukawa)
- Run 2/3: observation of ttH (Yukawa)
- HL-LHC: observation of $H \rightarrow \mu\mu$ (2nd gen Yukawa)
- HL-LHC: Higgs width \rightarrow SM $\pm 50\%$ (BSM constraint)
- HL-LHC: $H \rightarrow$ invisible $< 10\%$ (BSM constraint)
- HL-LHC: $gg \rightarrow HH?$ (Higgs potential)
- HL-LHC: Hcc coupling? (2nd gen Yukawa)



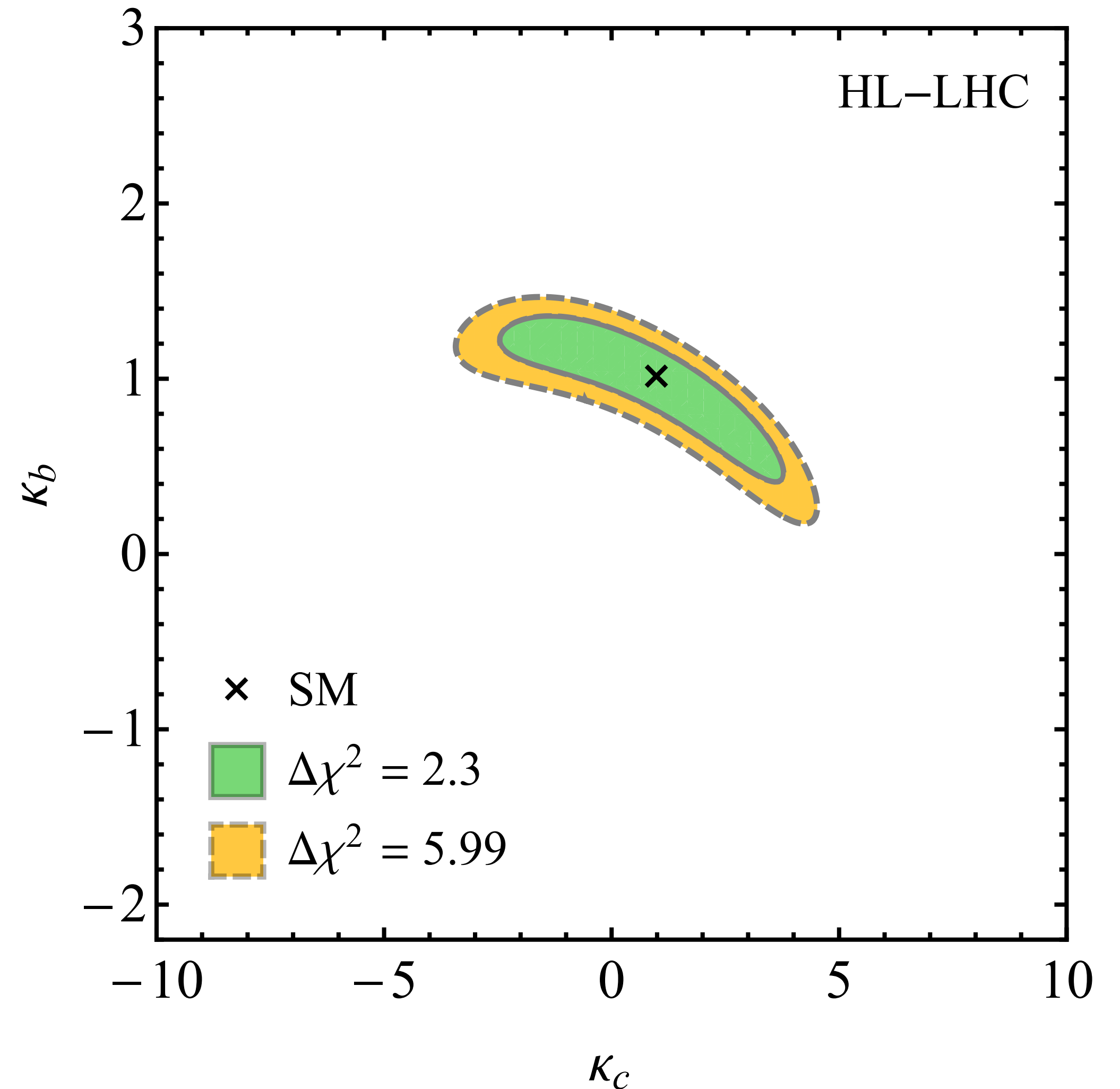
cf. talks at HL-LHC workshop

indirect constraints on Hcc

impact of modified Hcc coupling on Higgs+jet p_T dist

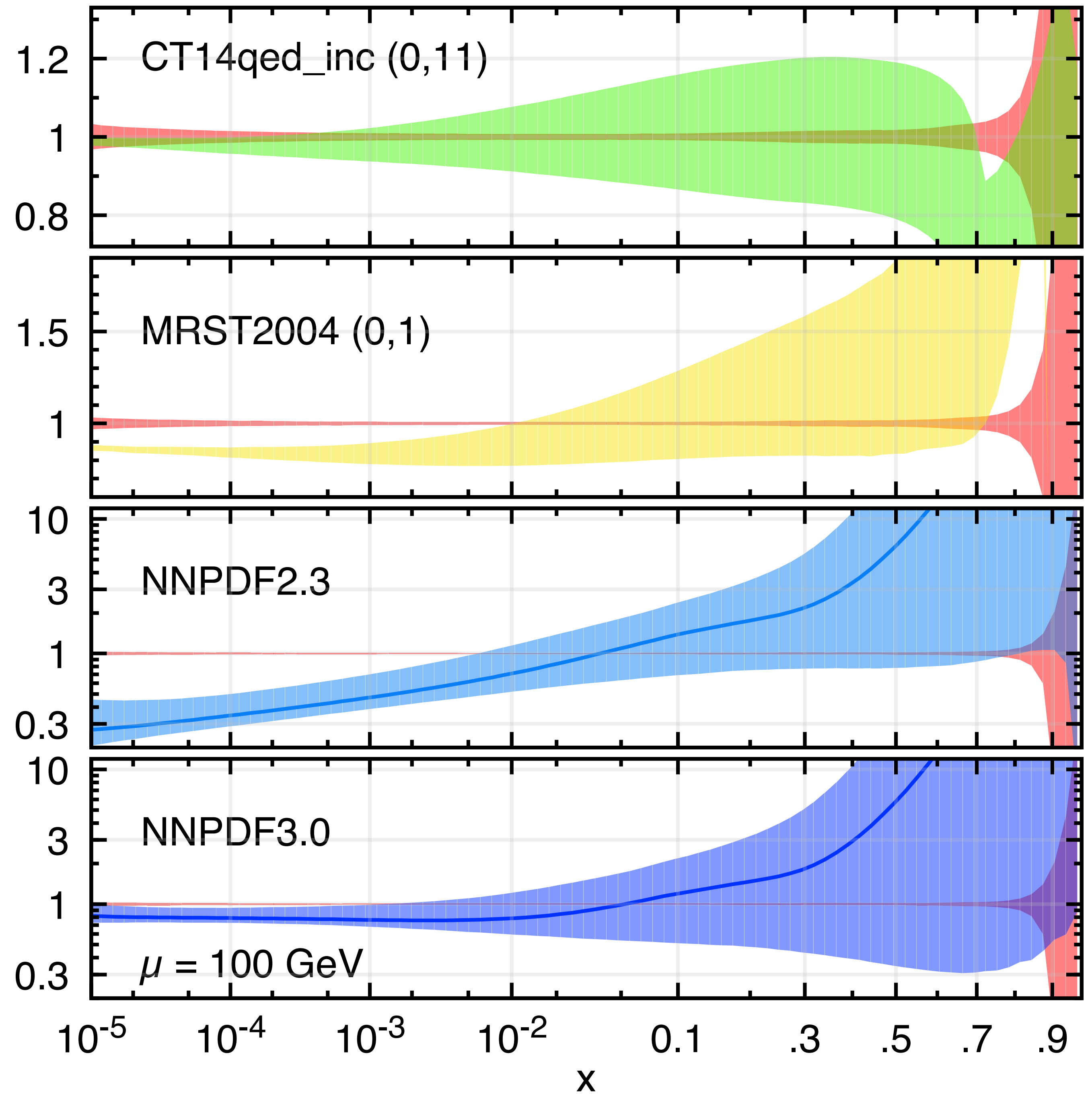


joint limits on κ_c & κ_b @ HL-LHC



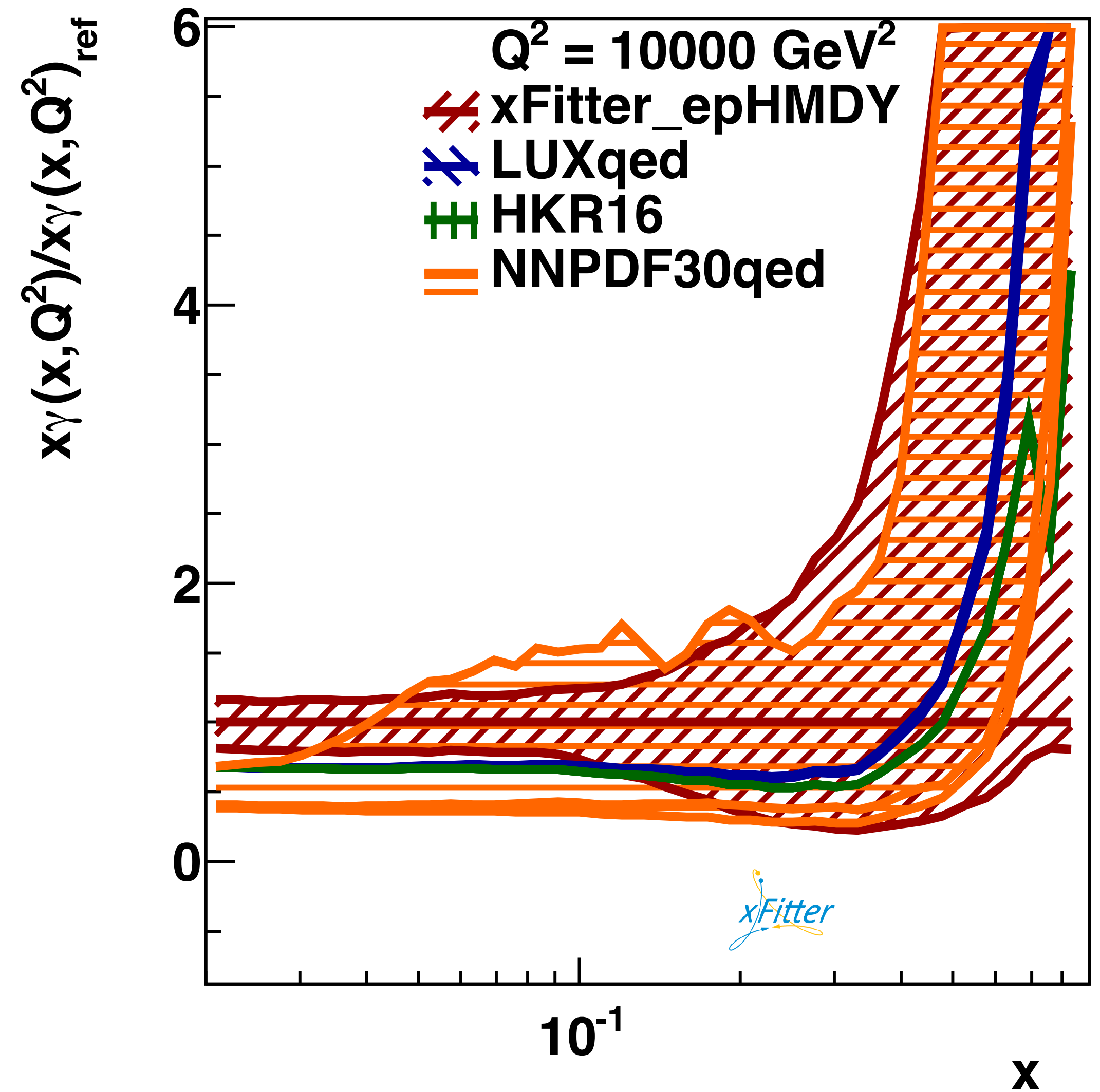
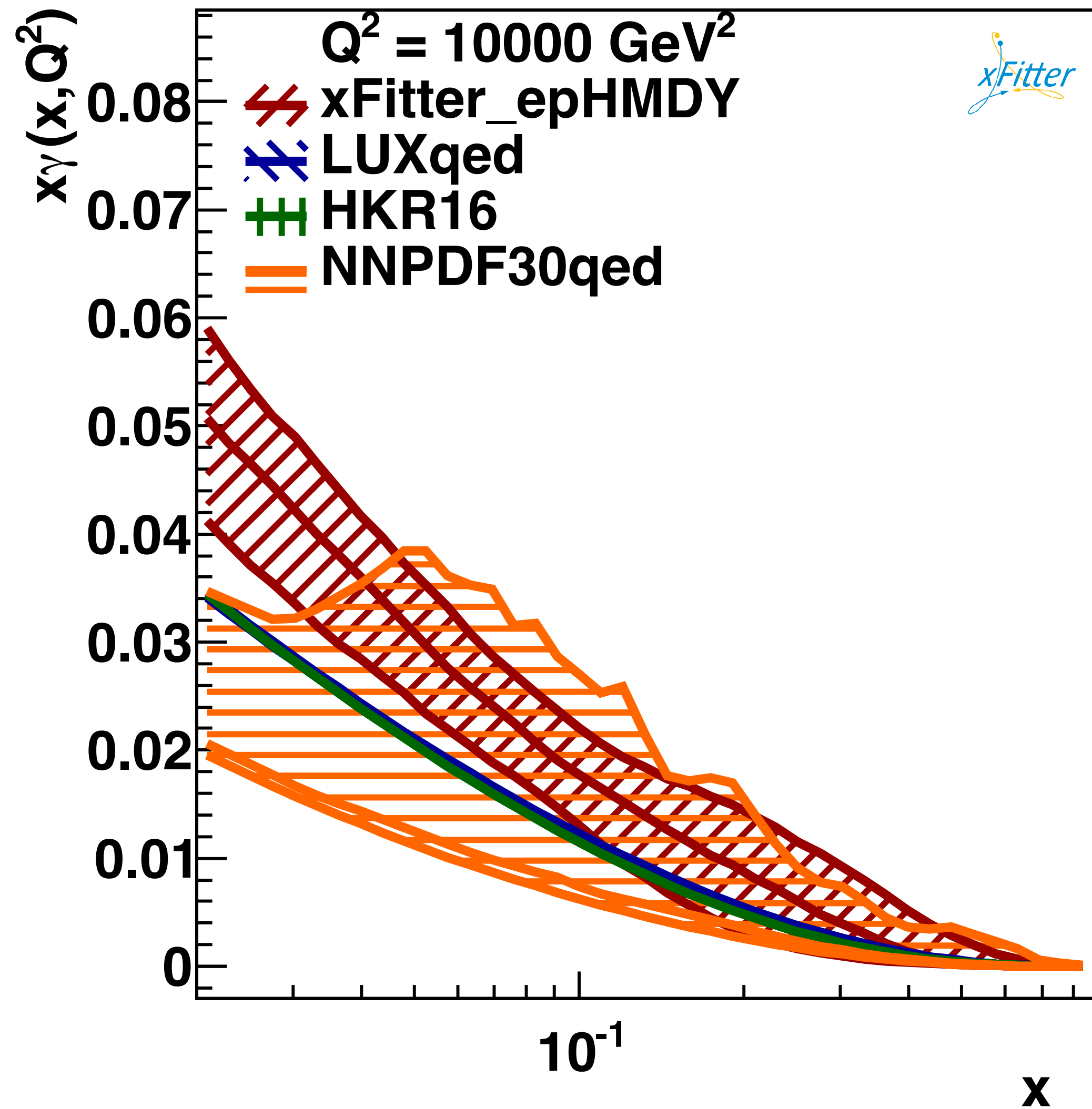
photon PDF

LUXqed v. other photon PDFs

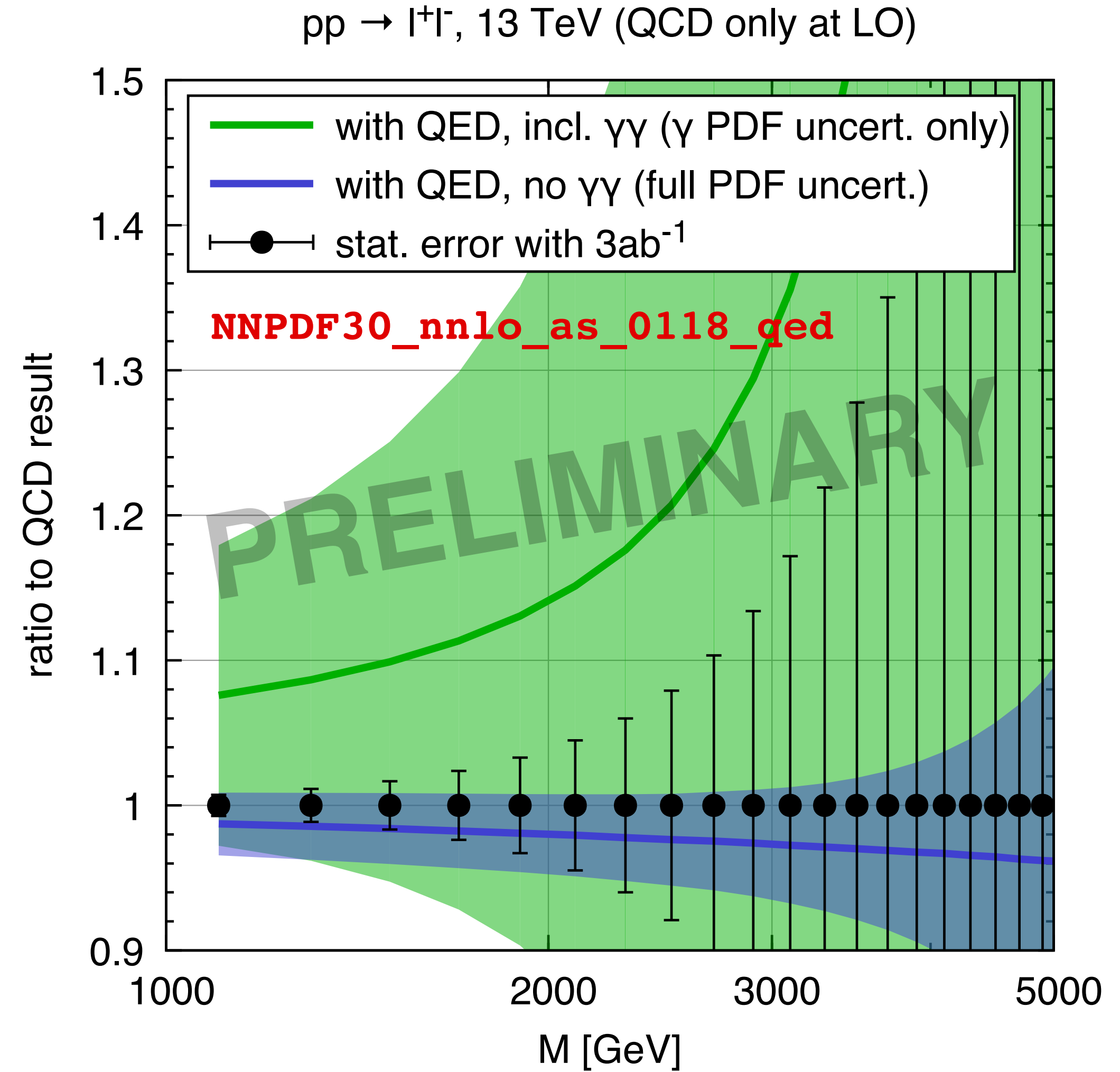
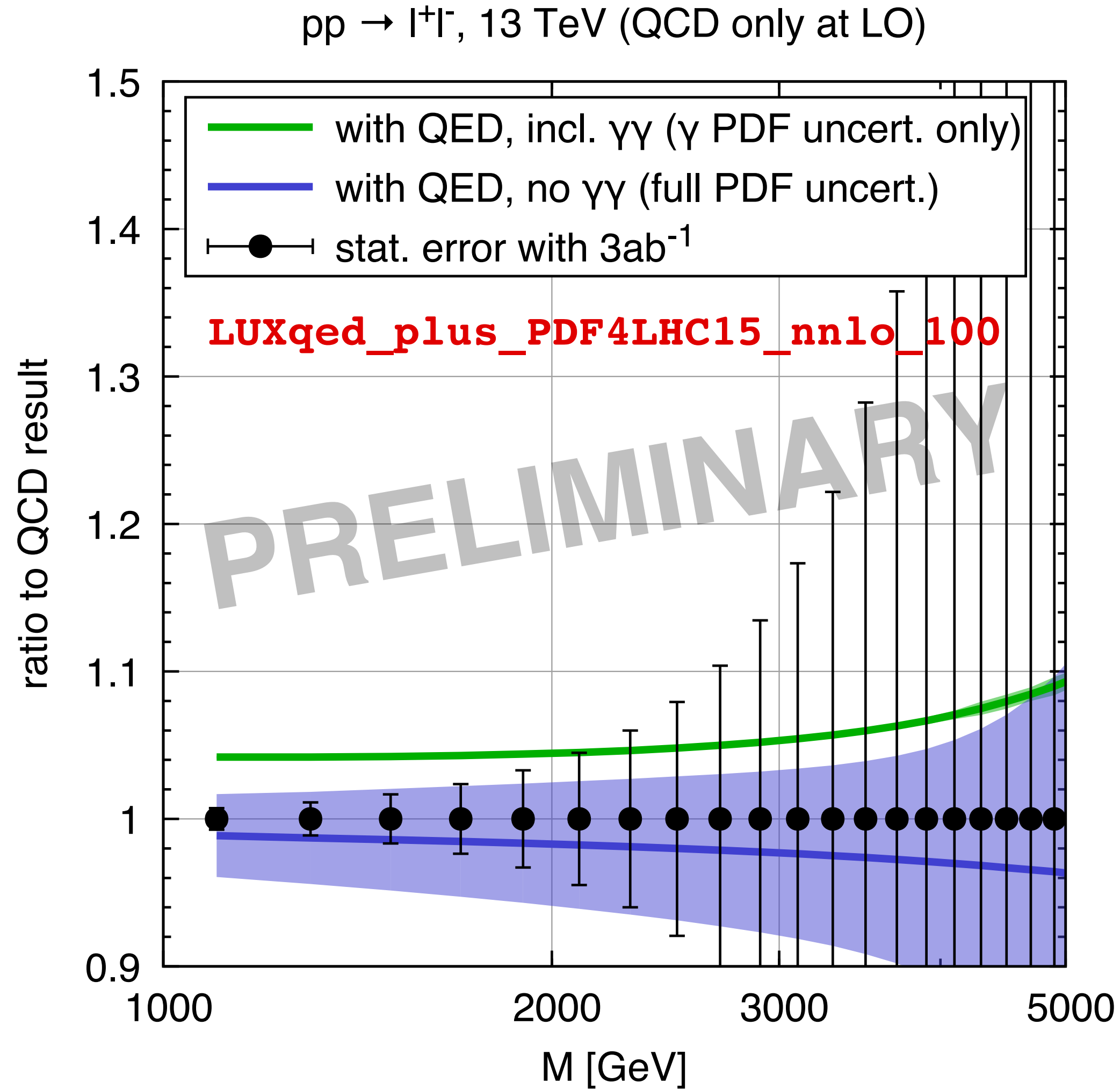


LUXqed v. a recent fit to Drell-Yan data

F. Giuli et al, 1701.08553



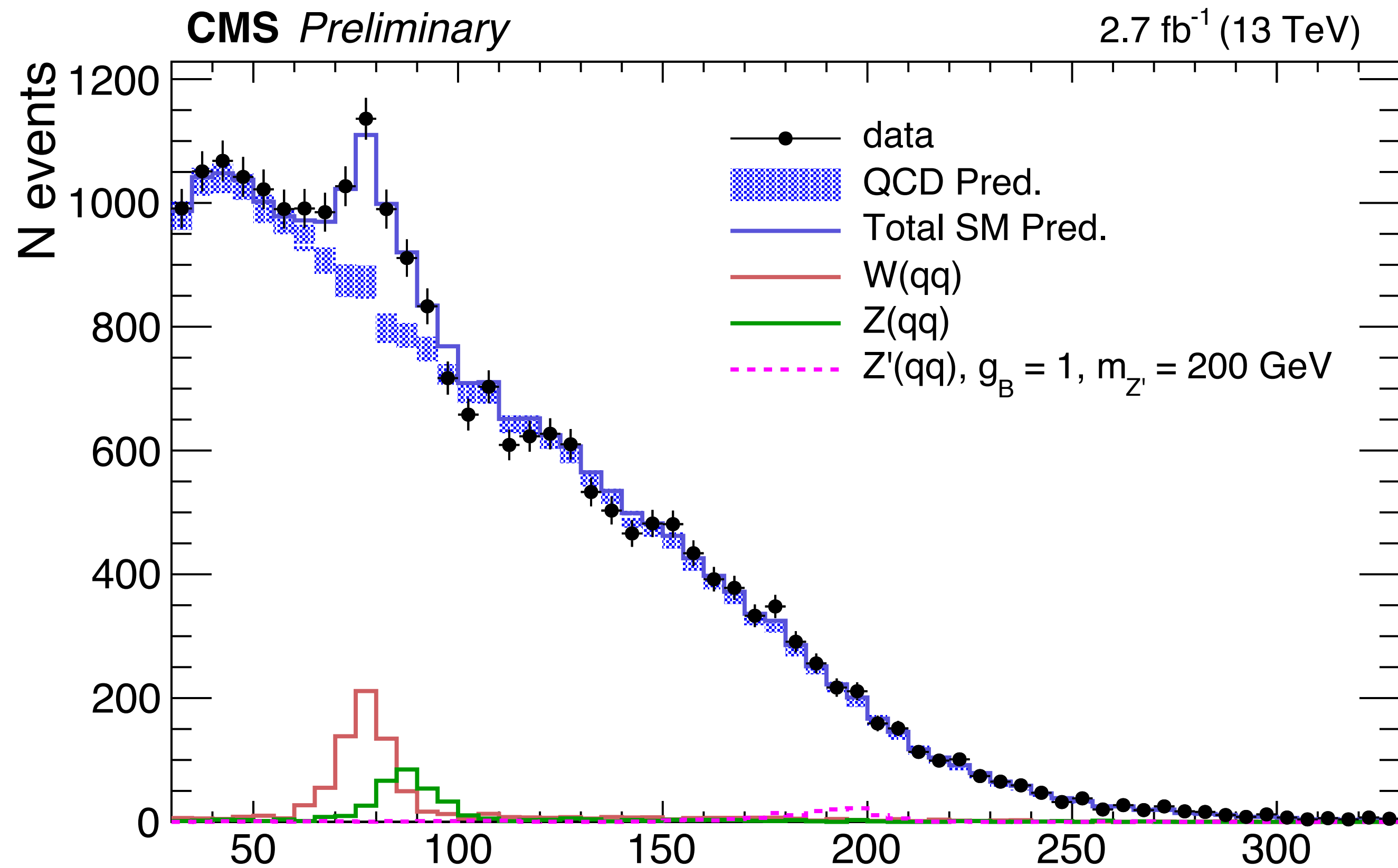
di-lepton spectrum with 3ab^{-1}



LUXQED photon has few % effect on di-lepton spectrum and negligible uncertainties

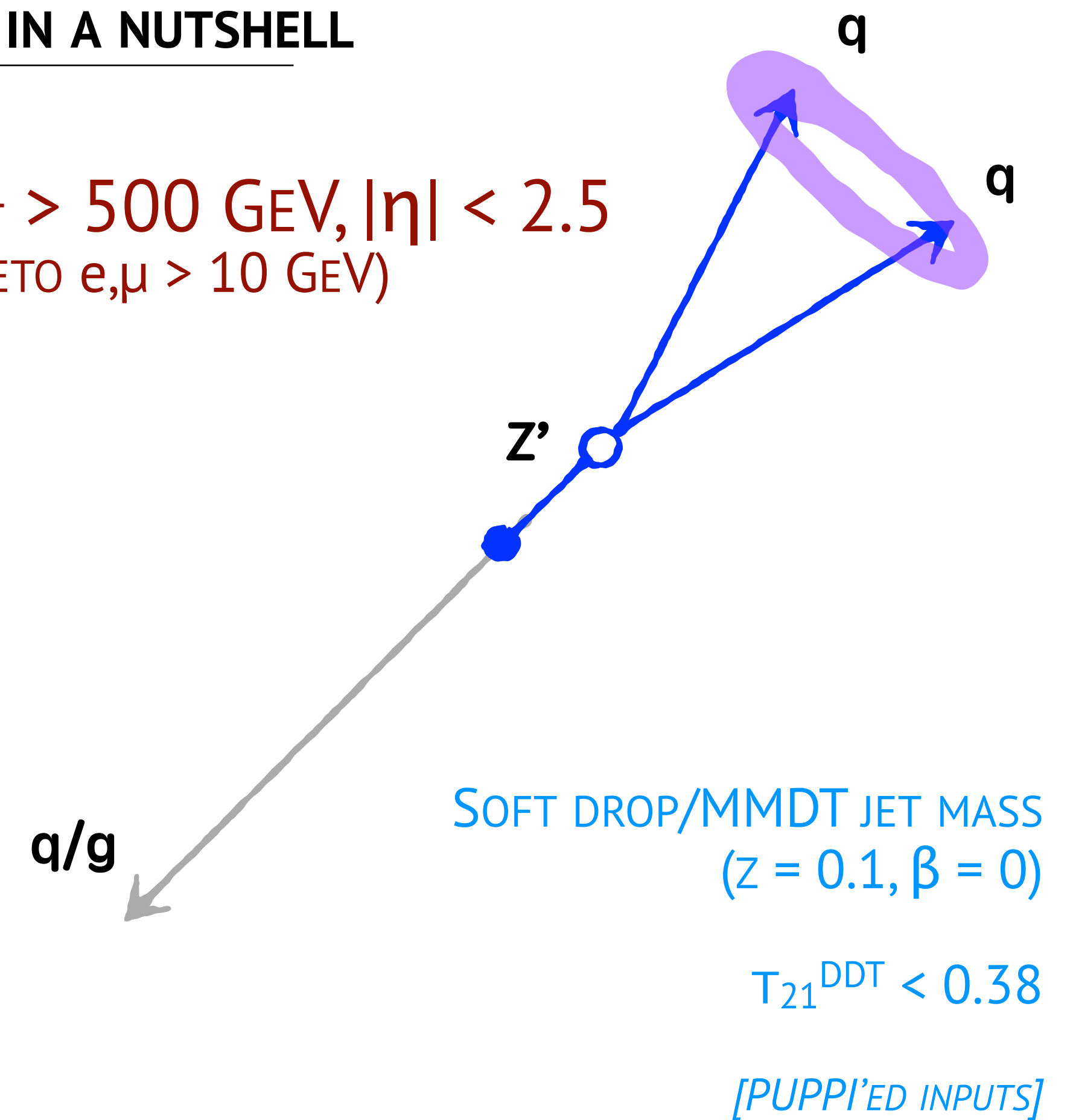
boosted hadronic decays

The potential of jet substructure — hadronic W & Z peaks



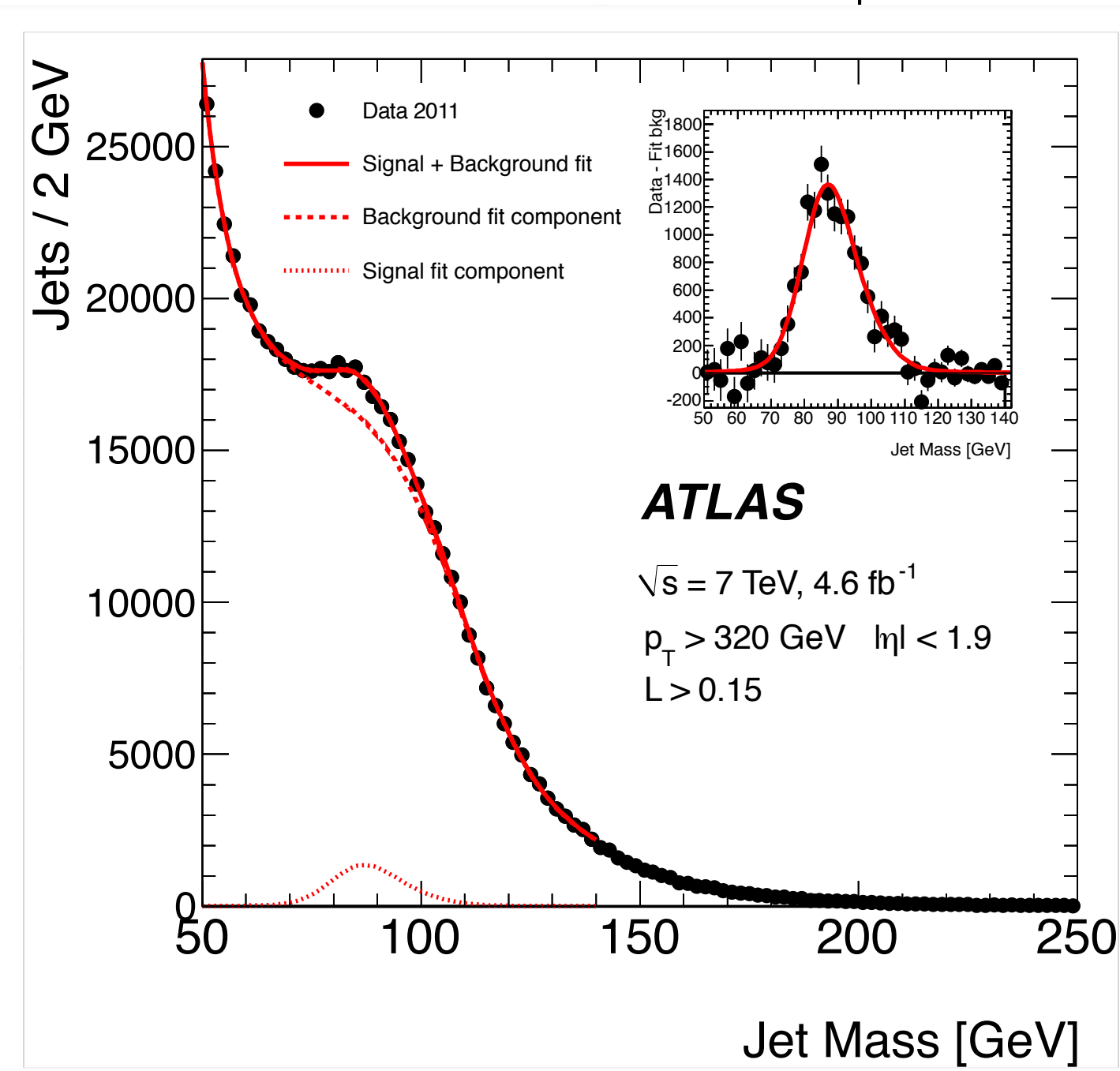
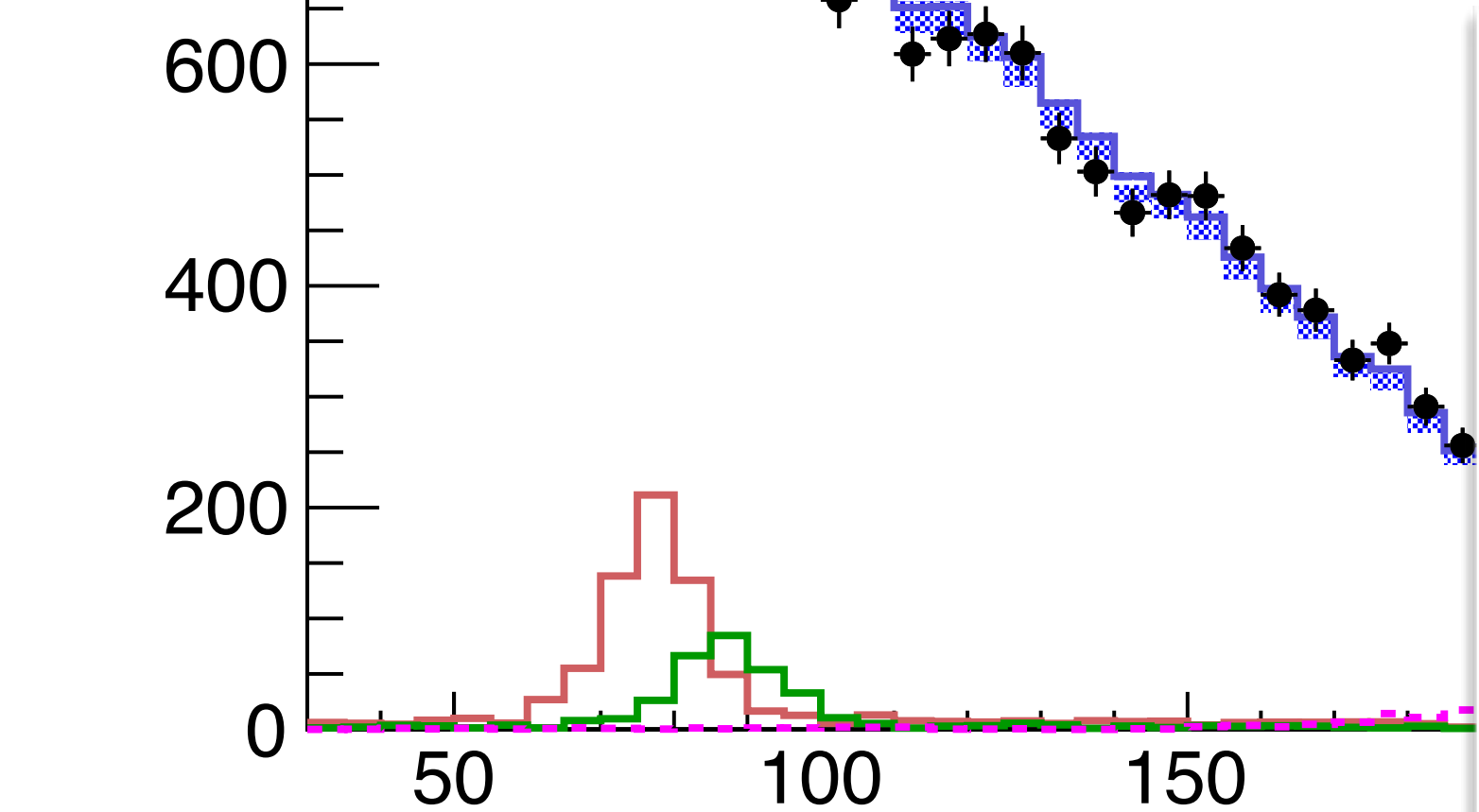
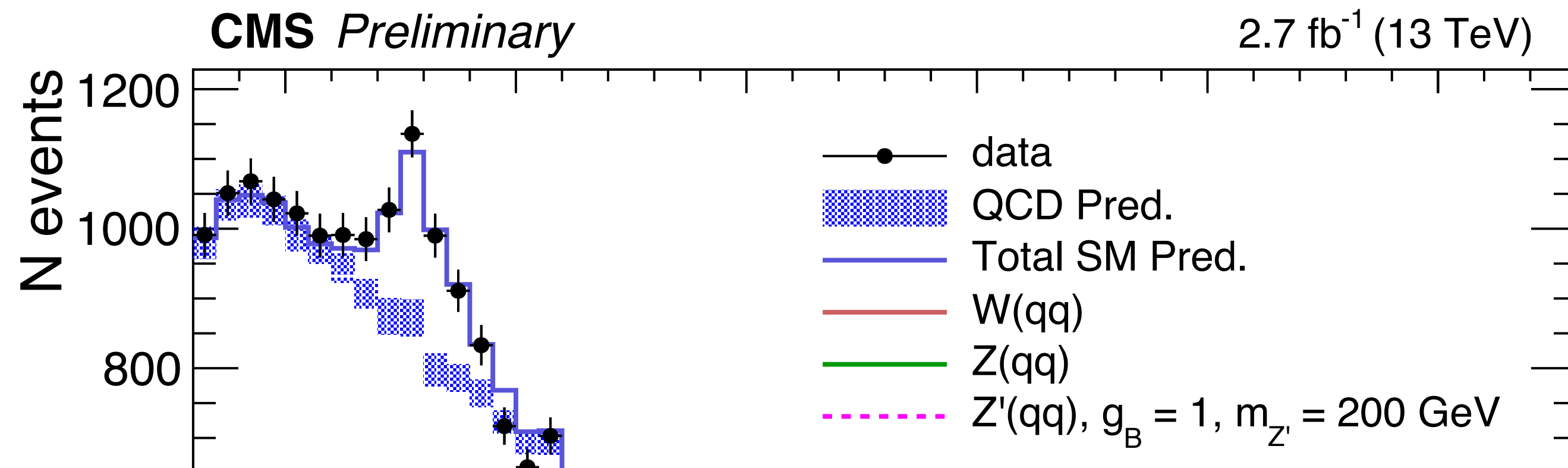
IN A NUTSHELL

$P_T > 500$ GeV, $|\eta| < 2.5$
(VETO $e, \mu > 10$ GeV)



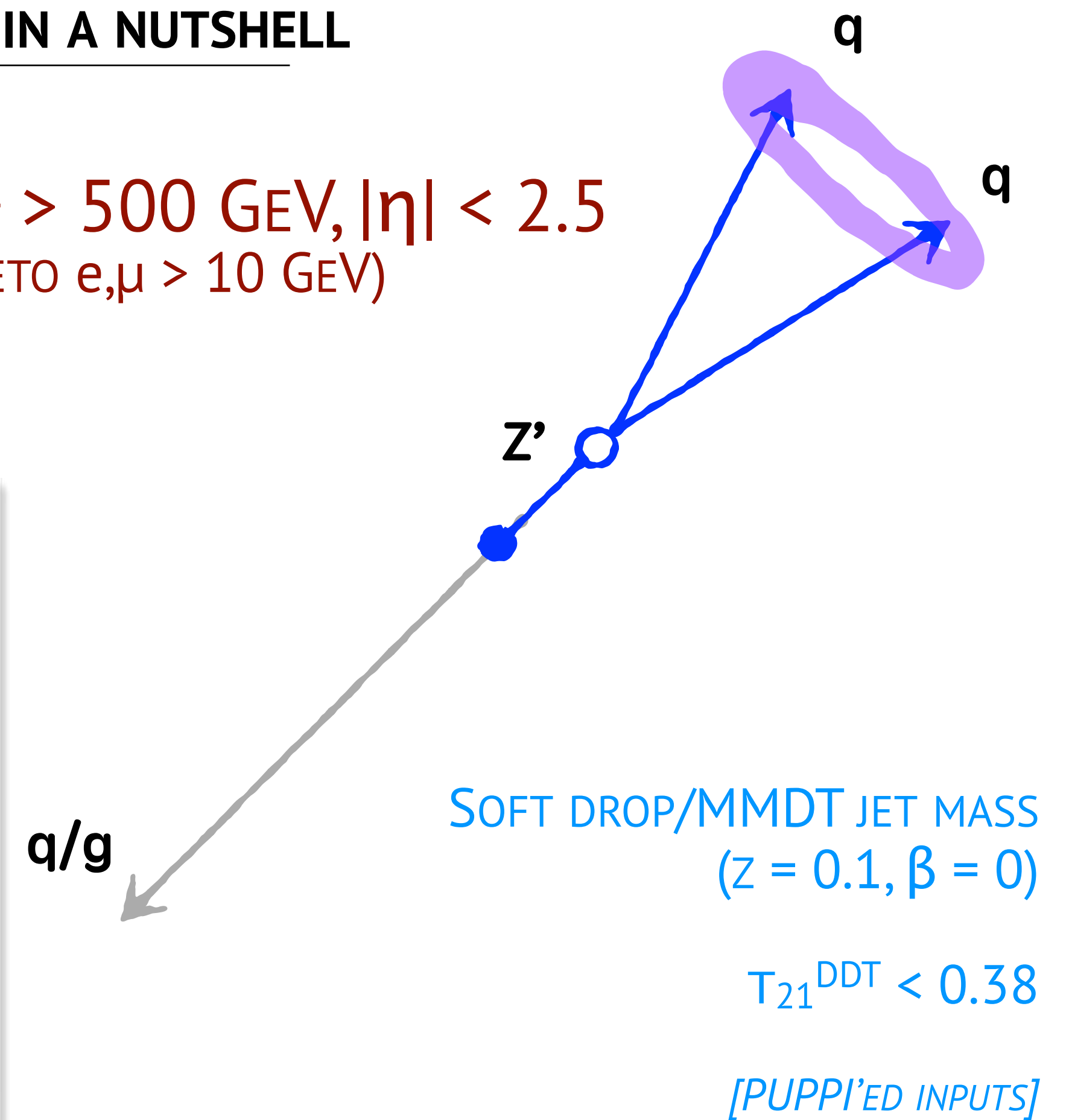
Nhan Tran @ Boost 2016

The potential of jet substructure — hadronic W & Z peaks



IN A NUTSHELL

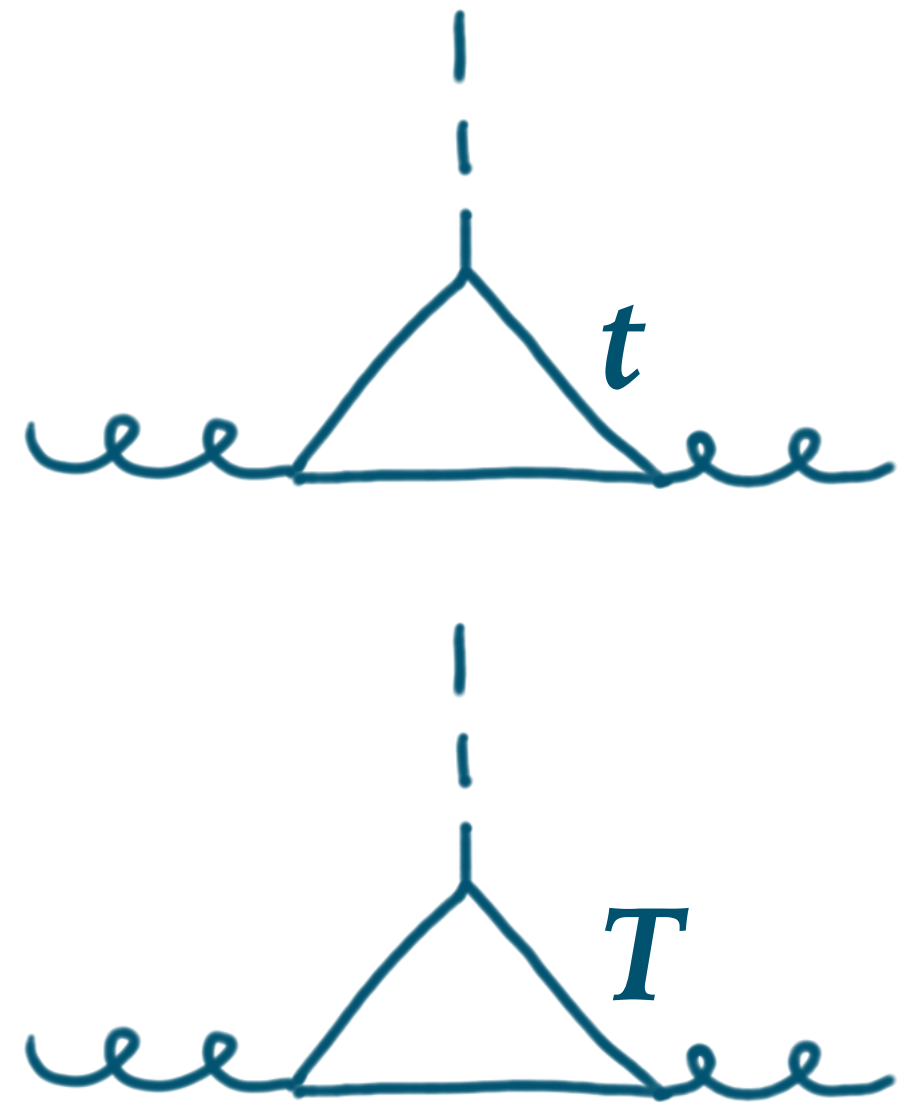
$p_T > 500$ GeV, $|\eta| < 2.5$
 (VETO $e, \mu > 10$ GeV)



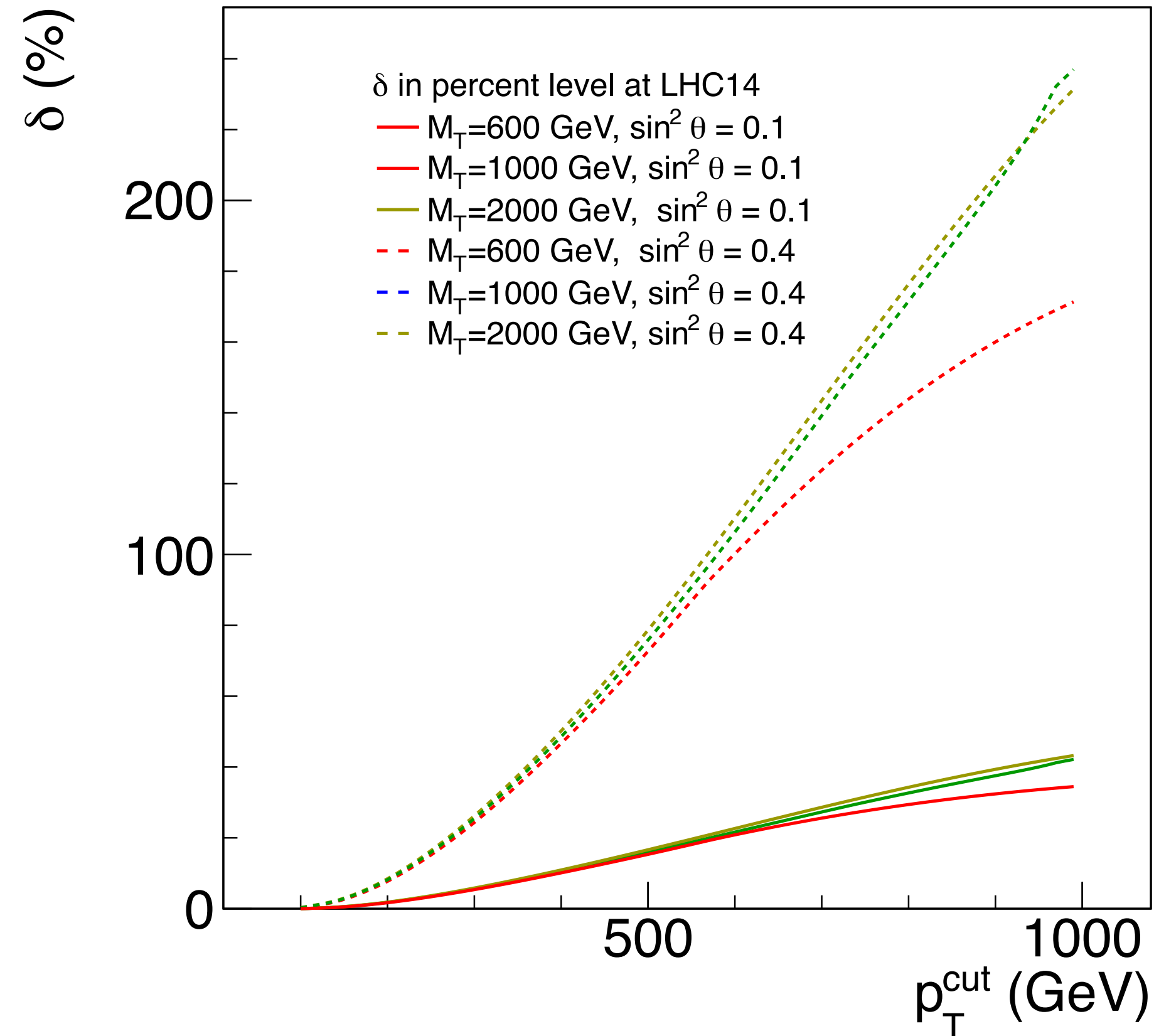
Nhan Tran @ Boost 2016

low v. high pT

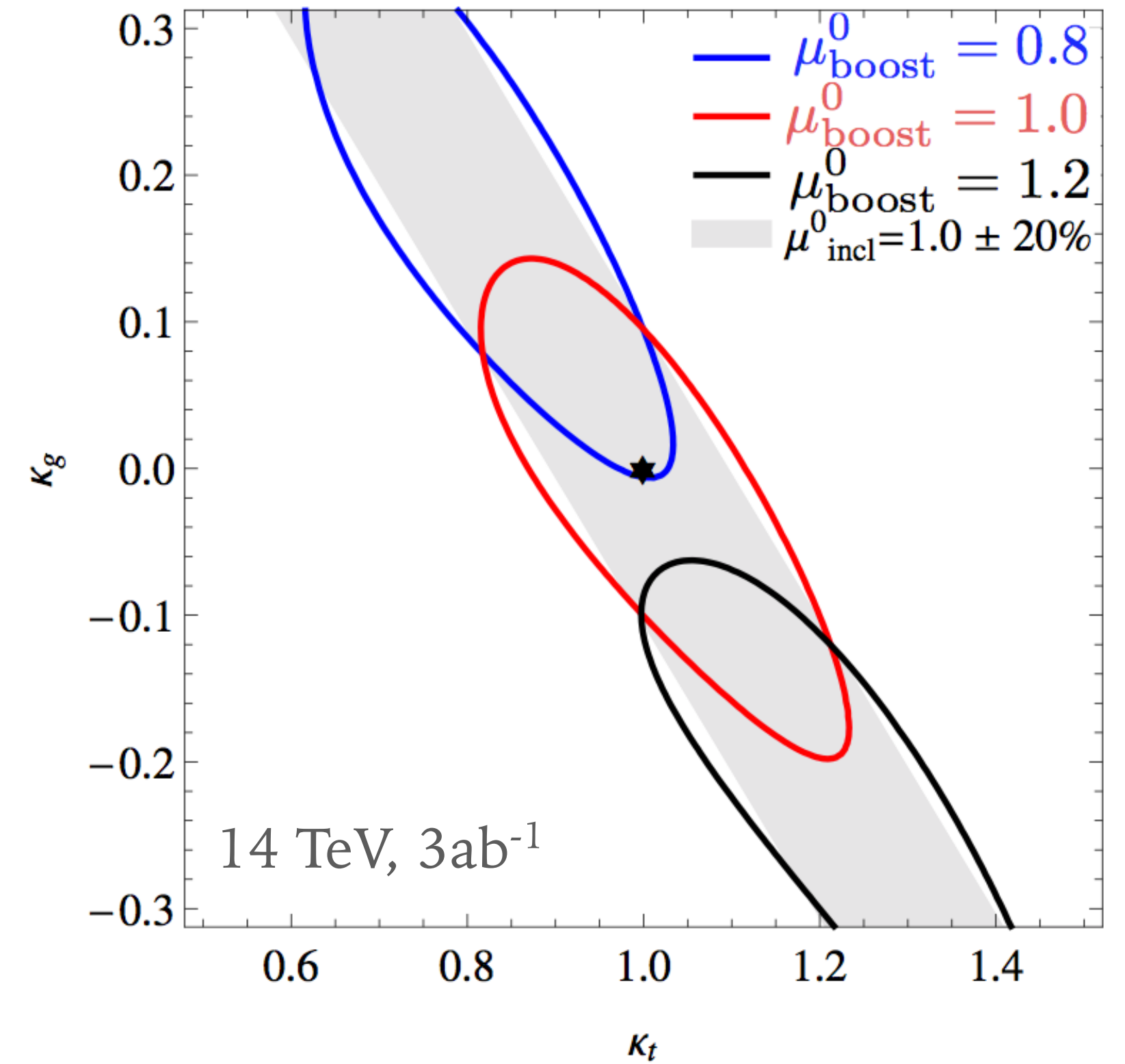
High-pt Higgs (e.g. to distinguish κ_g and κ_t)



Structure of loops is best probed by going to high p_T



A.Banfi, A.Martin, V.Sanz (2013)



C.Grojean et al. (2013)

see also Azatov, Paul (2013)

S.Dawson, I.Lewis, M.Zeng (2014)

what are experimental prospects?
are there any theory-issues to be solved?

VH PRODUCTION AT LARGE M(VH)

See also e.g.
 Biekötter, Knochel, Krämer, Liu, Riva,
 arXiv:1406.7320

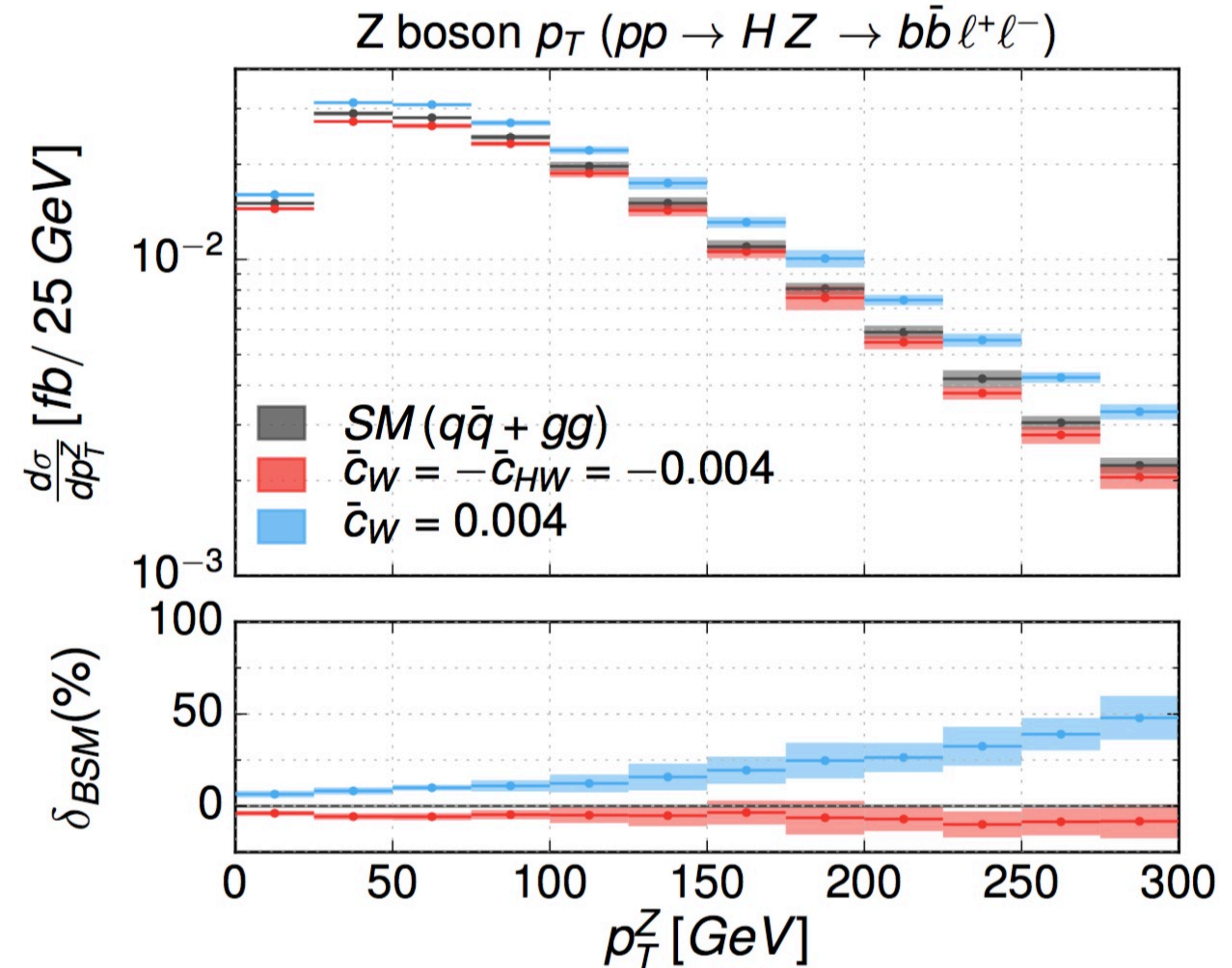
- Higher-dimension operators cause deviations that grow as, e.g.

$$\frac{\delta\sigma_{\text{dim-6}}}{\sigma} \sim \frac{p_T^2}{\Lambda^2}$$

- In some relevant range of p_T , Λ value to which you're sensitive grows as

$$\Lambda \sim (\text{Lumi})^{1/4}$$

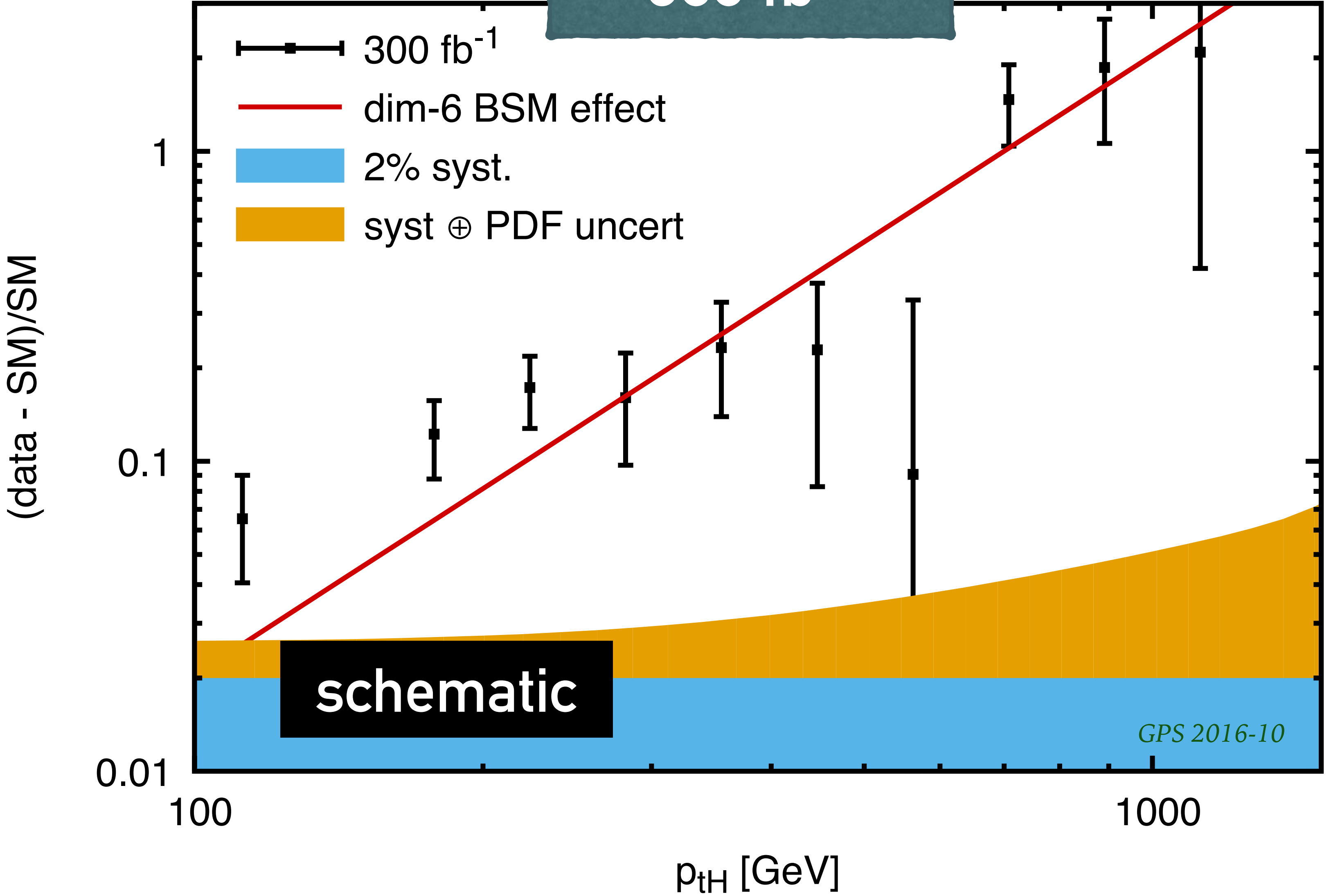
- that's faster than most direct searches (x100 in lumi \rightarrow x1.5 in reach for Z')



Mimasu, Sanz, Williams, arXiv:1512.02572v

WH at large Q^2 with dim-6 BSM effect

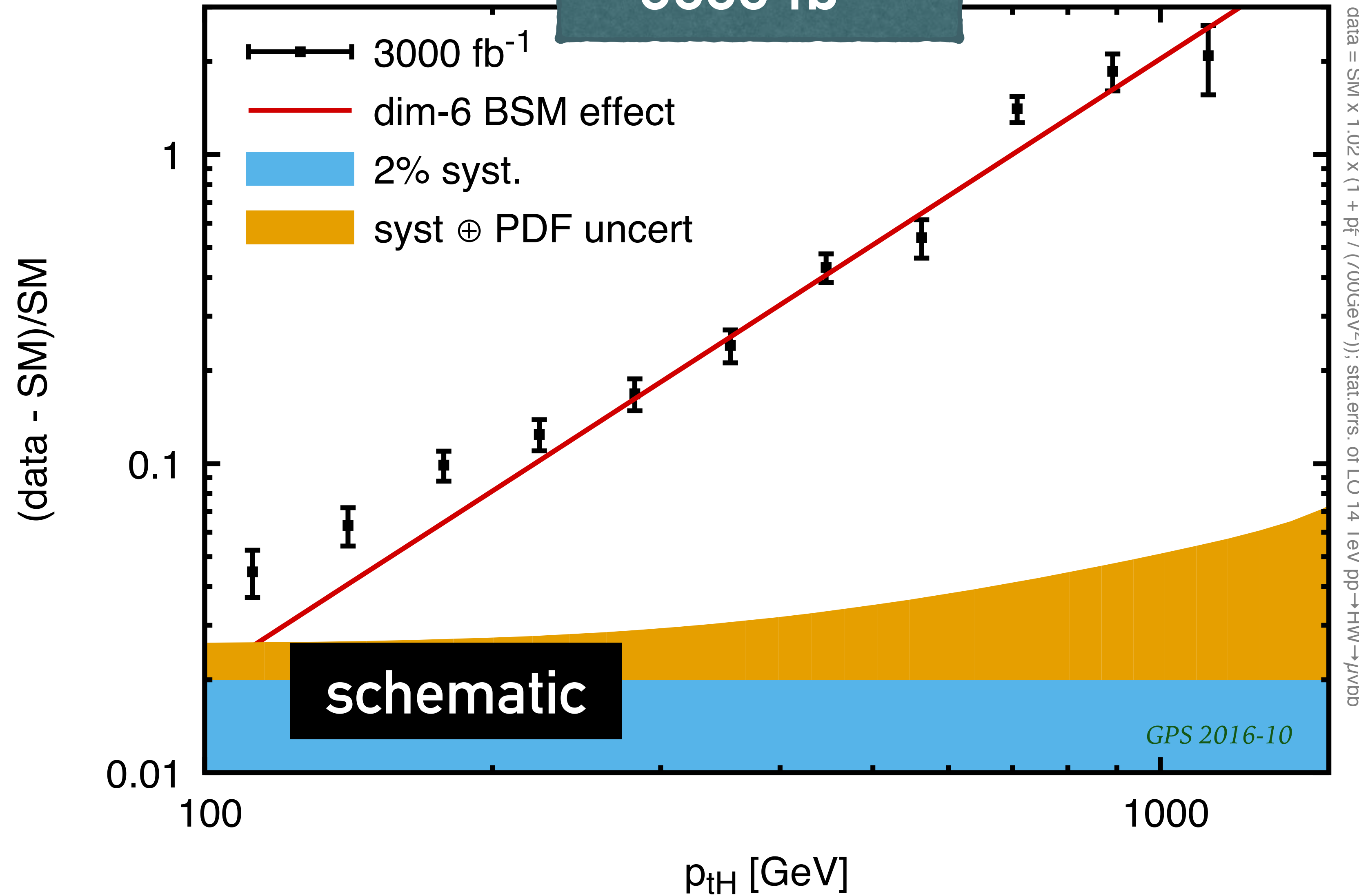
300 fb⁻¹



data = SM x 1.02 x (1 + p_{tH}² / (700 GeV²)); stat. errs. of LO 14 TeV pp → HW → μνbb

WH at large Q^2 with dim-6 BSM effect

3000 fb⁻¹

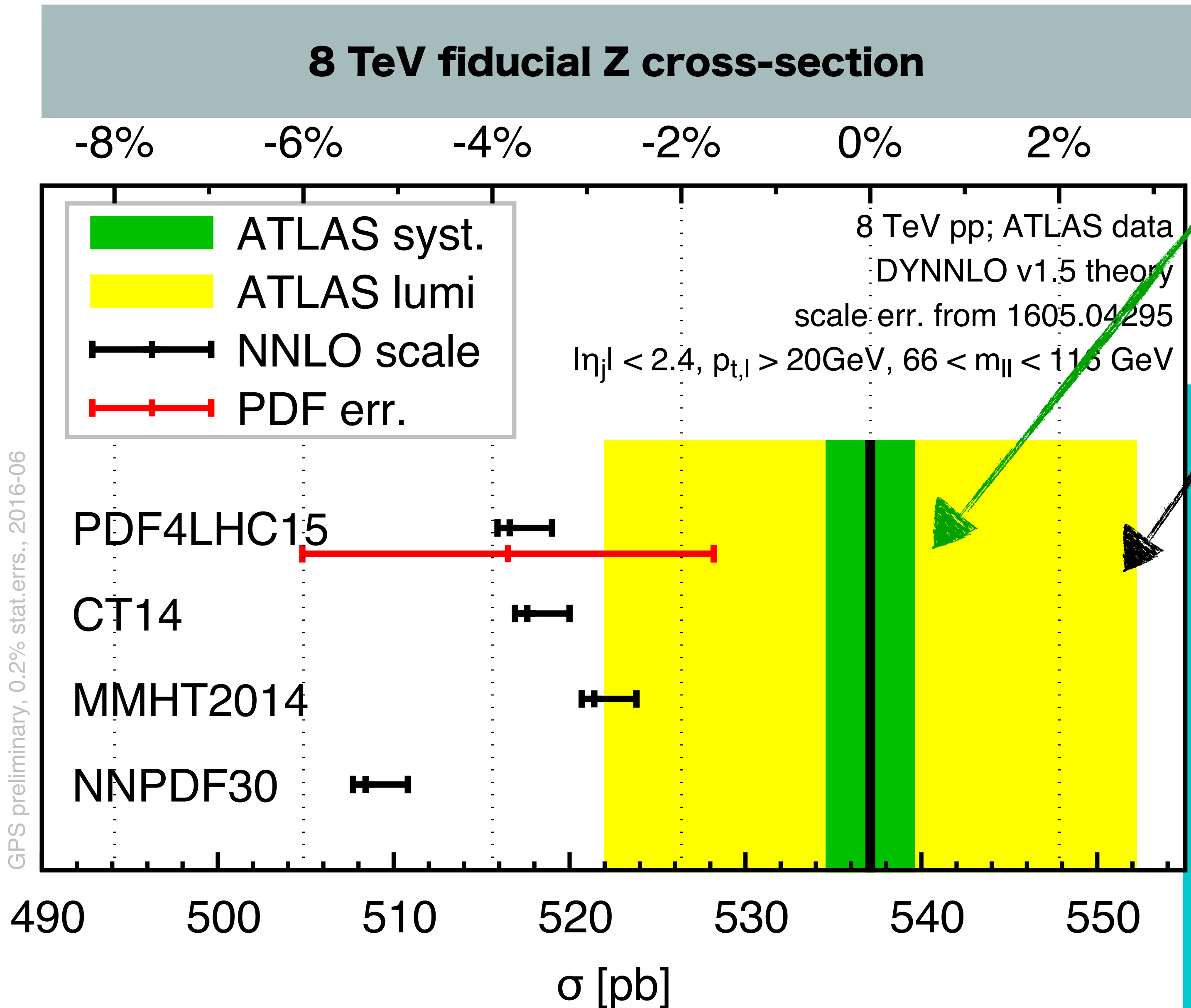


new physics isn't just a single number that's wrong (think $g-2$)

but rather a **distinct scaling pattern of deviation** ($\sim p_T^2$)

moderate and high p_T 's have similar statistical significance — so it's useful to understand whole p_T range

There are, however, issues. Notably in Z production



$\pm 0.45\%$ syst. **$\times 6$**

$\pm 2.8\%$ lumi

Up to 5% discrepancy with data

Experimental progress on luminosity determination may be the **keystone** for precision physics at LHC.

Are there hardware changes to HL-LHC that could help with lumi determination?