## New results for parton showers

Gavin P. Salam*
Rudolf Peierls Centre for Theoretical Physics $\mathcal{E}$ All Souls College, Oxford

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including work with
M. Dasgupta, F. Dreyer, K. Hamilton, P. Monni
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"big unanswered questions"
about fundamental particles \& their interactions (dark matter, matter-antimatter asymmetry, nature of dark energy, hierarchy of scales...) and
"big answerable questions"
(structure of Higgs sector, determining fundamental parameters of Lagrangian of particle physics)

## Higgs precision ( $H \rightarrow \gamma \gamma$ ) : optimistic estimate v. luminosity \& time

extrapolation of $\mu_{\mathrm{YY}}$ precision from $7+8 \mathrm{TeV}$ results

$1 \mathrm{fb}^{-1}=10^{14}$ collisions

## Higgs precision ( $H \rightarrow \gamma \gamma$ ) : optimistic estimate v. luminosity \& time

extrapolation of $\mu_{\mathrm{YY}}$ precision from 7+8 TeV results


The LHC has the statistical potential to take Higgs physics from "observation" to $1-2 \%$ precision

But only if we learn how to connect experimental observations with theory at that precision
$1 \mathrm{fb}^{-1}=10^{14}$ collisions

## how is all of this made quantitative?

whether new-physics searches, Higgs physics, or other SM studies

## UNDERLYING THEORY

$$
\begin{aligned}
& \mathcal{L}=-\frac{1}{q} F_{\mu \nu} F^{\mu \nu} \\
& +i F D_{\psi} \\
& +x_{i} y_{i,} y_{s} \phi+h d \\
& +\left|D_{m} \phi\right|^{2}-V(\phi)
\end{aligned}
$$

how do you make quantitative connection?

## EXPERIMENTAL DATA

|  | how do you make quantitative connection? |
| :---: | :---: |
| $\begin{aligned} & +x_{i} y_{1} y_{s}, \phi+h c \\ & +\left\|p_{n},\right\|^{2}-V(\phi) \end{aligned}$ |  |



## UNDERLYING THEORY

$$
\begin{aligned}
\mathcal{L} & =-\frac{1}{4} F_{\mu \nu} F^{\mu \nu} \\
& +i \neq D \psi \\
& +x_{i} y_{i j} \psi_{s} \phi+h_{L} \\
& +\left|D_{\mu} \phi\right|^{2}-V(\phi)
\end{aligned}
$$

how do you make quantitative connection?

through a chain

of experimental
and theoretical links
[in particular Quantum Chromodynamics (QCD)]

## What are the links?

ATLAS and CMS (big LHC expts.) have written 850 articles since 2014

## links $\equiv$ papers they cite

quantum chromodynamics (QCD) theory papers


knowing what goes into a collision i.e. proton structure [rich UK involvement']


## organising event information ("jets")


the question of organising information from hundreds of particles will come back later
predicting full particle structure that comes out of a collision
[rich UK involvementi]
fraction of ATLAS \& CMS papers that cite them


schematic view of key components of QCD predictions and Monte Carlo event simulation


## schematic view of key components of QCD predictions and Monte Carlo event simulation


schematic view of key components of QCD predictions and Monte Carlo event simulation pattern of particles in MC can be directly compared to pattern in experiment
general purpose Monte Carlo event generators: THE BIG 3


Herwig 7


Pythia 8


Sherpa 2
they do an amazing job of simulation vast swathes of data; collider physics would be unrecognisable without them




## Fundamental experimental calibrations (jets)



Jet energy scale, which feeds into hundreds of other measurements

Largest systematic errors (1-2\%) come from differences between MC generators
(here Sherpa v. Pythia)
$\rightarrow$ fundamental limit on LHC precision potential

## using full event information (quark/gluon tagging)



use more info $\rightarrow$
become more sensitive to MC limitations
up to $35 \%$ differences in MCs v. data
a concern given trend towards use of maximal info, e.g. with machine learning

## Matching with hard process is hitting a limit (e.g. Jäger, Karberg, Scheller 1812.05118)



# Limits effectiveness of current matching methods (here POWHEG) 

Parton structure also gets in way of better (NNLOPS) hard-process + shower matching schemes

## what is a parton shower?

## illustrate with dipole / antenna showers

Gustafson $\mathcal{E}$ Pettersson 1988, Ariadne 1992, main Sherpa $\mathcal{E}$ Pythia8 showers, option in Herwig7, Vincia shower $\mathcal{E}$ (partially) Deductor shower

## At its simplest


iteration of $2 \rightarrow 3$ (or $1 \rightarrow 2$ ) splitting kernel

## in practice: an evolution equation (in evolution scale v, e.g. 1/trans.mom.)



Start with q-qbar state.
Evolve a step in v and throw a random number to decide if state remains unchanged

$$
\frac{d P_{2}(v)}{d v}=-f_{2 \rightarrow 3}^{q \bar{q}}(v) P_{2}(v)
$$

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Evolve a step in v and throw a random number to decide if state remains unchanged
$\frac{d P_{2}(v)}{d v}=-f_{2 \rightarrow 3}^{q \bar{q}}(v) P_{2}(v)$

## in practice: an evolution equation (in evolution scale v, e.g. 1/trans.mom.)



Start with q-qbar state.
Evolve a step in vand throw a random number to decide if state remains unchanged

At some point, rand.numb. is such that state splits ( $2 \rightarrow 3$, i.e. emits gluon). Evolution equation changes

$$
\frac{d P_{3}(v)}{d v}=-\left[f_{2 \rightarrow 3}^{q g}(v)+f_{2 \rightarrow 3}^{g \bar{q}}(v)\right] P_{3}(v)
$$

gluon is part of two dipoles ( $q g, \bar{q} g$ )

## in practice: an evolution equation (in evolution scale v, e.g. 1/trans.mom.)


self-similar evolution
continues until it reaches a nonperturbative scale

## recent directions of parton-shower work?

1. including $2 \rightarrow 4$ (or $1 \rightarrow 3$ ) splittings
2. subleading colour corrections (dipole picture is large $\mathrm{N}_{\mathrm{C}}$ )
3. EW showers

## Including $1 \rightarrow 3$ splitings $(\equiv 2 \rightarrow 4)$

> Jadach et al, e.g. 1504.06849, 1606.01238
> Li \& Skands, 1611.00013
> Höche, Krauss \& Prestel, 1705.00982, Höche \& Prestel, 1705.00742, Dulat, Höche \& Prestel, 1805.03757

$$
\begin{aligned}
D_{j i}^{(0)}(z, \mu)= & \delta_{i j} \delta(1-z) \\
D_{j i}^{(1)}(z, \mu)= & -\frac{1}{\varepsilon} P_{j i}^{(0)}(z) \\
D_{j i}^{(2)}(z, \mu)= & -\frac{1}{2 \varepsilon} P_{j i}^{(1)}(z)+\frac{\beta_{0}}{4 \varepsilon^{2}} P_{j i}^{(0)}(z)+\frac{1}{2 \varepsilon^{2}} \int_{z}^{1} \frac{\mathrm{~d} x}{x} P_{j k}^{(0)}(x) P_{k i}^{(0)}(z / x) \\
& \leftrightarrow(\underbrace{}_{i} \underbrace{0^{600}}_{i}
\end{aligned}
$$

## Including $1 \rightarrow 3$ splittings



Dulat, Höche \& Prestel, 1805.03757

## Hierarchy of subleading colour corrections


cf. also work by Hatta \& Ueda, 1304.6930; Nagy \& Soper papers; some subleading colour also in DIRE2 work


## W emission affects only left-handed quarks

$\rightarrow$ strong polarisation of quarks in unpolarised proton (at high enough energies)

# what does a parton shower achieve? 

not just a question of ingredients,
but also the final result of assembling them together

Dasgupta, Dreyer, Hamilton, Monni \& GPS, 1805.09327

# what should a parton shower achieve? 

not just a question of ingredients,
but also the final result of assembling them together

Dasgupta, Dreyer, Hamilton, Monni \& GPS, 1805.09327

## it's a complicated issue...

> For a total cross section, e.g. for Higgs production, it's easy to talk about systematic improvements (LO, NLO, NNLO, ...). But they're restricted to that one observable

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> With a parton shower (+hadronisation) you produce a "realistic" full set of particles. You can ask questions of arbitrary complexity:
> the multiplicity of particles

- the total transverse momentum with respect to some axis (broadening)
- the angle of 3rd most energetic particle relative to the most energetic one [machine learning might "learn" many such features]


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- the total transverse momentum with respect to some axis (broadening)
- the angle of 3rd most energetic particle relative to the most energetic one [machine learning might "learn" many such features]
how can you prescribe correctness \& accuracy of the answer, when the questions you ask can be arbitrary?


## The standard answer so far

It's common to hear that showers are Leading Logarithmic (LL) accurate.
That language, widespread for multiscale problems, comes from analytical resummations. E.g. for (famous) "Thrust"

$$
T=\max _{\vec{n}_{T}} \frac{\sum_{i}\left|\vec{p}_{i} \cdot \vec{n}_{T}\right|}{\sum_{i}\left|\vec{p}_{i}\right|}
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$$



$$
\begin{aligned}
& \sigma\left(1-T<e^{-L}\right)=\sigma_{t o t} \exp \frac{\left[L g_{1}\left(\alpha_{s} L\right)\right.}{\mathrm{LL}}+\frac{g_{2}\left(\alpha_{s} L\right)}{\mathrm{NLL}}+\frac{\alpha_{s} g_{3}\left(\alpha_{s} L\right)}{\mathrm{NNLL}}+\frac{\left.\alpha_{s}^{2} g_{4}\left(\alpha_{s} L\right)+\cdots\right]}{\mathrm{N}^{2} L L} \\
& \quad\left[\alpha_{s} \ll 1, L \gg 1\right]
\end{aligned}
$$

## The standard answer so far

Sometimes you may see statements like "Following standard practice to improve the logarithmic accuracy of the parton shower, the soft enhanced term of the splitting functions is rescaled by $1+a_{s}(t) /(2 \pi) K$ "

## Questions:

1) Which is it? LL or better?
2) For what known observables does this statement hold?
3) What good is it to know that some handful of observables is LL (or whatever) when you want to calculate arbitrary observables?
4) Does LL even mean anything when you do machine learning?
5) Why only "LL" when analytic resummation can do so much better?

## Our proposal for "minimal" criteria for a shower

## Resummation

Establish logarithmic accuracy for all known classes of resummation:
> global event shapes (thrust, broadening, angularities, jet rates, energy-energy correlations, ...)
> non-global observables (cf. Banfi, Corcella \& Dasgupta, hep-ph/0612282)

- fragmentation / parton-distribution functions
> (multiplicity, cf. original Herwig angular-ordered shower from 1980's)


## Matrix elements

Establish in what sense iteration of (e.g. $2 \rightarrow 3$ ) splitting kernel reproduces $N$-particle tree-level matrix elements for any $N$.

## Examine two showers

- Pythia8 shower: because it's the most widely used
> DIRE shower ( 2015 version, with just $2 \rightarrow 3$ splitting), because it's unique in being available for two General Purpose MC programs (Pythia8 \& Sherpa2)

The results I'll talk about will be the same for both and they'll be limited to fixed order for simplicity (though it's easy enough to generalise to an all-order study)

## Matrix element for single emission: it's correct

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$$
\begin{aligned}
d \mathcal{P}_{\tilde{i} \tilde{j} \rightarrow i k j} & =d \mathcal{P} \\
& =\frac{2 \alpha_{\mathrm{S}} C}{\pi} \frac{d p_{\perp}}{p_{\perp}} d \eta \frac{e^{-2 \eta}}{1+e^{-2 \eta}}+\frac{2 \mathcal{P}_{\mathrm{j}}}{\pi} \frac{2 \alpha_{\mathrm{S}} C}{p_{\perp}} d \eta \frac{e^{2 \eta}}{1+e^{2 \eta}}
\end{aligned}
$$



## Matrix element for single emission: it's correct



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## Matrix element for single emission: it's correct



## Two emissions



## Two emissions

impact of gluon-2 emission on gluon-1 momentum


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## Two emissions


also noticed in 1992 by Andersson, Gustafson \& Sjogren $\rightarrow$ special "fudge" in Ariadne

## Two emissions matrix-element

## Phasespace map

$\stackrel{\ln p_{\perp}}{\longrightarrow} \eta$

analogous effect commented on by Nagy \& Soper for DY recoil, but wider relevance not appreciated?

## Two emissions matrix-element

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## Prevents shower from getting NLL accuracy for any e+e- event shape

Observable $\mathrm{NLL}_{\ln \Sigma}$ discrepancy

| $1-T$ | $0.116_{-0.004}^{+0.004} \bar{\alpha}^{3} L^{3}$ |
| :---: | :---: |
| vector $p_{t}$ sum | $-0.349_{-0.003}^{+0.003} \bar{\alpha}^{3} L^{3}$ |
| $B_{T}$ | $-0.0167335 \bar{\alpha}^{2} L^{2}$ |
| $y_{3}^{\mathrm{cam}}$ | $-0.18277 \bar{\alpha}^{2} L^{2}$ |
| $\mathrm{FC}_{1}$ | $-0.066934 \bar{\alpha}^{2} L^{2}$ |

numerically, coefficients are not large compared to other effects, cf.
$C M W \simeq 0.65 \bar{\alpha}^{2} L^{2}$
(because all these observables are quite inclusive)
but machine-learning uses all info - including large phasespace regions with 100\% deficiencies

## so far took $C_{F}=C_{A} / 2$, i.e. leading $N_{c}$ limit

In real life they're not equal $\&$ common choice for allocating them assigns $\mathrm{C}_{\mathrm{A}} / 2$ to large part of phasespace that is actually gluon emission from quark (i.e. $\mathrm{C}_{\mathrm{F}}$ )


$\underline{\text { Dipole radiation pattern }}$

$$
C_{F}
$$

$\ln p_{\perp}$



## Key observation \#3

CF v. CA/2 issues occurs over a large area $\rightarrow$ double (leading) log effects?

## another view of the colour issue

- The dipole shower phase space partitioning of $g_{2}$ 's radiation pattern is:

q
- Angular ordering implies a partitioning more like the following:



## impact on observables?

Has LL subleading- $\mathrm{N}_{\mathrm{C}}$ effect on 3-jet rates, thrust, but not for things like broadening, 2jet rate (which are physically close the evolution variable, i.e. transverse momentum). E.g. for thrust

$$
\delta \Sigma(L)=-\frac{1}{64} \bar{\alpha}^{2} L^{4}\left(\frac{C_{A}}{2 C_{F}}-1\right)
$$

## closing

## Conclusions

Parton showers are a crucial element in collider physics
Seeing many developments (subleading colour for non-global logarithms, multi-particle emission kernels, etc.)

But maybe we need to go back to foundations:
> improving parton showers is not just a question of better components (e.g. higher-order splitting kernels)
> question of how components are assembled is equally crucial
> we must identify \& state what a parton shower should be achieving
> new studies along these lines are teaching us important things about existing showers

## BACKUP

Constant evolution variable contours in the Lund plane


## two soft emissions : boost dipole partitions back into the event COM

- Consider we emitted soft gluon gı from hard qq, $\overline{\text { so }}$ we end up with a qgı $\overline{\text { and }}$ a gıq dipole:

- To get us from the event COM to the gıq dipole COM [blue line] requires a BIG BOOST $\rightarrow$
- Dipole partitioned at $\eta=0$ in that frame:
- 



- To get us back to the event COM from the giq dipole COM undo the same BIG BOOST $\leftarrow$

- In event COM partition comes out very close to $q$; instead of equidistant in angle between $\mathrm{g} \boldsymbol{\&} \mathrm{q}$
jet with $R=0.4, p_{t}=200 \mathrm{GeV}$

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## organise phasespace: Lund diagrams



## The path forward: collect 20-30x more collisions by ~2035


> Suppose we had a choice between

- HL-LHC ( $14 \mathrm{TeV}, 3 \mathrm{ab}^{-1}$ )
> or going to higher c.o.m. energy but limited to $80 \mathrm{fb}^{-1}$.
> How much energy would we need to equal the HL-LHC?

| today's <br> reach | HL-LHC <br> reach <br> $\left(13 \mathrm{TeV}, 80 f \mathrm{f}^{-1}\right)$ | energy needed <br> for same reach <br> with $80 f \mathrm{TeV} \mathrm{3ab}^{-1}$ |
| :---: | :---: | :---: |
| $4.7 \mathrm{TeV} \mathrm{SSM} \mathrm{Z'}$ | 6.7 TeV | 20 TeV |
| 2 TeV weakly <br> coupled Z' | 3.7 TeV | 37 TeV |
| 680 GeV <br> chargino | 1.4 TeV | 54 TeV |

## Hard processes: to 3rd order (NNLO) in perturbation theory strong coupling constant ( $a_{s}$ )



WH diff., Ferrera, Grazzini, Tramontano $\mathrm{Y}-\mathrm{\gamma}$, Catani et al.

Hj (partial), Boughezal et al.
ttbar total, Czakon, Fiedler, Mitov
Z- $\rangle$, Grazzini, Kallweit, Rathlev, Torre
jj (partial), Currie, Gehrmann-De Ridder, Glover, Pires
ZZ, Cascioli it et al.
ZH diff., Ferrera, Grazzini, Tramontano WW, Gehrmann et al.
ttbar diff., Czakon, Fiedler, Mitov
Z- Z , W- Y , Grazzini, Kallweit, Rathlev
-Hj , Boughezal et al.
-Wj, Boughezal, Focke, Liu, Petriello
 -Hj , Boughezal et al.
 MCFM at NNLO, Boughezal et al.

## $\mathrm{p}_{\mathrm{tz}}$, Gehrmann-De Ridder et al.

MCFM at NNLO, Berger, Gao, C.-Yuan, Zhu
MCFM at NNLO, de Florian et al.
$\mathrm{p}_{\mathrm{tH}}$, MCFM at NNLO, Chen et al.
$\mathrm{p}_{\mathrm{tz}}$, Gehrmann-De Ridder et al.
jj, Currie, Glover, Pires
as of April 2017
yX, Campbell, Ellis, Williams
yj, Campbell, Ellis, Williams

