New results for parton showers

Gavin P. Salam*

Rudolf Peierls Centre for Theoretical Physics & All Souls College, Oxford

including work with M. Dasgupta, F. Dreyer, K. Hamilton, P. Monni

* on leave from CERN and CNRS





Annual Theory Meeting IPPP Durham, 17/12/2018

at colliders, you can probe

"big unanswered questions"

about fundamental particles & their interactions

(dark matter, matter-antimatter asymmetry,

nature of dark energy, hierarchy of scales...)

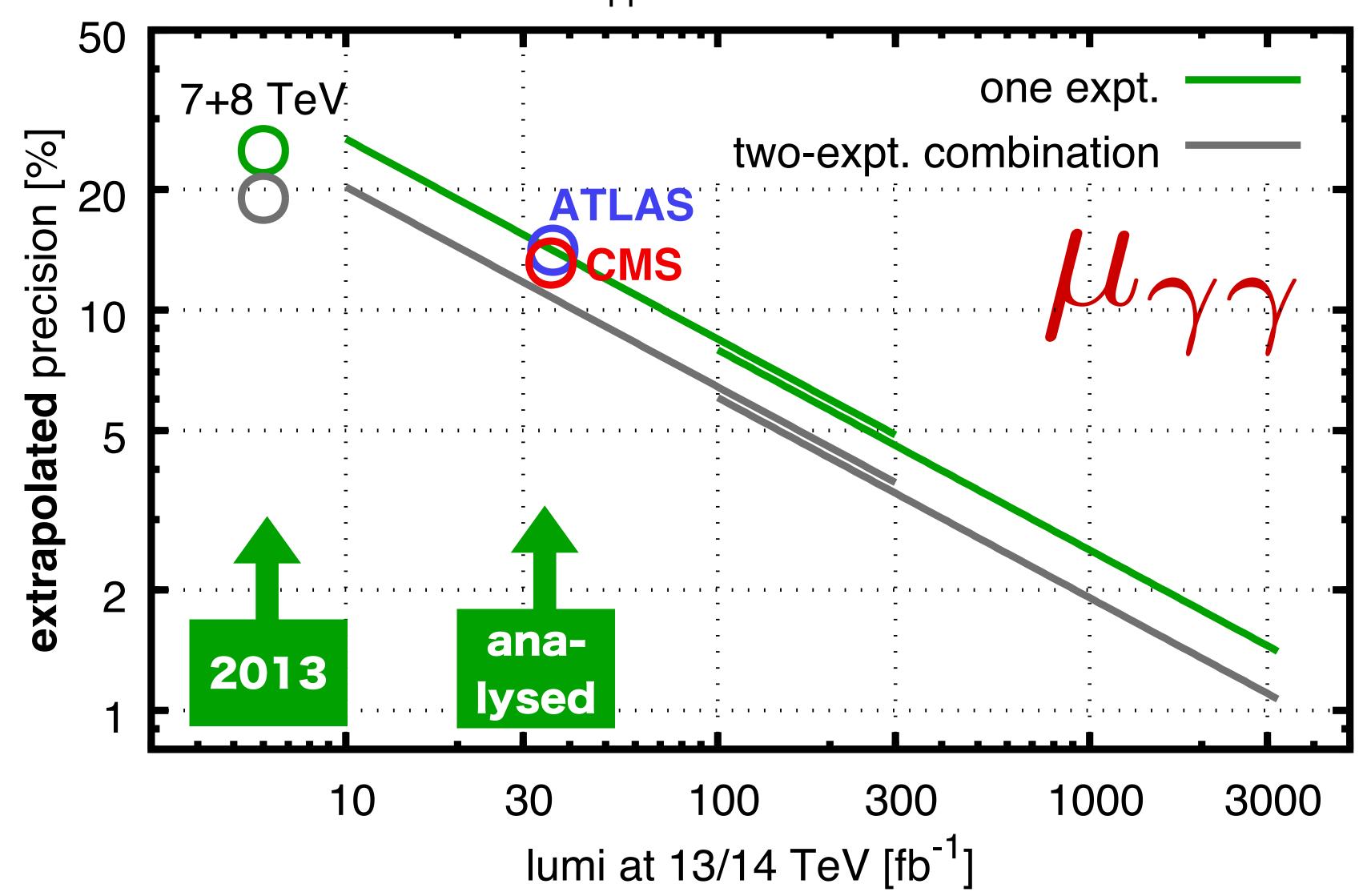
and

"big answerable questions"

(structure of Higgs sector, determining fundamental parameters of Lagrangian of particle physics)

Higgs precision (H $\rightarrow \gamma\gamma$) : optimistic estimate v. luminosity & time

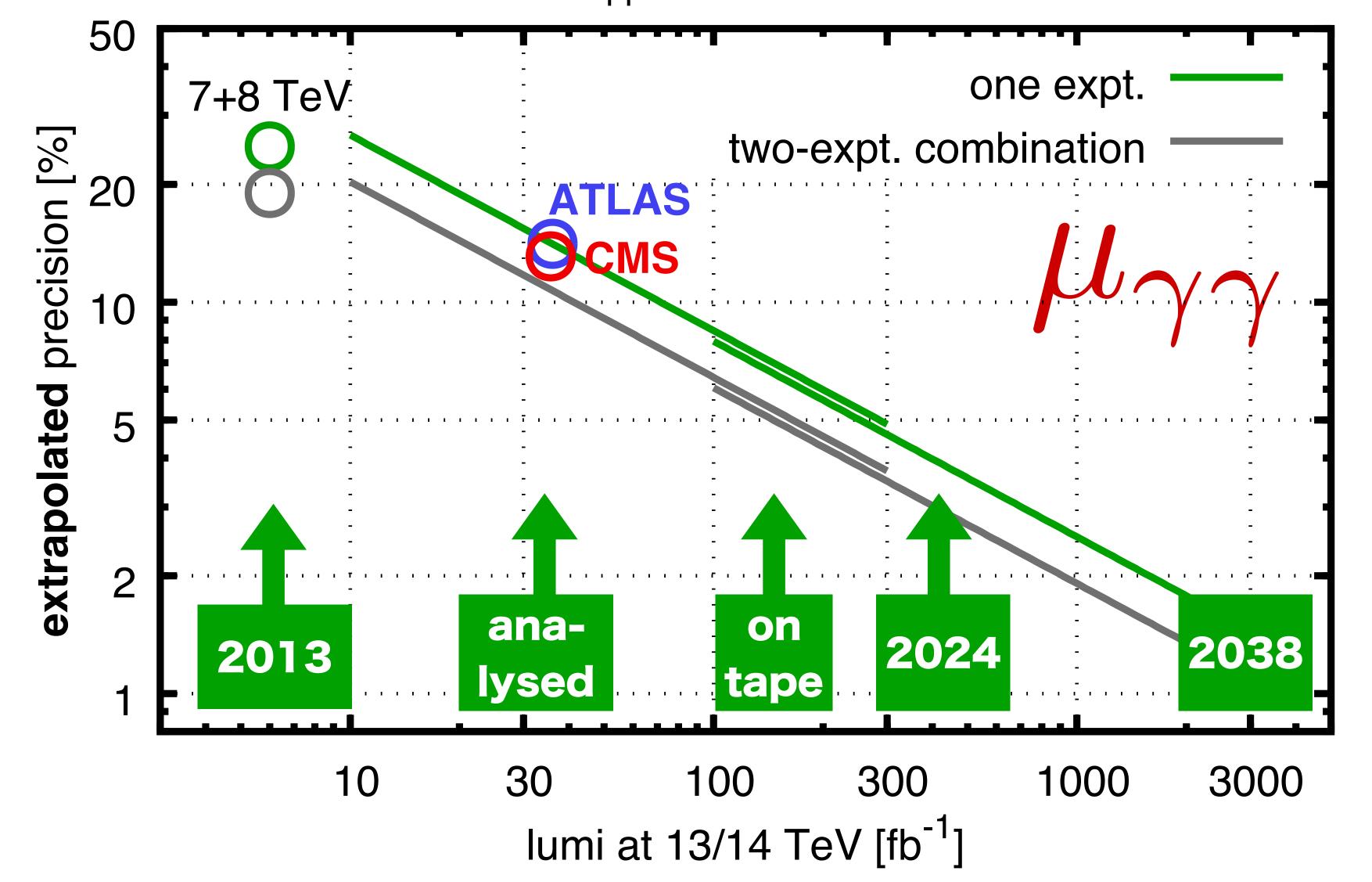




1 fb⁻¹ = 10^{14} collisions

Higgs precision (H $\rightarrow \gamma\gamma$) : optimistic estimate v. luminosity & time





The LHC has the statistical potential to take Higgs physics from "observation" to 1–2% precision

But only if we learn how to connect experimental observations with theory at that precision

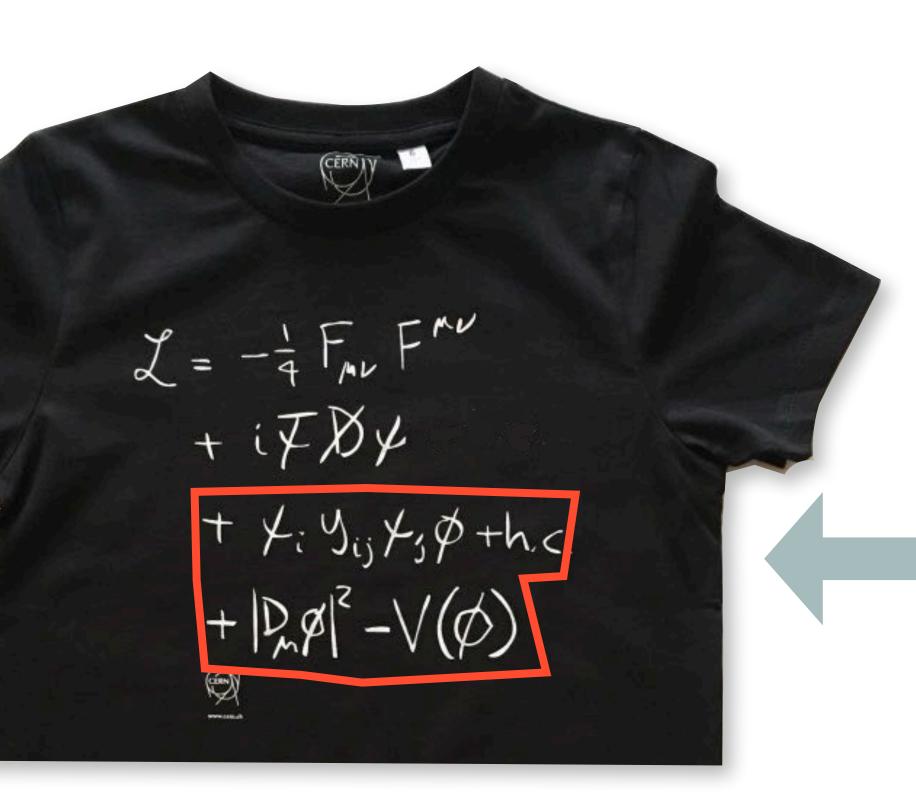
 $1 \text{ fb}^{-1} = 10^{14} \text{ collisions}$

how is all of this made quantitative?

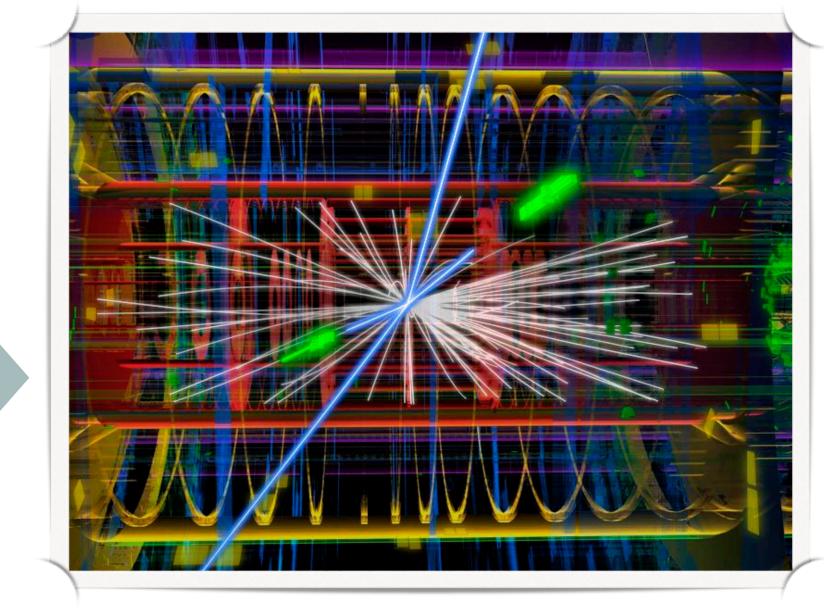
whether new-physics searches, Higgs physics, or other SM studies

UNDERLYING THEORY

EXPERIMENTAL DATA

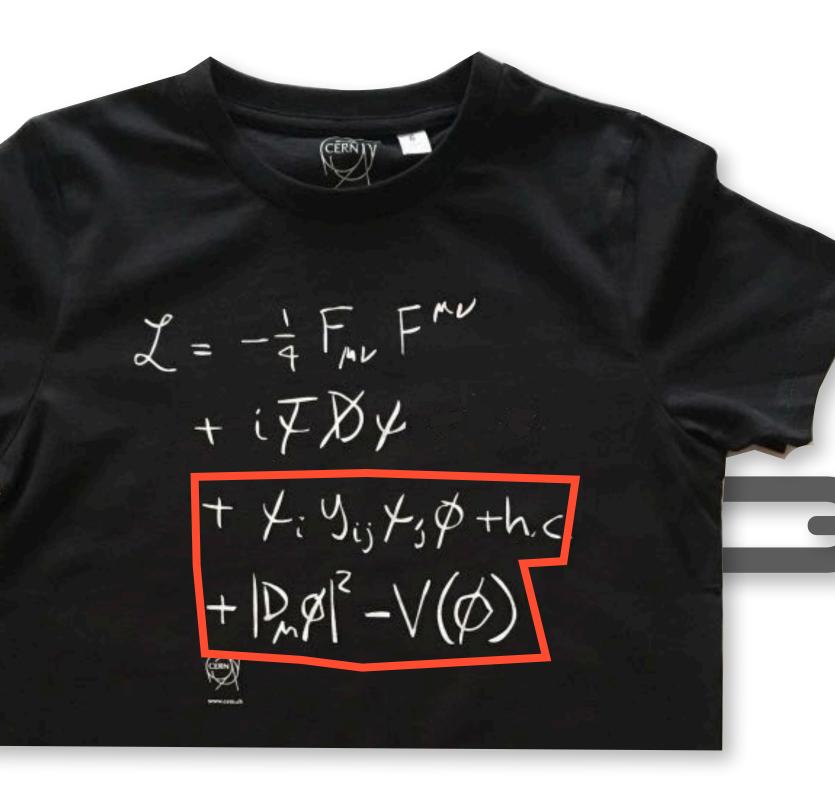


how do you make quantitative connection?



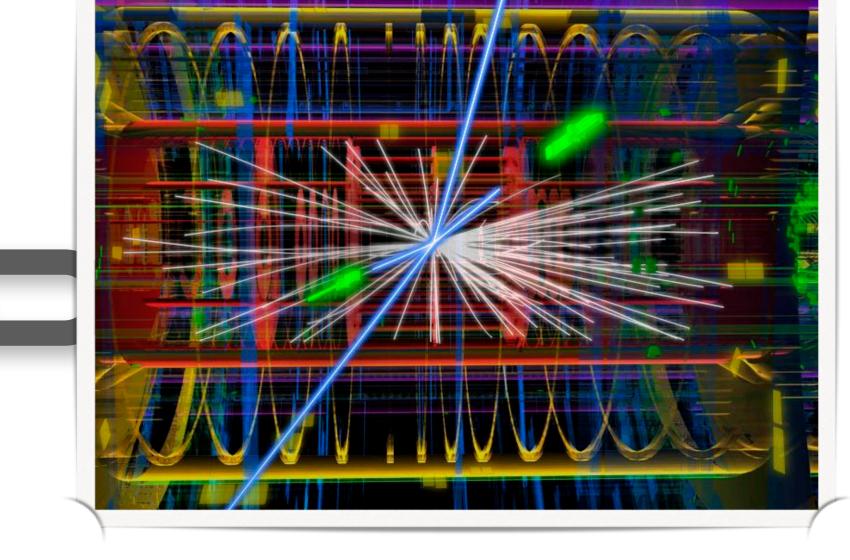
UNDERLYING THEORY

EXPERIMENTAL DATA

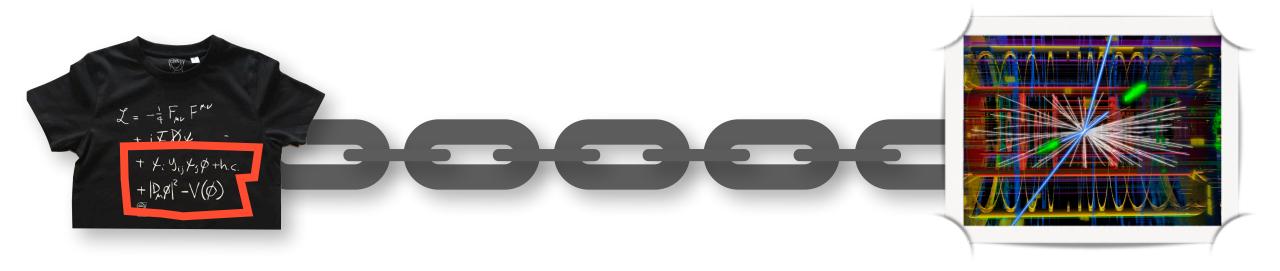


how do you make quantitative connection?

through a chain of experimental and theoretical links

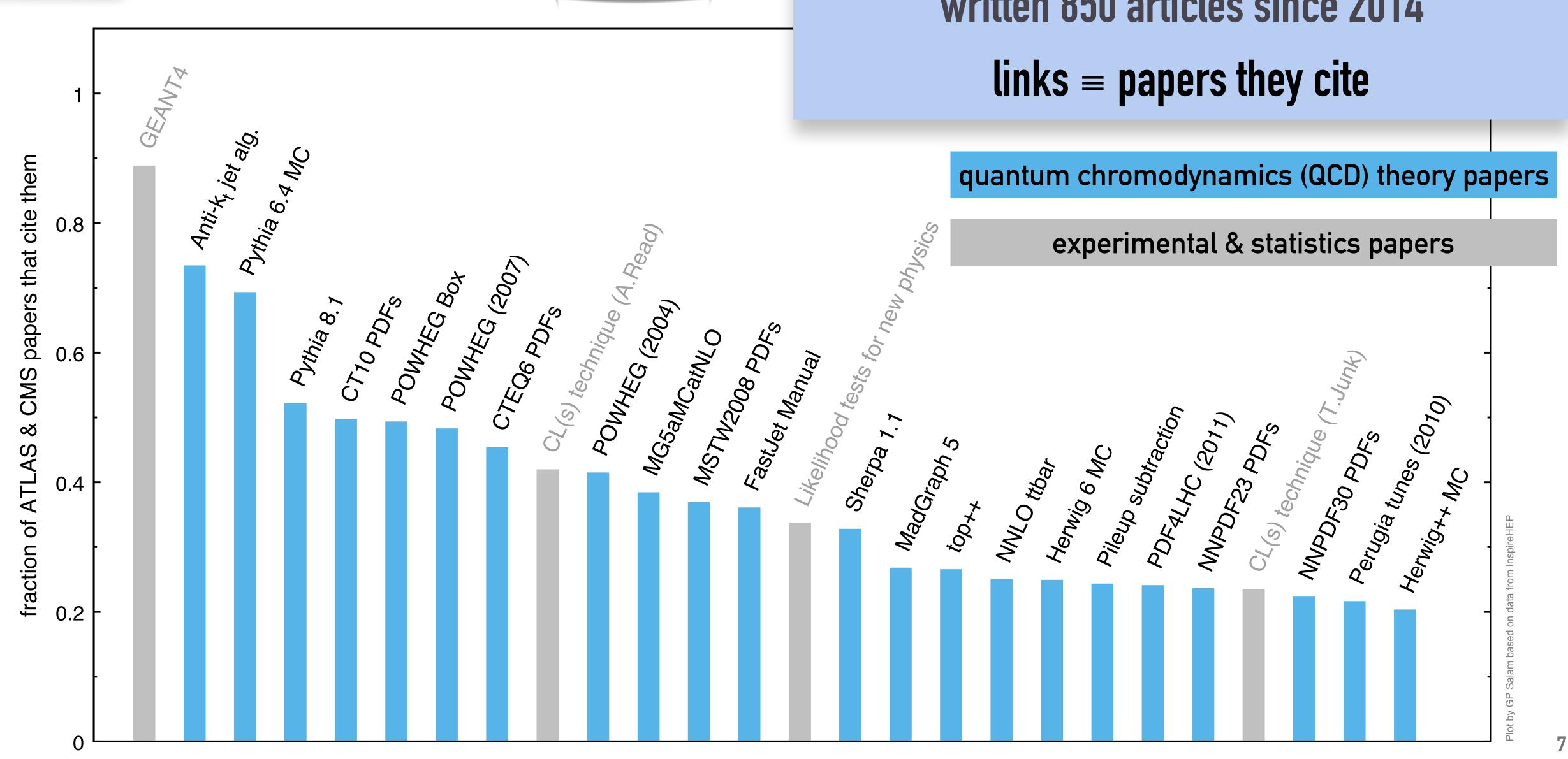


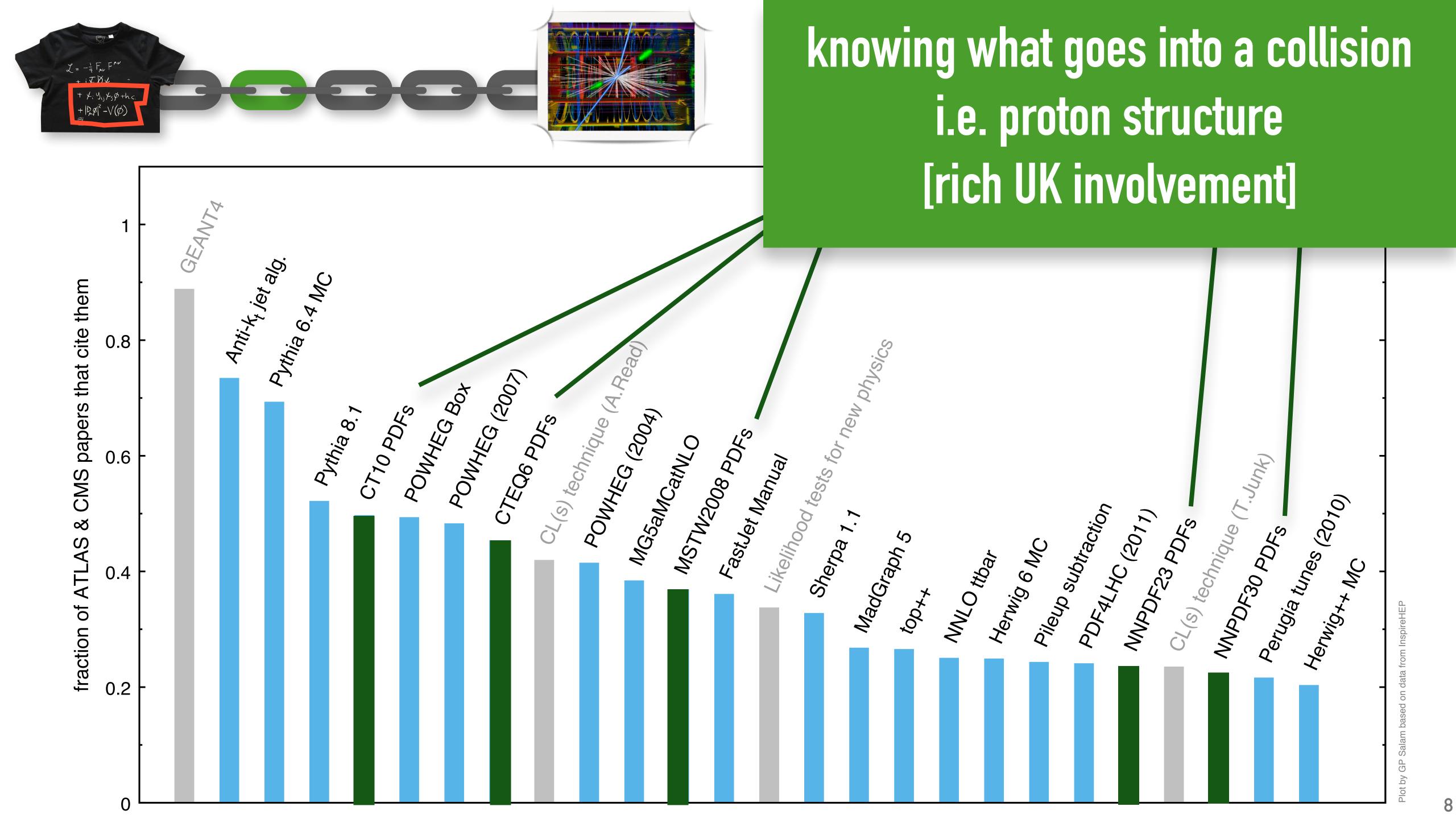
[in particular Quantum Chromodynamics (QCD)]

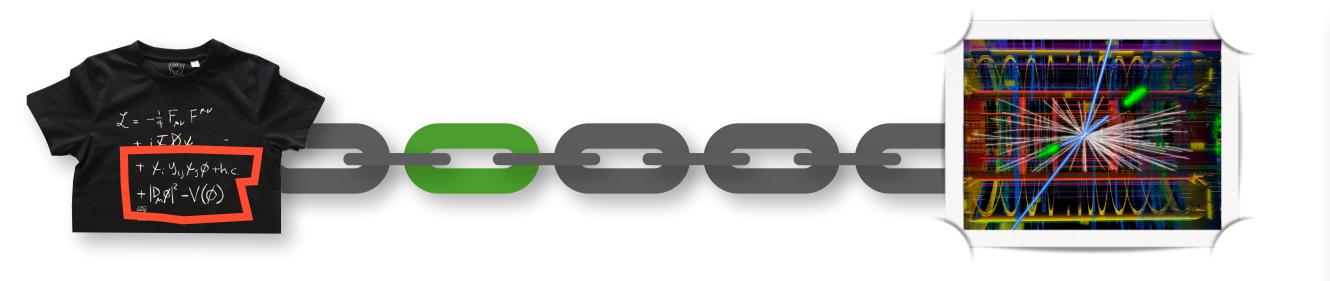


What are the links?

ATLAS and CMS (big LHC expts.) have written 850 articles since 2014

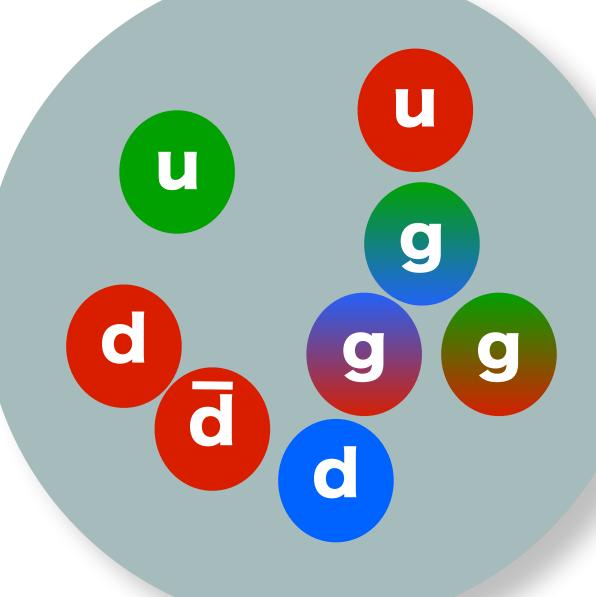


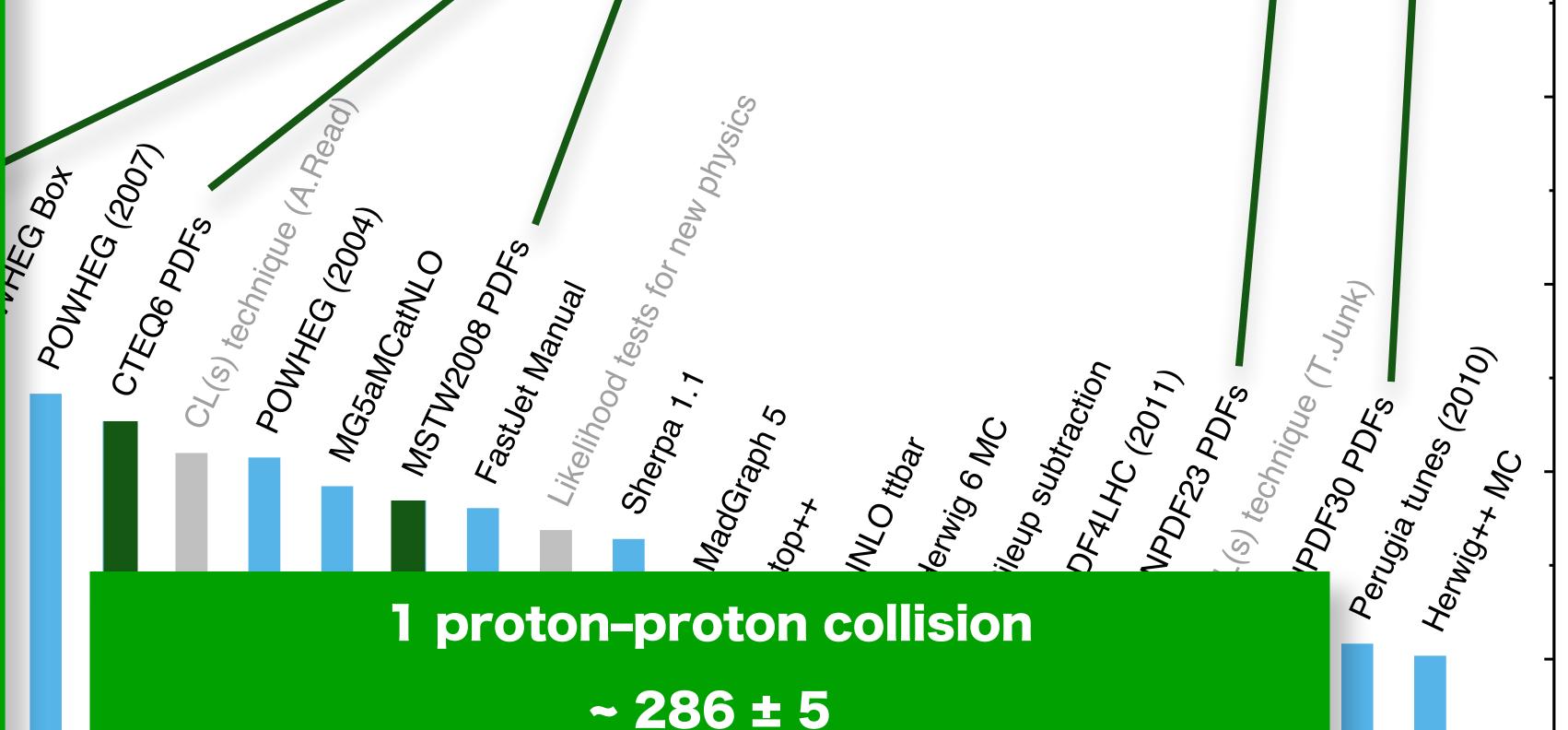




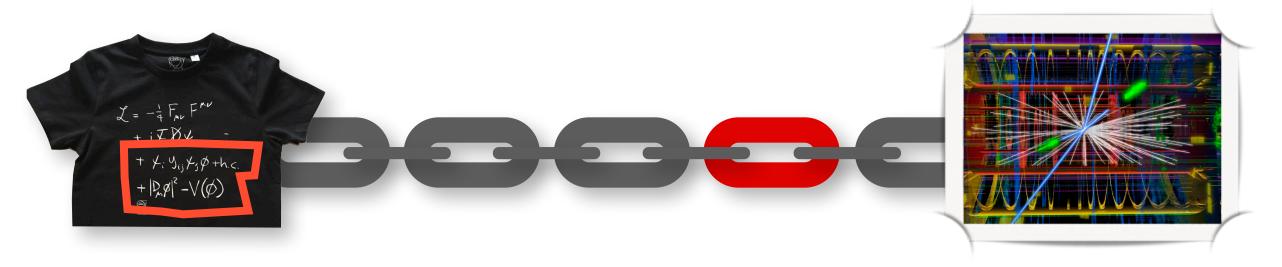
knowing what goes into a collision i.e. proton structure [rich UK involvement]

PROTON

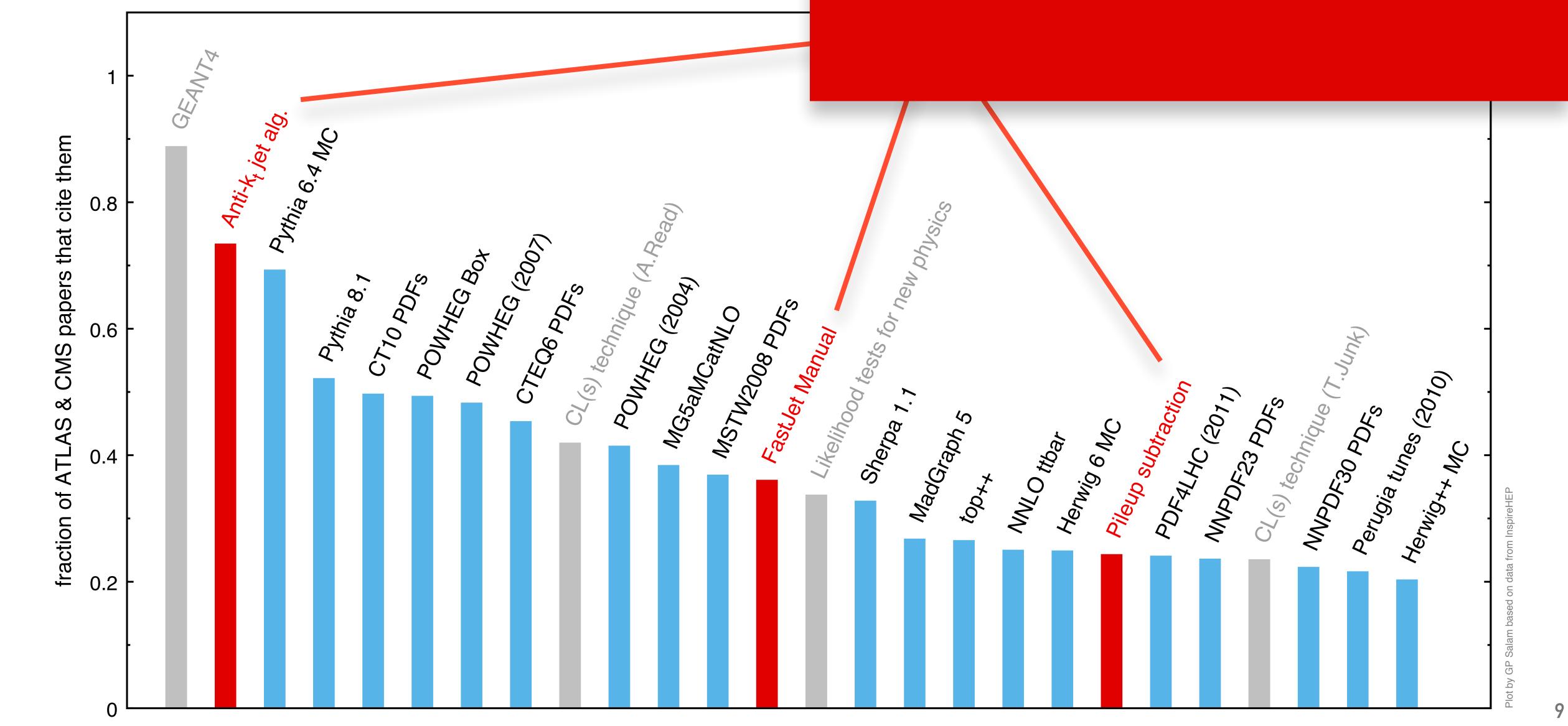


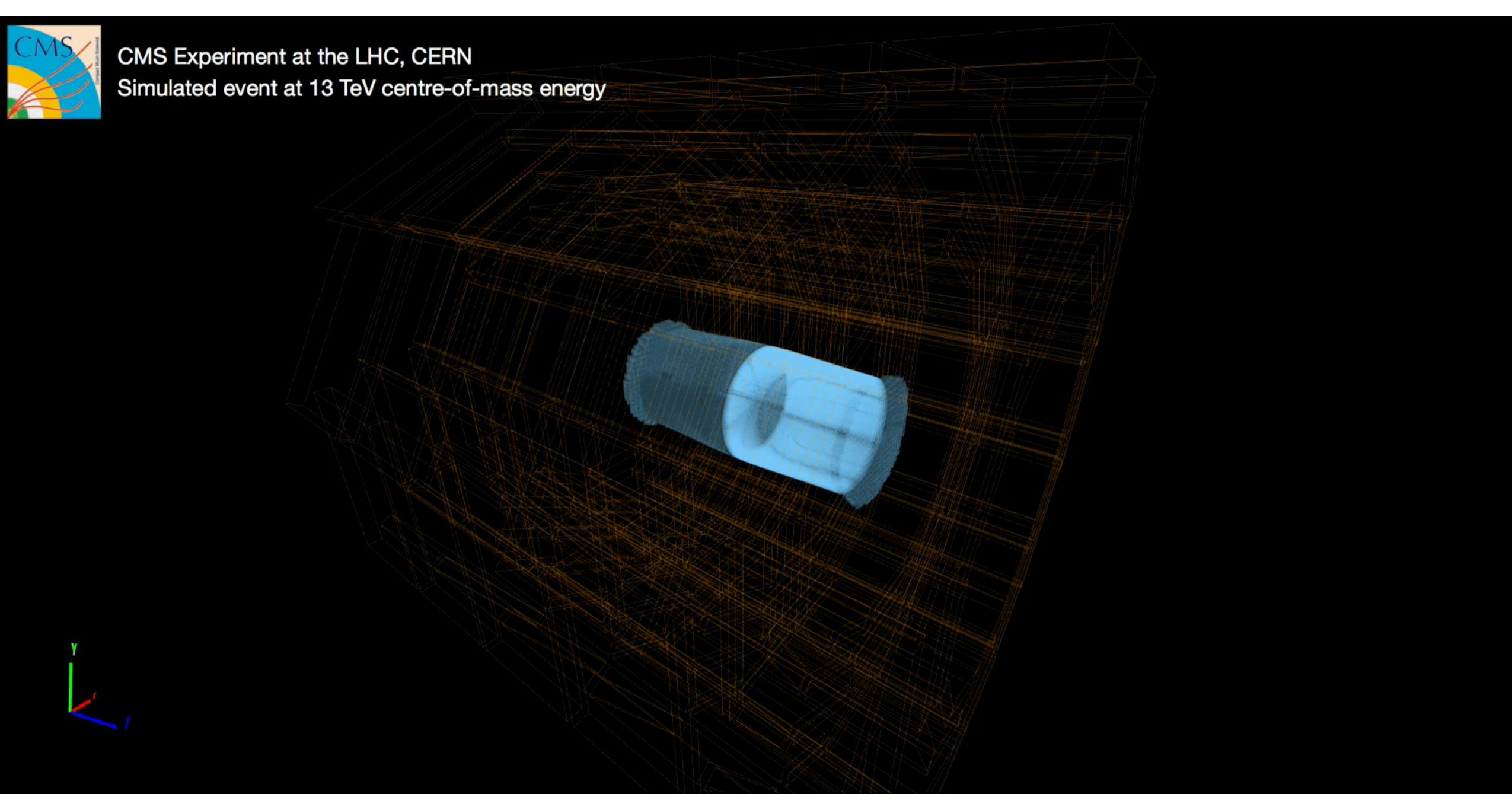


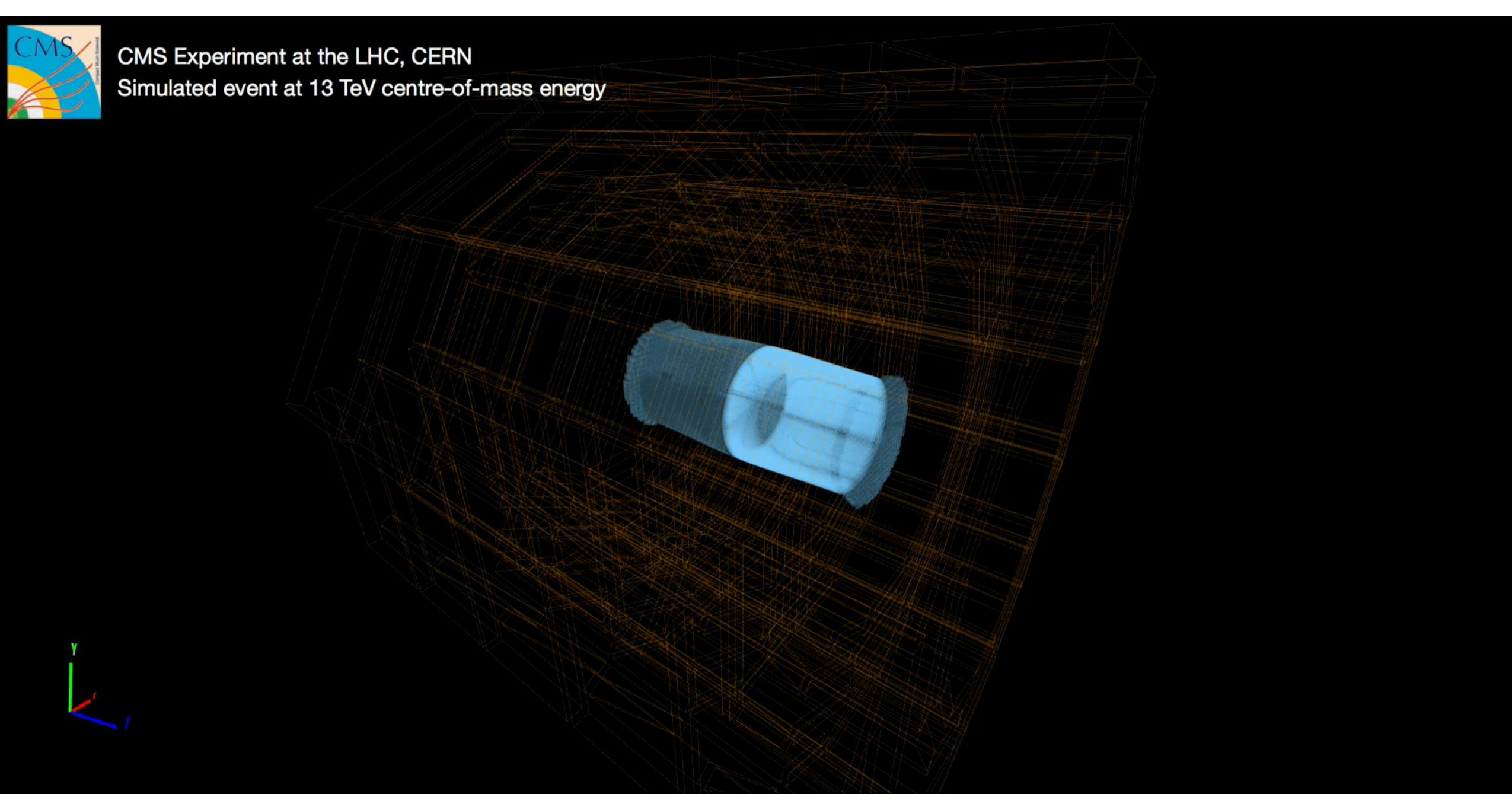
gluon-gluon collisions around the Higgs mass

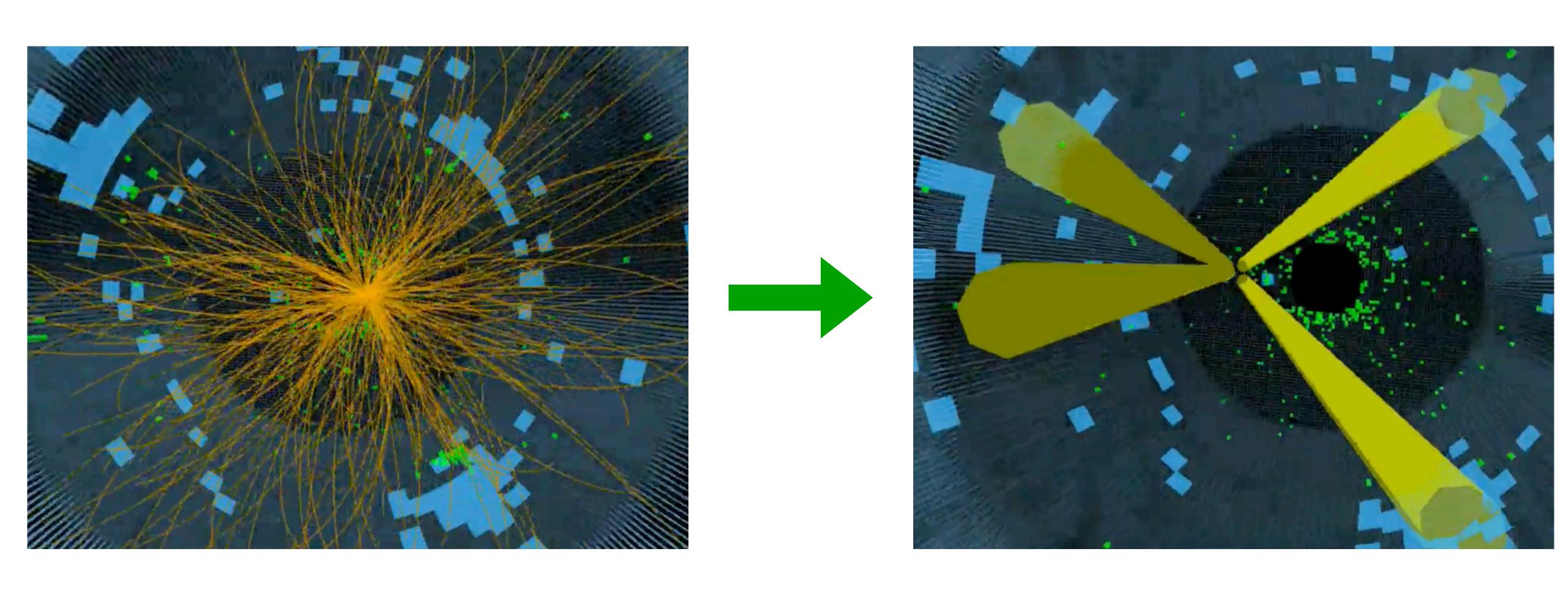


organising event information ("jets")

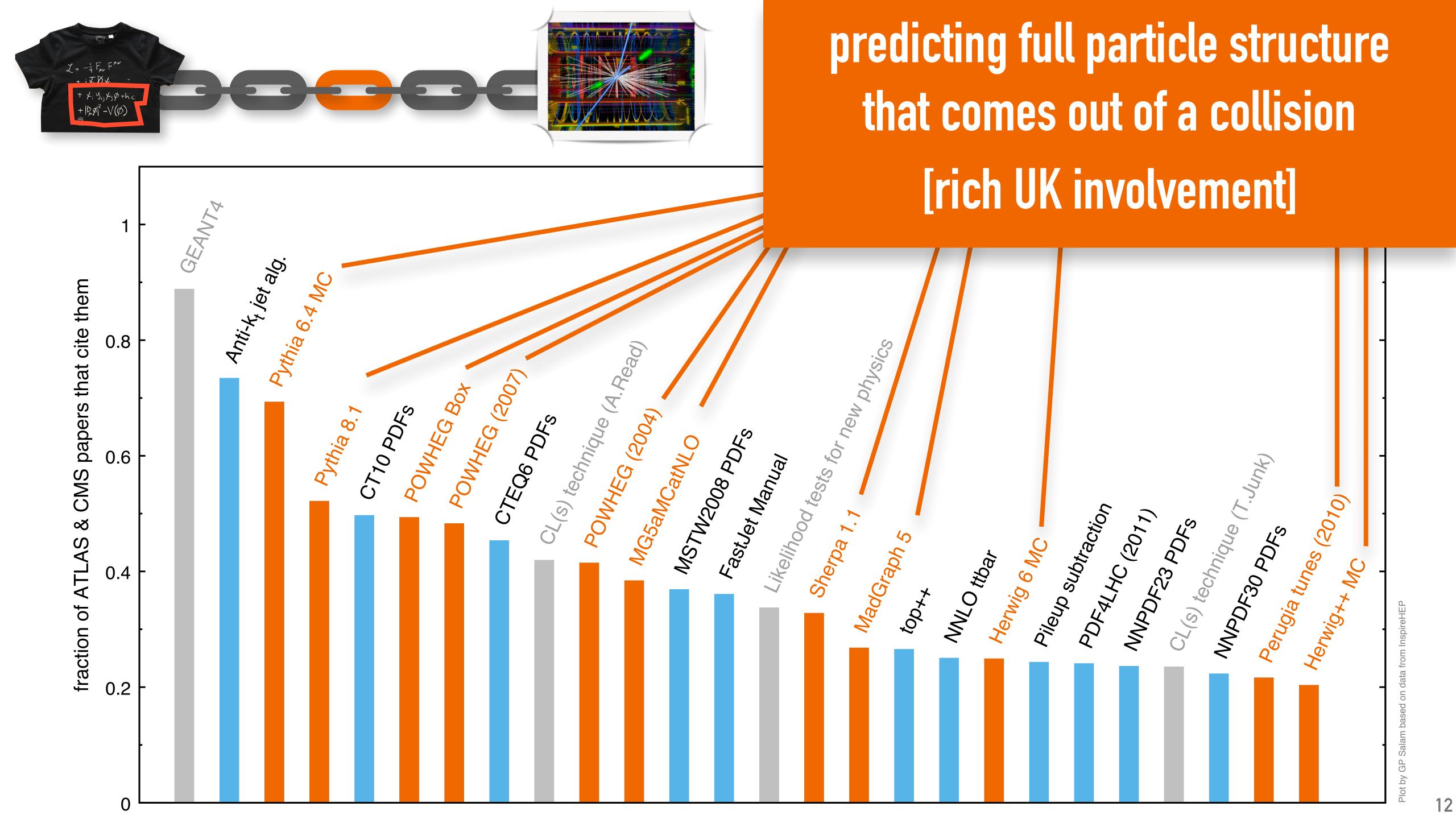






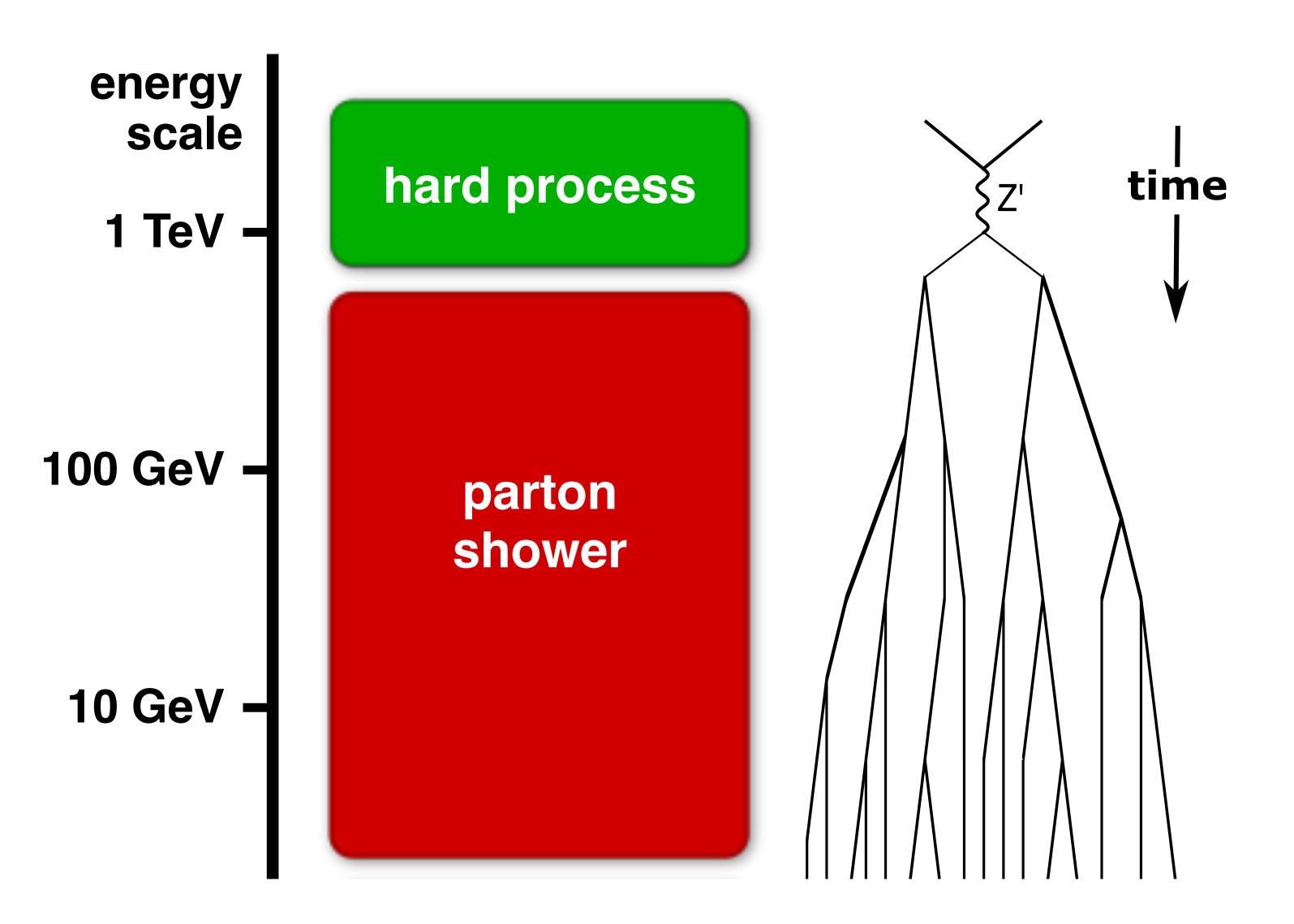


the question of organising information from hundreds of particles will come back later

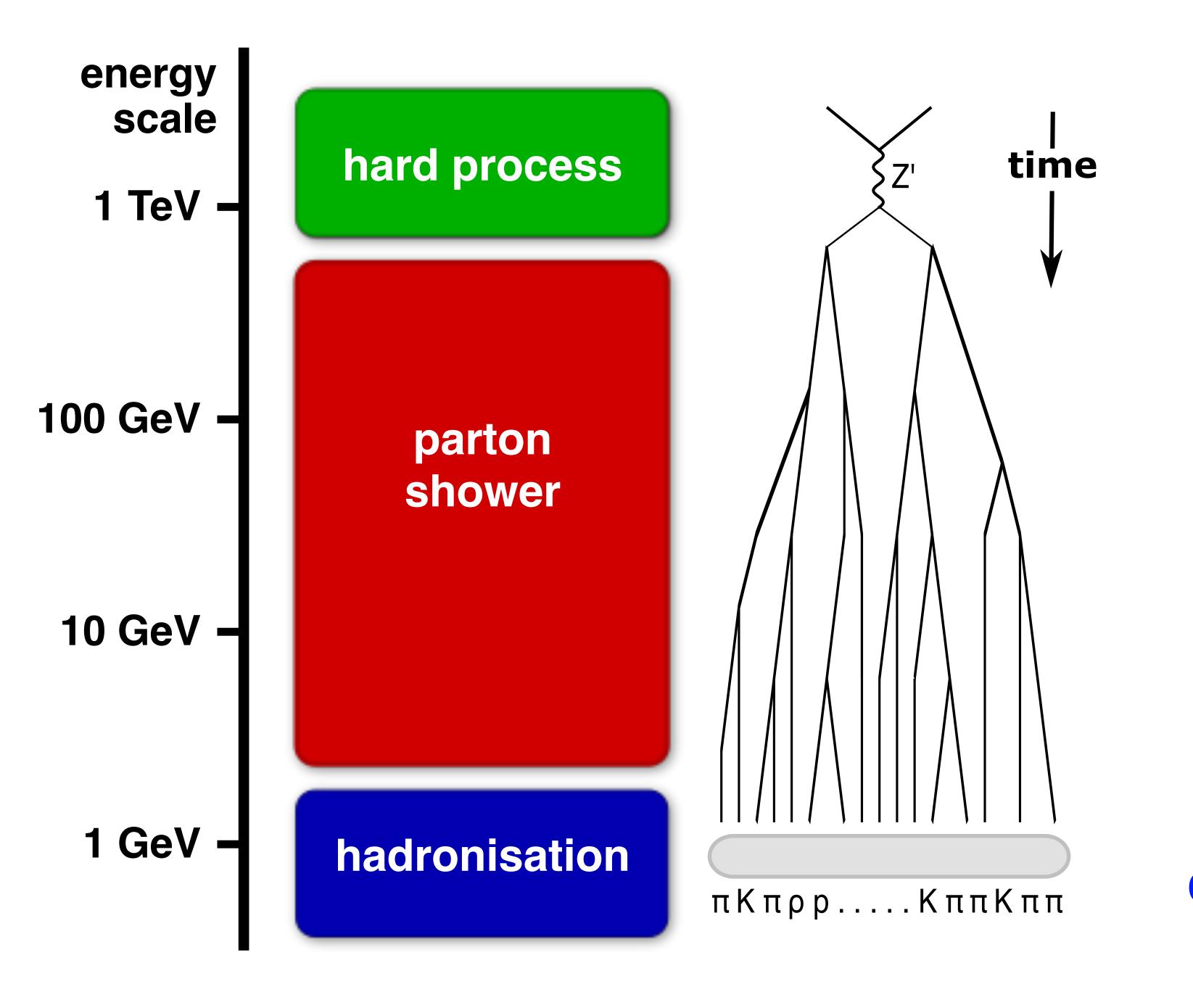




schematic view of key components of QCD predictions and Monte Carlo event simulation



schematic view of key components of QCD predictions and Monte Carlo event simulation



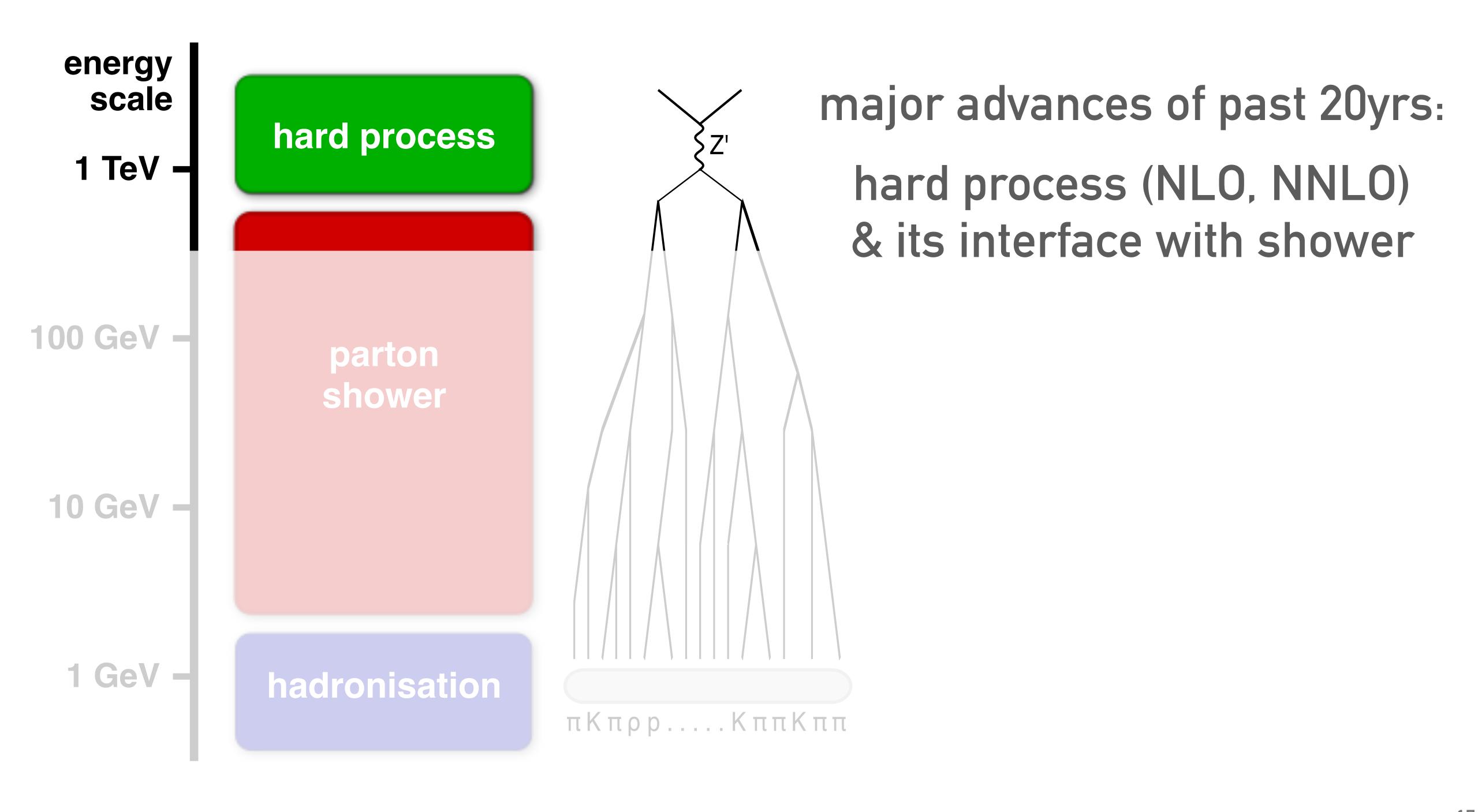
schematic view of key components of QCD predictions and Monte Carlo event simulation

pattern of particles in MC can be directly compared to pattern in experiment

general purpose Monte Carlo event generators: THE BIG 3



they do an amazing job of simulation vast swathes of data; collider physics would be unrecognisable without them



energy scale

1 TeV -

100 GeV -

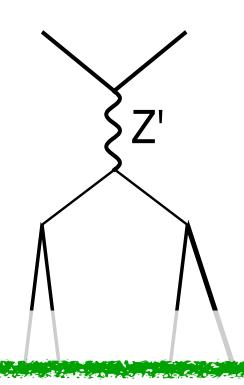
10 GeV -

1 GeV

hard process

parton shower

hadronisation



major advances of past 20yrs:

hard process (NLO, NNLO) & its interface with shower



MadGraph5_aMC@NLO

MC@NLO (in Herwig&Sherpa)

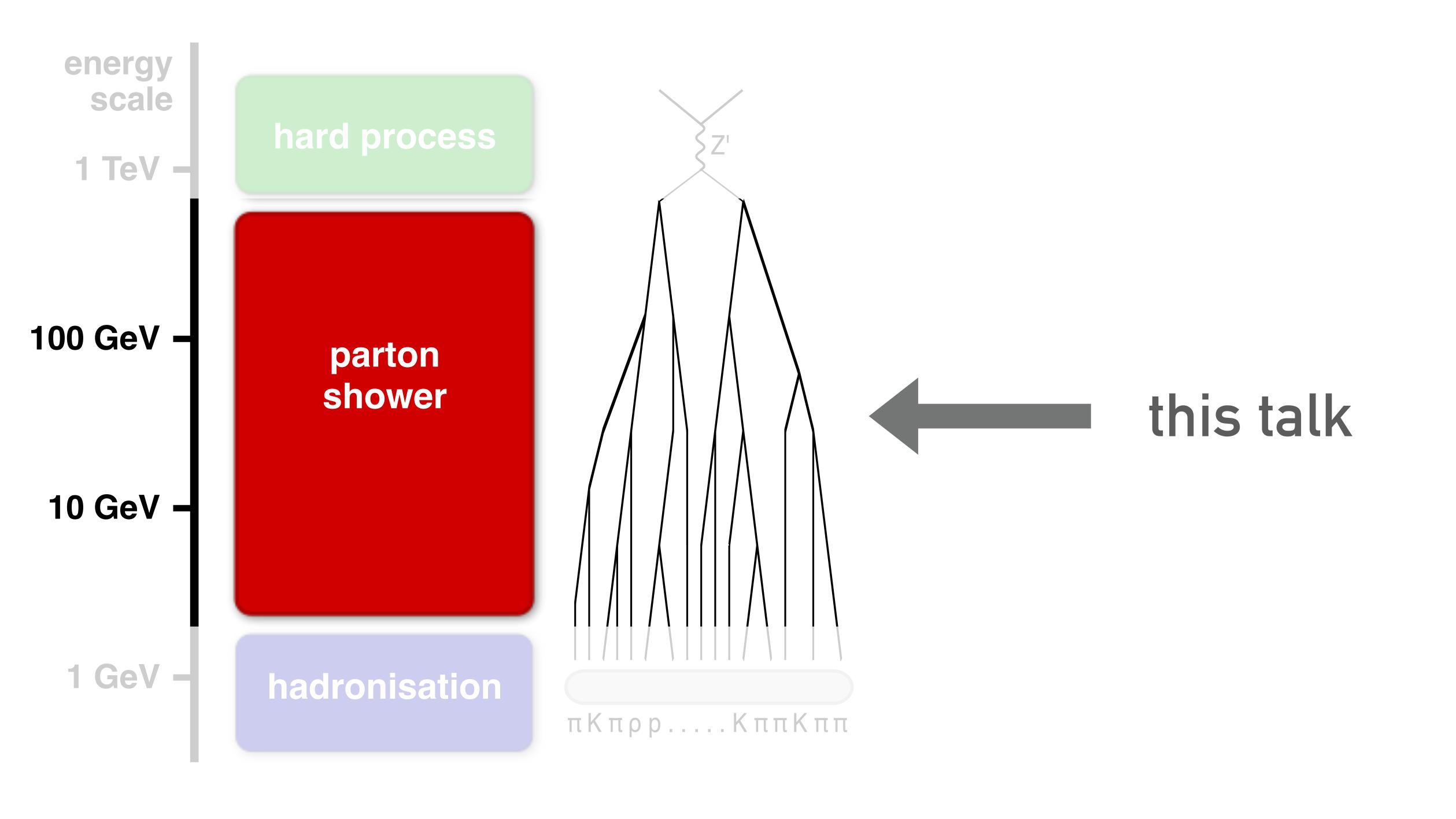
MINLO

UNNLOPS

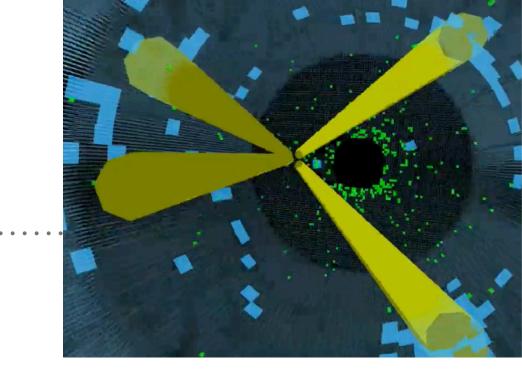
MLM, CKKW Vincia, FxFx

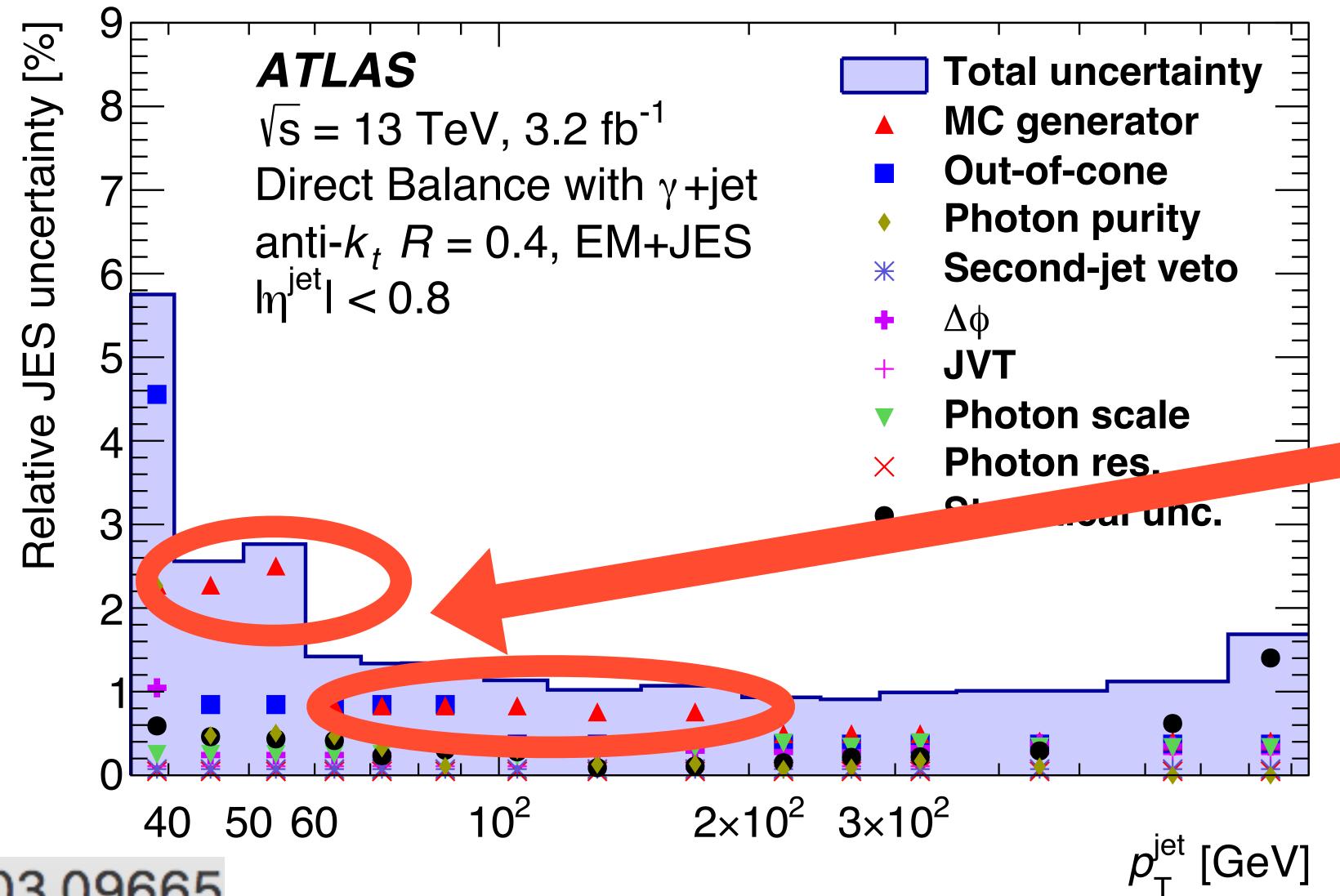






Fundamental experimental calibrations (jets)





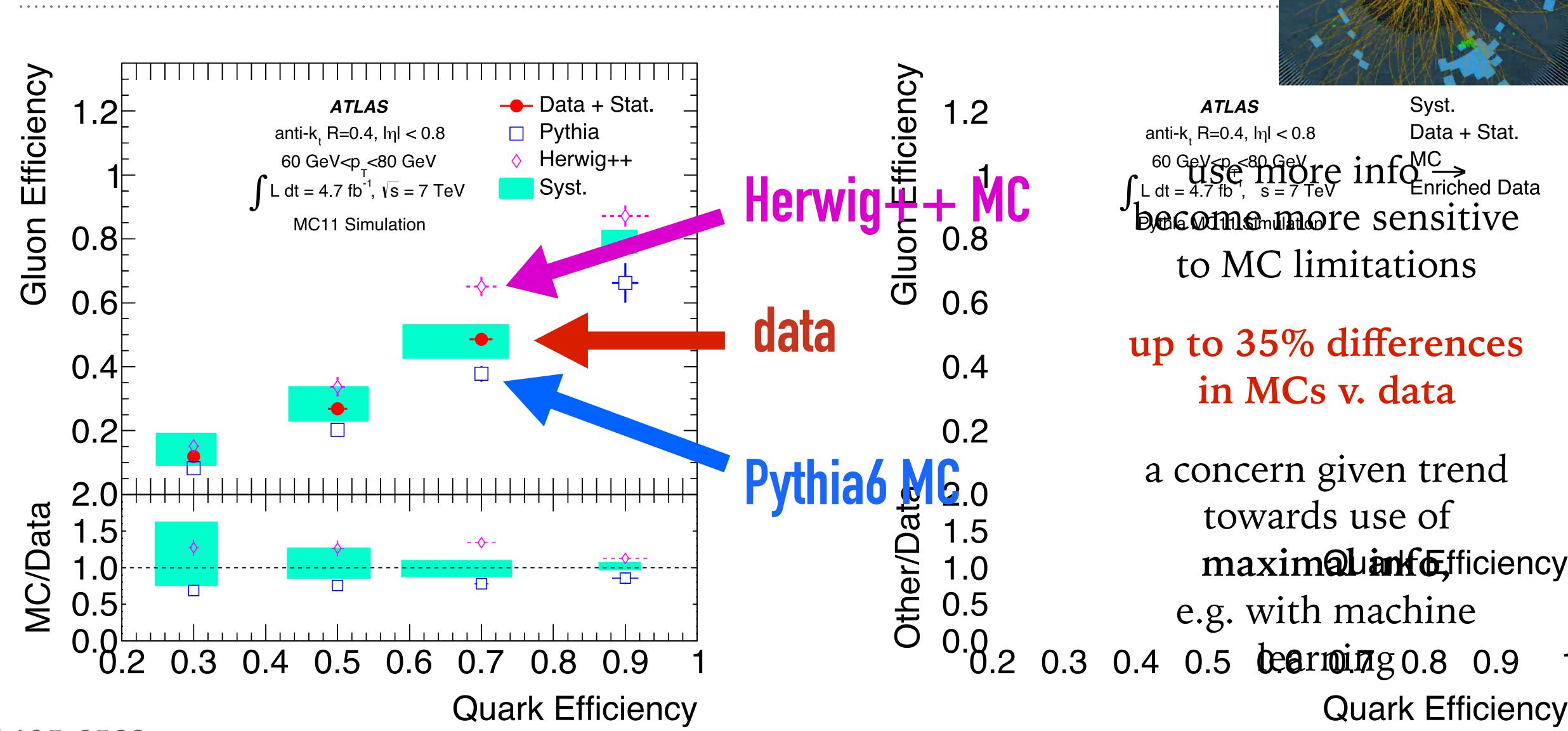
Jet energy scale, which feeds into hundreds of other measurements

Largest systematic errors (1–2%) come from differences between MC generators

(here Sherpa v. Pythia)

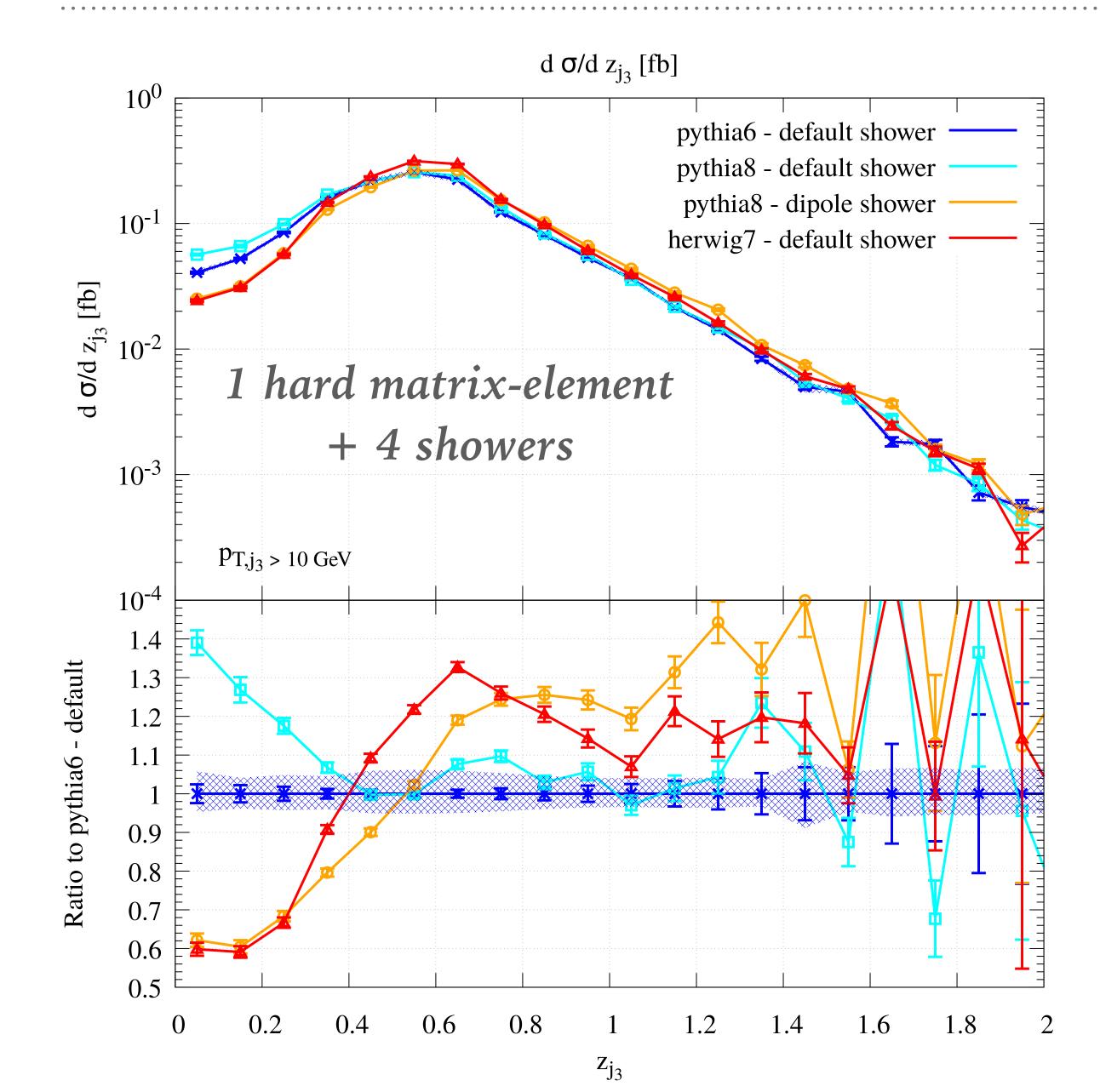
→ fundamental limit on LHC precision potential

using full event information (quark/gluon tagging)



1405.6583

Matching with hard process is hitting a limit (e.g. Jäger, Karlberg, Scheller 1812.05118)



Limits effectiveness of current matching methods (here POWHEG)

Parton structure also gets in way of better (NNLOPS) hard-process + shower matching schemes

what is a parton shower?

illustrate with dipole / antenna showers

Gustafson & Pettersson 1988, Ariadne 1992, main Sherpa & Pythia8 showers, option in Herwig7, Vincia shower & (partially) Deductor shower

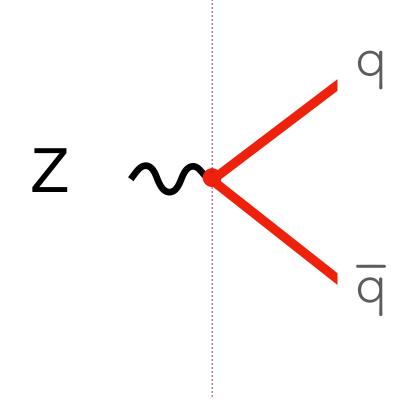
At its simplest

$$\sum_{n=0}^{\infty} \prod_{i=1}^{n} \left(\left\langle + \right\rangle \right) = \cdots$$

iteration of $2\rightarrow 3$ (or $1\rightarrow 2$) splitting kernel

Start with q-qbar state.

Evolve a step in v and throw a random number to decide if state remains unchanged



V

$$\frac{dP_2(v)}{dv} = -f_{2\to 3}^{q\bar{q}}(v) P_2(v)$$

V

Start with q-qbar state.

Evolve a step in v and throw a random number to decide if state remains unchanged

$$\frac{dP_2(v)}{dv} = -f_{2\to 3}^{q\bar{q}}(v) P_2(v)$$

V0

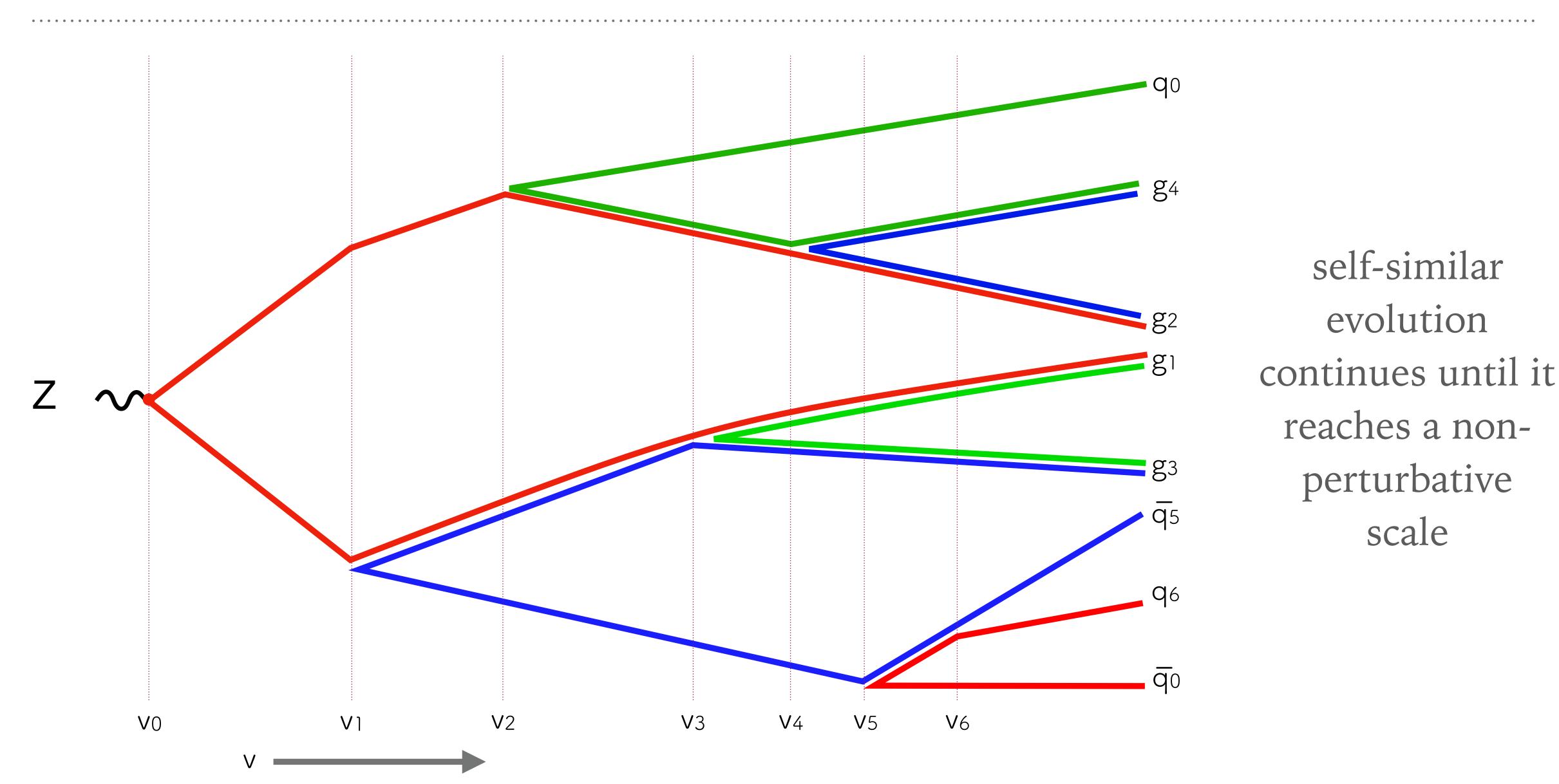
Start with q-qbar state.

Evolve a step in v and throw a random number to decide if state remains unchanged

At some point, rand.numb. is such that **state splits** $(2\rightarrow 3$, i.e. emits gluon). Evolution equation changes

$$\frac{dP_3(v)}{dv} = -\left[f_{2\to 3}^{qg}(v) + f_{2\to 3}^{g\bar{q}}(v)\right] P_3(v)$$

gluon is part of two dipoles $(qg, \bar{q}g)$



recent directions of parton-shower work?

- 1. including $2 \rightarrow 4$ (or $1 \rightarrow 3$) splittings
- 2. subleading colour corrections (dipole picture is large N_C)
- 3. EW showers

Including $1 \rightarrow 3$ splittings ($\equiv 2 \rightarrow 4$)

- ➤ Li & Skands, 1611.00013
- ➤ Jadach et al, e.g. 1504.06849, 1606.01238 ➤ Höche, Krauss & Prestel, 1705.00982,

 Höche & Prestel, 1705.00742, Dulat, Höche & Prestel, 1805.03757

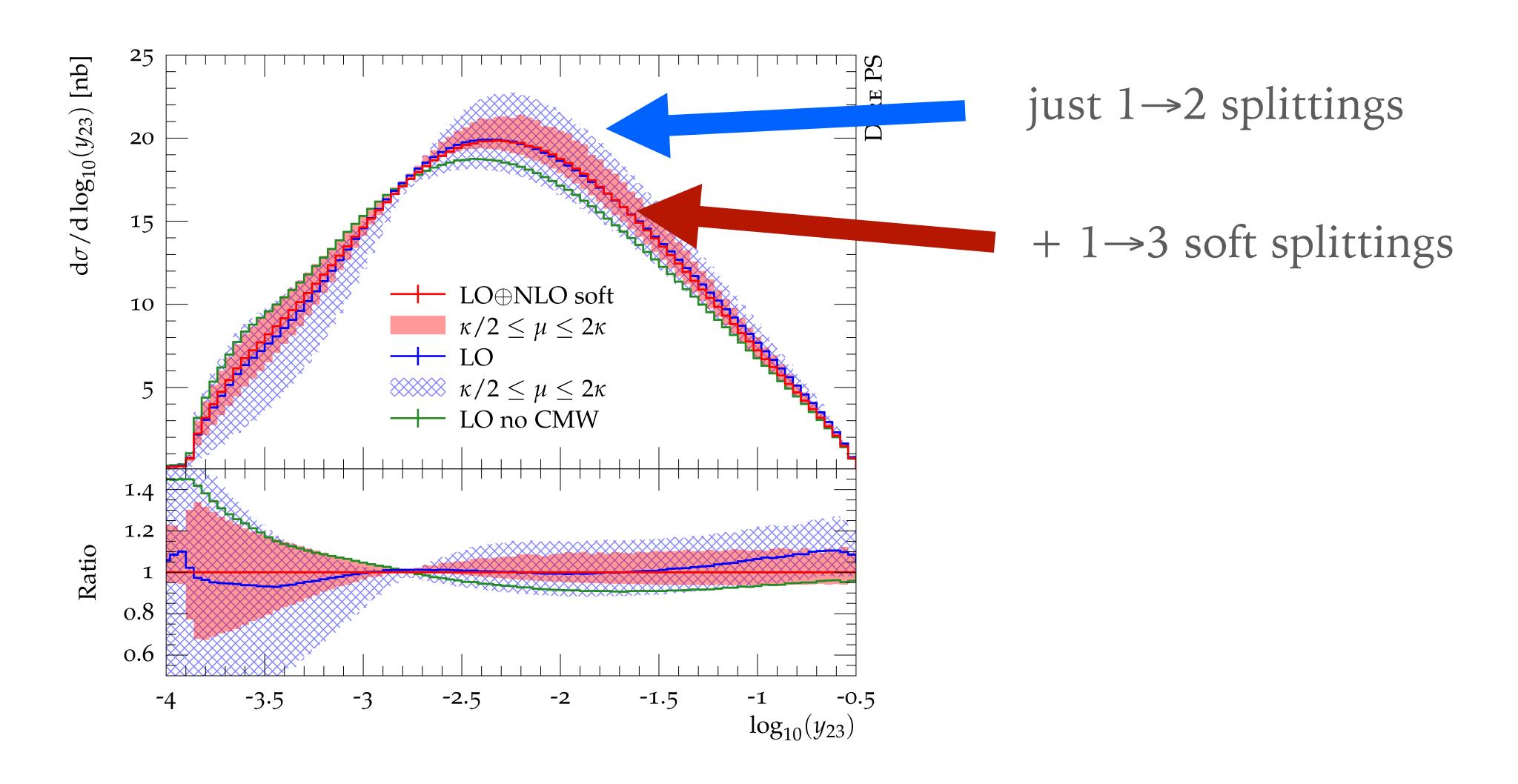
$$D_{ji}^{(0)}(z,\mu) = \delta_{ij}\delta(1-z) \qquad \leftrightarrow \qquad \qquad \downarrow_{j} \qquad z \qquad / \qquad \downarrow_{i} \qquad 1$$

$$D_{ji}^{(1)}(z,\mu) = -\frac{1}{\varepsilon}P_{ji}^{(0)}(z) \qquad \leftrightarrow \qquad \downarrow_{i} \qquad \downarrow_{z} \qquad / \qquad \downarrow_{i} \qquad 1$$

$$D_{ji}^{(2)}(z,\mu) = -\frac{1}{2\varepsilon}P_{ji}^{(1)}(z) + \frac{\beta_{0}}{4\varepsilon^{2}}P_{ji}^{(0)}(z) + \frac{1}{2\varepsilon^{2}}\int_{z}^{1}\frac{\mathrm{d}x}{x}P_{jk}^{(0)}(x)P_{ki}^{(0)}(z/x)$$

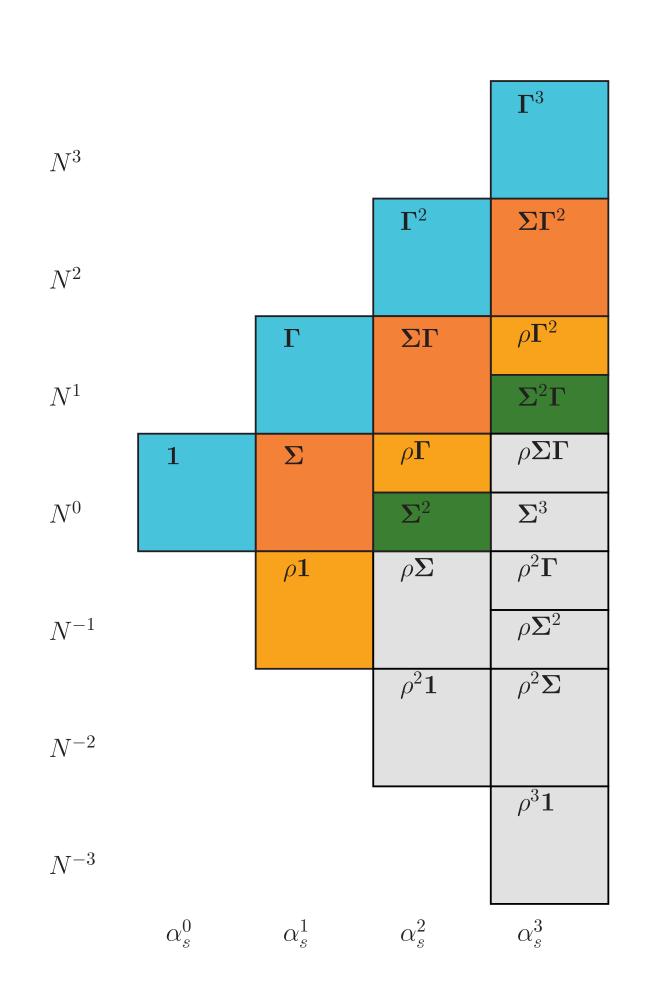
$$\leftrightarrow \left(\bigvee_{i} \bigvee_{j} \sum_{z} + \bigvee_{i} \bigvee_{j} \bigvee_{z} \right) / \bigvee_{i} \bigvee_{i} \bigvee_{j} \bigvee_{i} \bigvee_{j} \bigvee_{j} \bigvee_{i} \bigvee_{j} \bigvee_{j} \bigvee_{i} \bigvee_{j} \bigvee_{i} \bigvee_{j} \bigvee_{i} \bigvee_{j} \bigvee_{j} \bigvee_{i} \bigvee_{j} \bigvee_{i}$$

Including 1 → 3 splittings



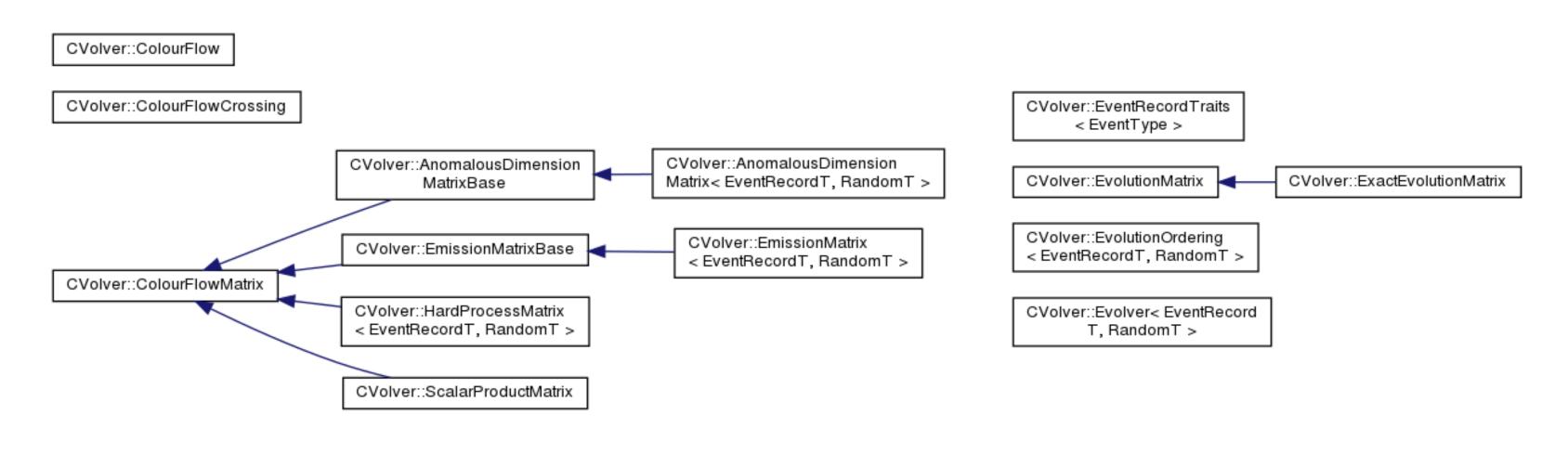
Dulat, Höche & Prestel, 1805.03757

Hierarchy of subleading colour corrections



Angeles, De Angelis, Forshaw, Plätzer, Seymour – JHEP 05 (2018) 044 Gieseke, Kirchgaesser, Plätzer, Siodmok – arXiv:1808.06770 De Angelis, Forshaw, Plätzer – arXiv:181y.xxxxx Forshaw, Holguin, Plätzer – arXiv:181y.xxxxx Plätzer, Sjödahl, Thorén, arXiv:1808.00332

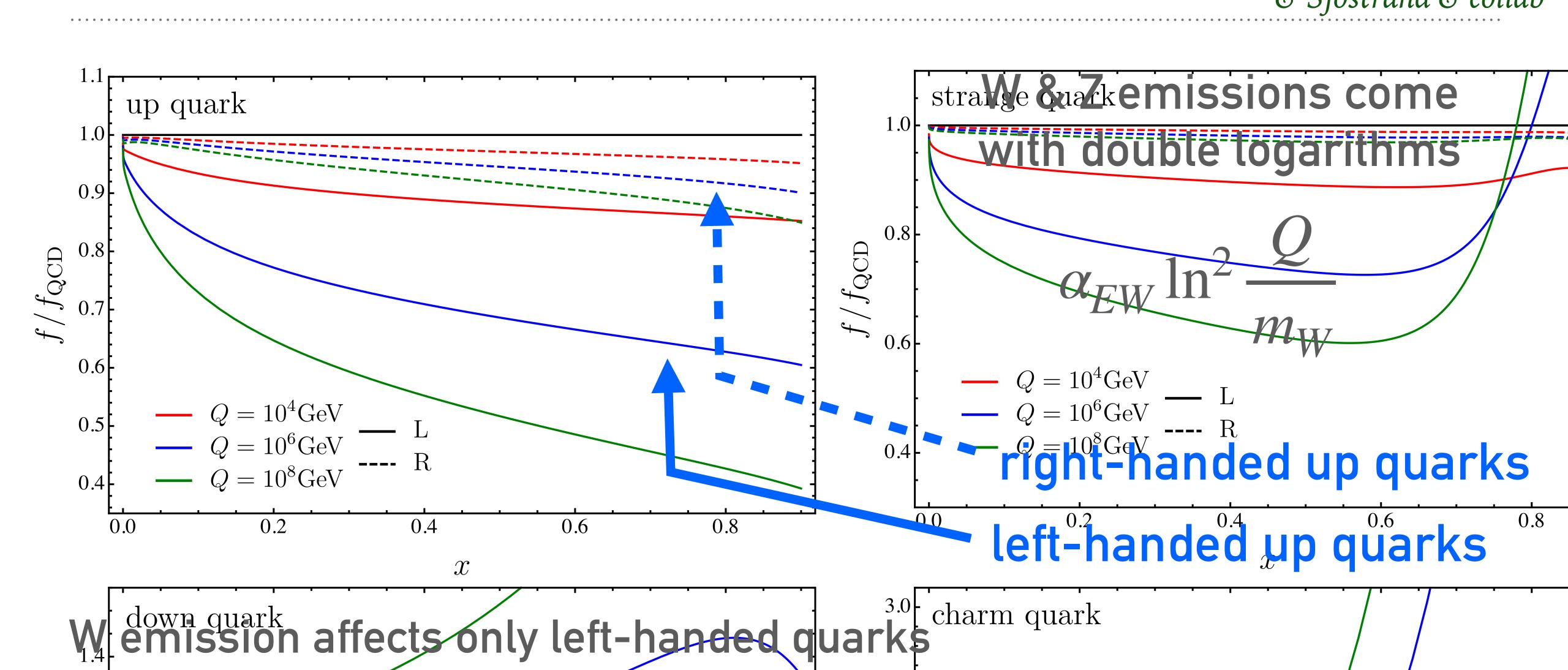
$$\mathbf{A}_n(E) = \mathbf{V}(E, E_n) \mathbf{D}_n \mathbf{A}_{n-1}(E_n) \mathbf{D}_n^{\dagger} \mathbf{V}^{\dagger}(E, E_n) \theta(E - E_n)$$



Plugin approach can accommodate anything from (N)GLs to full parton showers.

cf. also work by Hatta & Ueda, 1304.6930; Nagy & Soper papers; some subleading colour also in DIRE2 work

EW showers (esp. beyond LHC)



strong polarisation of quarks in unpolarised proton (at high enough energies)

32

what does a parton shower achieve?

not just a question of ingredients, but also the final result of assembling them together

Dasgupta, Dreyer, Hamilton, Monni & GPS, 1805.09327

what should a parton shower achieve?

not just a question of ingredients, but also the final result of assembling them together

Dasgupta, Dreyer, Hamilton, Monni & GPS, 1805.09327

it's a complicated issue...

➤ For a total cross section, e.g. for Higgs production, it's easy to talk about systematic improvements (LO, NLO, NNLO, ...). But they're restricted to that one observable

it's a complicated issue...

- ➤ For a total cross section, e.g. for Higgs production, it's easy to talk about systematic improvements (LO, NLO, NNLO, ...). But they're restricted to that one observable
- ➤ With a parton shower (+hadronisation) you produce a "realistic" full set of particles. You can ask questions of arbitrary complexity:
 - > the multiplicity of particles
 - > the total transverse momentum with respect to some axis (broadening)
 - ➤ the angle of 3rd most energetic particle relative to the most energetic one [machine learning might "learn" many such features]

it's a complicated issue...

- ➤ For a total cross section, e.g. for Higgs production, it's easy to talk about systematic improvements (LO, NLO, NNLO, ...). But they're restricted to that one observable
- ➤ With a parton shower (+hadronisation) you produce a "realistic" full set of particles. You can ask questions of arbitrary complexity:
 - > the multiplicity of particles
 - > the total transverse momentum with respect to some axis (broadening)
 - ➤ the angle of 3rd most energetic particle relative to the most energetic one [machine learning might "learn" many such features]

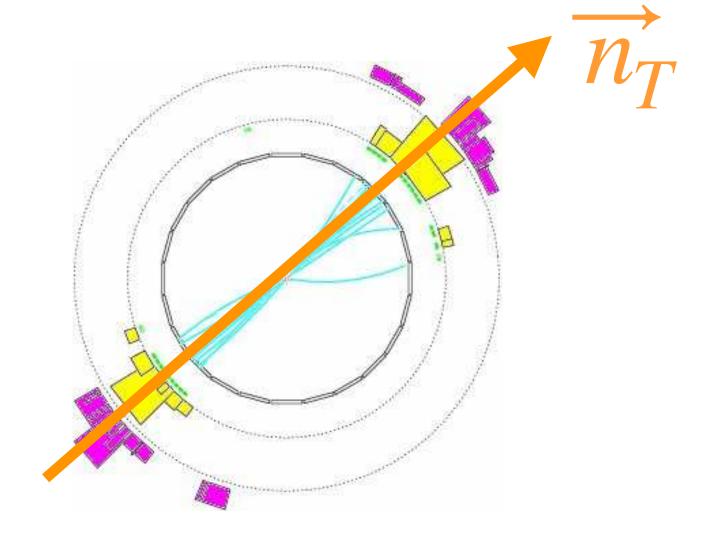
how can you prescribe correctness & accuracy of the answer, when the questions you ask can be arbitrary?

The standard answer so far

It's common to hear that showers are Leading Logarithmic (LL) accurate.

That language, widespread for multiscale problems, comes from analytical resummations. E.g. for (famous) "Thrust"

$$T = \max_{\vec{n}_T} \frac{\sum_{i} |\vec{p}_{i}.\vec{n}_{T}|}{\sum_{i} |\vec{p}_{i}|}$$



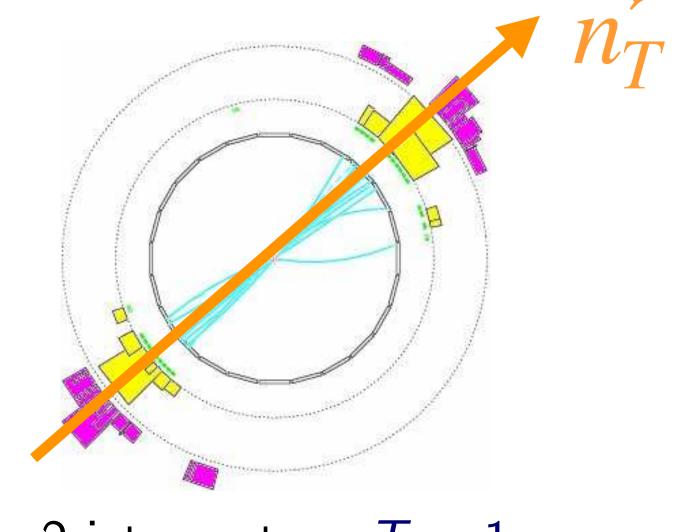
2-jet event: $T \simeq 1$

The standard answer so far

It's common to hear that showers are Leading Logarithmic (LL) accurate.

That language, widespread for multiscale problems, comes from analytical resummations. E.g. for (famous) "Thrust"

$$T = \max_{\vec{n}_T} \frac{\sum_{i} |\vec{p}_{i}.\vec{n}_{T}|}{\sum_{i} |\vec{p}_{i}|}$$



2-jet event: $T \simeq 1$

$$\sigma(1-T< e^{-L}) = \sigma_{tot} \exp\left[Lg_1(\alpha_s L) + g_2(\alpha_s L) + \alpha_s g_3(\alpha_s L) + \alpha_s^2 g_4(\alpha_s L) + \cdots\right]$$

$$[\alpha_s \ll 1, L \gg 1]$$
LL
NNLL
NNLL
N3LL

Catani, Trentadue, Turnock & Webber '93

Becher & Schwartz '08 -

The standard answer so far

Sometimes you may see statements like "Following standard practice to improve the logarithmic accuracy of the parton shower, the soft enhanced term of the splitting functions is rescaled by $1+a_s(t)/(2\pi)K$ "

Questions:

- 1) Which is it? LL or better?
- 2) For what known observables does this statement hold?
- 3) What good is it to know that some handful of observables is LL (or whatever) when you want to calculate arbitrary observables?
- 4) Does LL even mean anything when you do machine learning?
- 5) Why only "LL" when analytic resummation can do so much better?

Our proposal for "minimal" criteria for a shower

Resummation

Establish logarithmic accuracy for all known classes of resummation:

- > global event shapes (thrust, broadening, angularities, jet rates, energy-energy correlations, ...)
- non-global observables (cf. Banfi, Corcella & Dasgupta, hep-ph/0612282)
- ➤ fragmentation / parton-distribution functions
- > (multiplicity, cf. original Herwig angular-ordered shower from 1980's)

Matrix elements

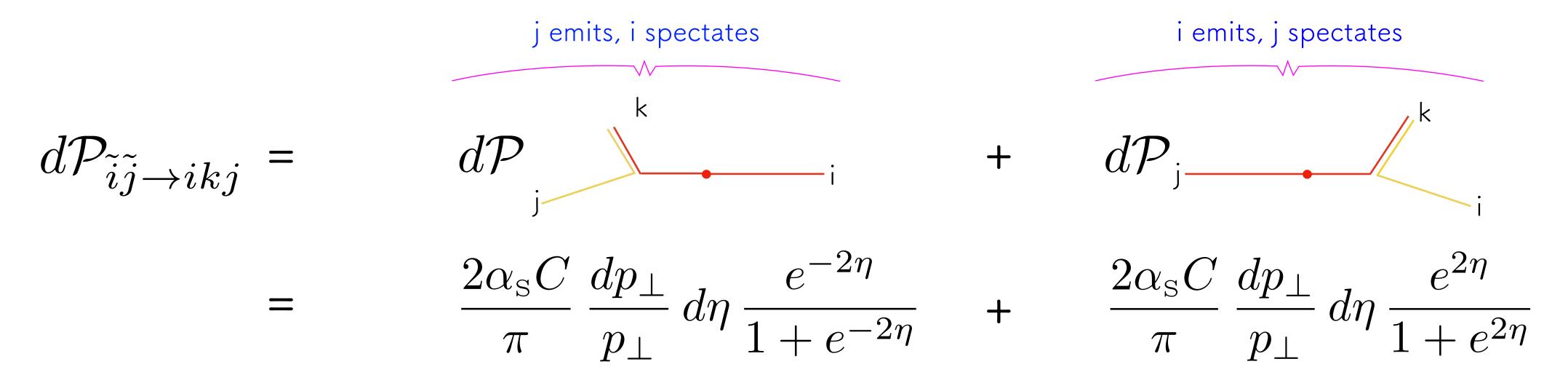
Establish in what sense iteration of (e.g. $2\rightarrow 3$) splitting kernel reproduces *N*-particle tree-level matrix elements *for any N*.

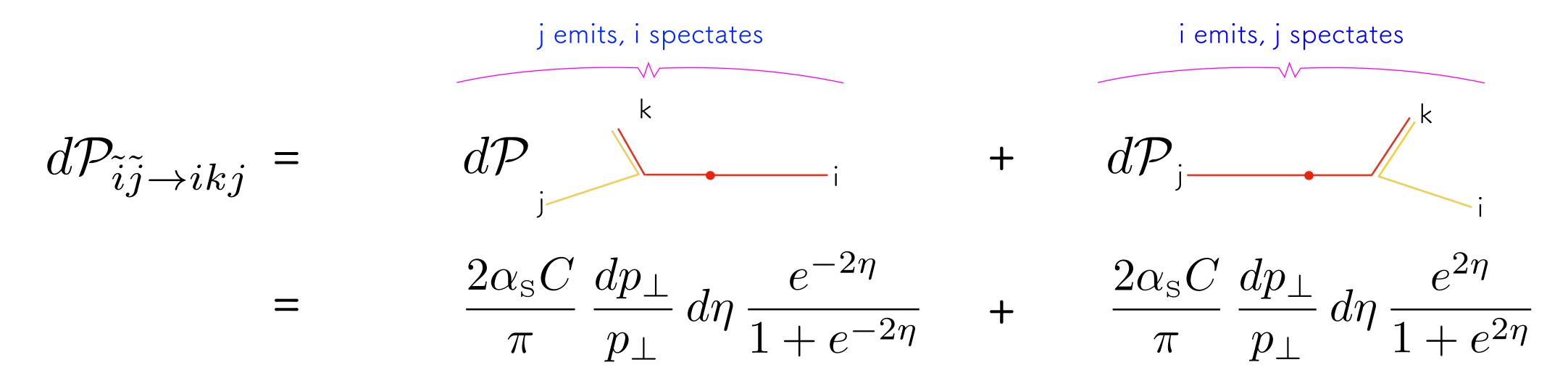
Examine two showers

- > Pythia8 shower: because it's the most widely used
- ▶ DIRE shower (2015 version, with just 2→3 splitting), because it's unique in being available for two General Purpose MC programs (Pythia8 & Sherpa2)

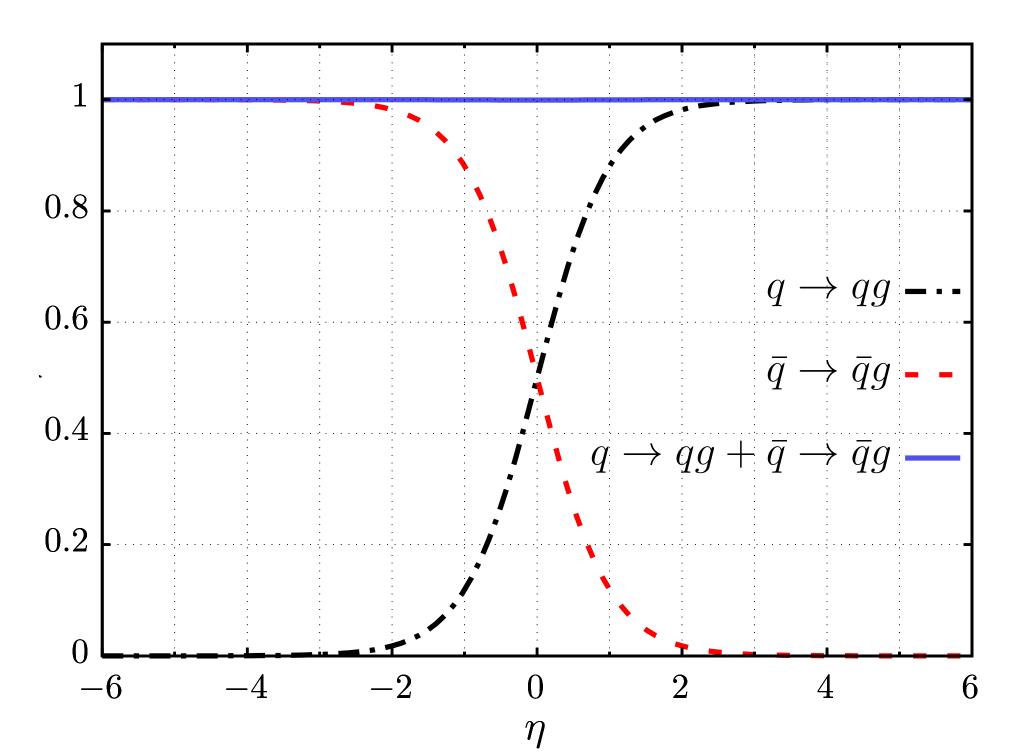
The results I'll talk about will be the same for both

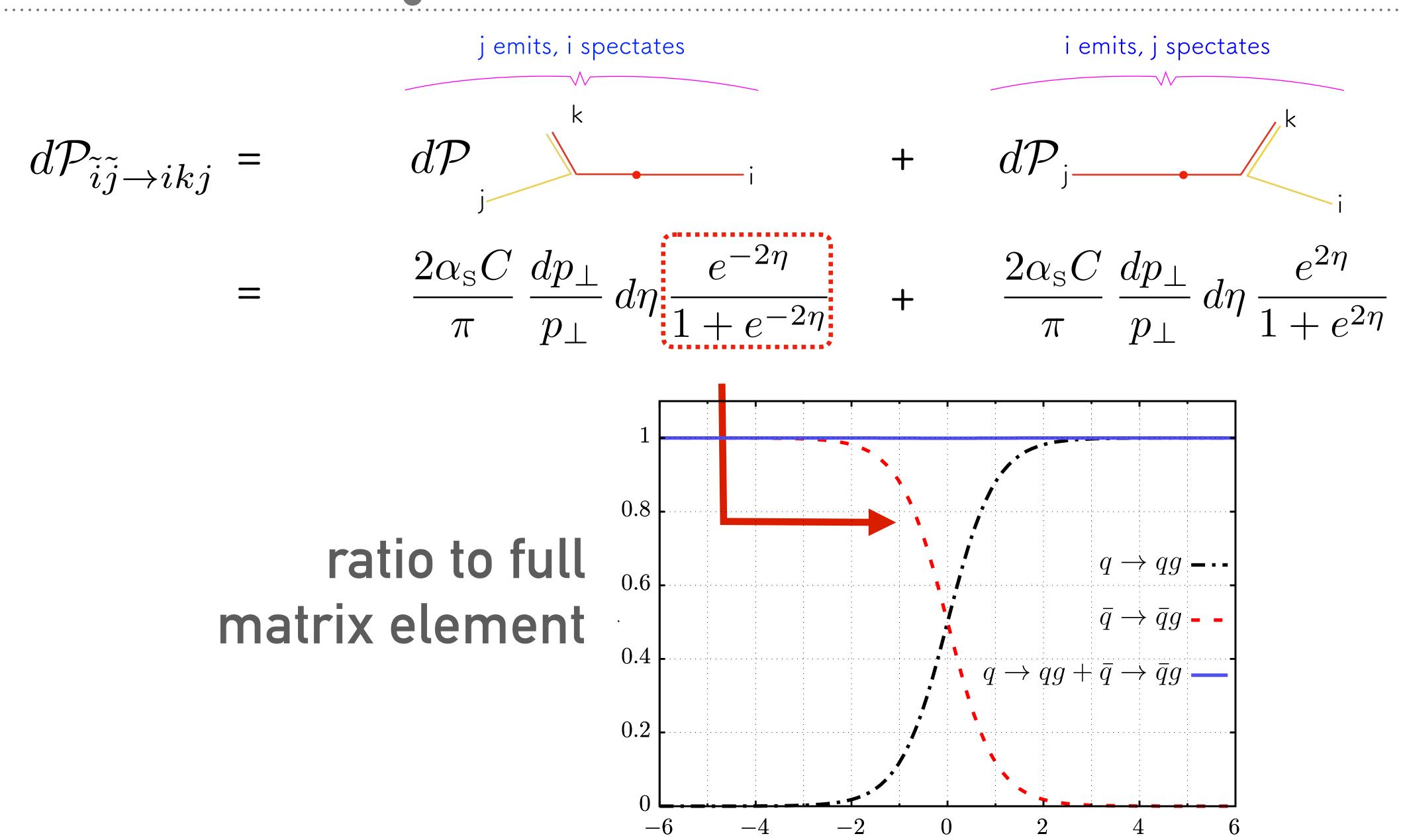
and they'll be limited to fixed order for simplicity (though it's easy enough to generalise to an all-order study)

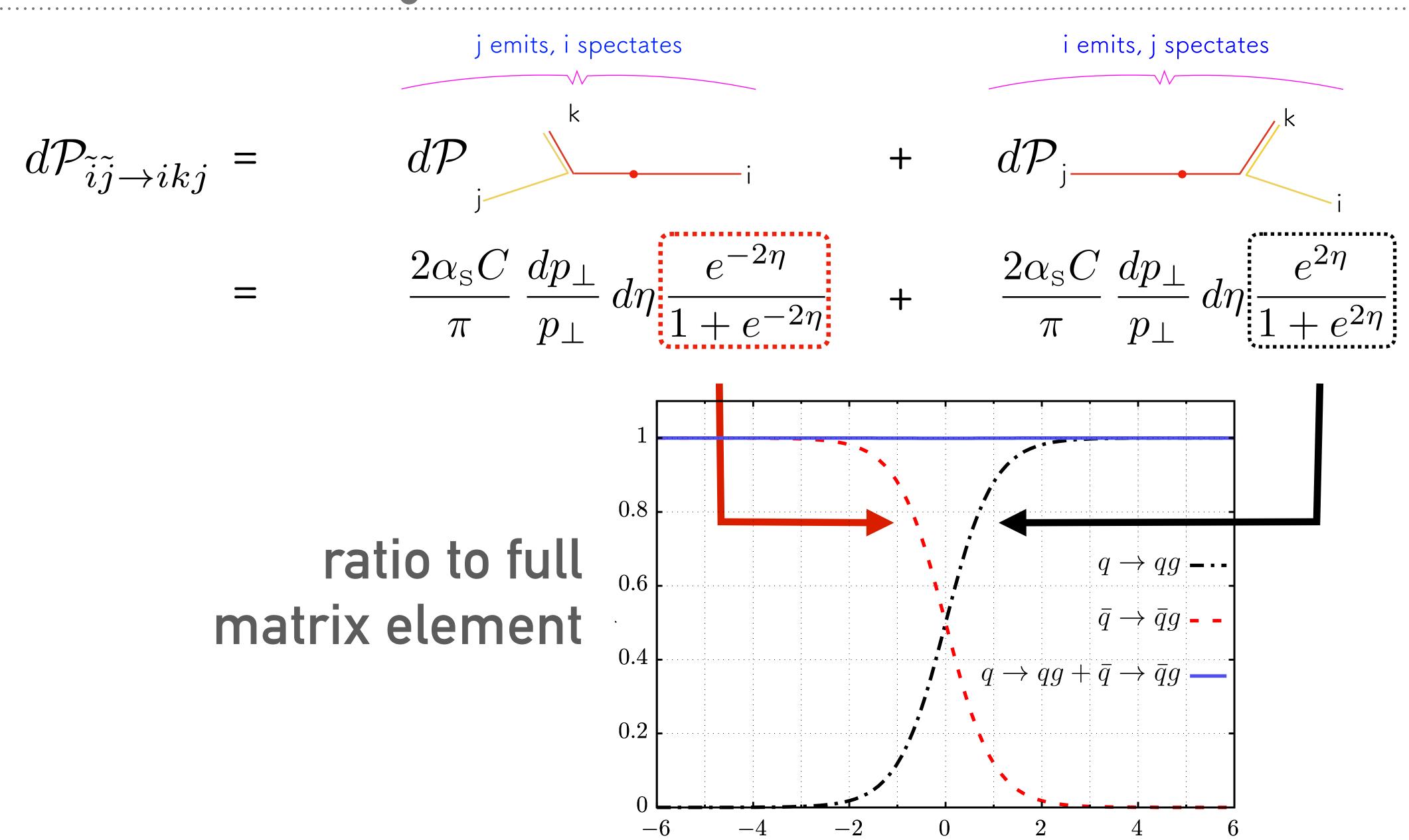


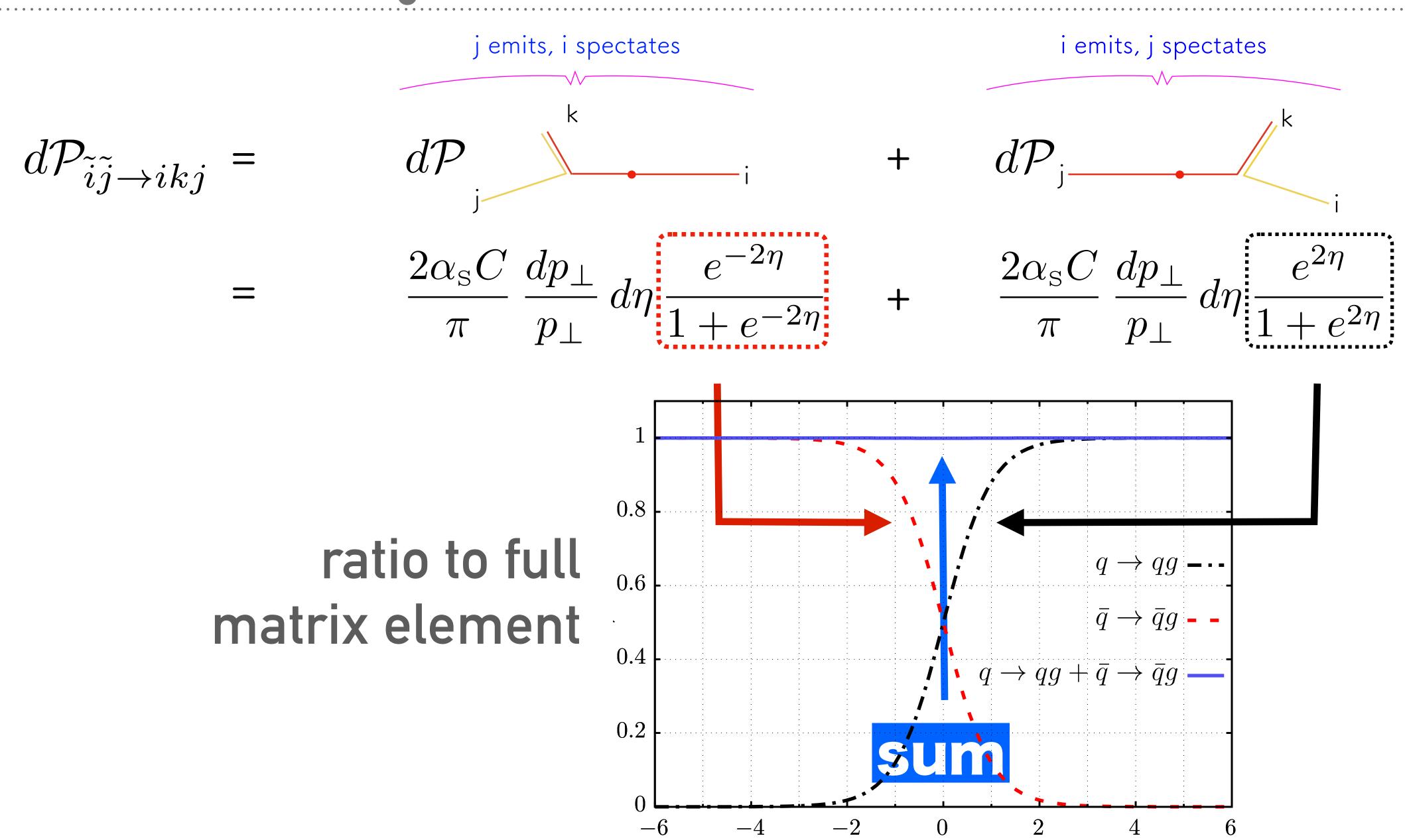


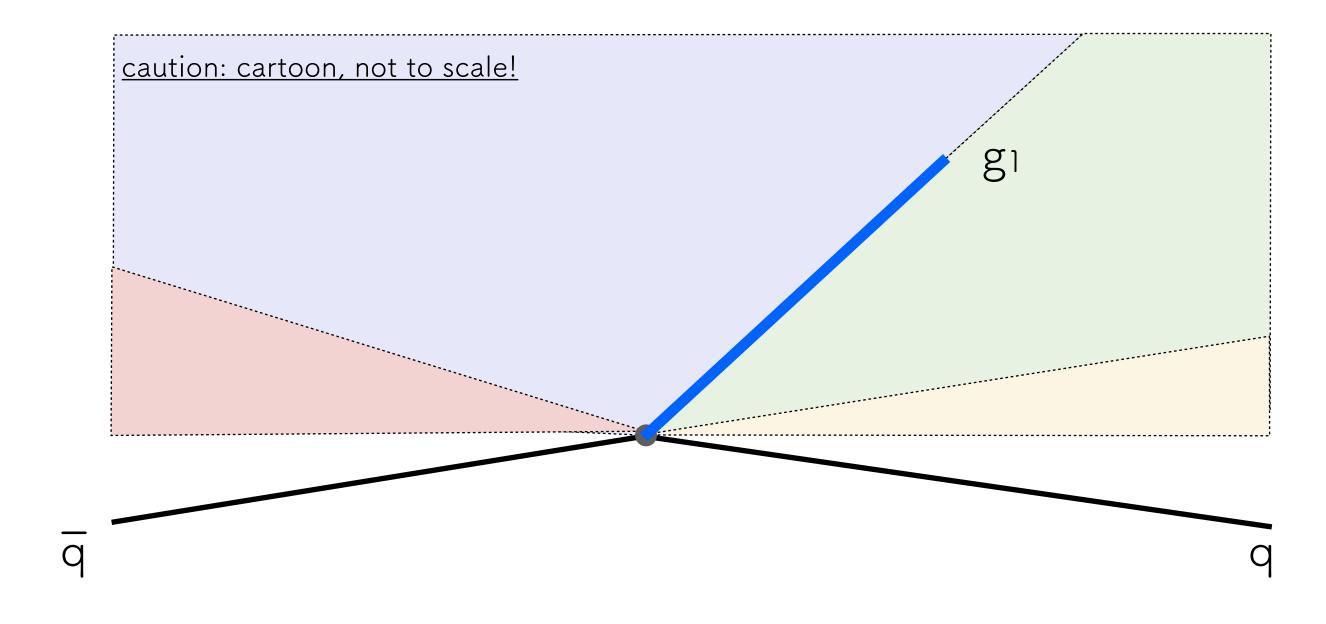
ratio to full matrix element

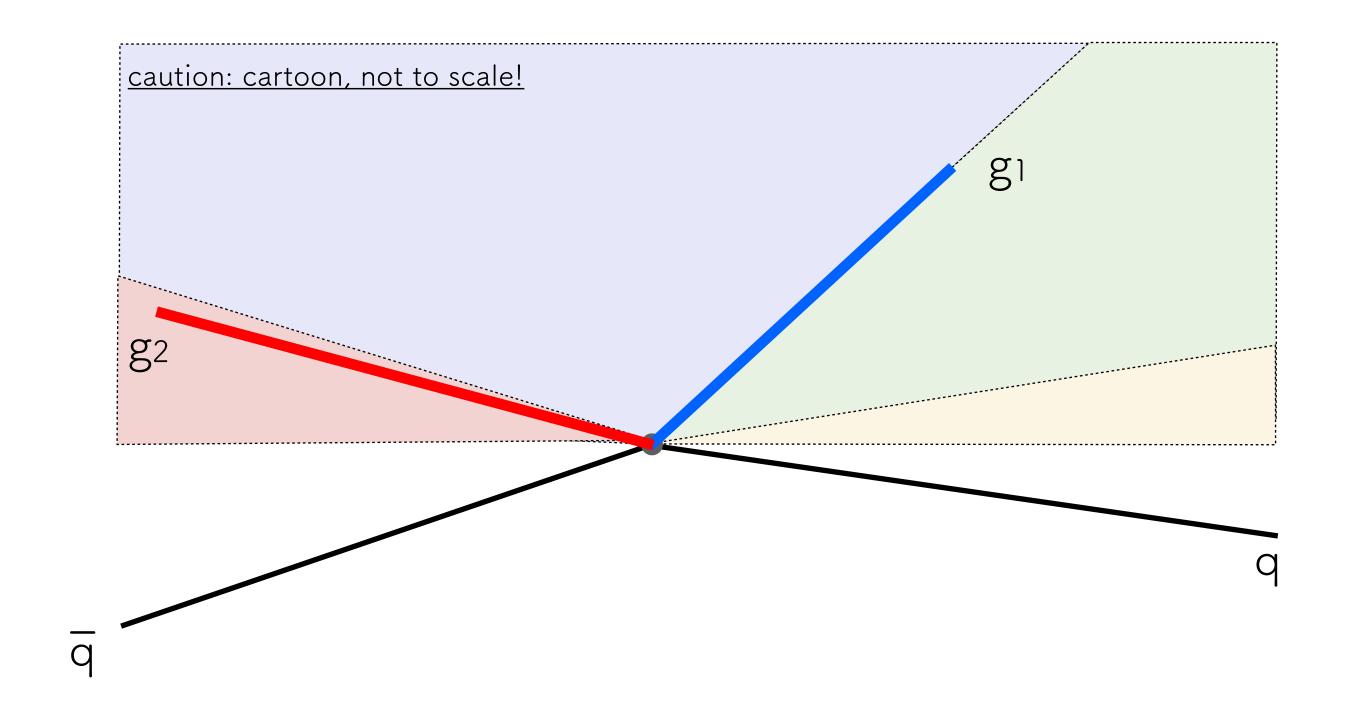


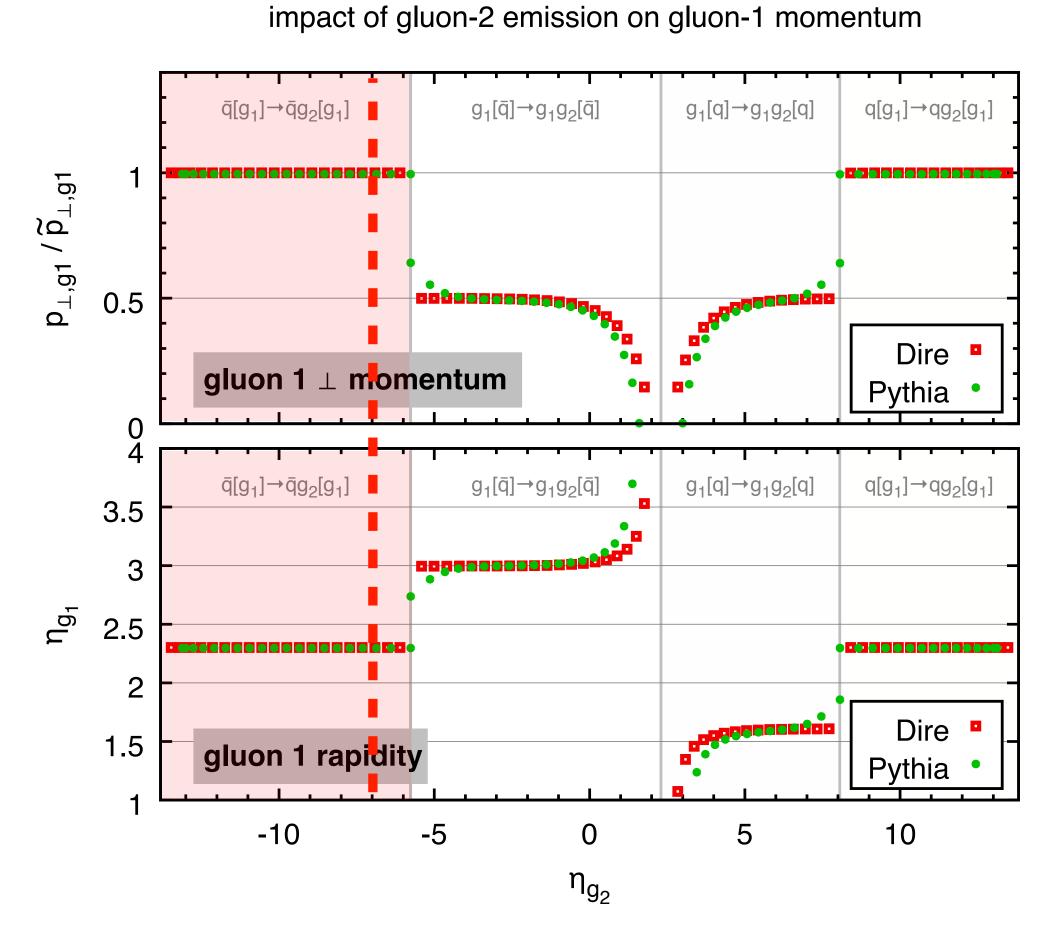


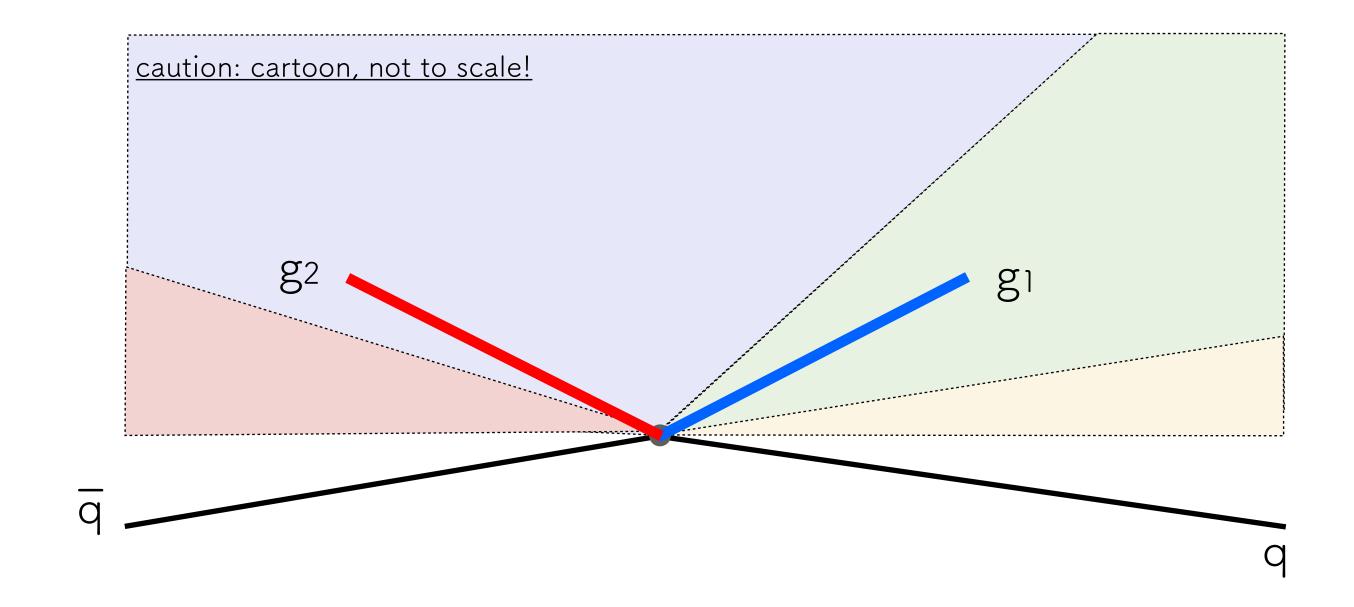




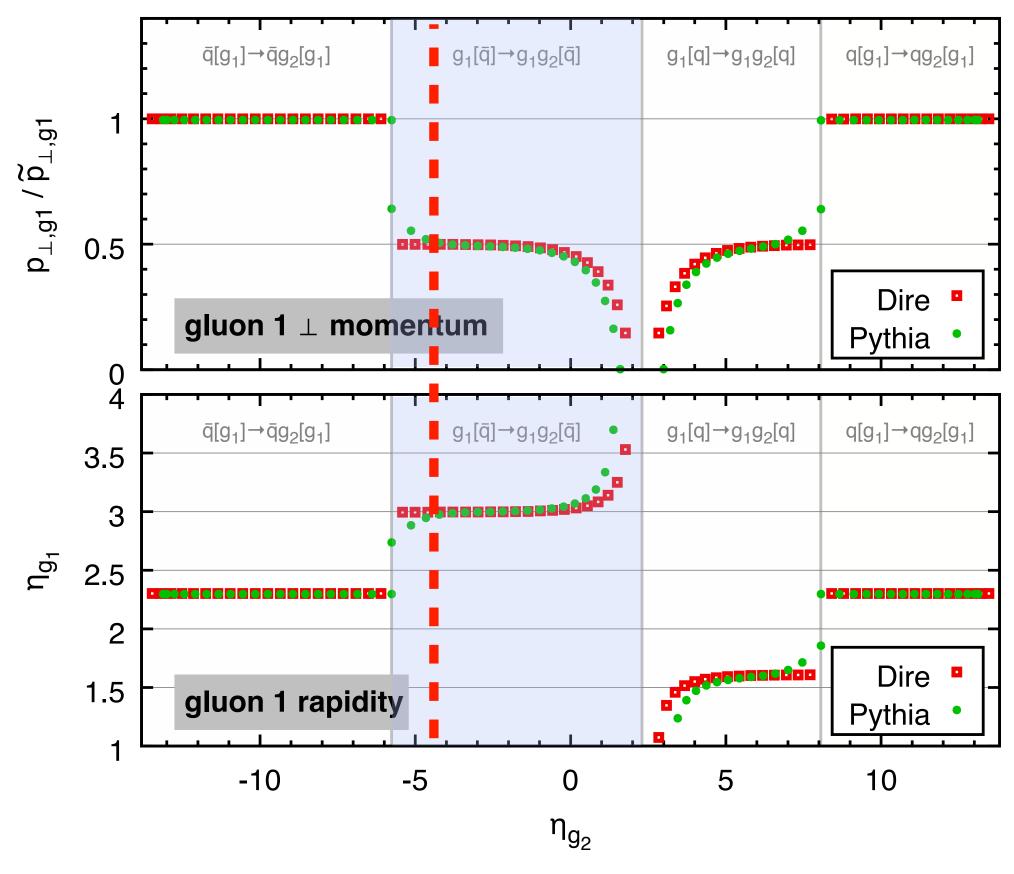


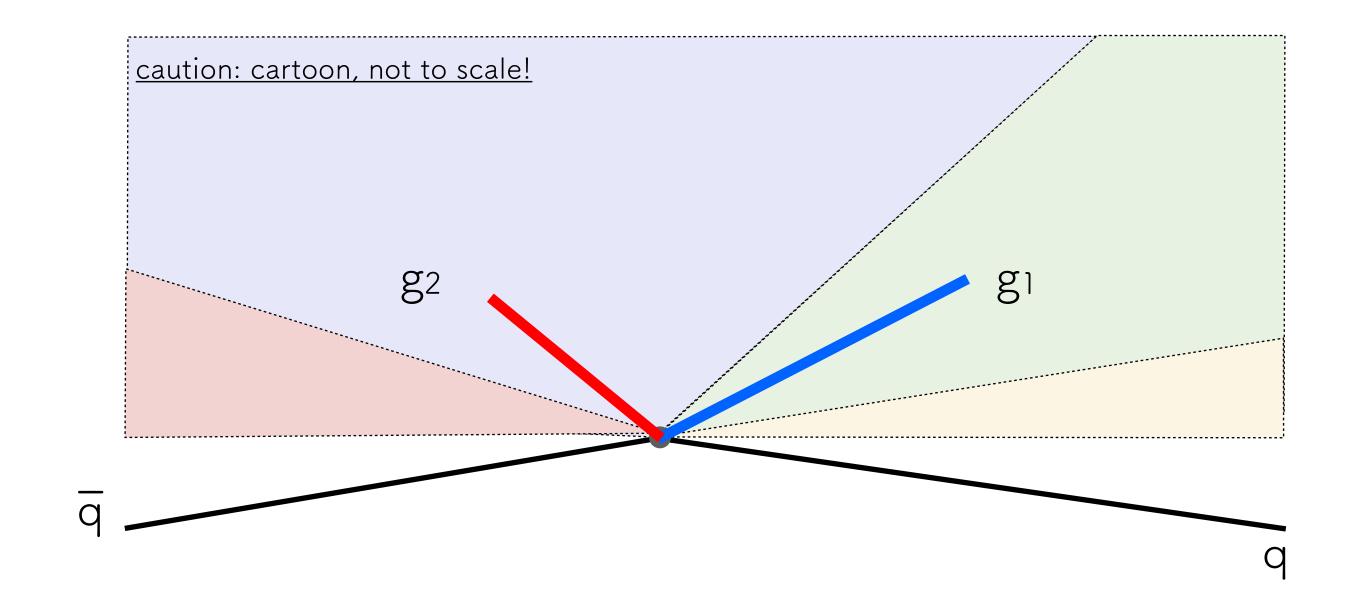




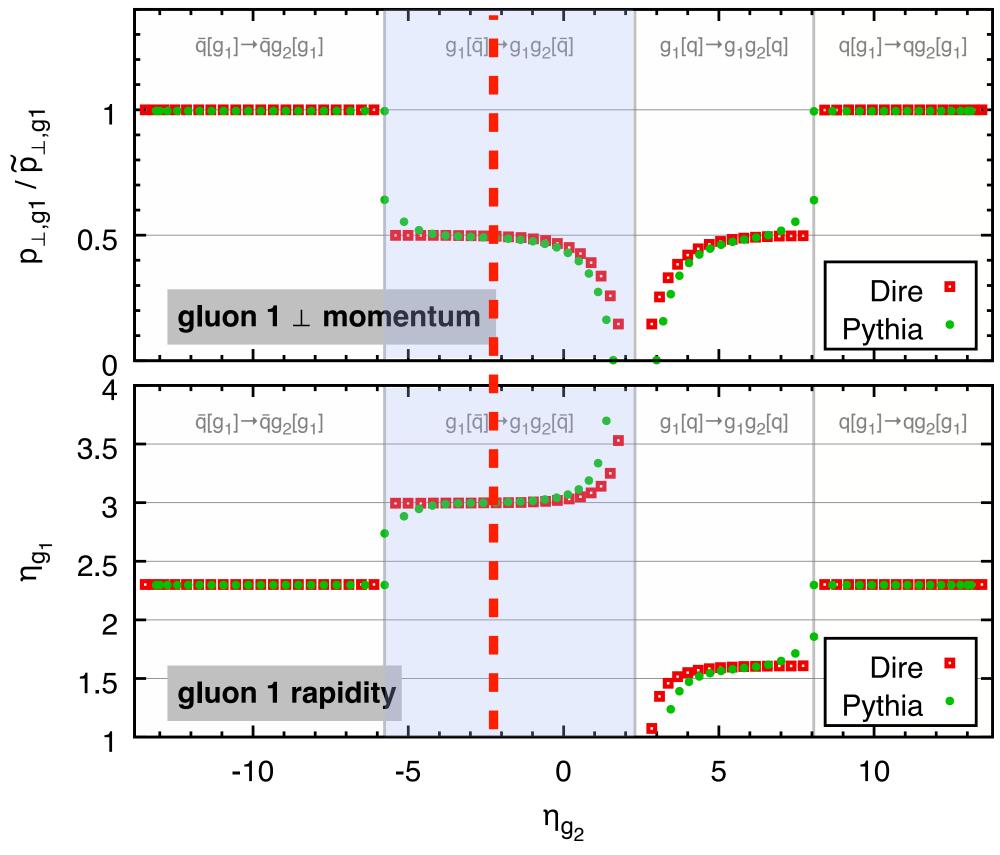


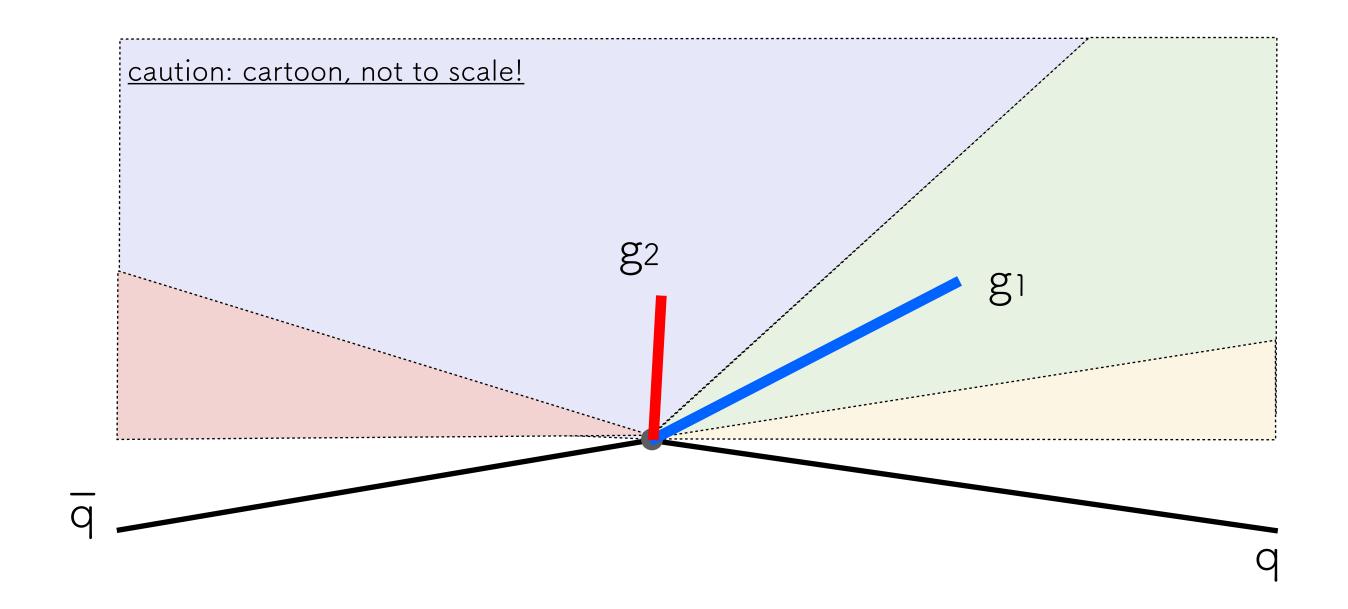
impact of gluon-2 emission on gluon-1 momentum



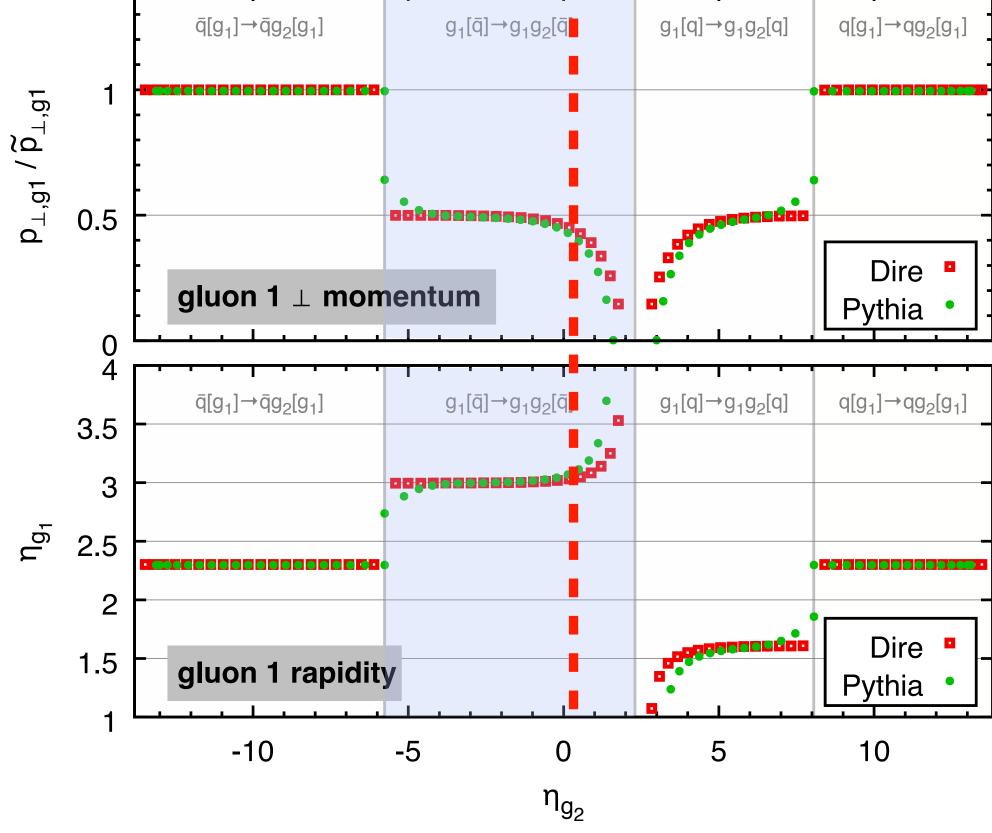


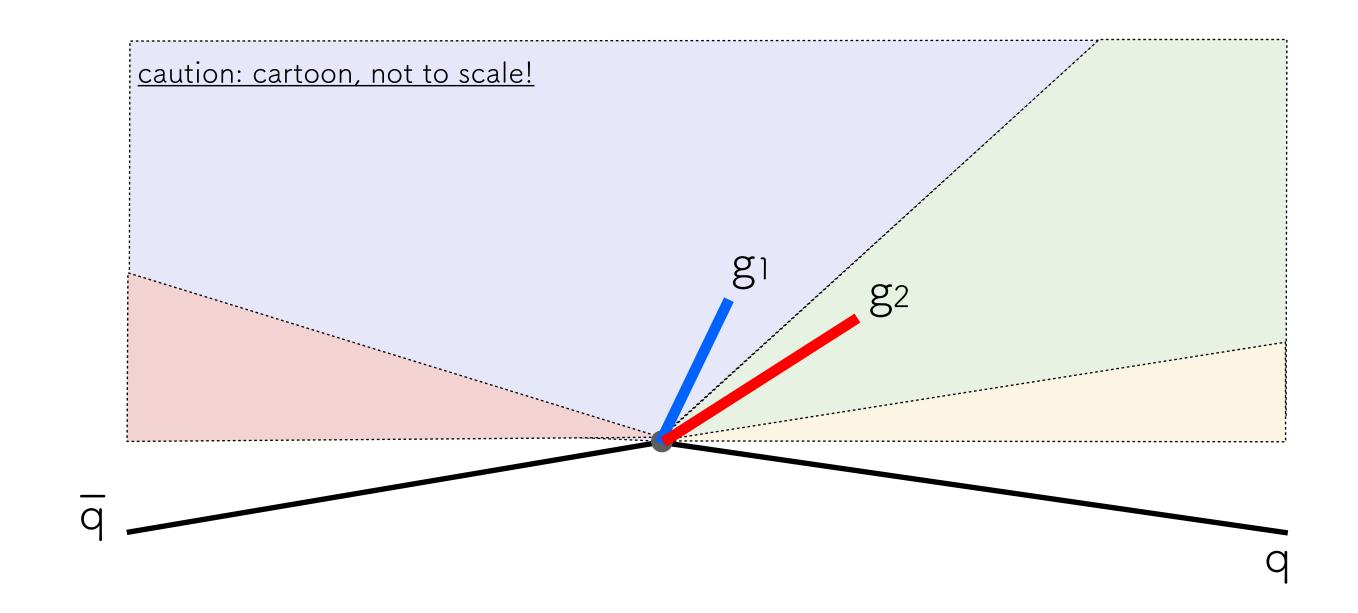
impact of gluon-2 emission on gluon-1 momentum

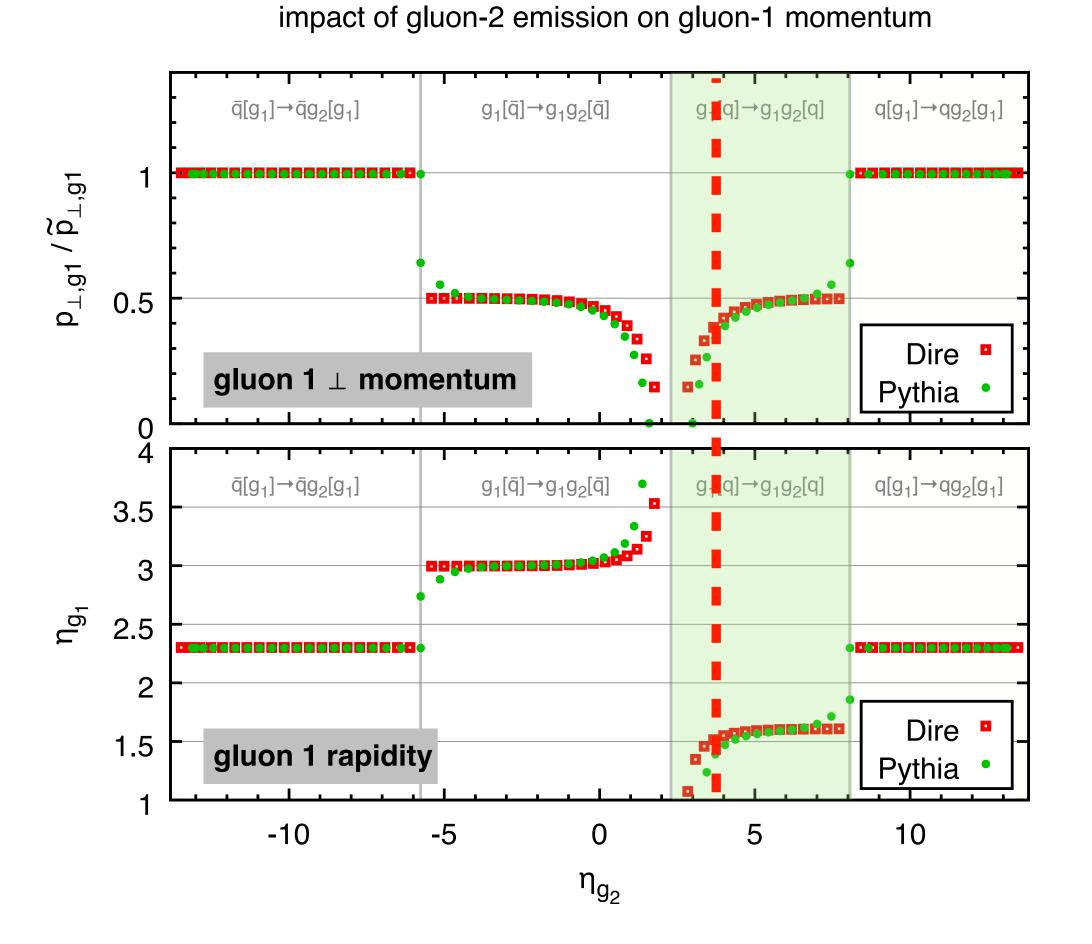


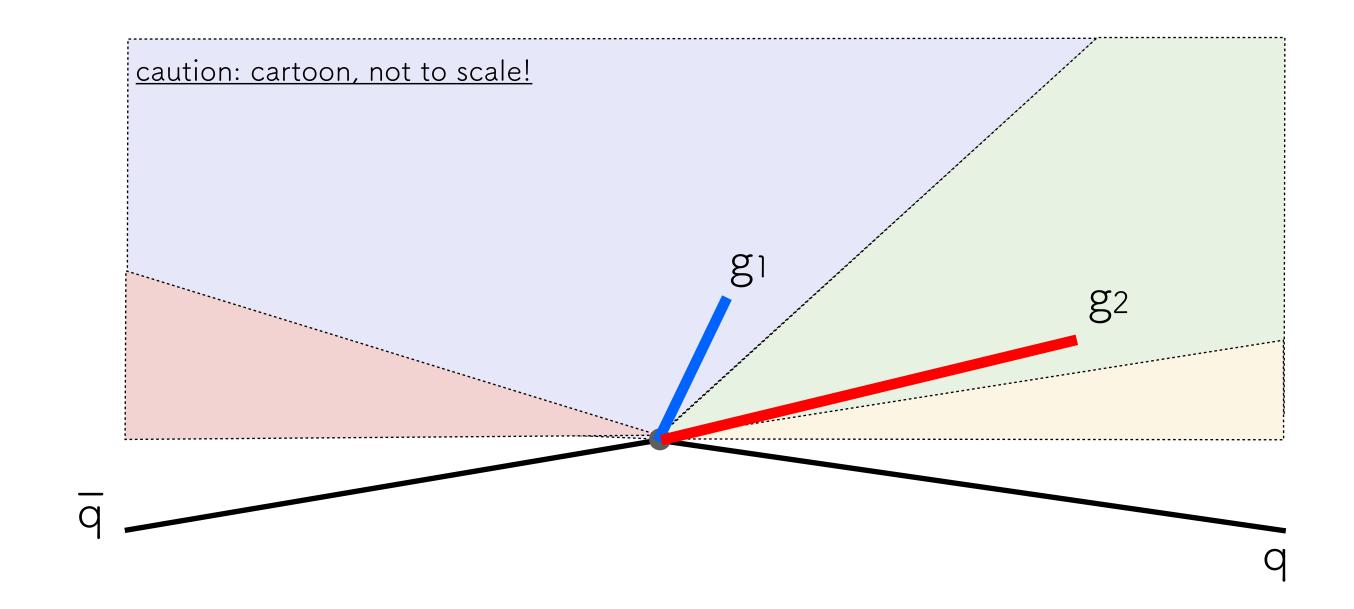


impact of gluon-2 emission on gluon-1 momentum $\bar{q}[g_1] \rightarrow \bar{q}g_2[g_1]$ $g_1[\bar{q}] \rightarrow g_1g_2[\bar{q}]$ $g_1[q] \rightarrow g_1g_2[q]$ $q[g_1] \rightarrow qg_2[g_1]$

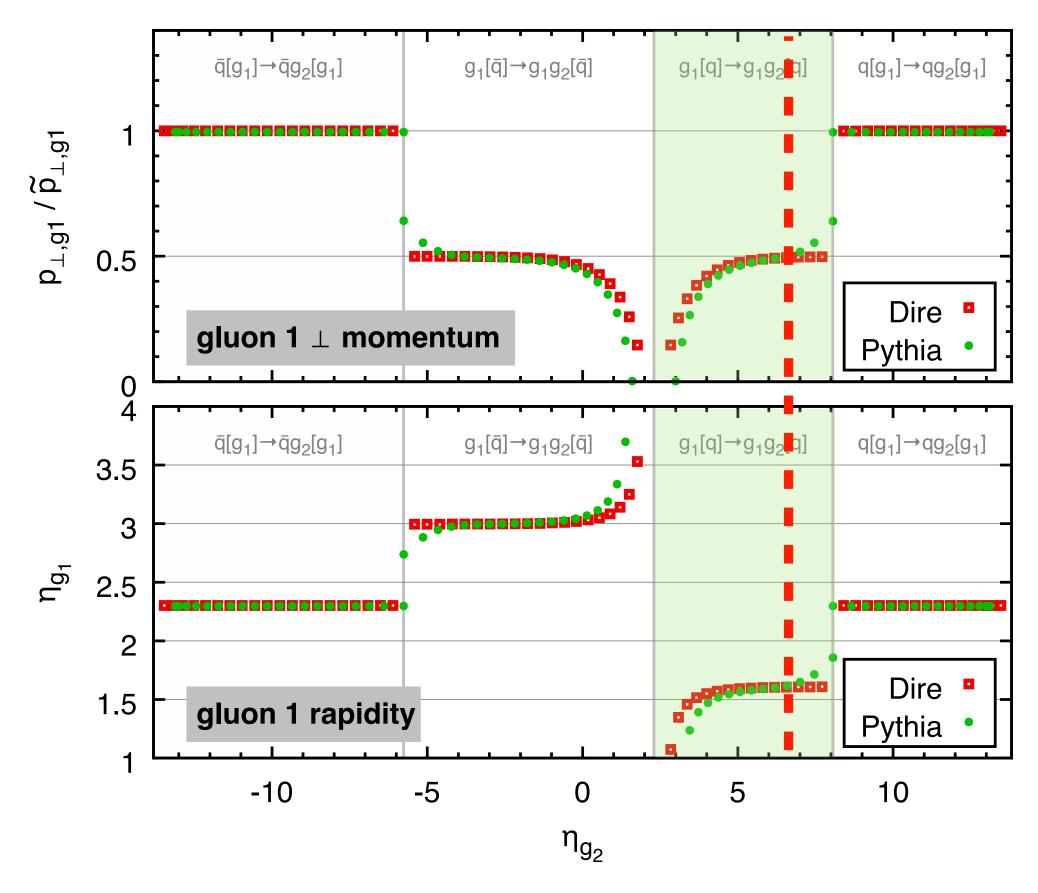


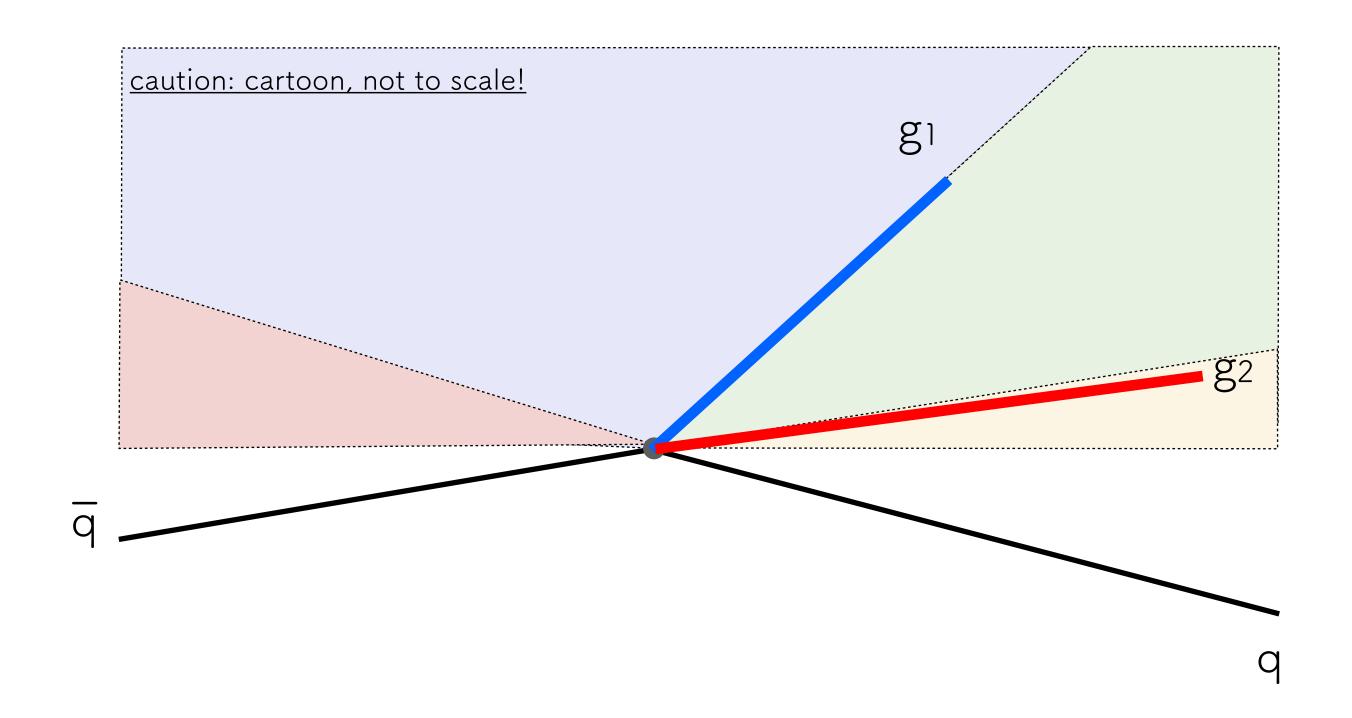


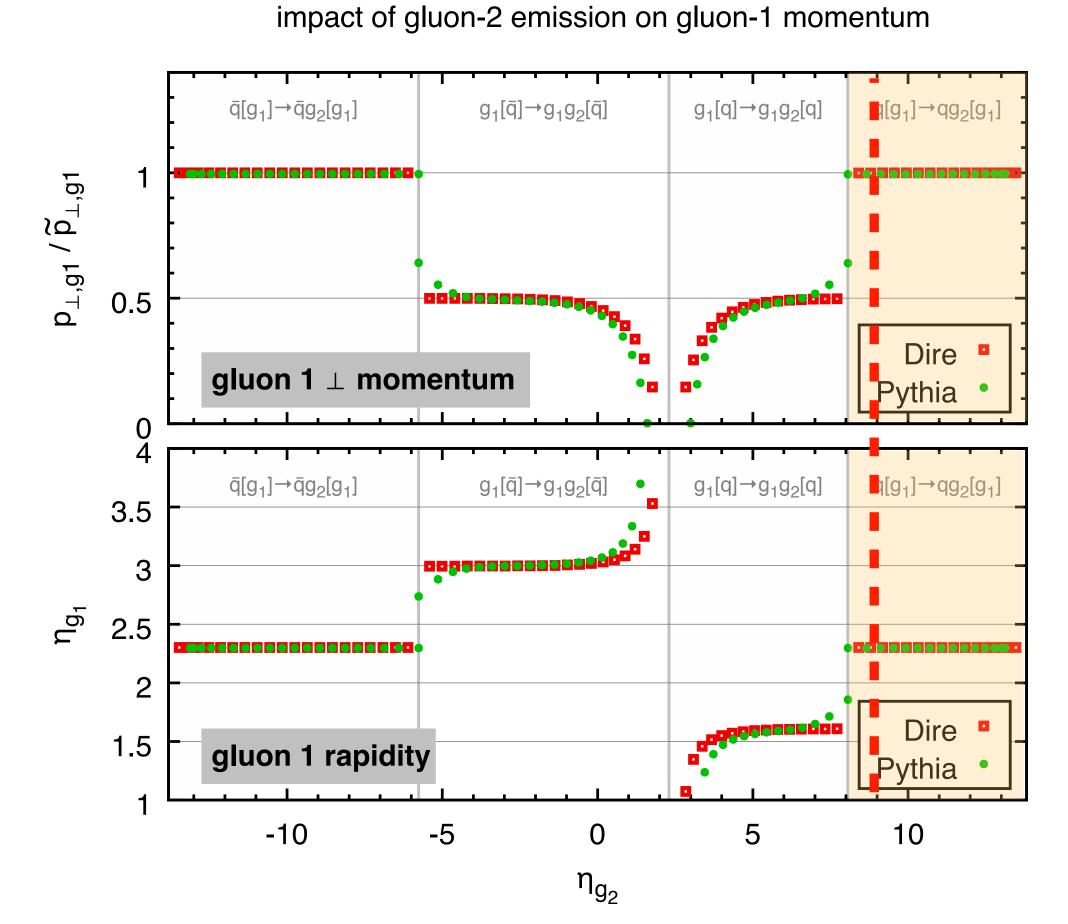


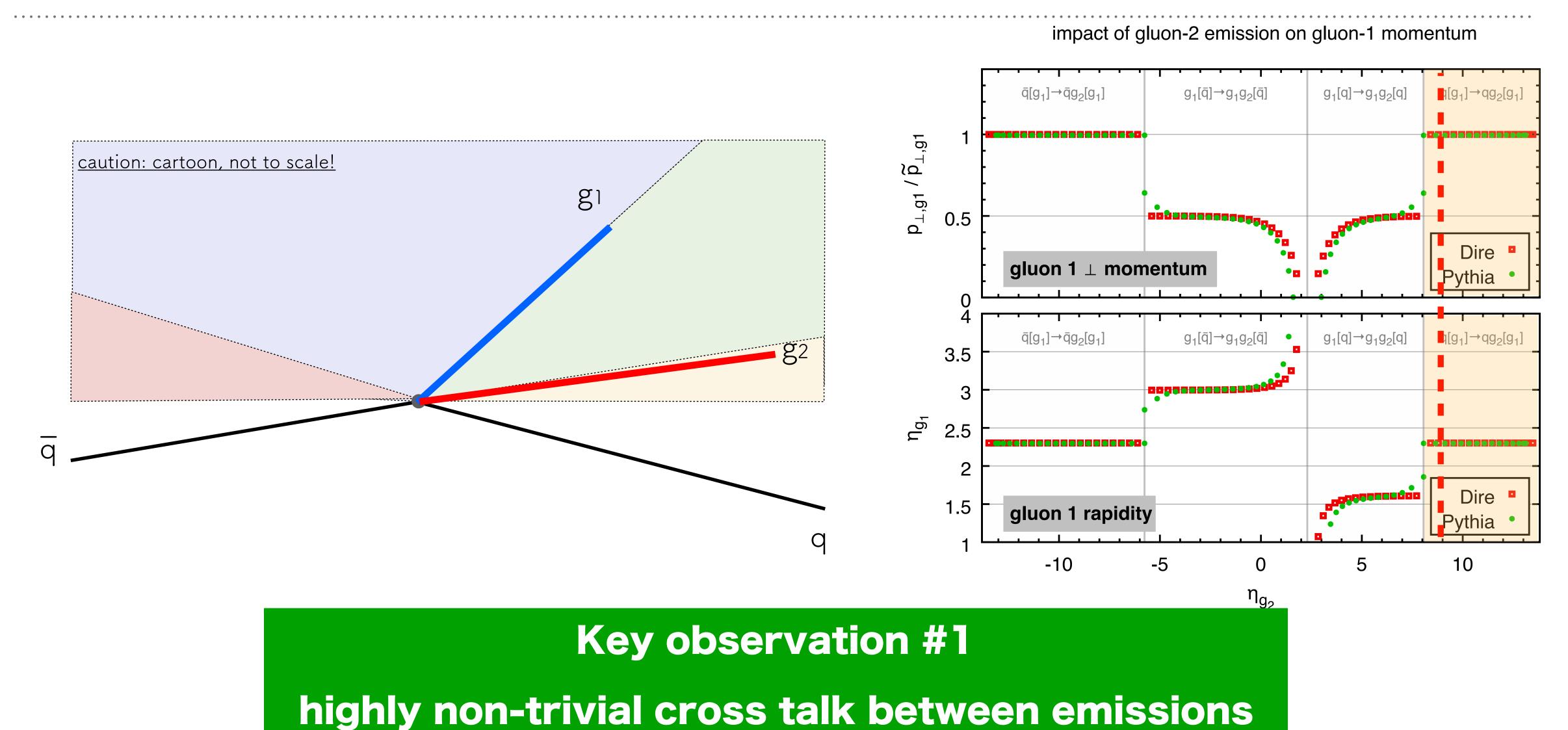


impact of gluon-2 emission on gluon-1 momentum



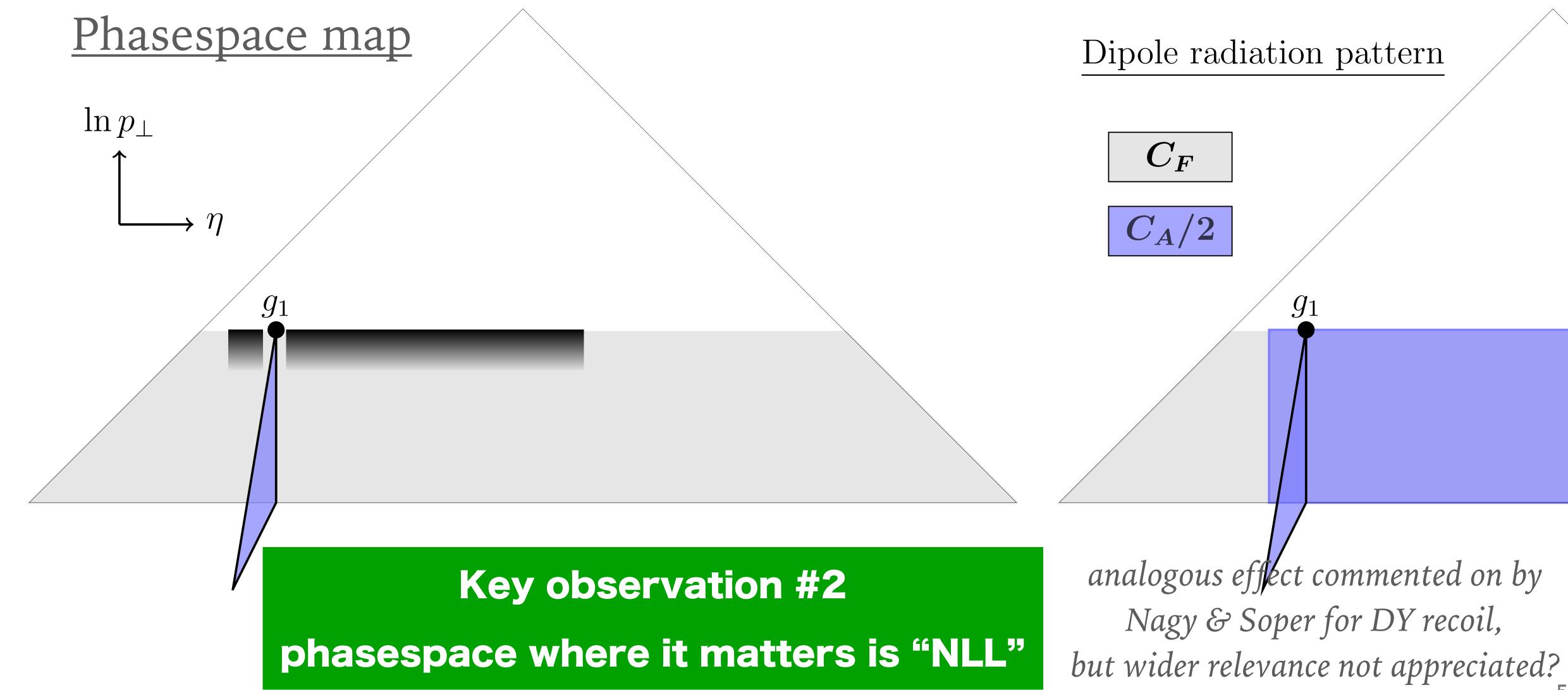




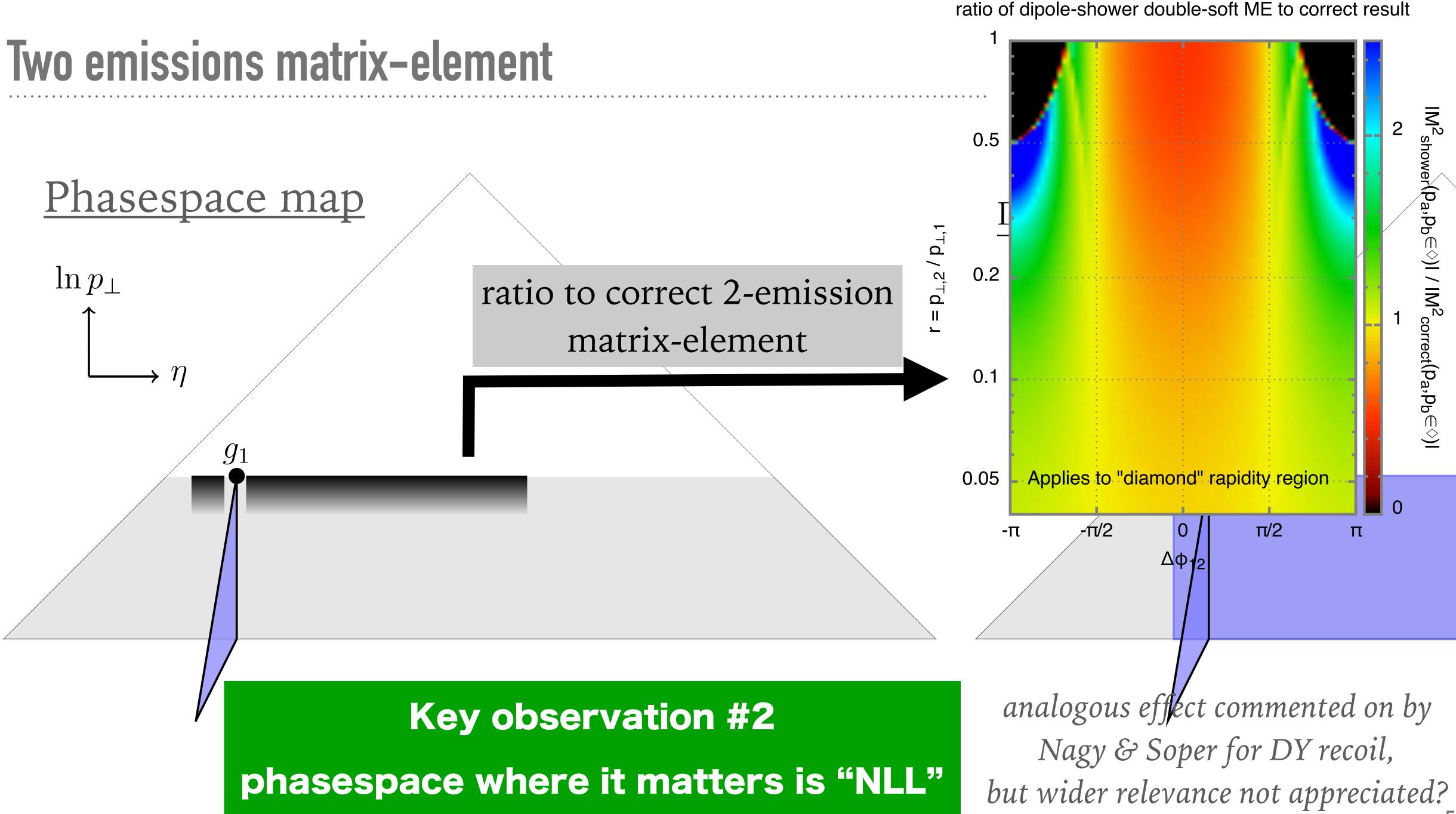


also noticed in 1992 by Andersson, Gustafson & Sjogren → special "fudge" in Ariadne

Two emissions matrix-element



50



Prevents shower from getting NLL accuracy for any e+e- event shape

	$NLL_{ln}\sum discrepancy$
1 - T	$0.116^{+0.004}_{-0.004}\bar{\alpha}^3 L^3$
vector p_t sum	$-0.349^{+0.003}_{-0.003}\bar{\alpha}^3 L^3$
B_T	$-0.0167335\bar{\alpha}^2 L^2$
$y_3^{ m cam}$	$-0.18277\bar{\alpha}^2L^2$
FC_1	$-0.066934ar{lpha}^2 L^2$

numerically, coefficients are not large compared to other effects, cf.

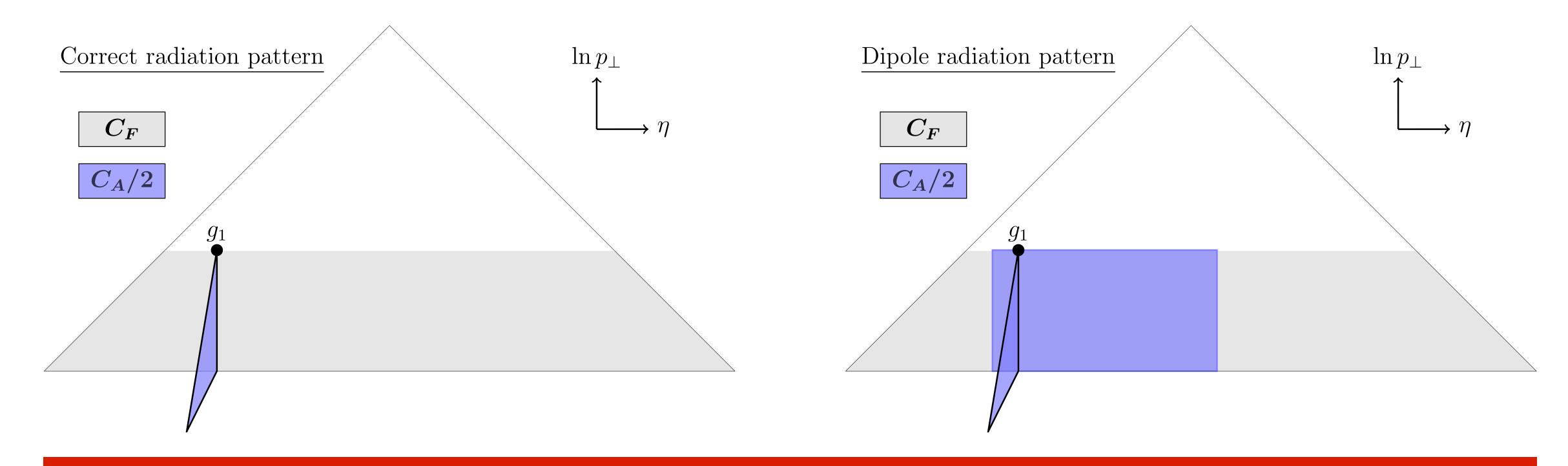
 $CMW \simeq 0.65\bar{\alpha}^2 L^2$

(because all these observables are quite inclusive)

but machine-learning uses all info — including large phasespace regions with 100% deficiencies

so far took $C_F = C_A/2$, i.e. leading N_C limit

In real life they're not equal & common choice for allocating them assigns $C_A/2$ to large part of phasespace that is actually gluon emission from quark (i.e. C_F)

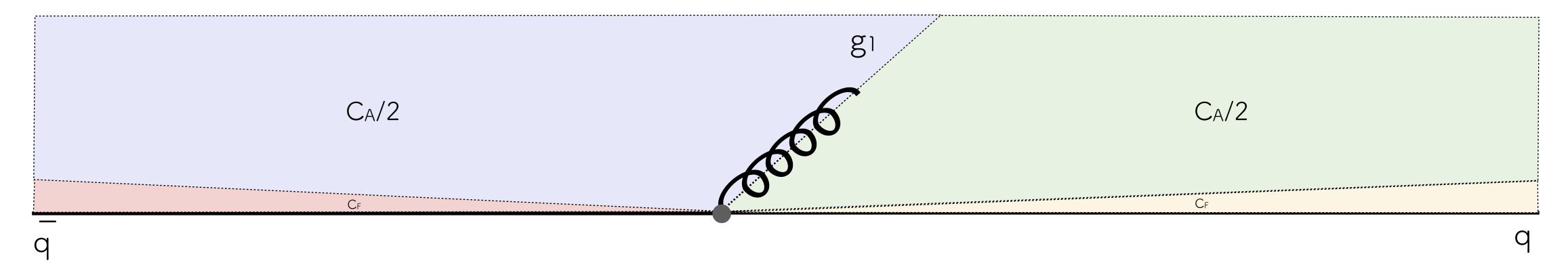


Key observation #3

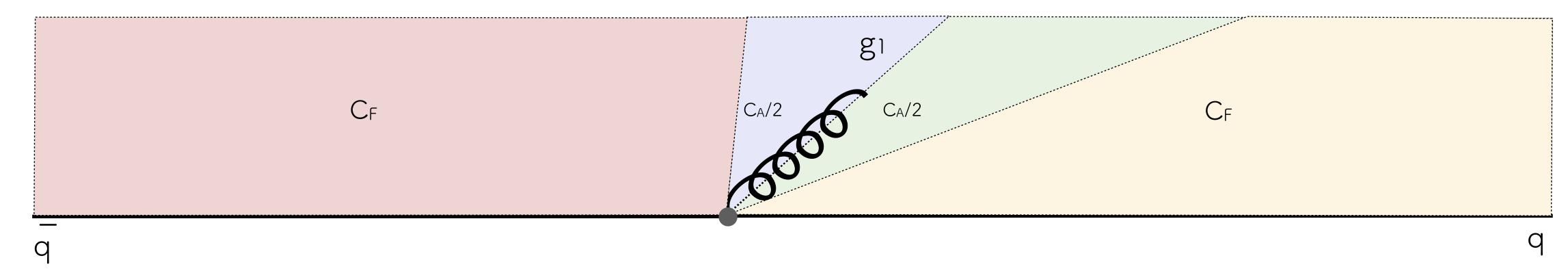
C_F v. C_A/2 issues occurs over a large area → double (leading) log effects?

another view of the colour issue

The dipole shower phase space partitioning of g₂'s radiation pattern is:



• Angular ordering implies a partitioning more like the following:



impact on observables?

Has LL subleading- N_C effect on 3-jet rates, thrust, but *not for* things like broadening, 2-jet rate (which are physically close the evolution variable, i.e. transverse momentum).

E.g. for thrust

$$\delta\Sigma(L) = -\frac{1}{64}\bar{\alpha}^2 L^4 \left(\frac{C_A}{2C_F} - 1\right)$$

Closing

Conclusions

Parton showers are a crucial element in collider physics

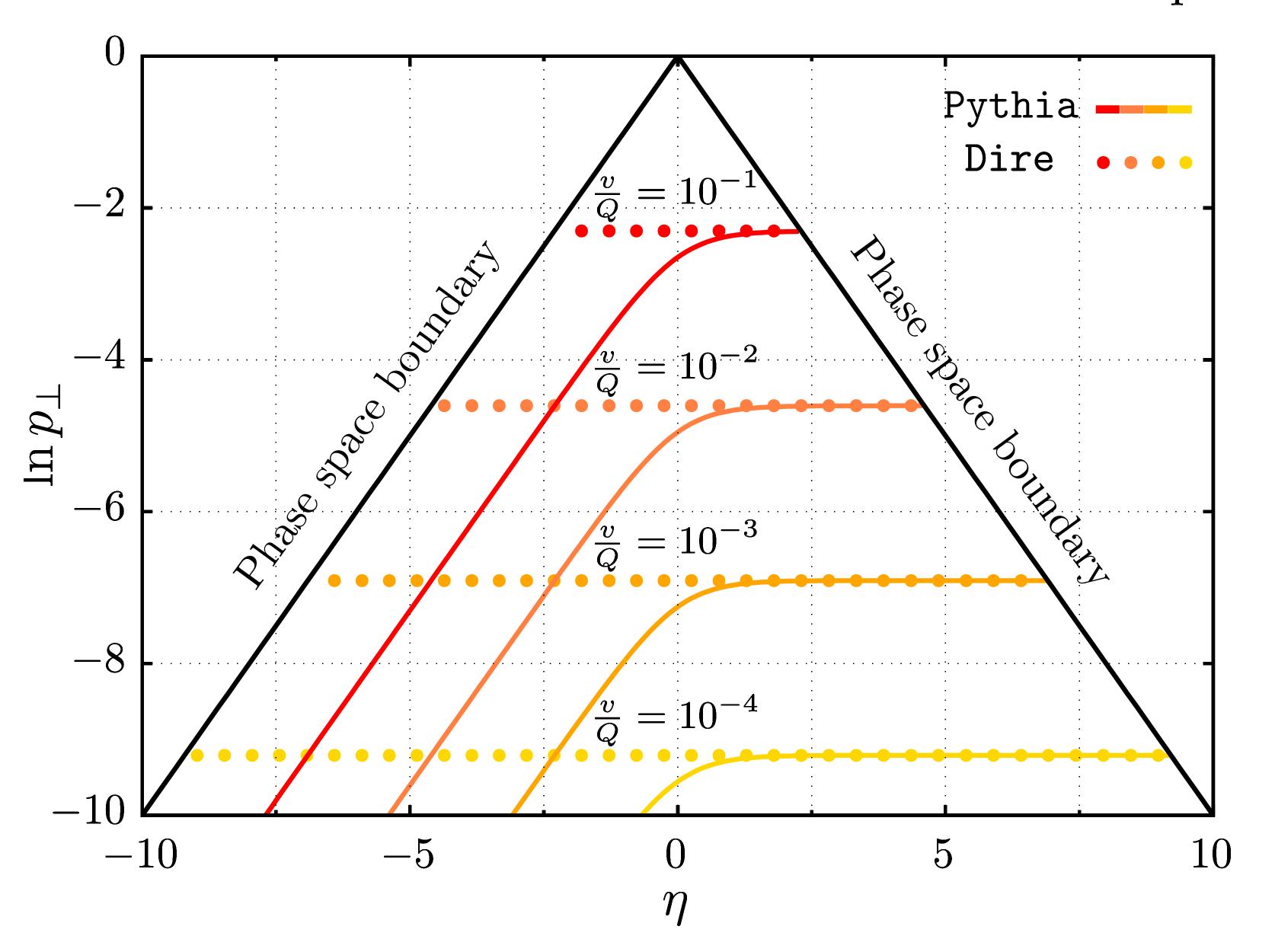
Seeing many developments (subleading colour for non-global logarithms, multi-particle emission kernels, etc.)

But maybe we need to go back to foundations:

- improving parton showers is not just a question of better components (e.g. higher-order splitting kernels)
- > question of how components are assembled is equally crucial
- > we must identify & state what a parton shower should be achieving
- new studies along these lines are teaching us important things about existing showers

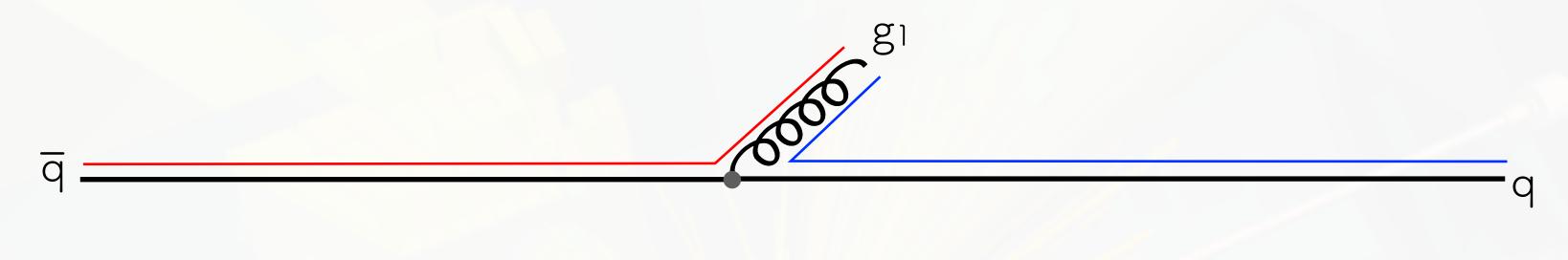
BACKUP

Constant evolution variable contours in the Lund plane

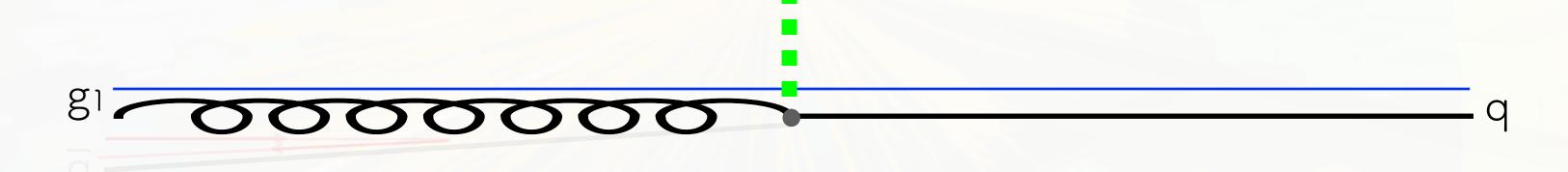


two soft emissions: boost dipole partitions back into the event COM

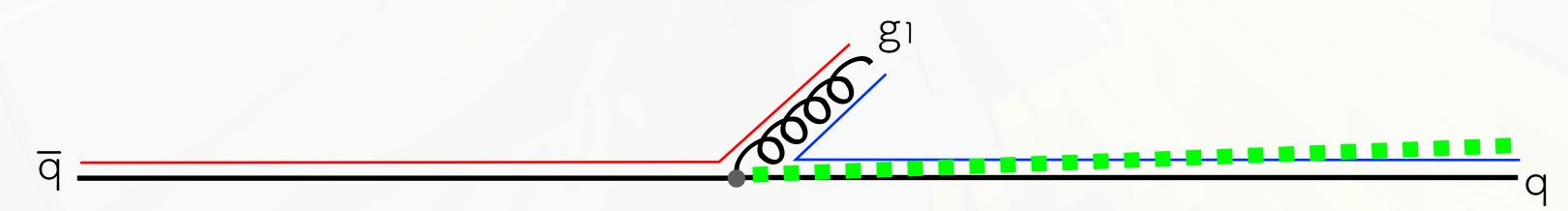
• Consider we emitted **soft gluon g1** from **hard** qq, $s\bar{o}$ we end up with a qg1 \bar{a} nd a g1q dipole:



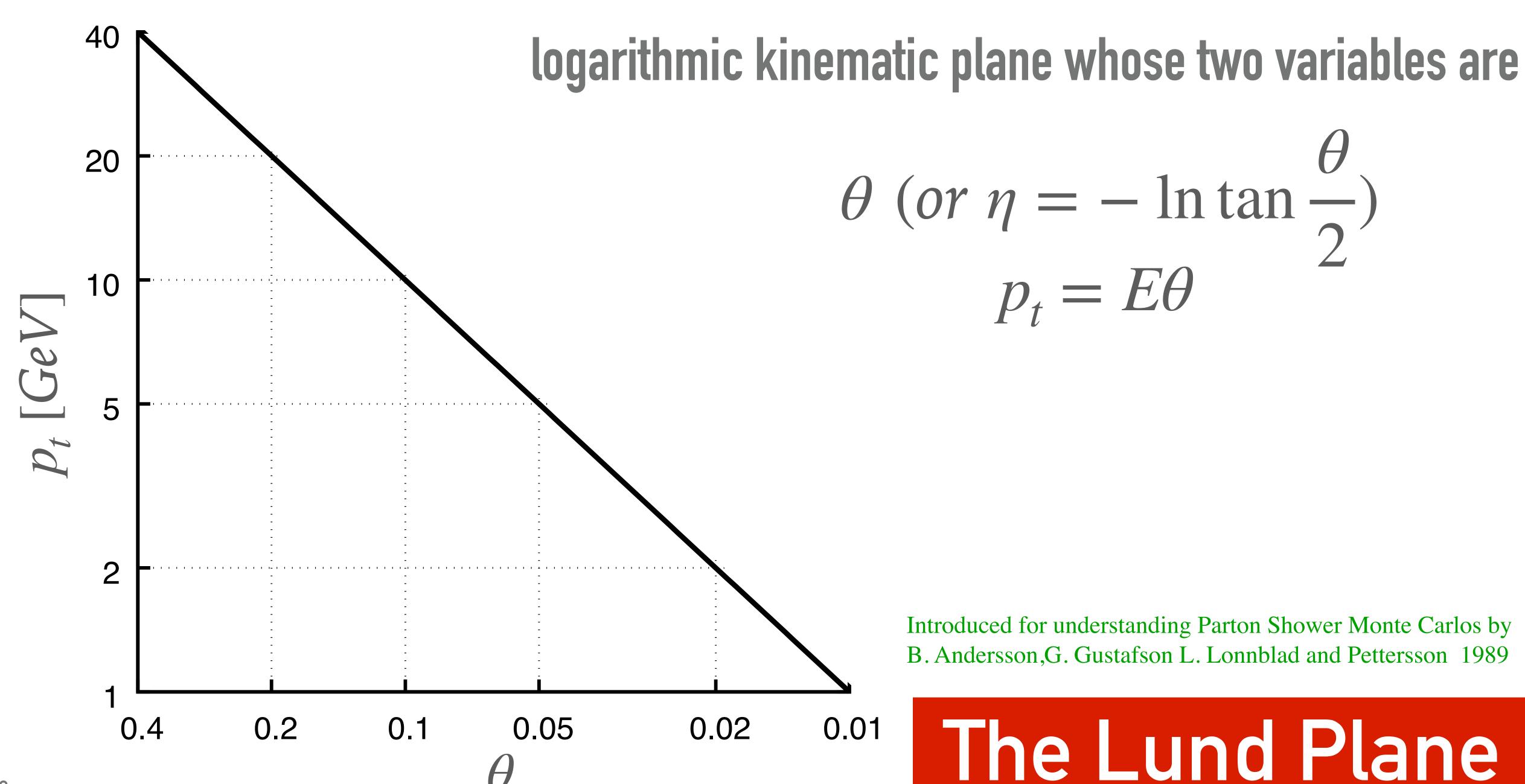
- $^{\circ}$ To get us from the event COM to the g1q dipole COM (blue line) requires a **BIG BOOST** \rightarrow
- Dipole partitioned at $\eta = 0$ in that frame:



° To get us back to the event COM from the g1q dipole COM undo the same **BIG BOOST** ←

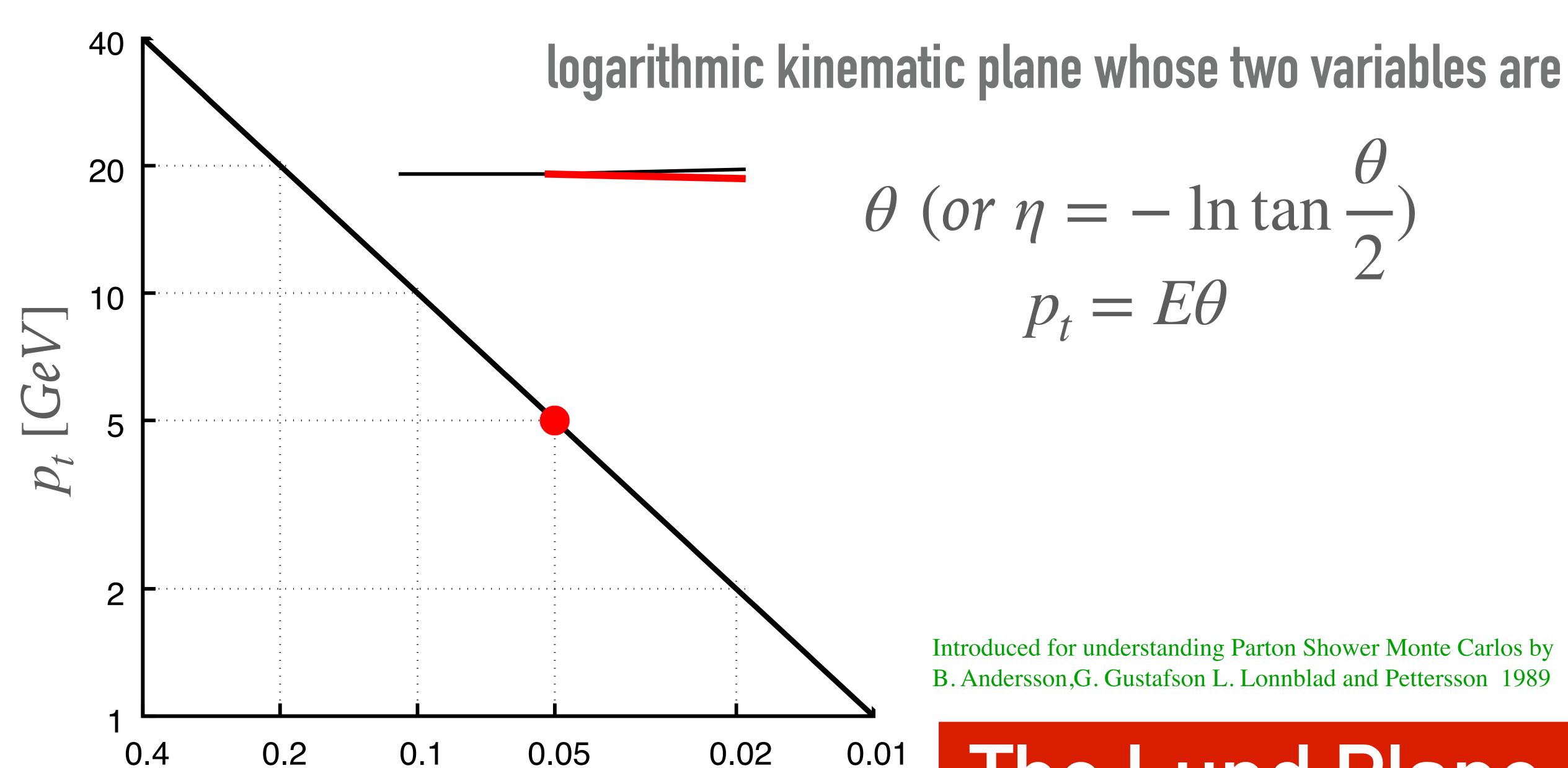


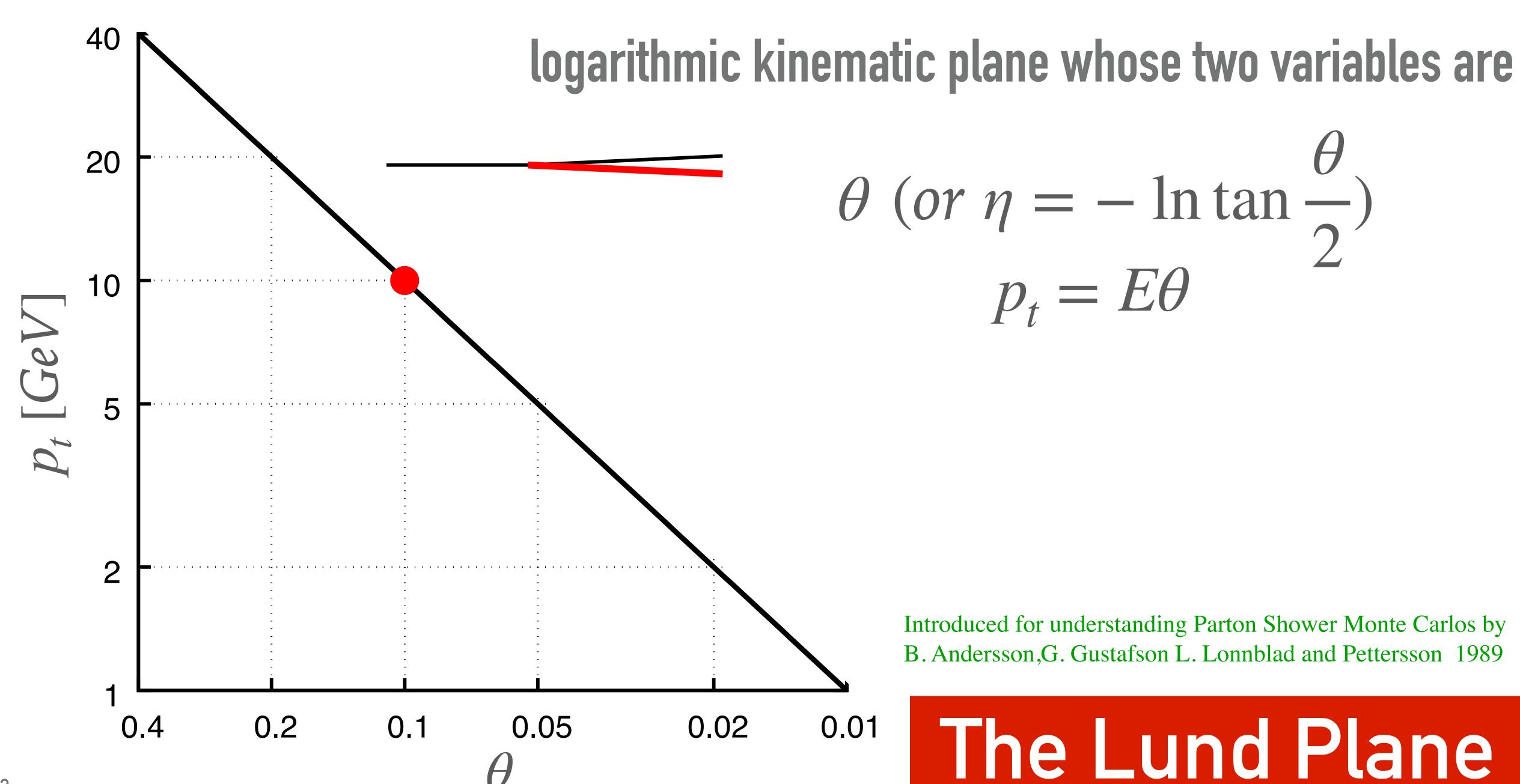
• In event COM partition comes out very close to q; instead of equidistant in angle between g1 & q



$\theta \text{ (or } \eta = -\ln \tan \frac{\theta}{2})$ $p_t = E\theta$

Introduced for understanding Parton Shower Monte Carlos by B. Andersson, G. Gustafson L. Lonnblad and Pettersson 1989

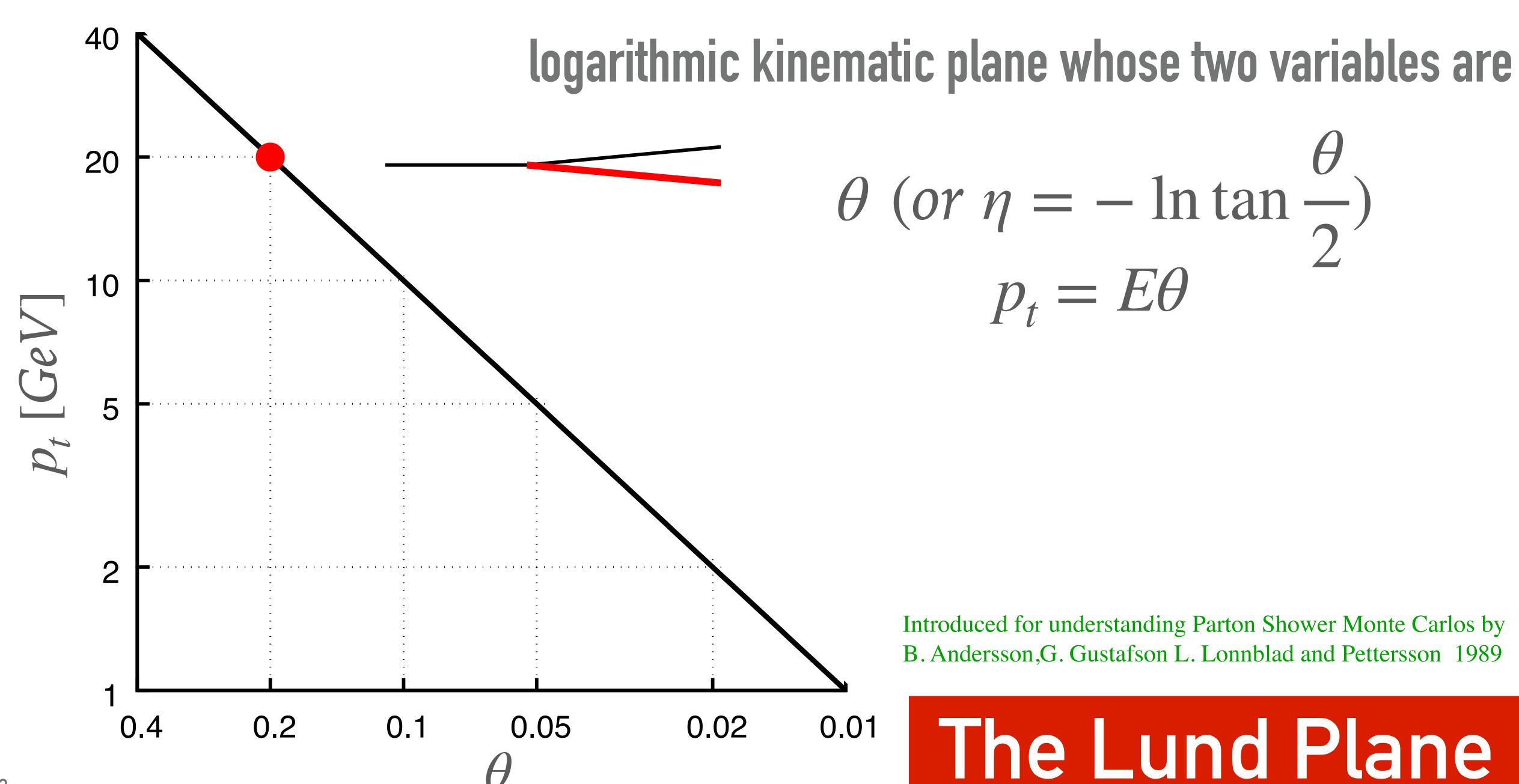




$$\theta \text{ (or } \eta = -\ln \tan \frac{\theta}{2})$$

$$p_t = E\theta$$

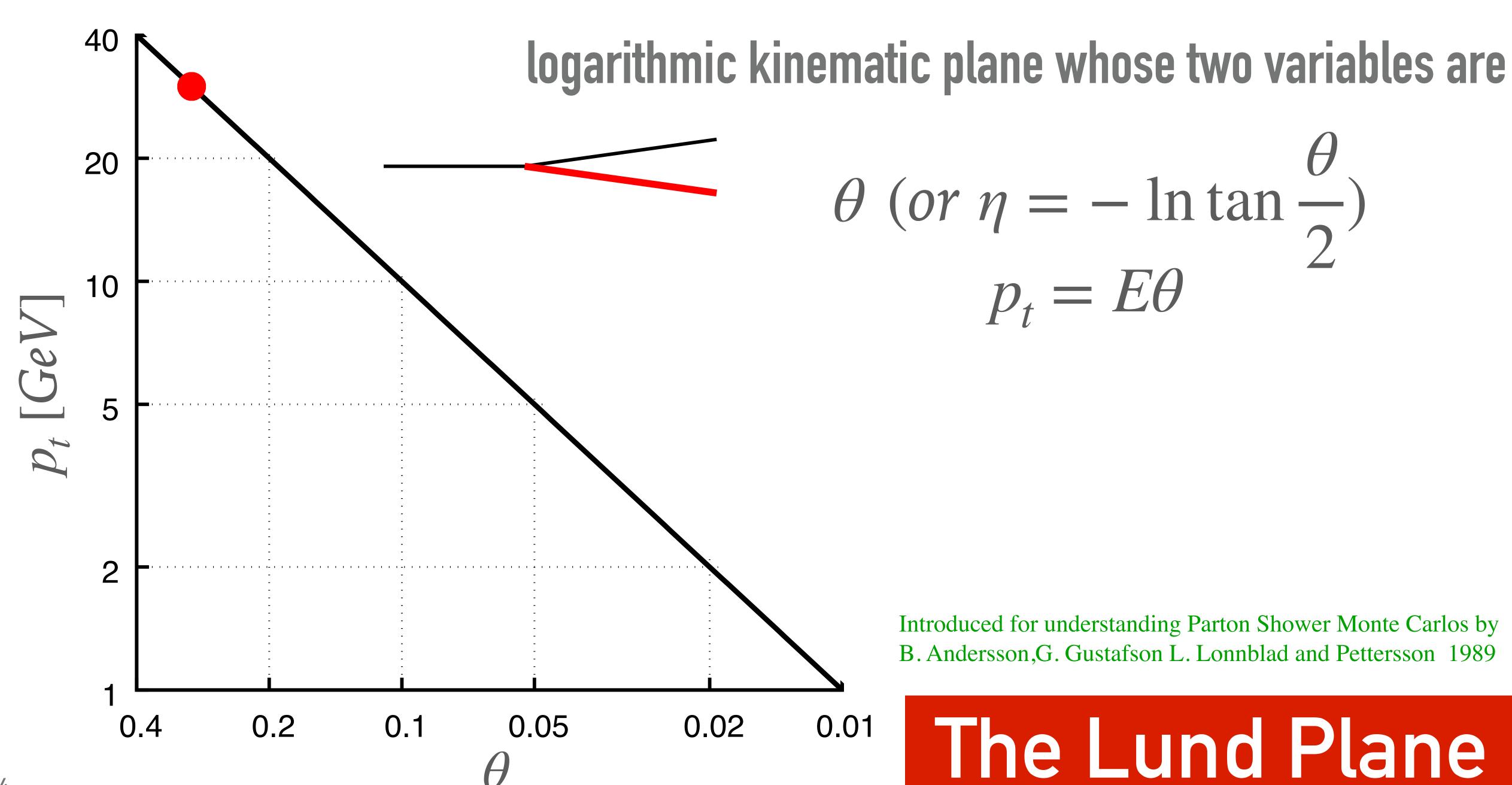
Introduced for understanding Parton Shower Monte Carlos by B. Andersson, G. Gustafson L. Lonnblad and Pettersson 1989



$$\theta \text{ (or } \eta = -\ln \tan \frac{\theta}{2})$$

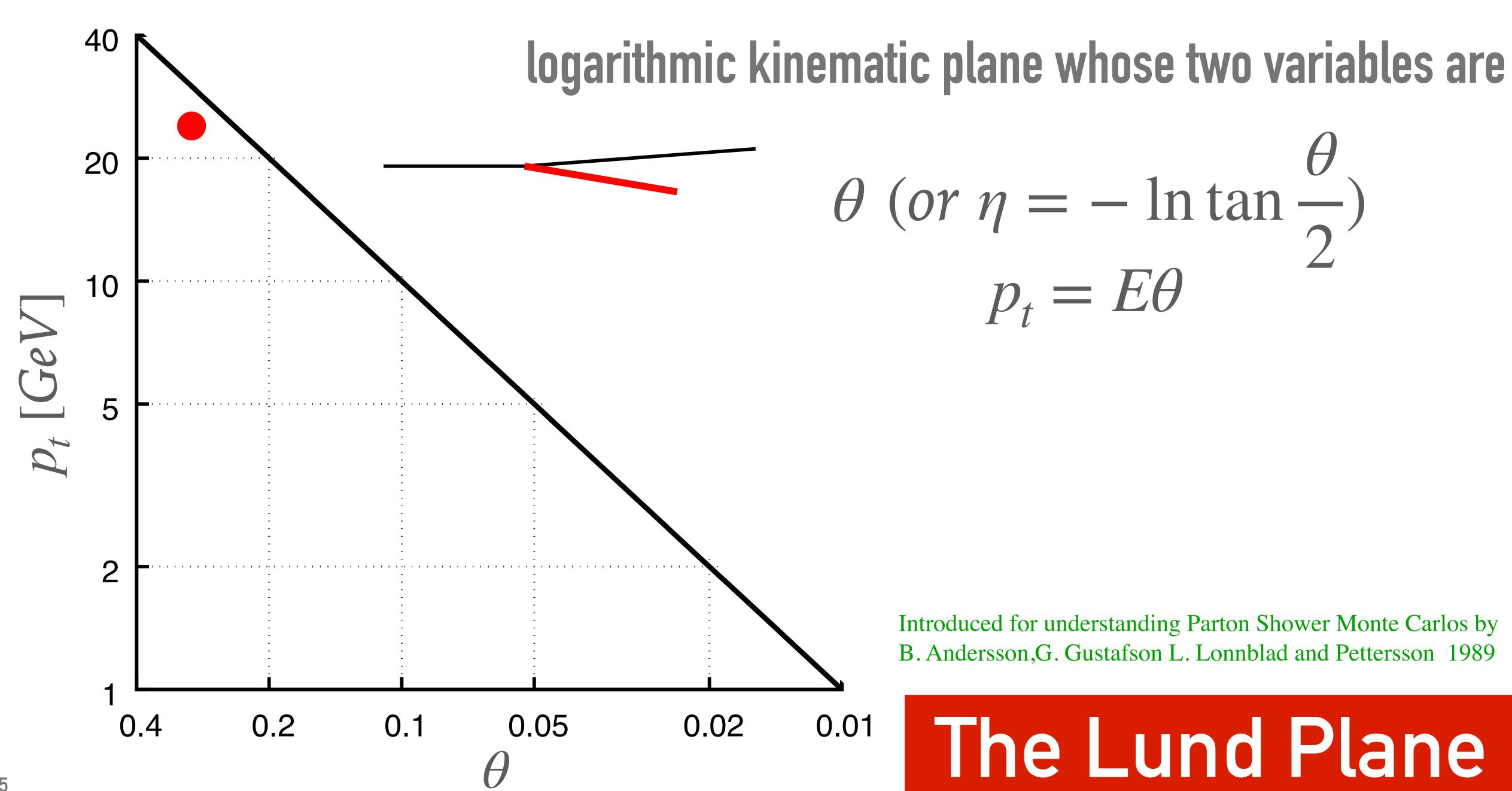
$$p_t = E\theta$$

Introduced for understanding Parton Shower Monte Carlos by B. Andersson, G. Gustafson L. Lonnblad and Pettersson 1989



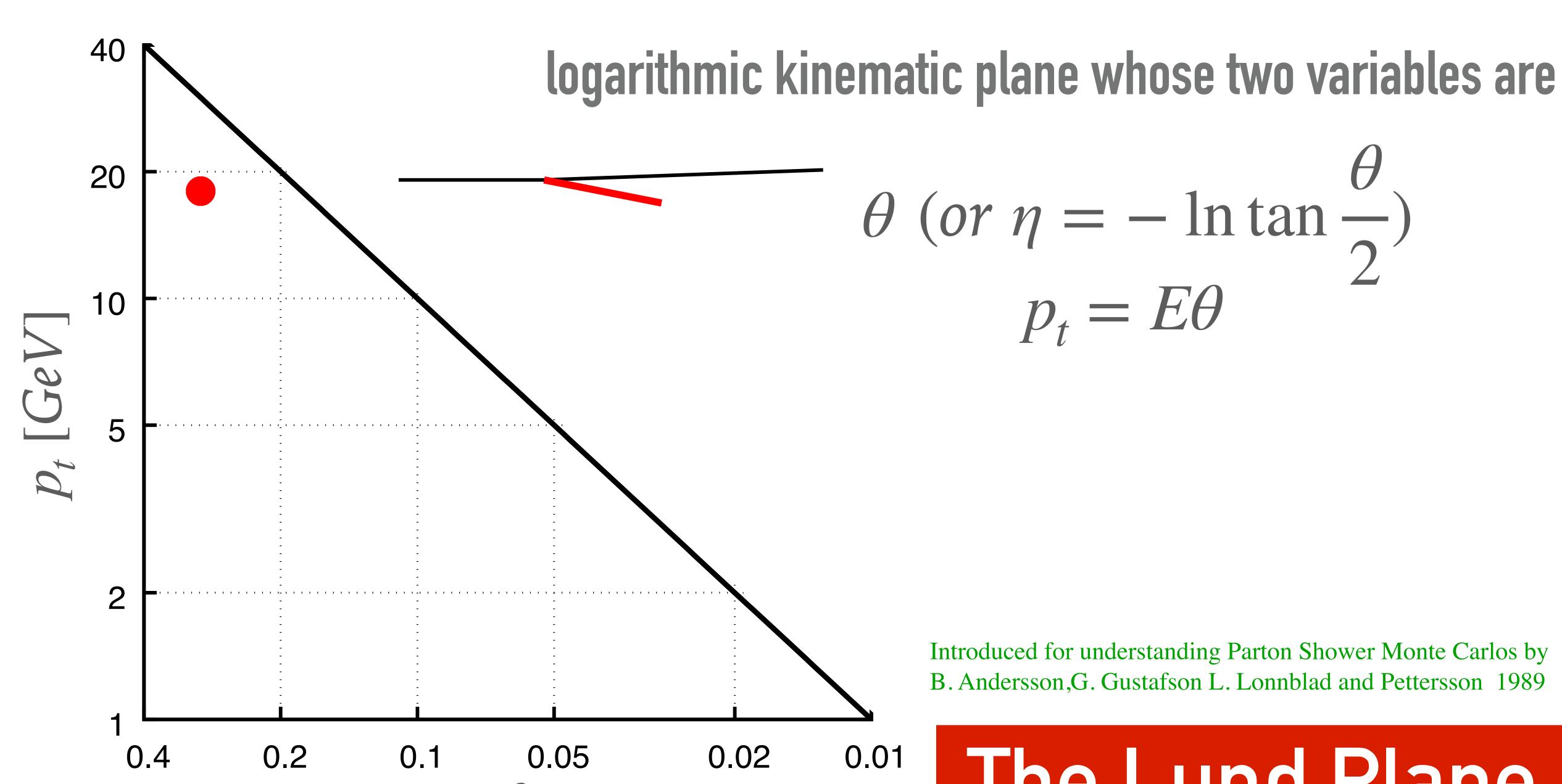
$\theta \text{ (or } \eta = -\ln \tan \frac{\theta}{2})$ $p_t = E\theta$

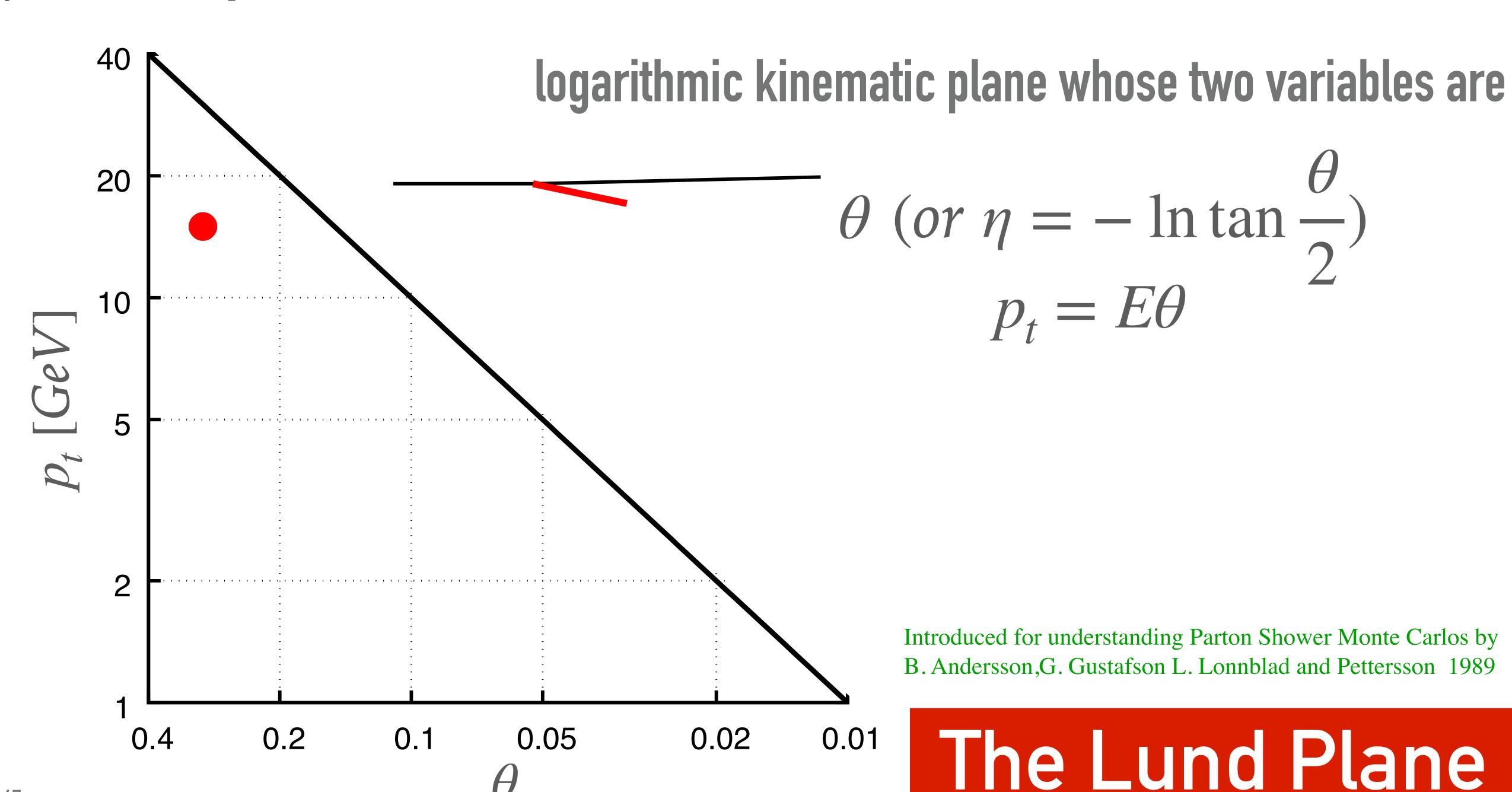
Introduced for understanding Parton Shower Monte Carlos by B. Andersson, G. Gustafson L. Lonnblad and Pettersson 1989

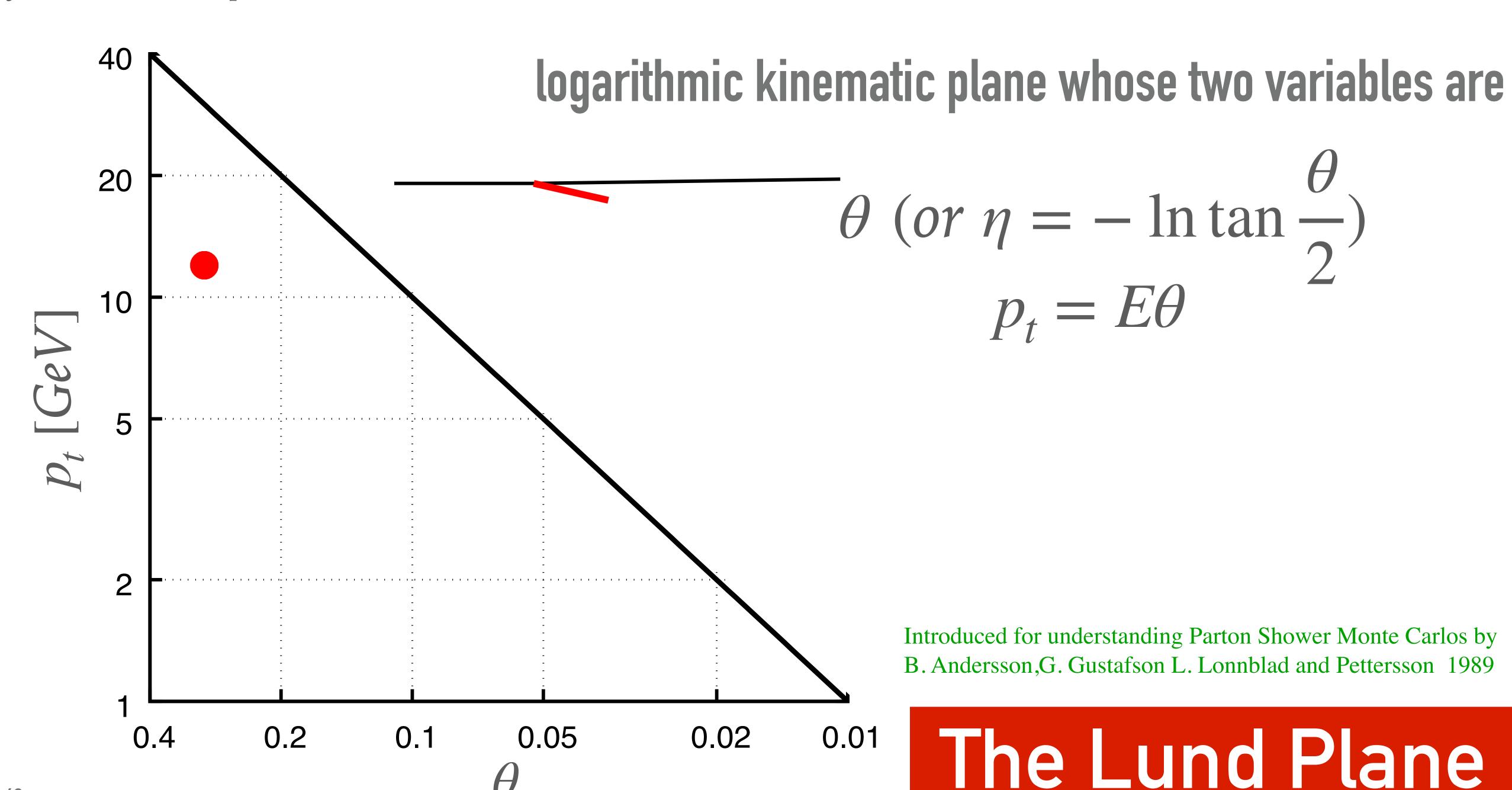


$\theta \text{ (or } \eta = -\ln \tan \frac{\theta}{2})$ $p_t = E\theta$

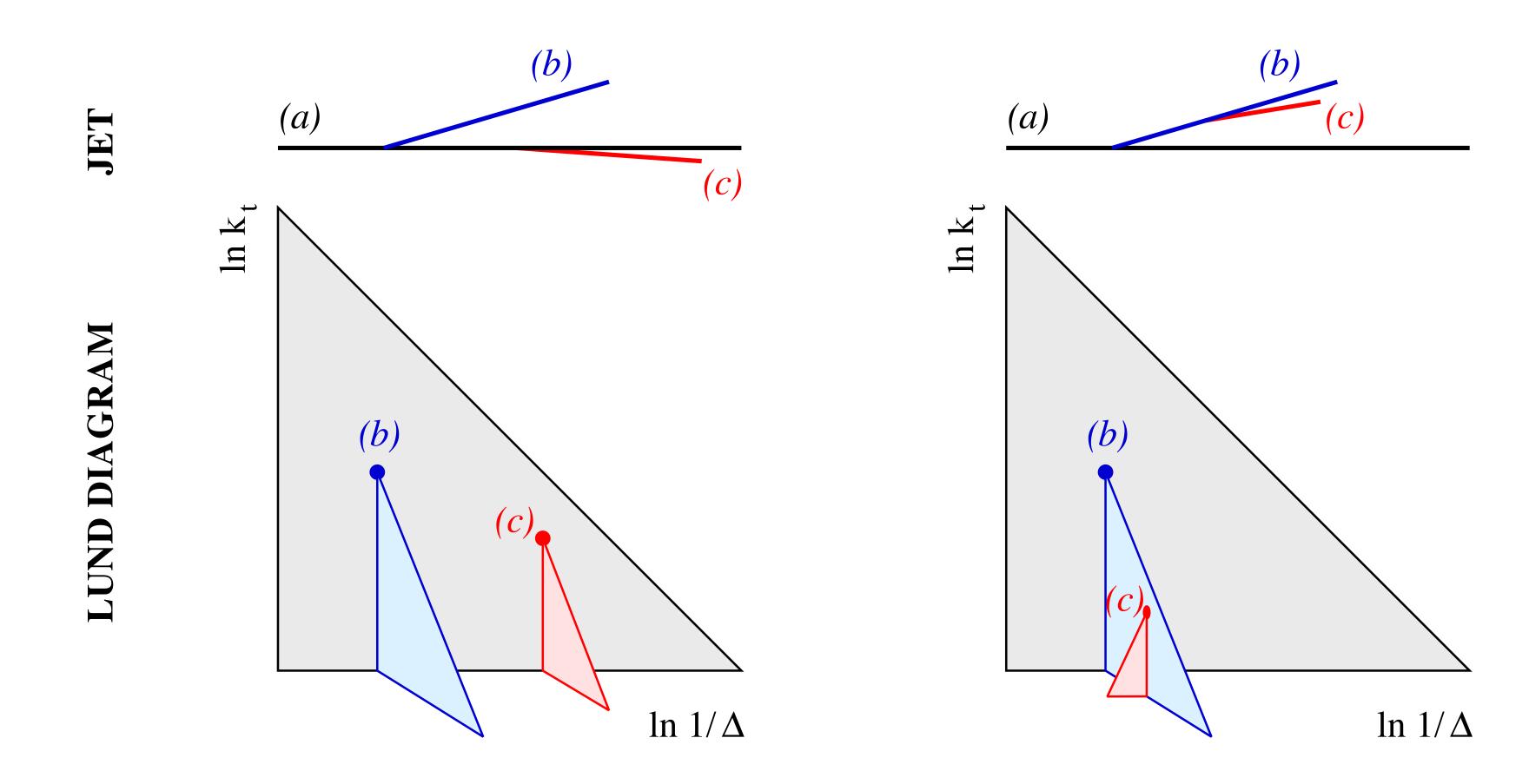
Introduced for understanding Parton Shower Monte Carlos by B. Andersson, G. Gustafson L. Lonnblad and Pettersson 1989



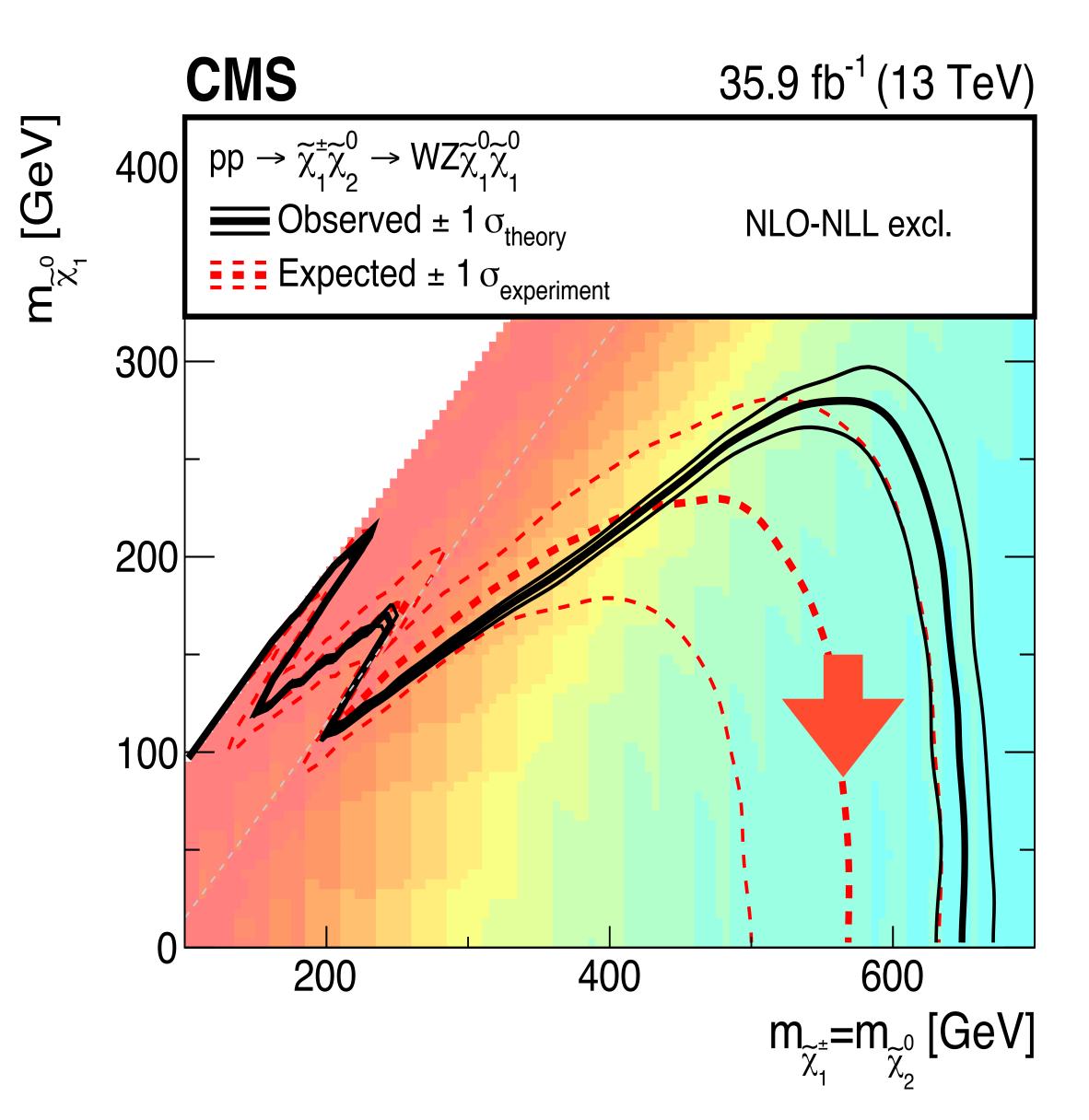




organise phasespace: Lund diagrams



The path forward: collect 20-30x more collisions by ~ 2035



- > Suppose we had a choice between
 - ➤ HL-LHC (14 TeV, 3ab-1)
 - ➤ or going to higher c.o.m. energy but limited to 80fb⁻¹.
- ➤ How much energy would we need to equal the HL-LHC?

today's reach (13 TeV, 80fb ⁻¹)	HL-LHC reach (14 TeV 3ab ⁻¹)	energy needed for same reach with 80fb ⁻¹
4.7 TeV SSM Z'	6.7 TeV	20 TeV
2 TeV weakly coupled Z'	3.7 TeV	37 TeV
680 GeV chargino	1.4 TeV	54 TeV

Hard processes: to 3rd order (NNLO) in perturbation theory strong coupling constant (a_s)

