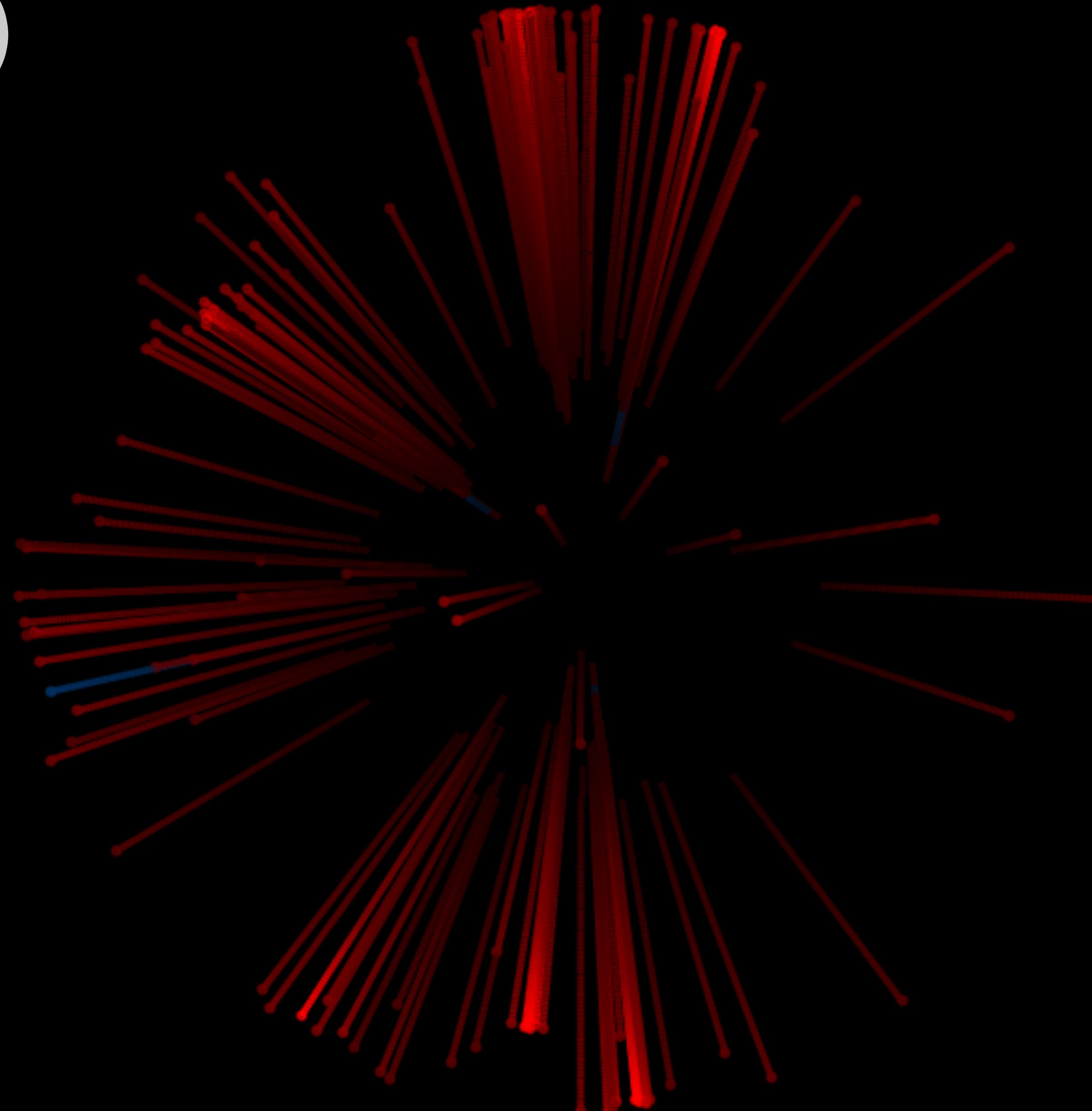
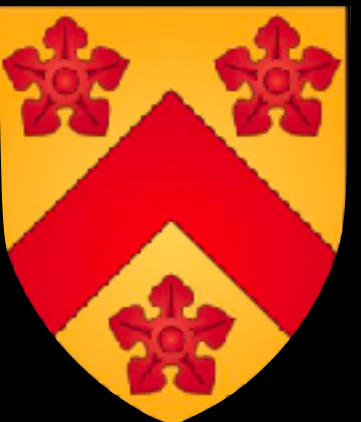


Jets through the LHC era (a personal view)

5th KEK-PH, Jet Physics
via Zoom
30 November 2021



Gavin Salam
Rudolf Peierls Centre for
Theoretical Physics
& All Souls College, Oxford



The context of this talk: LHC physics (colour-coded by directly-probed energy scales)

**Standard-model
physics
(QCD & electroweak)**

100 MeV - 4 TeV

top-quark physics

170 GeV - 0(TeV)

Higgs physics

125 GeV - 500 GeV

**direct new-particle
searches**

100 GeV - 8 TeV

**flavour physics
(bottom & some charm)**

1 - 5 GeV

heavy-ion physics

100 MeV - 500 GeV

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Key high-energy physics goals (my view)

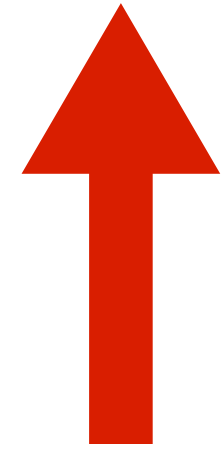
1. Establish the structure of the Higgs sector of the SM
2. Search for signs of physics beyond the SM, direct (incl. dark matter candidates, SUSY, etc.) and indirect
3. Measure SM parameters, proton structure (PDFs), establish theory-data comparison methods, etc.

The Lagrangian and Higgs interactions: two out of three qualitatively new!

$$\mathcal{L}_{\text{SM}} = \dots + |D_{\mu}\phi|^2 + \psi_i y_{ij} \psi_j \phi - V(\phi)$$

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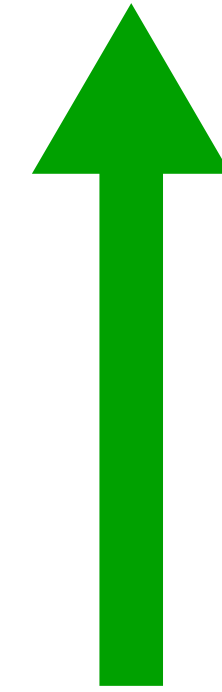
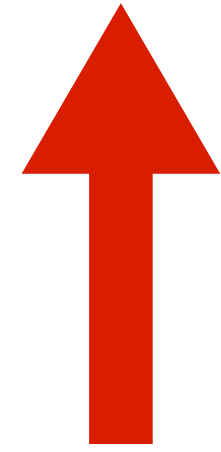
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like those in QED, QCD, EW,
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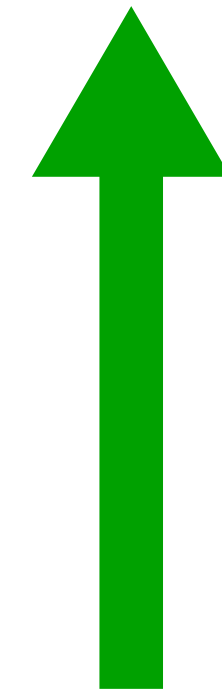
Yukawa interactions. Responsible for fermion masses, and induces “fifth force” between fermions. **Direct study started only in 2018!**

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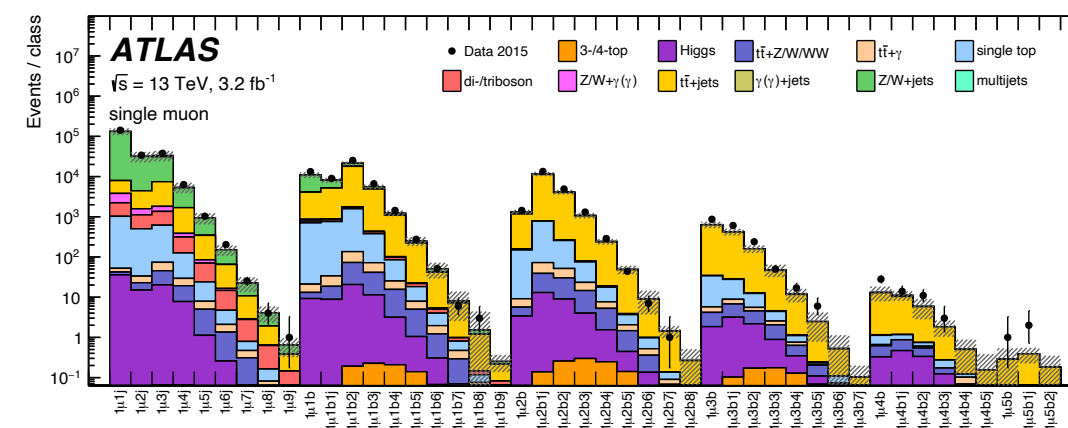
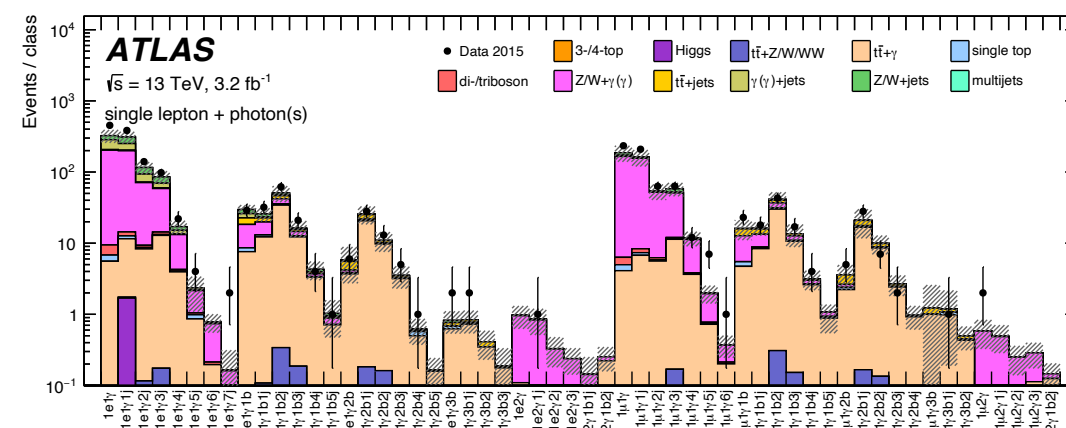
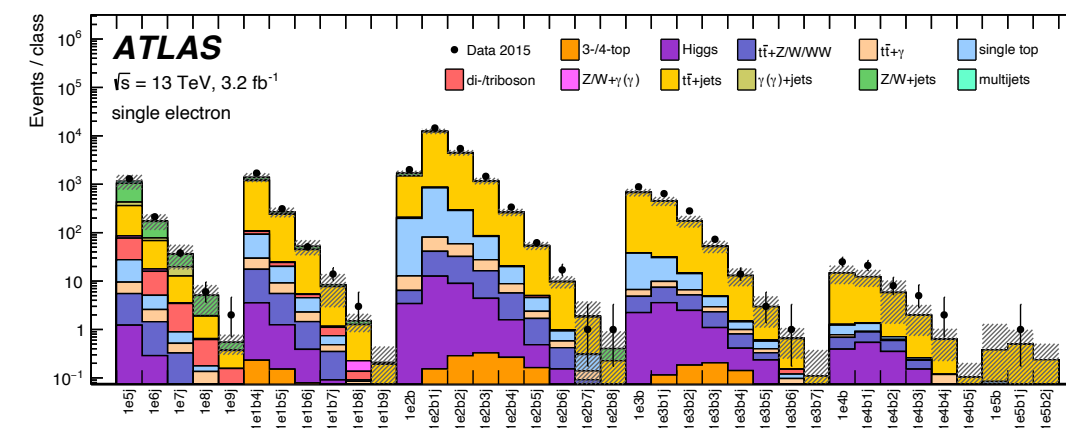
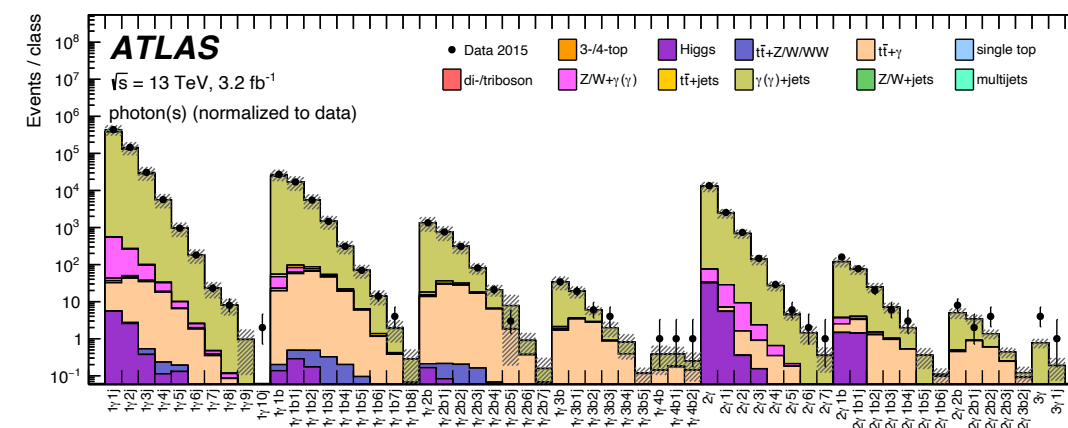
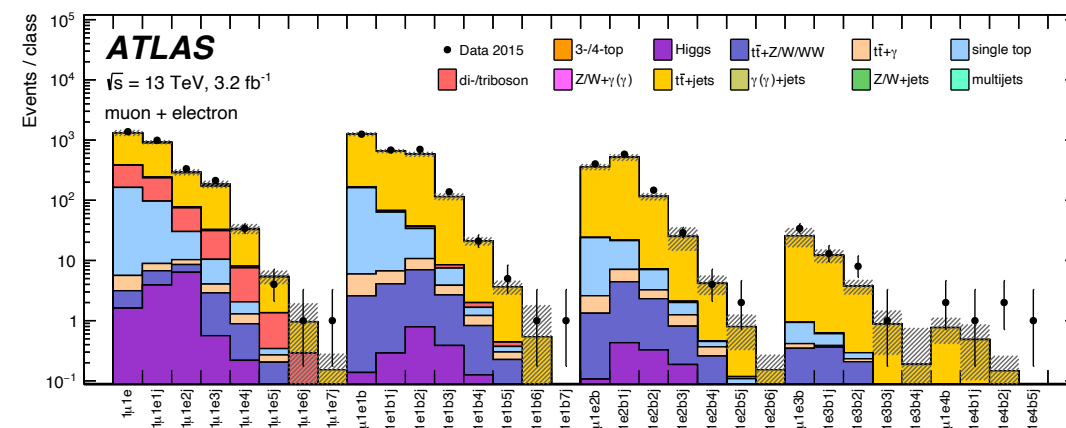
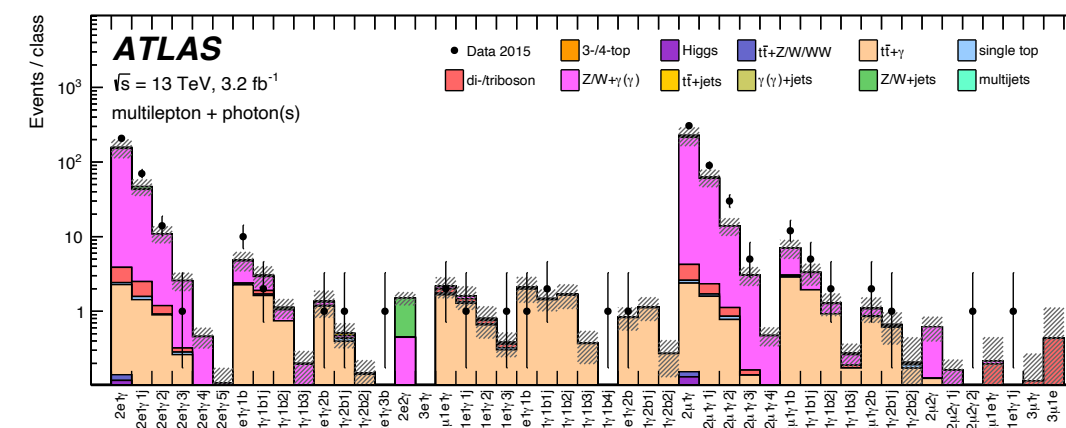
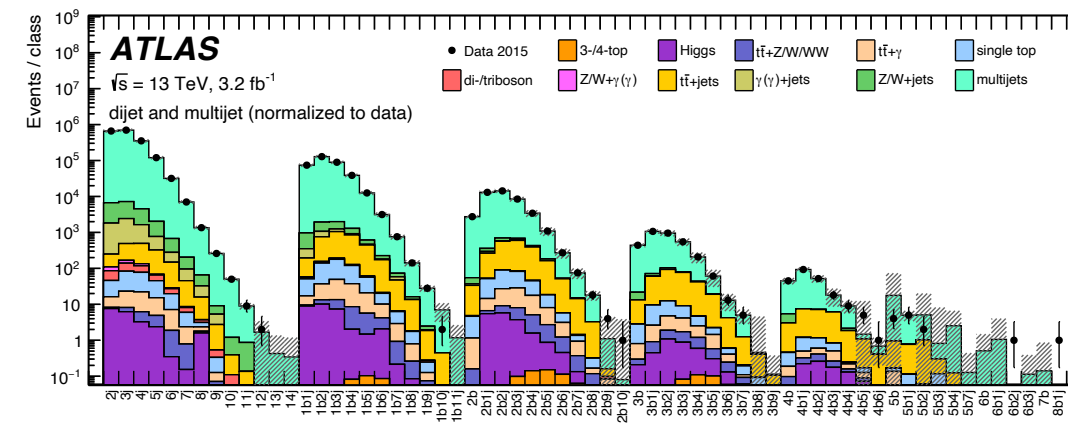
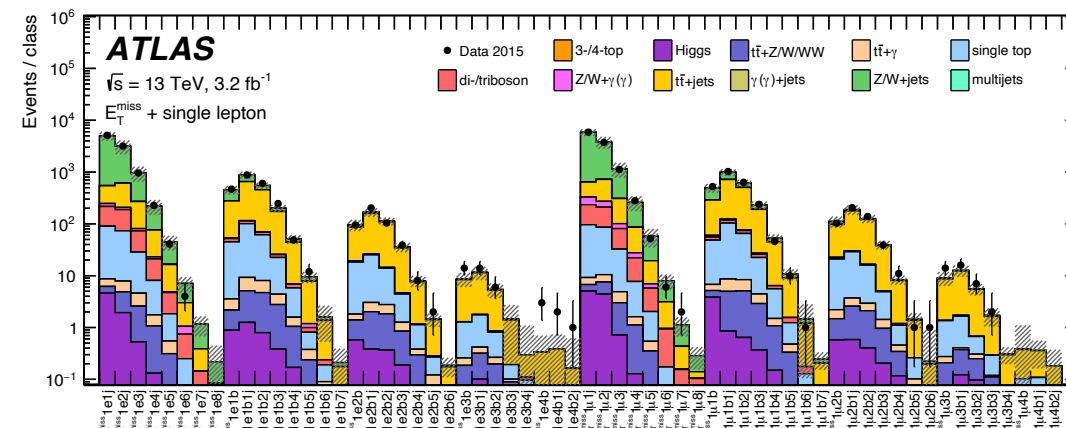
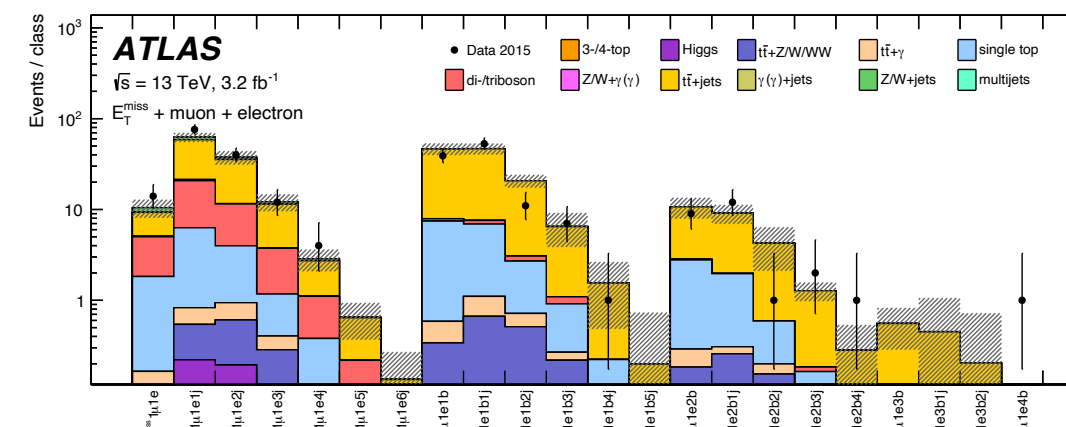
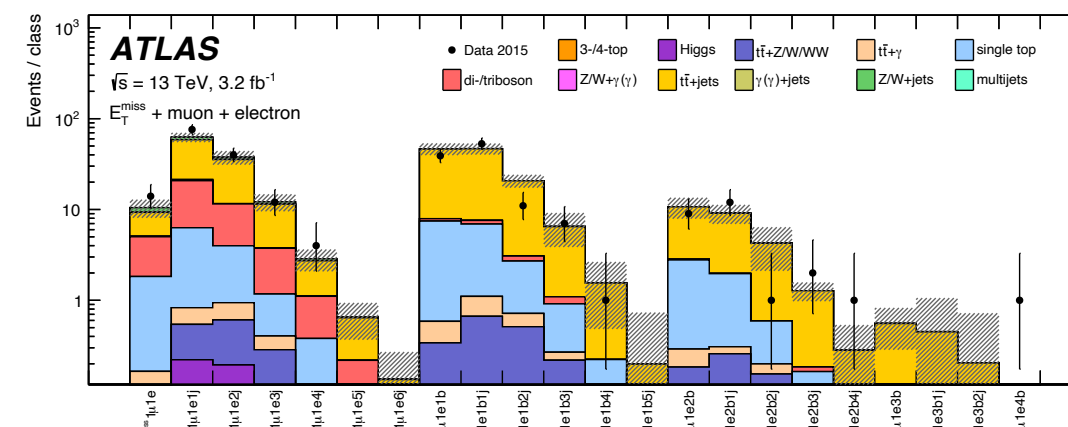
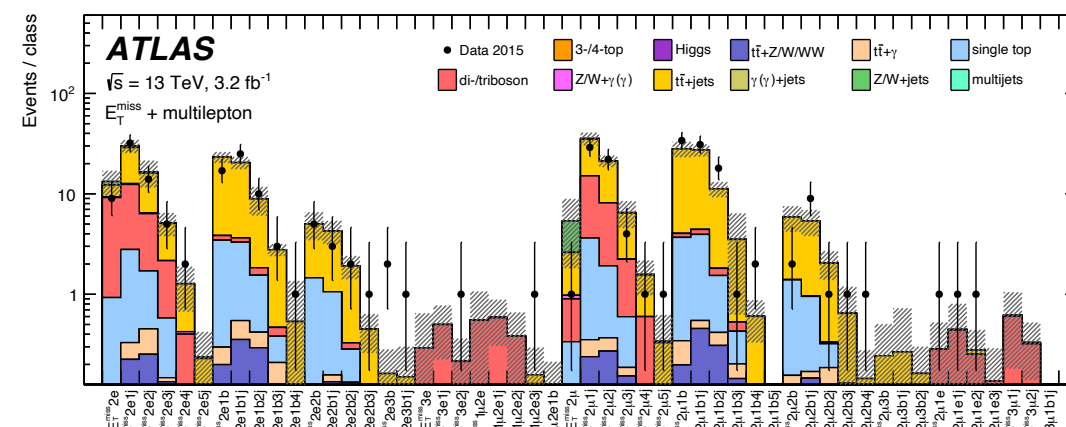
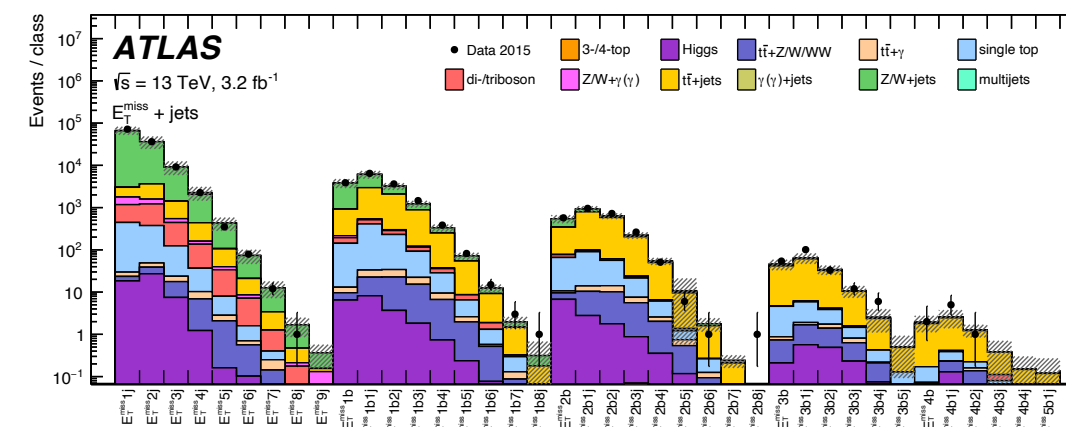
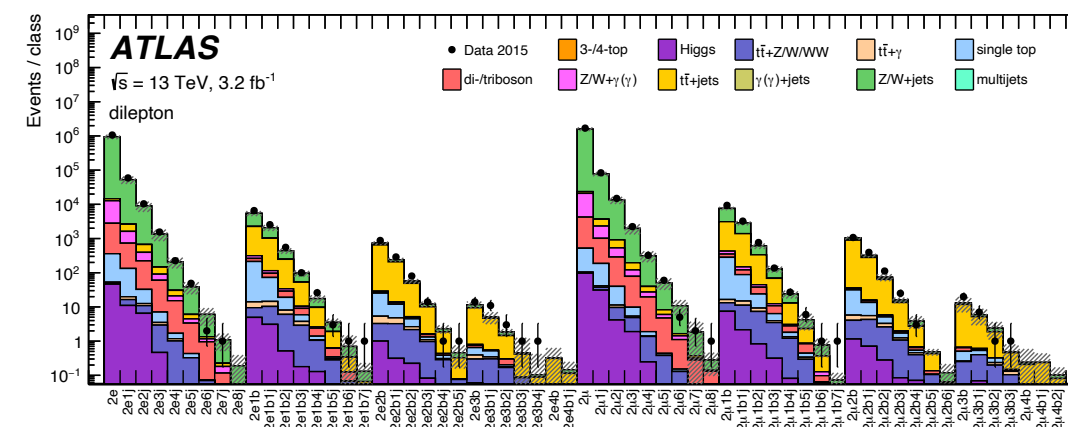
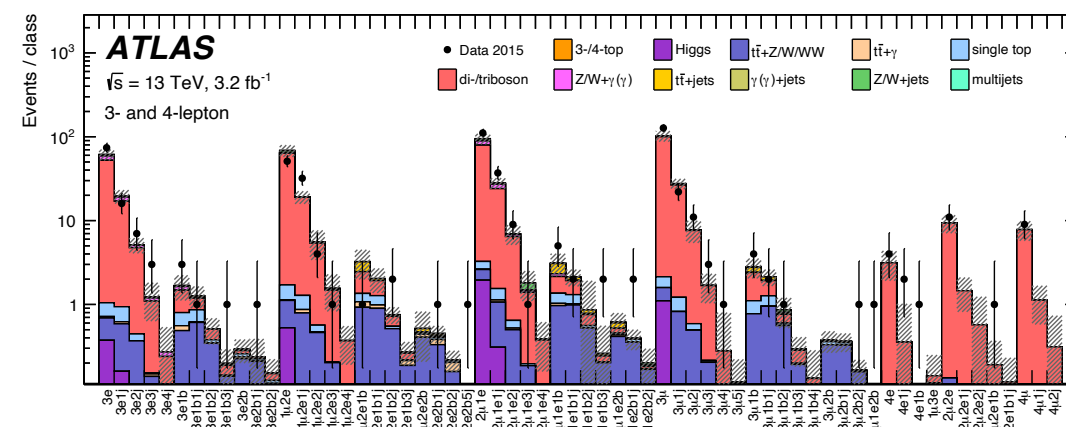


Yukawa interactions. Responsible for fermion masses, and induces “fifth force” between fermions. **Direct study started only in 2018!**



Higgs potential → self-interaction (“sixth?” force between scalars). Holds the SM together. **Unobserved**

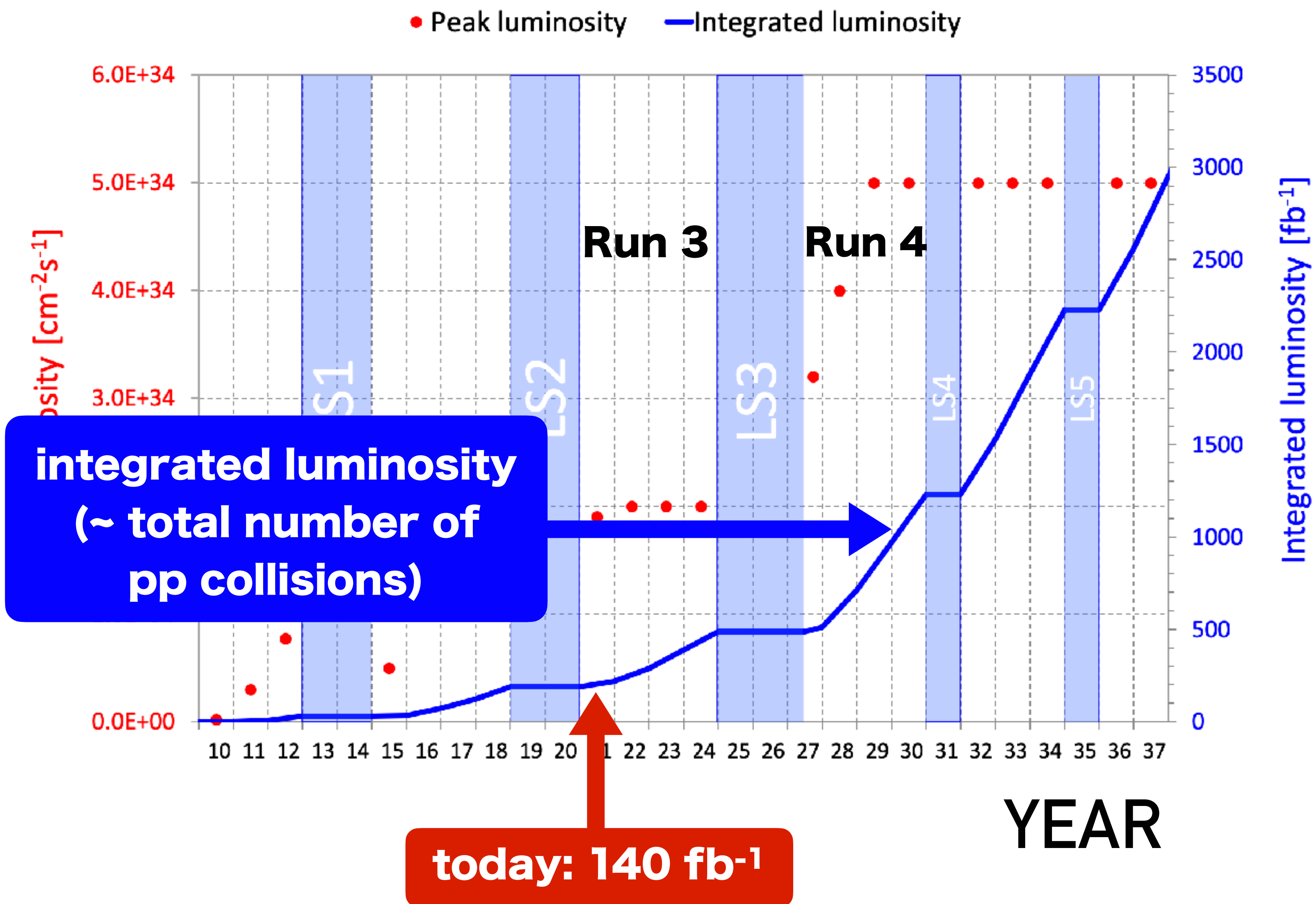
Broadband searches (here an example with 704 event classes, >36000 bins)



Just one illustration
out of many searches
at the LHC

ATLAS, arXiv:1807.07447
13 TeV, 3.2 fb⁻¹
General search

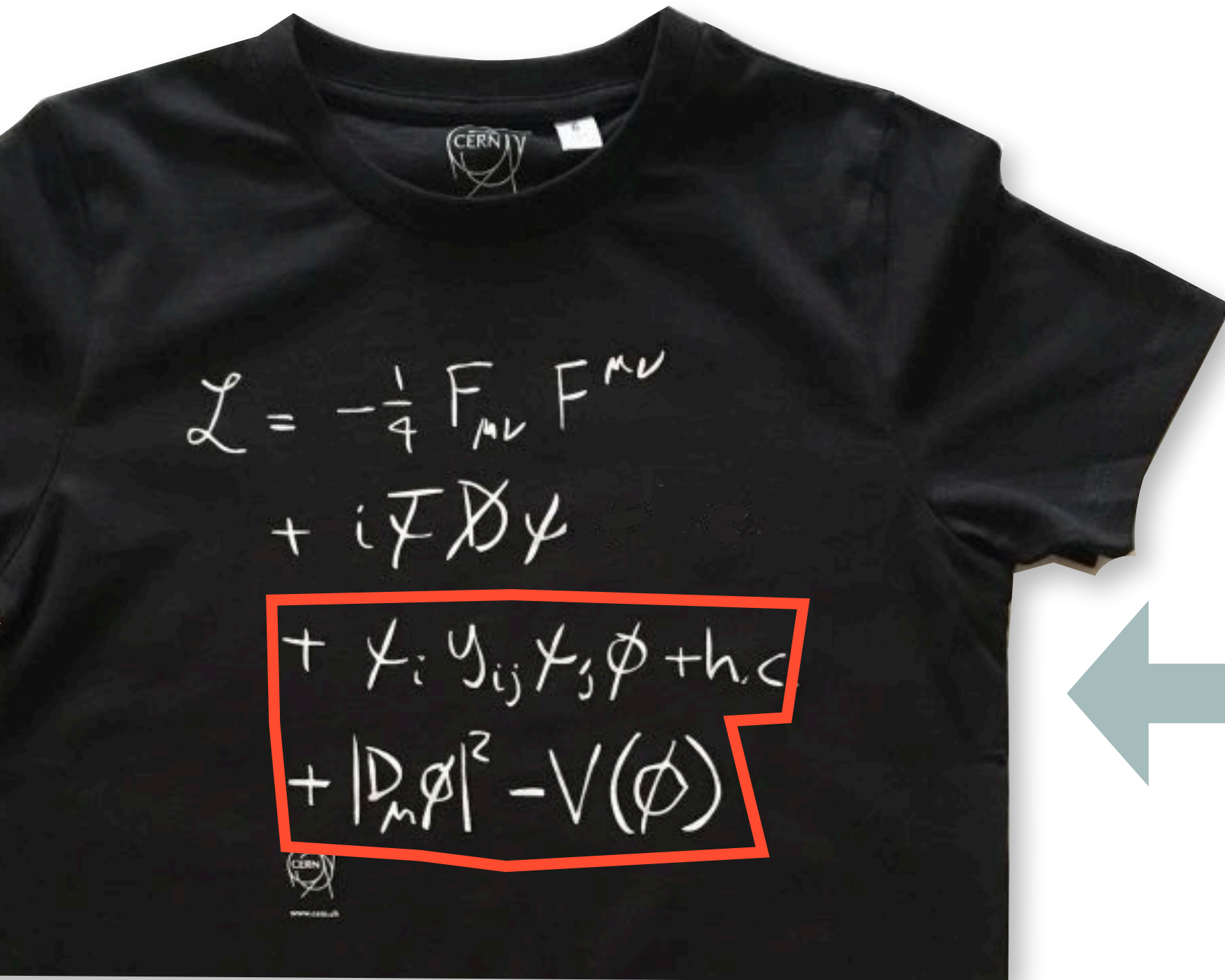
LHC luminosity v. time



year	lumi (fb^{-1})	
2020	140	
2025	450	($\times 3$)
2030	1200	($\times 8$)
2037	3000	($\times 20$)

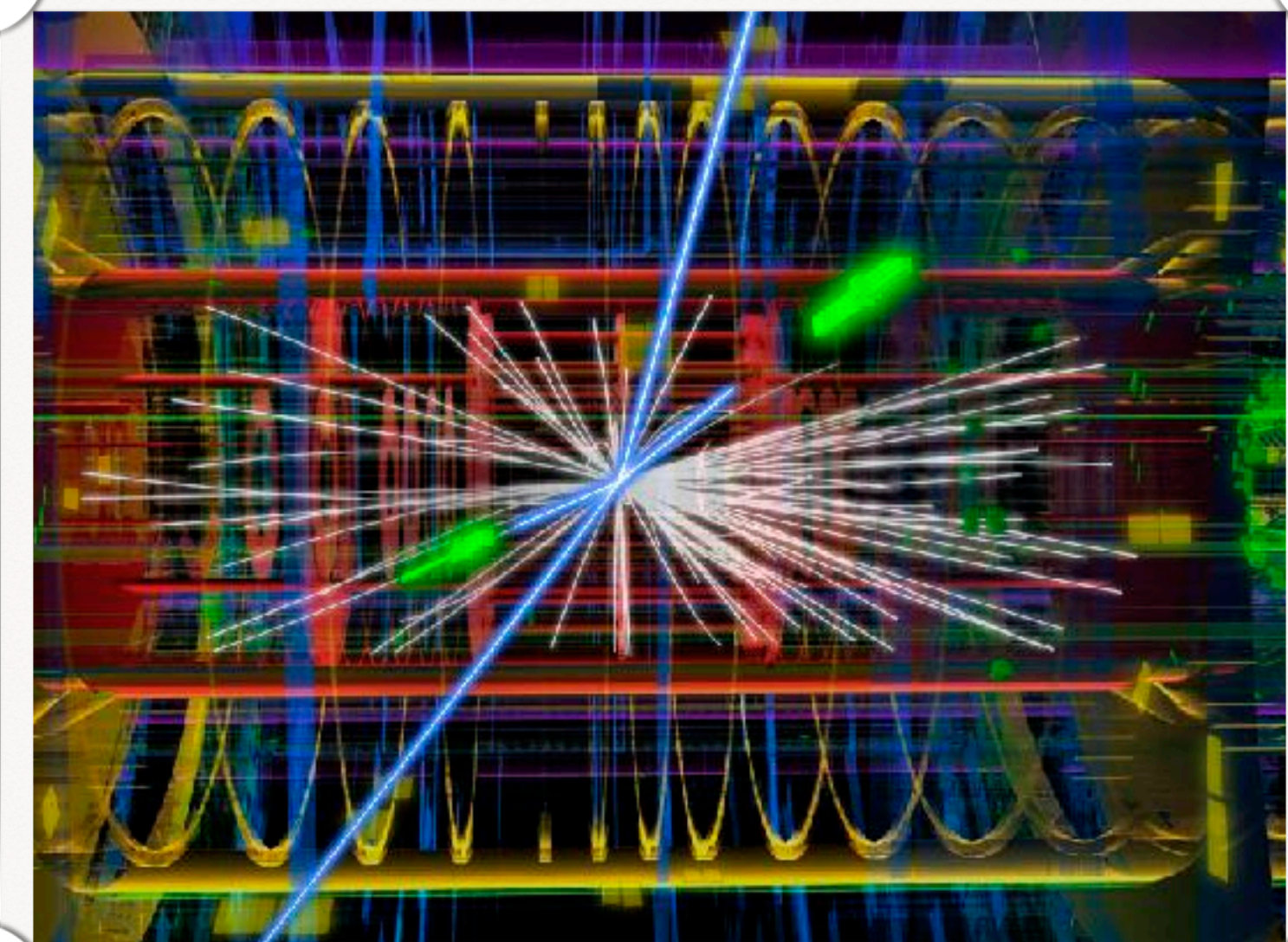
95% of collisions still to be delivered

UNDERLYING THEORY

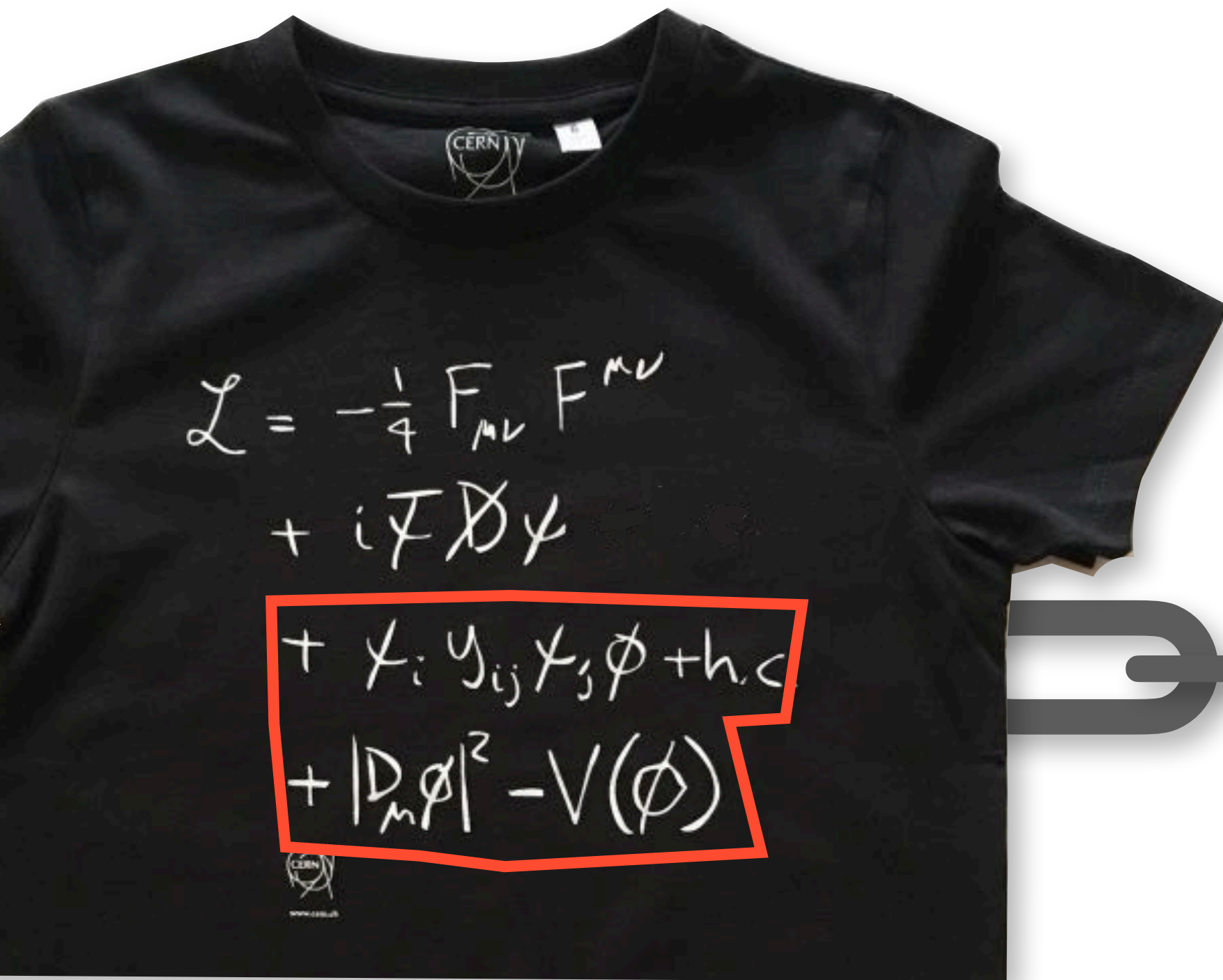


*how do you make
quantitative
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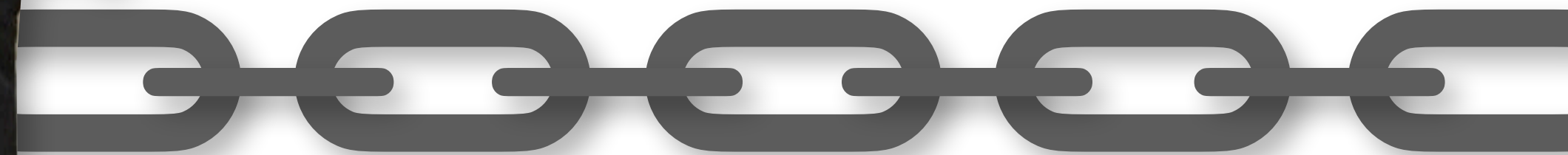
EXPERIMENTAL DATA



UNDERLYING THEORY



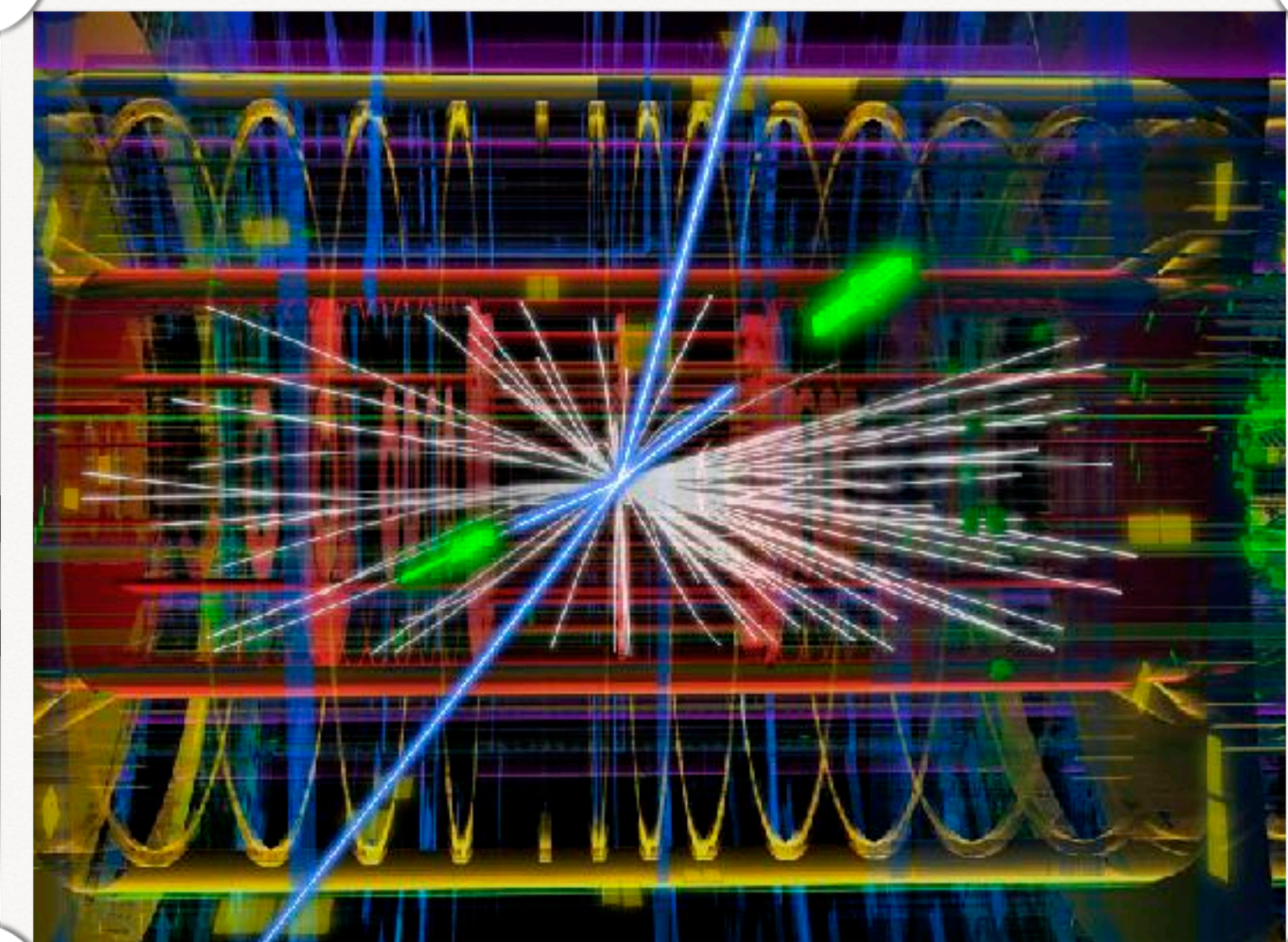
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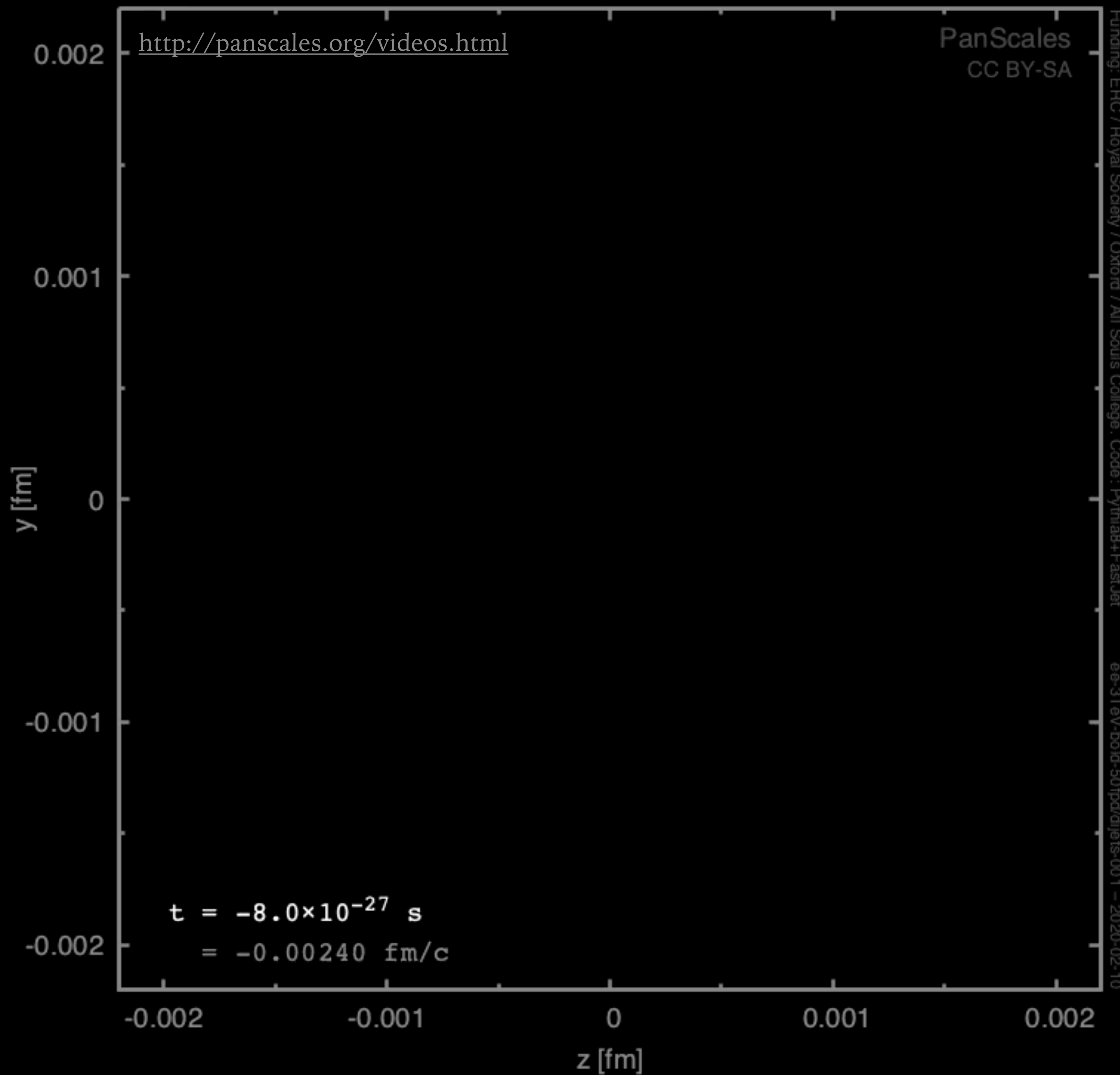


*through a chain
of experimental
and theoretical links*

[in particular Quantum Chromodynamics (QCD)]

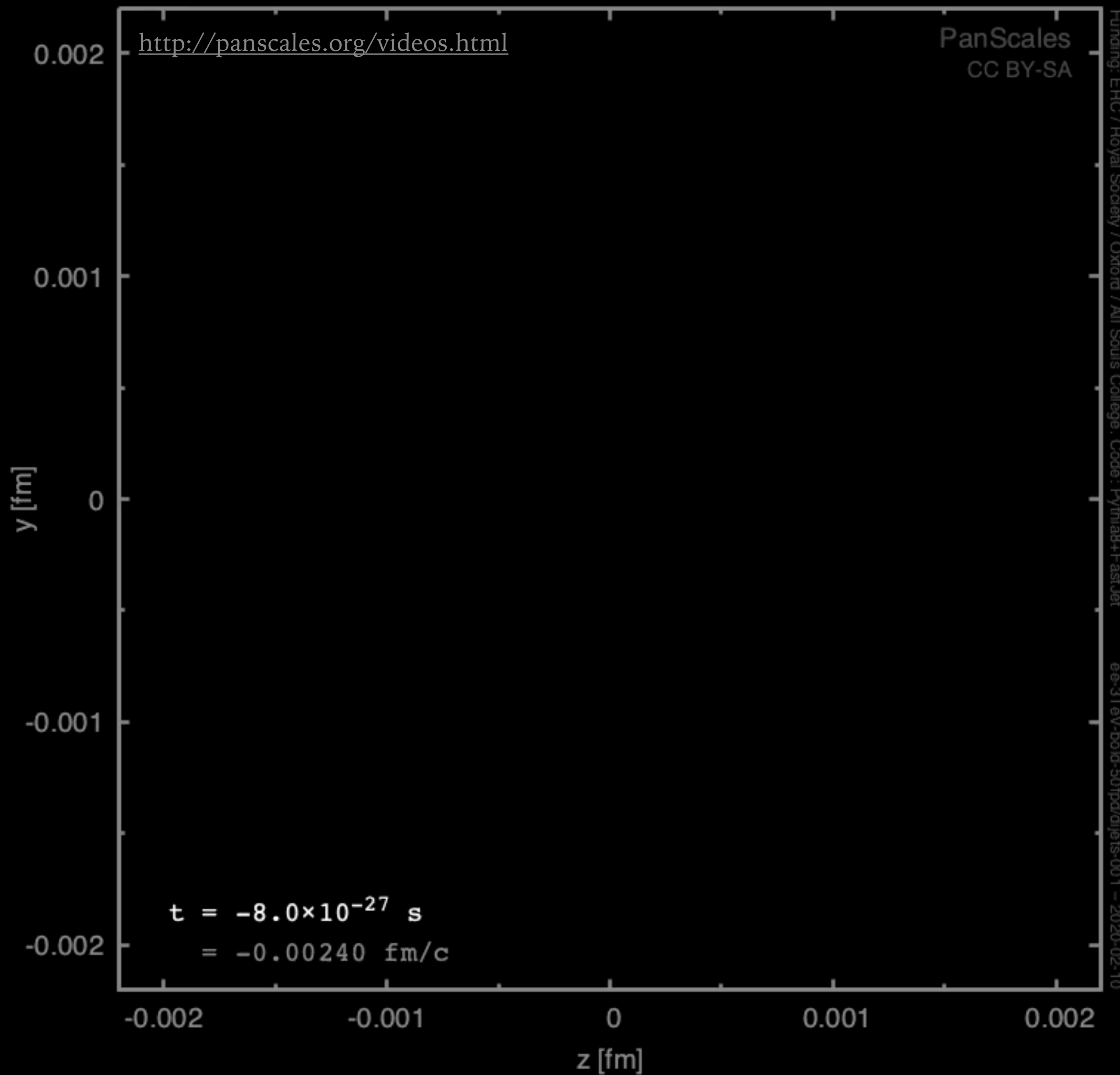
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- incoming beam particle
- intermediate particle (quark or gluon)
- final particle (hadron)

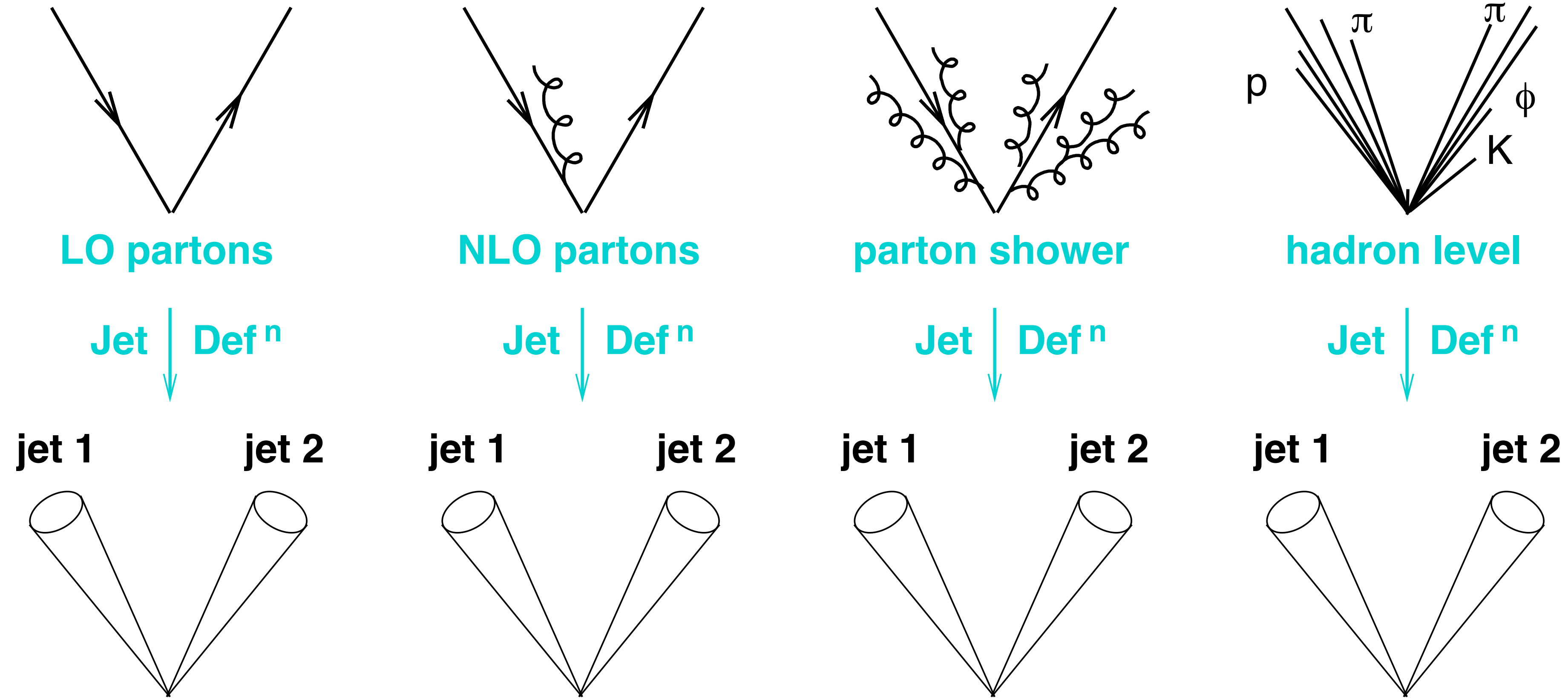
Event evolution spans 7 orders of magnitude in space-time



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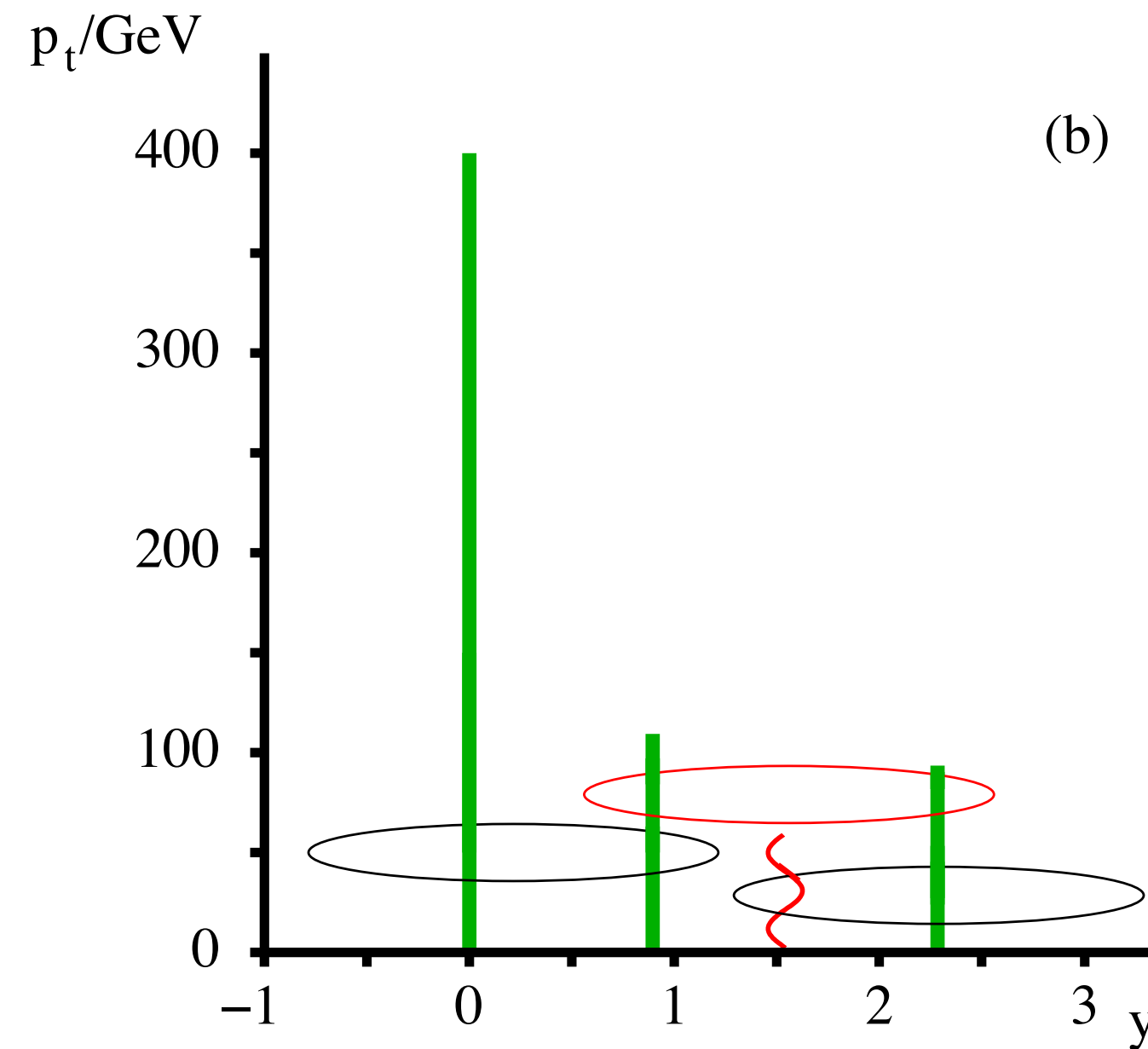
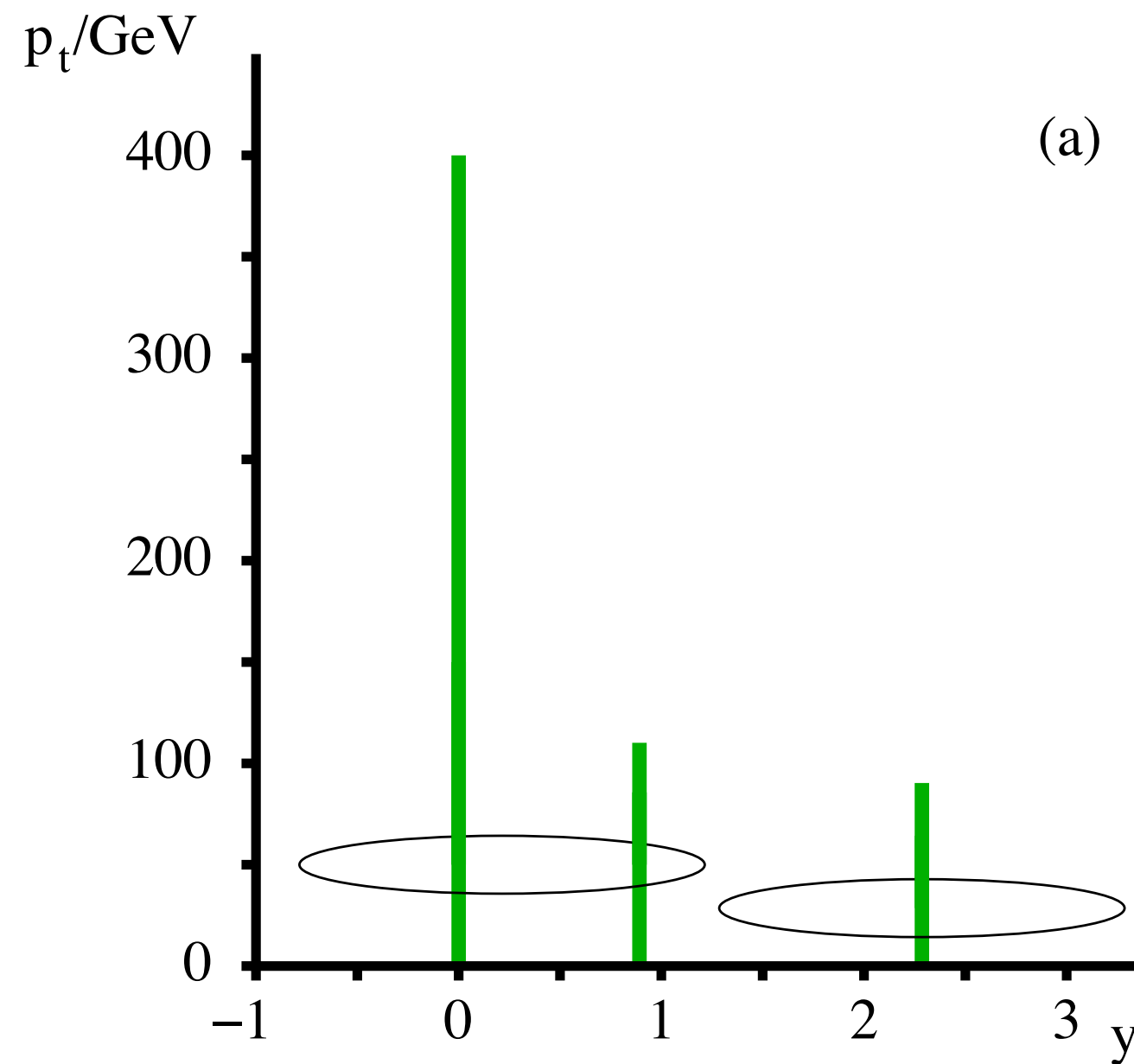
Jet finding projects high-dim info to low number of dimensions in a robust, reproducible way



Projection to jets should be resilient to QCD effects

step back ~ 15 years

pre-LHC hadron-collider jet algorithms (cone algorithms) were infrared unsafe

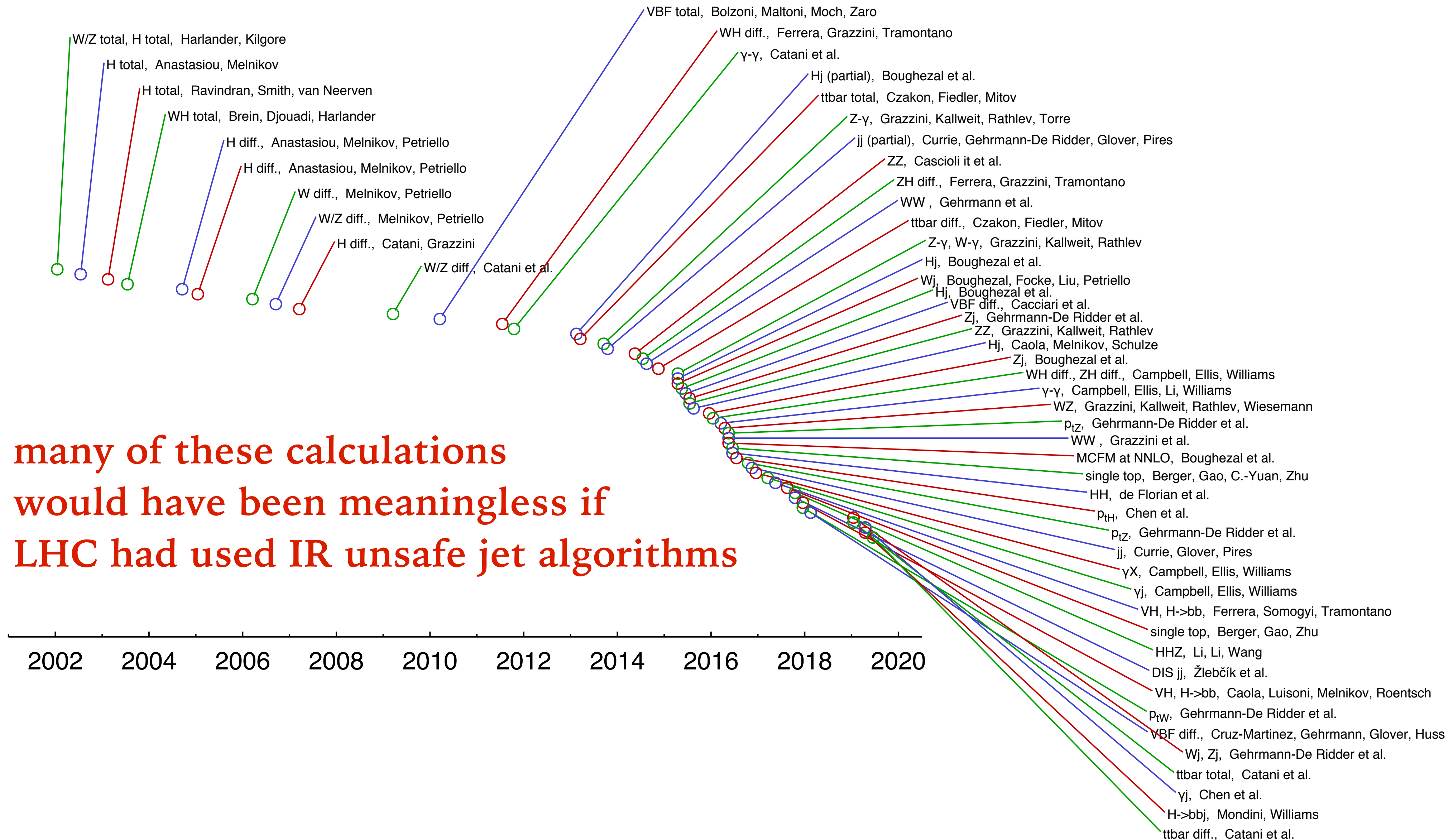


infrared unsafety =
 strong sensitivity to low-
 momentum perturbations of
 event structure
 → uncancelled ∞ in
 perturbative QCD calculations

Observable	1st miss cones at	Last meaningful order
Inclusive jet cross section	NNLO	NLO
$W/Z/H + 1$ jet cross section	NNLO	NLO
3 jet cross section	NLO	LO
$W/Z/H + 2$ jet cross section	NLO	LO
jet masses in 3 jets, $W/Z/H + 2$ jets	LO	none

GPS & Soyez
[arXiv:0704.0292](https://arxiv.org/abs/0704.0292)

yet we were on the cusp of a revolution in precision QCD (NNLO) calculations



possible solution

“sequential recombination” k_t algorithm

Catani, Dokshitzer, Seymour & Webber '93
Ellis & Soper '93

1. for all particle pairs, i, j , find smallest of

$$d_{ij} = \min(p_{ti}^2, p_{tj}^2) \frac{\Delta R_{ij}^2}{R^2}, \quad d_{iB} = p_{ti}^2$$

2. if d_{ij} recombine particles into one
3. if d_{iB} declare i to be a jet

Accept all jets above some some threshold
transverse momentum

Marunouchi, Tokyo (c. 2009)



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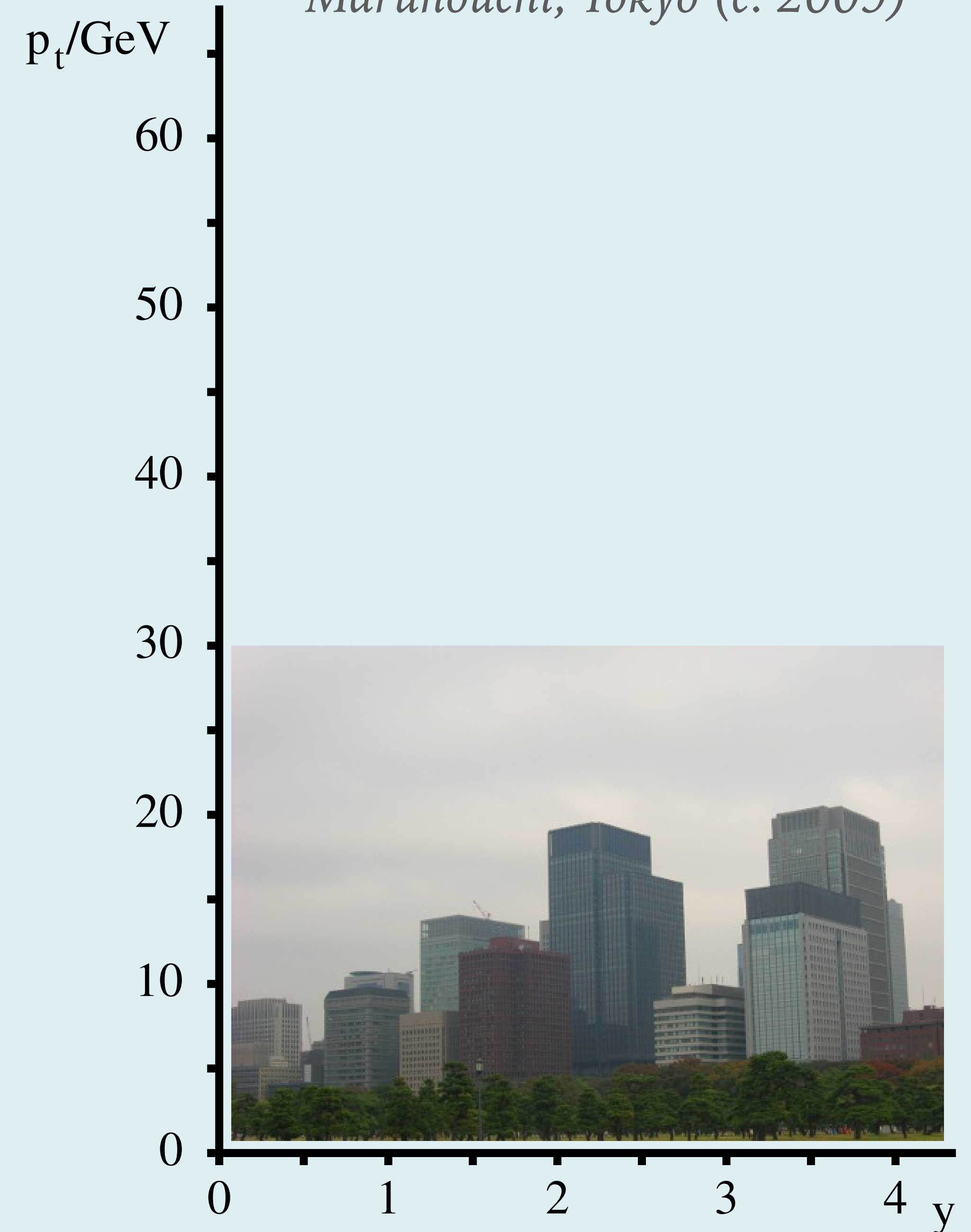
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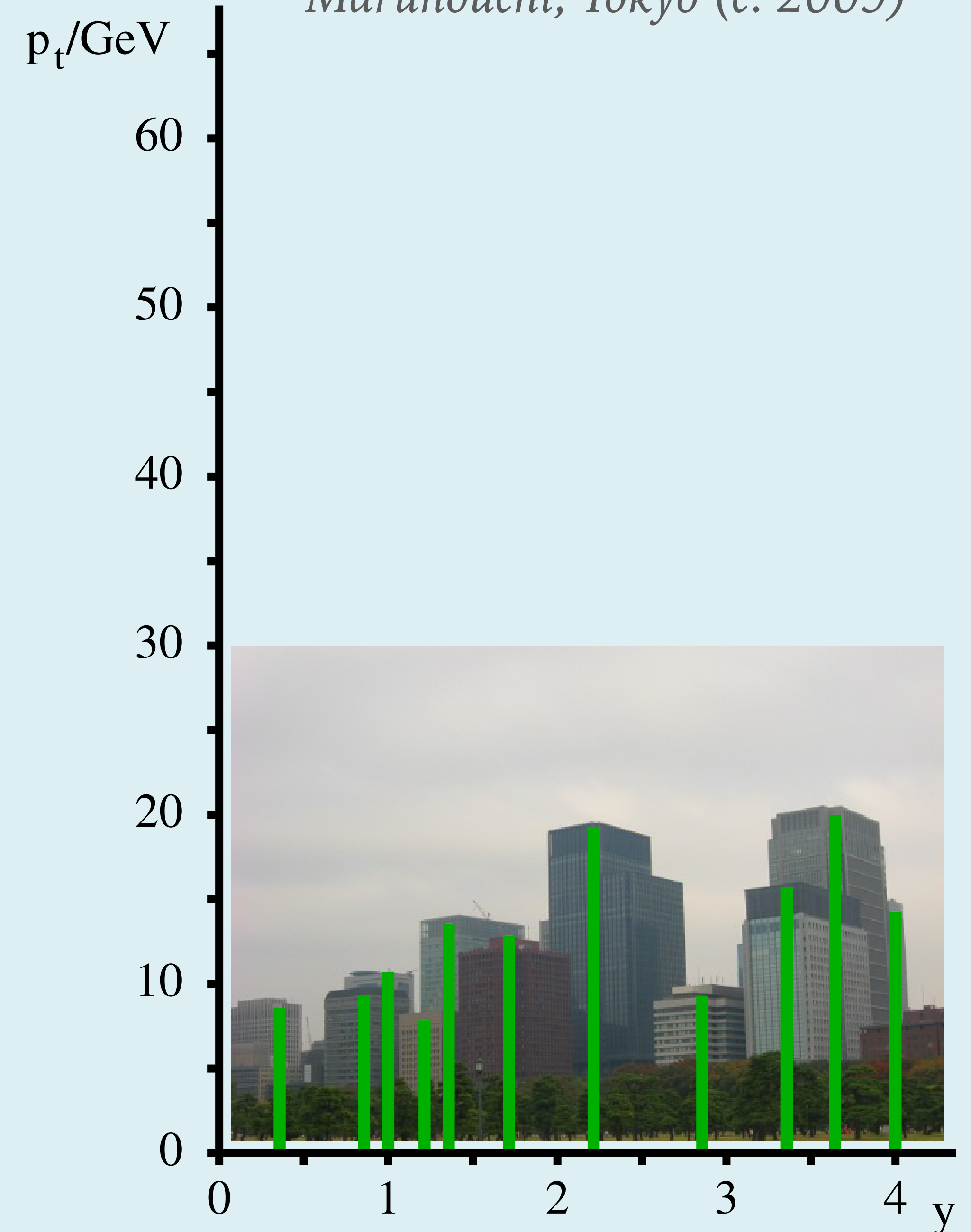
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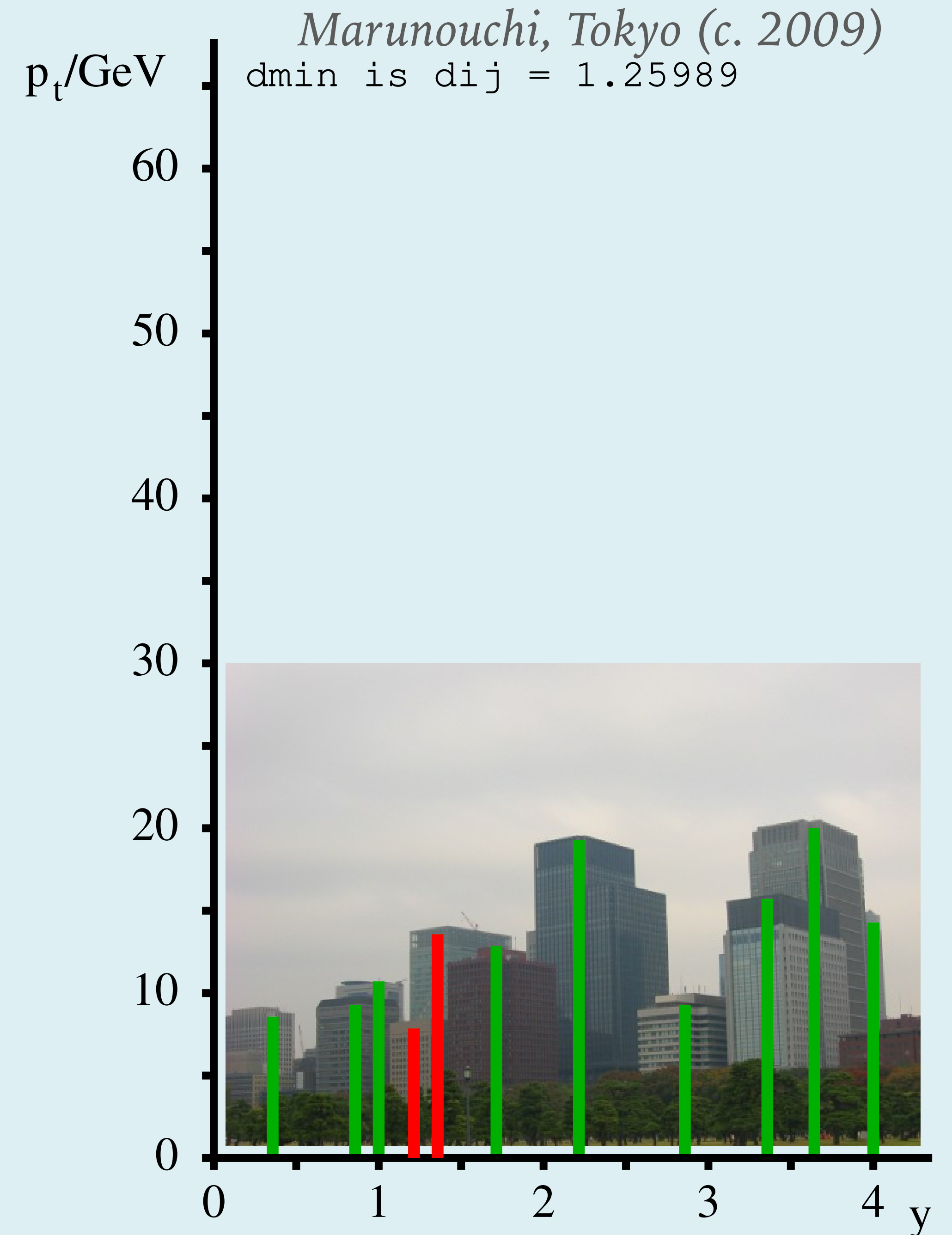
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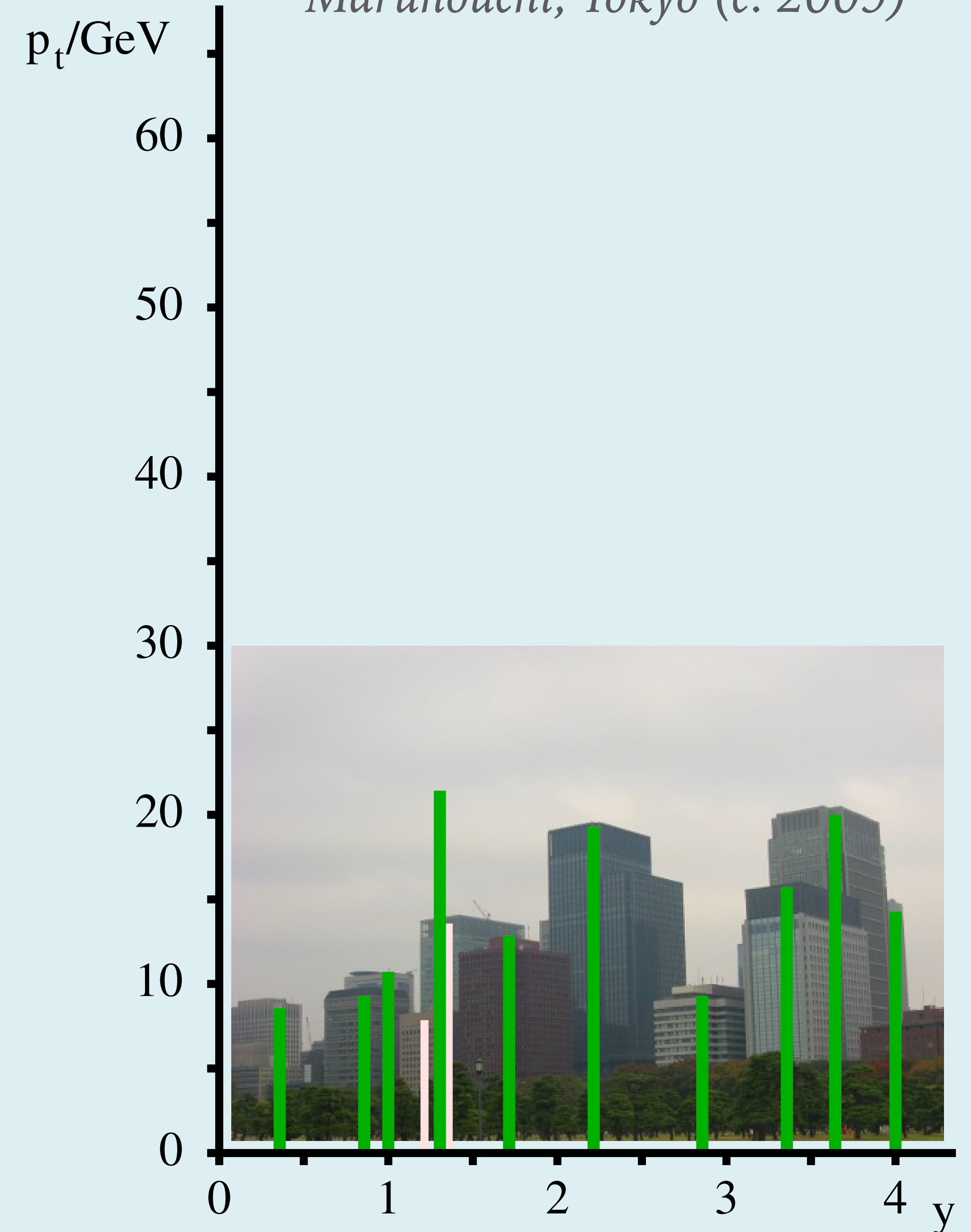
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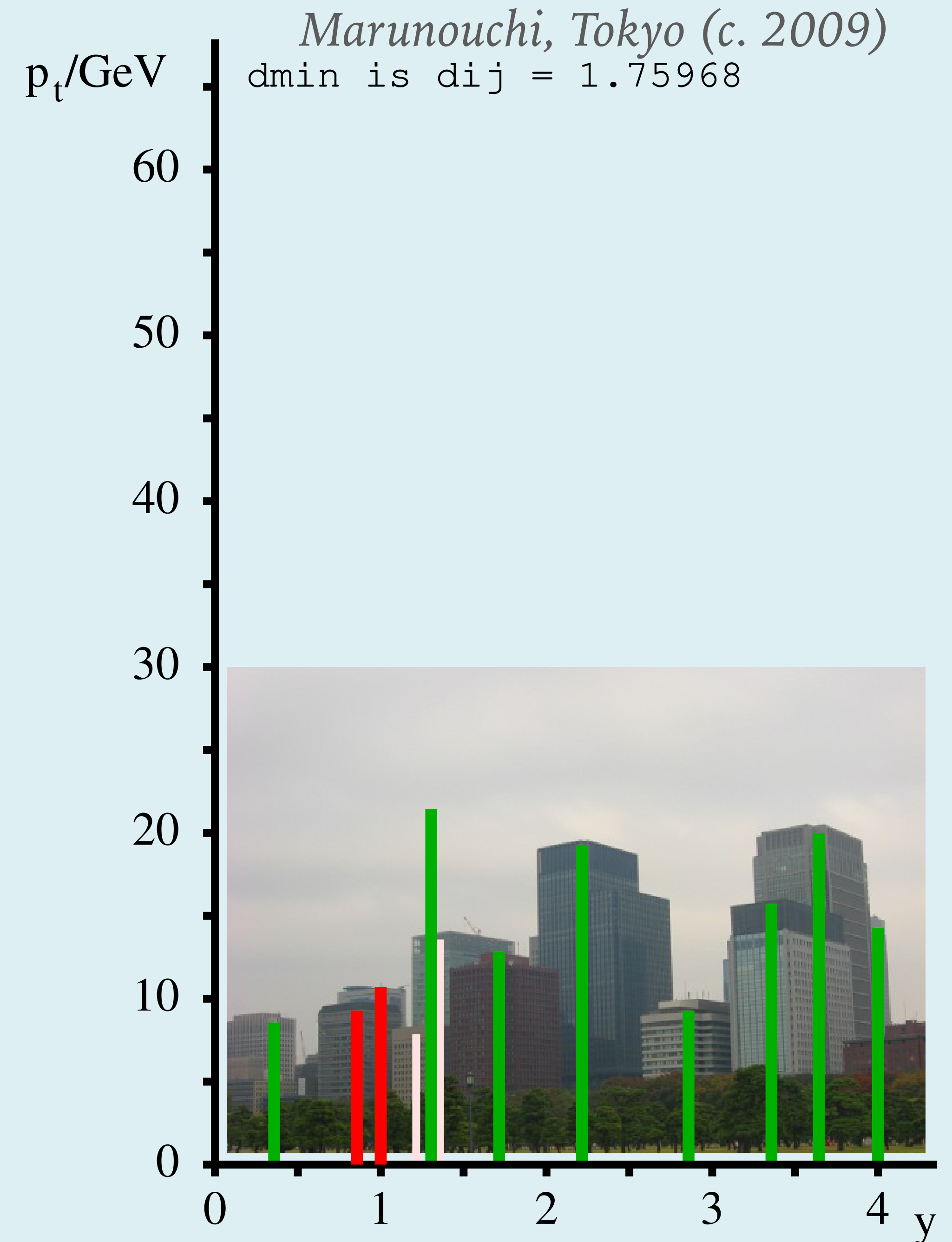
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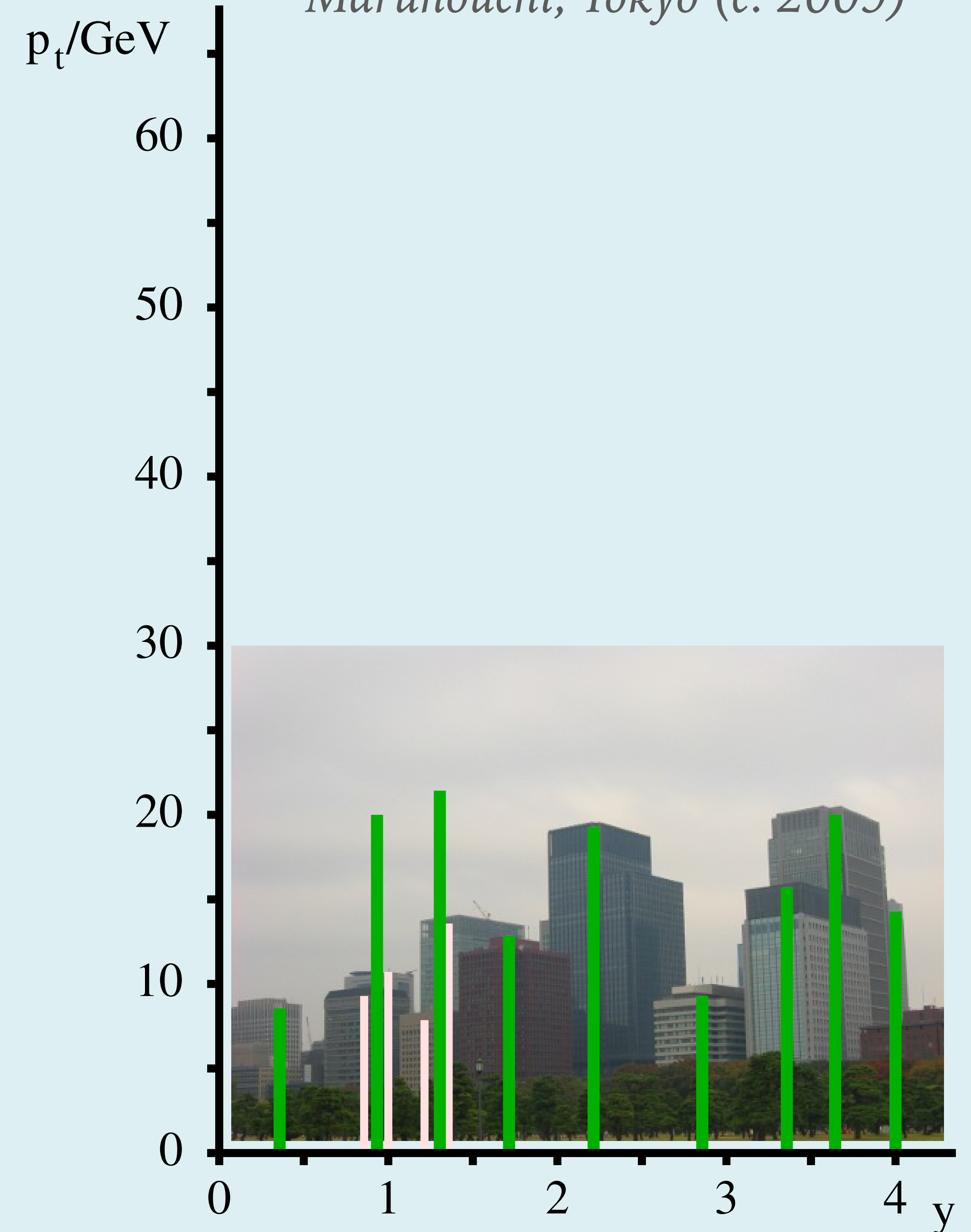
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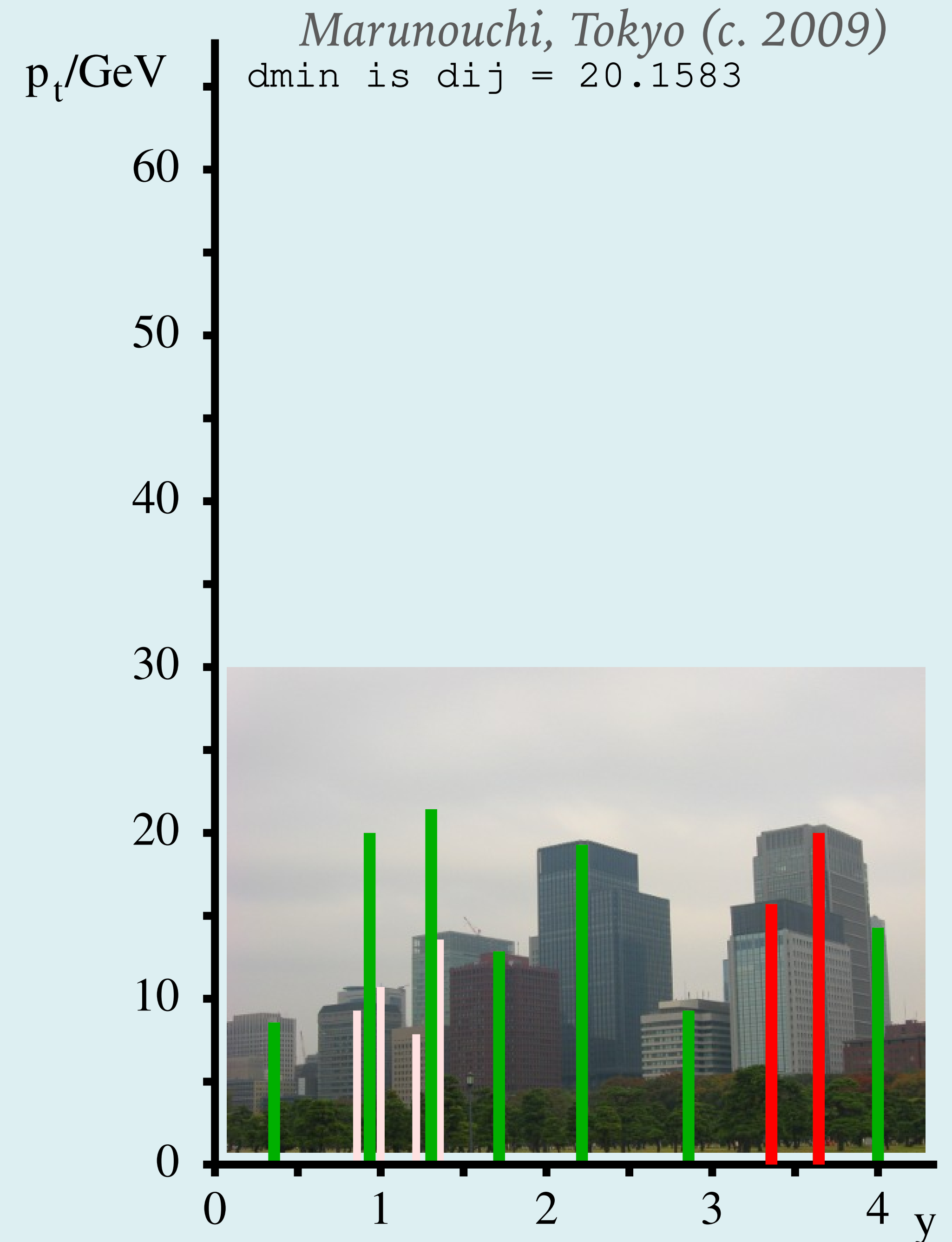
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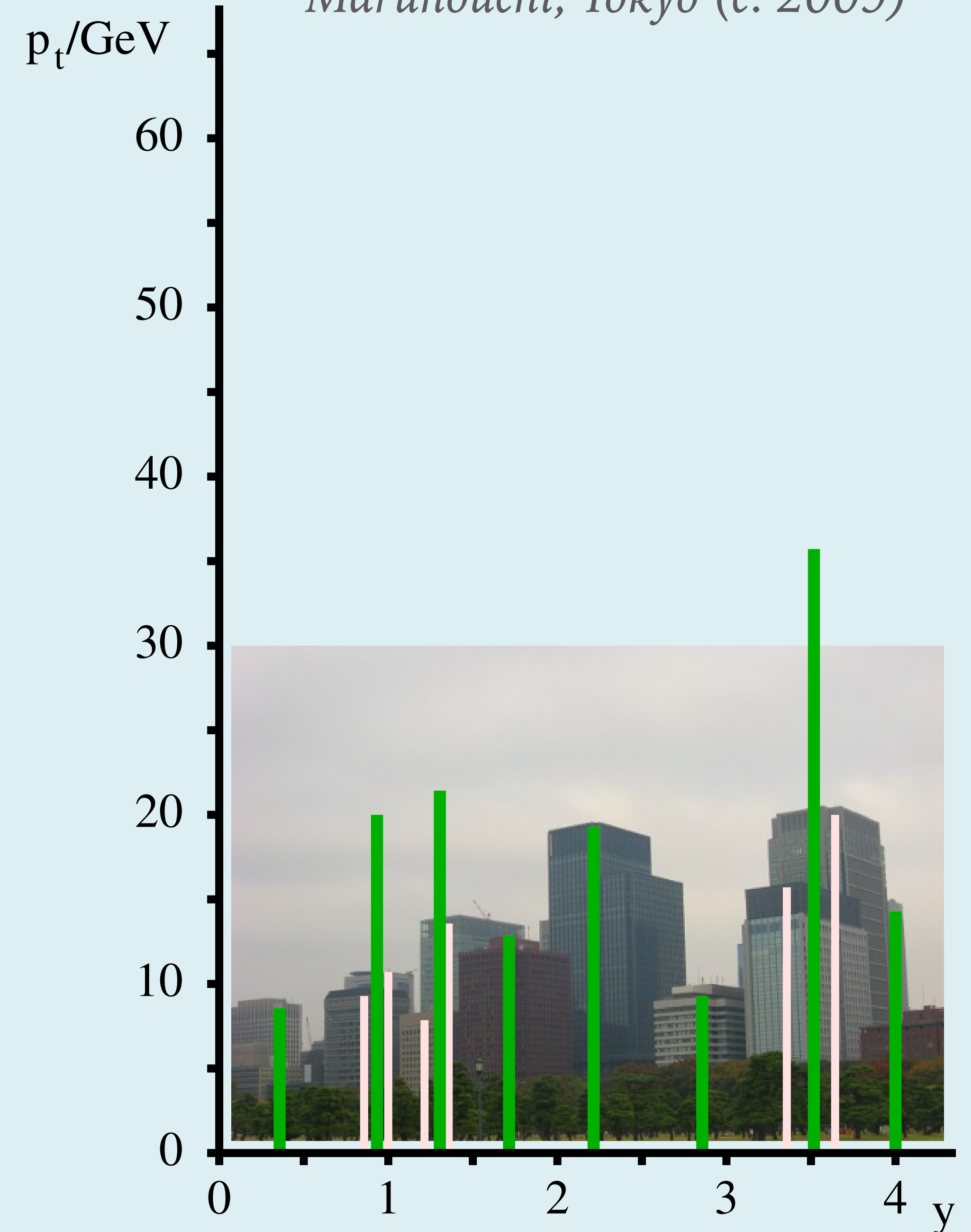
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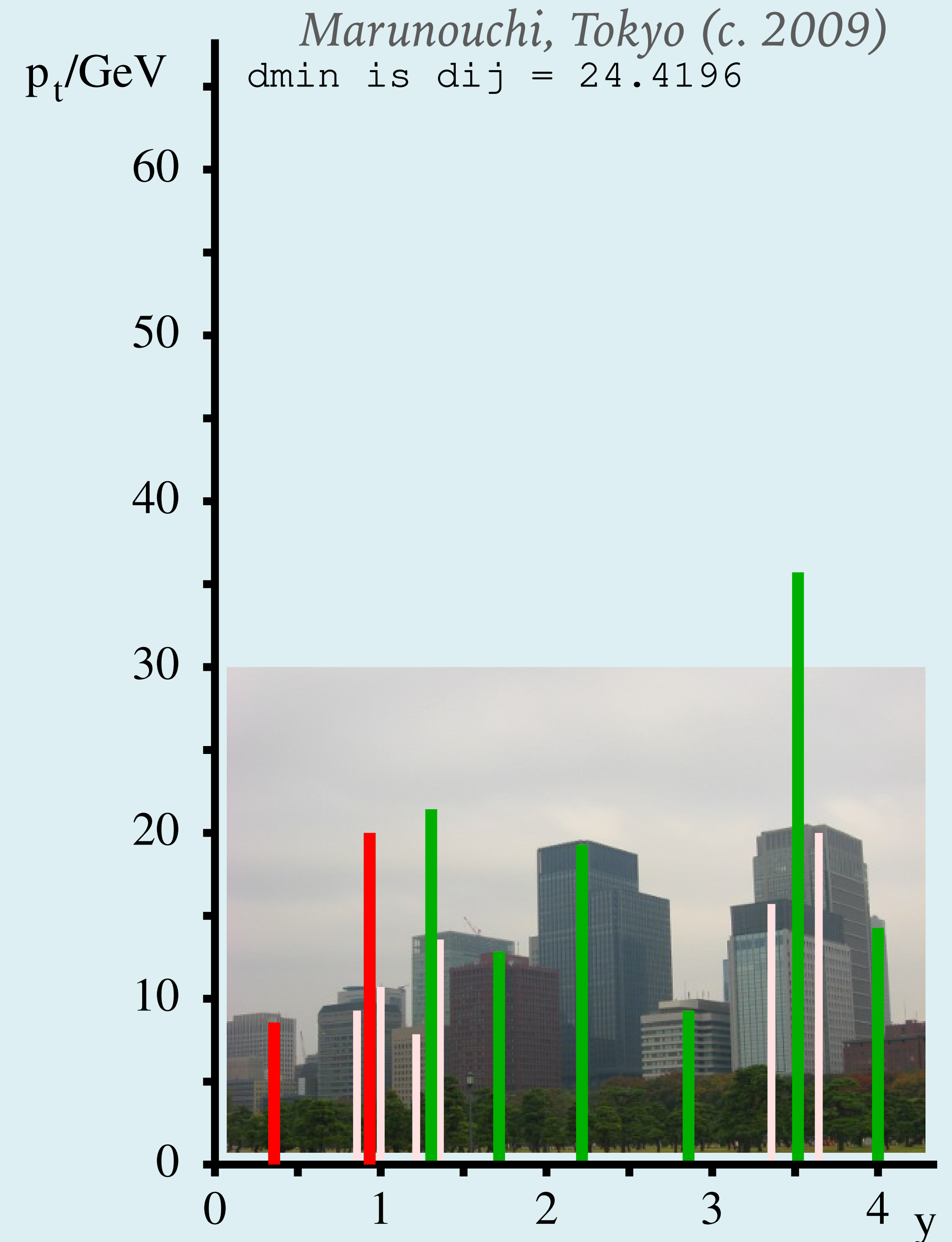
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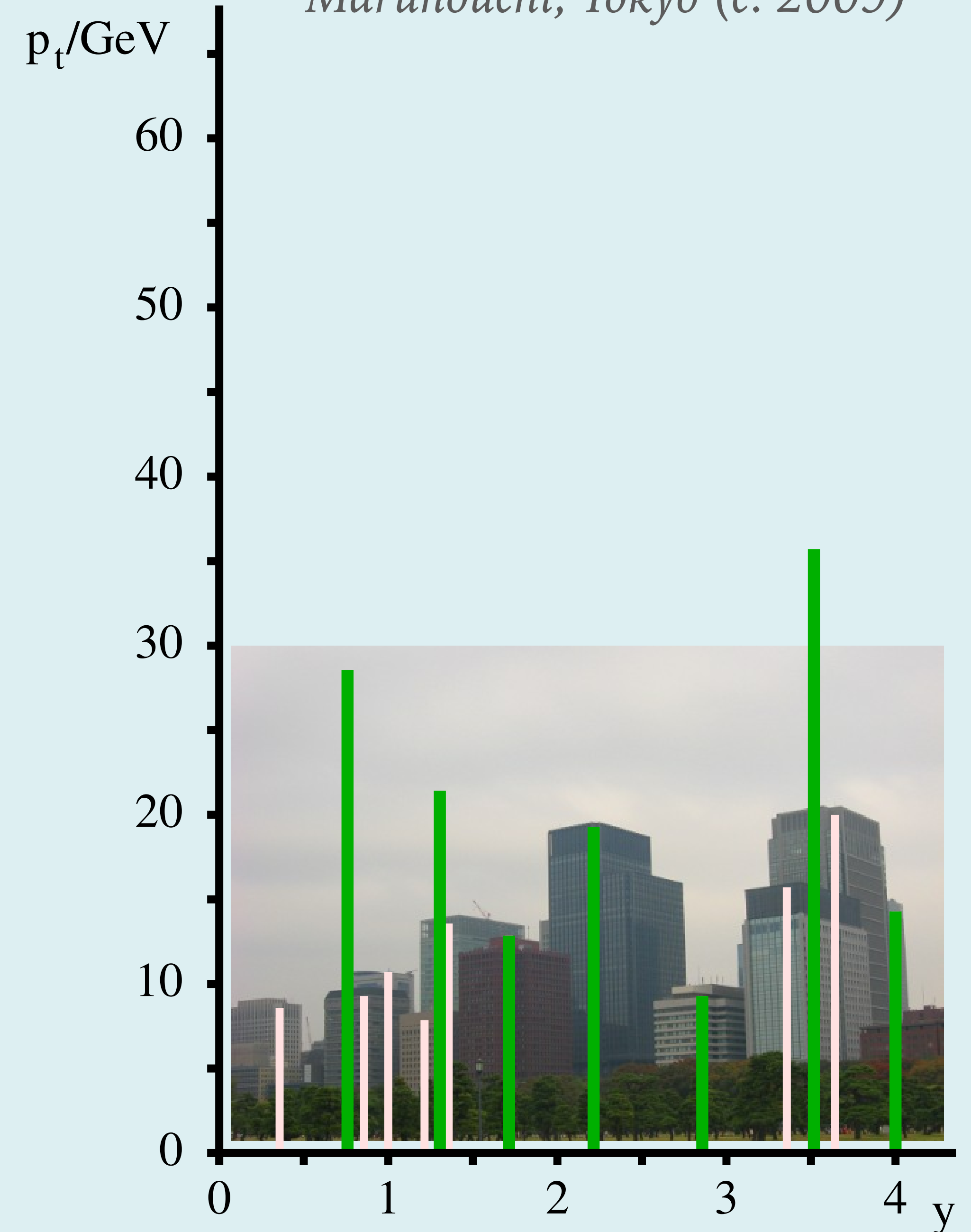
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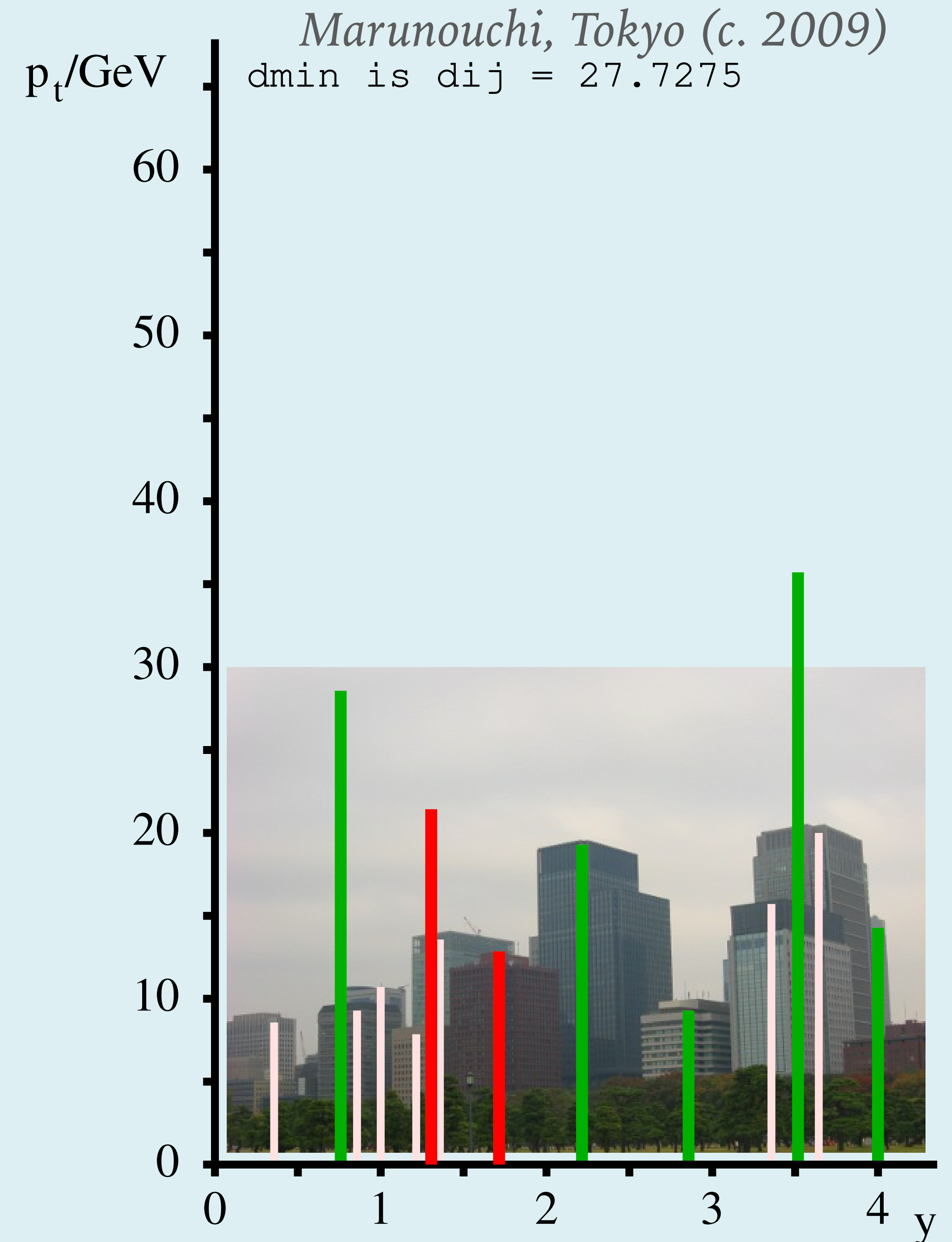
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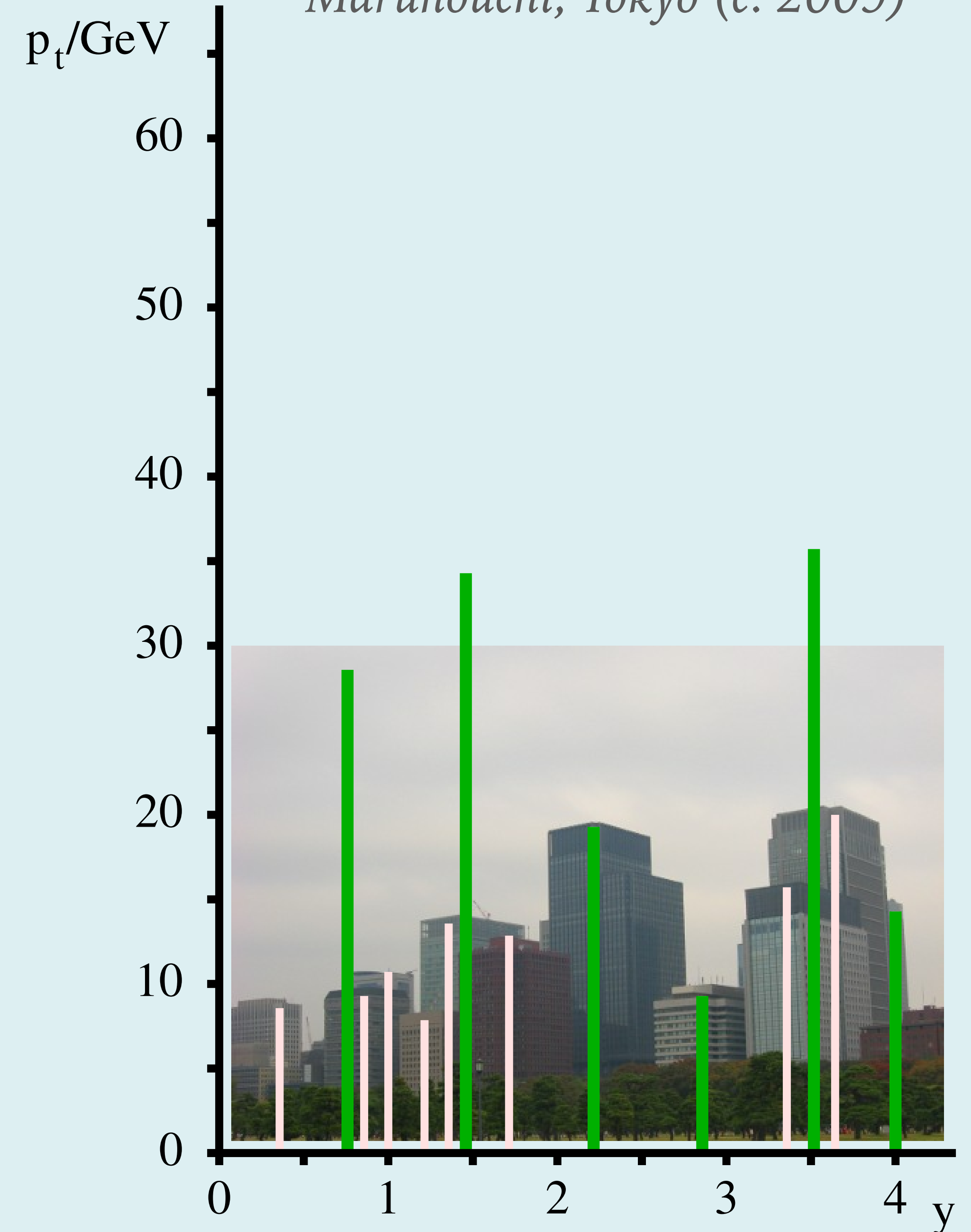
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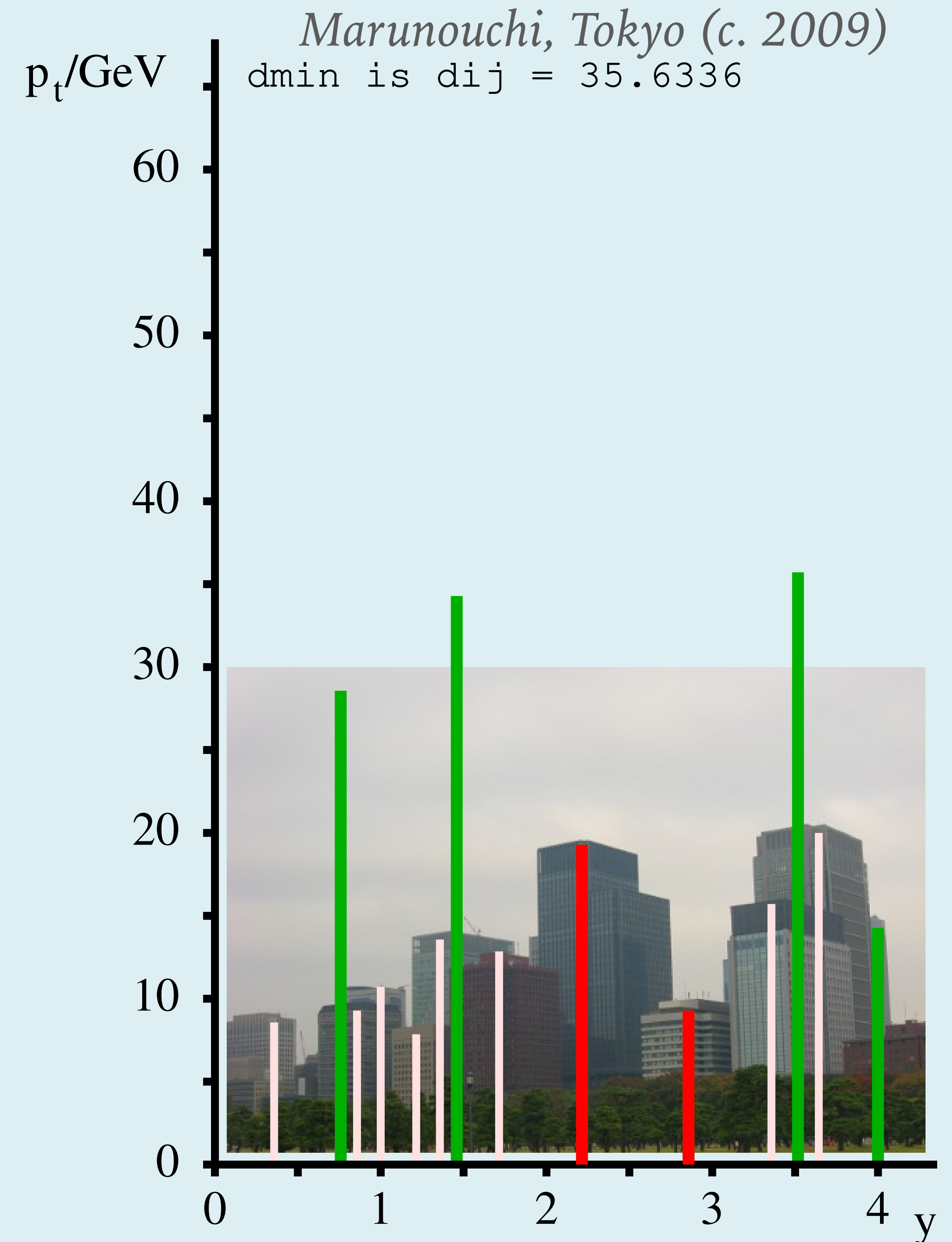
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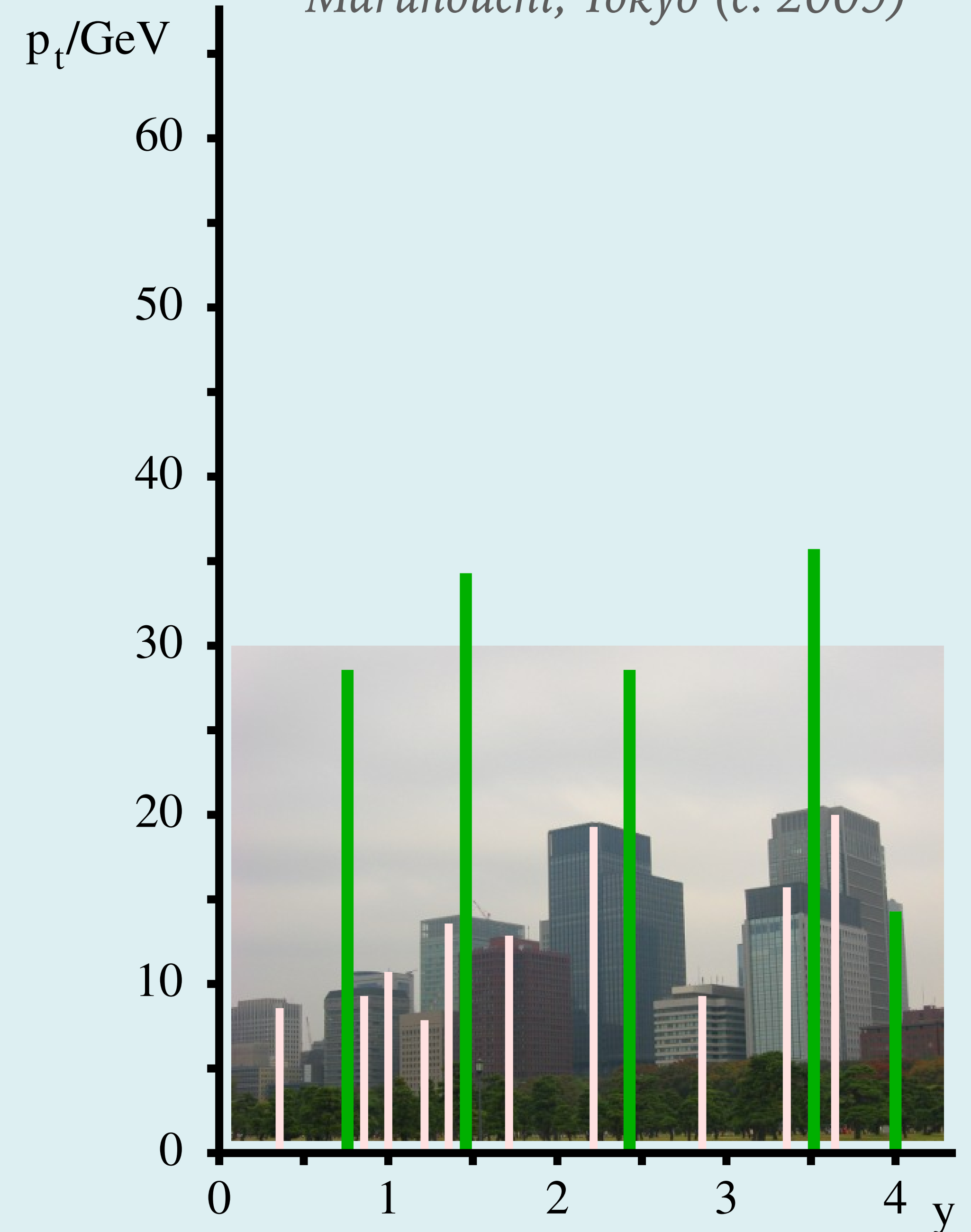
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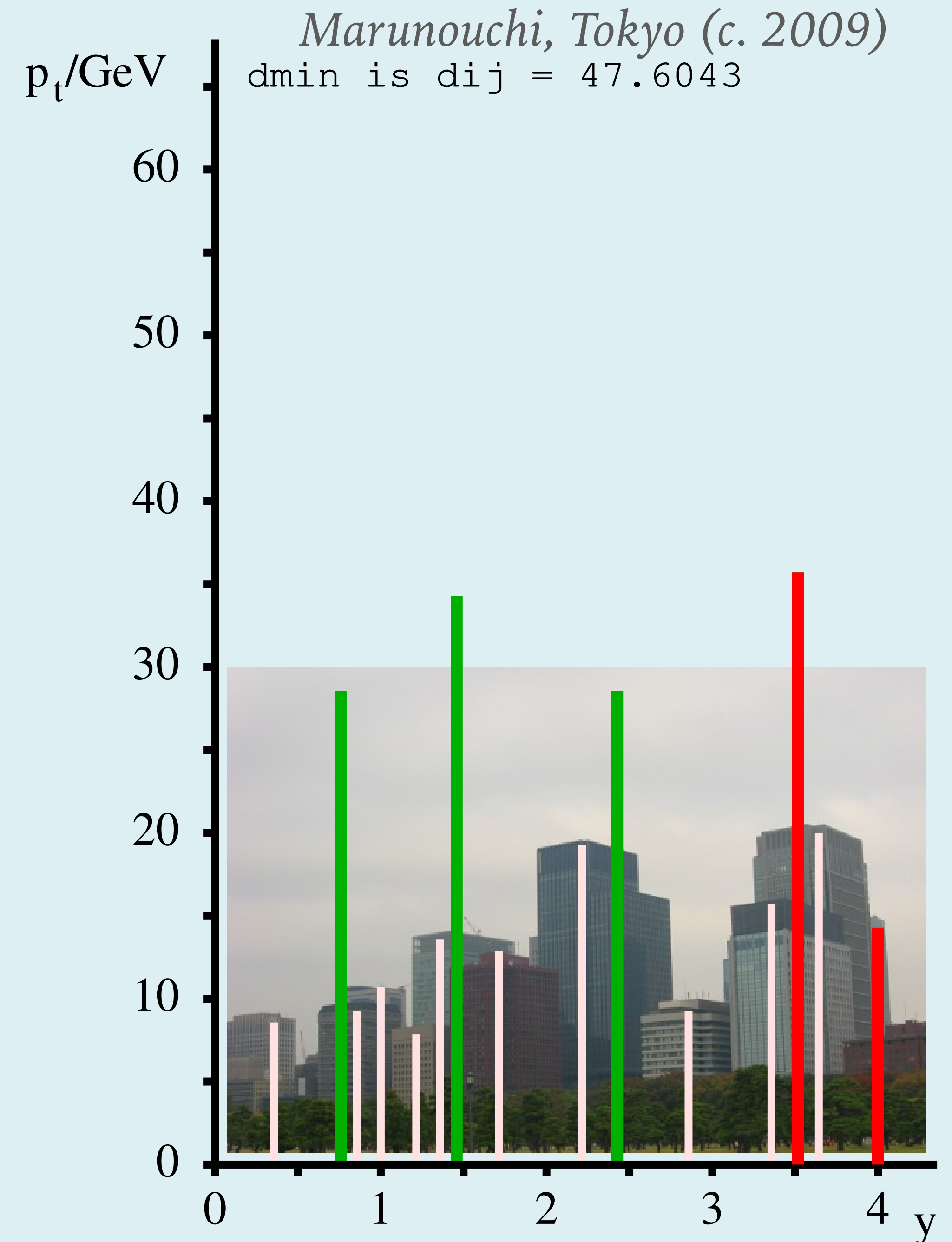
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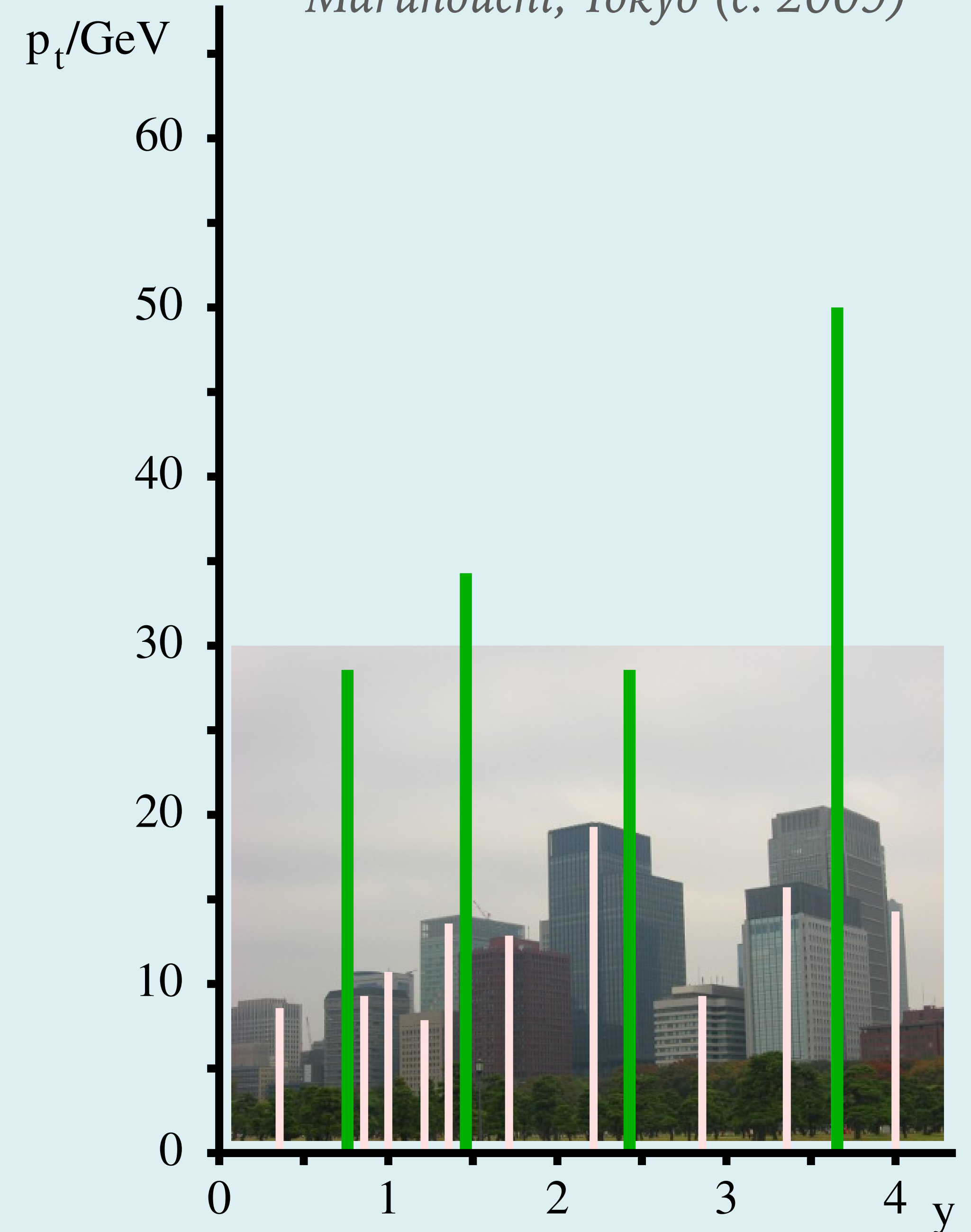
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2. if d_{ij} recombine particles into one
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Accept all jets above some some threshold
transverse momentum

Marunouchi, Tokyo (c. 2009)



possible solution

“sequential recombination” k_t algorithm

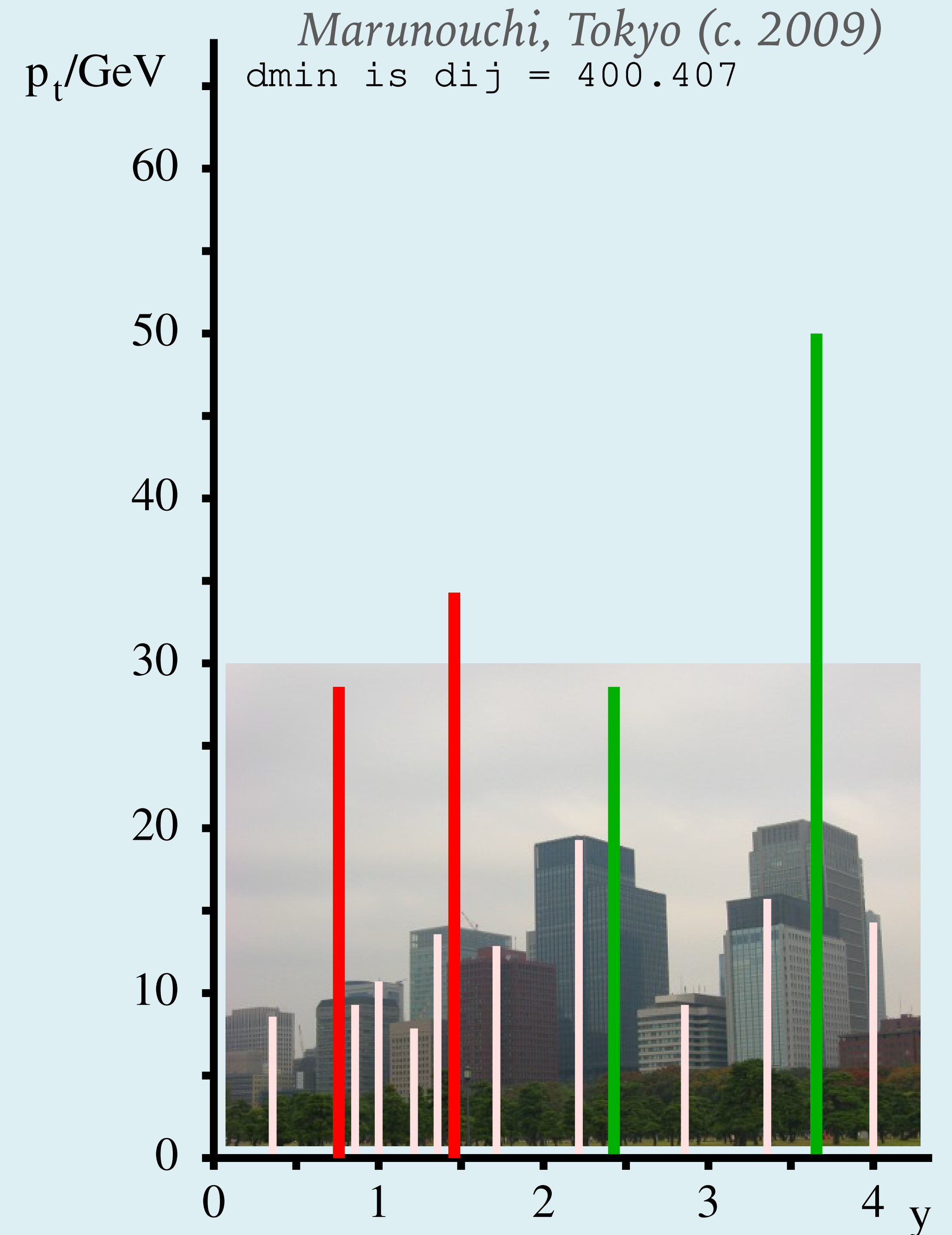
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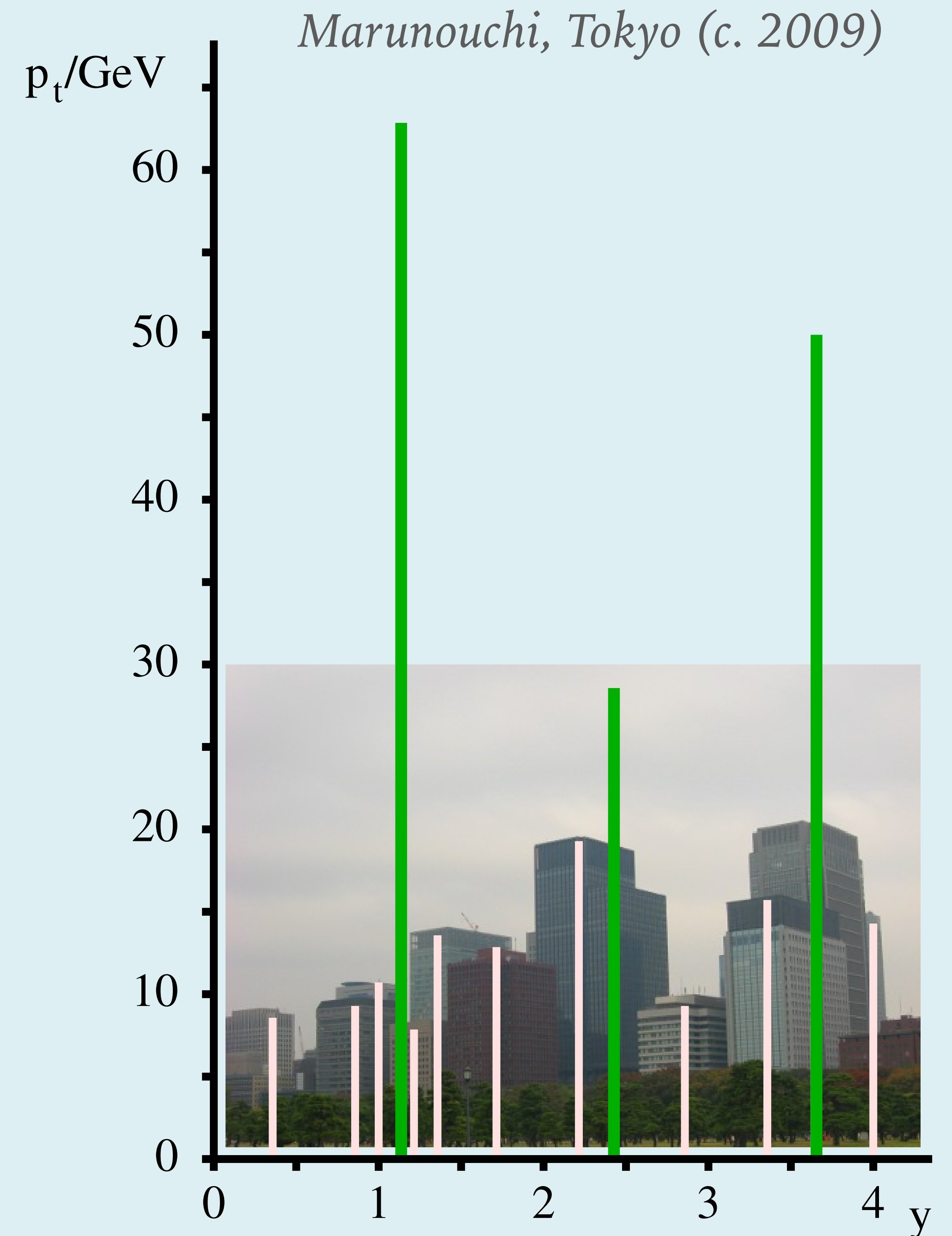
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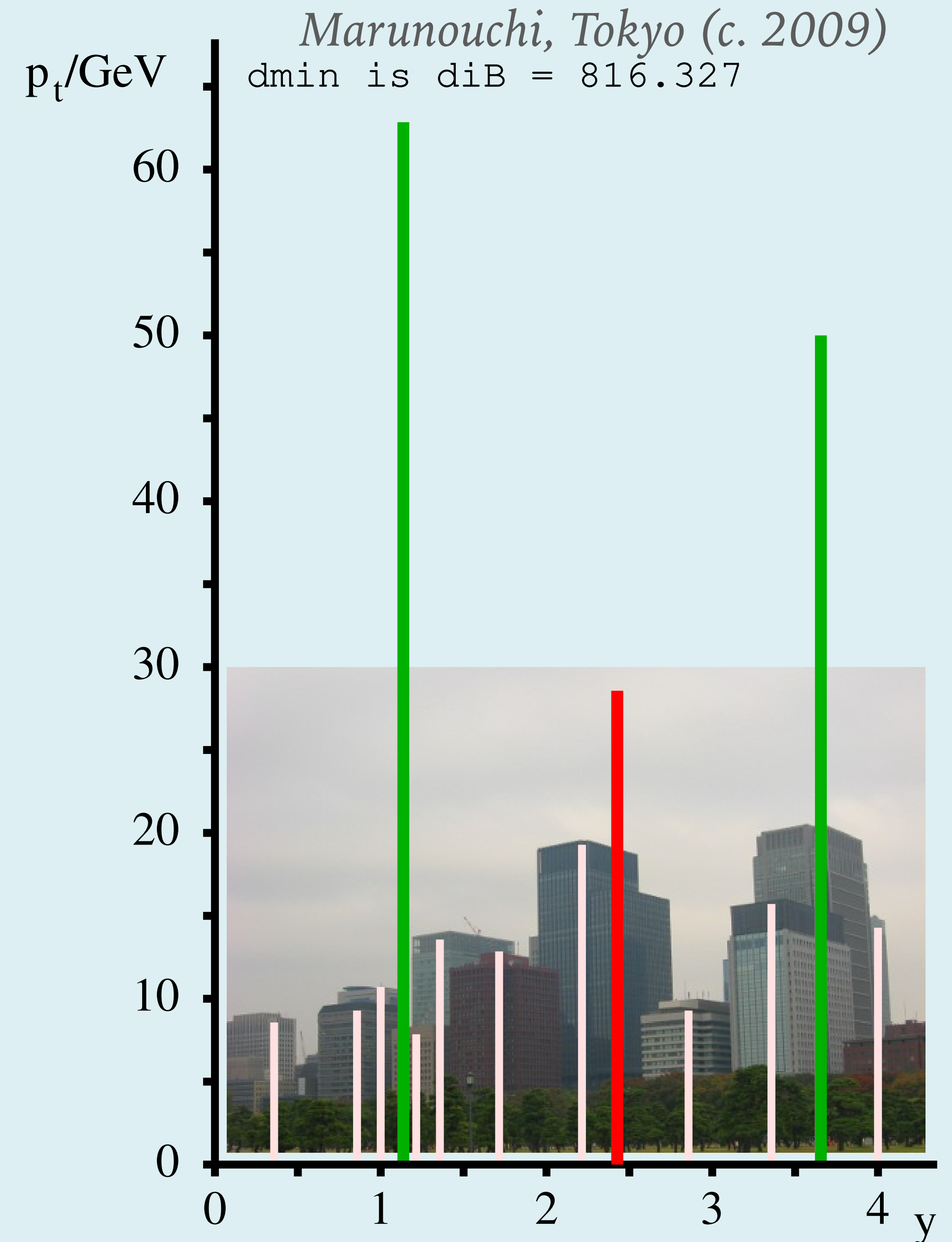
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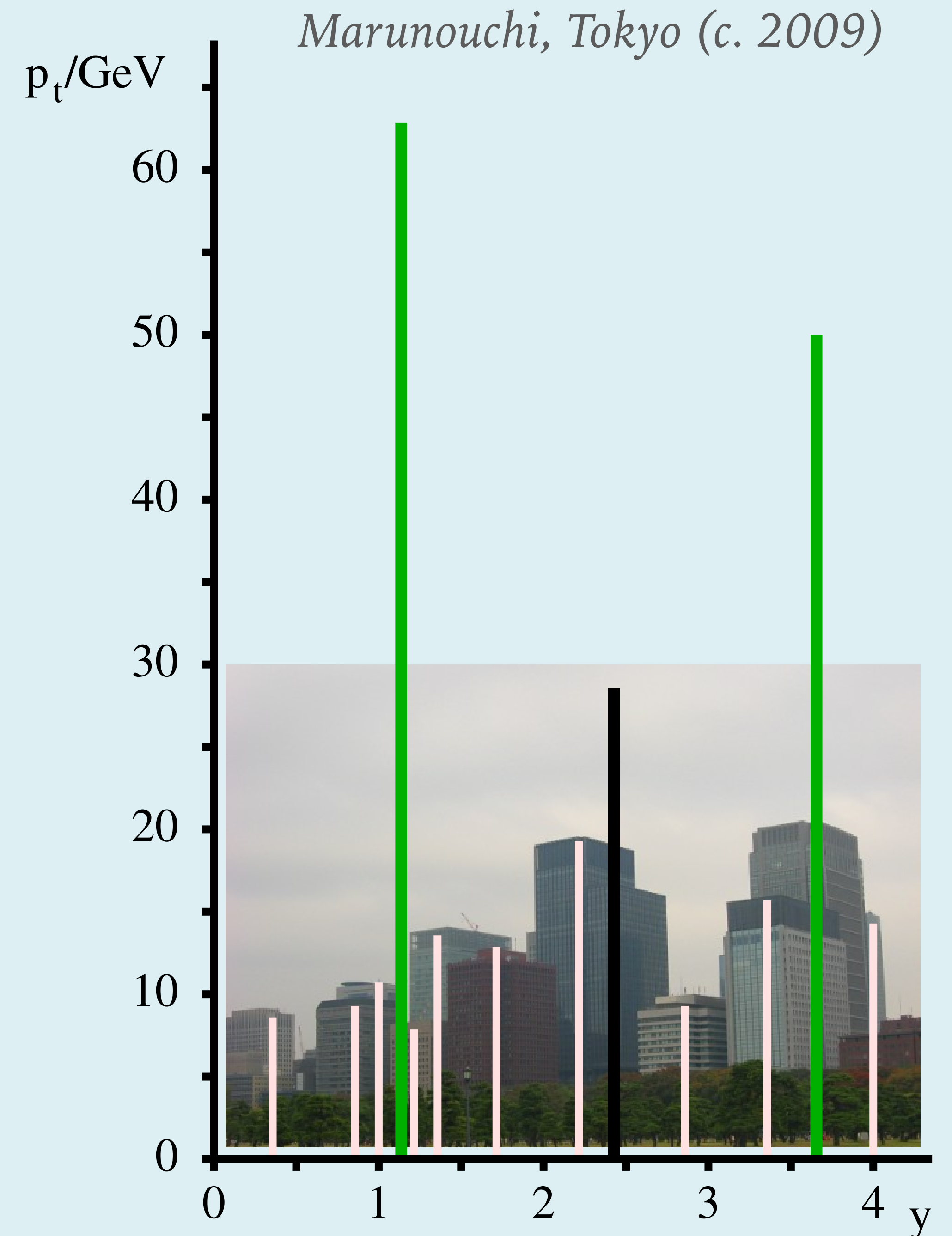
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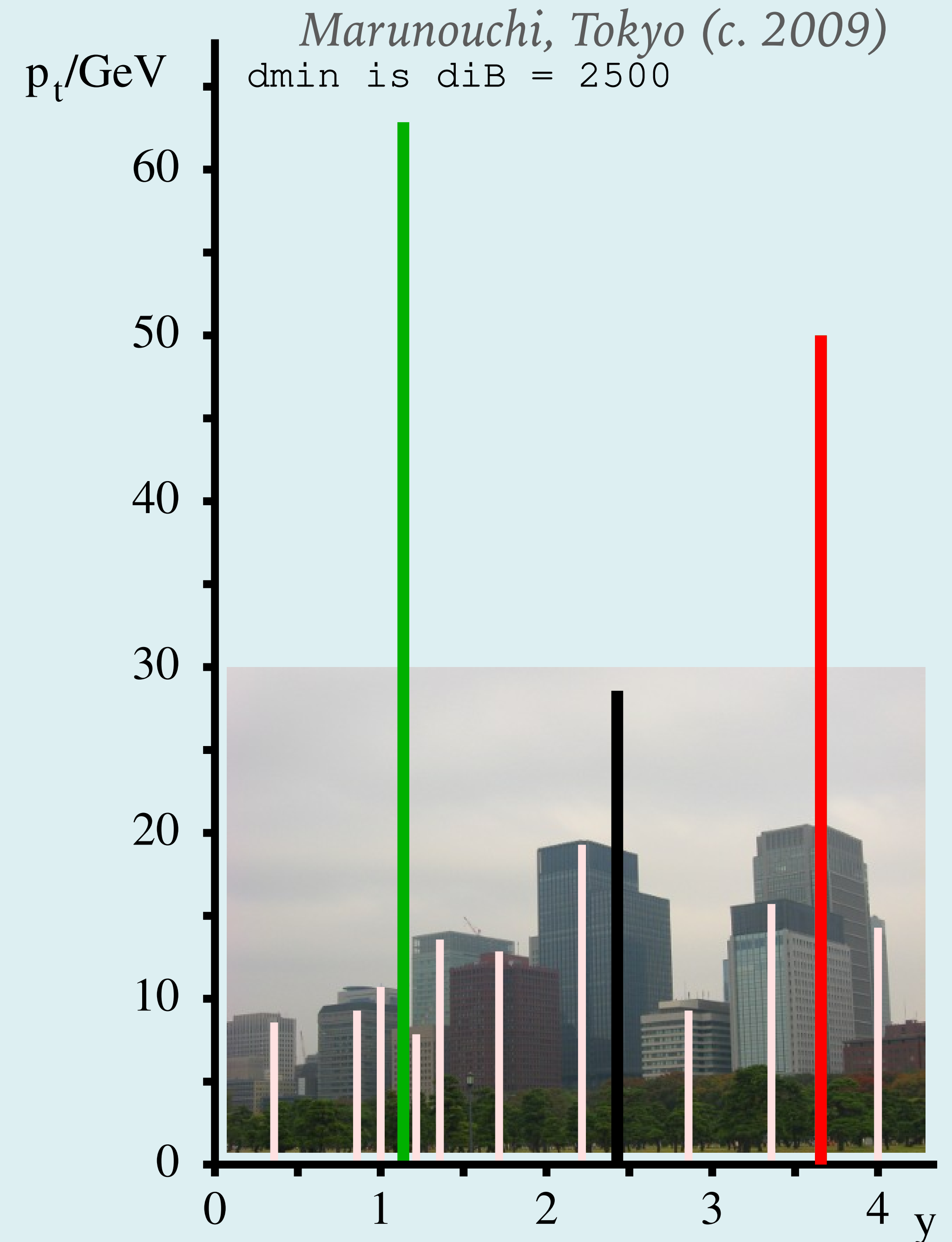
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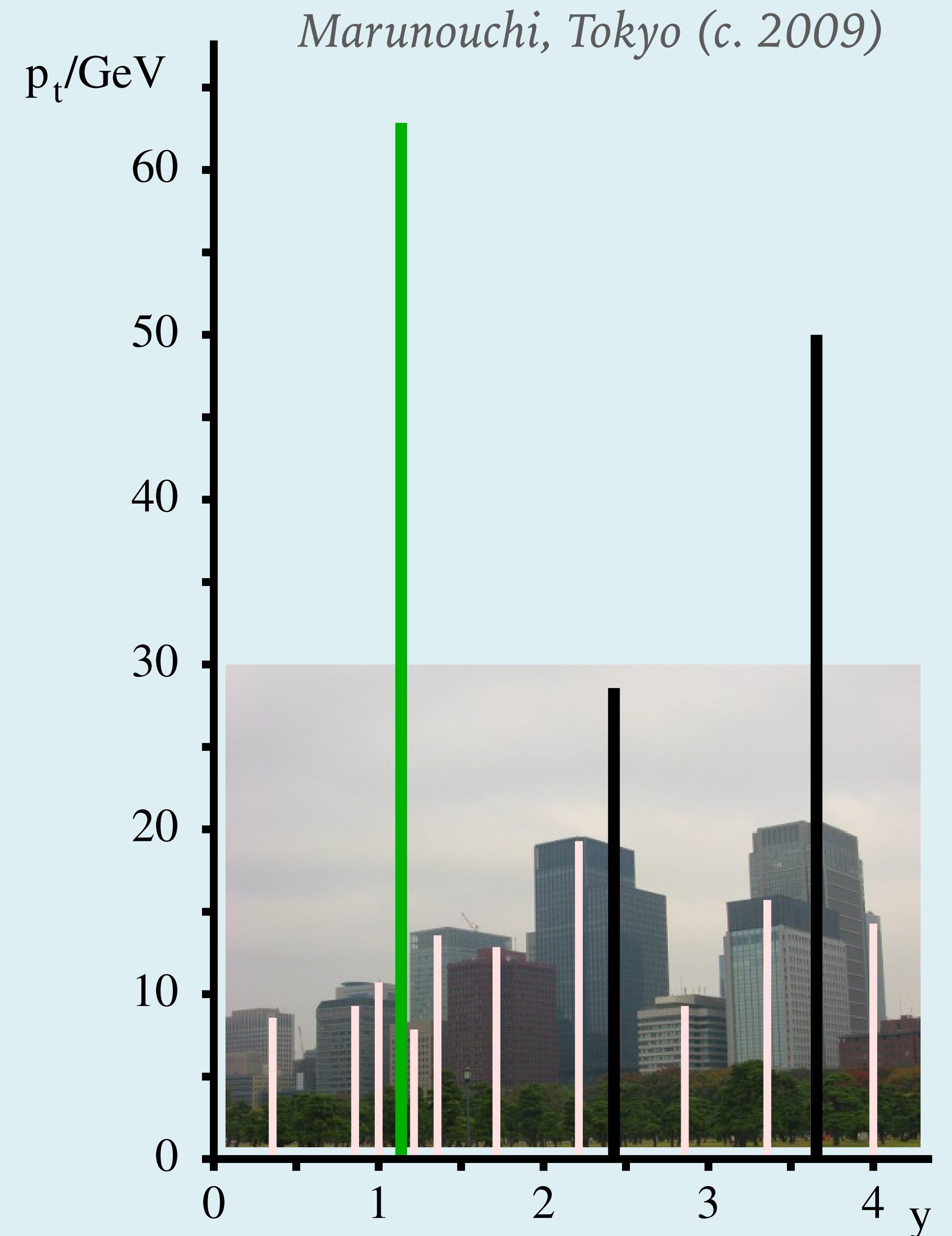
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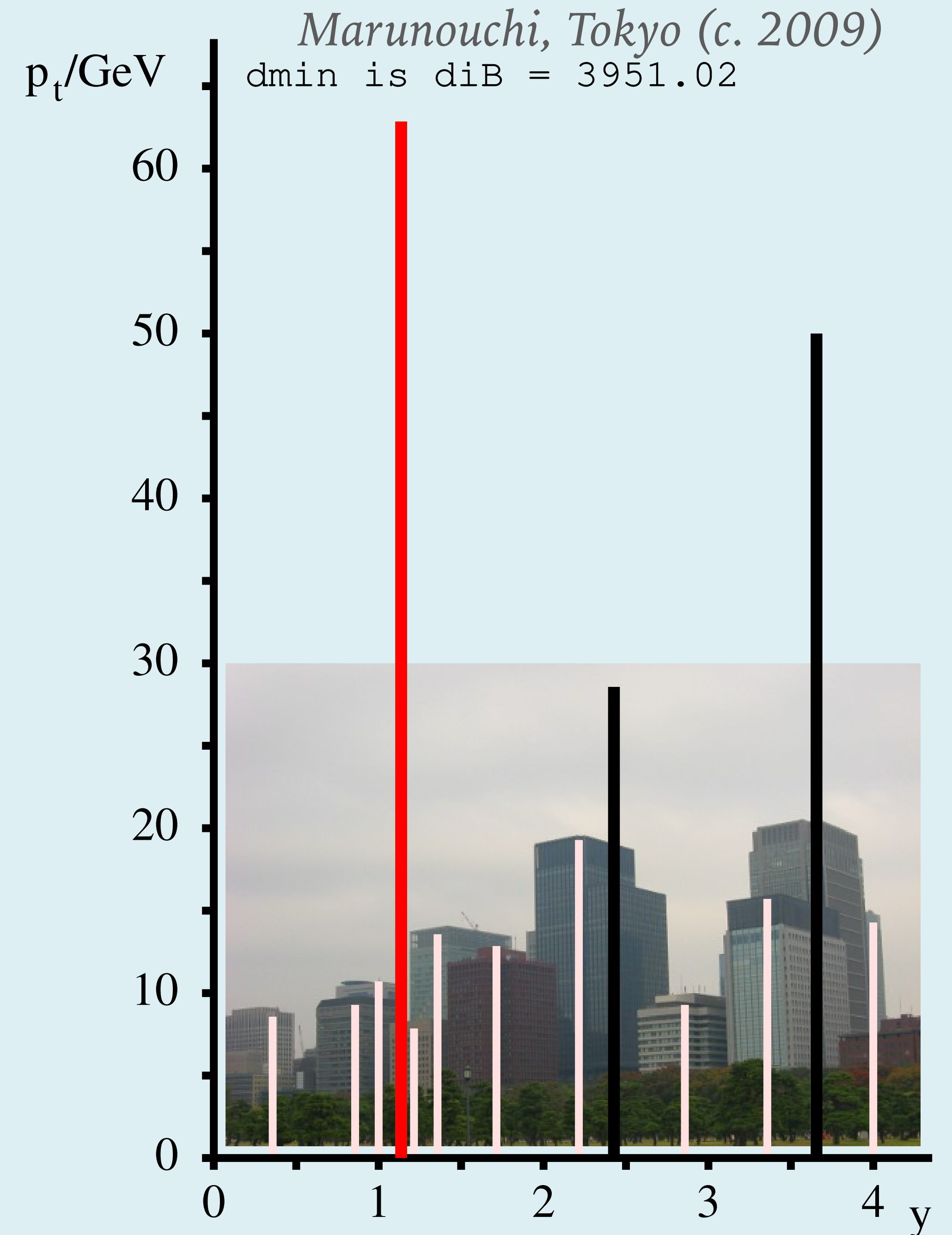
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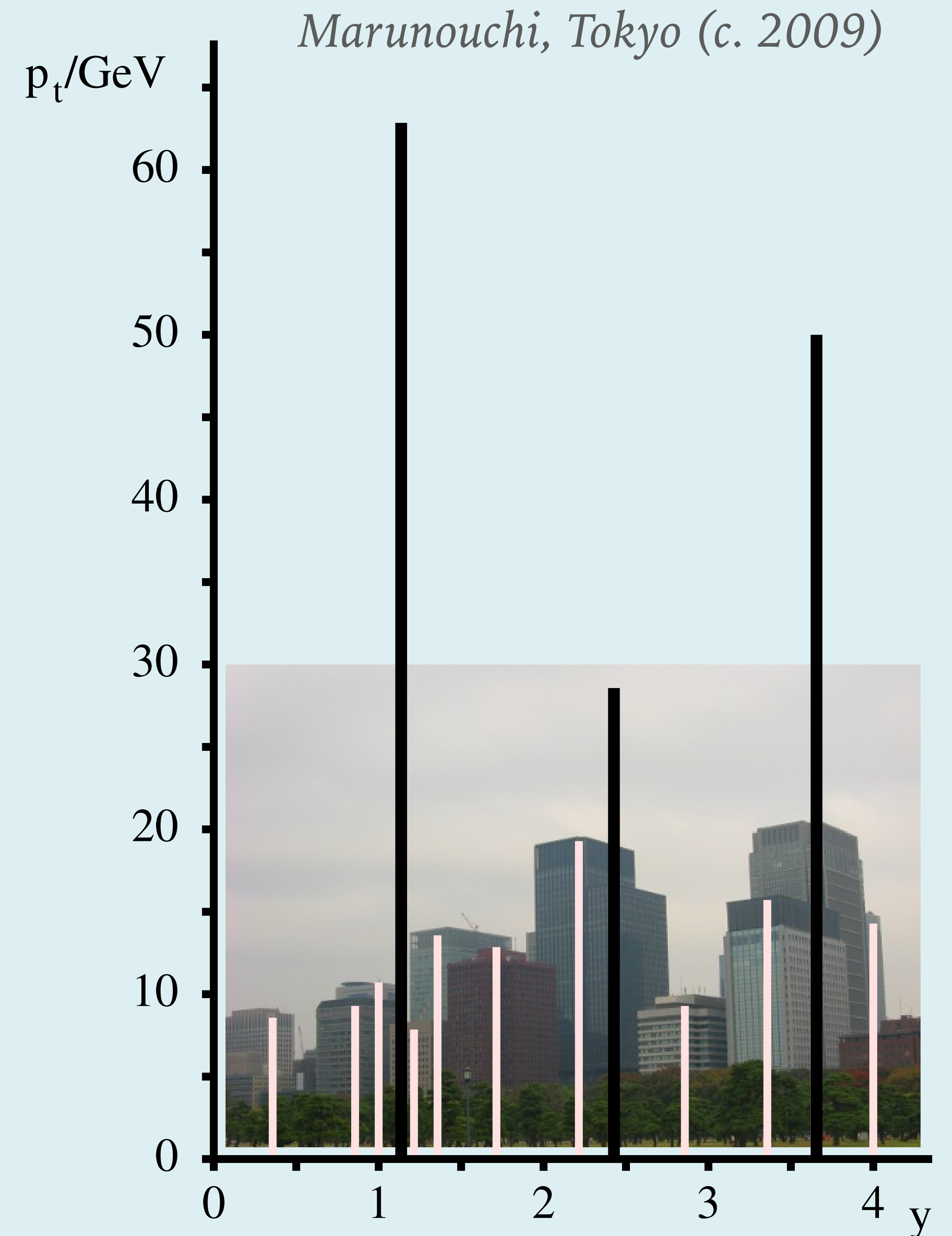
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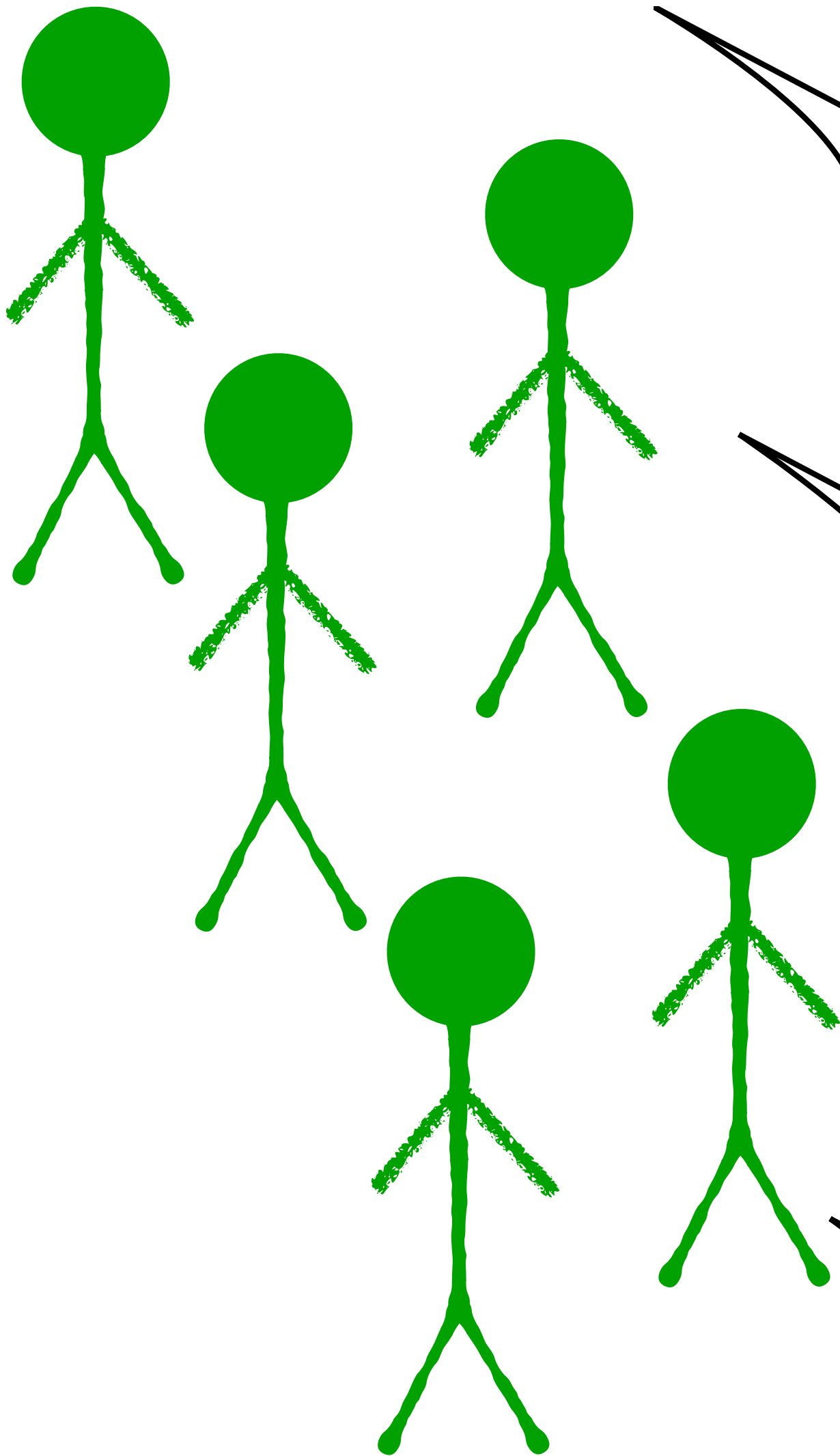
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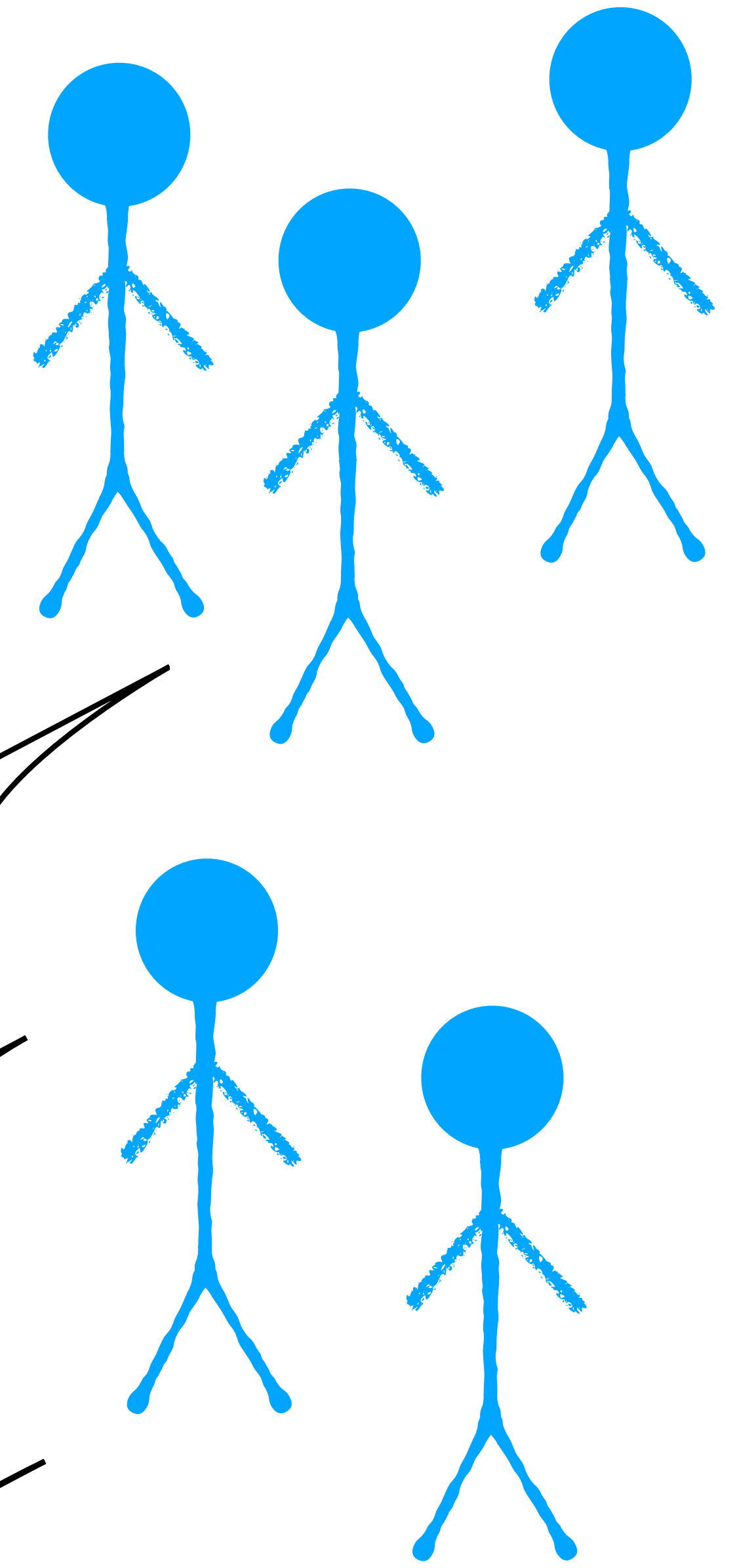


c. 2005



many theorists

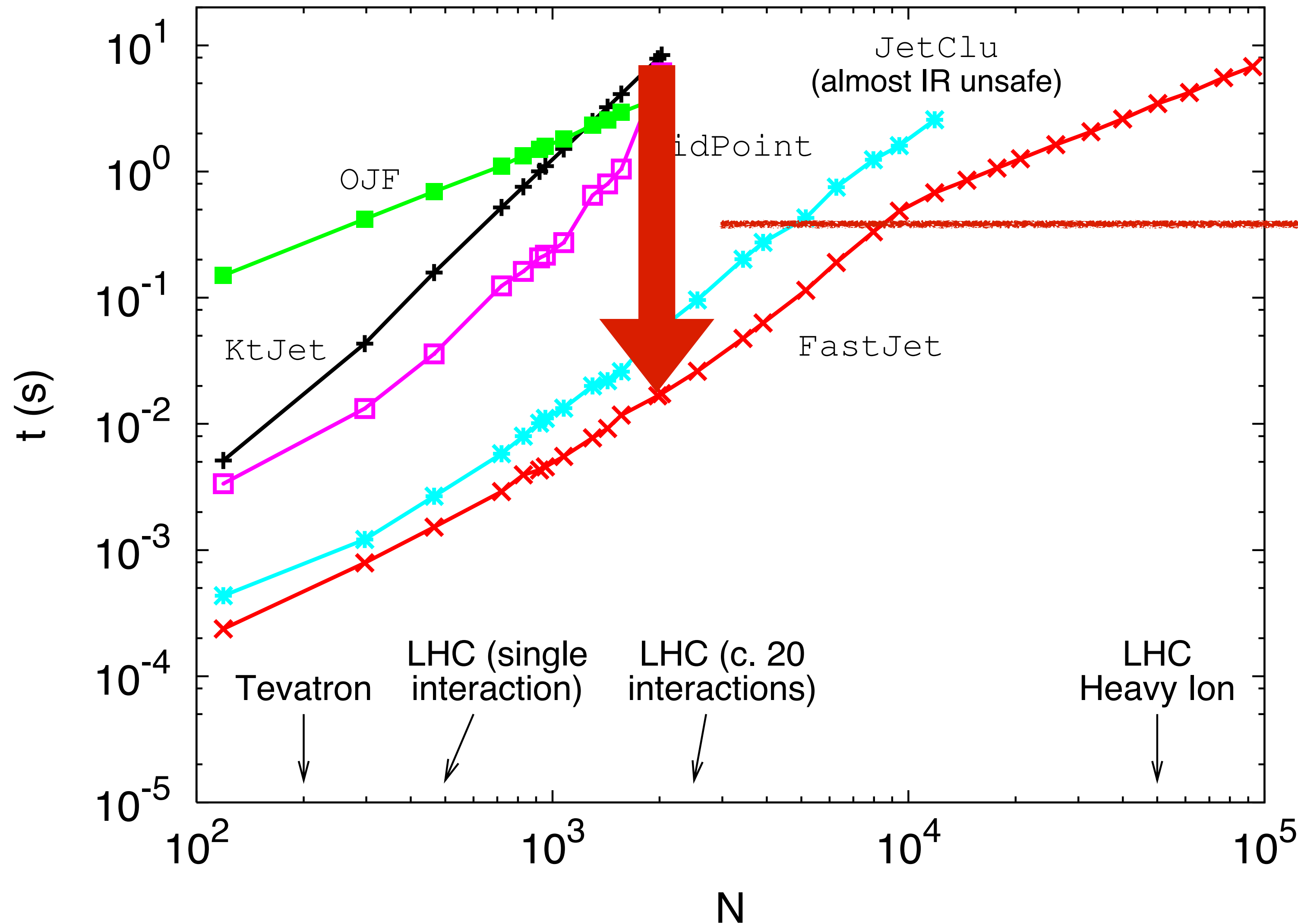
- kt adapts to the jet structure
- the cone is too rigid
- cone has big hadronisation corrections
- cone is infrared unsafe



many experimenters

- the cone gives nice conical jets
- kt's a vacuum cleaner
- kt's too irregular I can't correct for pileup
- kt's too slow

advance #1: computational speed for IR safe “sequential recombination” jet algorithms

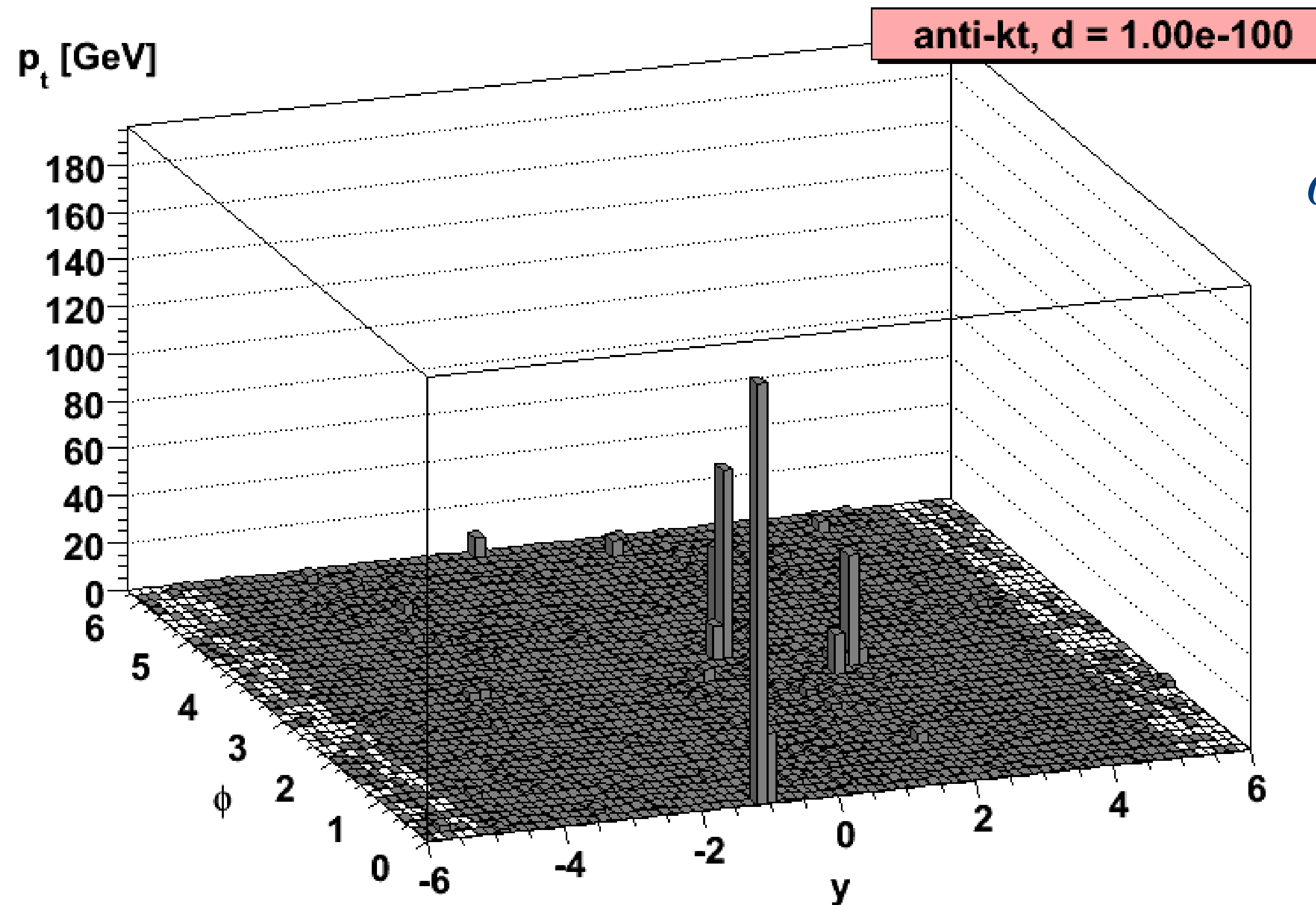


factor of ~ 1000 speedup, using the “FastJet lemma”

Exploits underlying geometric information to speed clustering from N^3 up to to $N \ln N$

Cacciari & GPS
hep-ph/0512210

advance #2: change clustering distances ($p_t \rightarrow 1/p_t$)

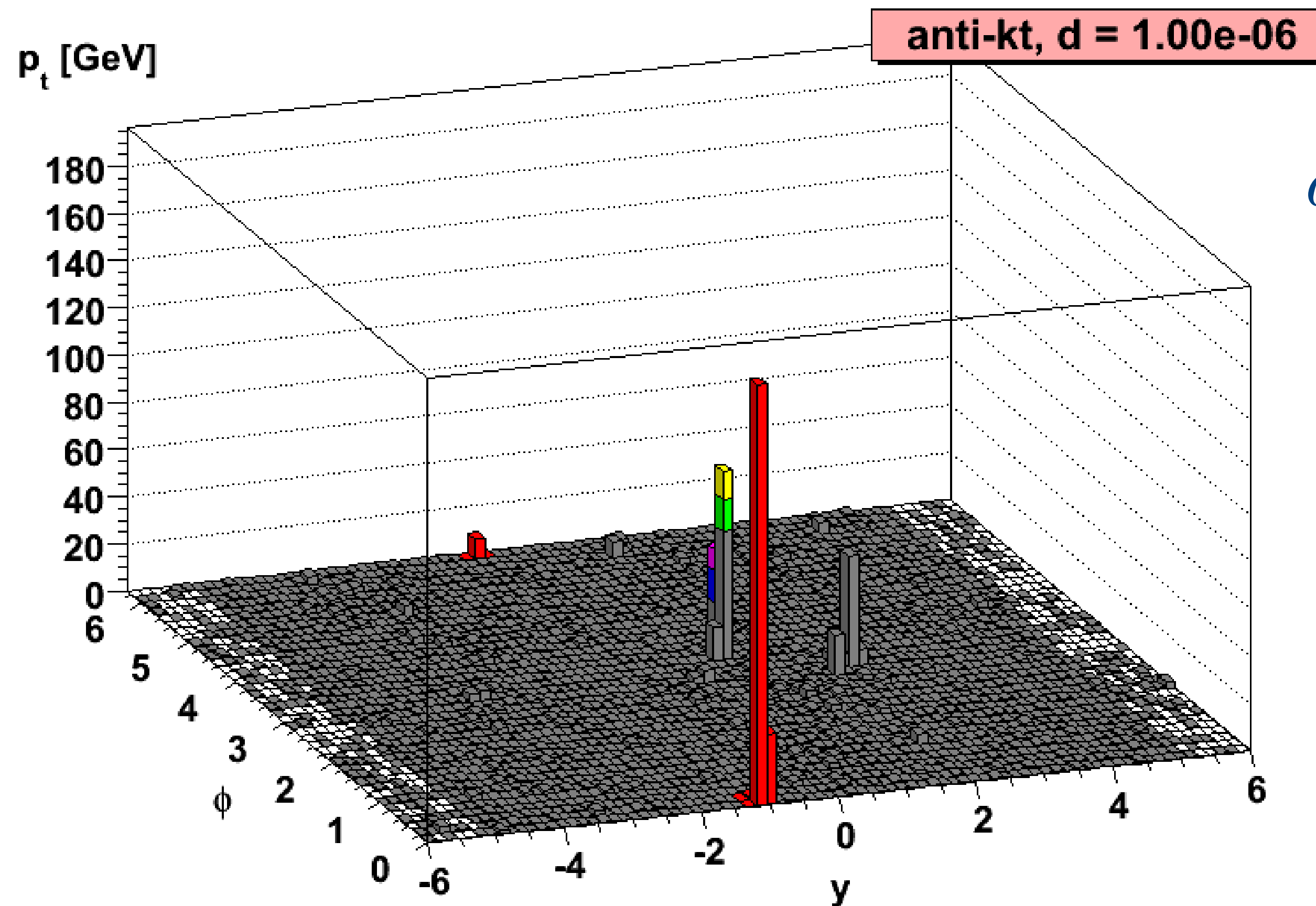


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Clustering grows
around hard cores
Gives **circular jets**

Cacciari, GPS & Soyez
0802.1189

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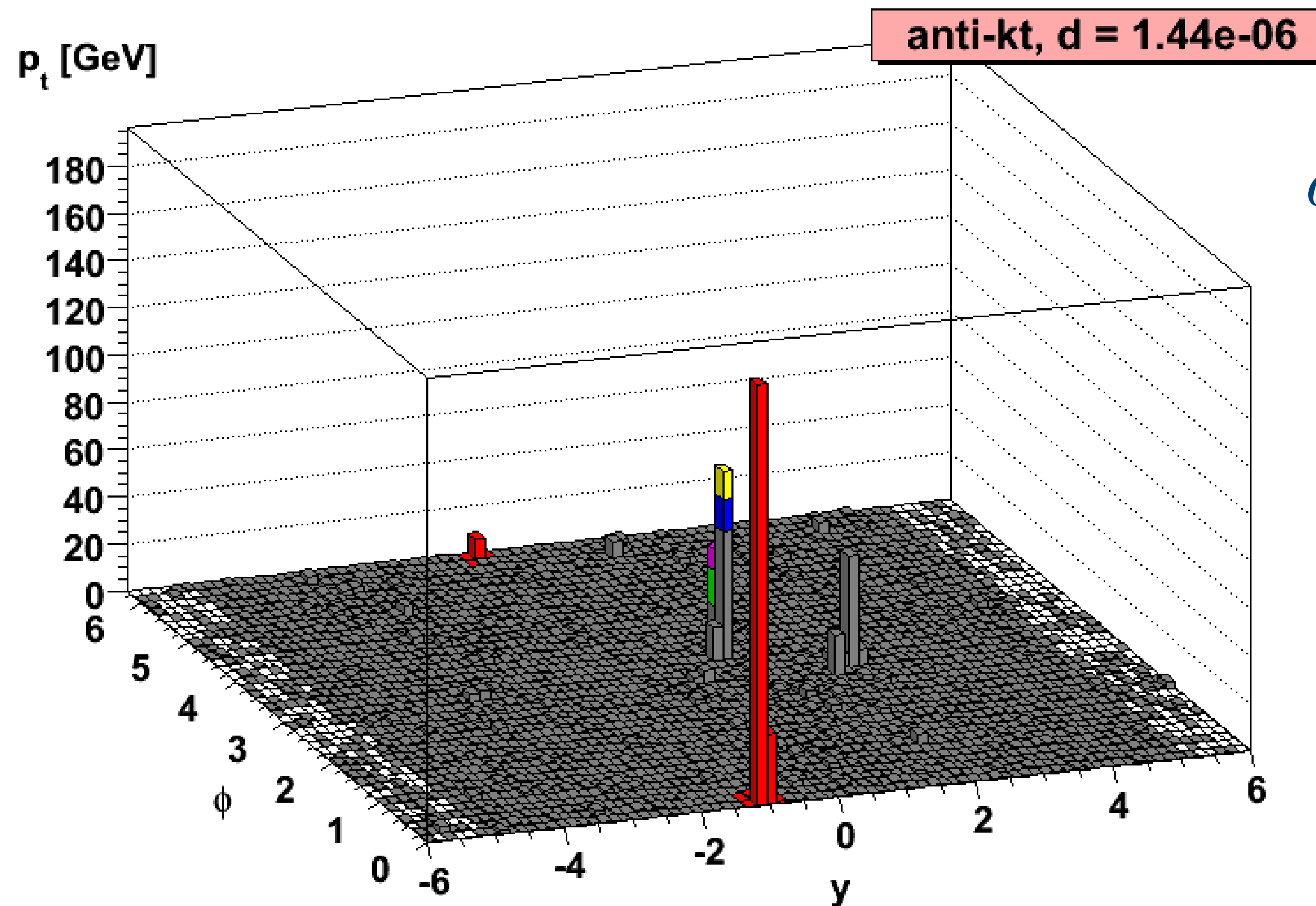


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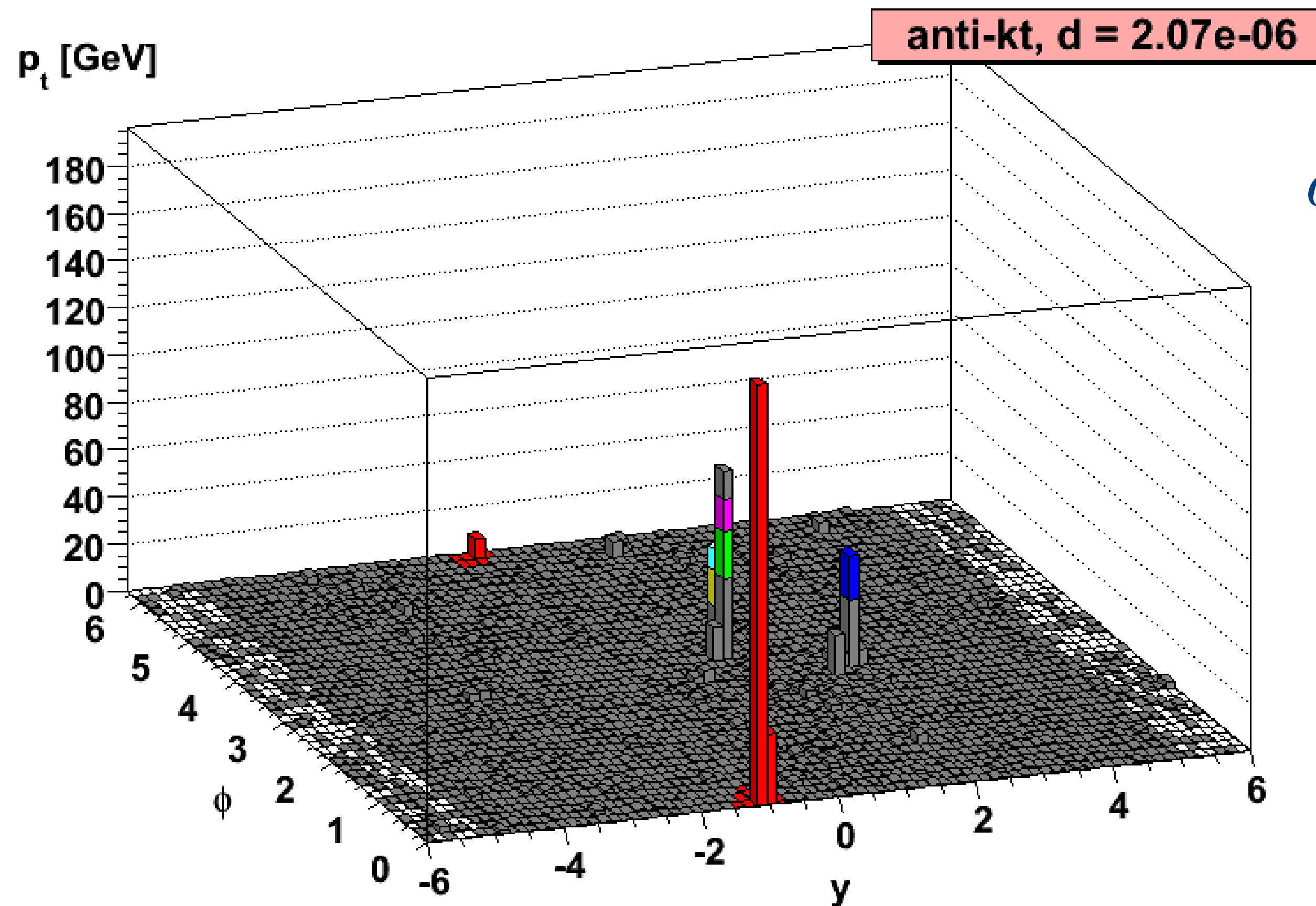
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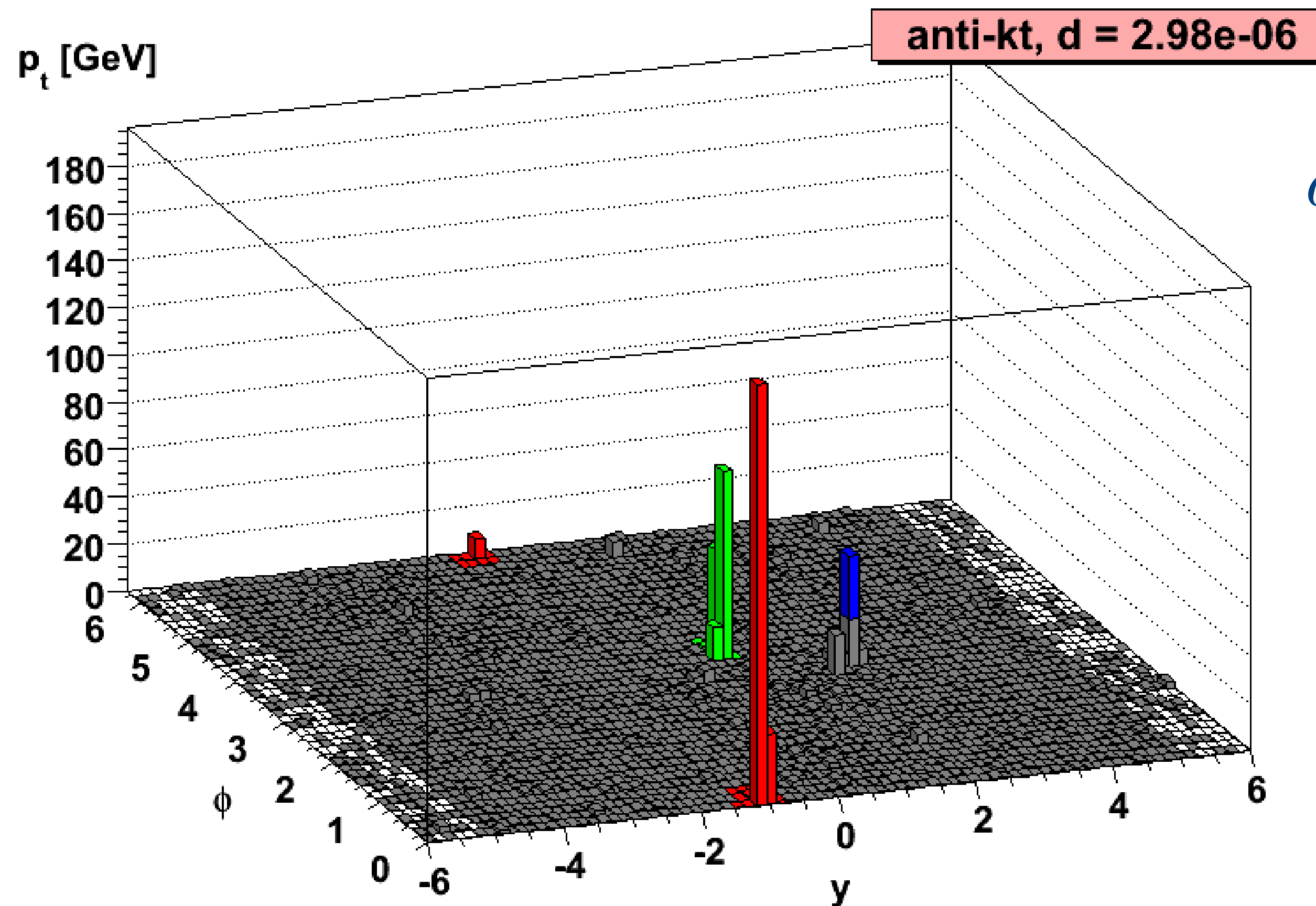


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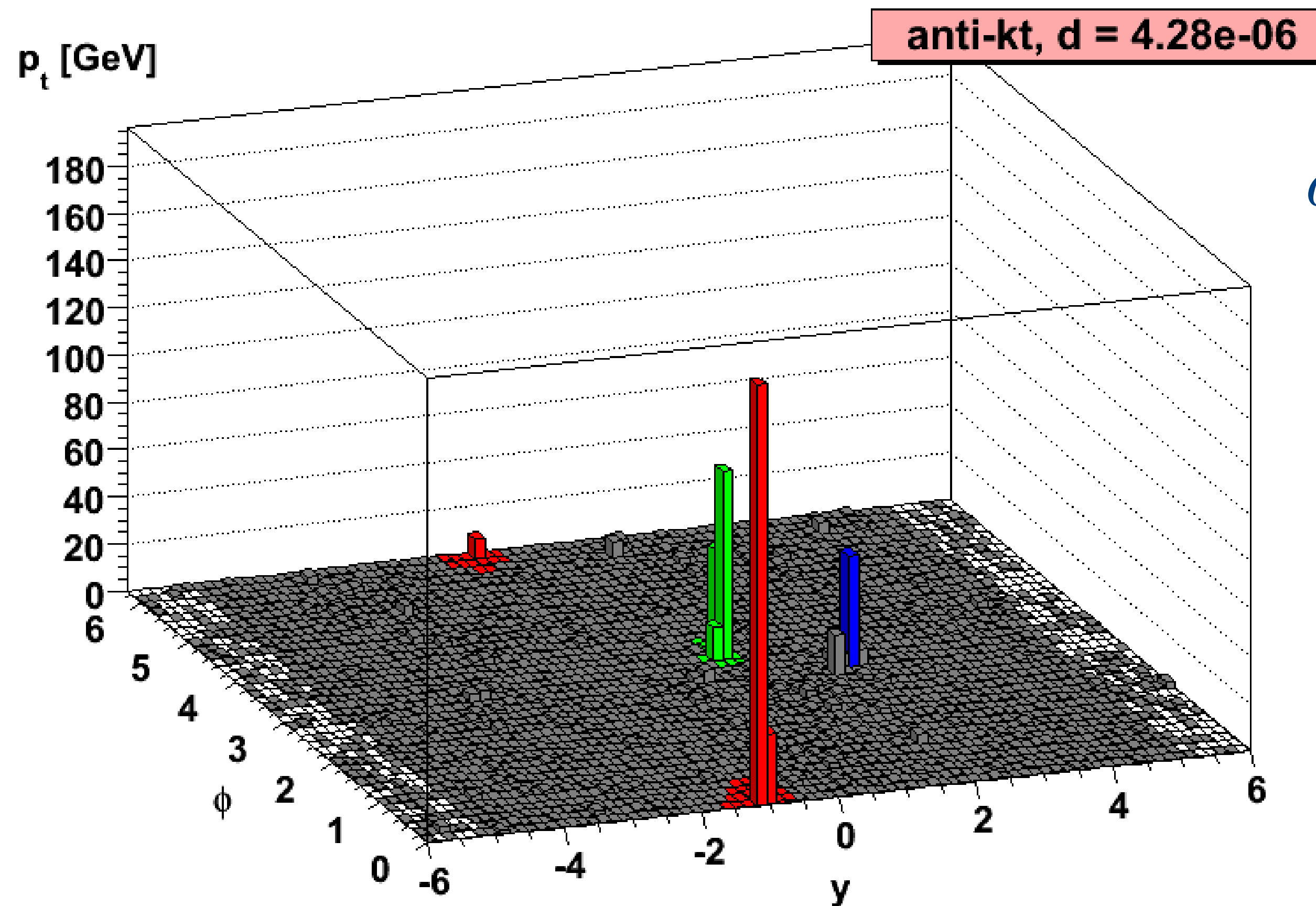
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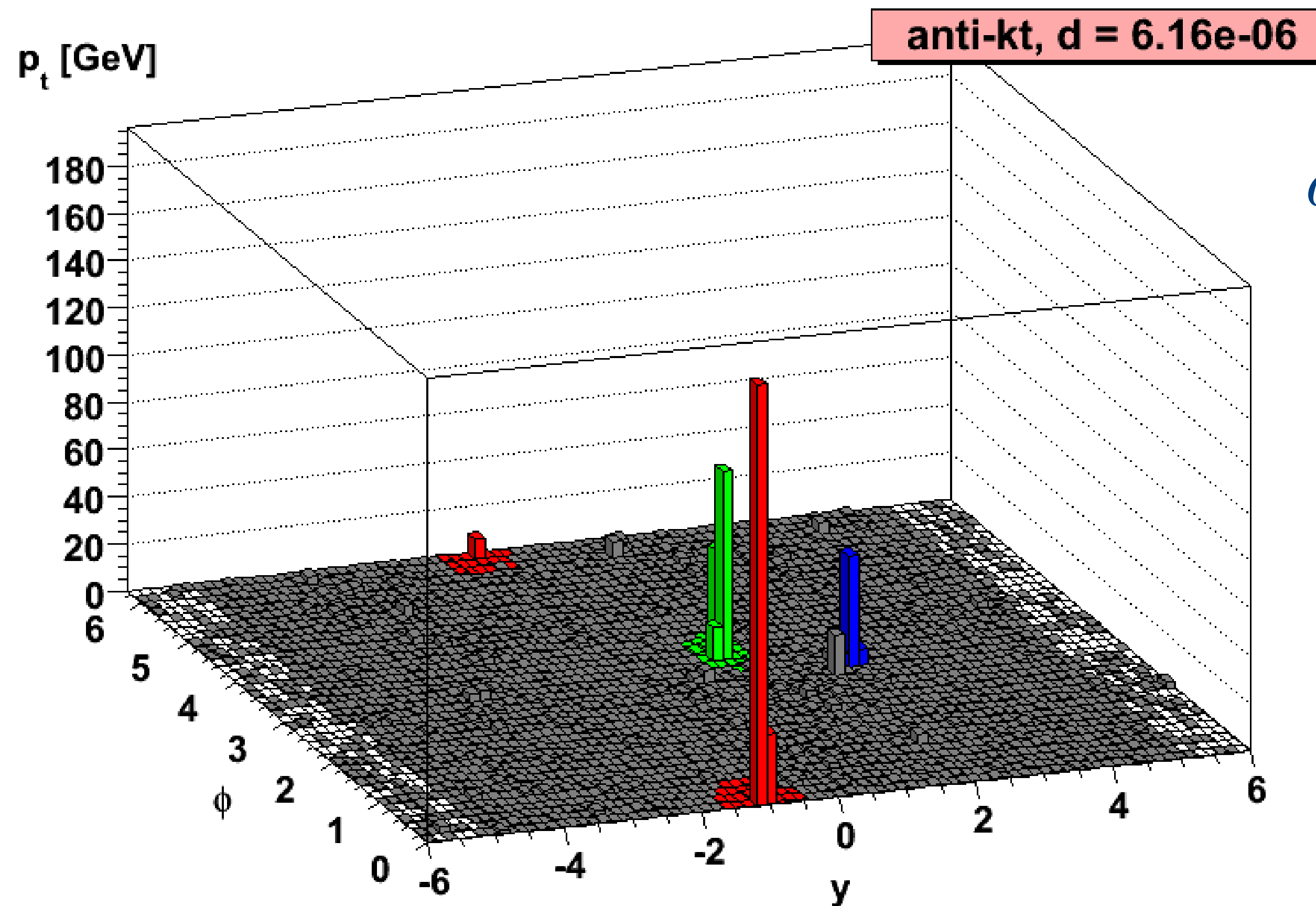
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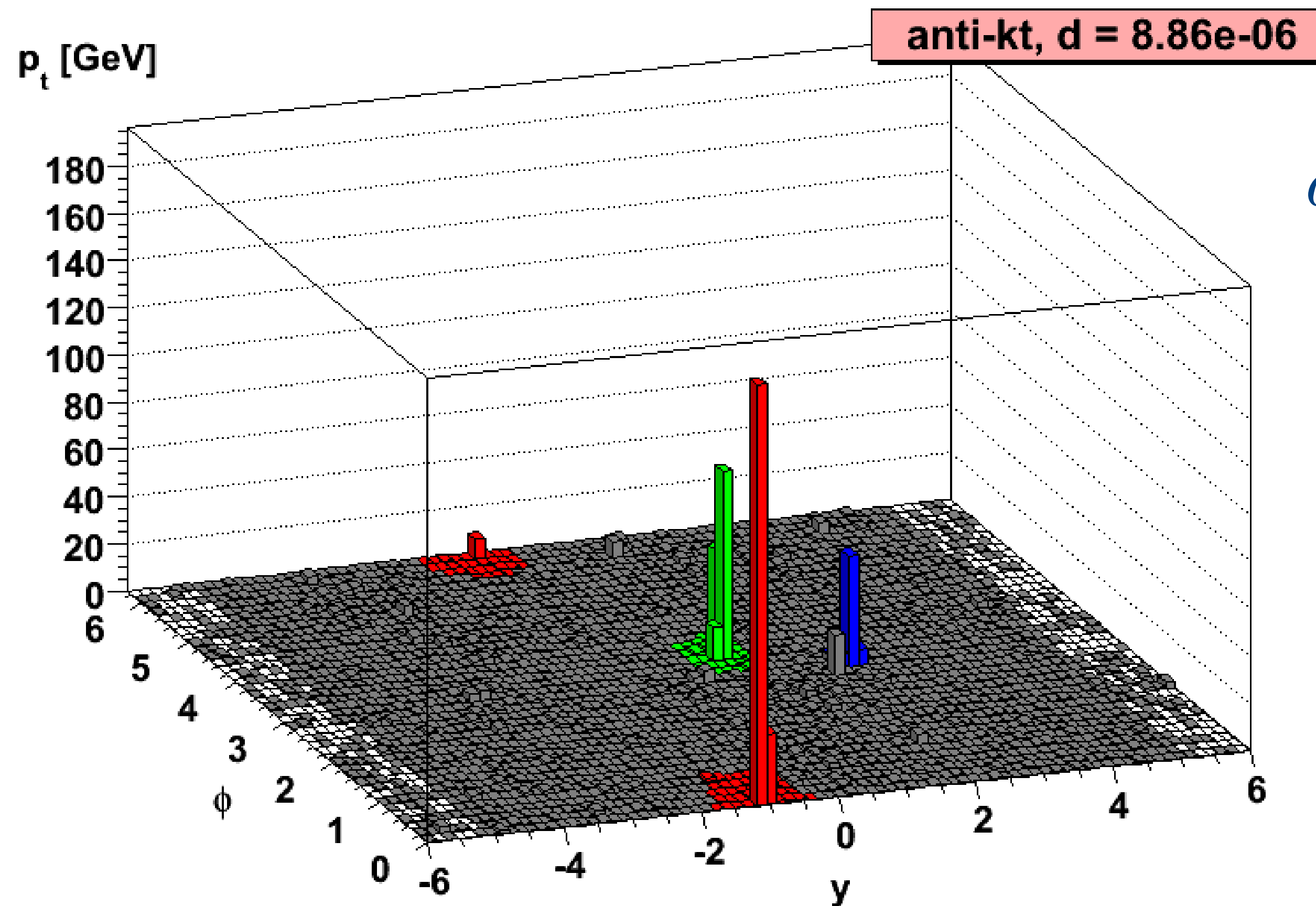
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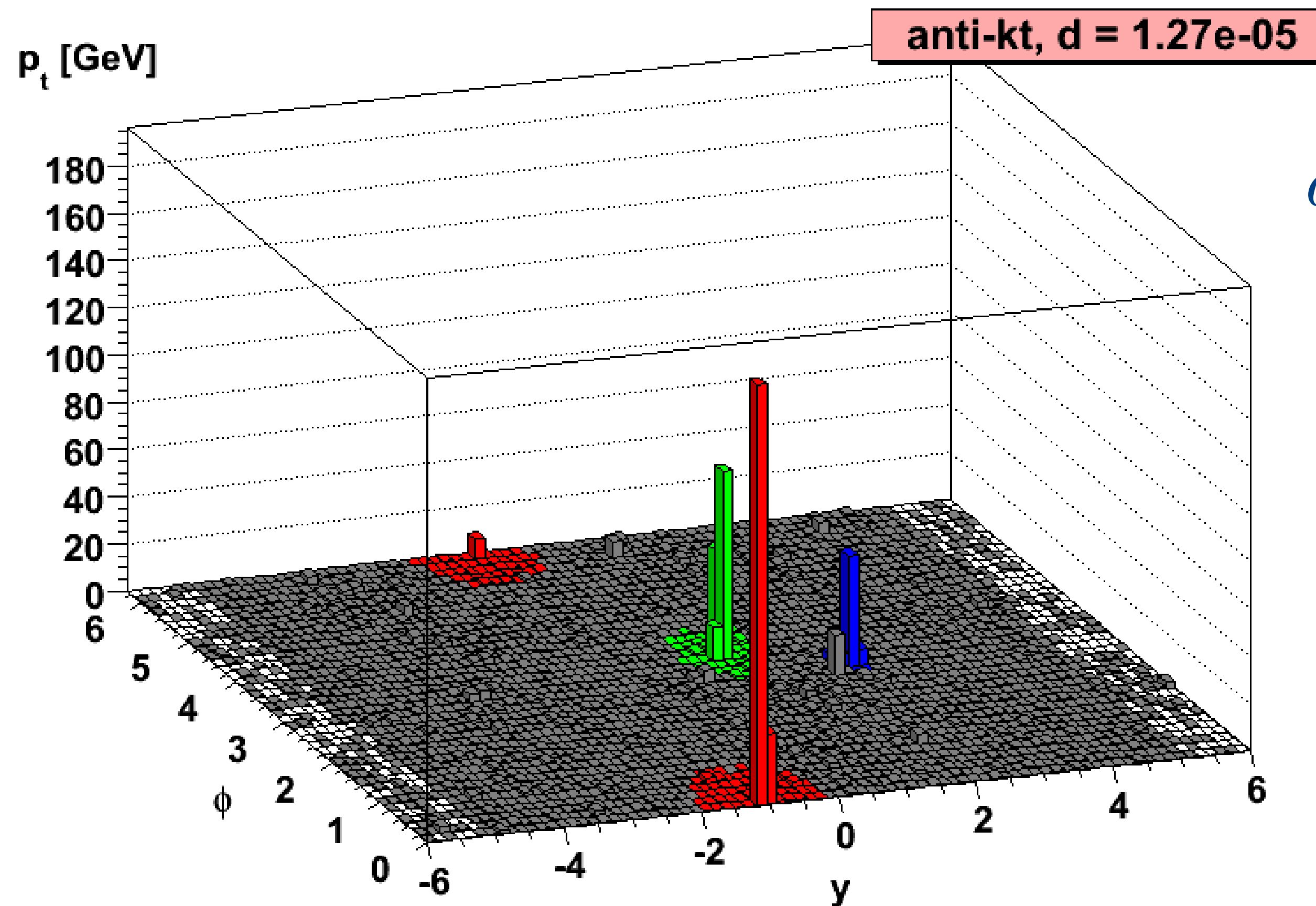


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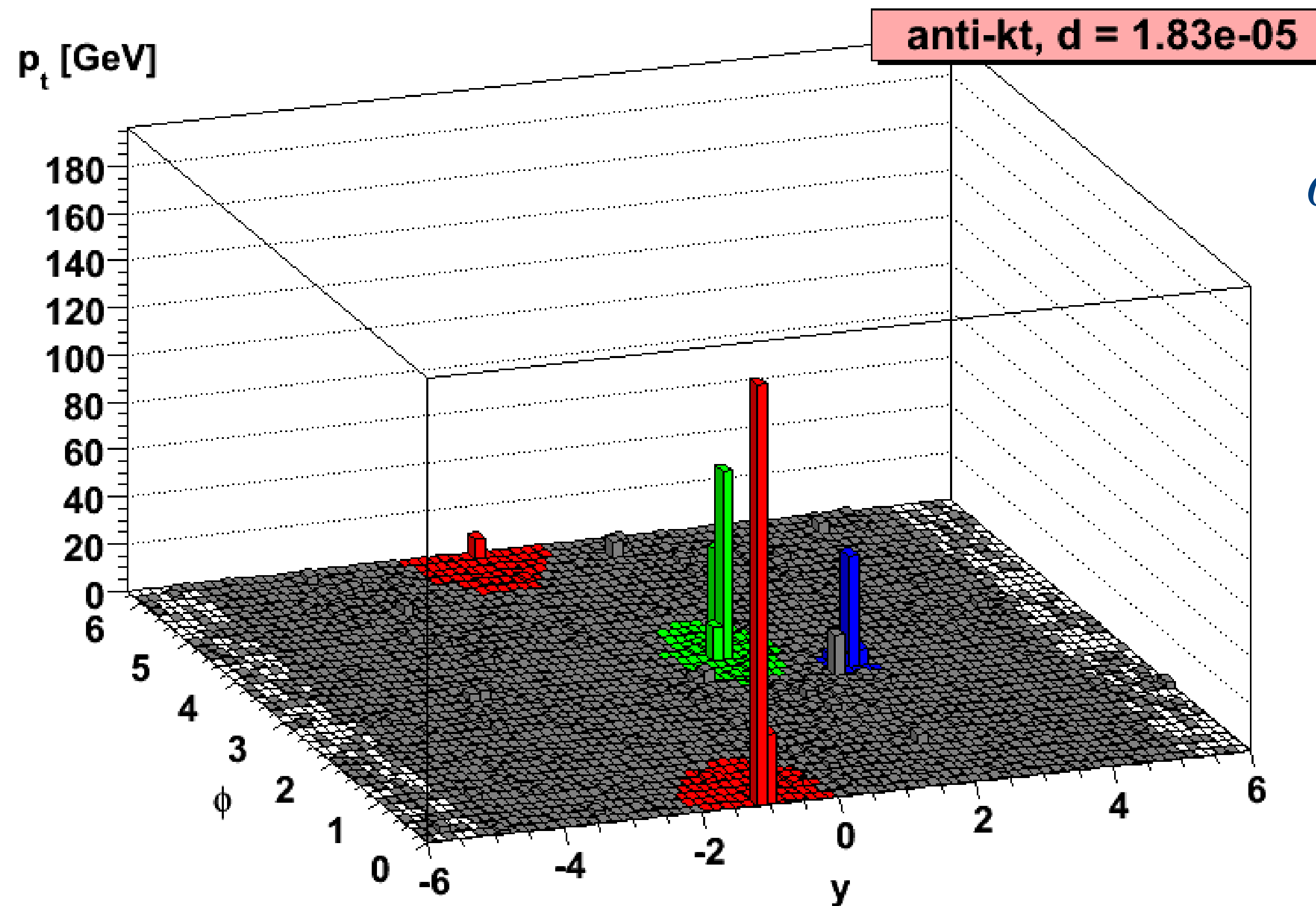
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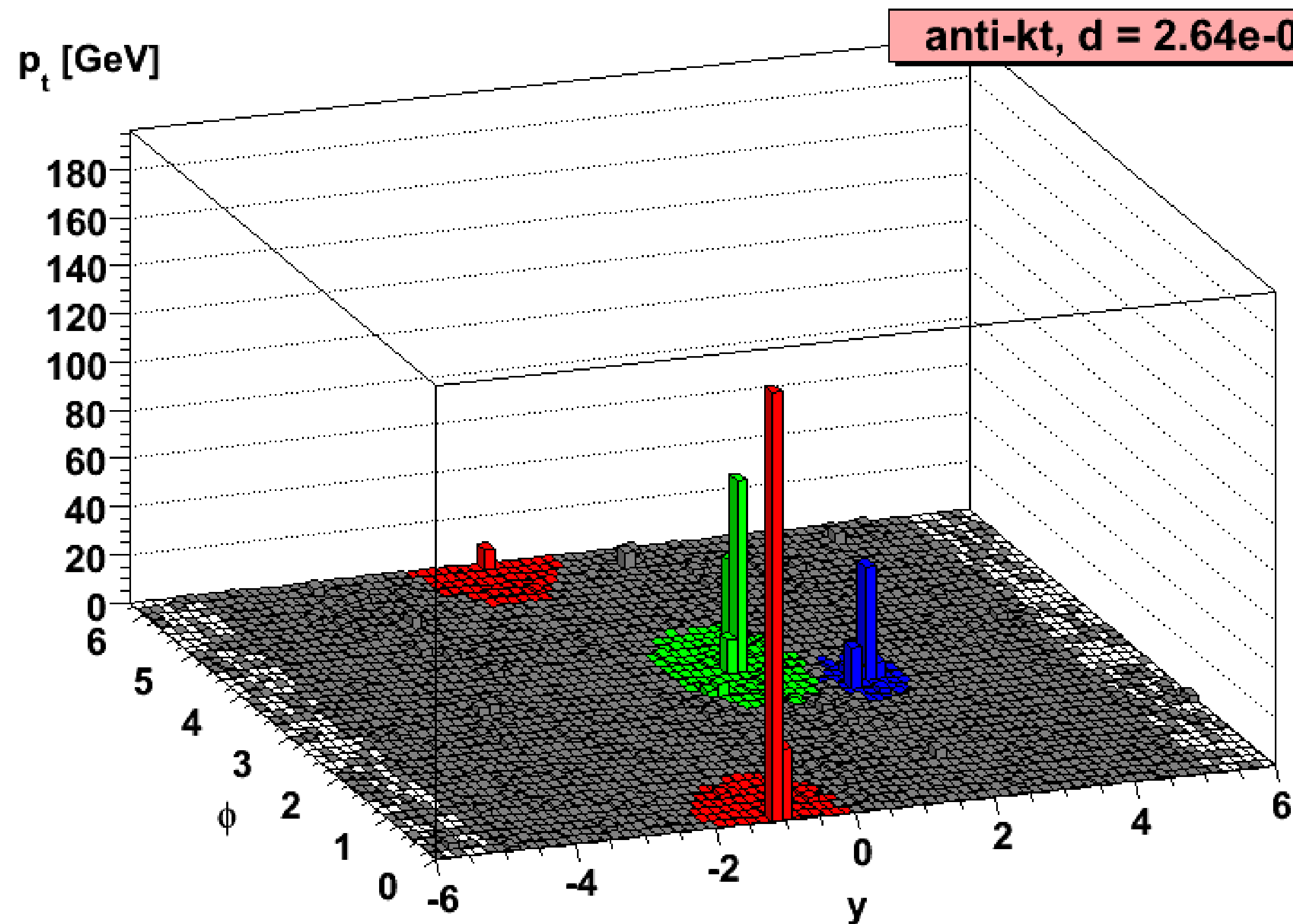
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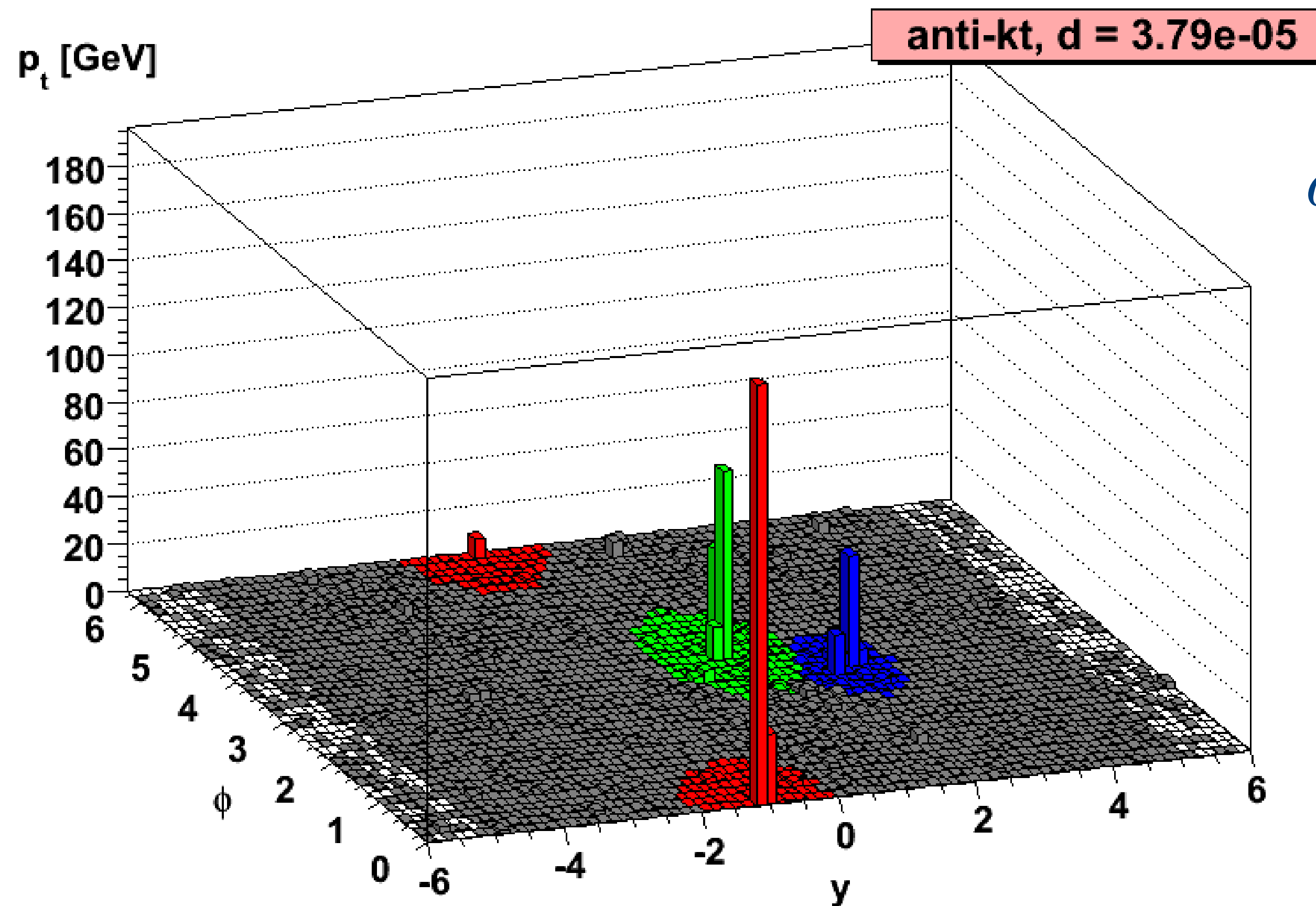
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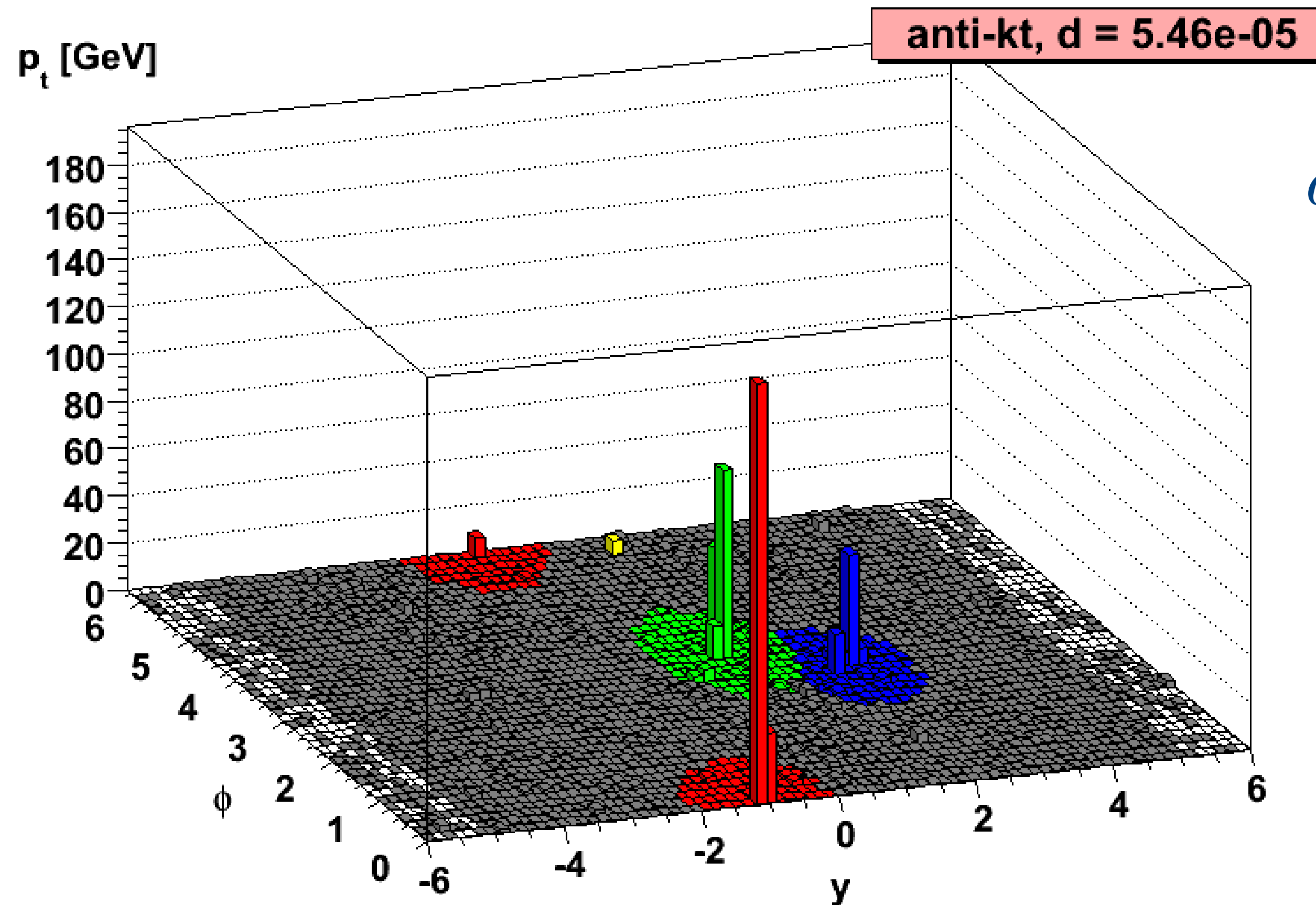


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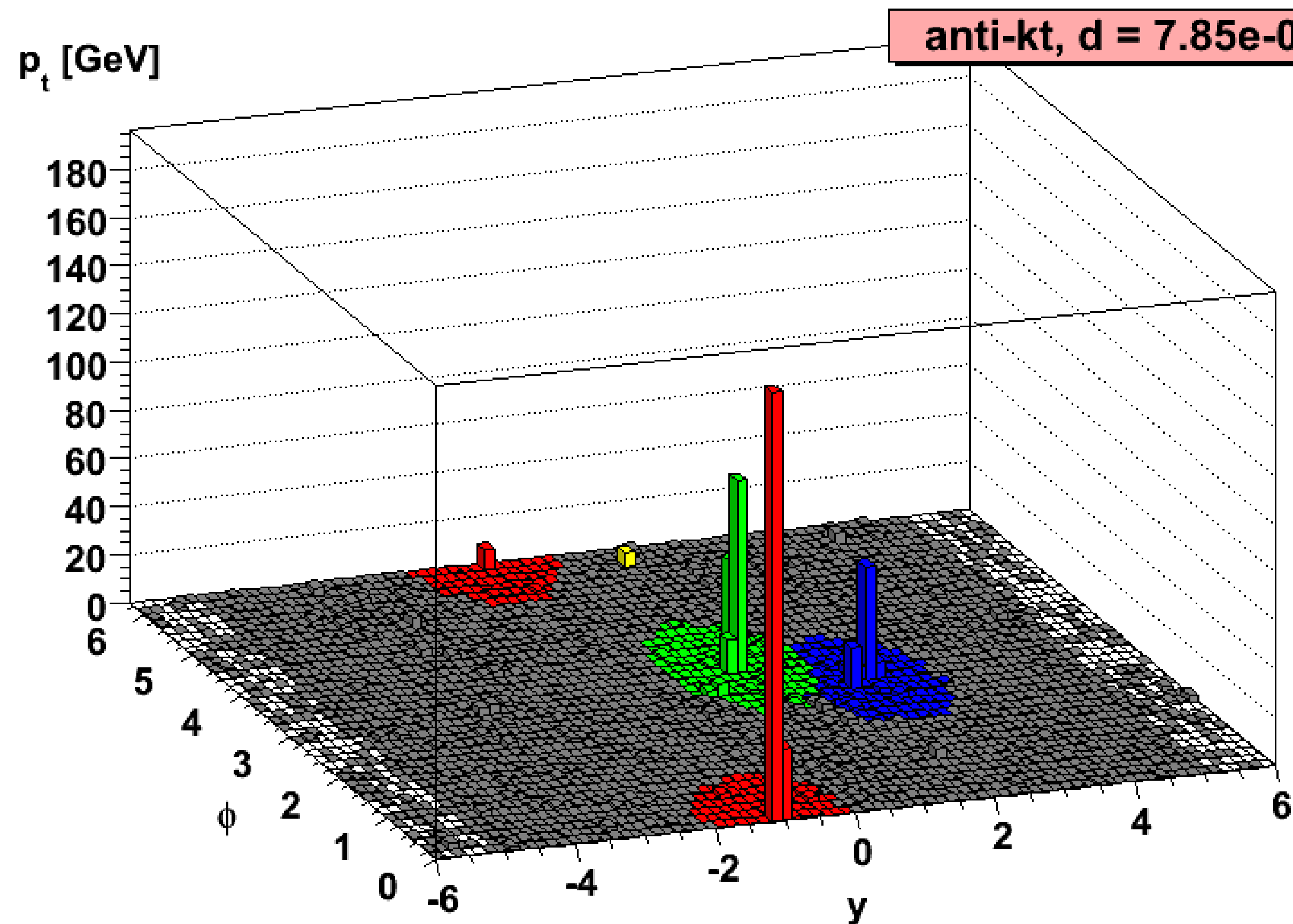
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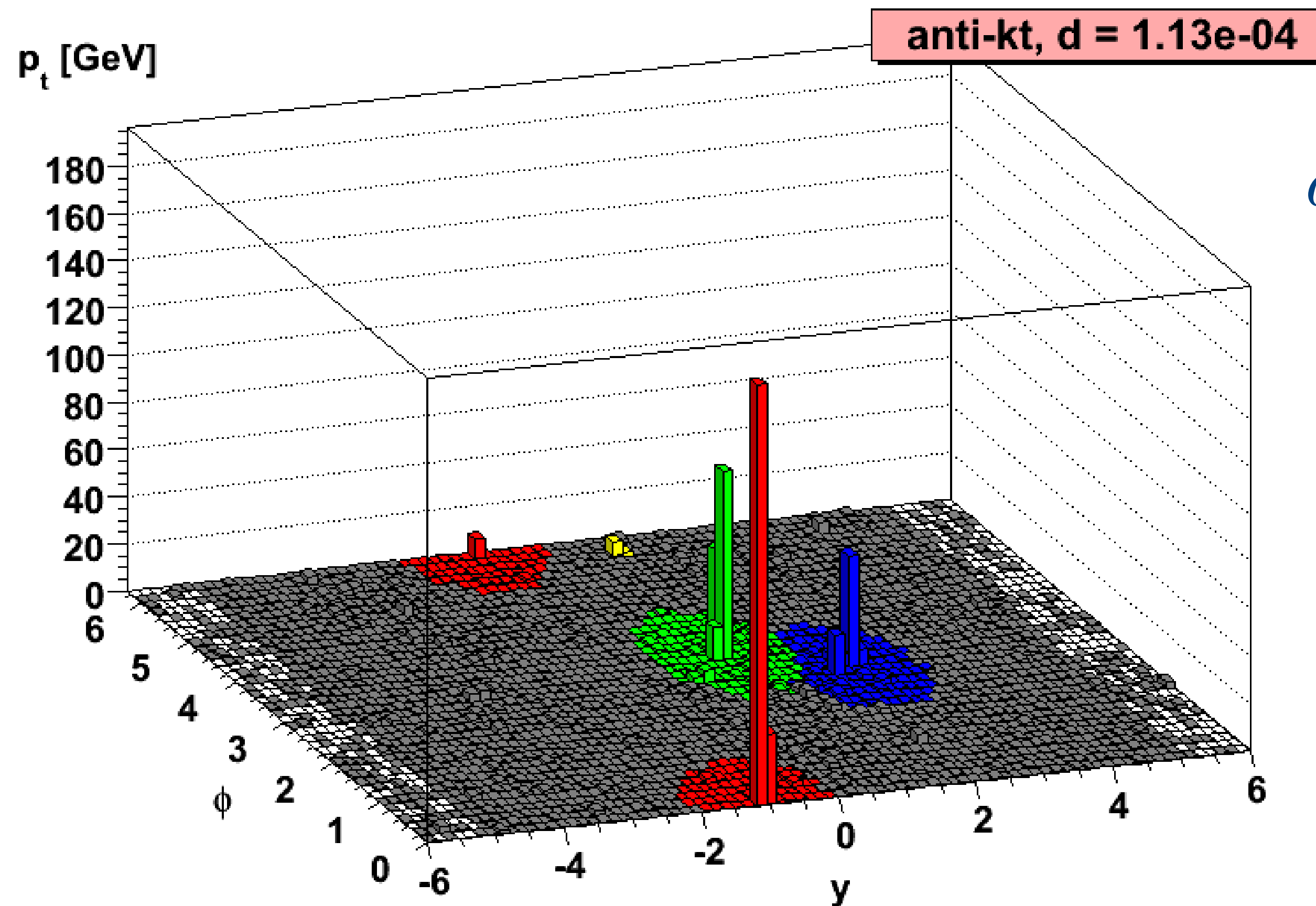


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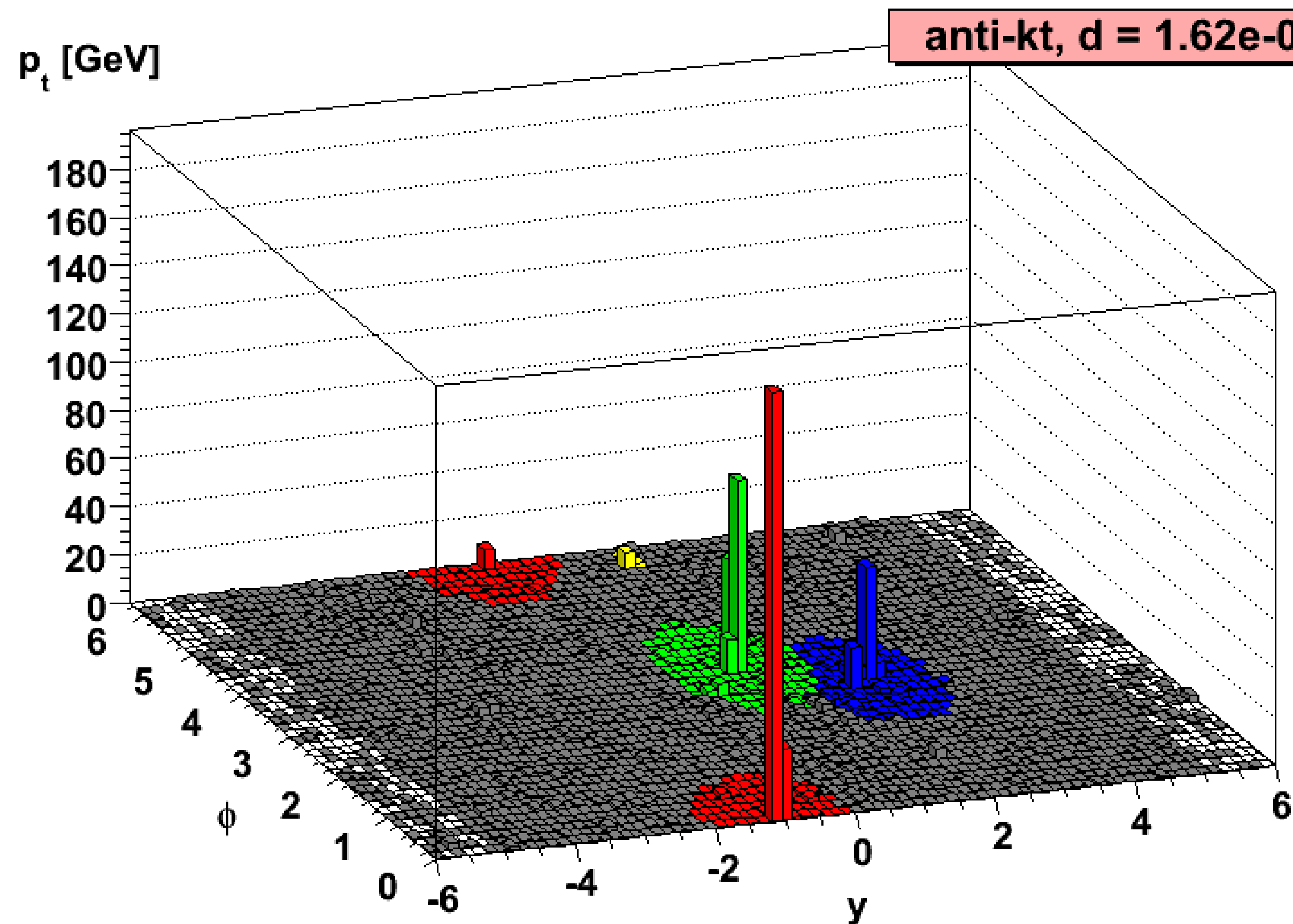
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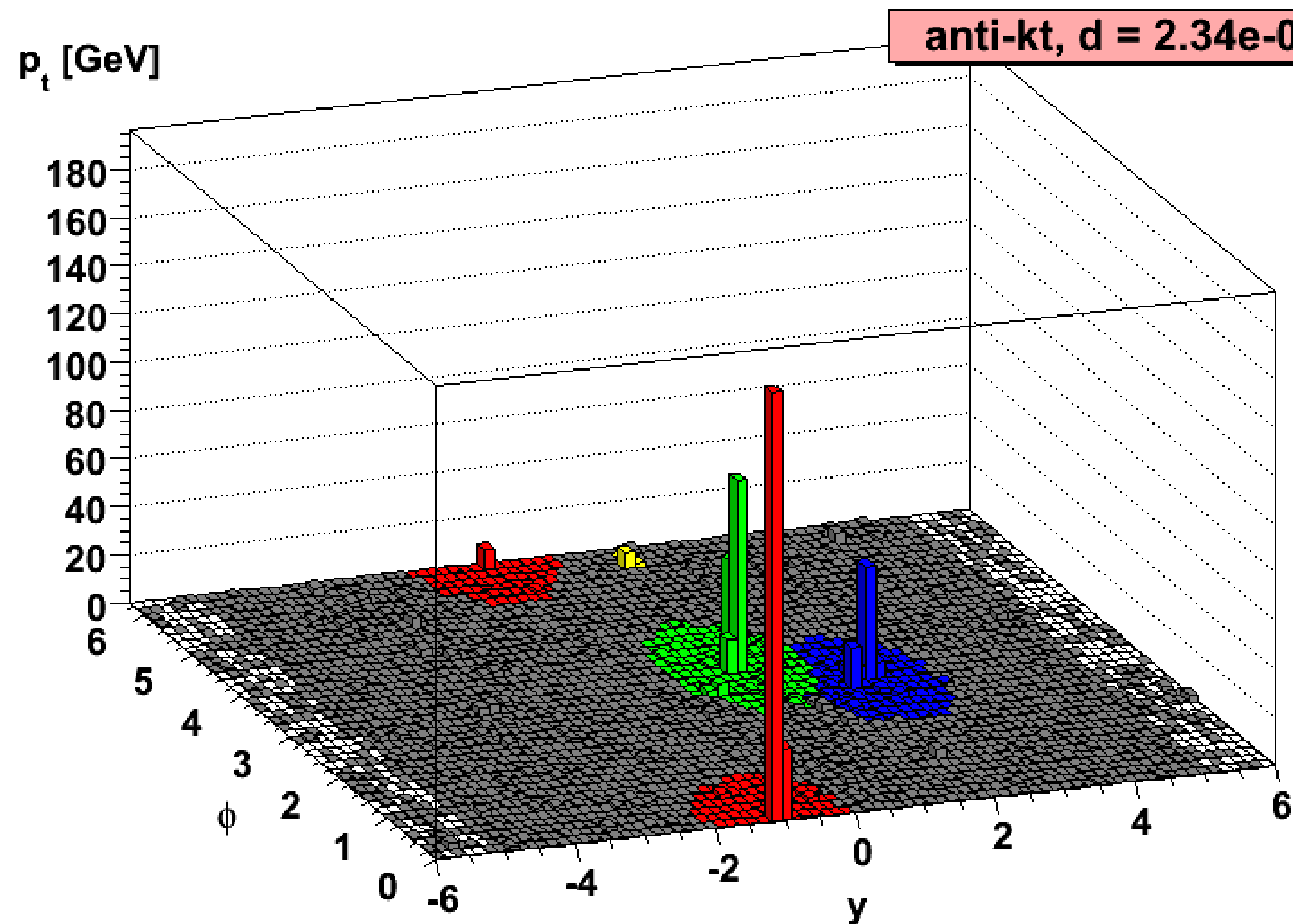


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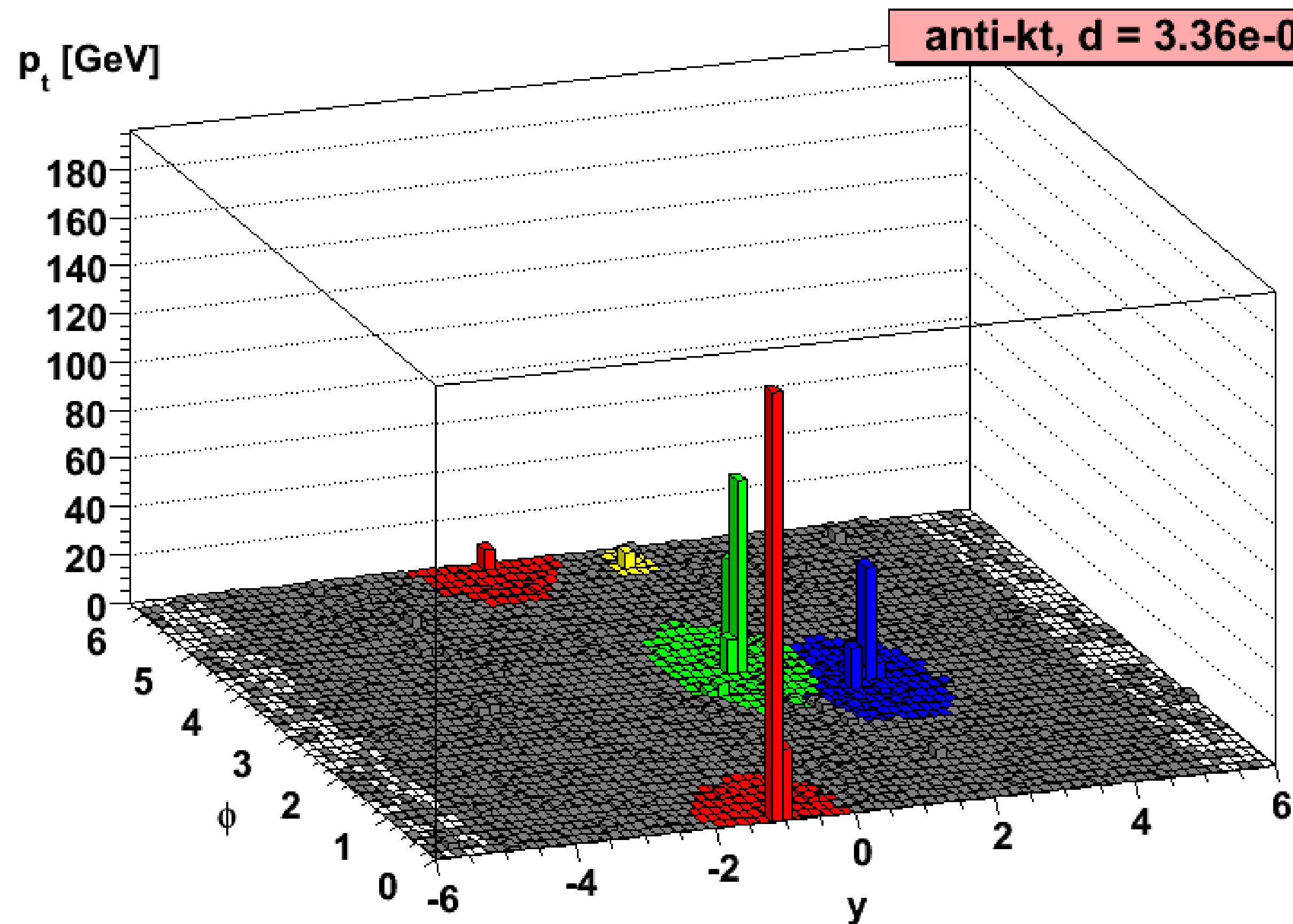


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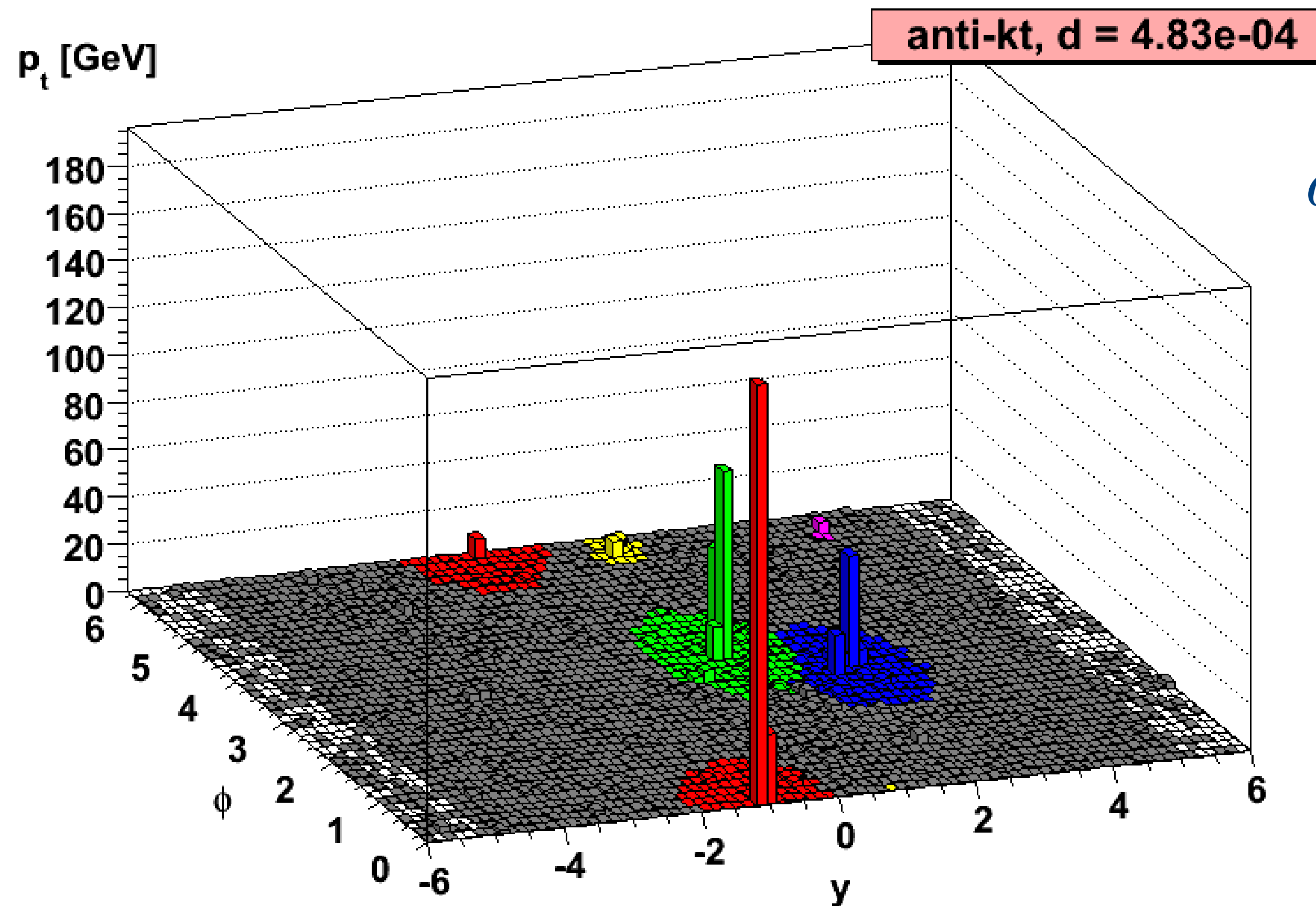


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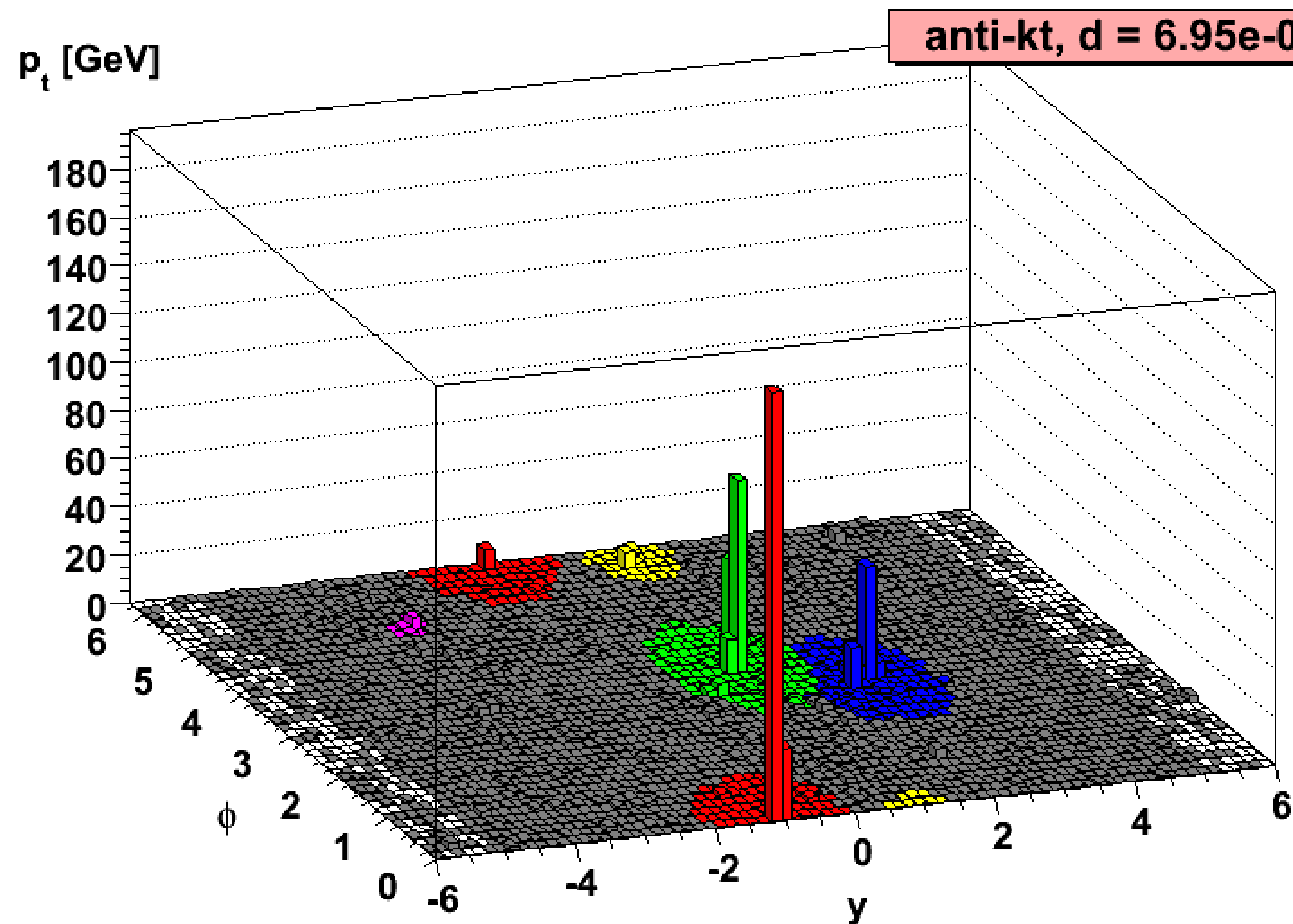


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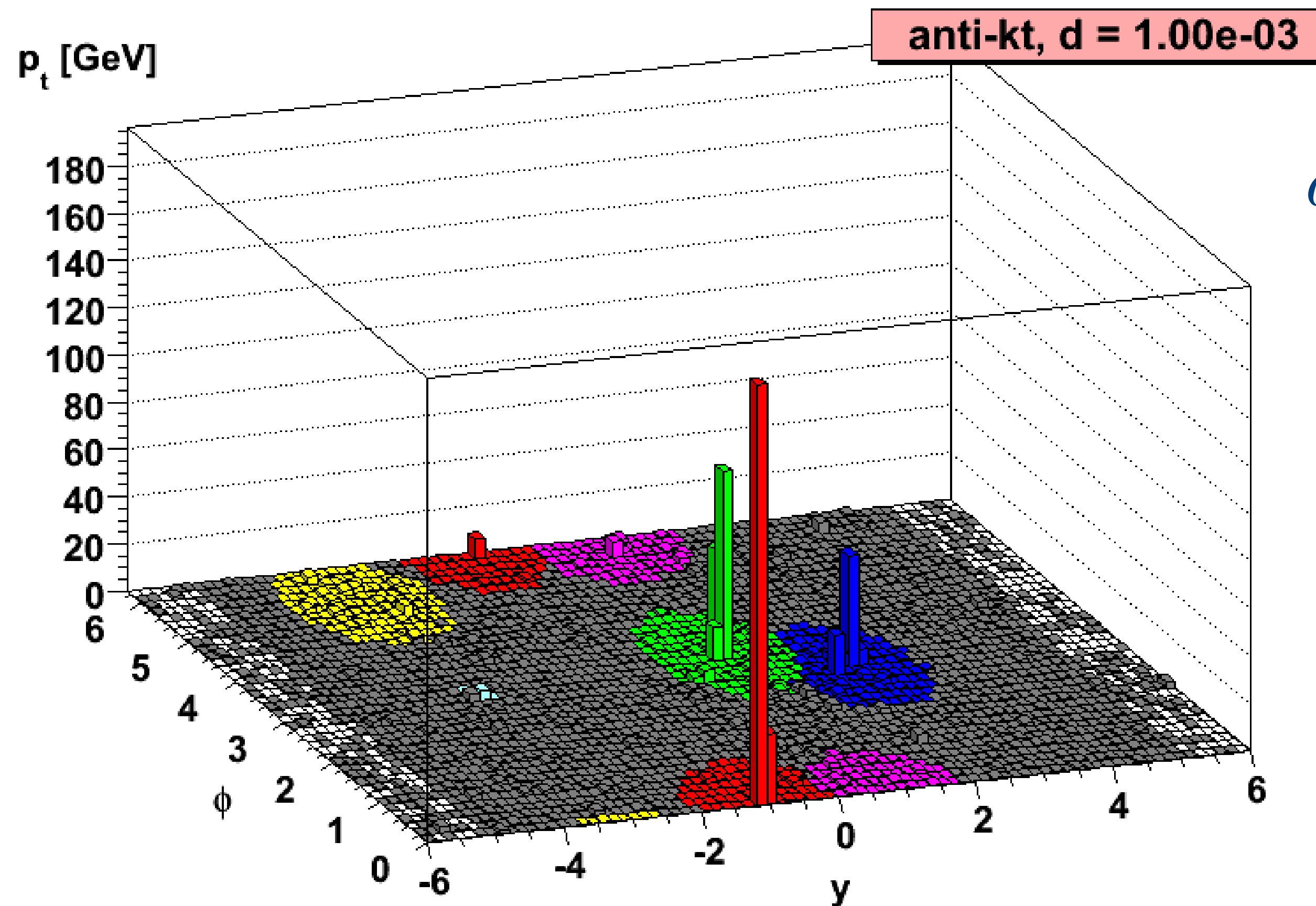
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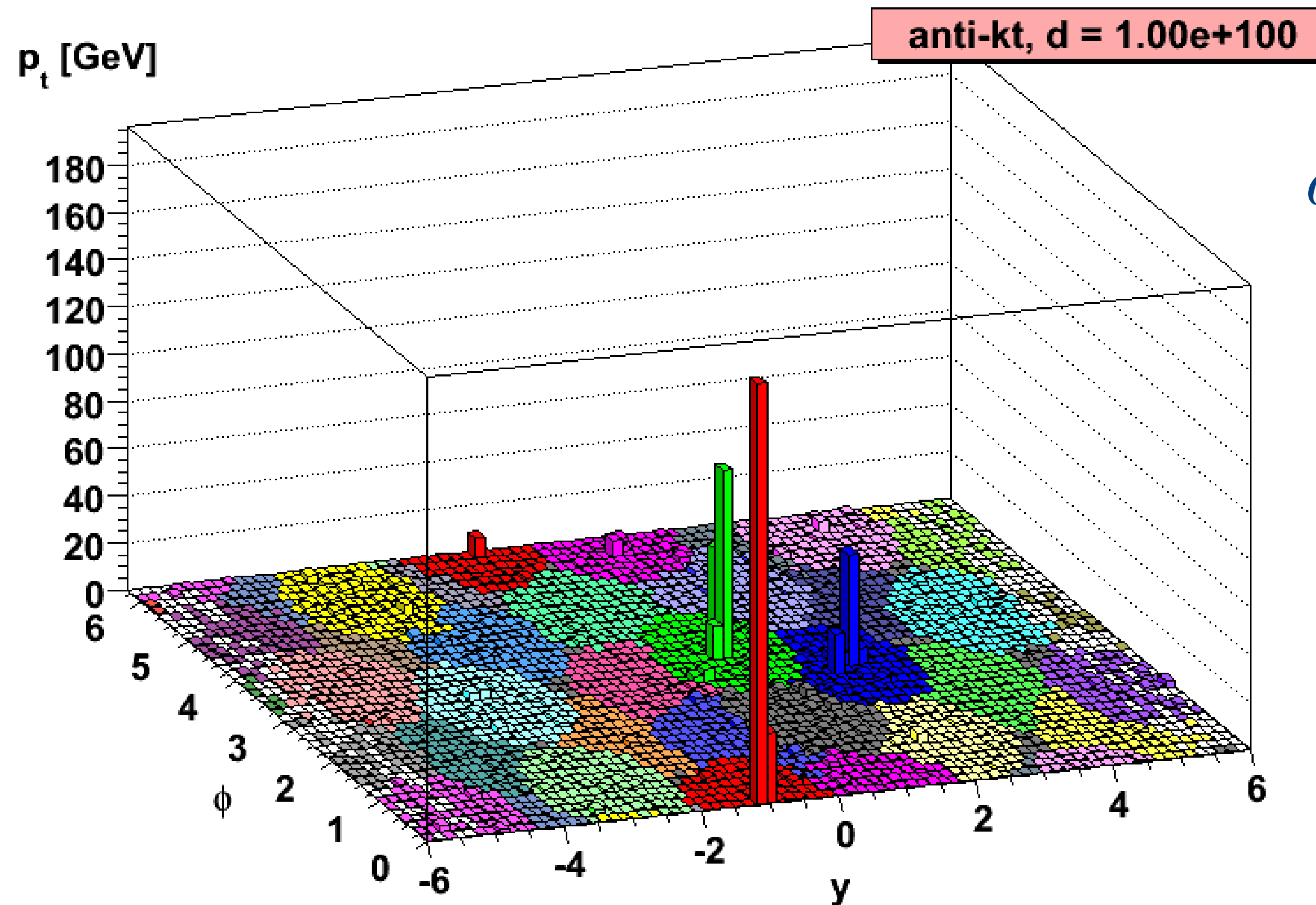


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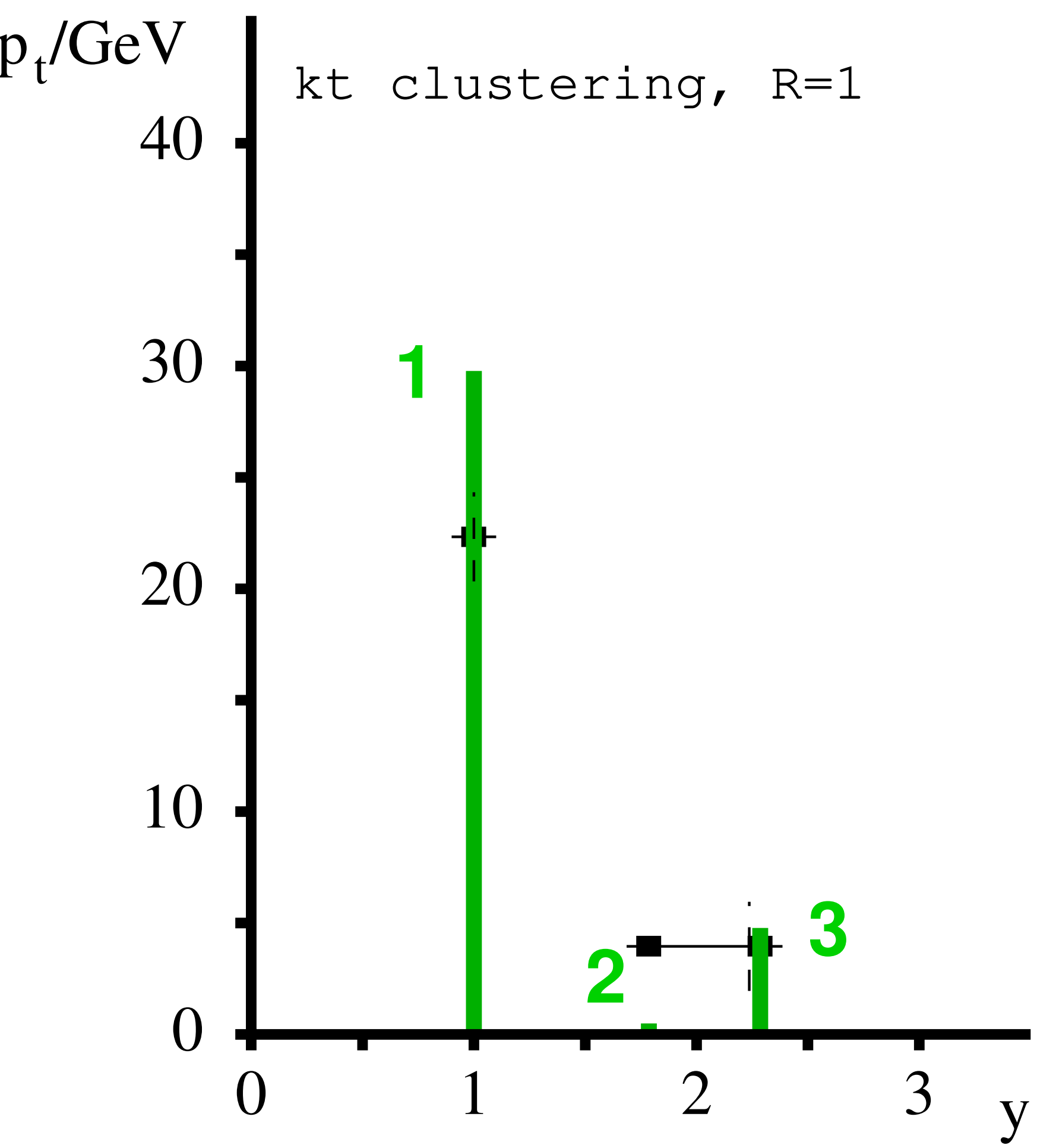
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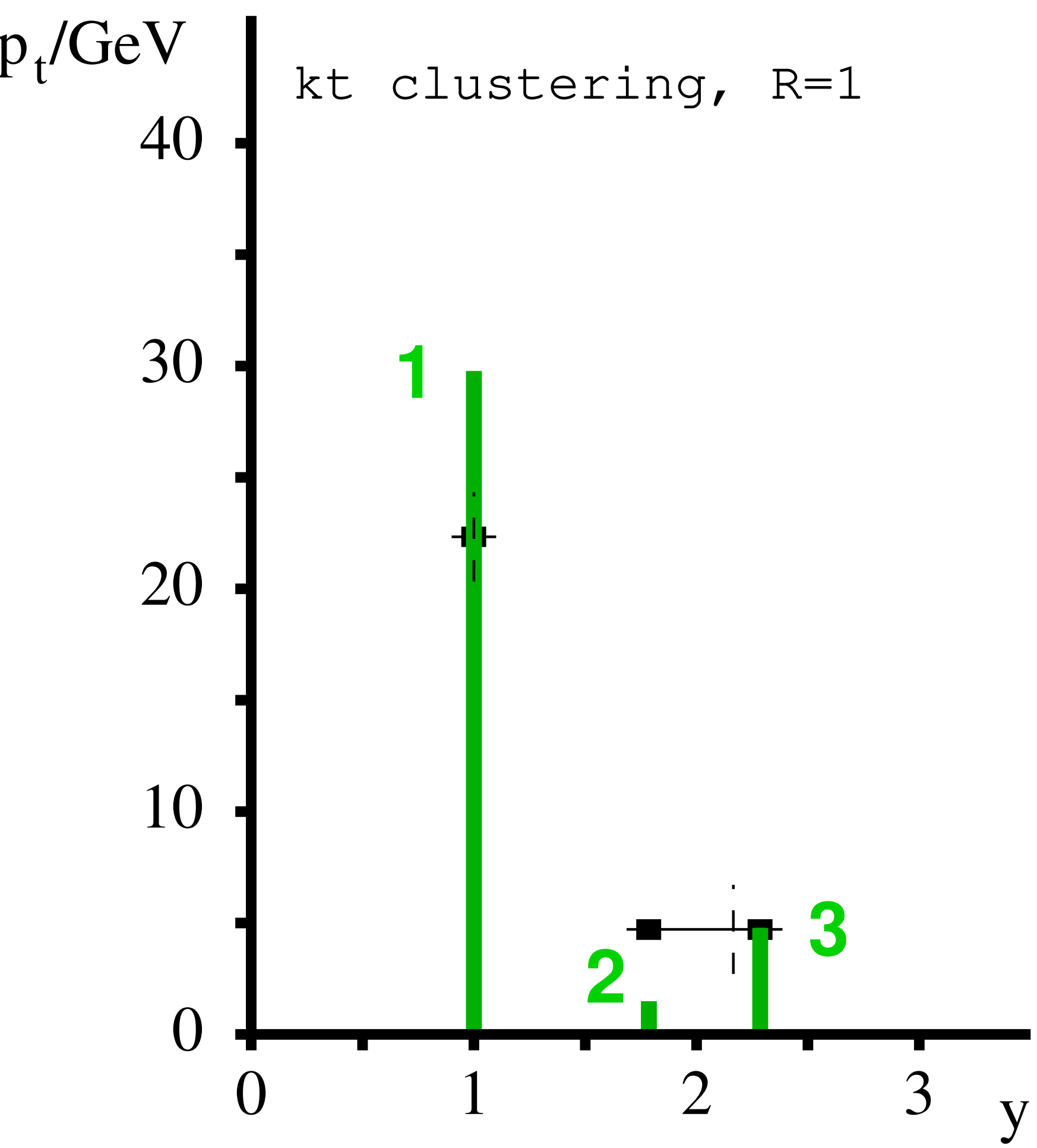


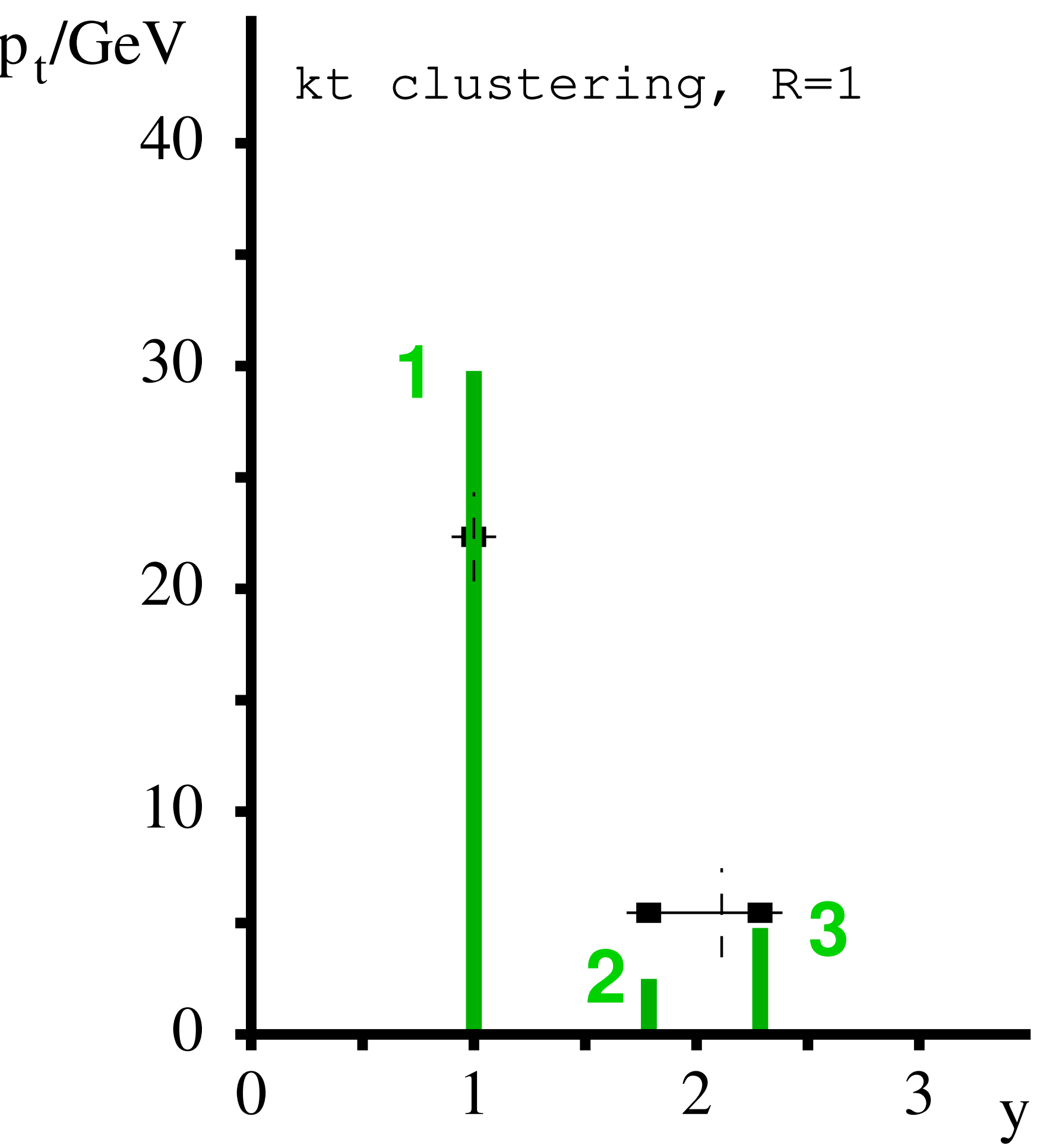
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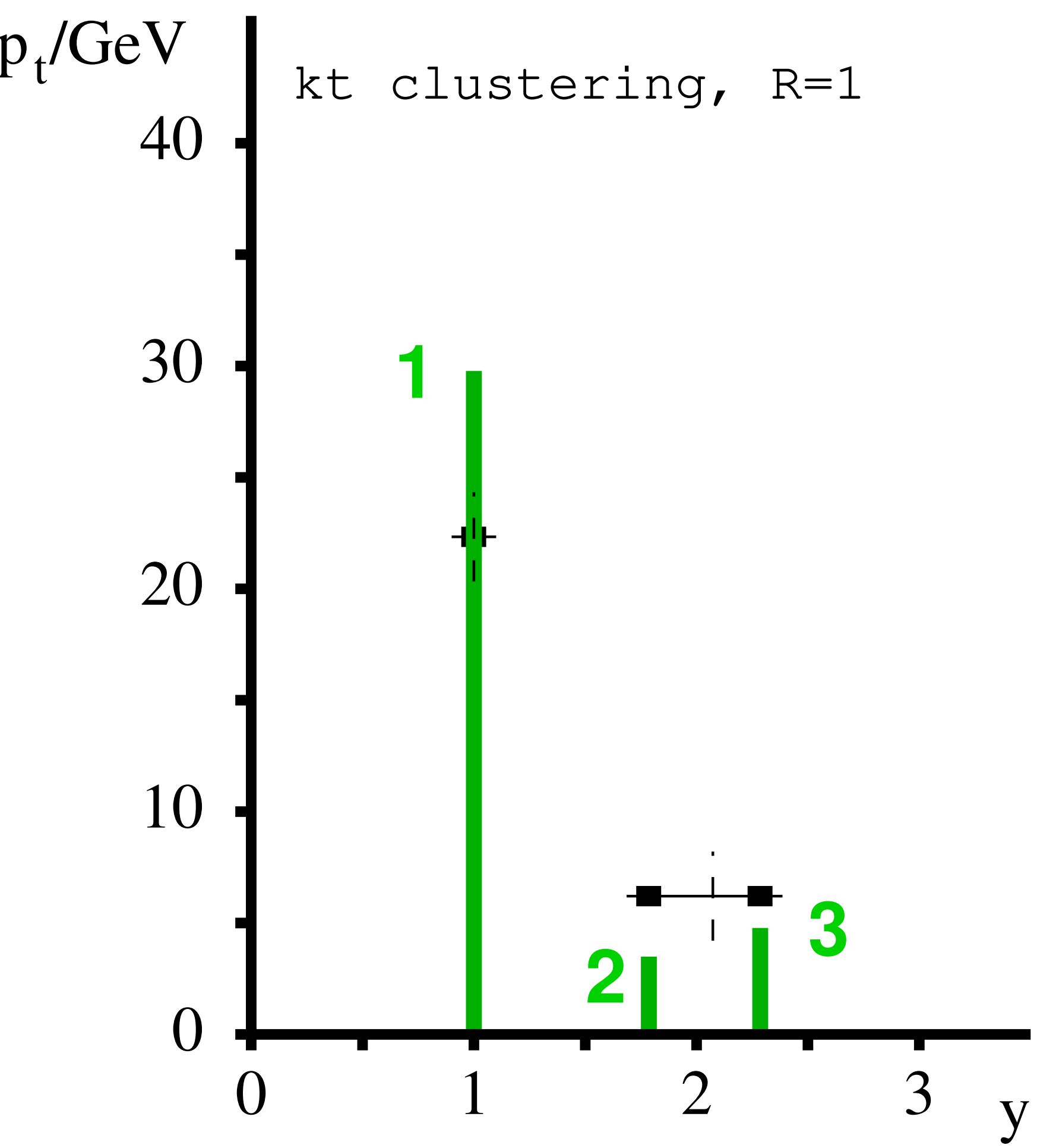
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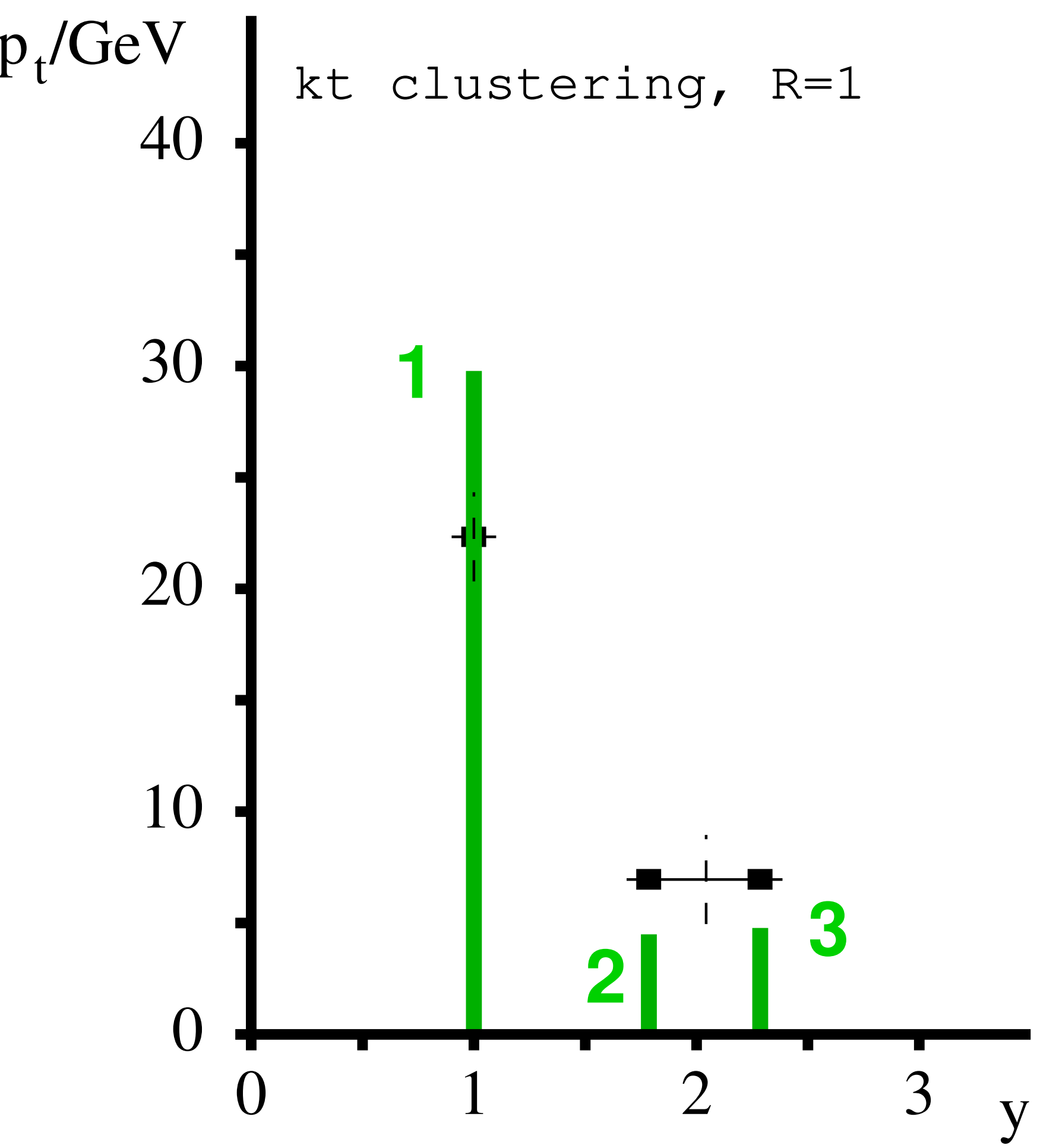
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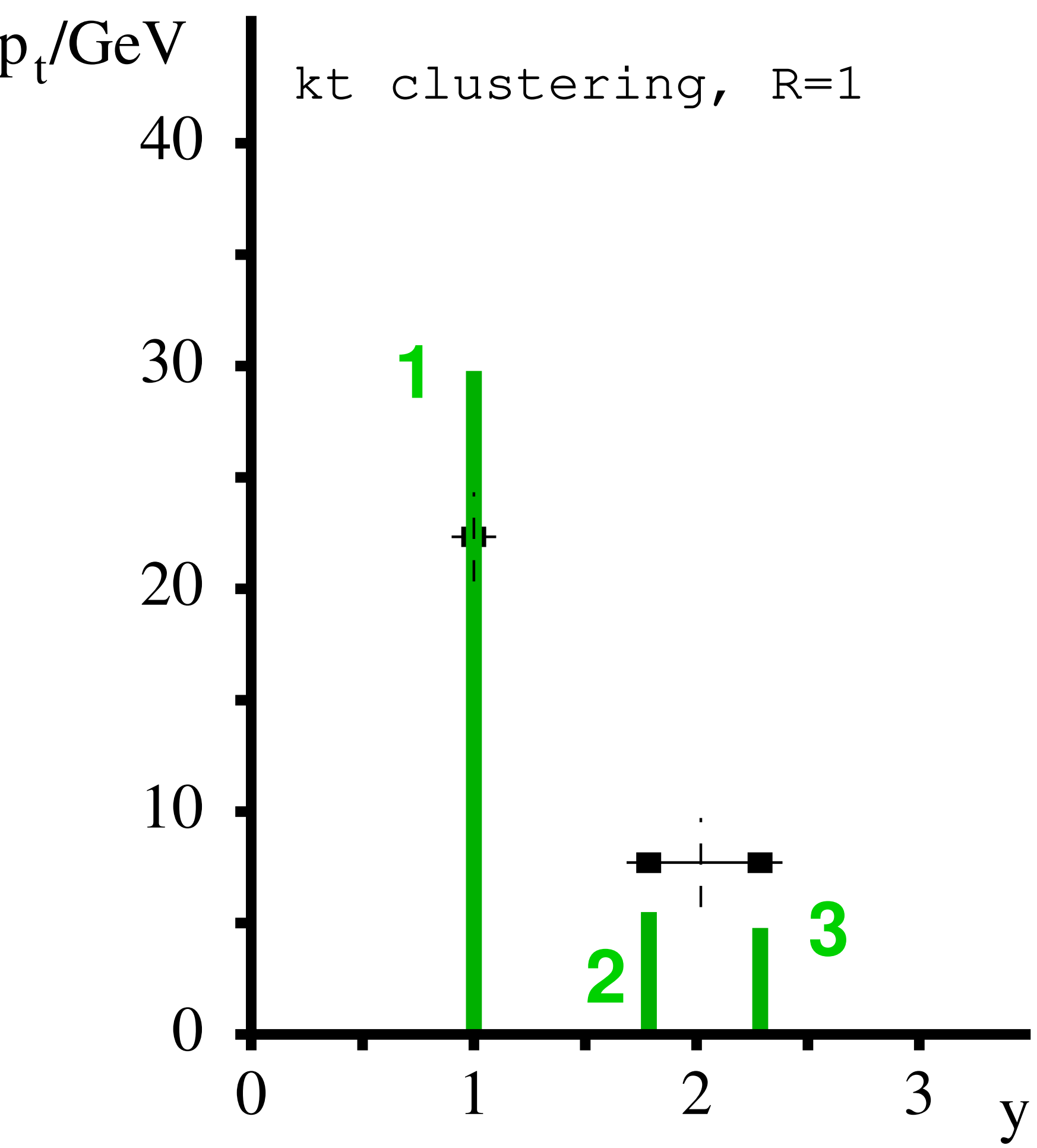


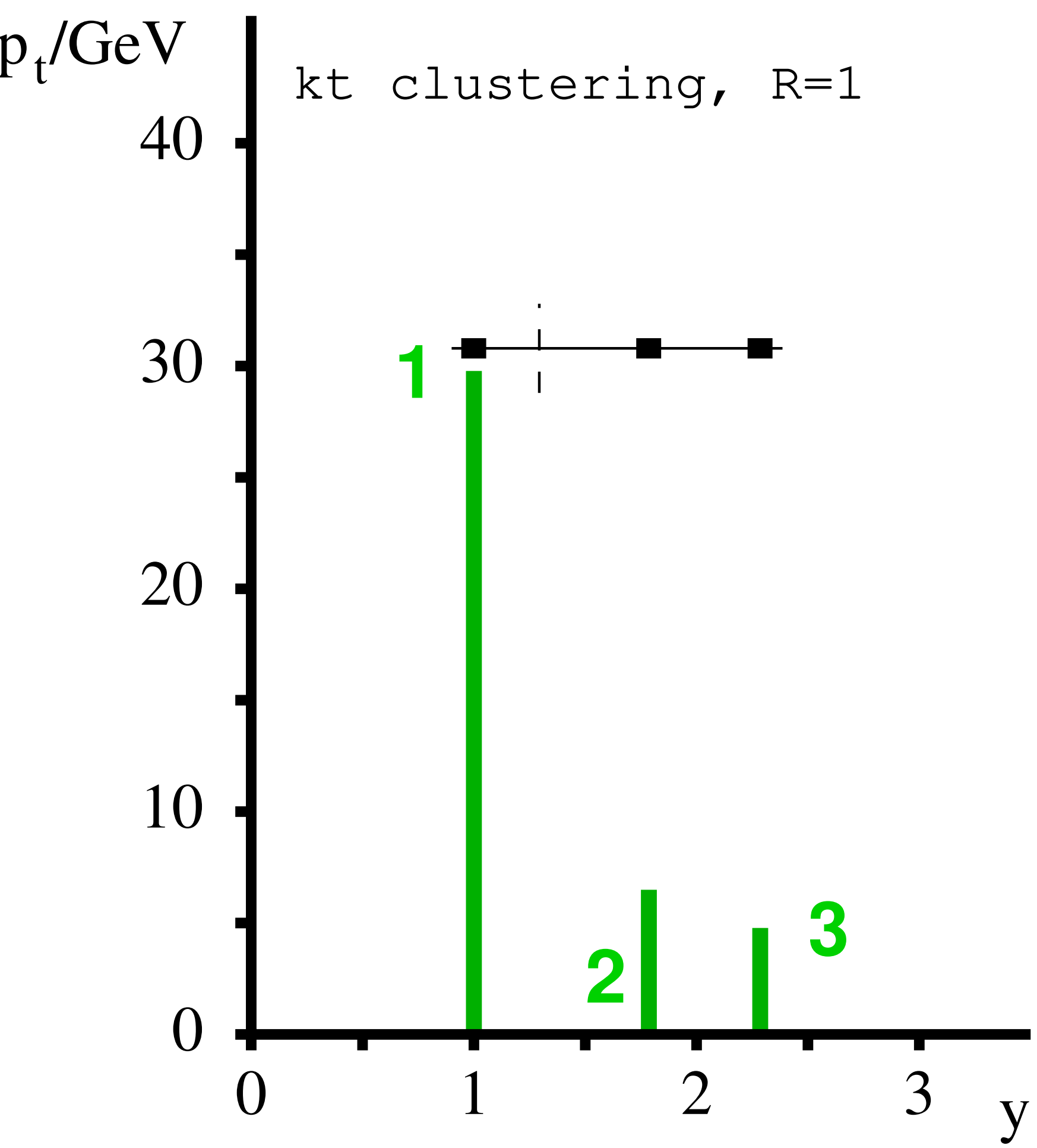


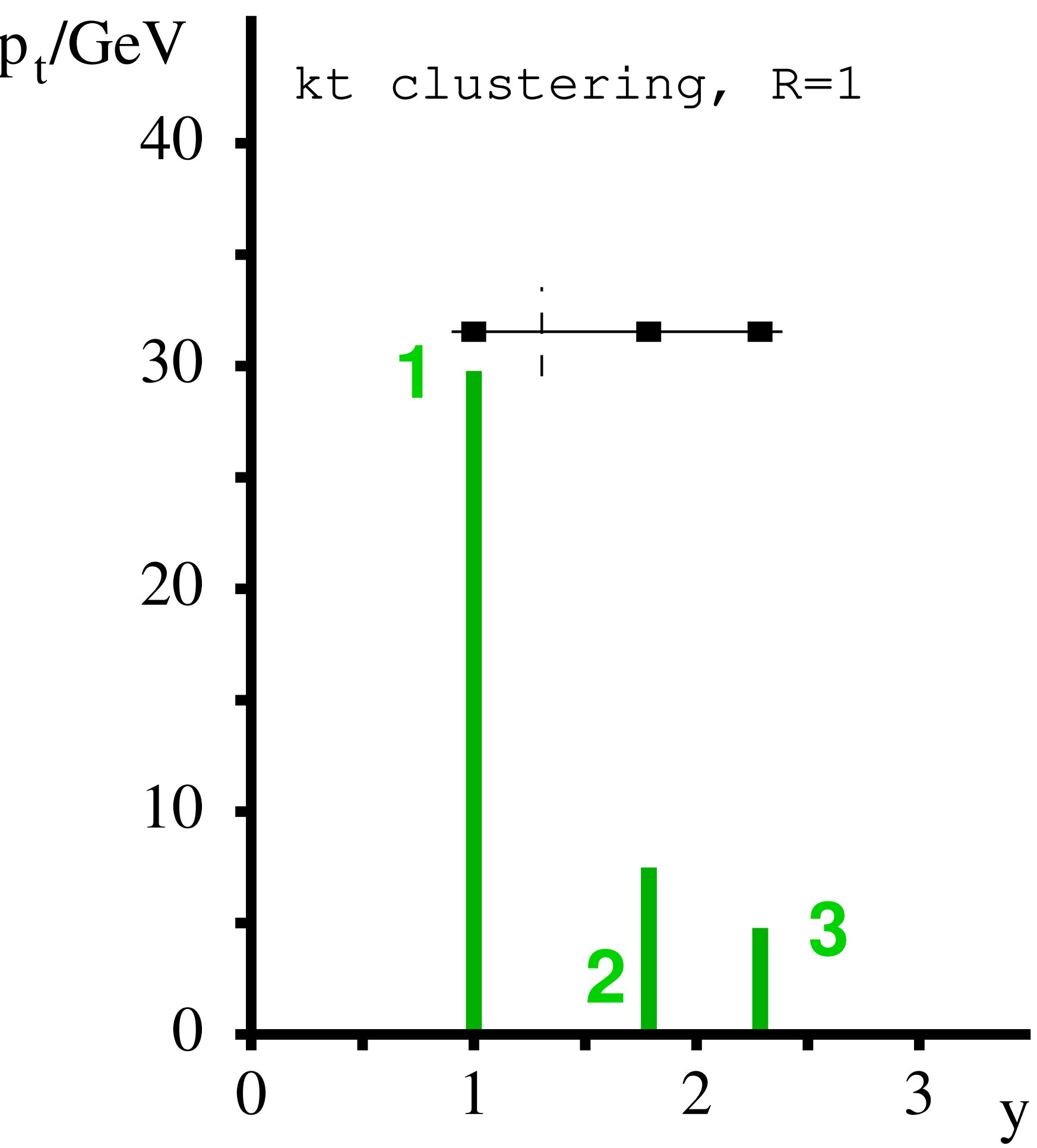


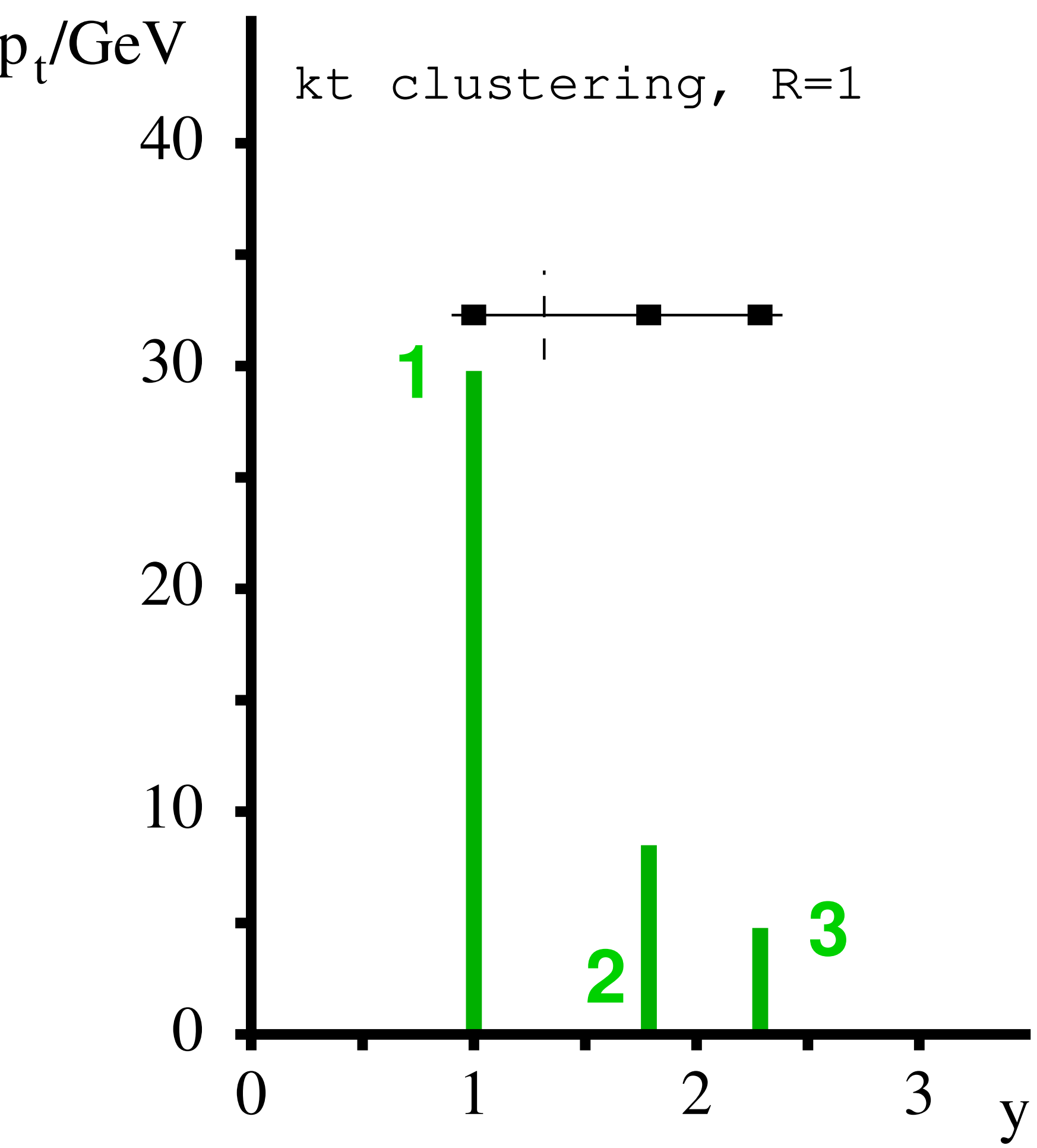


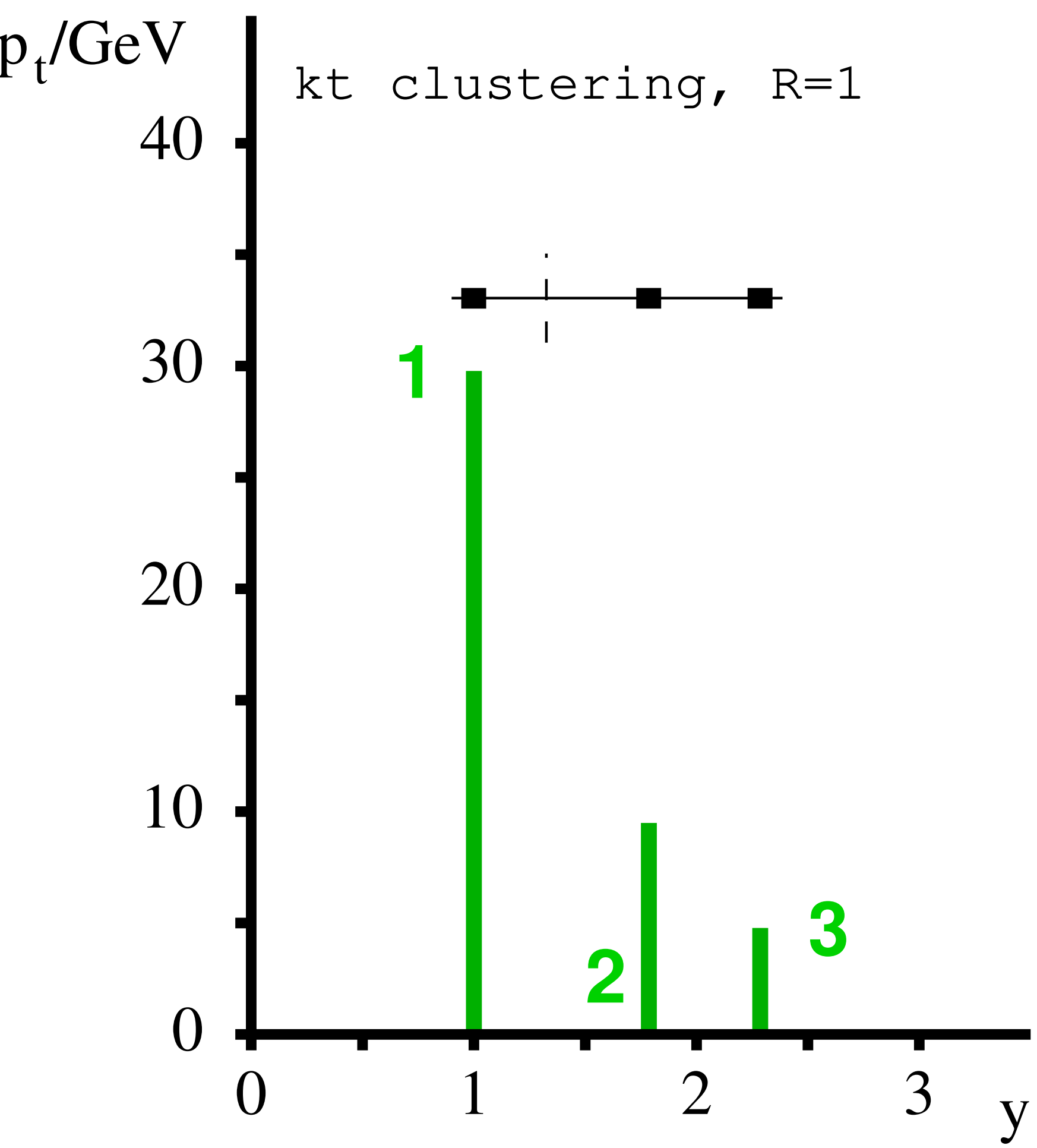


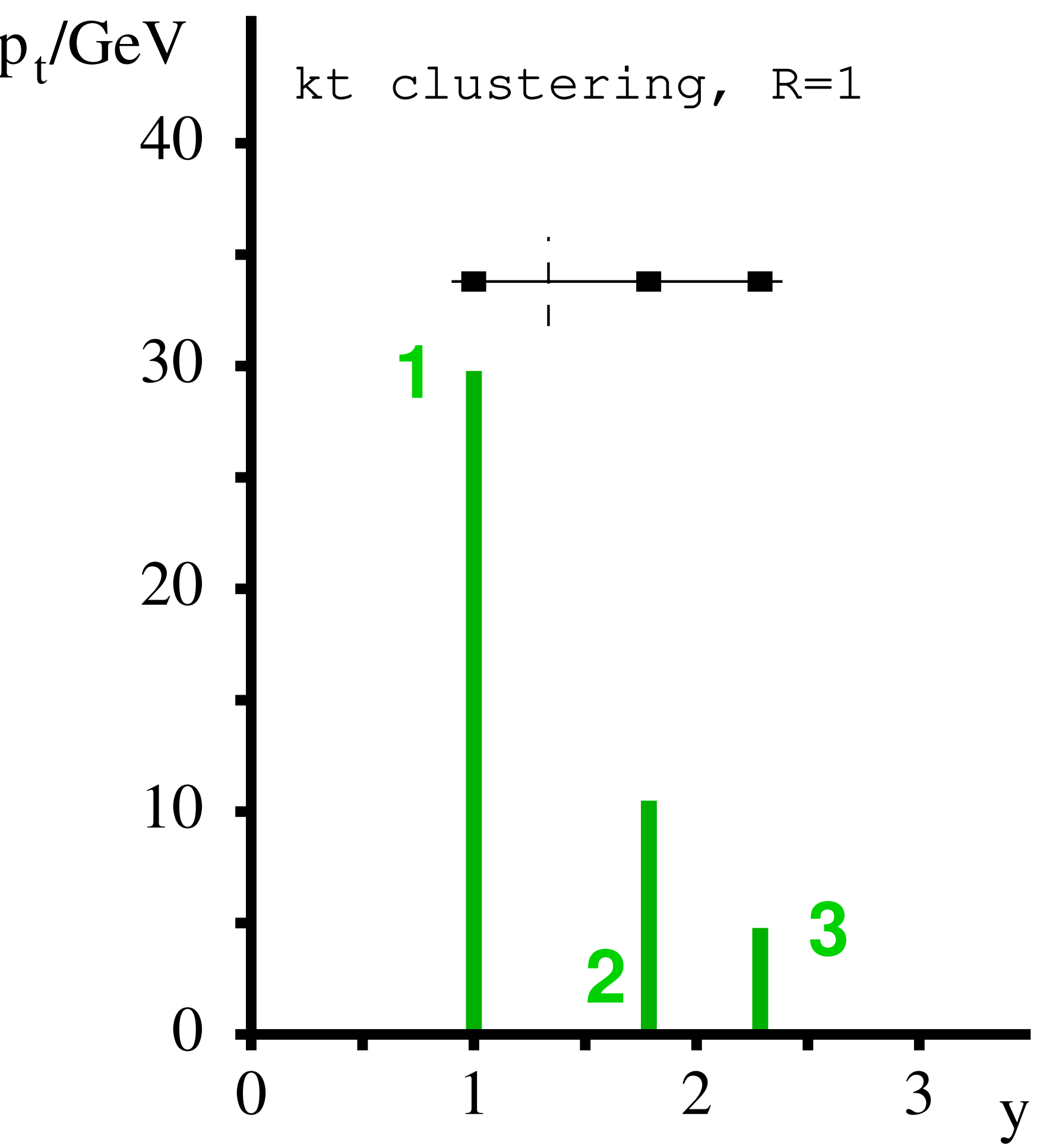


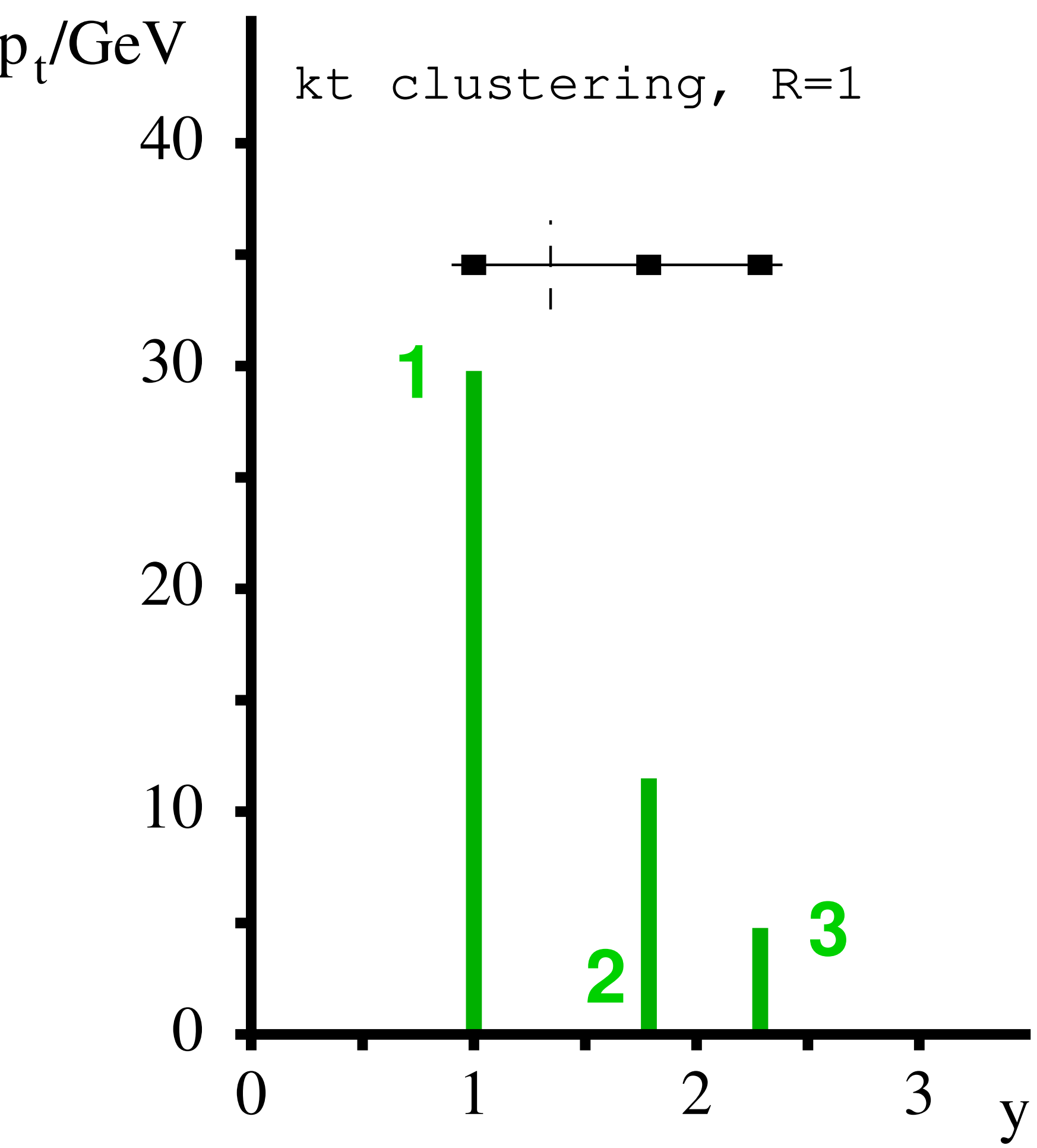


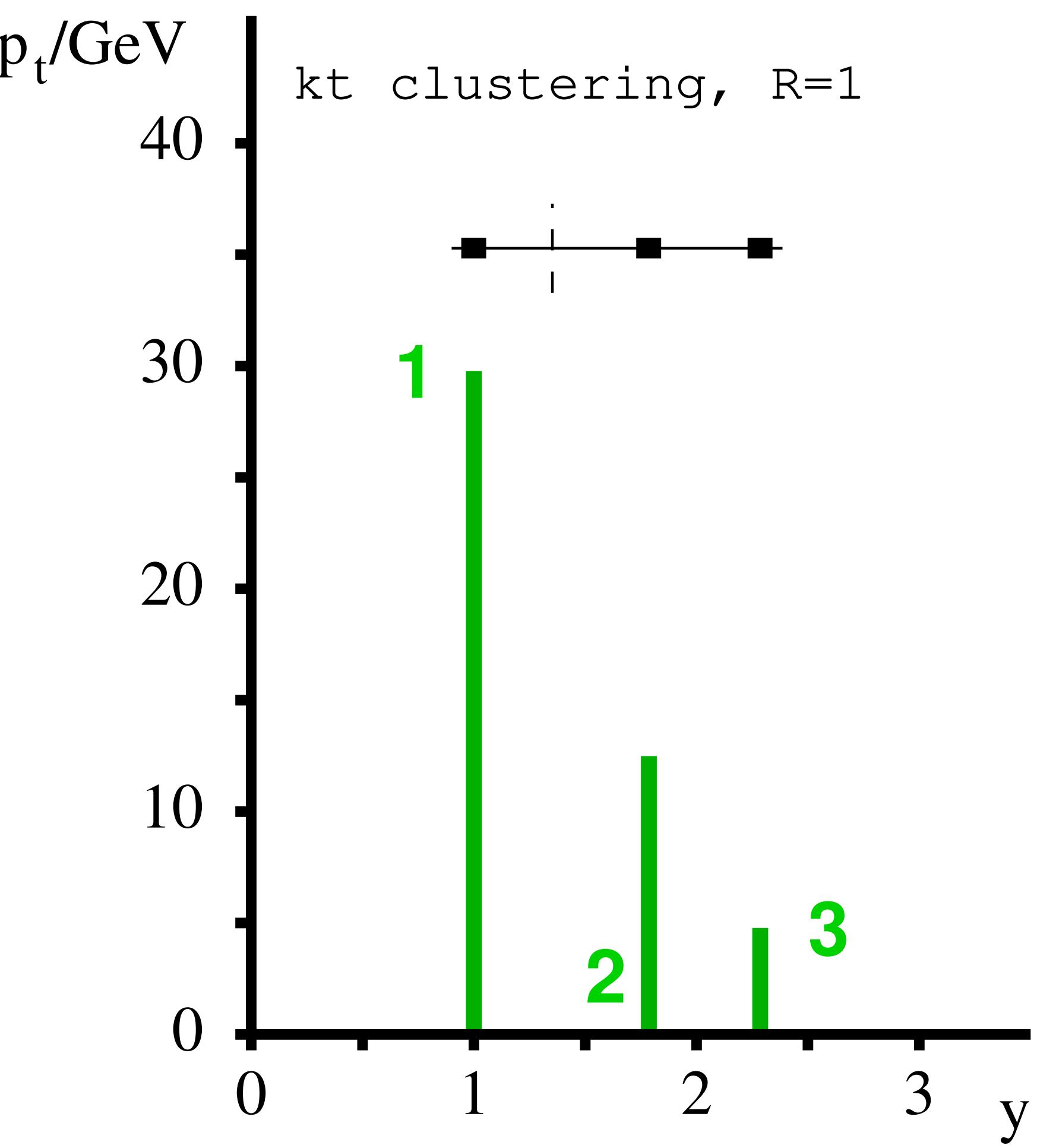


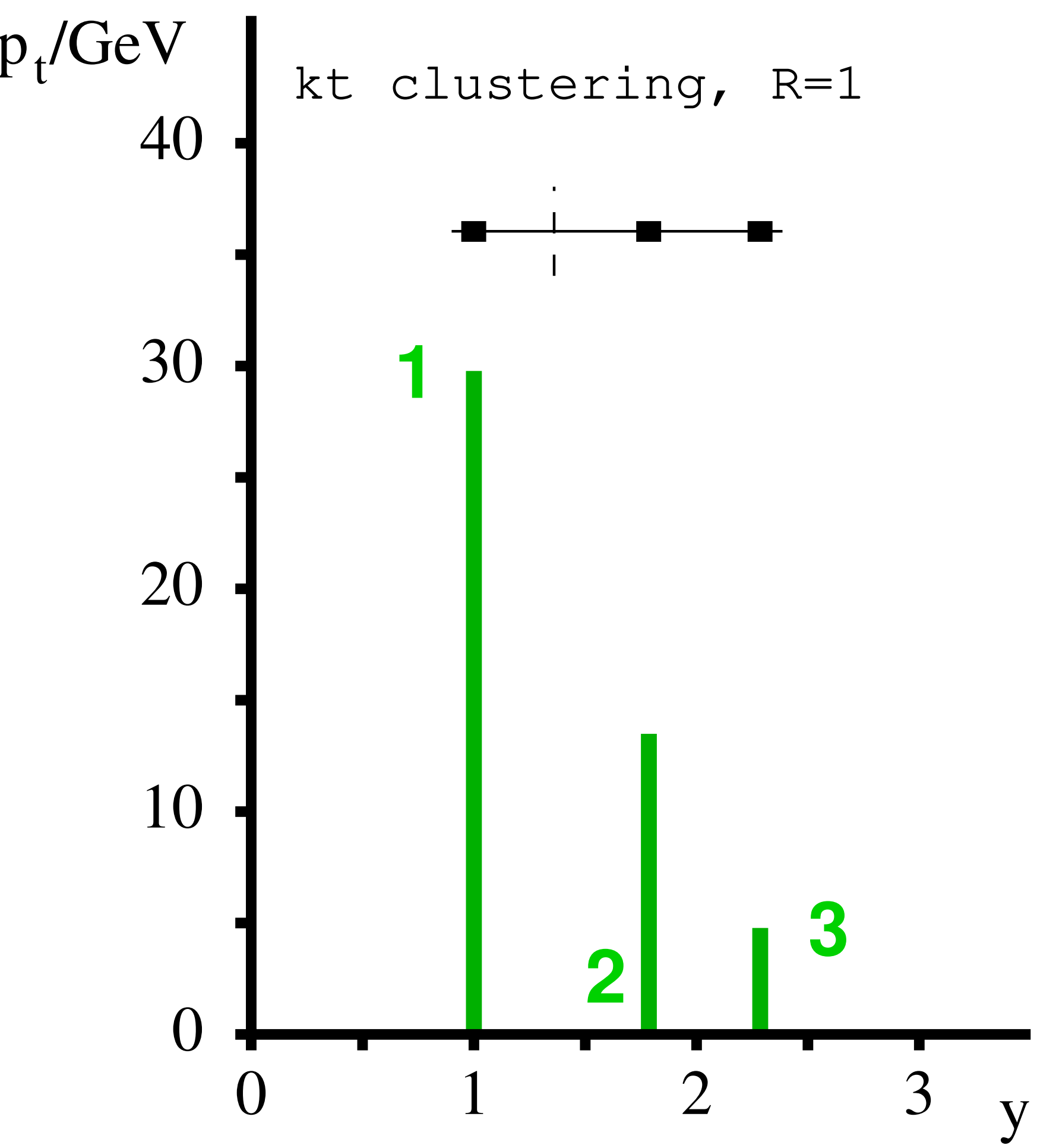


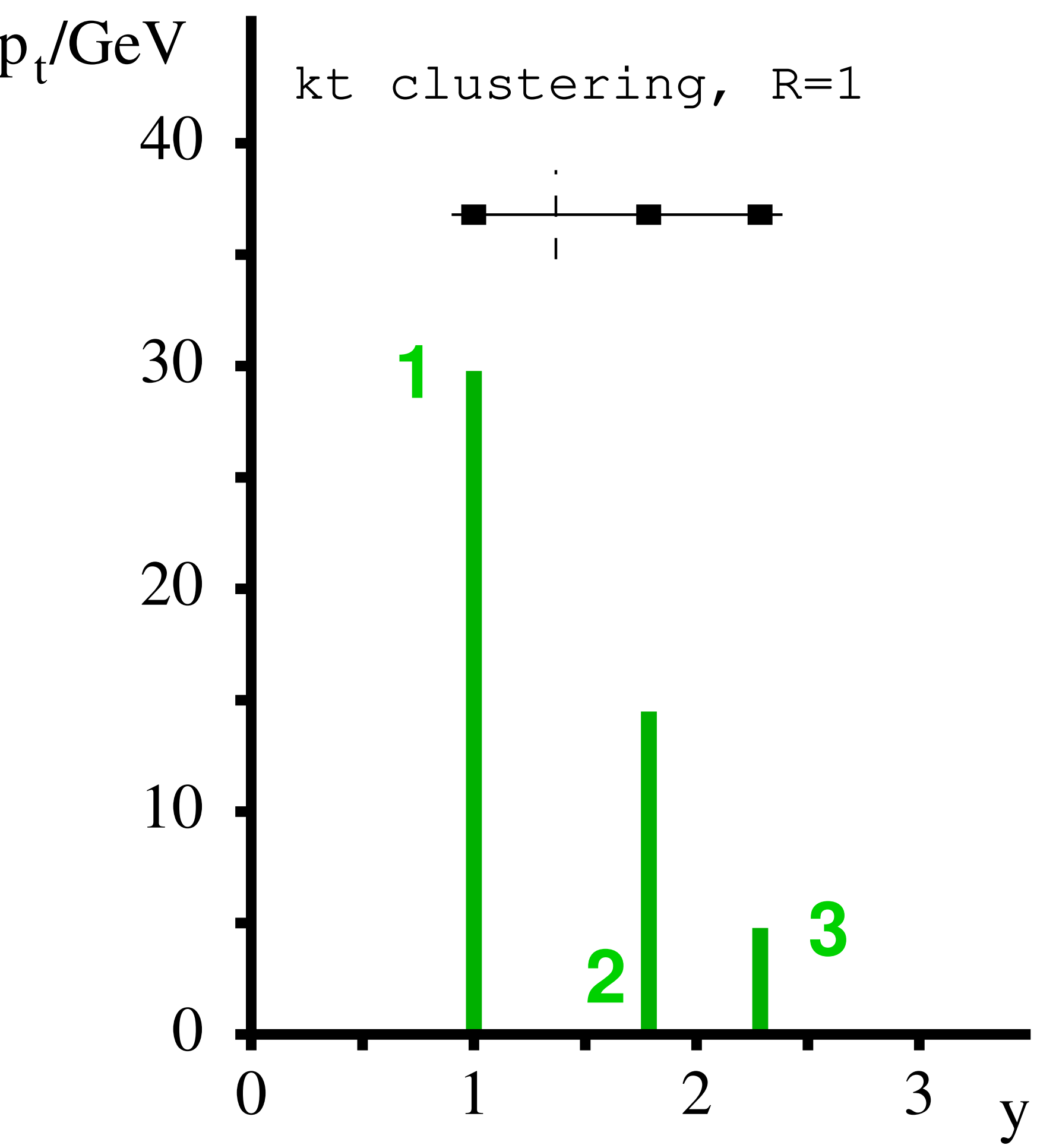


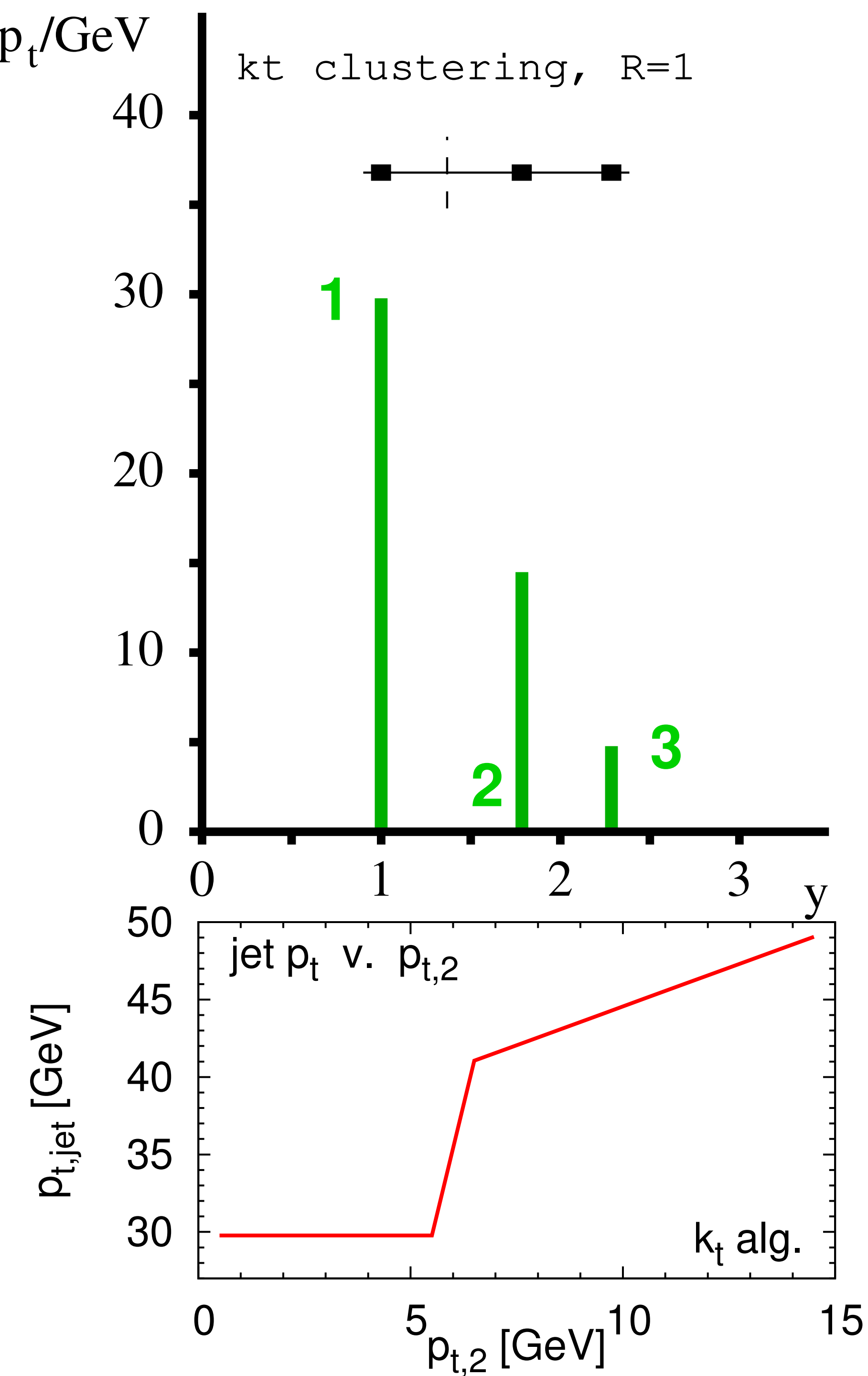


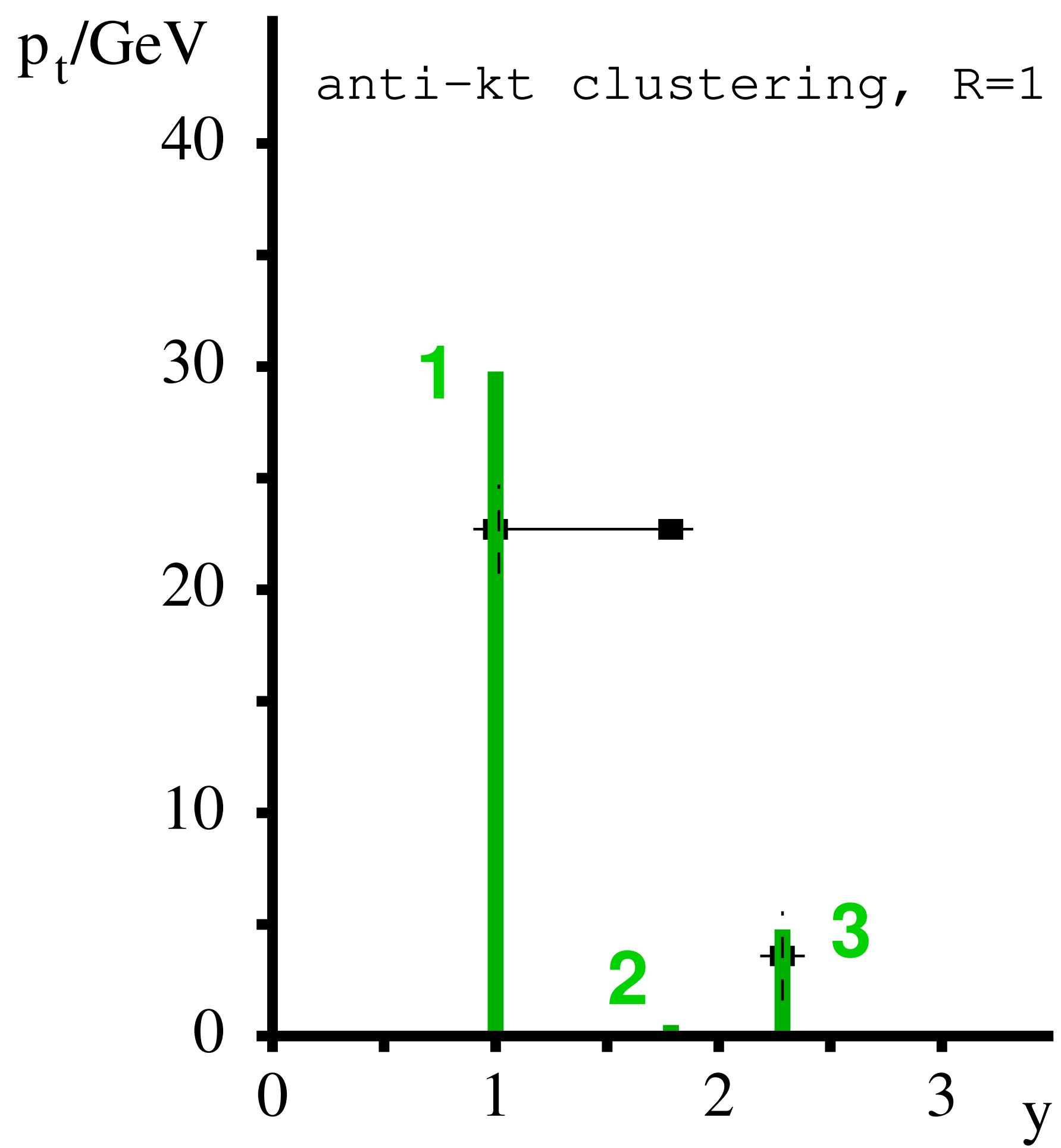
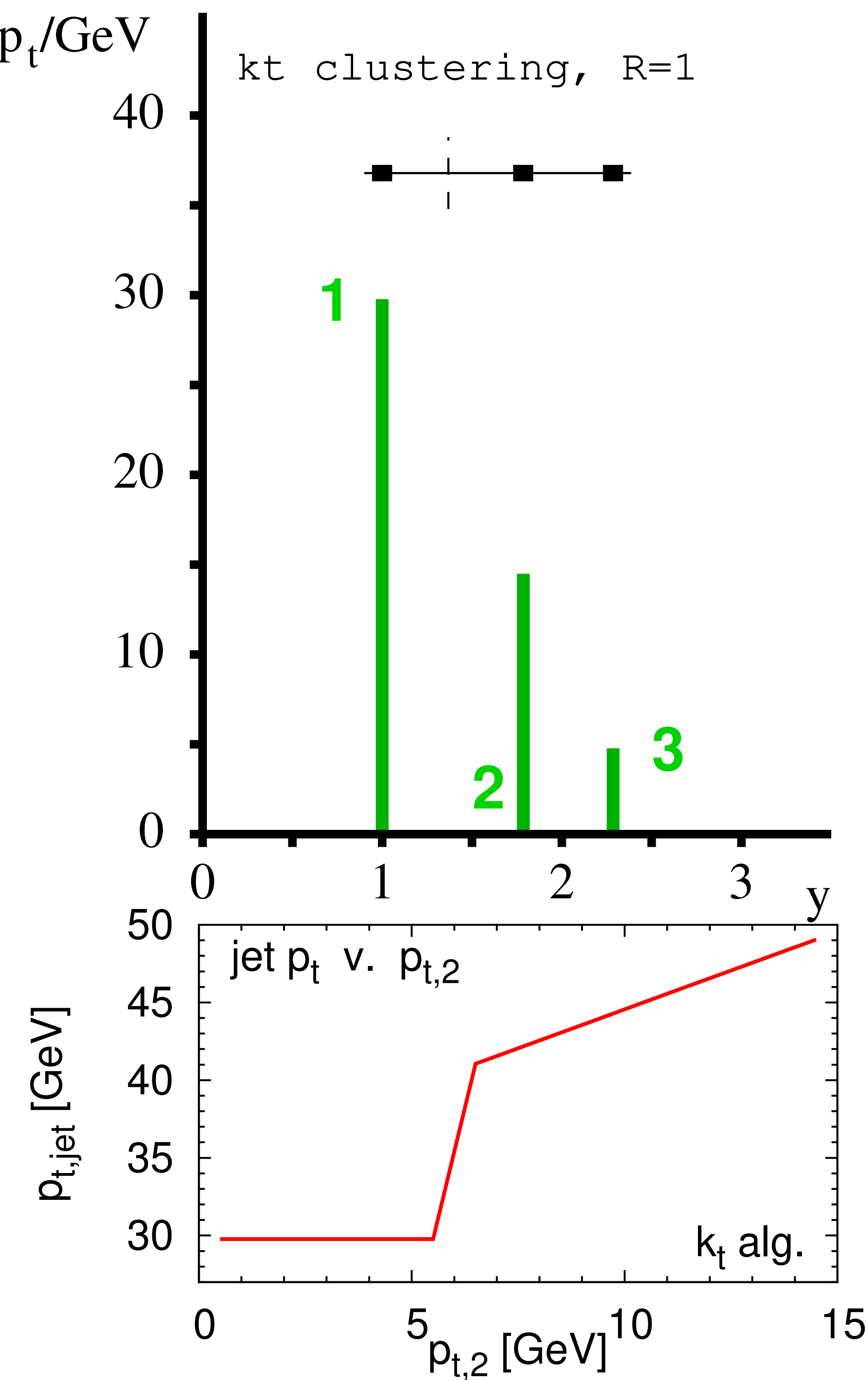


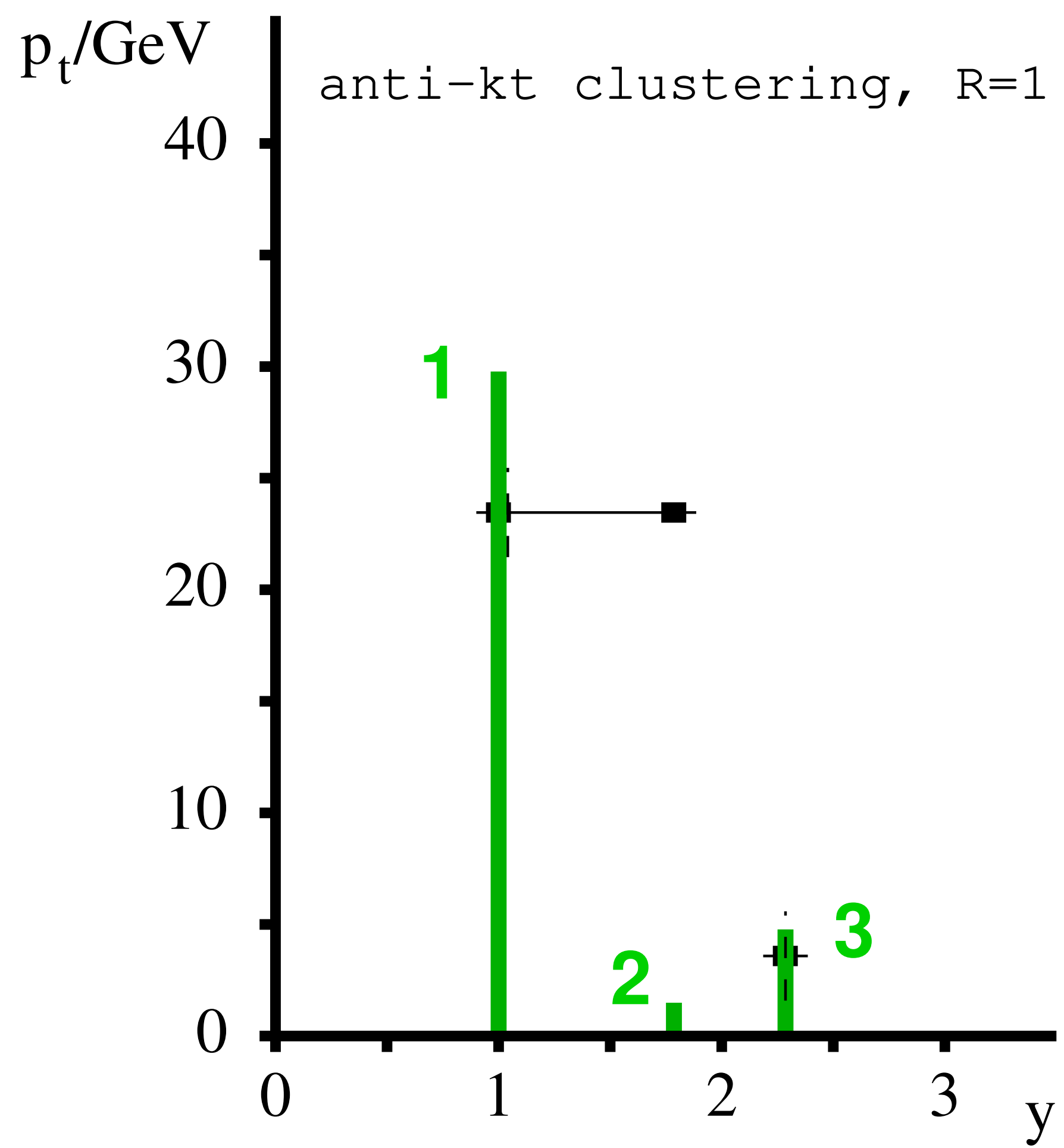
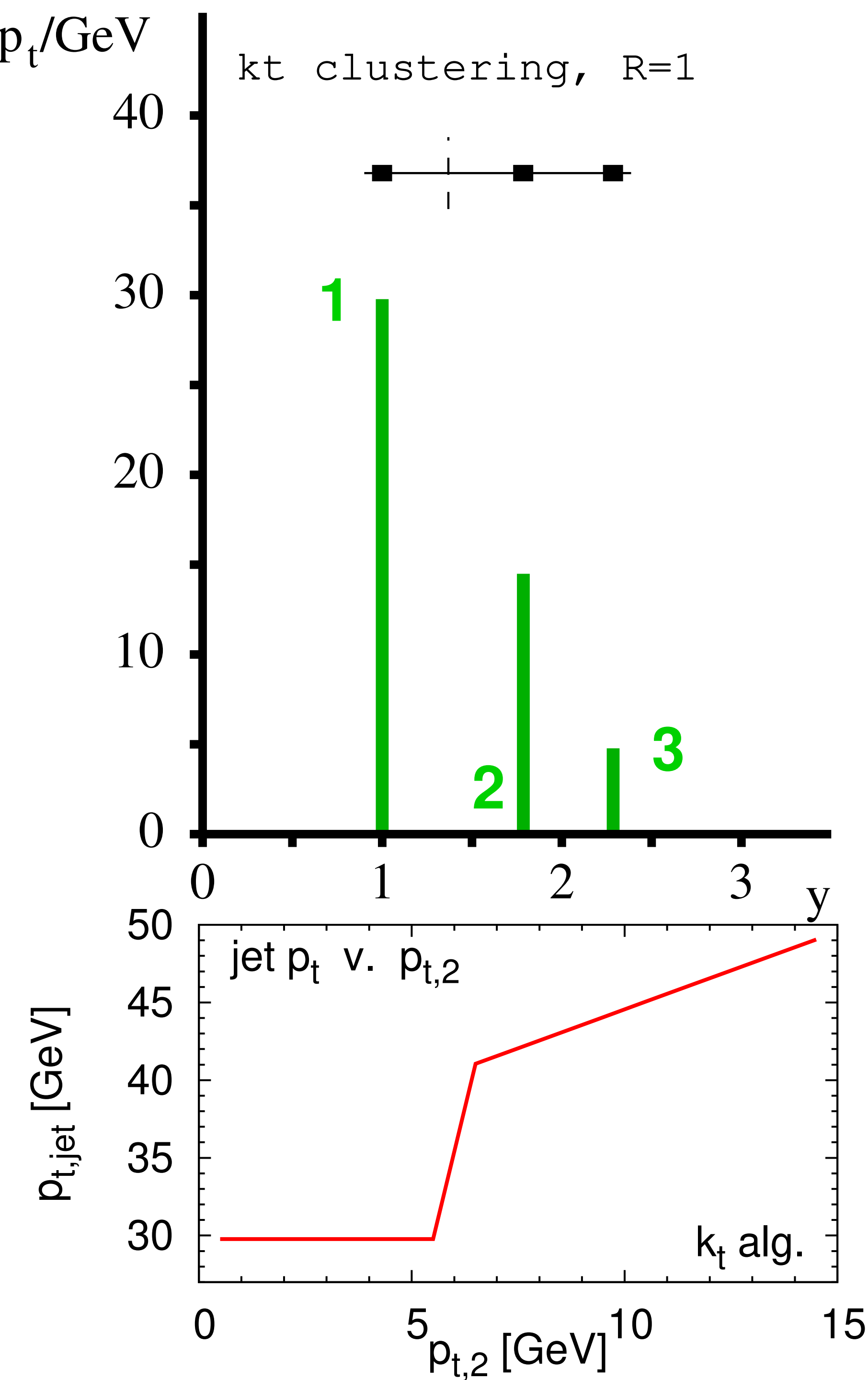


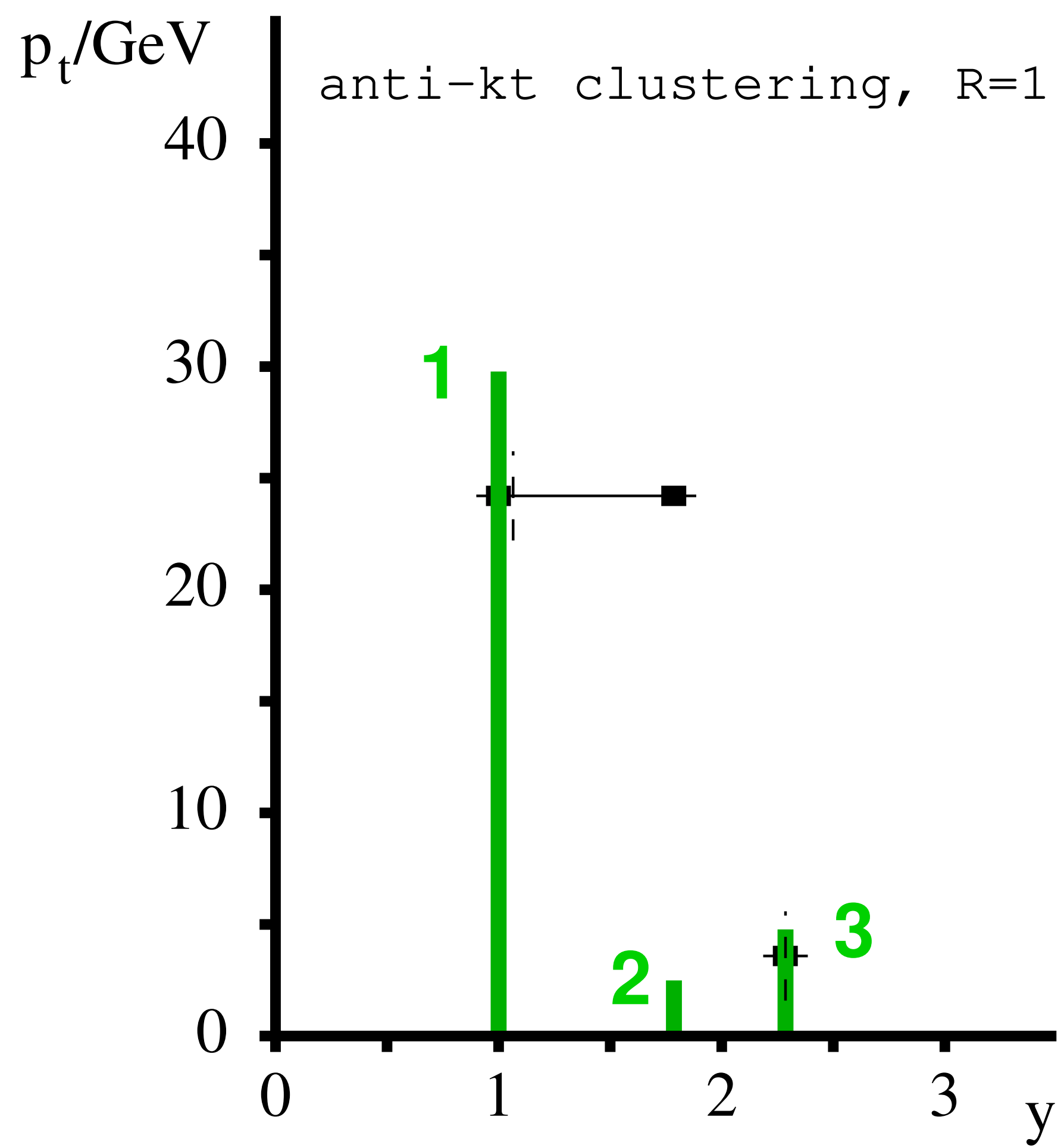
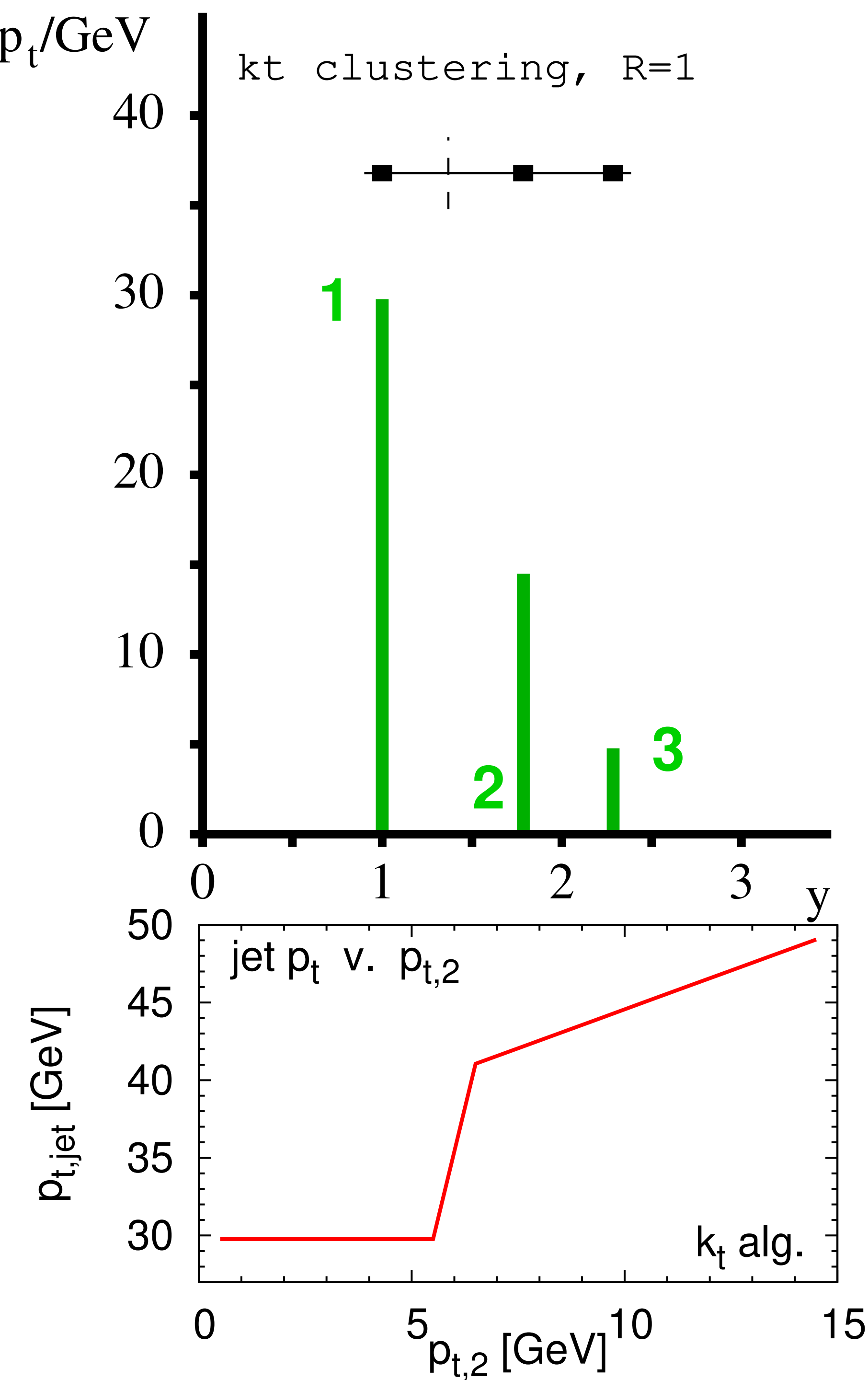


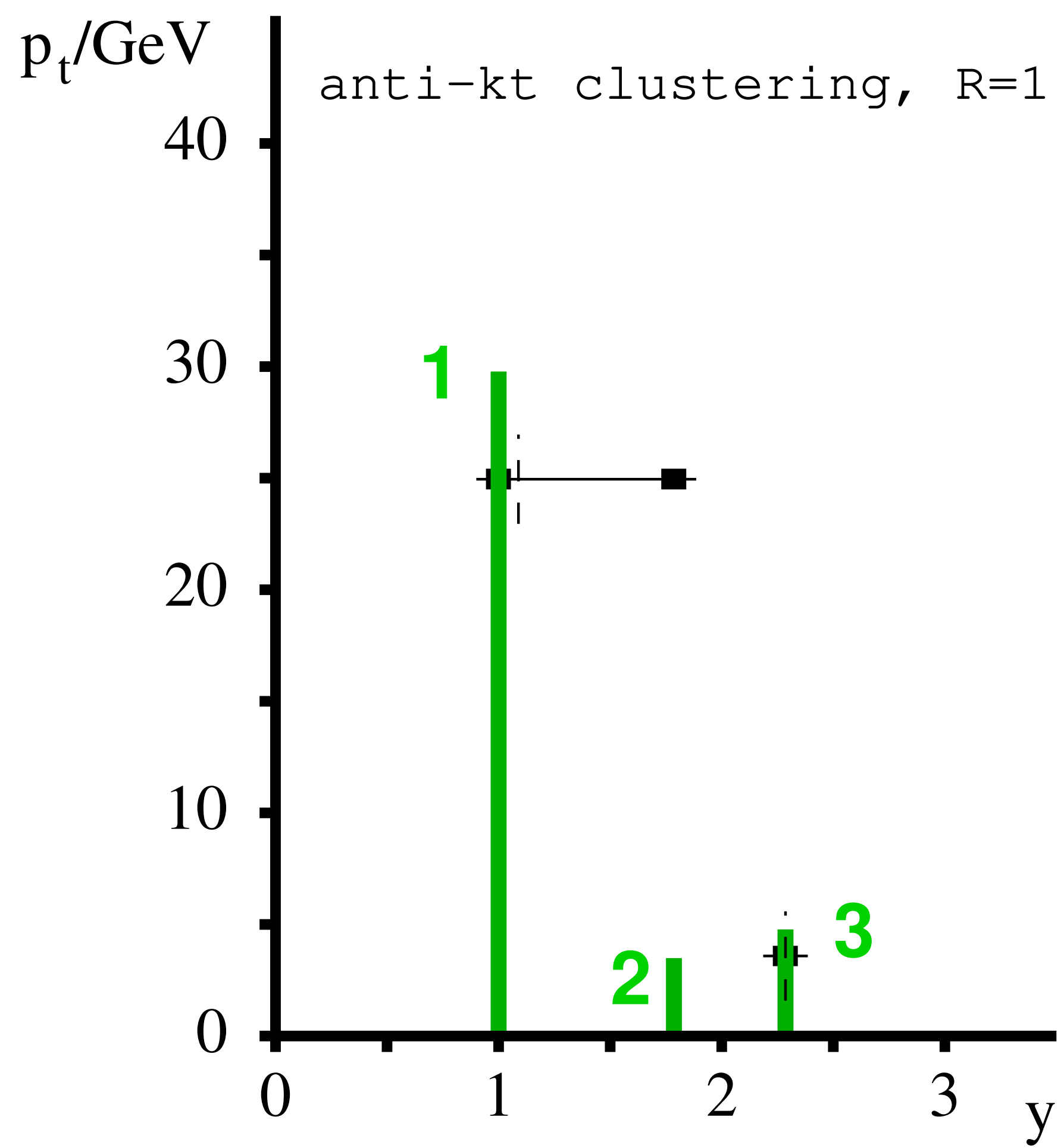
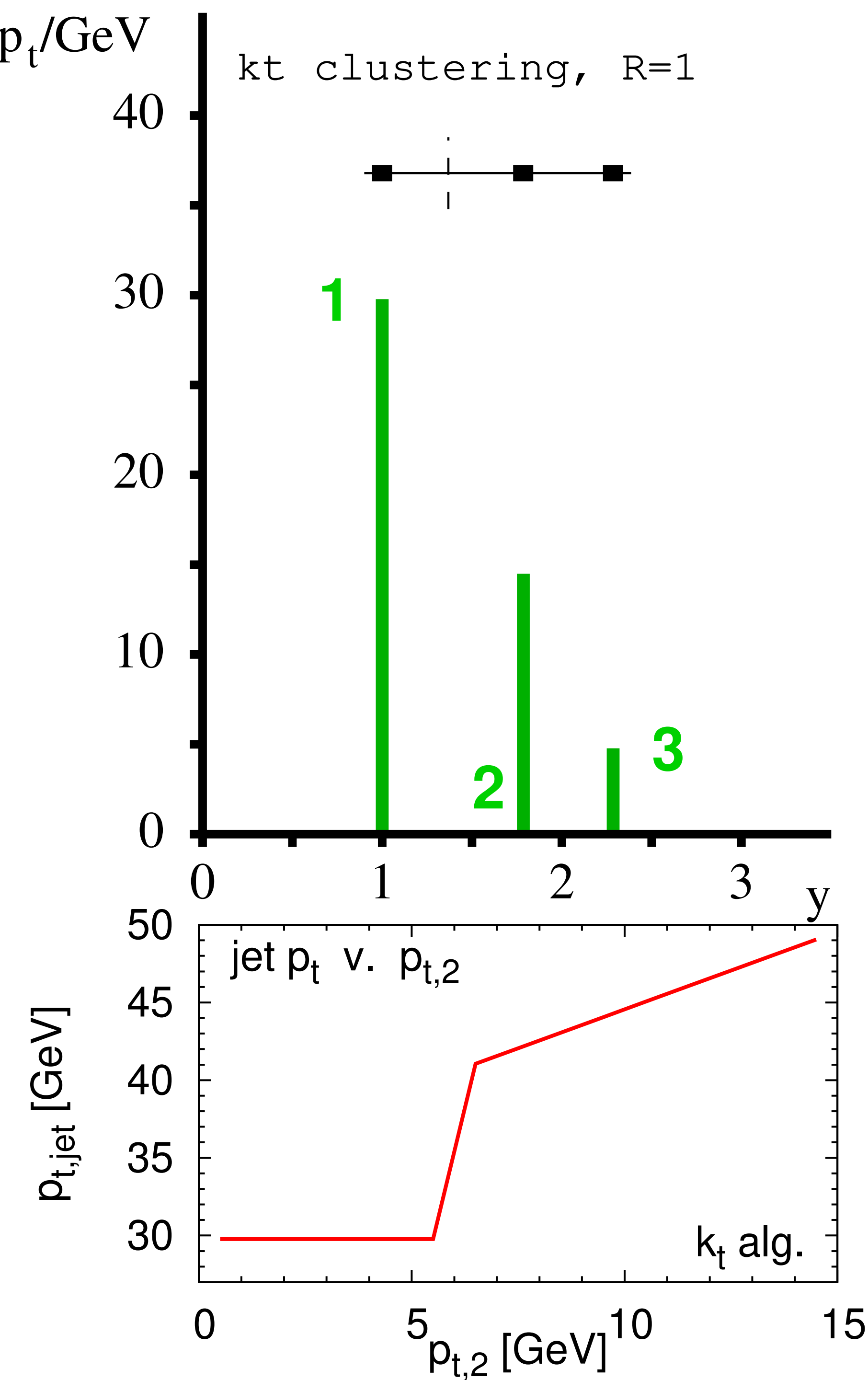


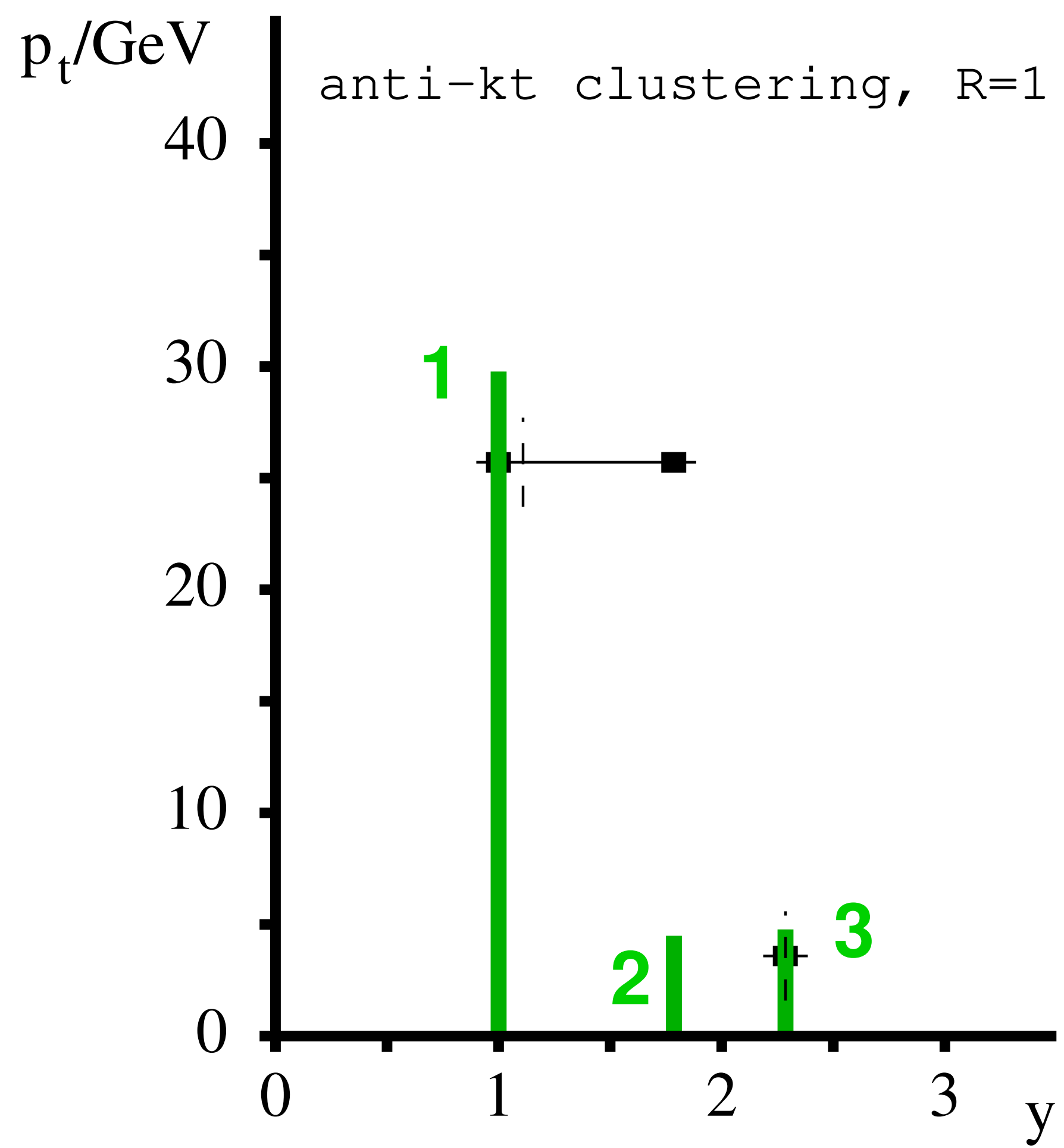
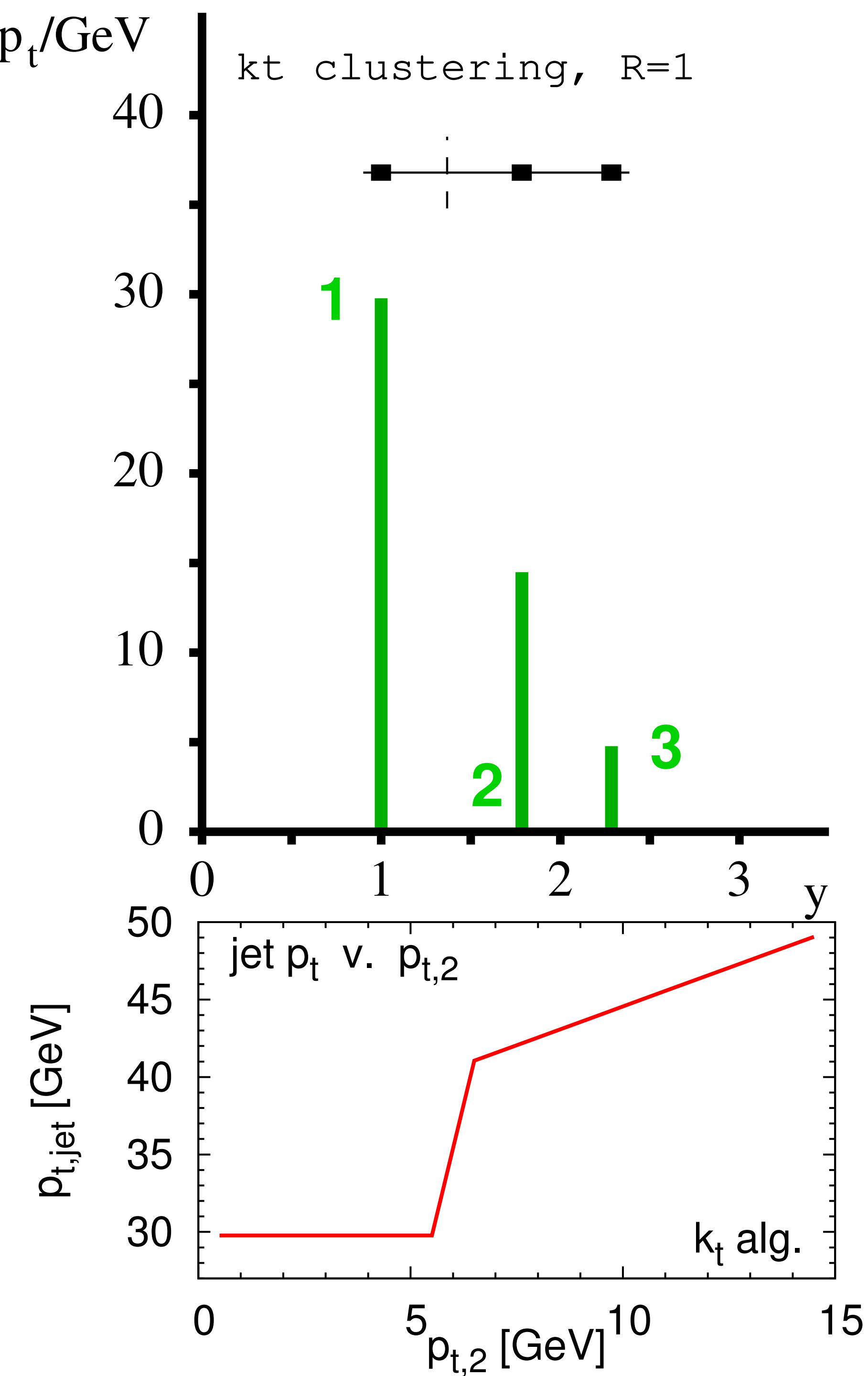


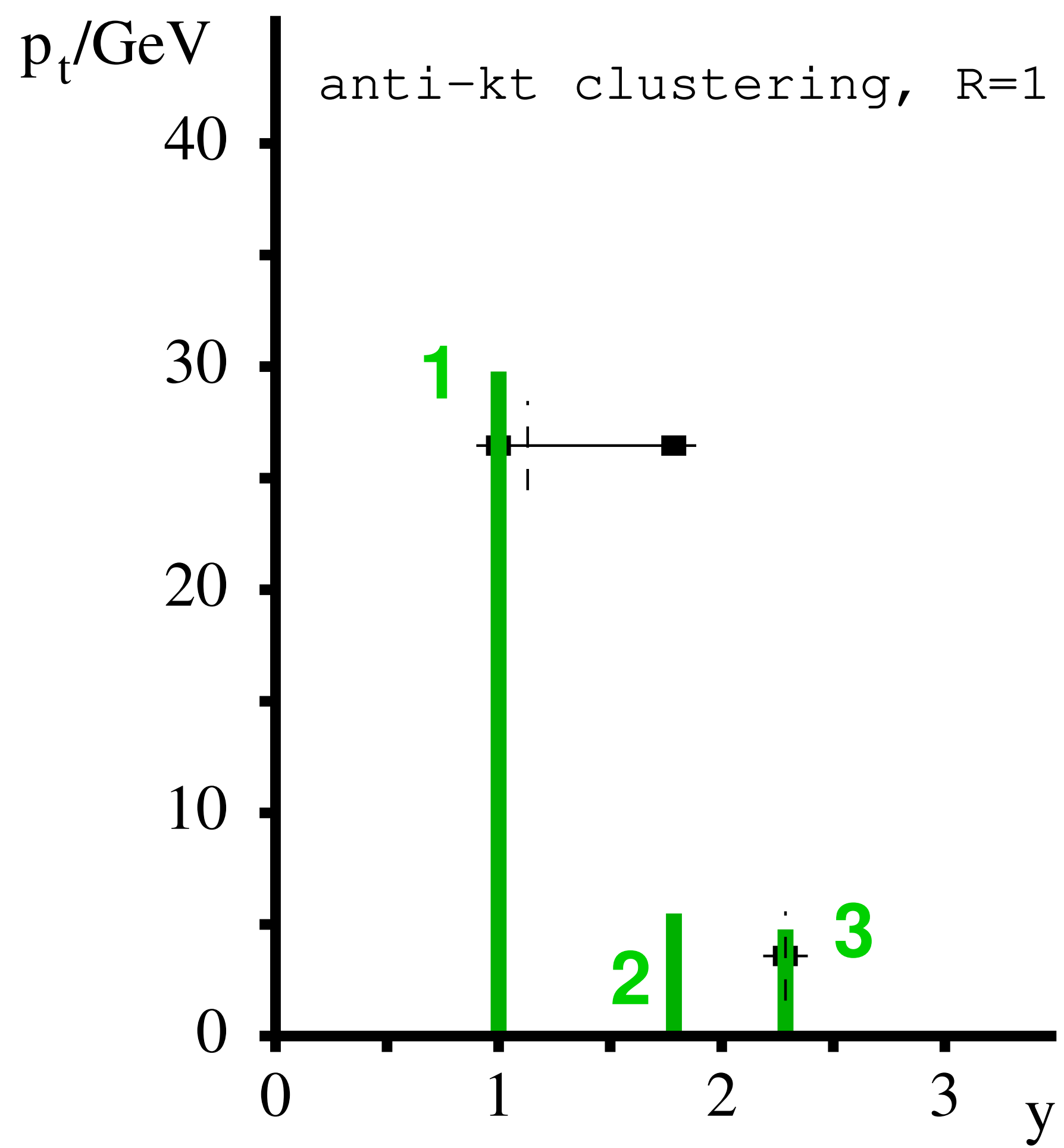
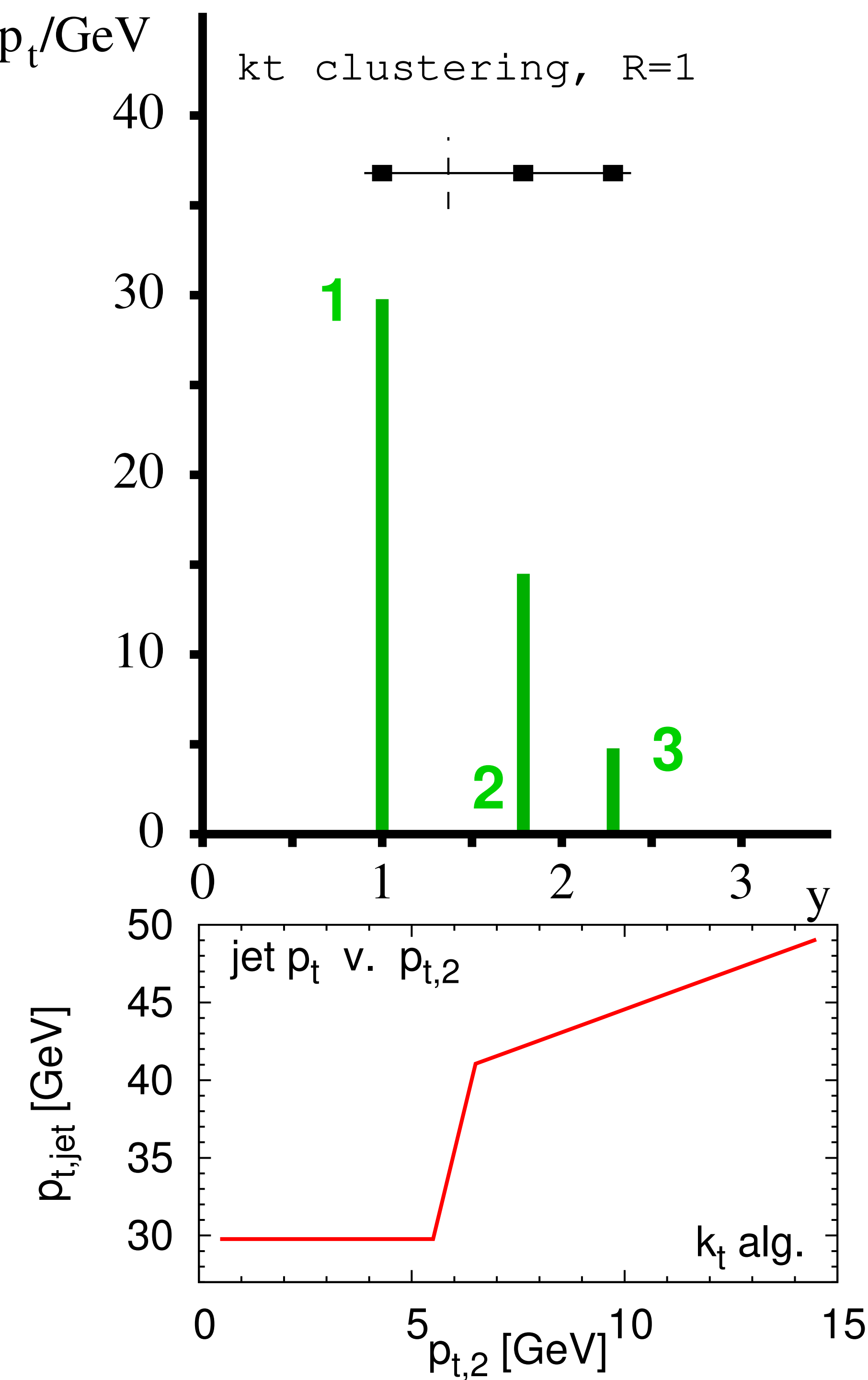


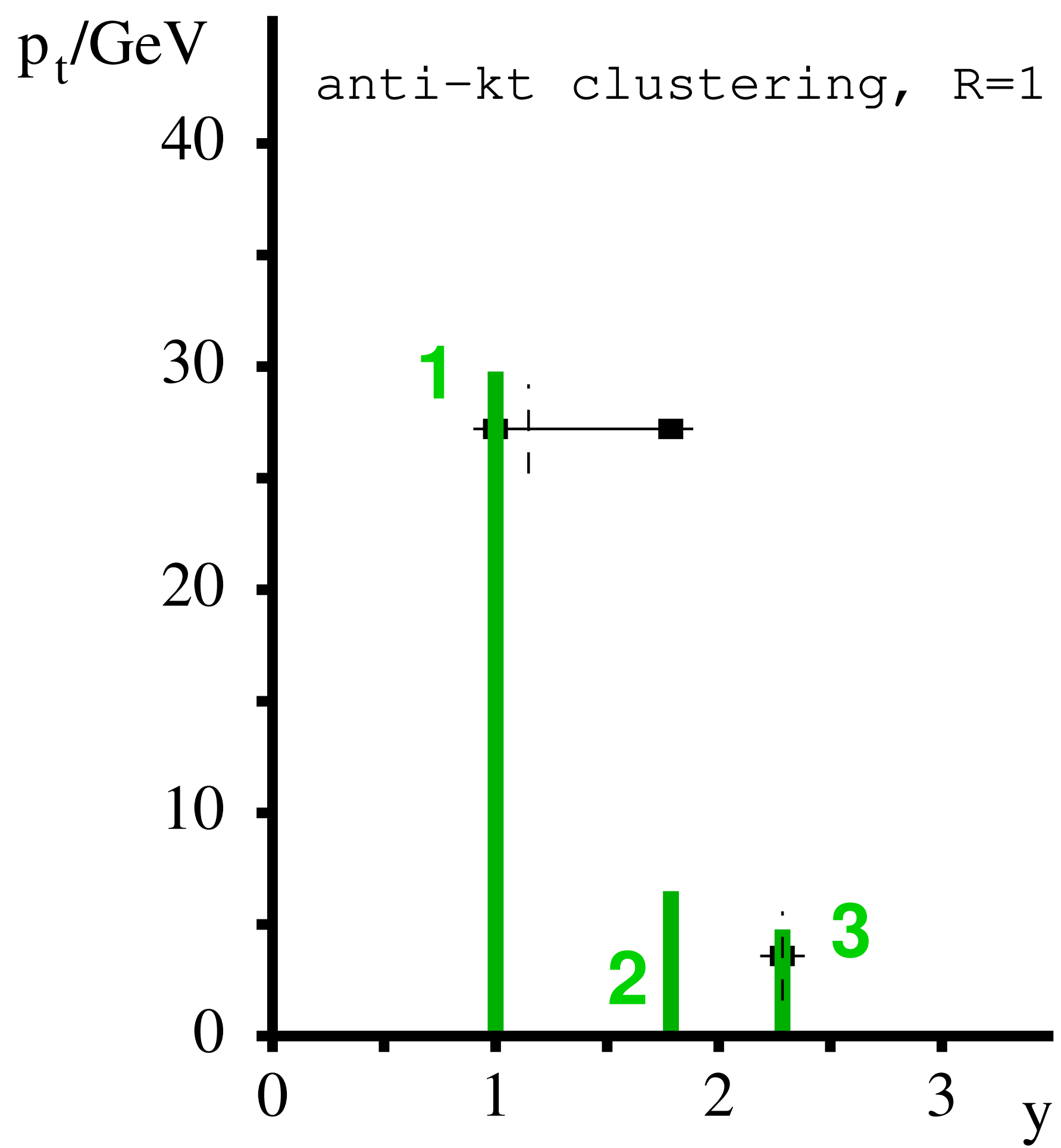
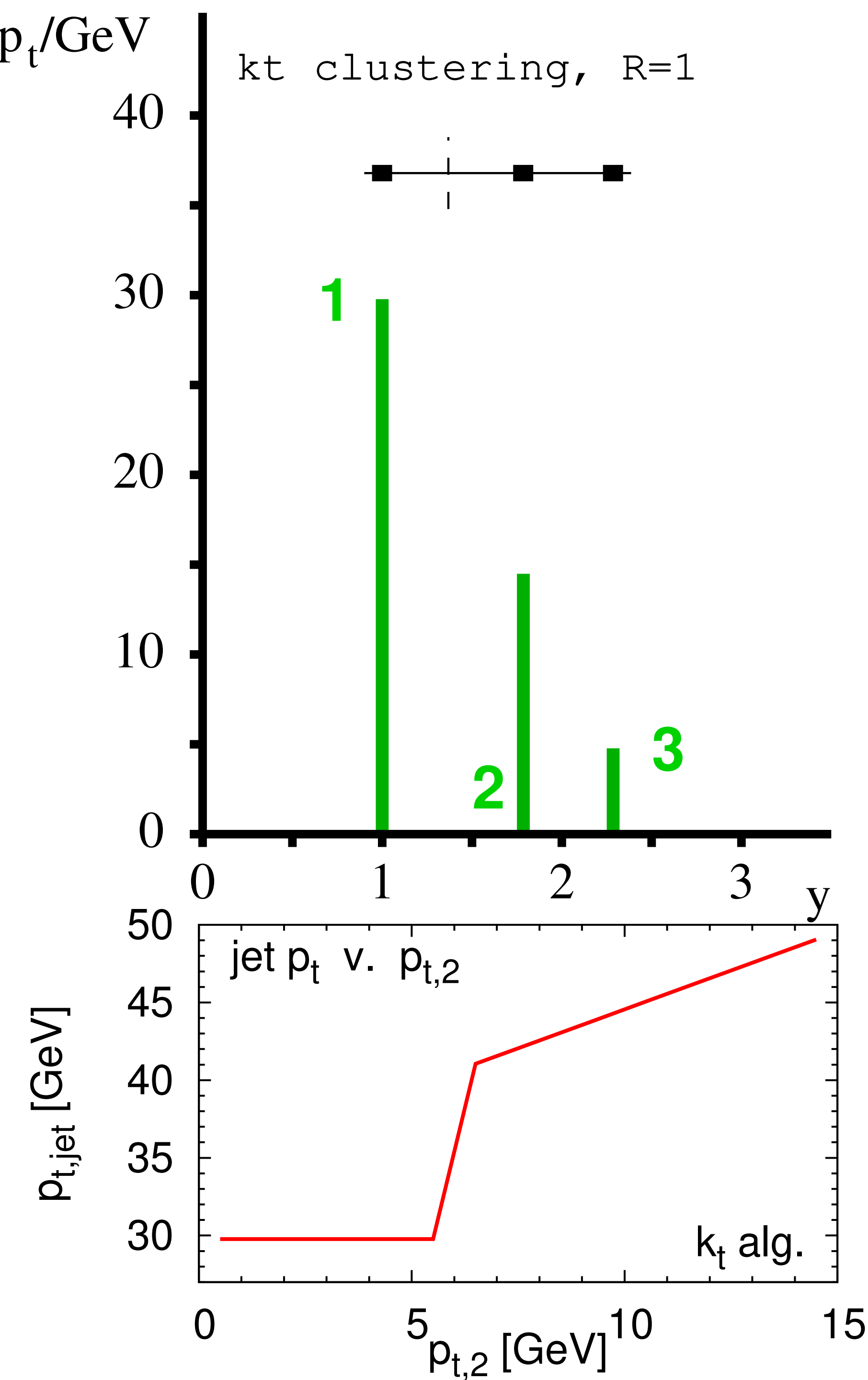


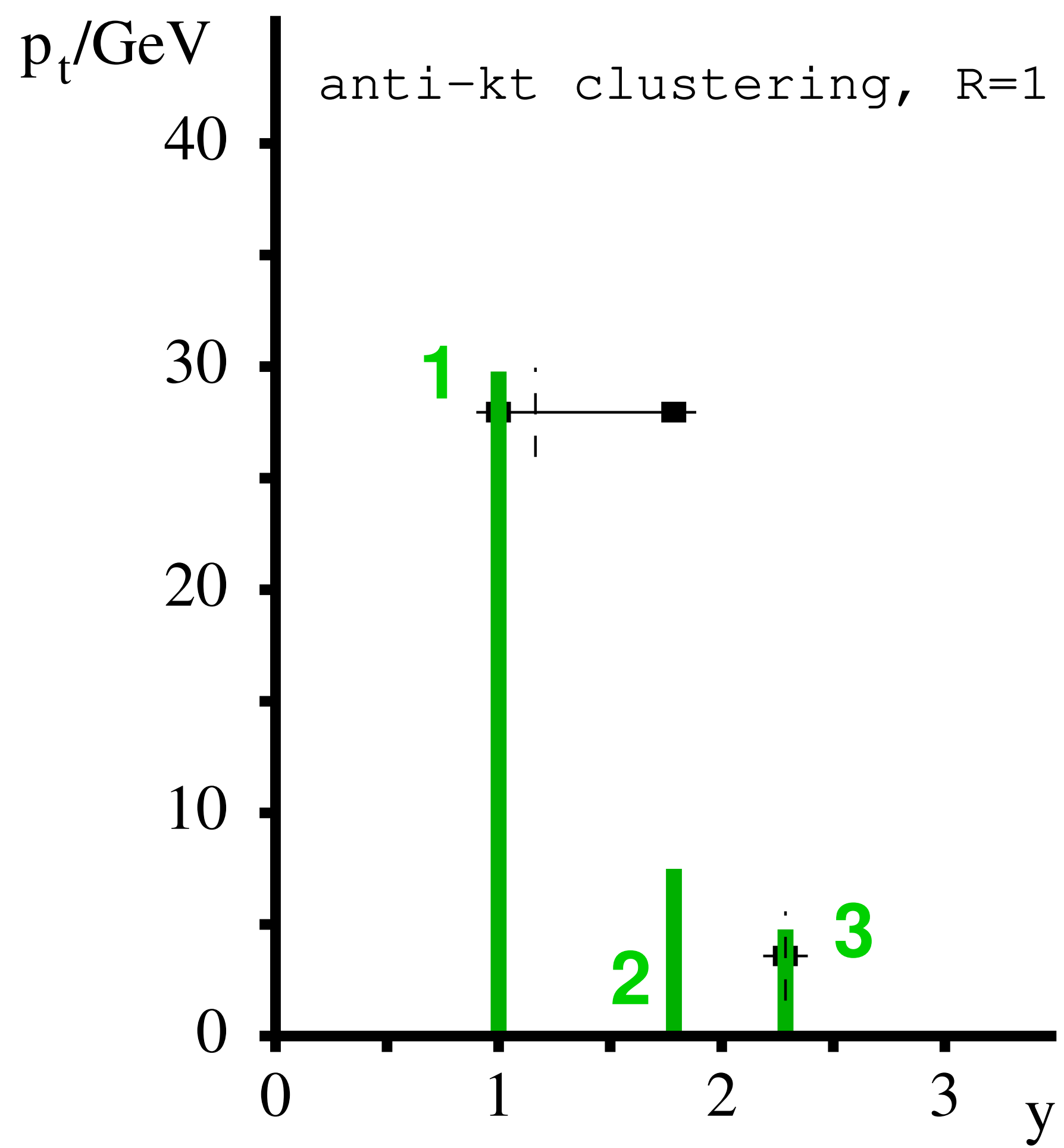
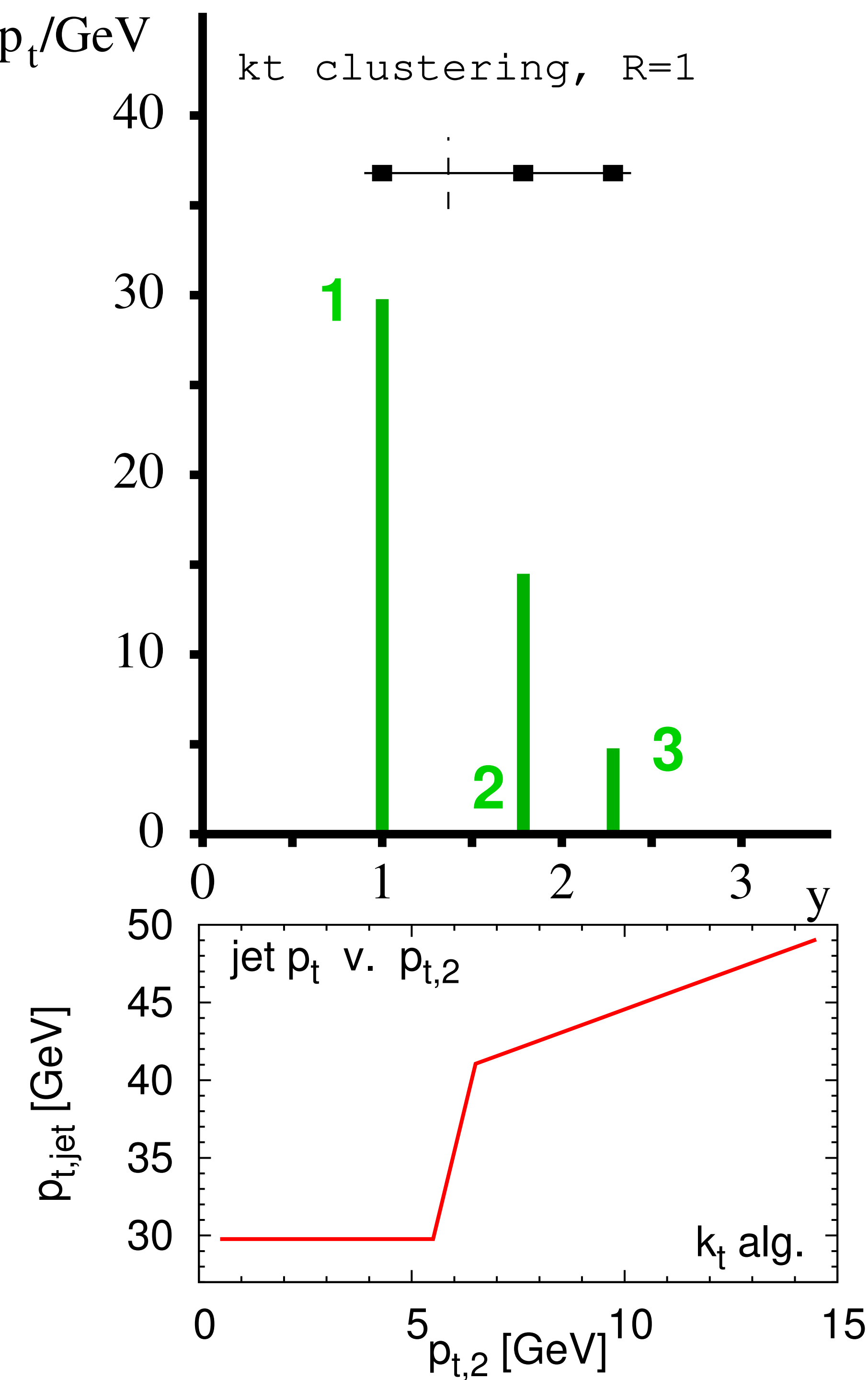


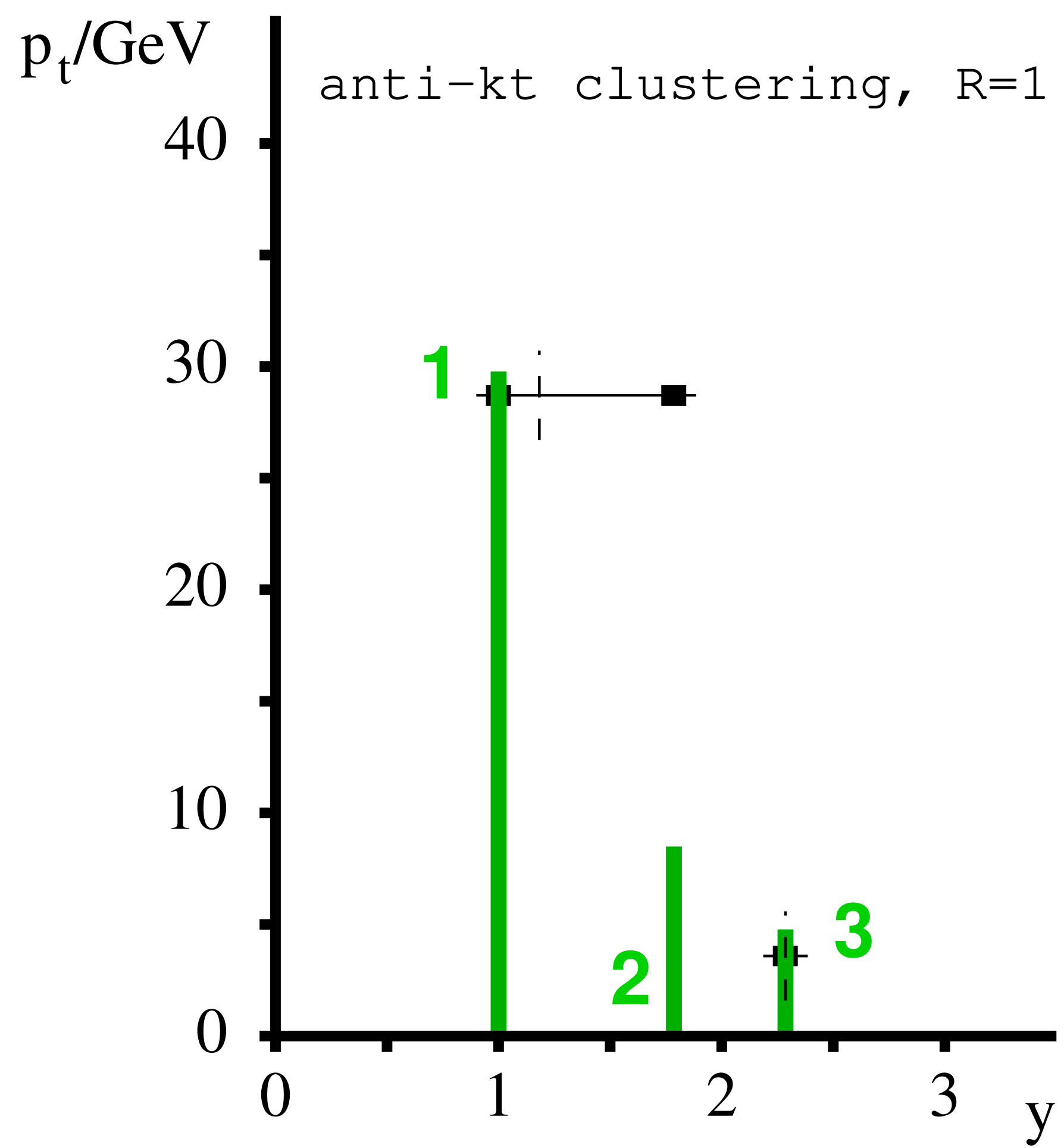
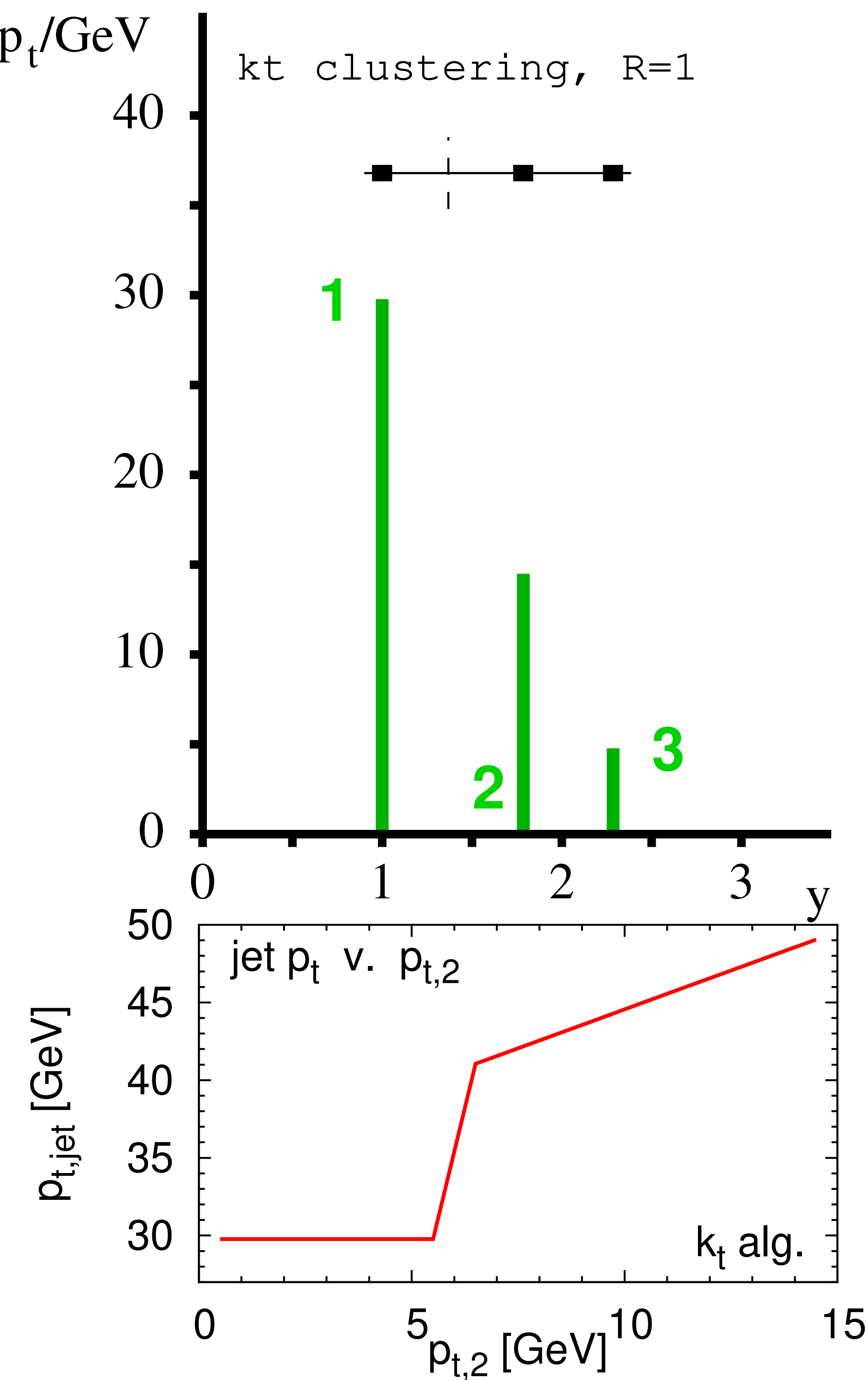


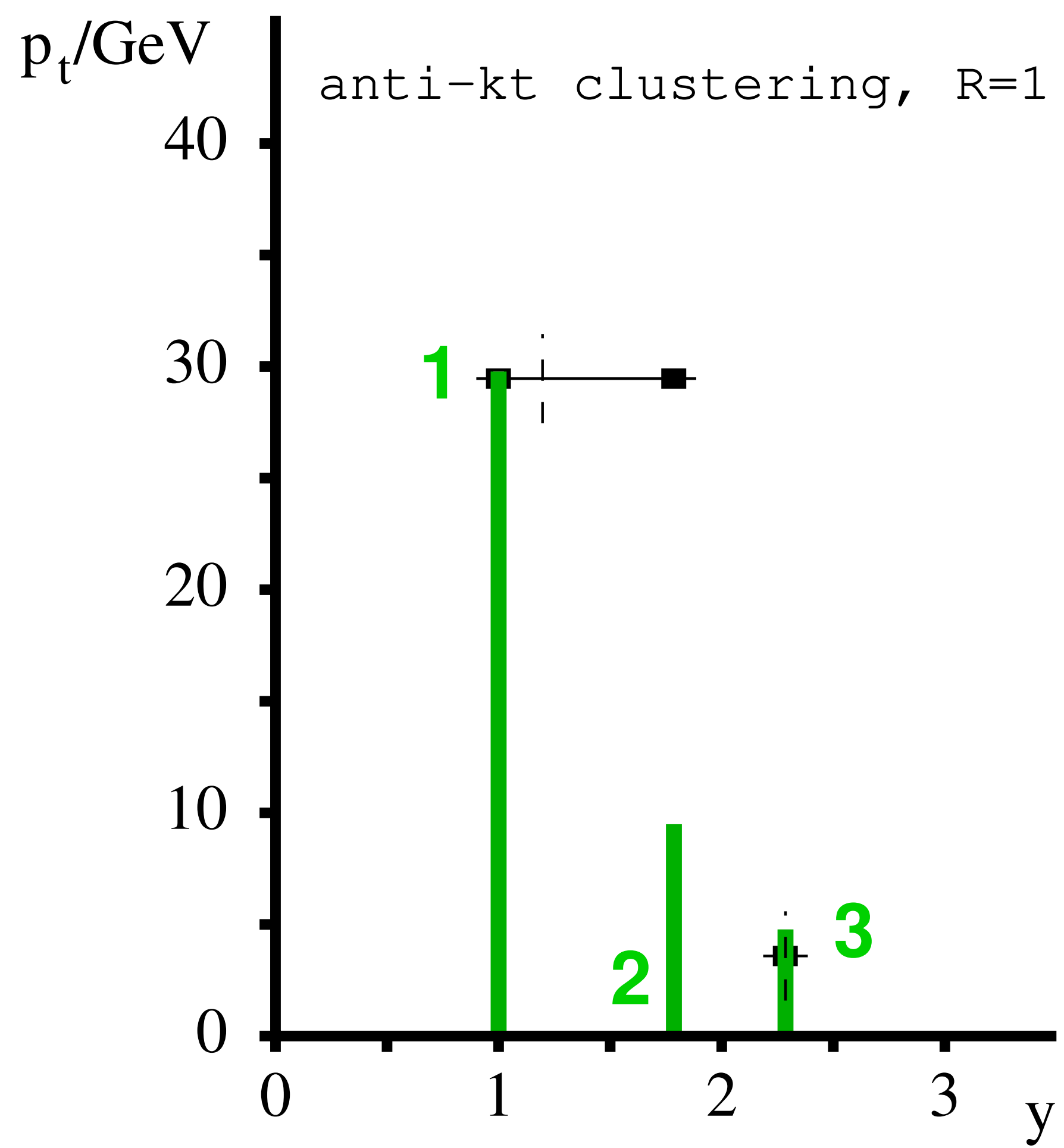
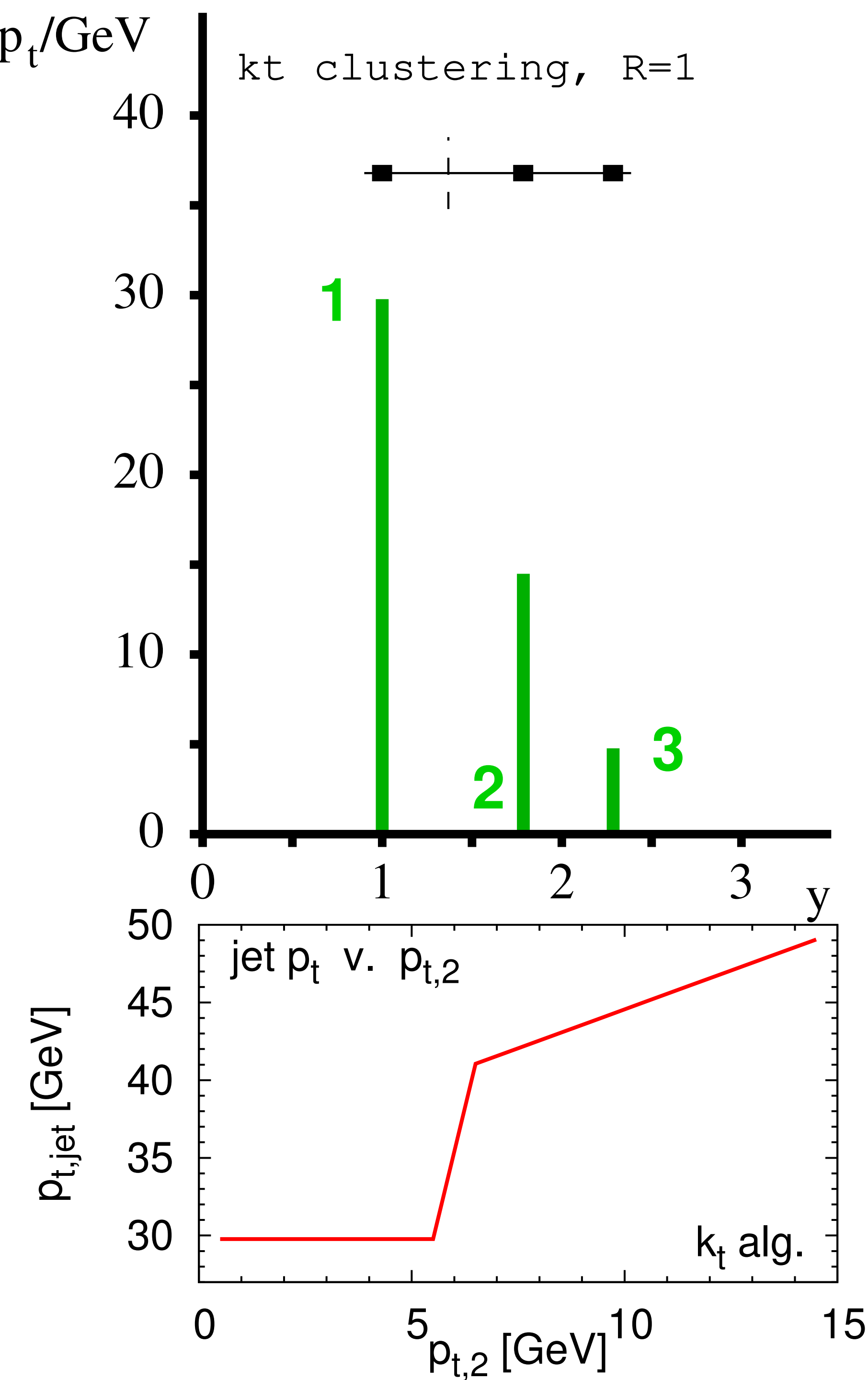


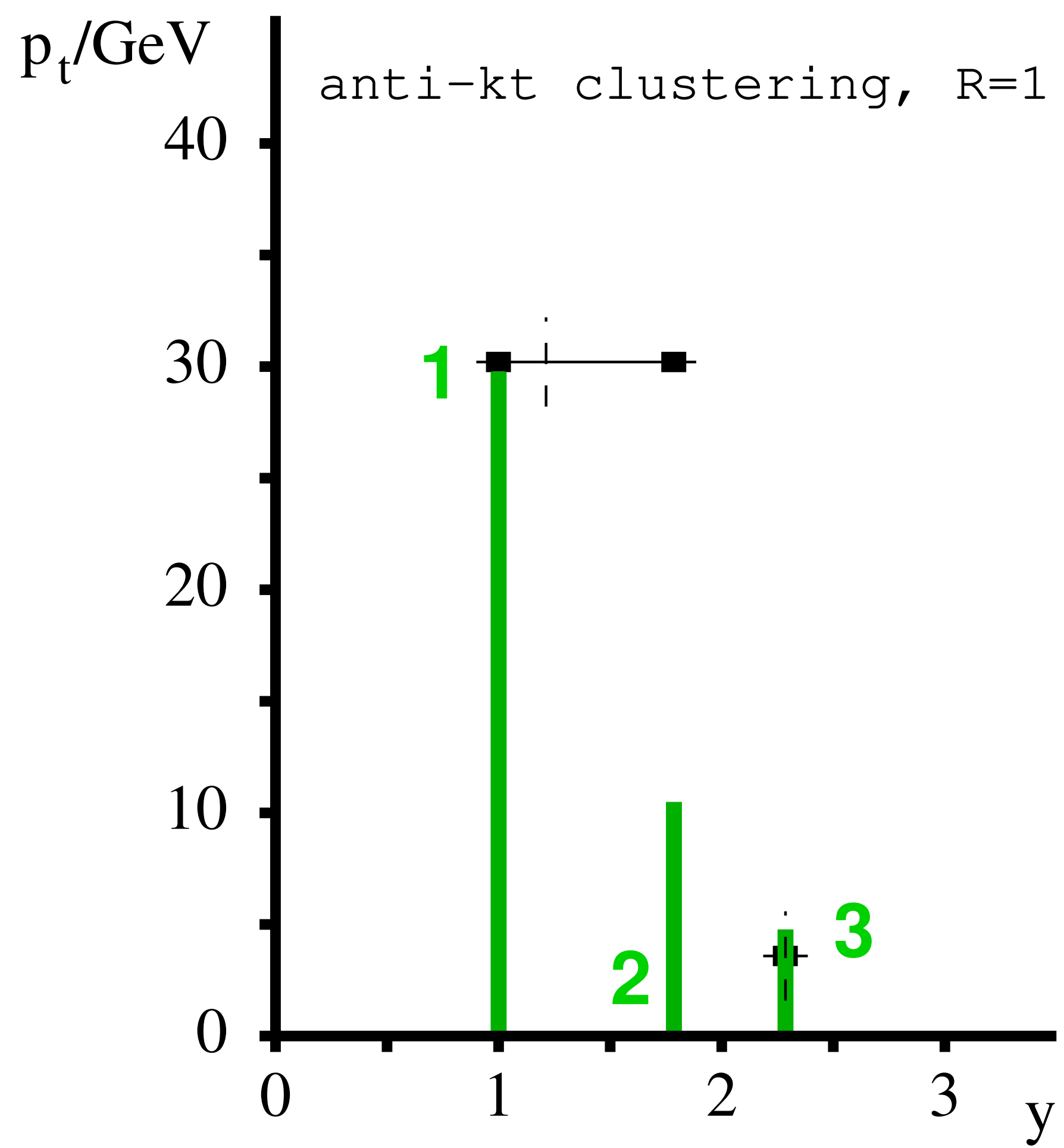
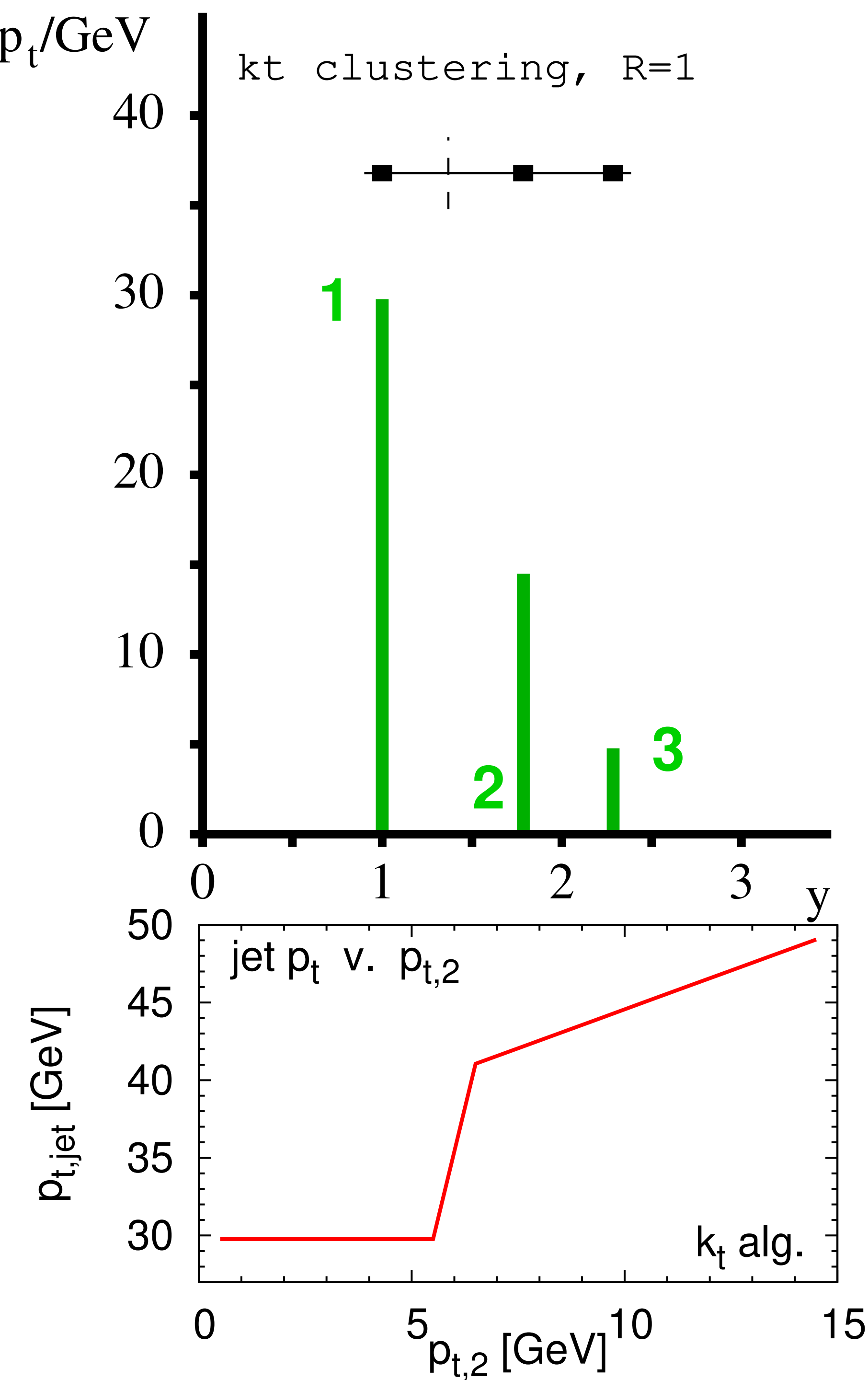


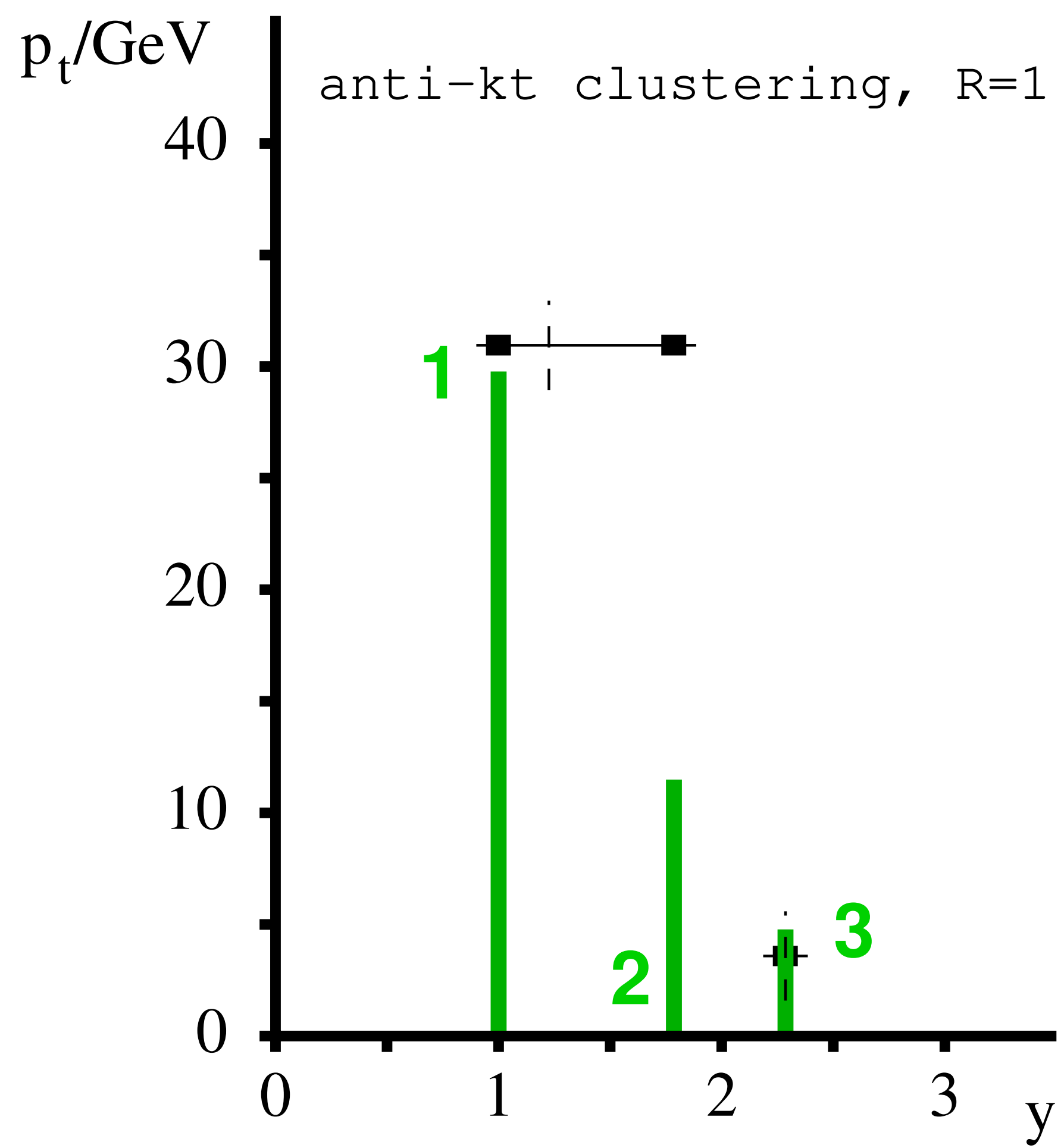
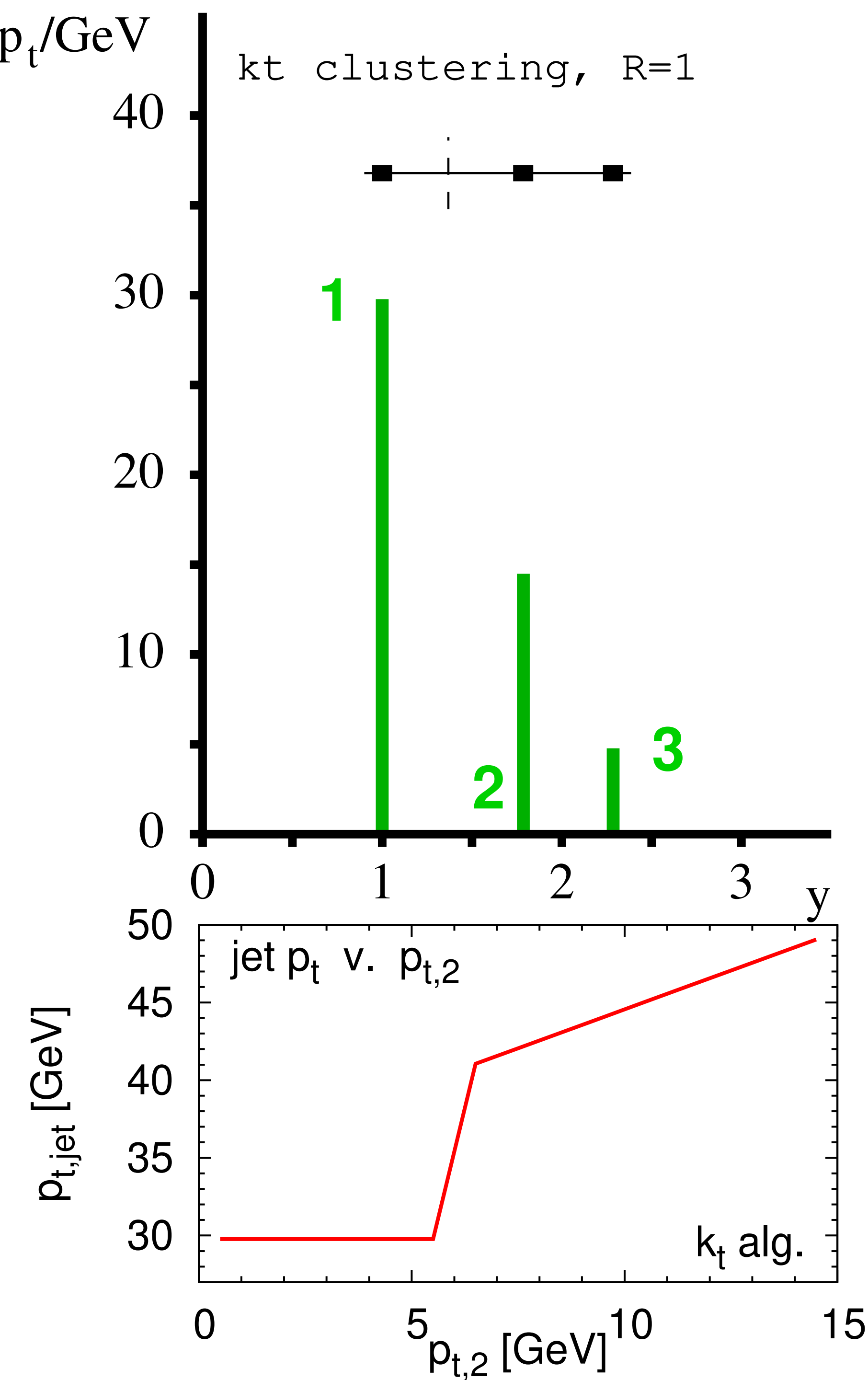


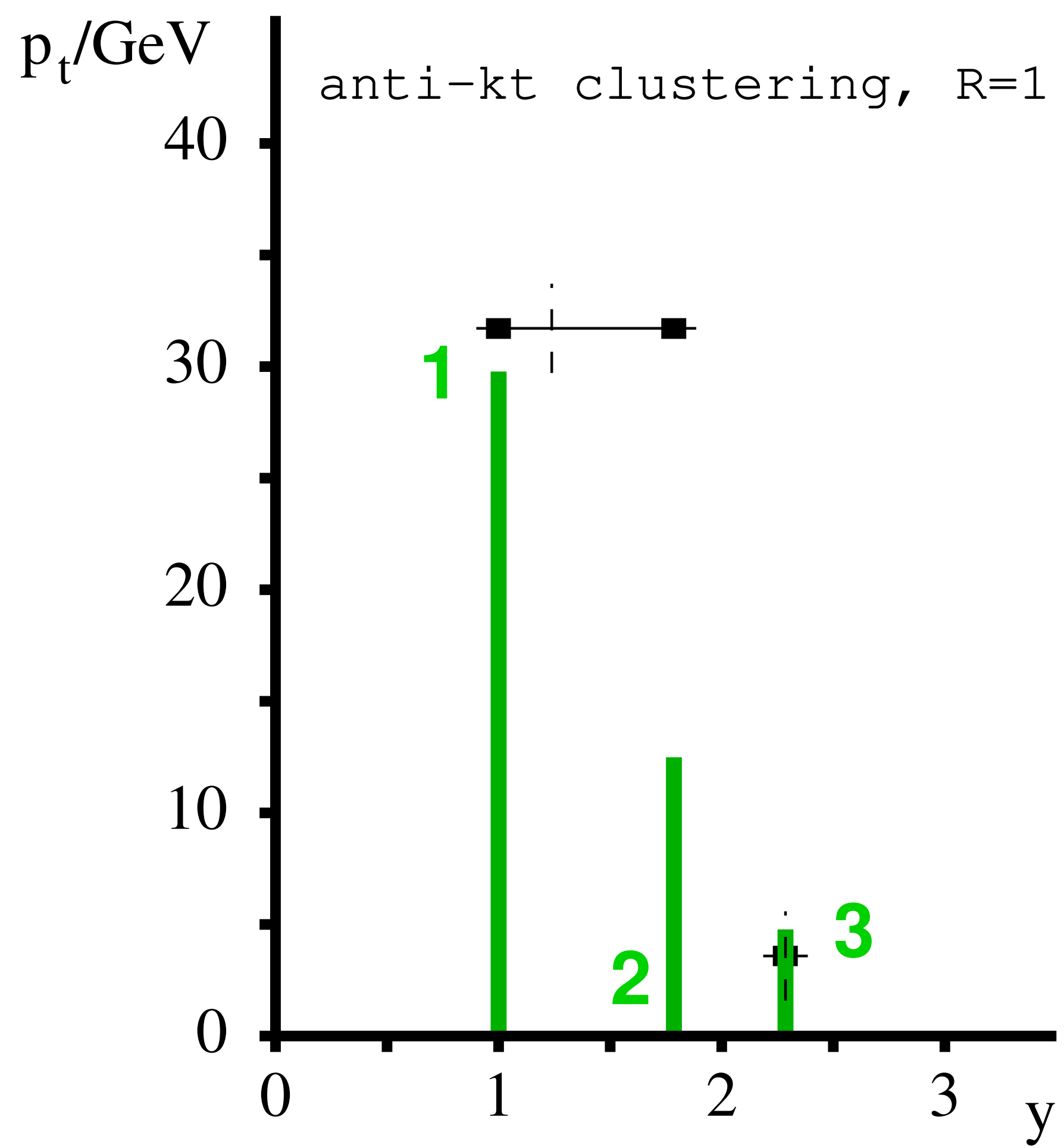
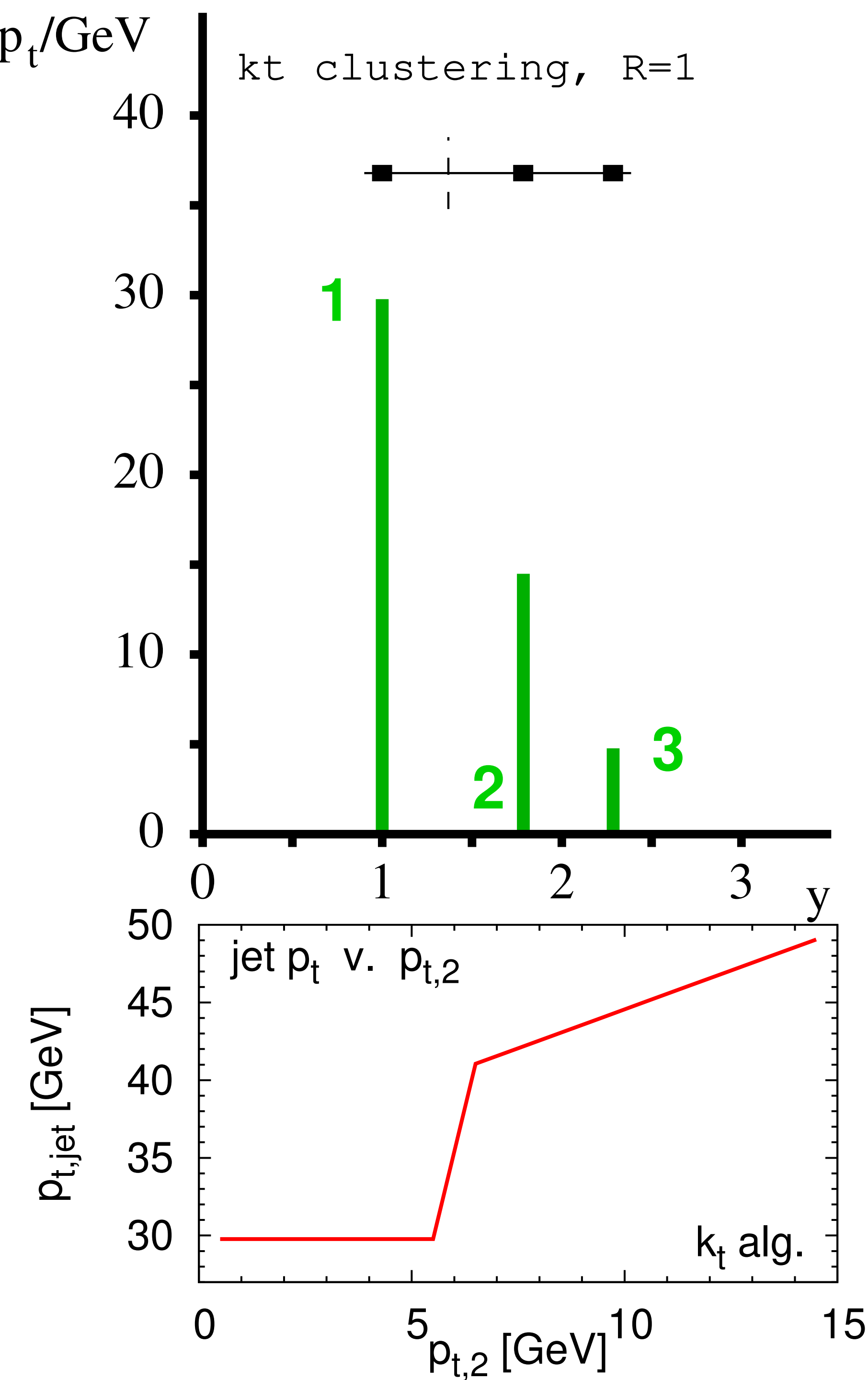


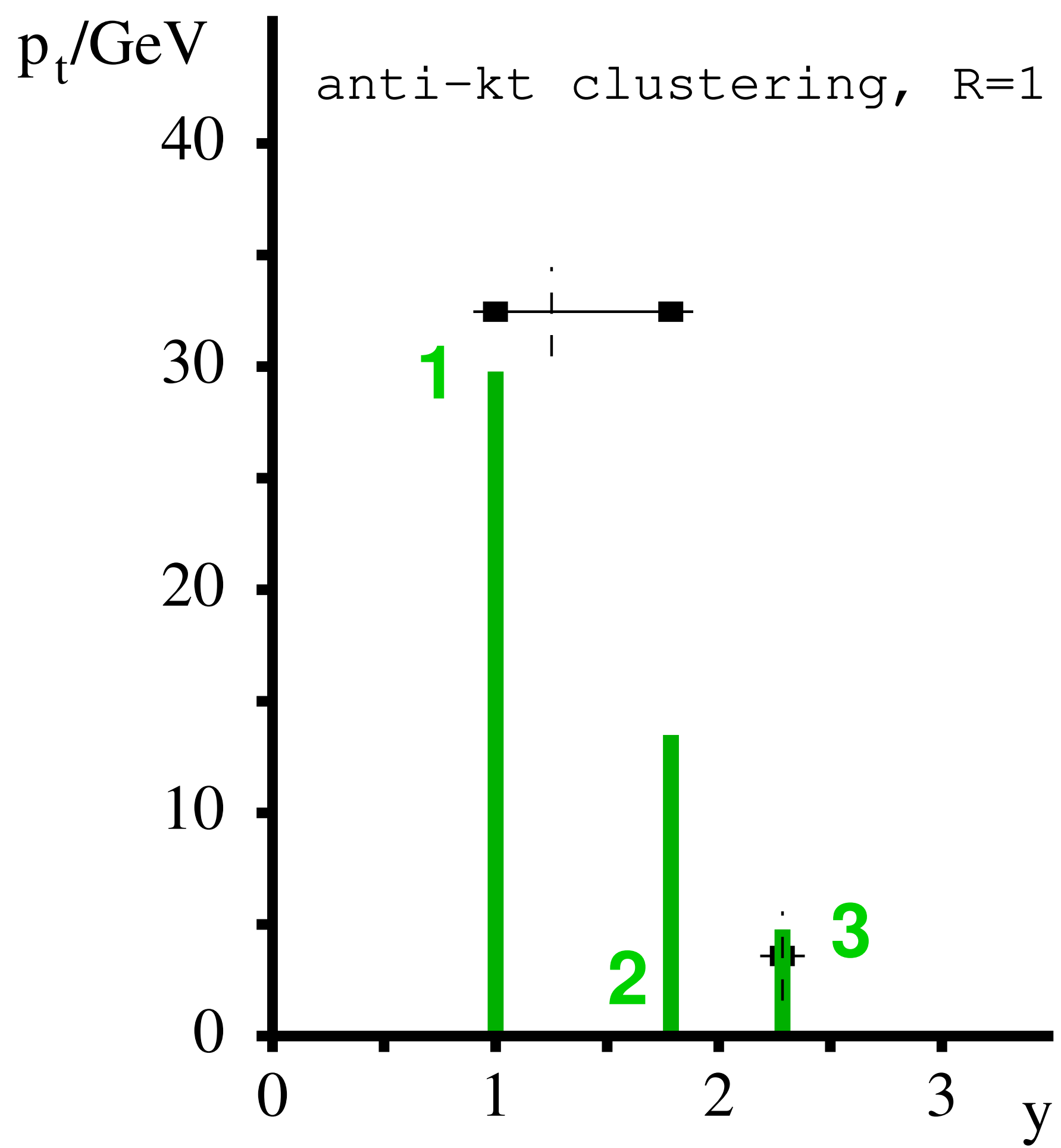
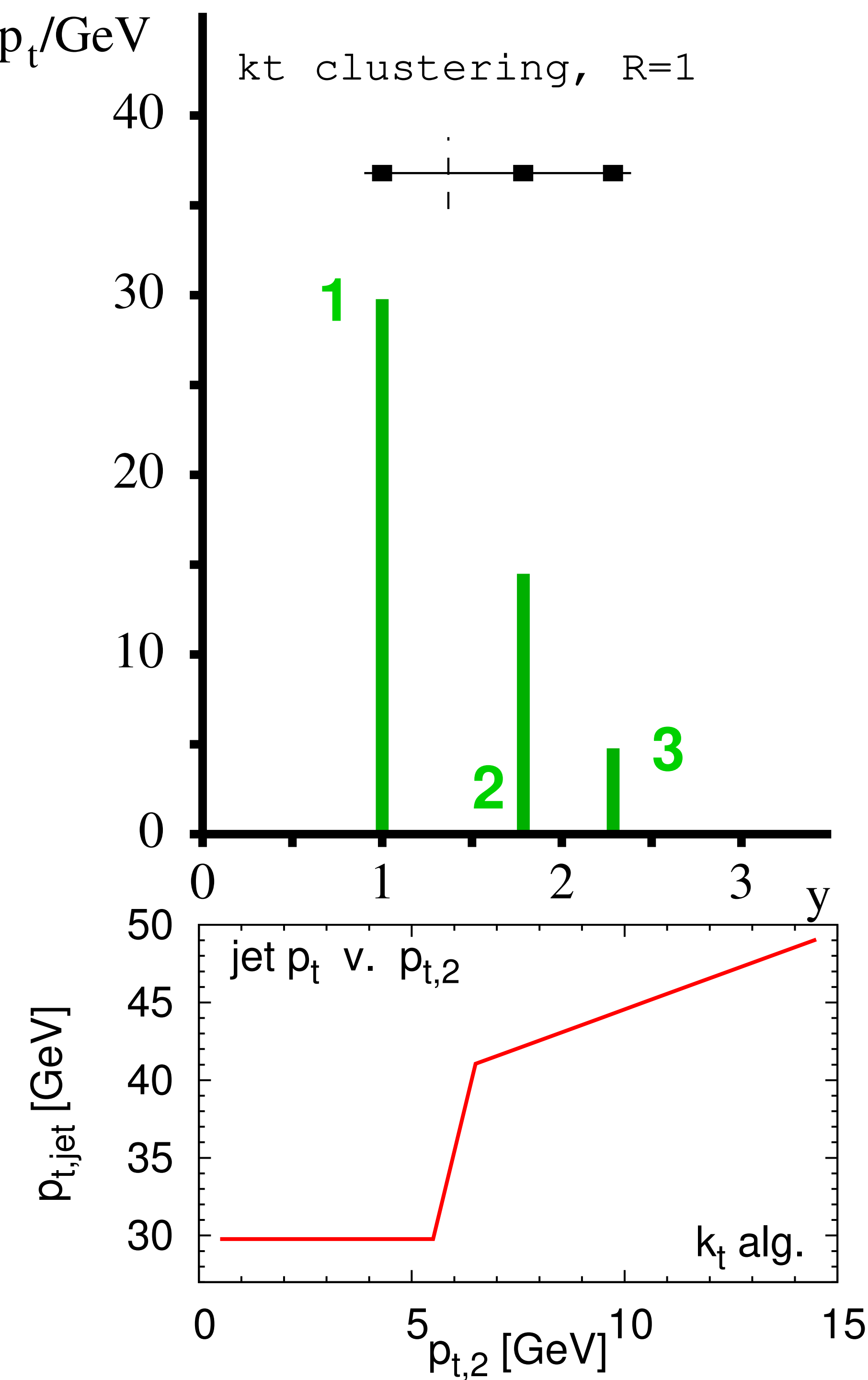


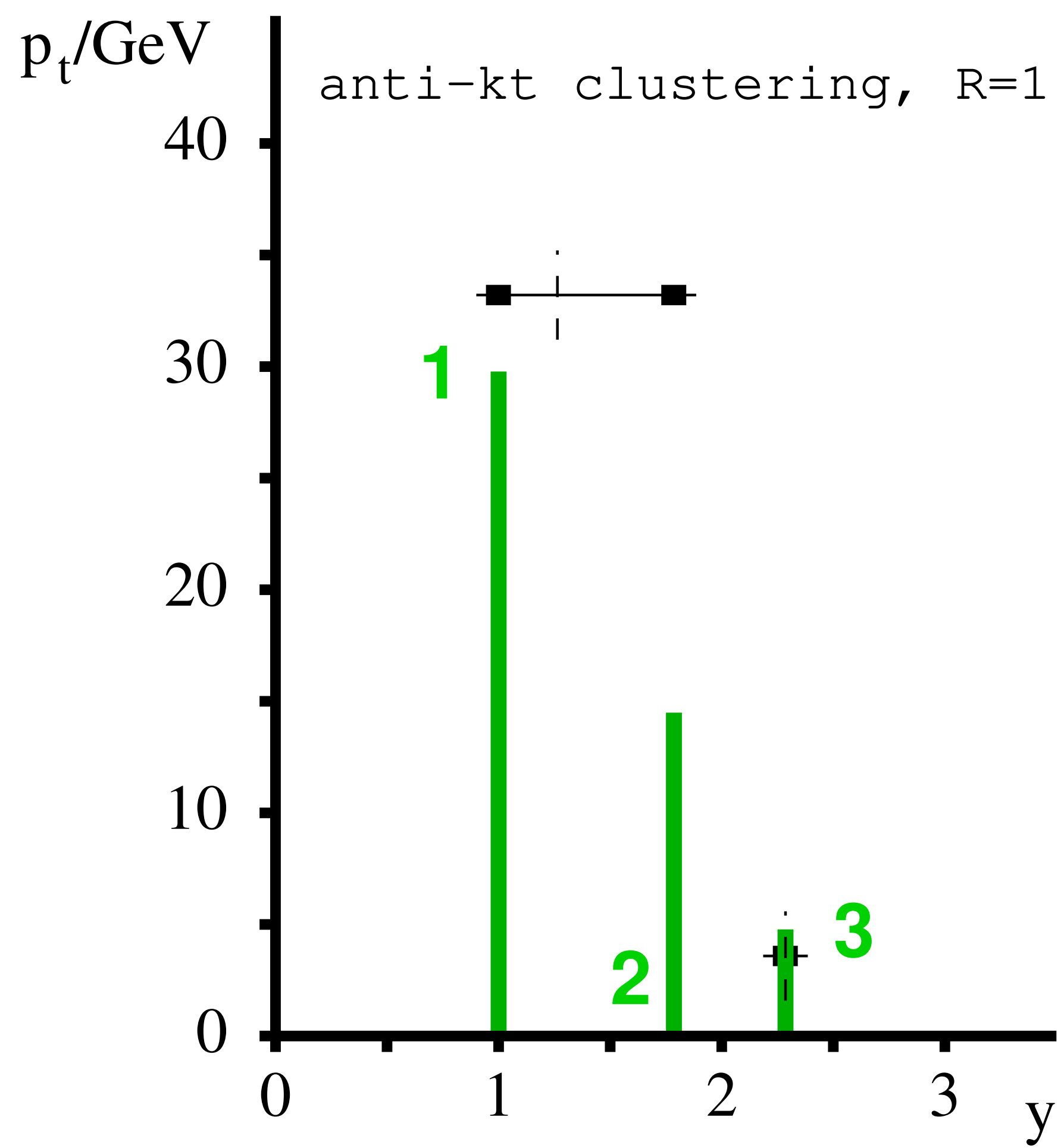
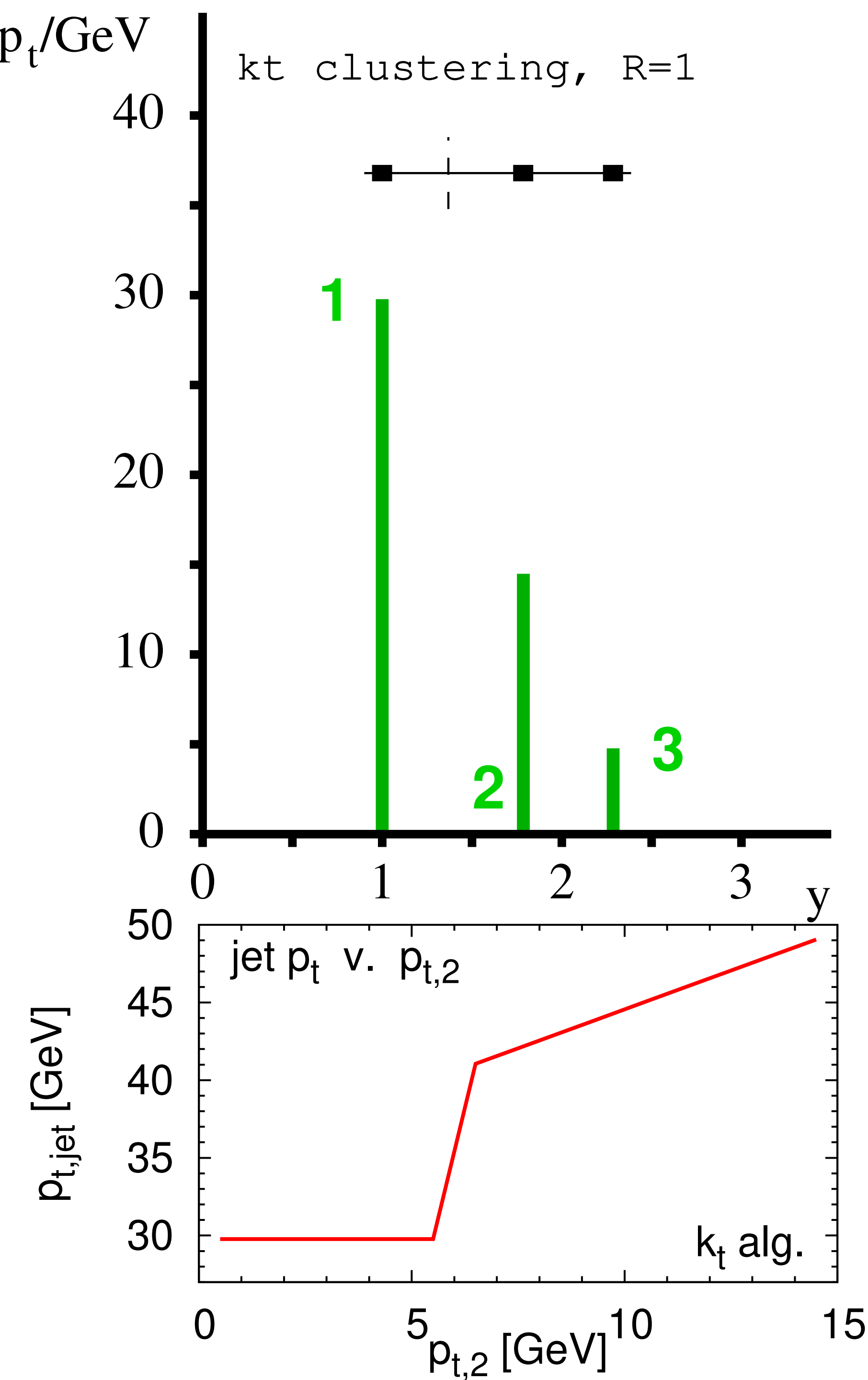


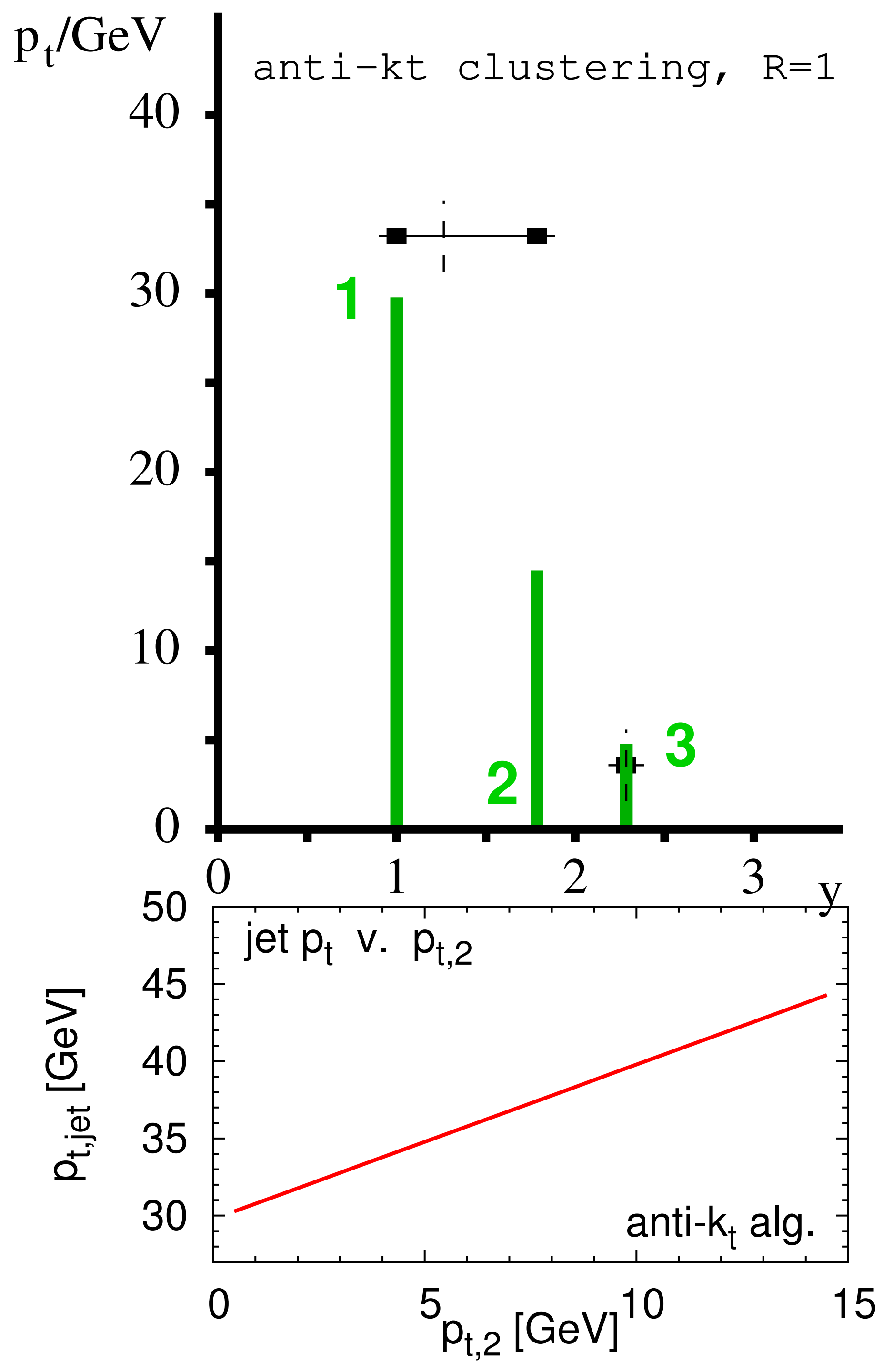
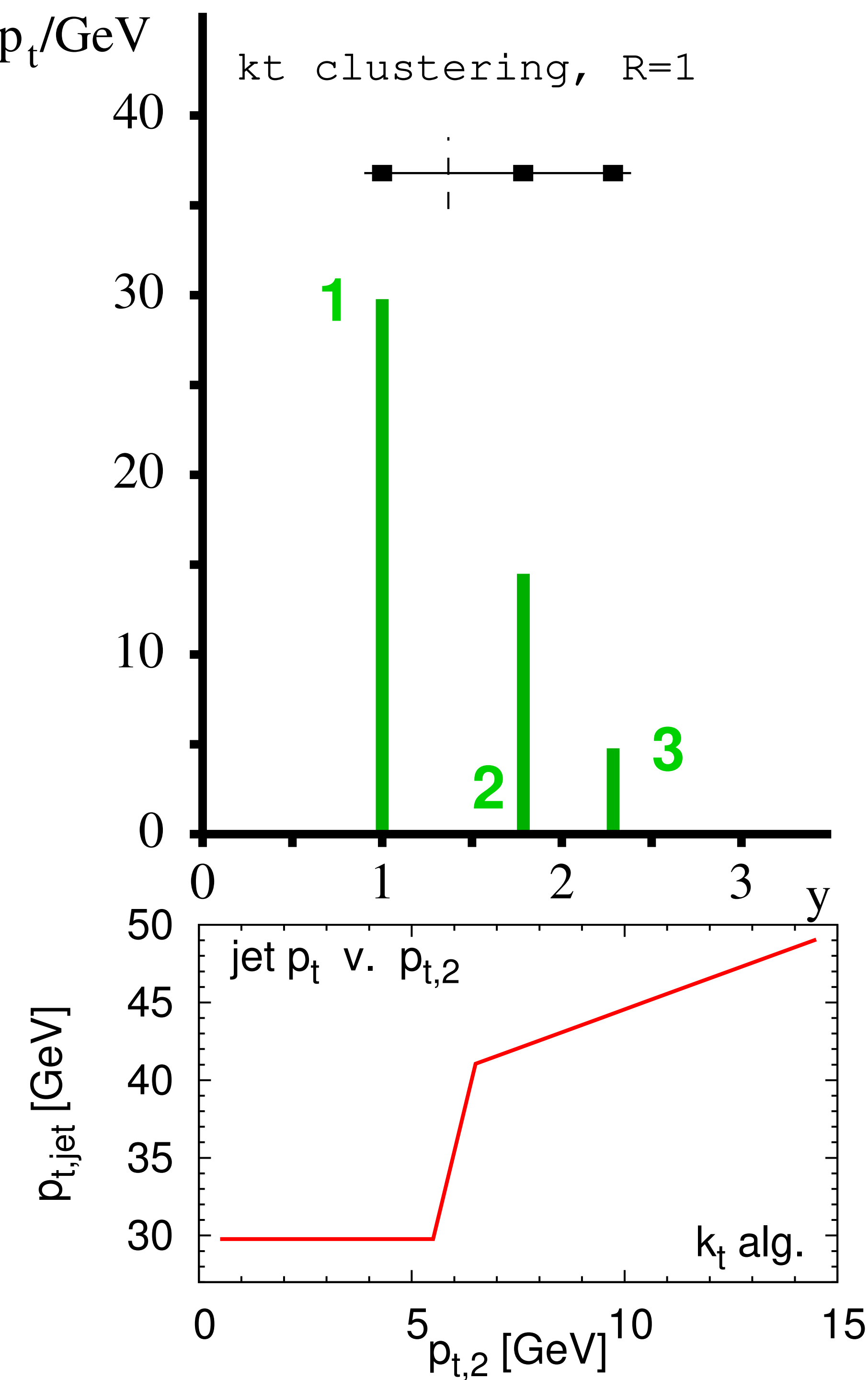


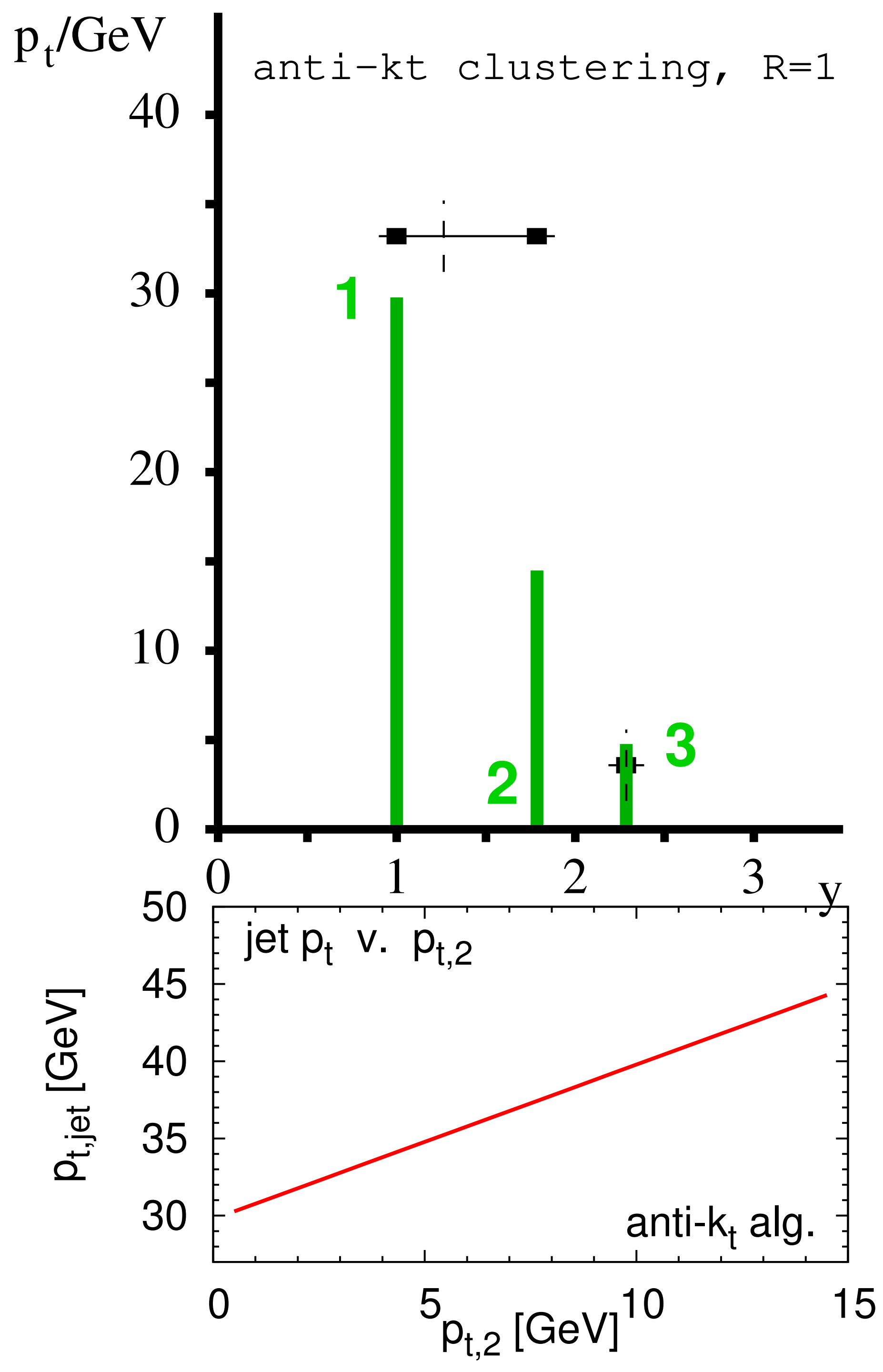
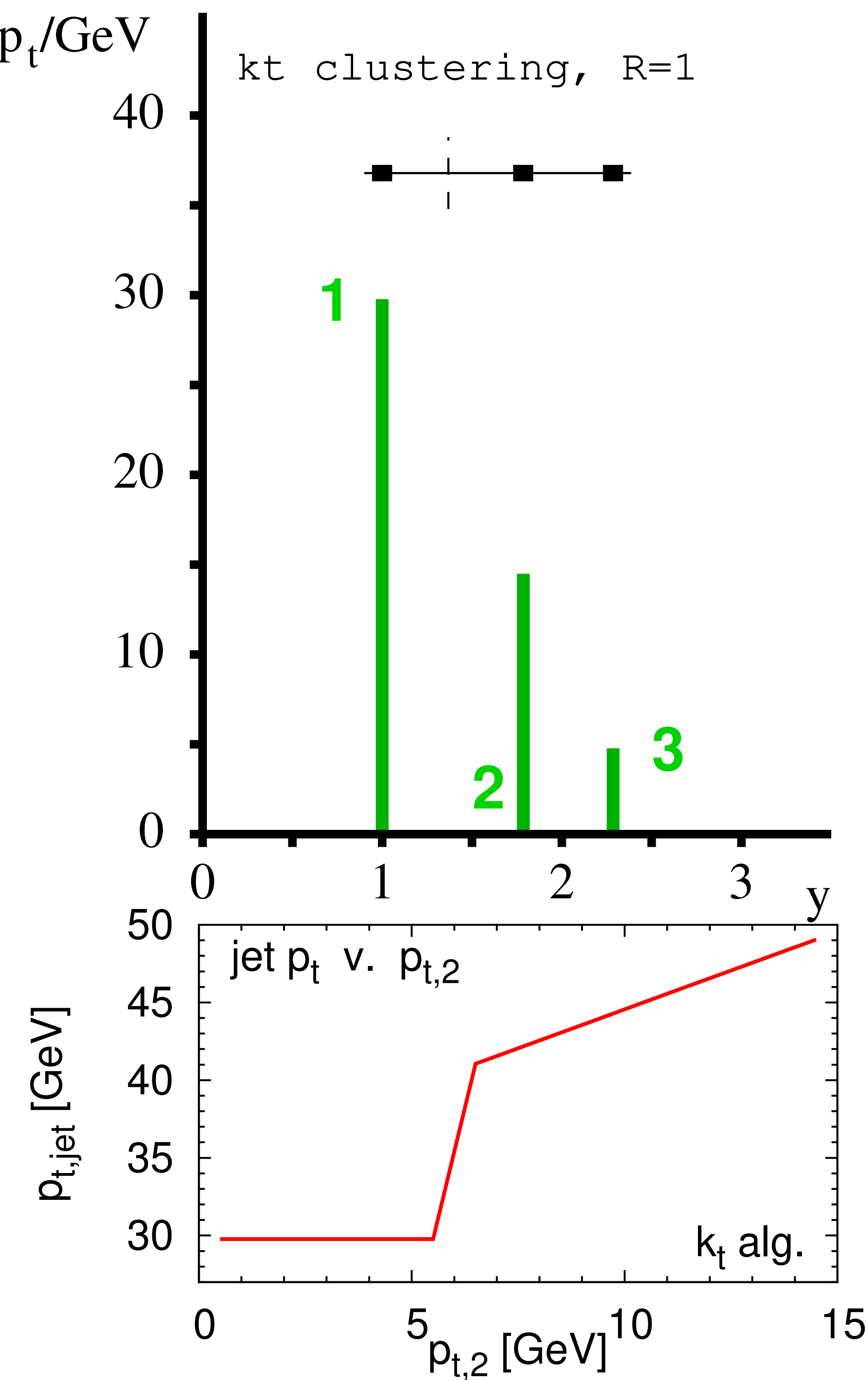






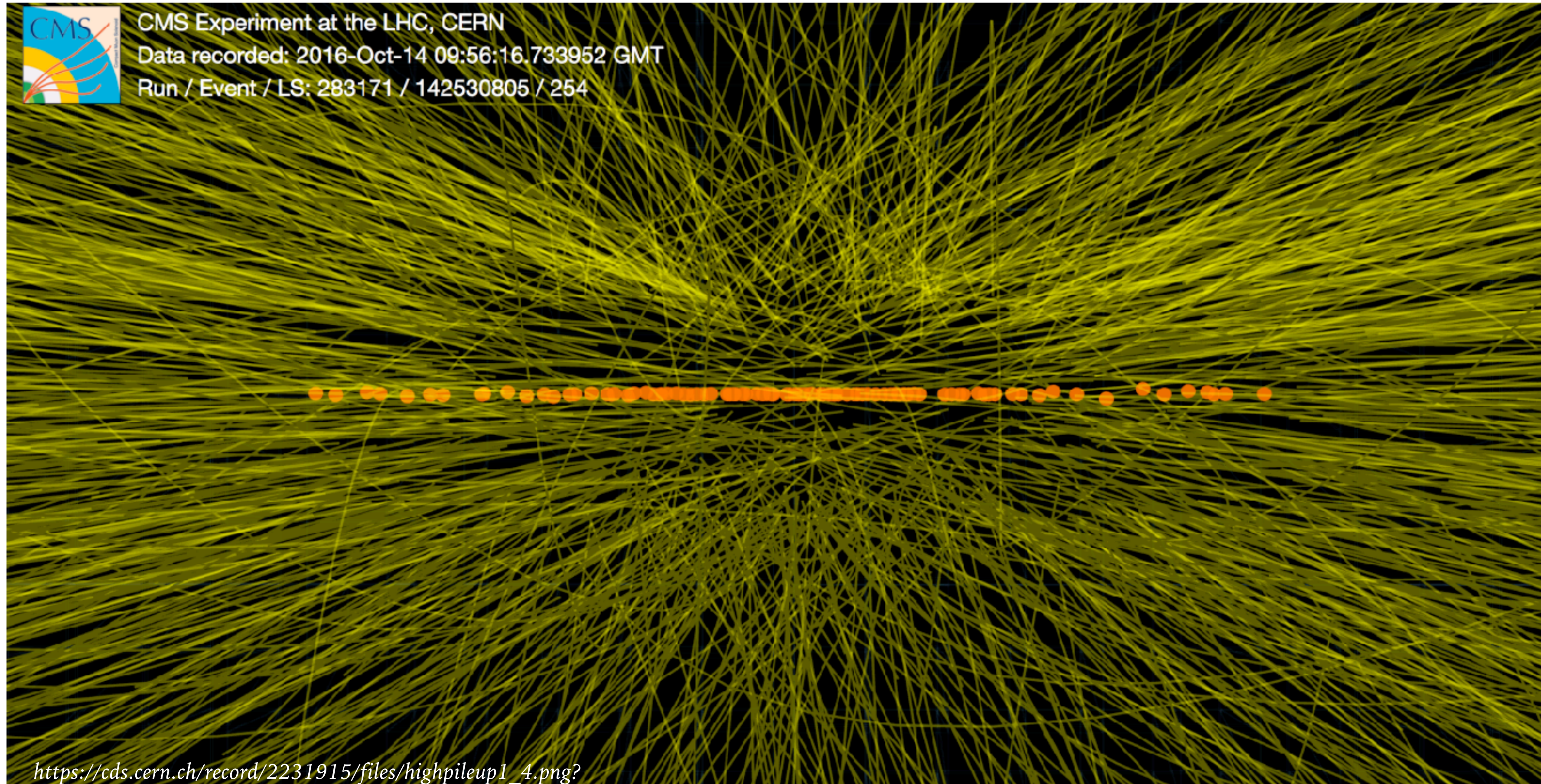






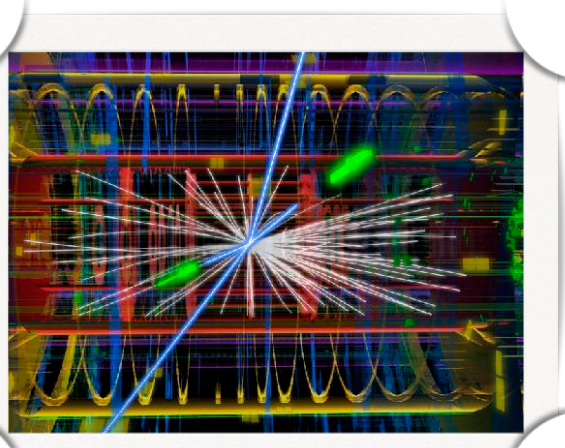
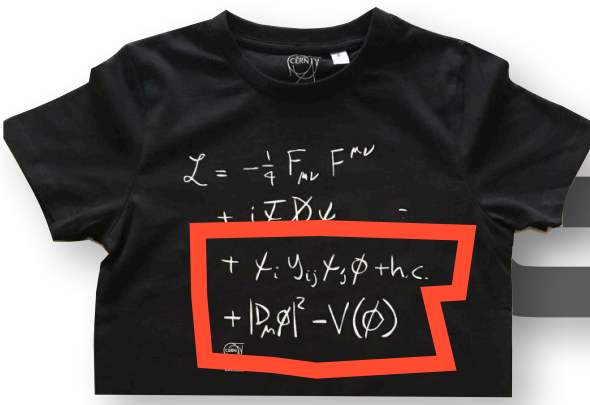
anti- k_t jet
algorithm gives
jet momentum
with better
linearity than k_t
algorithm
(or many other
jet algorithms)

advance #3: removal of pileup (many simultaneous pp collisions)

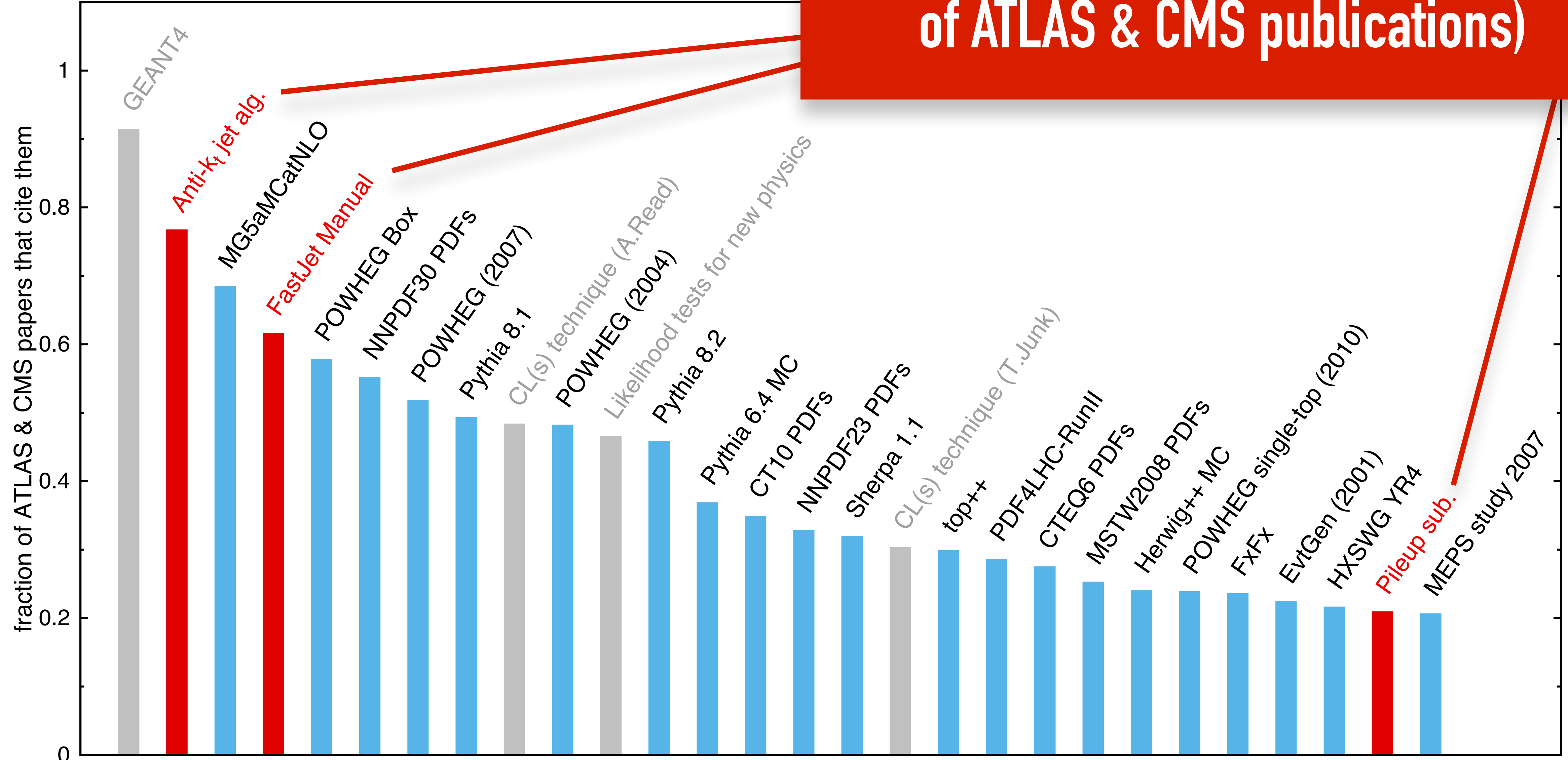


Nowadays many methods used (area-subtraction, particle/unified flow objects, PUPPI, soft-killer, ...) and machine-learning likely to play increasing role

Beyond scope of today's talk



those 3 advances are central to LHC physics today (e.g. anti- k_t used in $>70\%$ of ATLAS & CMS publications)



Plot by GP Salam based on data from InspireHEP

looking inside jets — basics

most jet finding based on correspondence

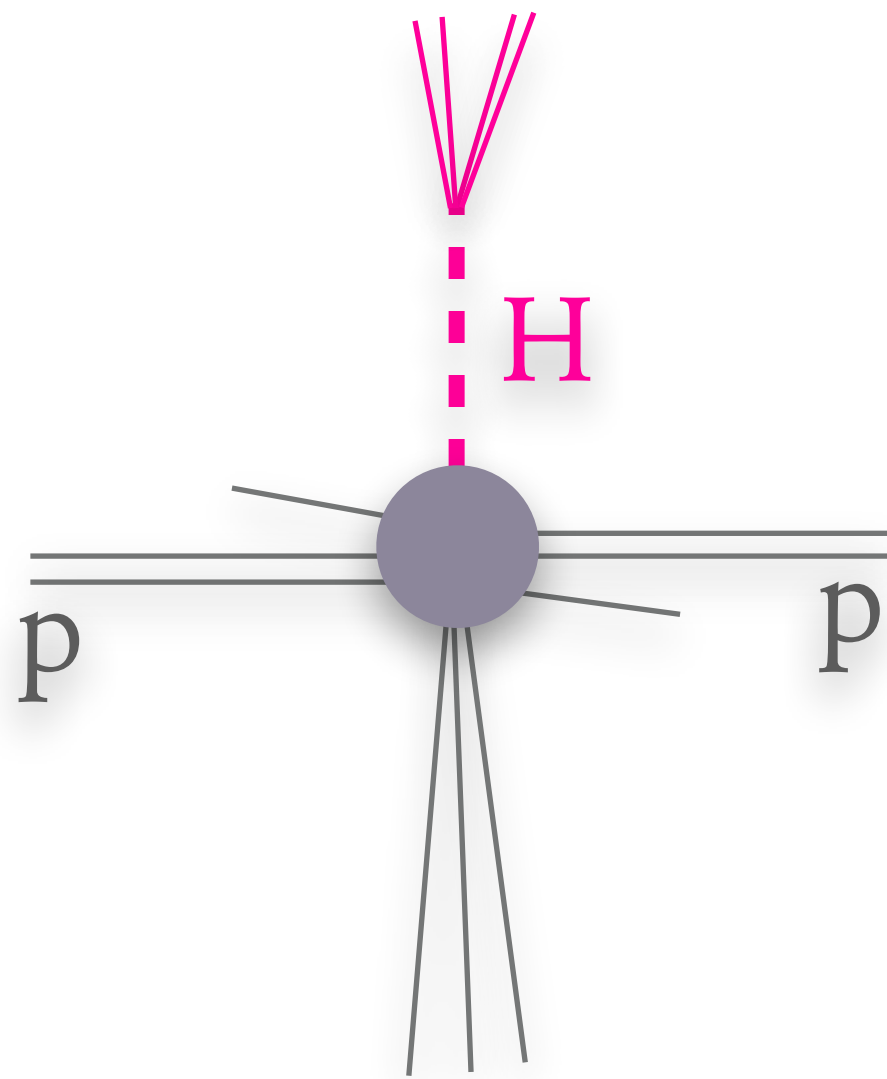
1 jet = 1 hard QCD parton (quark or gluon)

LHC forces us to go beyond that regime

high p_T Higgs & [SD] jet mass

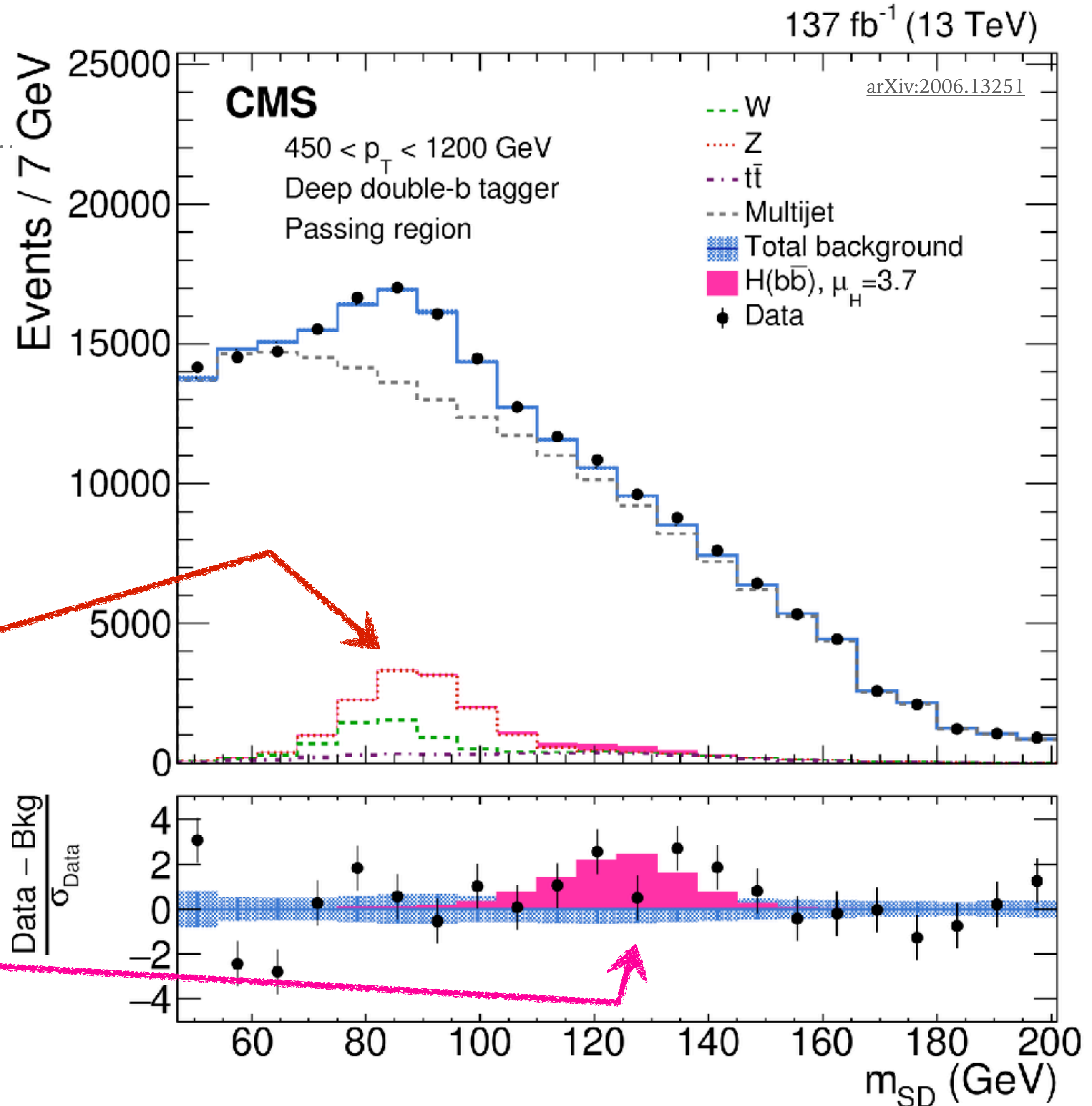
We wouldn't trust electromagnetism if we'd only tested it at one length/momentum scale.

New Higgs interactions need testing at both low and (here) high momenta.



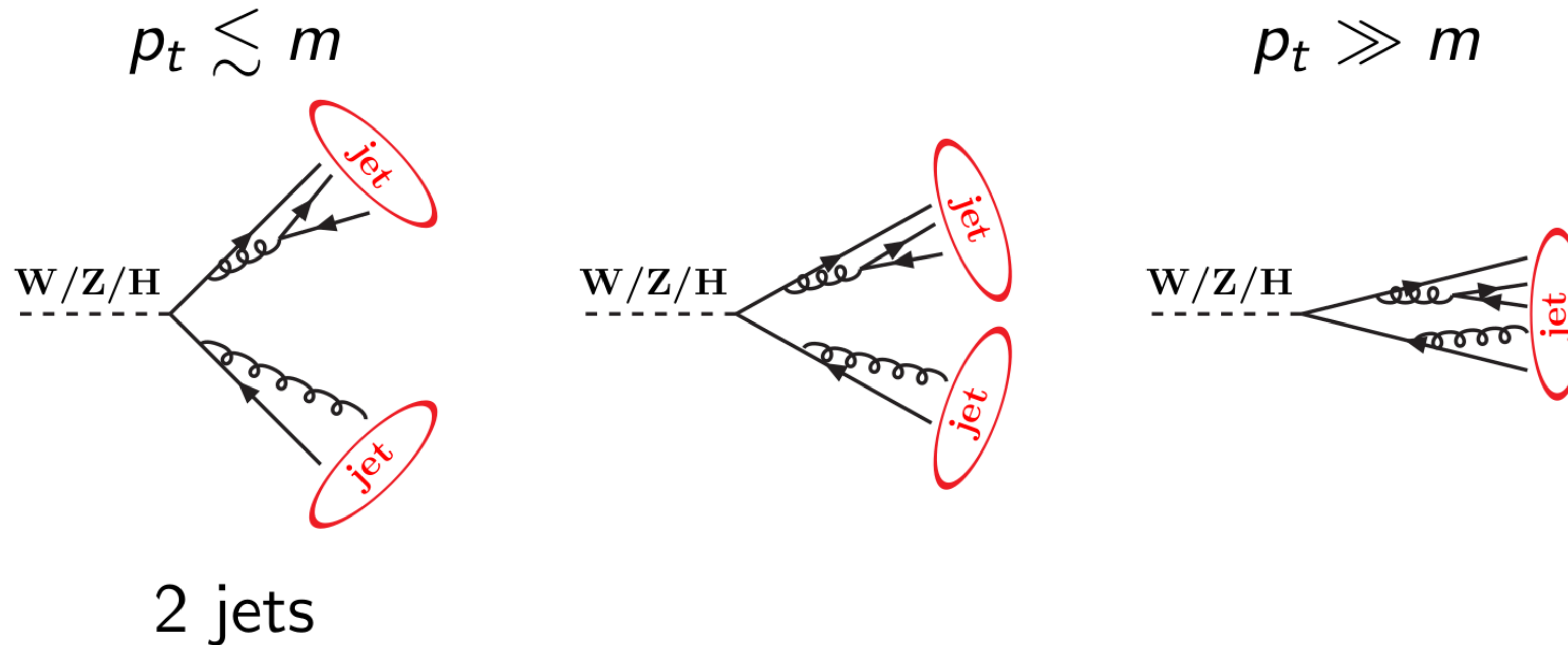
high- p_T
 $Z \rightarrow b\bar{b}$

high- p_T
 $H \rightarrow b\bar{b}$
 (2.5 σ)?

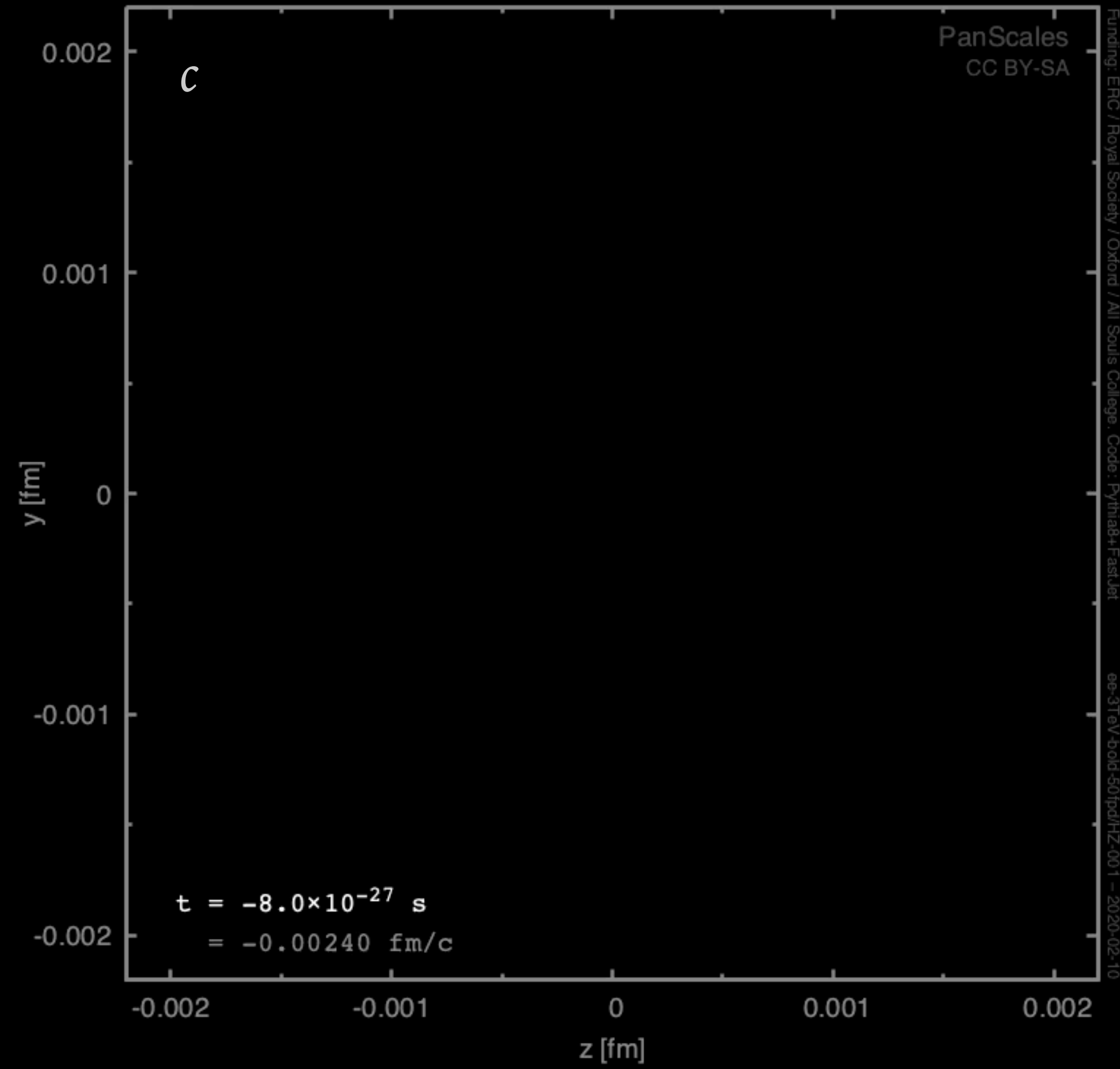
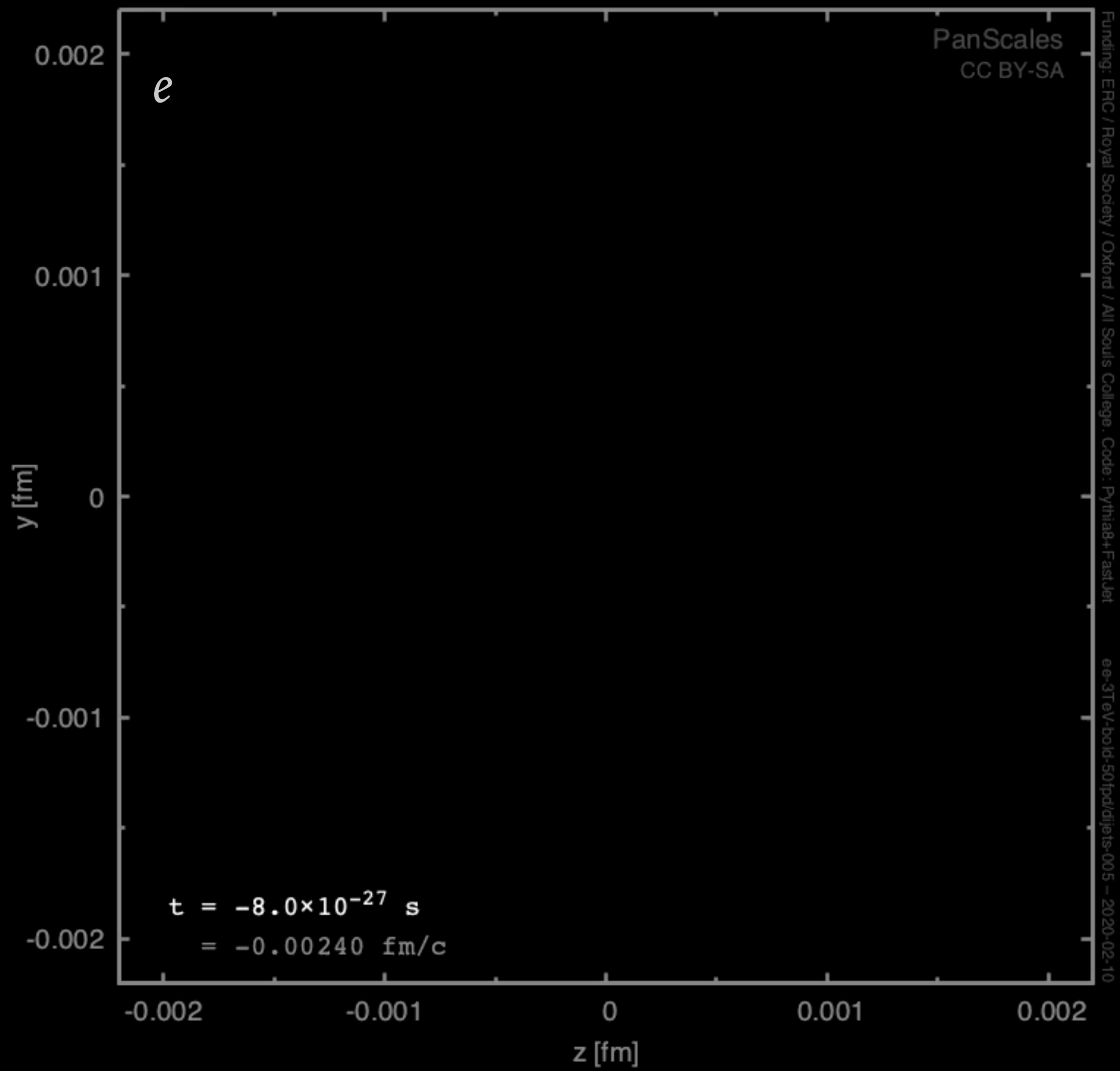


Jet substructure for boosted hadronic W/Z/H/t etc. decays

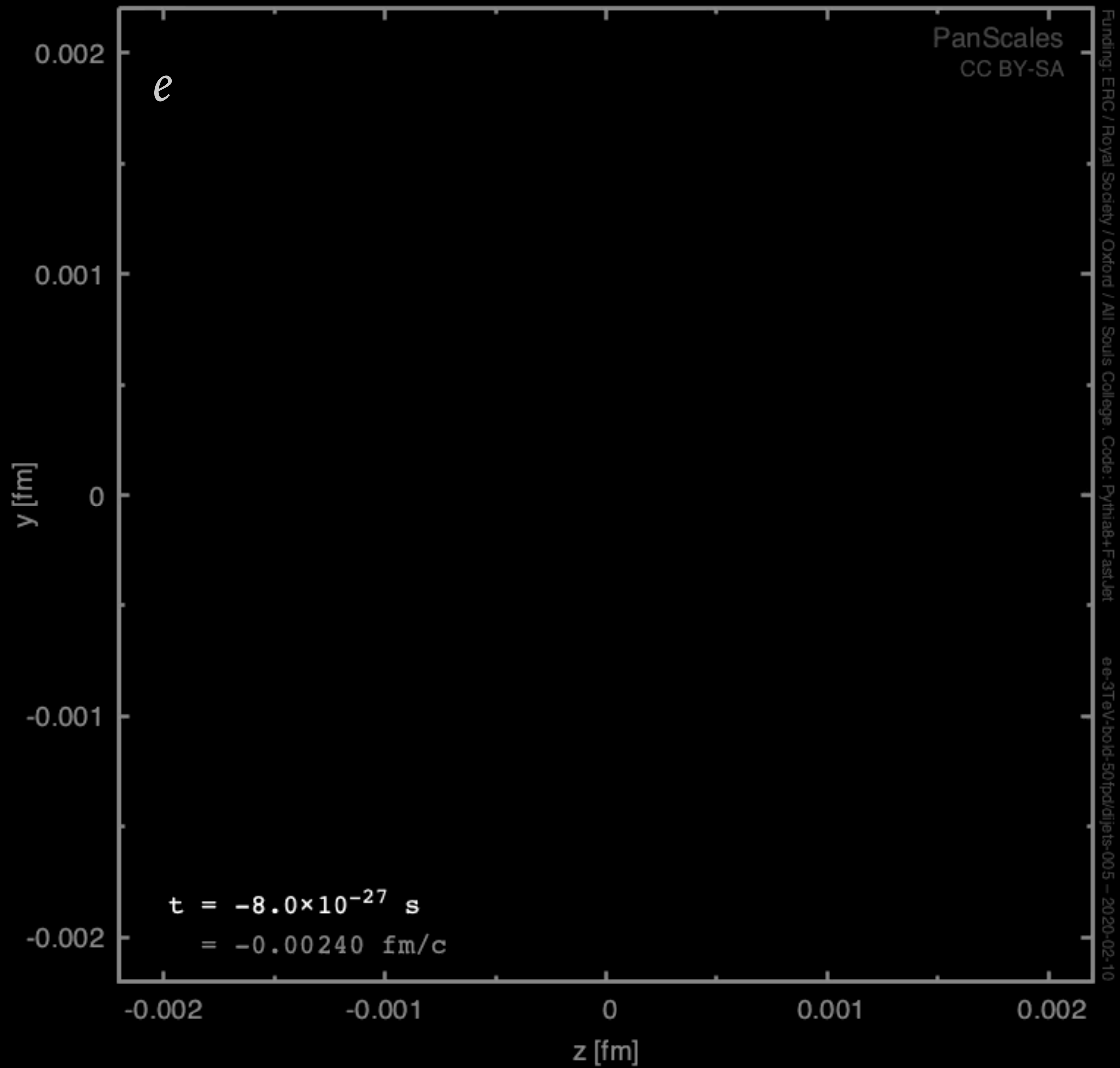
- ▶ At LHC energies, EW-scale particles (W/Z/t...) are often produced with $p_t \gg m$, leading to **collimated decays**.
- ▶ Hadronic decay products are thus often **reconstructed into single jets**.



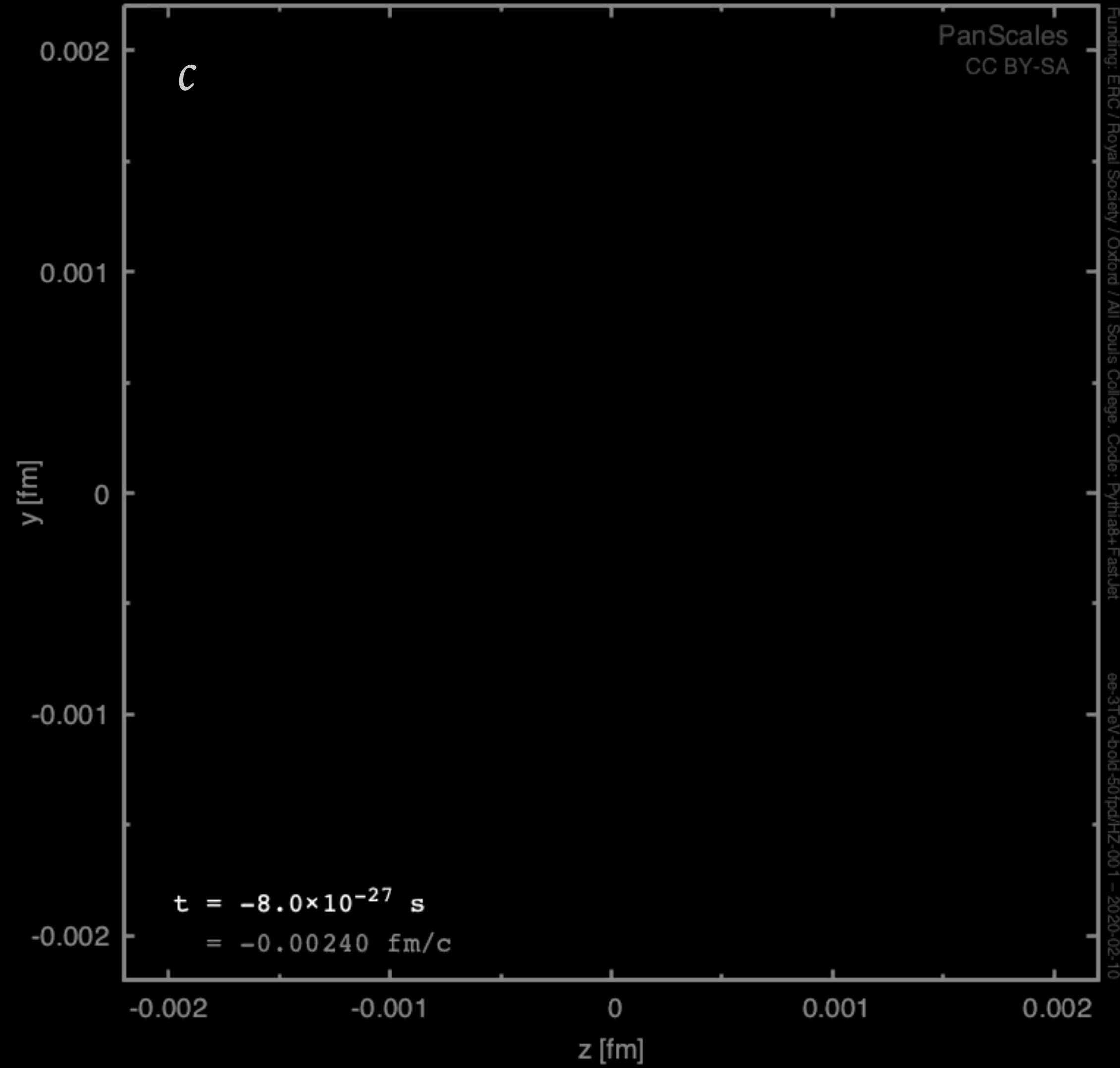
[Figure by G. Soyez]



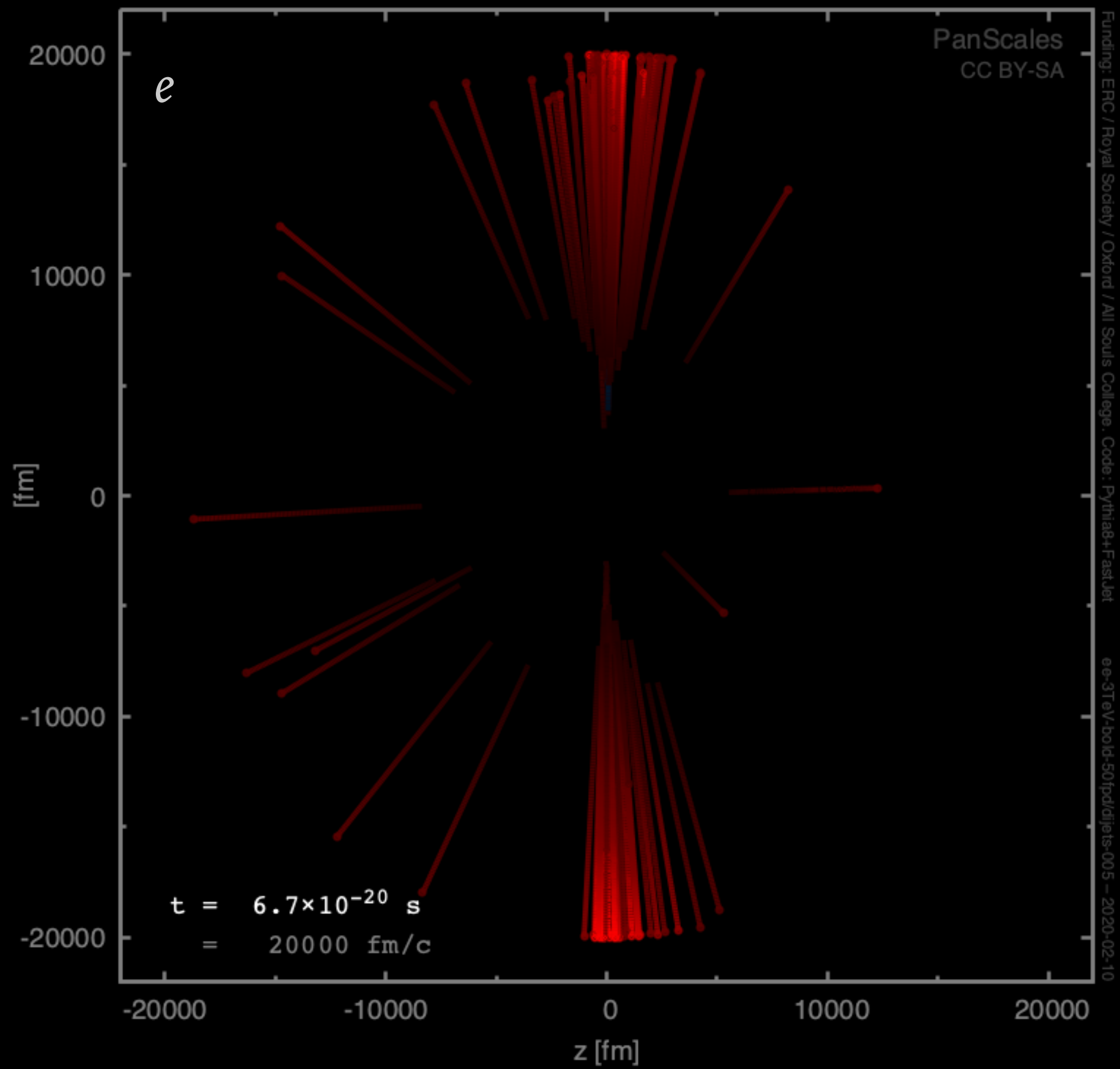
pure QCD event



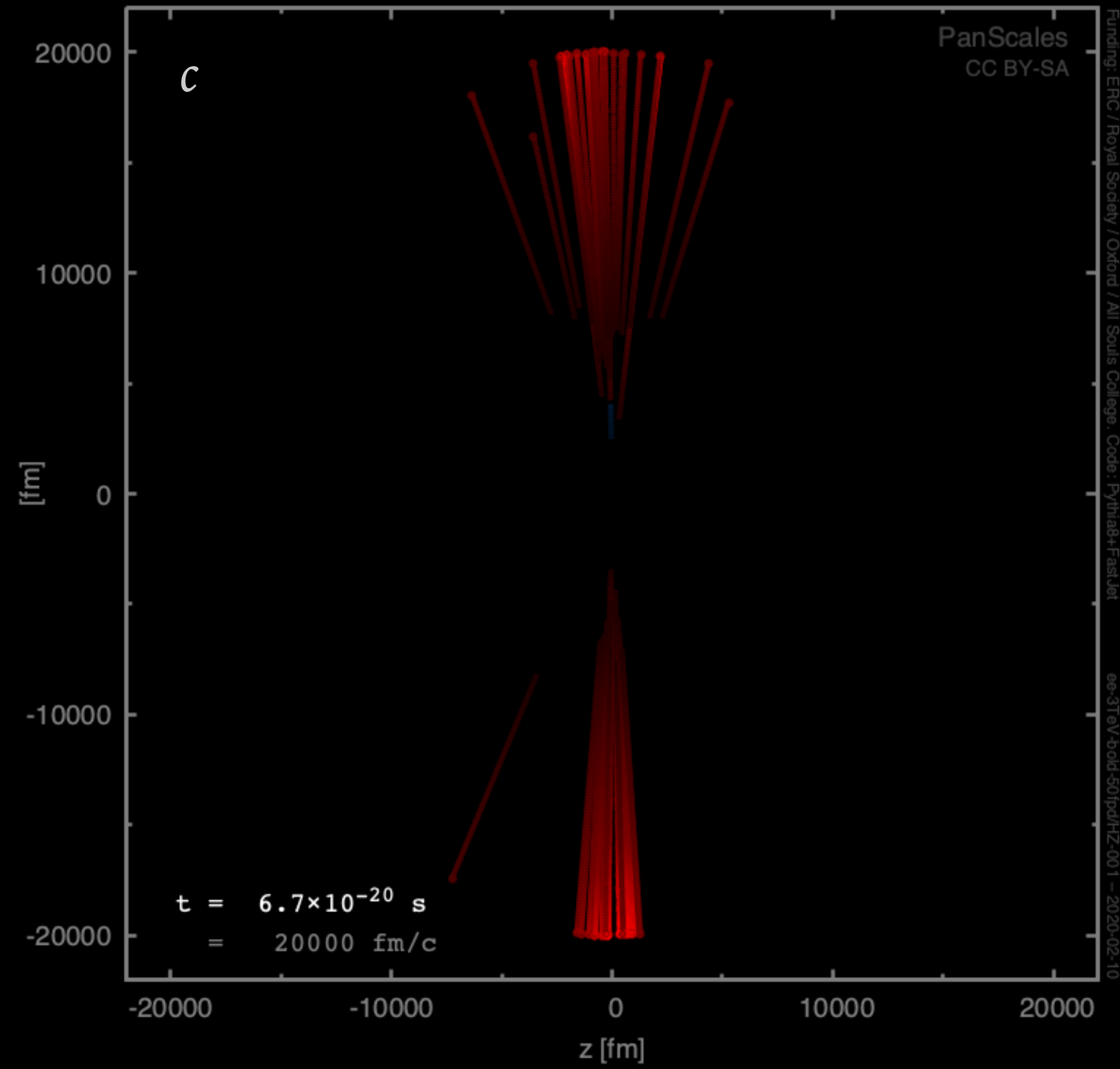
event with Higgs & Z boson decays



pure QCD event



event with Higgs & Z boson decays



Searches for new particles using cone and cluster jet algorithms: A Comparative study

Michael H. Seymour (Lund U.). Jun 1993. 23 pp.

Published in *Z.Phys. C62* (1994) 127-138

LU-TP-93-8

DOI: [10.1007/BF01559532](https://doi.org/10.1007/BF01559532)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)

[KEK scanned document](#)

[Detailed record](#) - [Cited by 179 records](#) 100+

“As a simple example (in fact the only way in which we use sub-jets in this paper), one could cluster the event until there is exactly one jet remaining-this is then the hardest jet. Then **one could recluster only those particles that ended up in the hardest jet until there are exactly two jets**-these are then the sub-jets corresponding to the hardest emission within the hardest jet.”

“Then we recluster only those particles that ended up in the hardest jet, using a radius $R=\alpha R_{jj}$,”

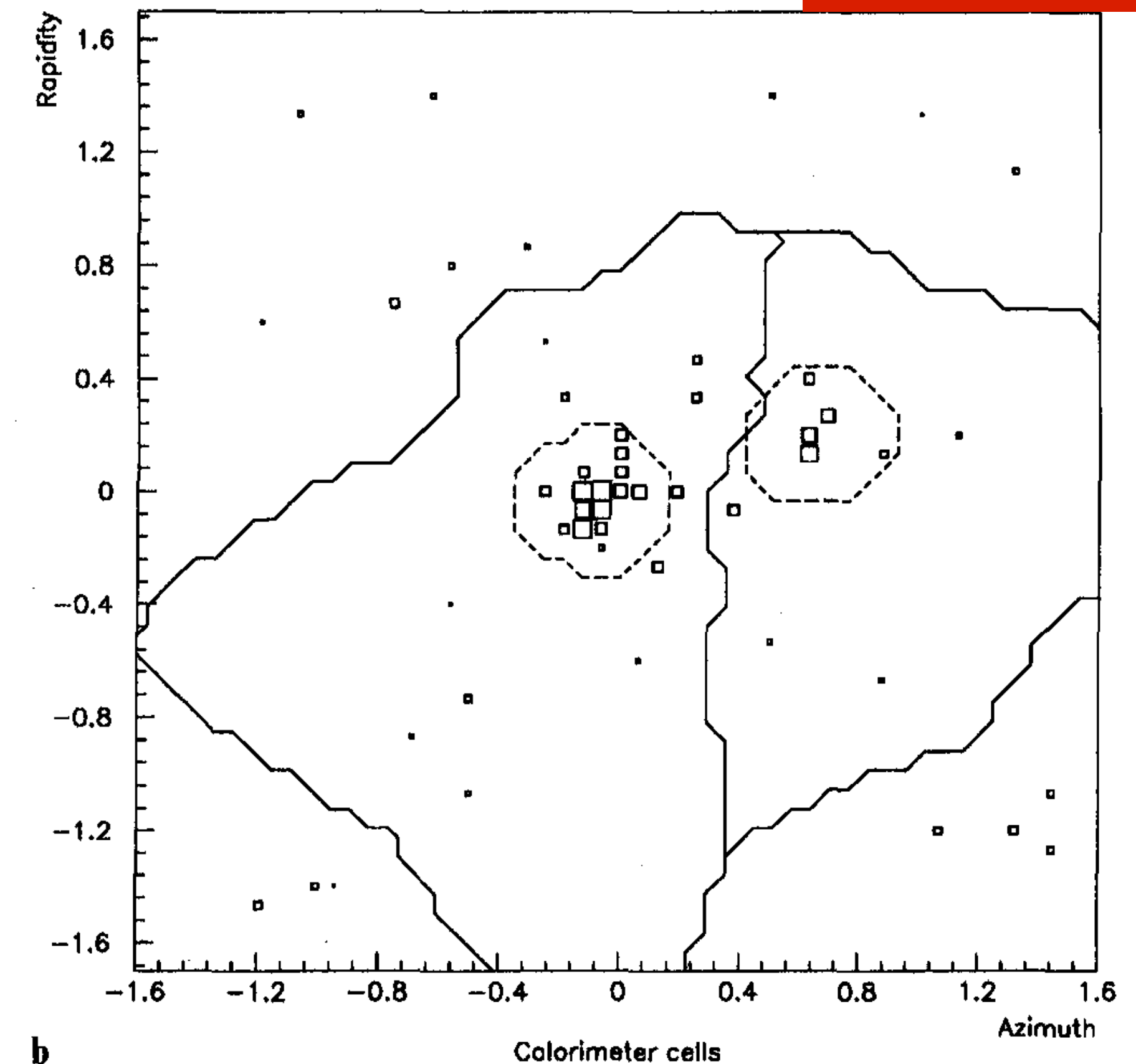


Fig. 2. A hadronic W decay, as seen at calorimeter level, **a** without, and **b** with, particles from the underlying event. Box sizes are logarithmic in the cell energy, lines show the borders of the sub-jets for infinitely soft emission according to the cluster (solid) and cone (dashed) algorithms

WW scattering at the CERN LHC

J.M. Butterworth (University Coll. London), B.E. Cox, Jeffrey R. Forshaw (Manchester U.). Jan 2002. 29 pp.

Published in **Phys.Rev. D65 (2002) 096014**

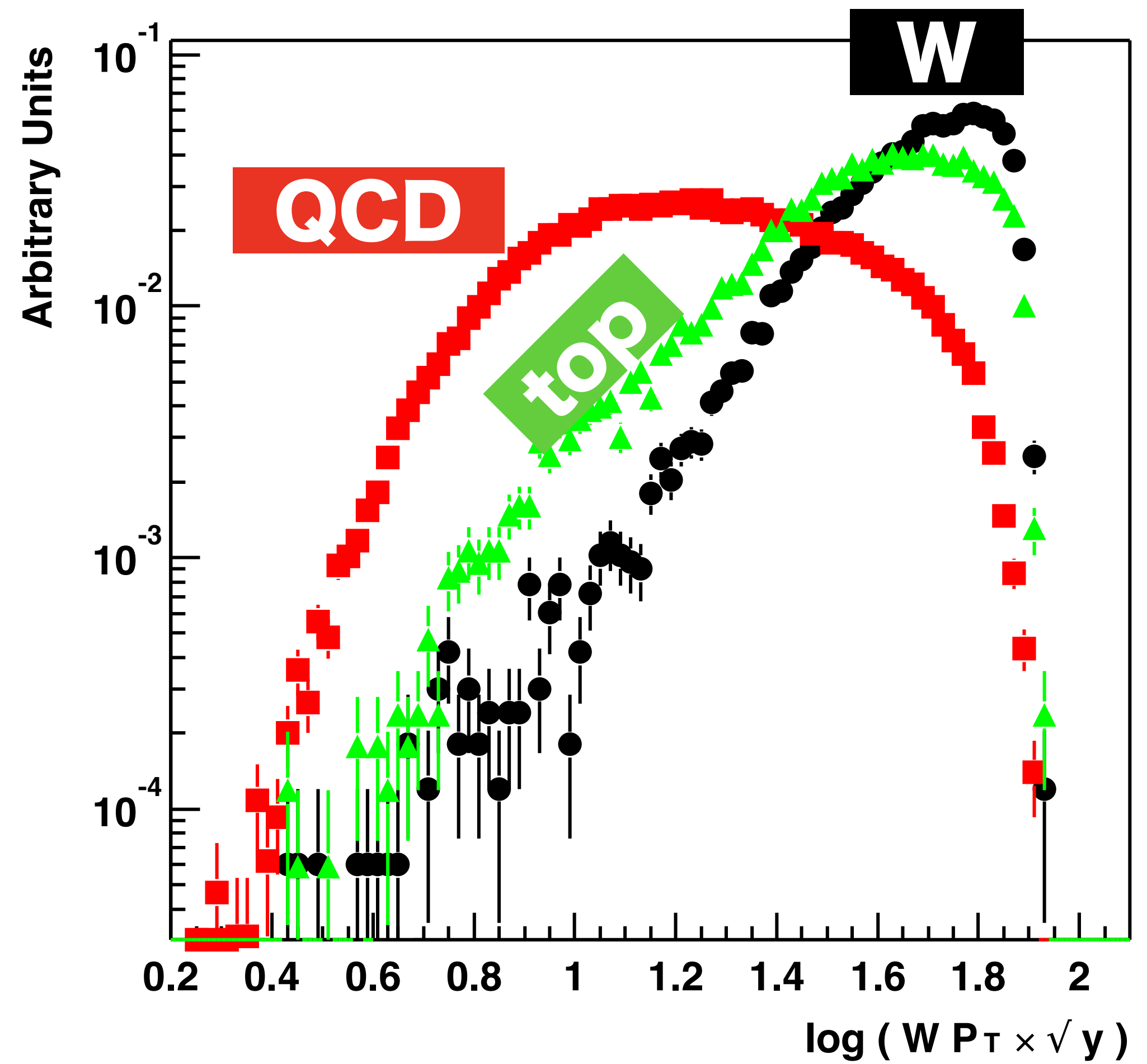
MC-TH-01-13, MAN-HEP-01-05, UCL-HEP-2001-06

DOI: [10.1103/PhysRevD.65.096014](https://doi.org/10.1103/PhysRevD.65.096014)

e-Print: [hep-ph/0201098](https://arxiv.org/abs/hep-ph/0201098) | [PDF](#)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[CERN Document Server](#); [ADS Abstract Service](#)

[Detailed record](#) - [Cited by 299 records](#) 250+



$$\log (W P_T \times \sqrt{y}) = \frac{1}{2} \log(d_{12})$$

this analysis we develop a new technique. The extra pieces of information gained from the subjet decomposition are the y cut at which the subjets are defined and the four-vectors of the subjets. For a genuine W decay the expectation is that the scale at which the jet is resolved into subjets (i.e. yp_T^2) will be $\mathcal{O}(M_W^2)$. The distribution of $\log(p_T \sqrt{y})$ is shown in Figure 12(d). The scale of the splitting is indeed high in the signal and softer in the $W + \text{jets}$ background, where the hadronic W is in general a QCD jet rather than a genuine second W . A cut is applied at $1.6 < \log(p_T \sqrt{y}) < 2.0$. The effect of this cut is

Better jet clustering algorithms

Yuri L. Dokshitzer (Milan U.), G.D. Leder, S. Moretti, B.R. Webber (Cambridge U.). Jul 1997. 33 pp.

Published in **JHEP 9708 (1997) 001**

CAVENDISH-HEP-97-06

DOI: [10.1088/1126-6708/1997/08/001](https://doi.org/10.1088/1126-6708/1997/08/001)

e-Print: [hep-ph/9707323](https://arxiv.org/abs/hep-ph/9707323) | [PDF](#)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) |

[EndNote](#)

[ADS Abstract Service](#)

[Detailed record](#) - [Cited by 890 records](#) 500+

Hadronization corrections to jet cross-sections in deep inelastic scattering

M. Wobisch (Aachen, Tech. Hochsch.), T. Wengler (Heidelberg U.). Apr 1998. 10 pp.

Published in **In *Hamburg 1998/1999, Monte Carlo generators for HERA physics* 270-279**

PITHA-99-16

To be published in the proceedings of Conference: [C98-04-27](#), p.270-279

[Proceedings](#)

e-Print: [hep-ph/9907280](https://arxiv.org/abs/hep-ph/9907280) | [PDF](#)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)

[ADS Abstract Service](#)

[Detailed record](#) - [Cited by 477 records](#) 250+

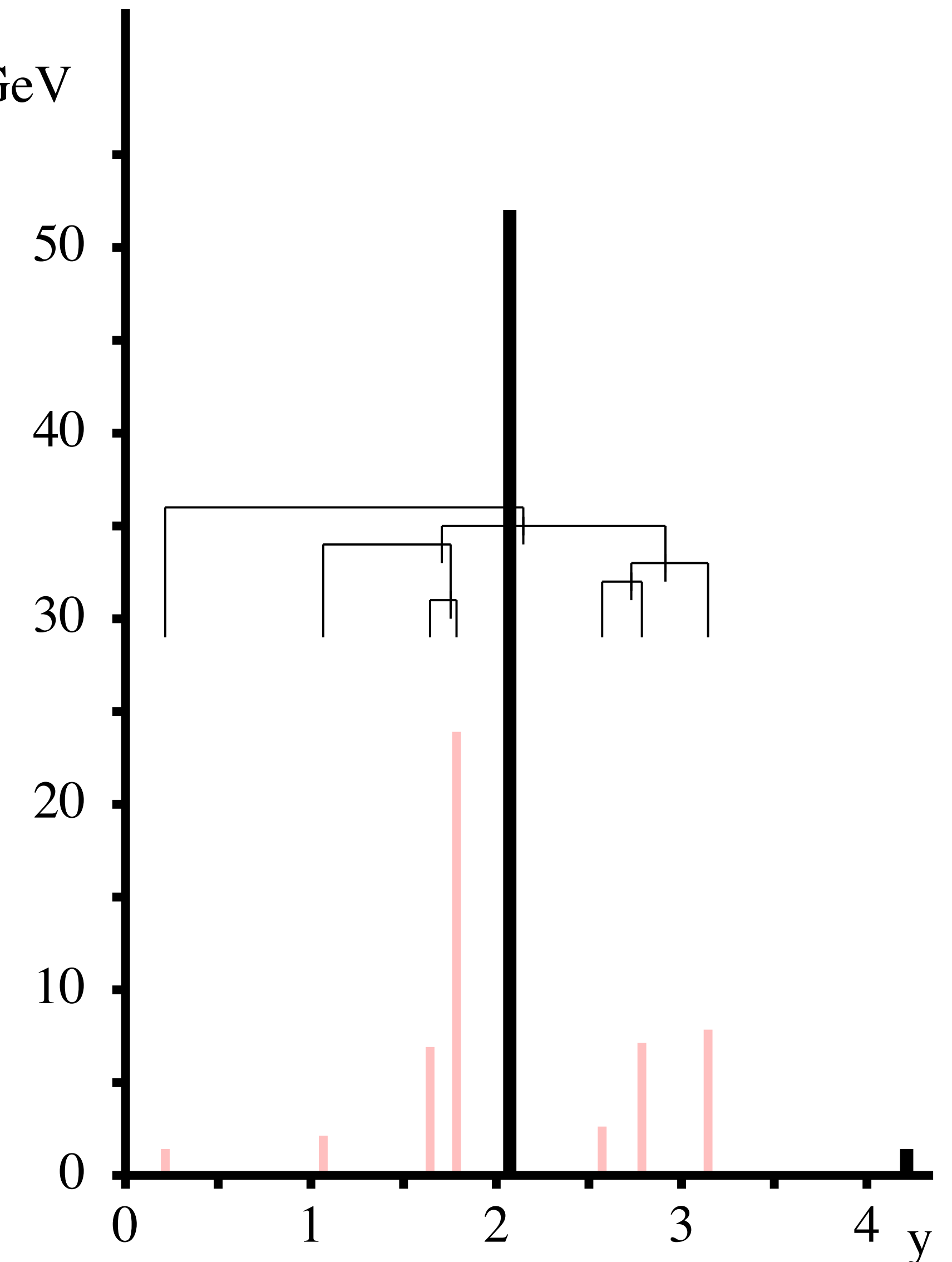
ABSTRACT: We investigate modifications to the k_{\perp} -clustering jet algorithm which preserve the advantages of the original Durham algorithm while reducing non-perturbative corrections and providing better resolution of jet substructure. We find that a simple change in the sequence of clustering (combining smaller-angle pairs first), together with the ‘freezing’ of soft resolved jets, has beneficial effects.

the Cambridge / Aachen (C/A) jet algorithm

1. Identify pair of particles, i & j , with smallest ΔR_{ij}
2. If $\Delta R_{ij} < R$ (jet radius parameter)
 - A. recombine i & j into a single particle
 - B. loop back to step 1
3. Otherwise, stop the clustering

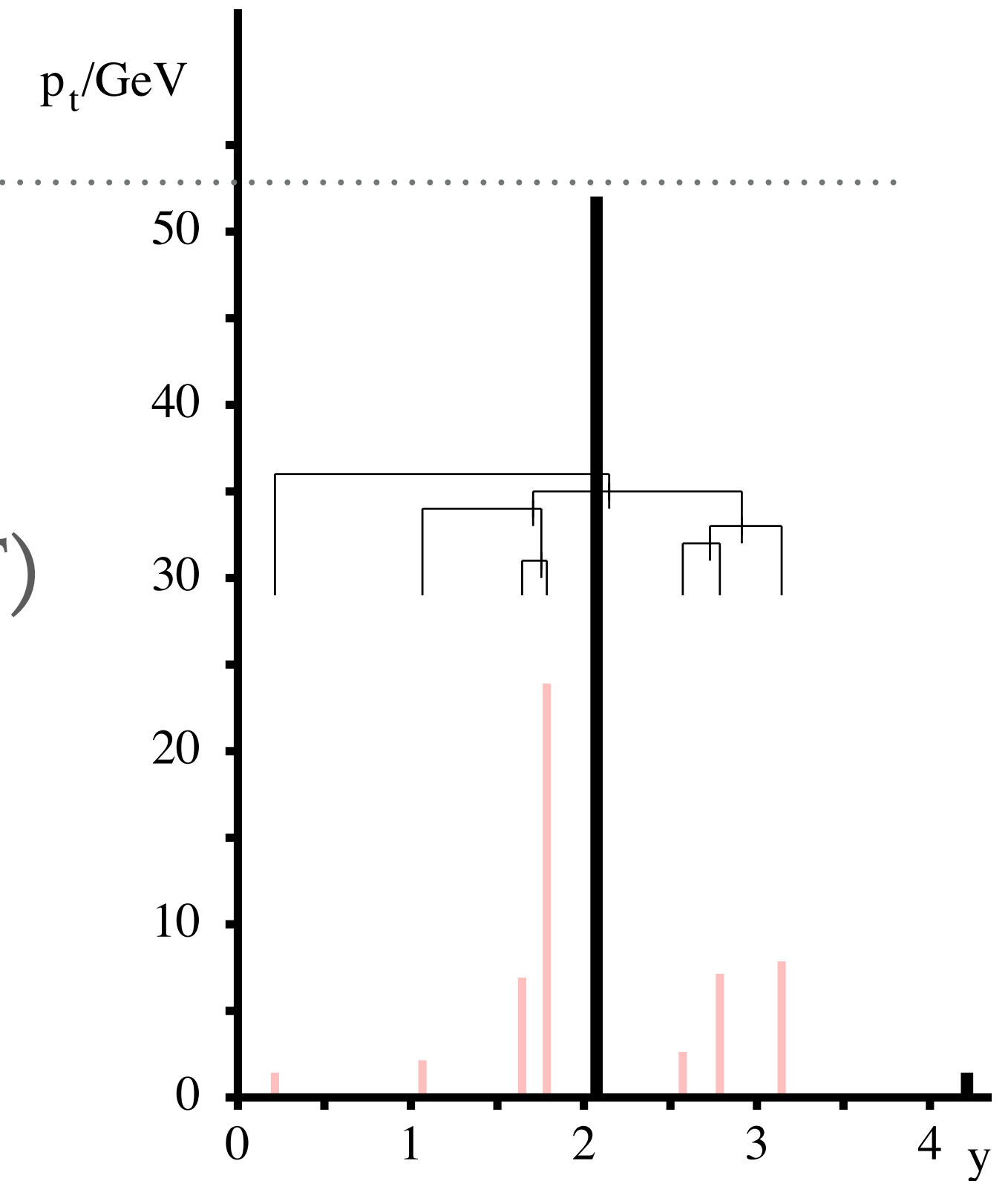
Dokshitzer, Leder, Moretti & Webber '97
Wobisch & Wengler '98

Cambridge/Aachen



A sequence of jet substructure tools taggers

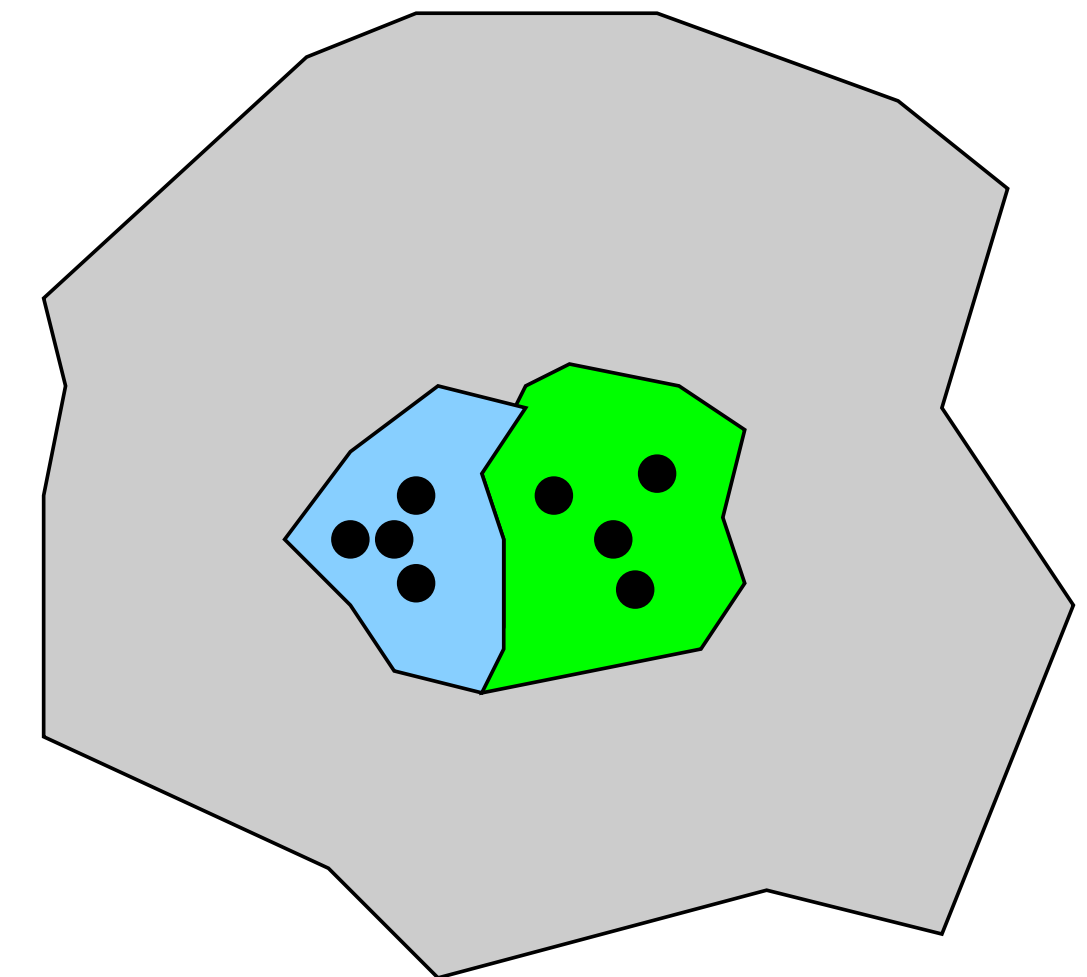
- **2008**: Mass-Drop Tagger (C/A declustering with a k_t/m cut)
[Butterworth, Davison, Rubin, GPS, [arXiv:0802.2470](https://arxiv.org/abs/0802.2470)]
- **2013**: Soft Drop, $\beta=0$, aka modified mass-drop tagger (mMDT)
[Dasgupta, Fregoso, Marzani, GPS, [arXiv:1307.0007](https://arxiv.org/abs/1307.0007)]
- **2014**: Soft Drop, $\beta \neq 0$
[Larkoski, Marzani, Soyez, Thaler, [arXiv:1402.2657](https://arxiv.org/abs/1402.2657)]



1. Undo last clustering of C/A jet into subjects 1, 2

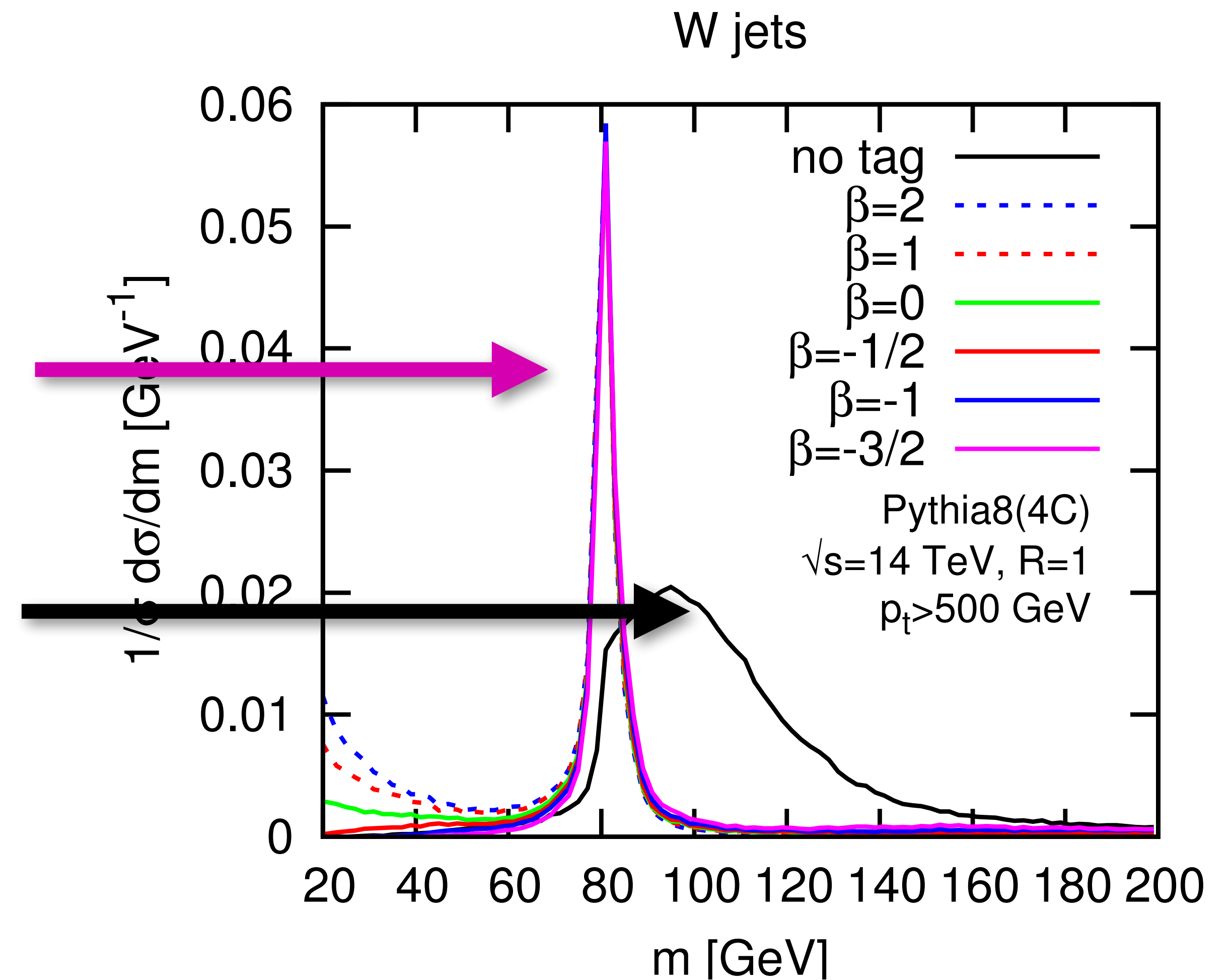
2. Stop if $z = \frac{\min(p_{t1}, p_{t2})}{p_{t1} + p_{t2}} \left(\frac{\Delta R_{12}}{R} \right)^\beta > z_{\text{cut}}$

3. Else discard softer branch, repeat step 1 with harder branch



SoftDrop: action on signal (e.g. W/H/Z)

Soft Drop signal jet mass
raw signal (plain jet mass)



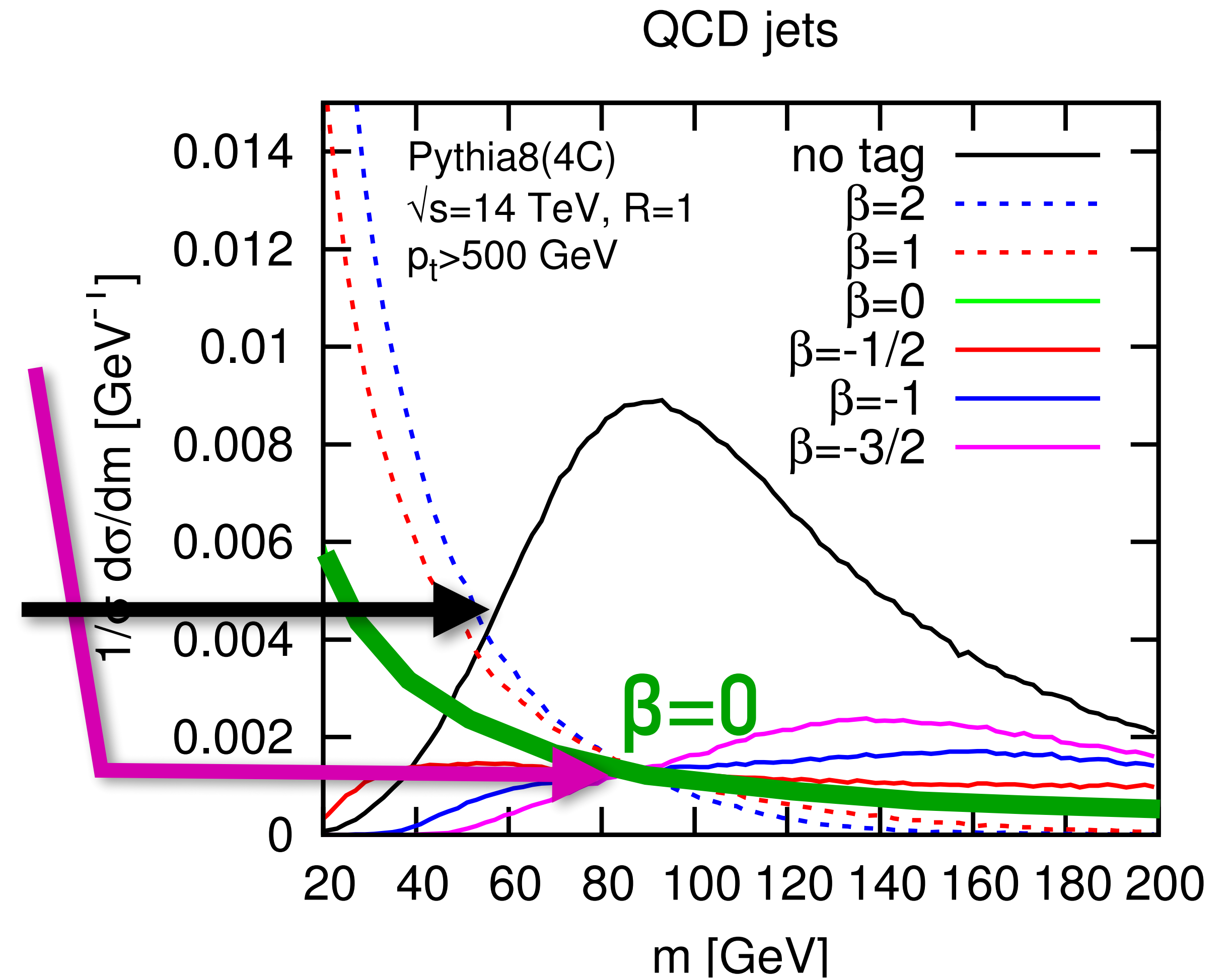
NB: z_{cut} chosen to keep signal efficiency fixed at 35% for all β

Plots from Larkoski, Marzani, Soyez & Thaler [arXiv:1402.2657](https://arxiv.org/abs/1402.2657)

SoftDrop: action on background (quark/gluon-induced jets)

Soft Drop background jet mass

raw background (plain jet mass)

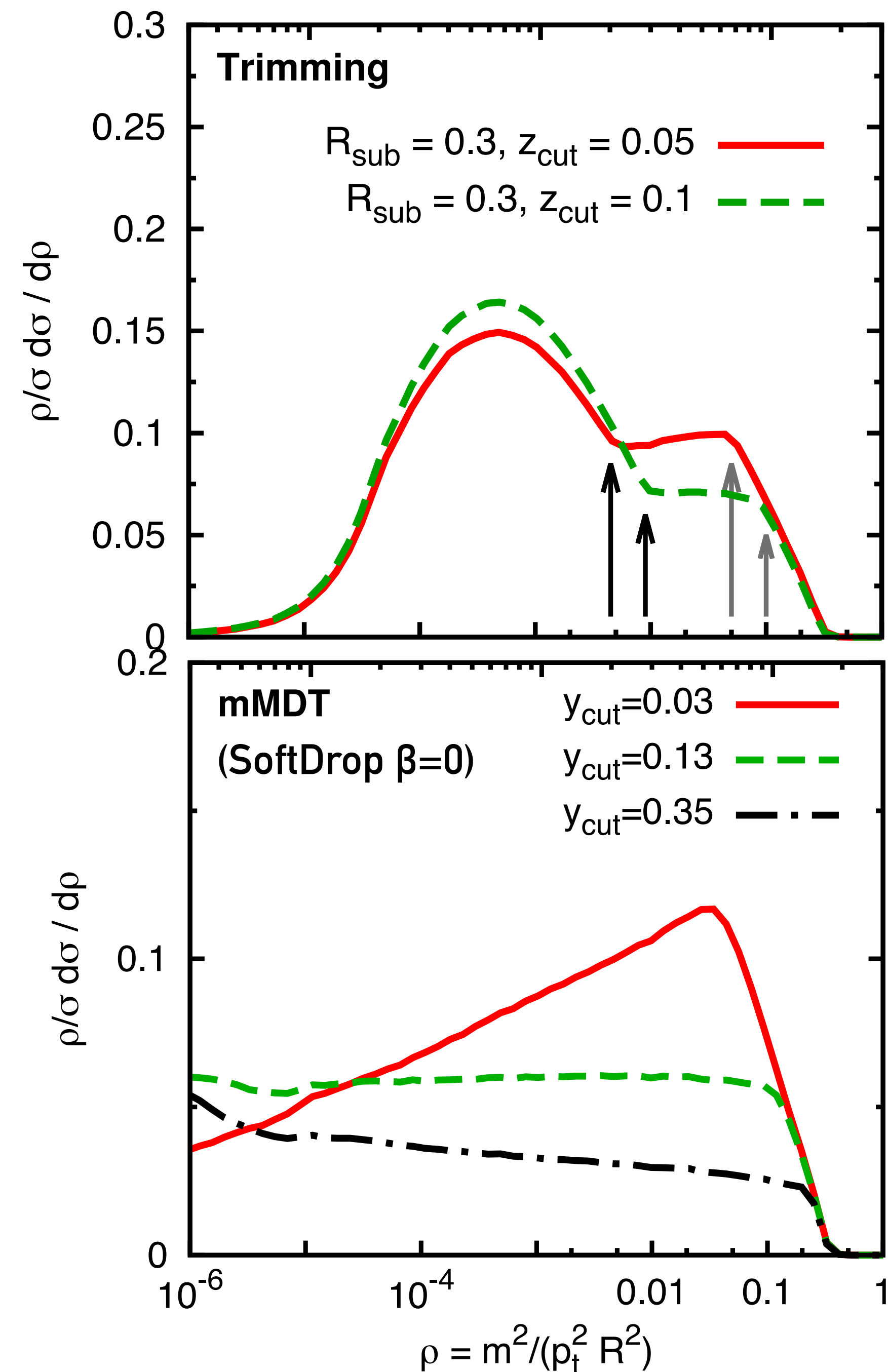


NB: z_{cut} chosen to keep signal efficiency fixed at 35% for all β

Plots from Larkoski, Marzani, Soyez & Thaler [arXiv:1402.2657](https://arxiv.org/abs/1402.2657)

For comparison: trimming sculpts background much more

- Trimming has three structures, induced by
 - z_{cut}
 - R_{sub}
 - Sudakov peak
- In comparison: just one structure in mMDT/SoftDrop (z_{cut})



SoftDrop $\beta=0$ (\equiv mMDT) has particularly simple QCD structure

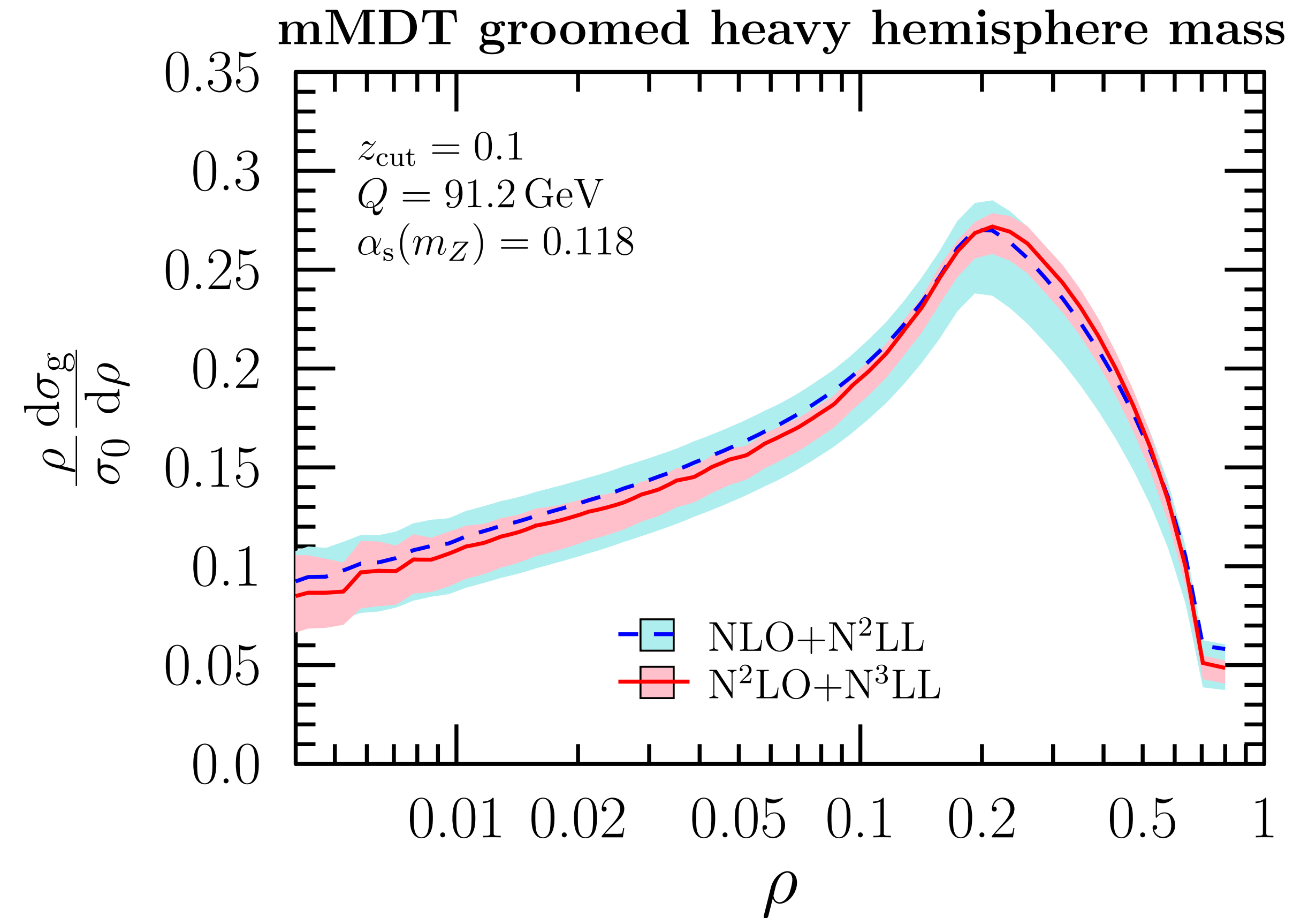
- most jet mass definitions involve **double-logarithmic** terms

$$(\alpha_s \ln^2 p_t/m)^n$$

- mMDT/SoftDrop ($\beta = 0$) has only **single logarithms**

$$(\alpha_s \ln p_t/m)^n$$

- simplicity \rightarrow most accurately calculated single-jet substructure observable



*Kardos, Larkoski, Trocsanyi, [arXiv:2002.00942](https://arxiv.org/abs/2002.00942) (small z_{cut})
other calc. approaches, see: [Anderle et al, arXiv:2007.10355](https://arxiv.org/abs/2007.10355)*

how much **more** info is there inside jets?

so far we examined use of hard 2-prong structure

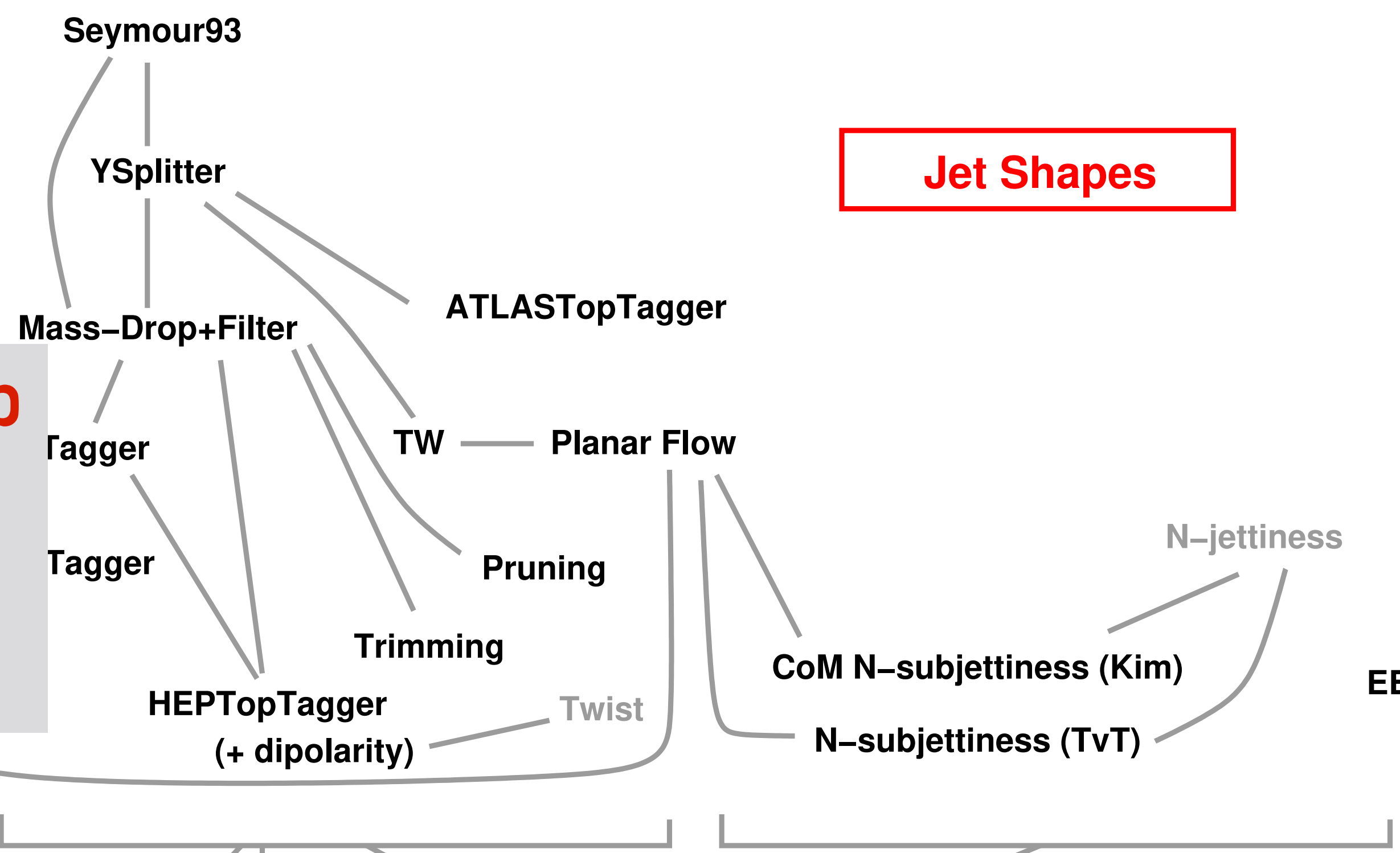
pp jet substructure field is full of activity

c. 2018

Jet Declustering

Matrix-Element

modified mass drop
soft drop
iterated soft drop
recursive soft drop



Jet Shapes

Degree	Connected Multigraphs
$d = 0$	
$d = 1$	
$d = 2$	
$d = 3$	
$d = 4$	
$d = 5$	

$C_n, D_n, v e_n^{(\beta)}, M_n, N_n, U_n, EFPs$

EEC

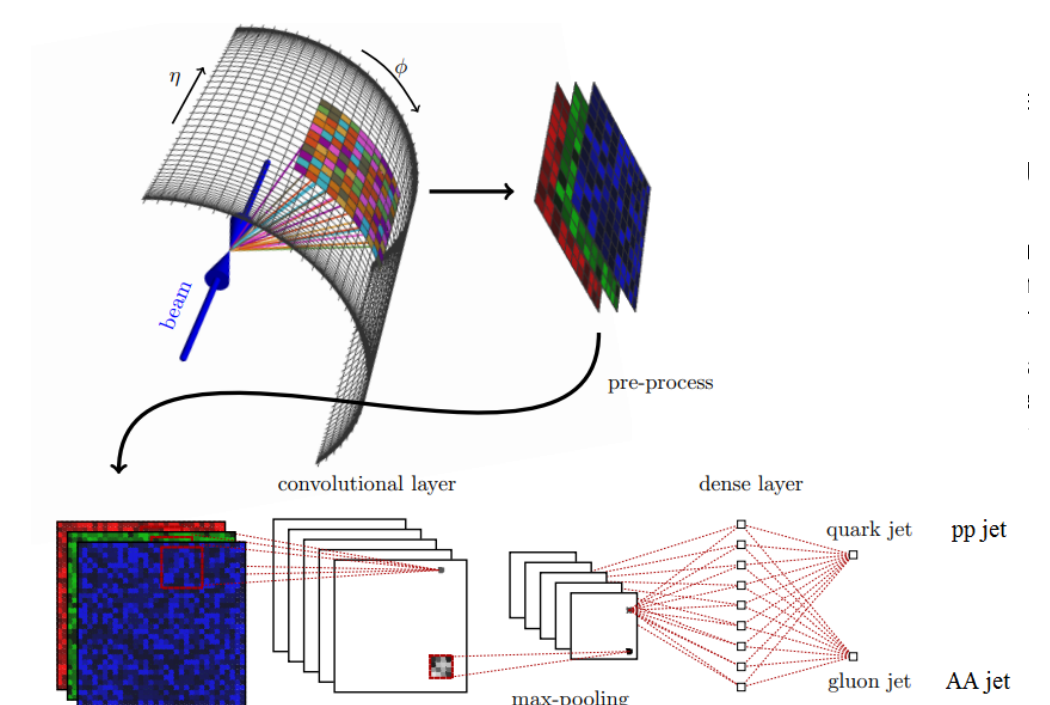
Shower Deconstruction

Qjets

Multi-variate tagger

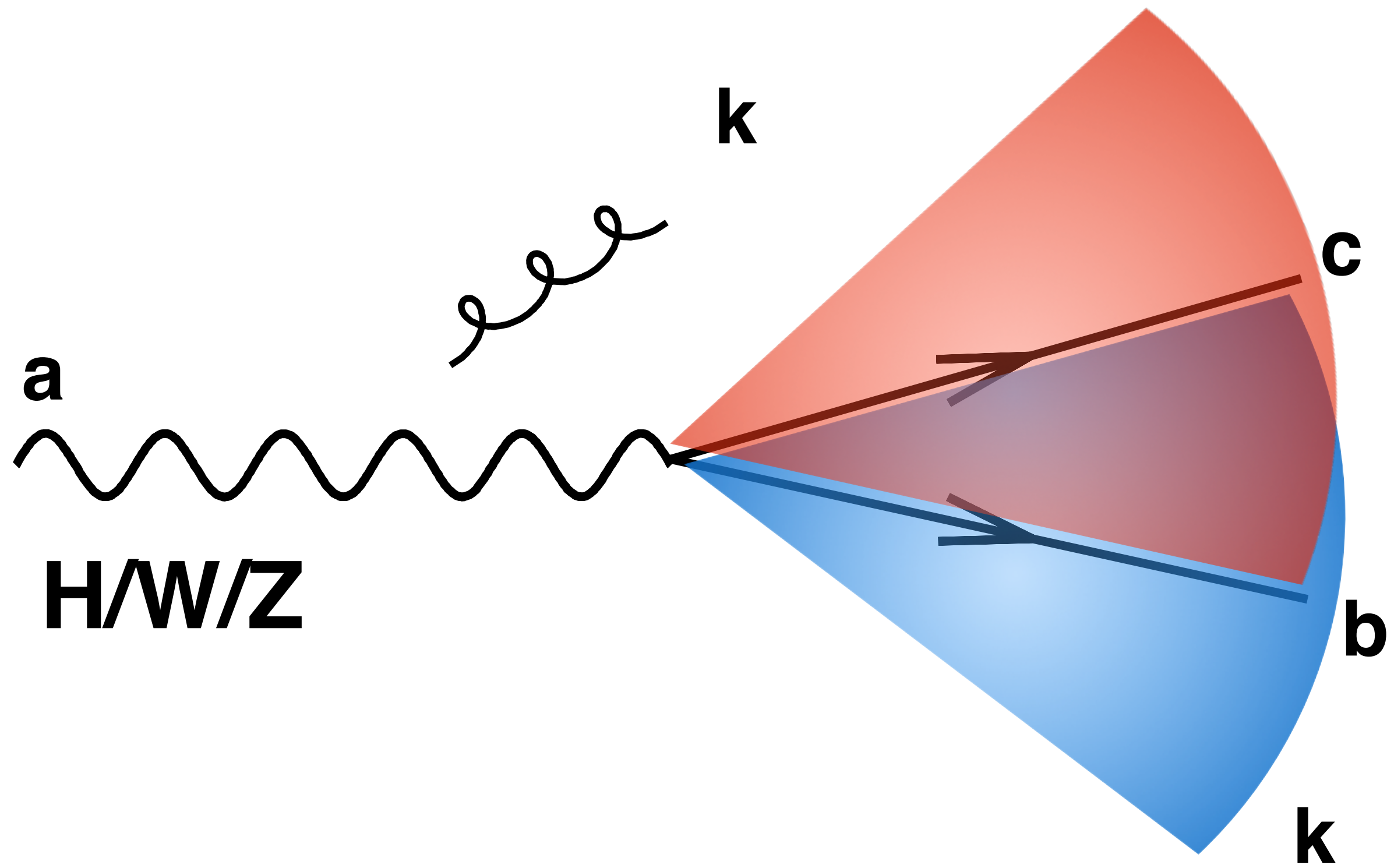
classification without labels
weak supervision

machine learning
DNN, CNN,
RNN, LSTM, etc

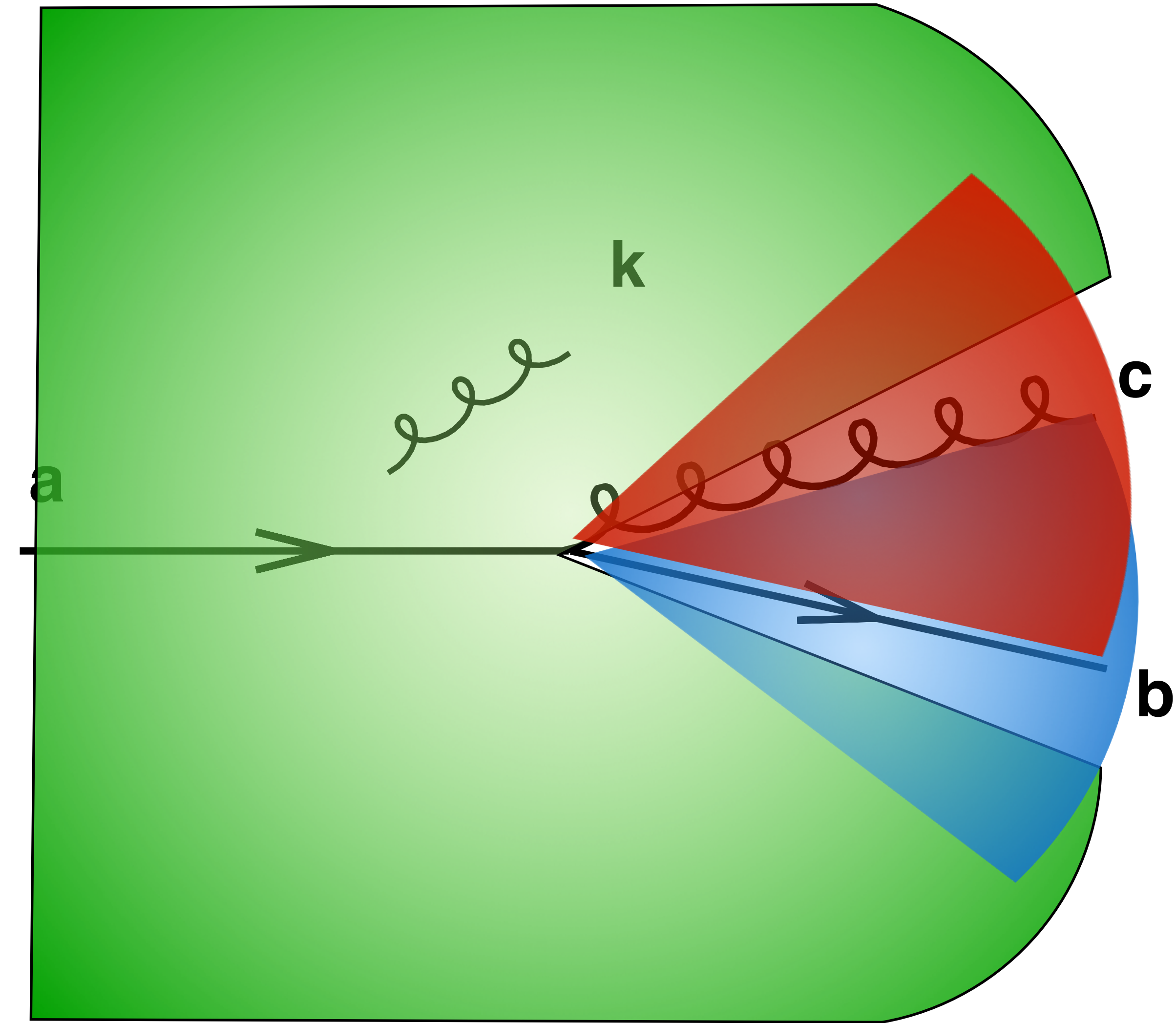


etc.

signal vs. background radiation patterns (first practical exploitation, Thaler & van Tilburg, N-subjettiness, [1011.2268](#))



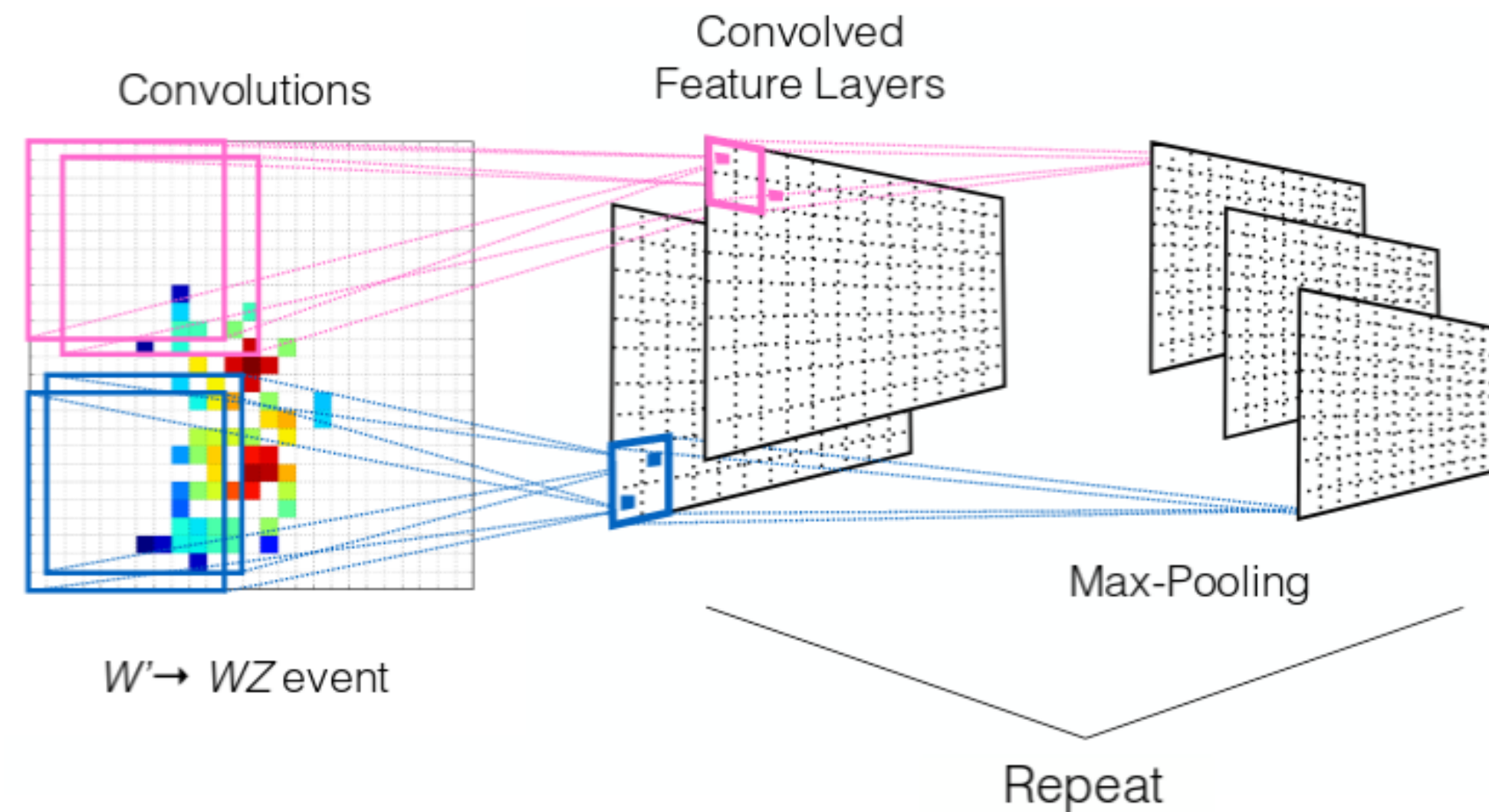
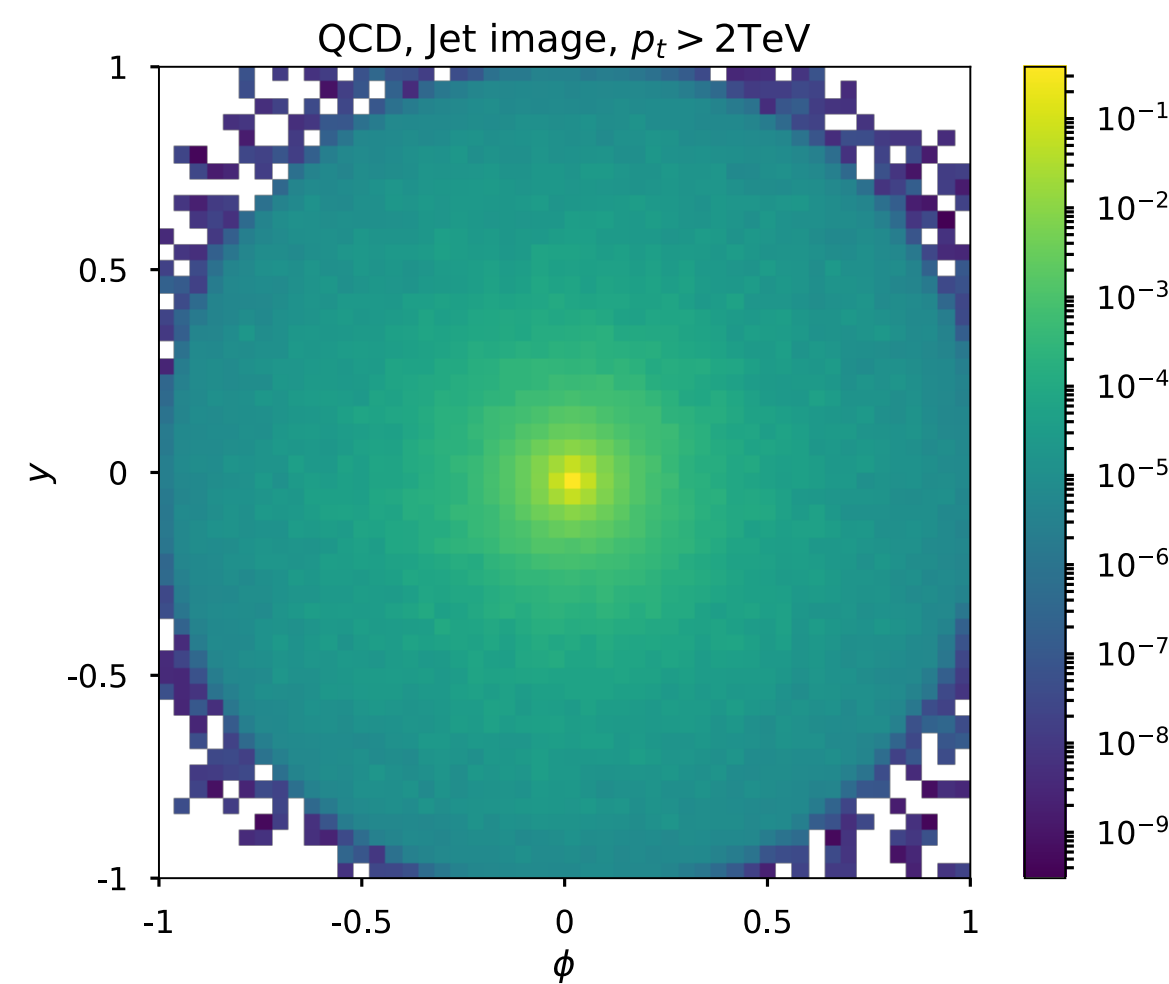
signal ($H/W/Z \rightarrow q\bar{q}$)



background ($q \rightarrow qqg$)

Machine learning and jet/event structure

- ▶ Project a jet onto a fixed $n \times n$ pixel image in rapidity-azimuth, where each pixel intensity corresponds to the momentum of particles in that cell.
- ▶ Can be used as input for classification methods used in computer vision, such as deep convolutional neural networks.

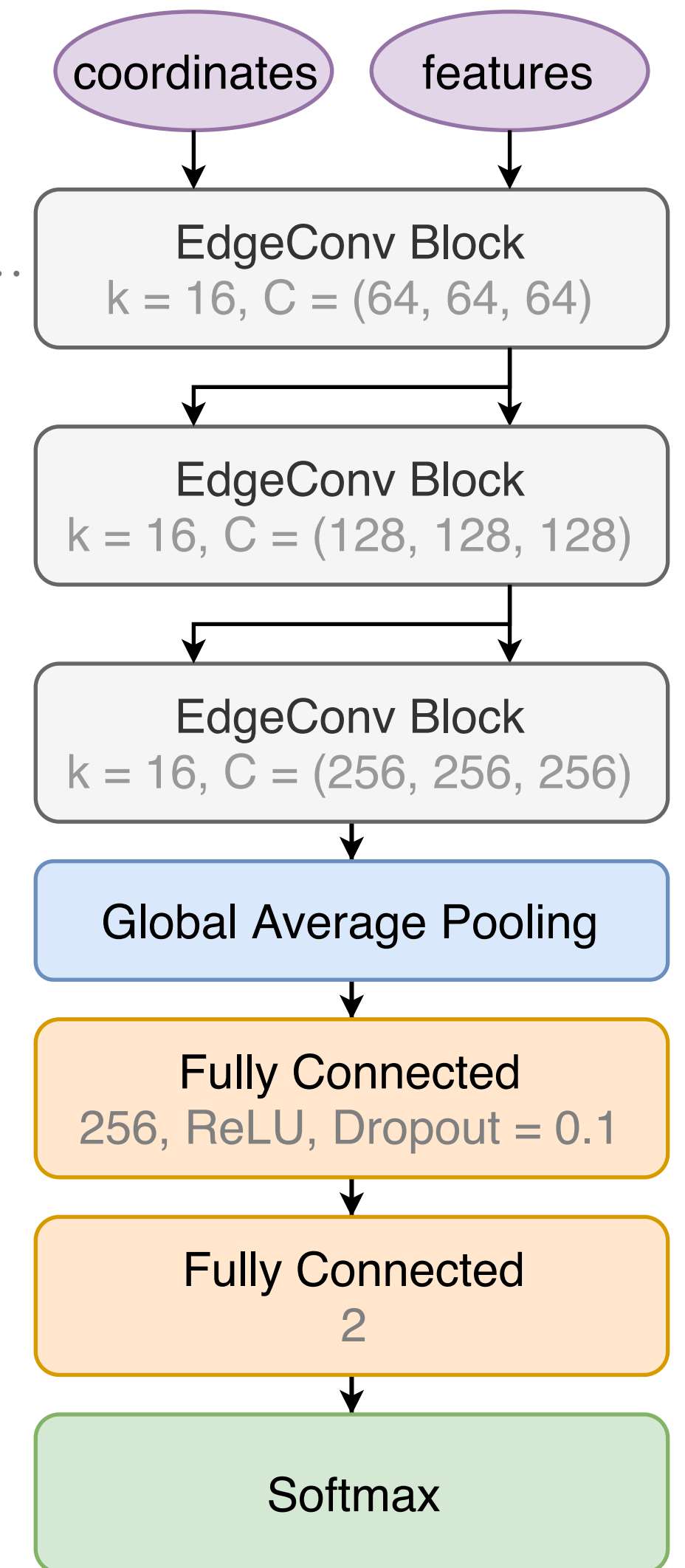


[Cogan, Kagan, Strauss, Schwartzman [JHEP 1502 \(2015\) 118](#)]

[de Oliveira, Kagan, Mackey, Nachman, Schwartzman [JHEP 1607 \(2016\) 069](#)]



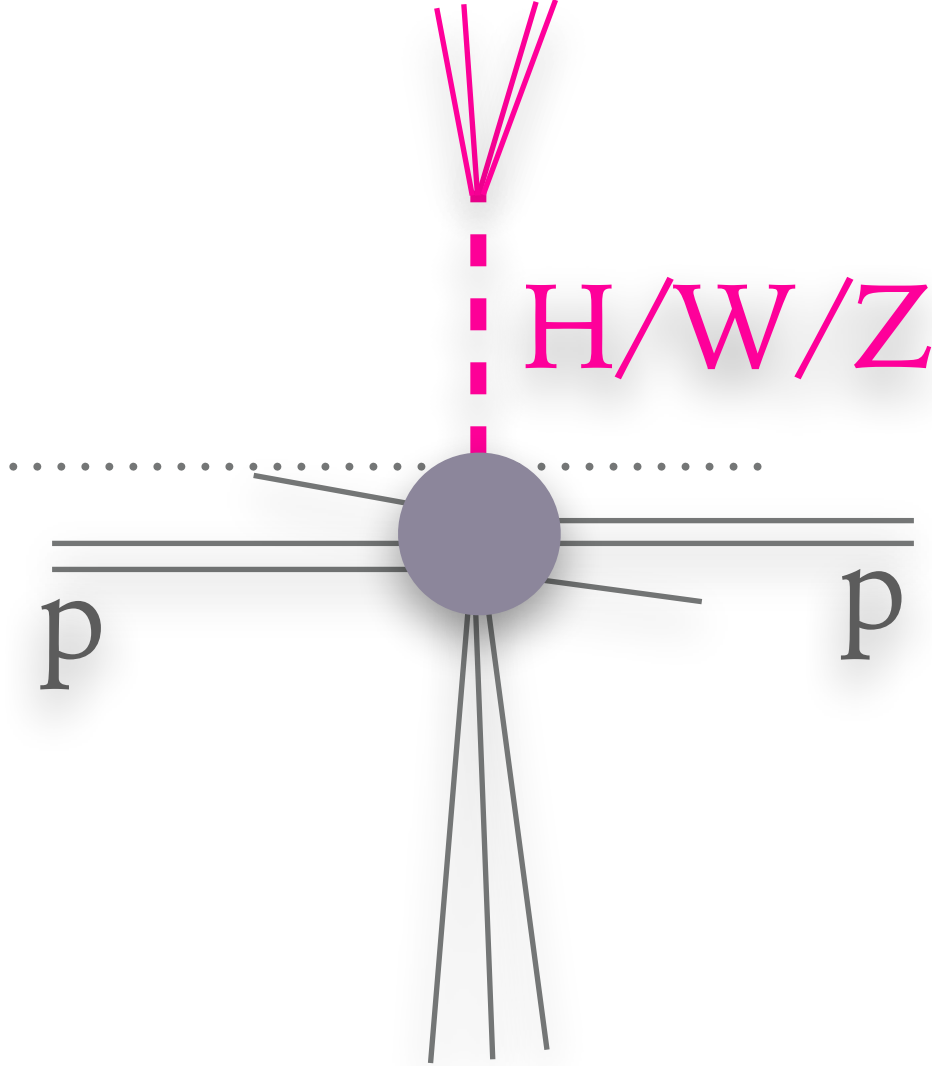
2021 Young Experimental Physicist Prize EPS HEPP prize



(a) ParticleNet

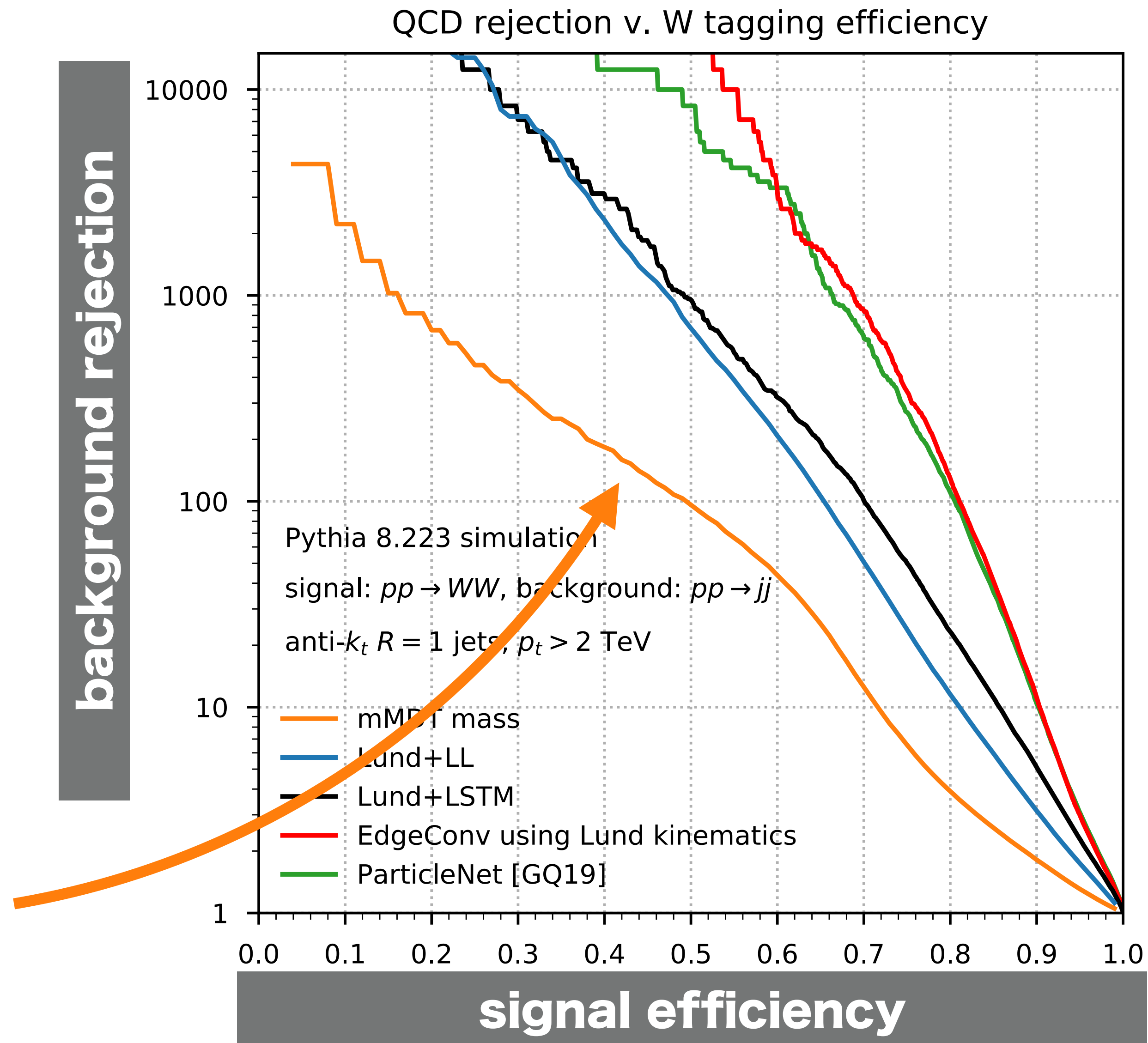
Qu & Guskos,
[arXiv:1902.08570](#)

using full jet/event information for H/W/Z-boson tagging

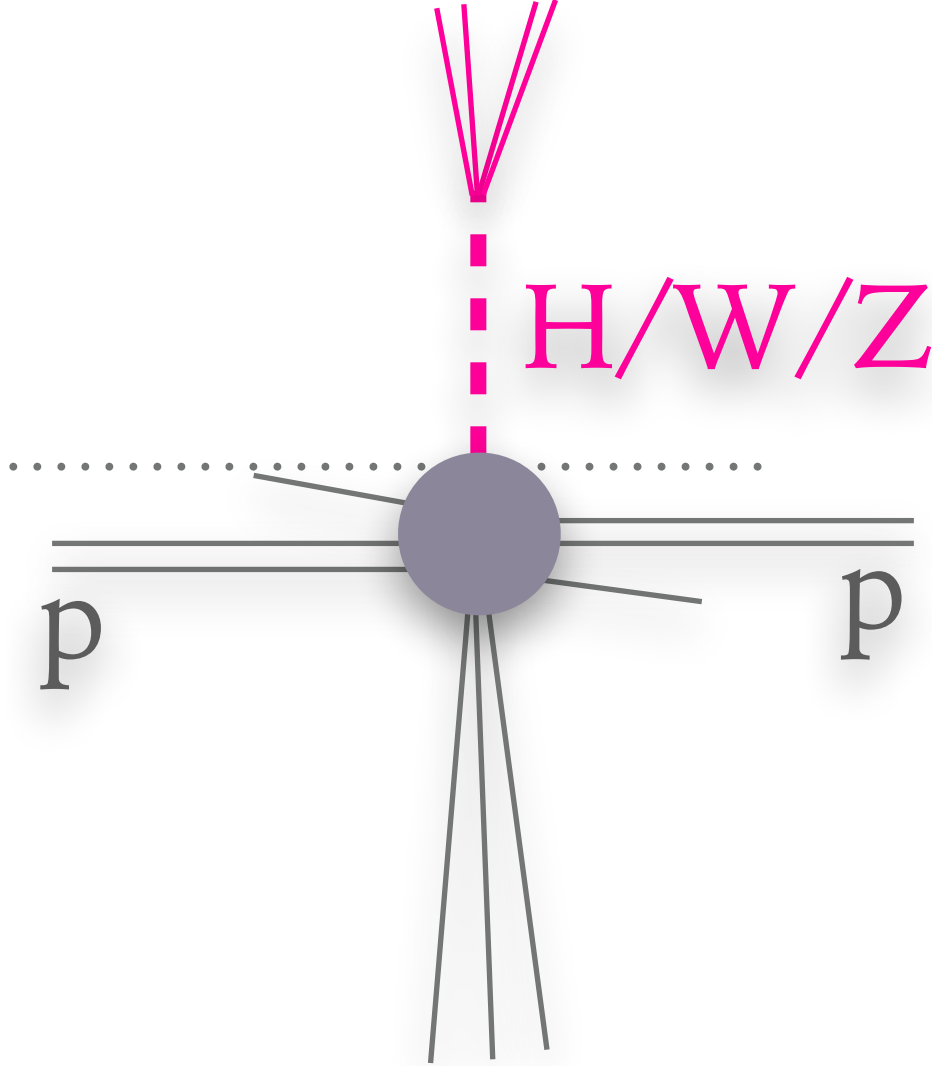


adapted from
Dreyer & Qu
2012.08526

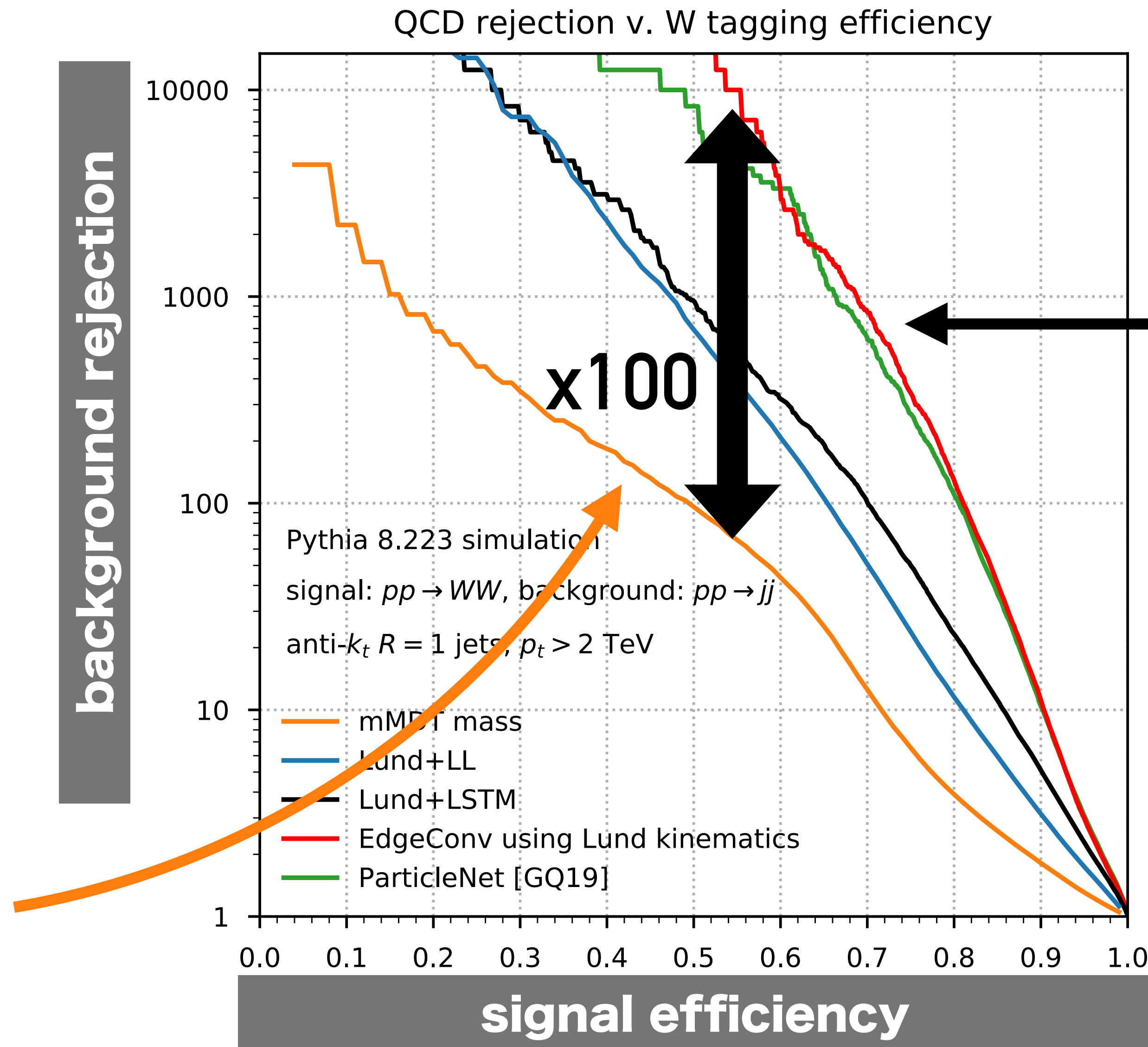
QCD rejection with
just jet mass
(SD/mMDT)
i.e. 2008 tools &
their 2013/14
descendants



using full jet/event information for H/W/Z-boson tagging



adapted from
Dreyer & Qu
2012.08526



QCD rejection
with use of full jet
substructure
(2021 tools)
100x better

QCD rejection with
just jet mass
(SD/mMDT)
i.e. 2008 tools &
their 2013/14
descendants

First started to be exploited
by Thaler & Van Tilburg with
“N-subjettiness” (2010/11)

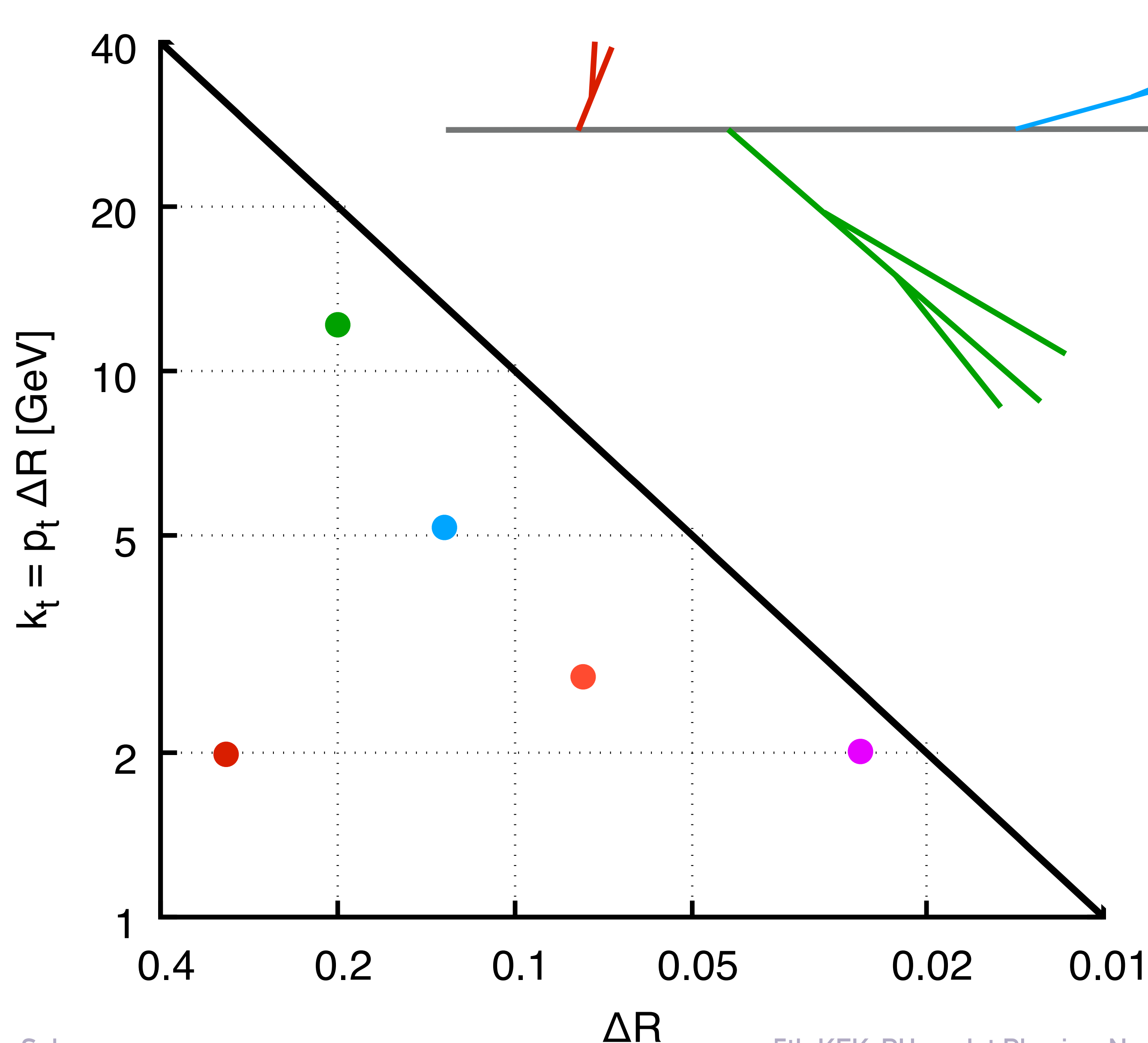
can we trust machine learning? A question of confidence in the training...

“

Unless you are highly confident in the information you have about the markets, you may be better off ignoring it altogether

*- Harry Markowitz (1990 Nobel Prize in Economics)
[via S Gukov]*

can we organise phase space to work for tagging and validation?



decluster particles in C/A alg. at successively smaller angles:
at each step record $\theta(= \Delta R)$, k_t
Lund plane & declustering

simple and robust

B. Andersson, G. Gustafson,
L. Lonnblad and Pettersson 1989
Dreyer, GPS & Soyez, [1807.04758](#)
& 5th heavy-ion workshop @ CERN, [1808.03689](#)

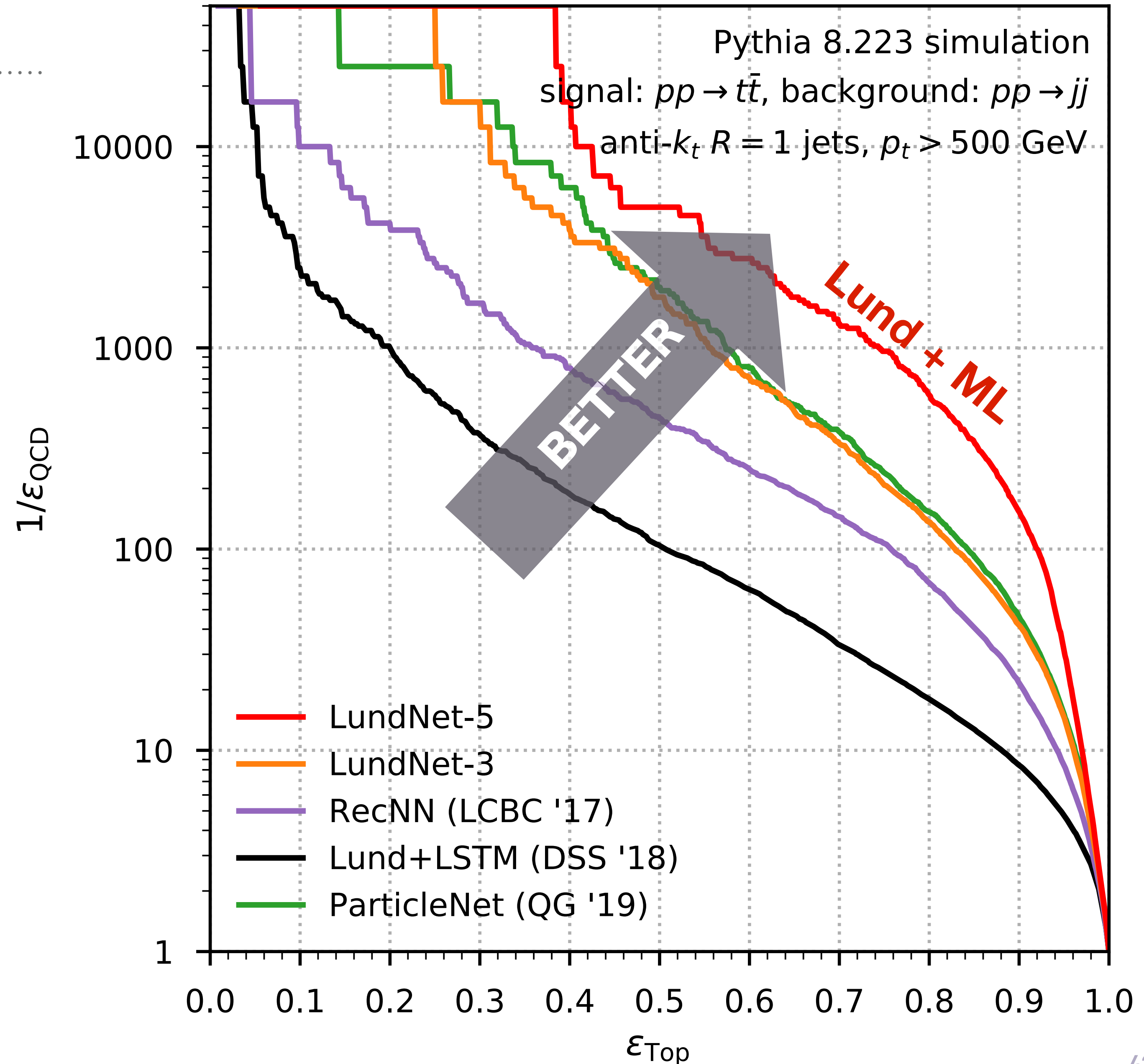
Use as input to machine-learning

Dreyer & Qu, [arXiv:2012.08526](https://arxiv.org/abs/2012.08526)

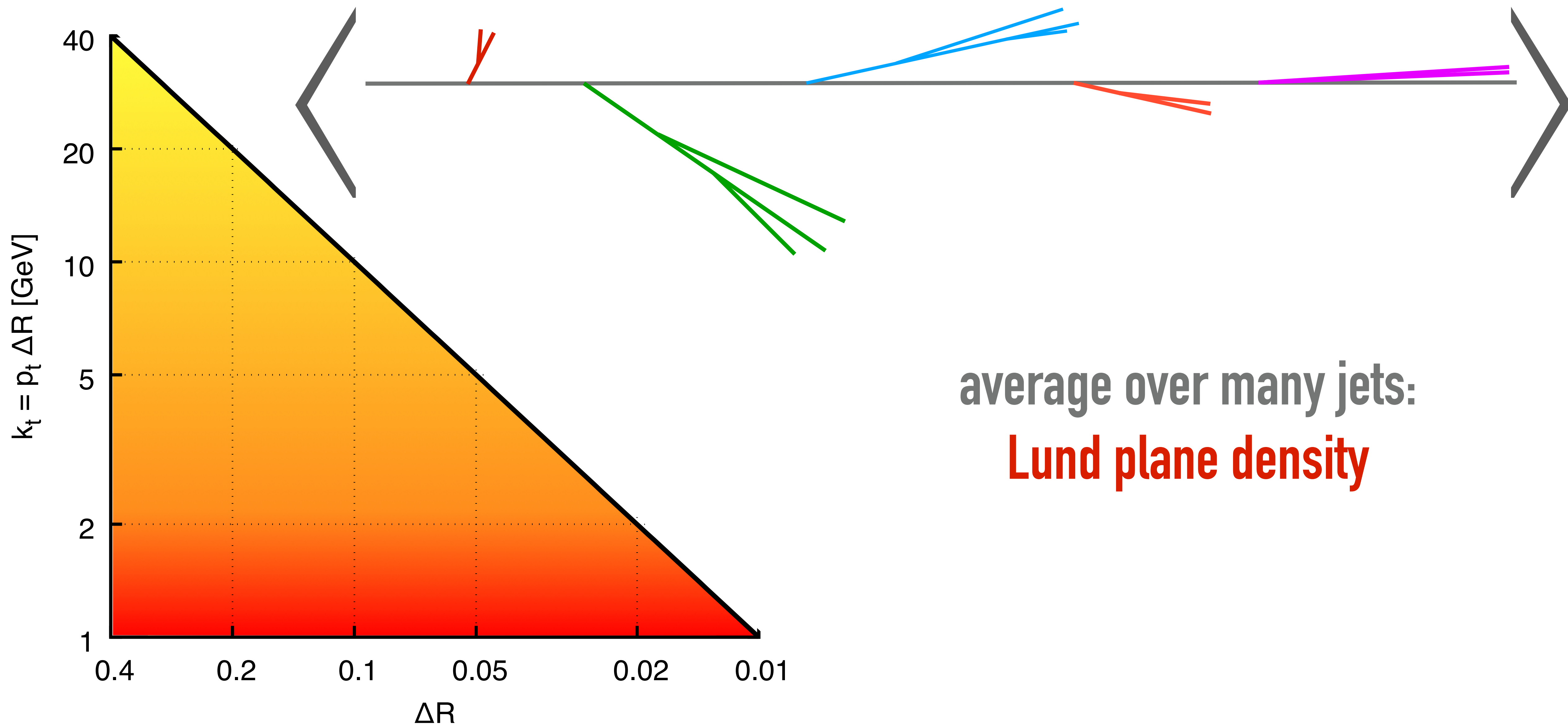
- ML with Lund inputs gives signal/background separation as good as, or better than other methods
- faster to train and to use
- these advantages probably come because the Lund diagram frames the physically relevant info in a way that makes it easier for machines to “learn”

	Number of parameters	Training time [ms/sample/epoch]	Inference time [ms/sample]
LundNet	395k	0.472	0.117
ParticleNet	369k	3.488	1.036
Lund+LSTM	67k	0.424	0.131

QCD rejection v. Top tagging efficiency



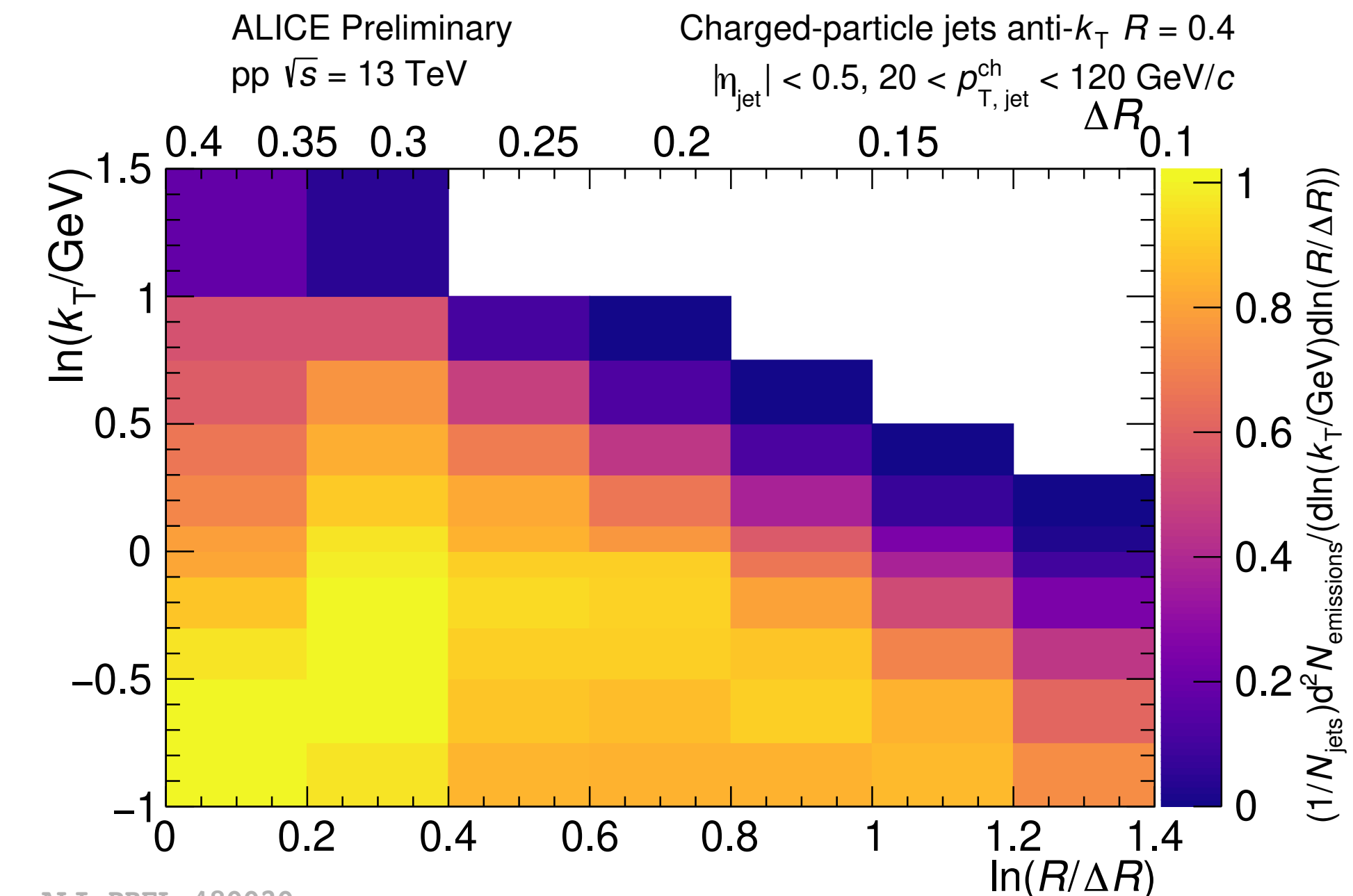
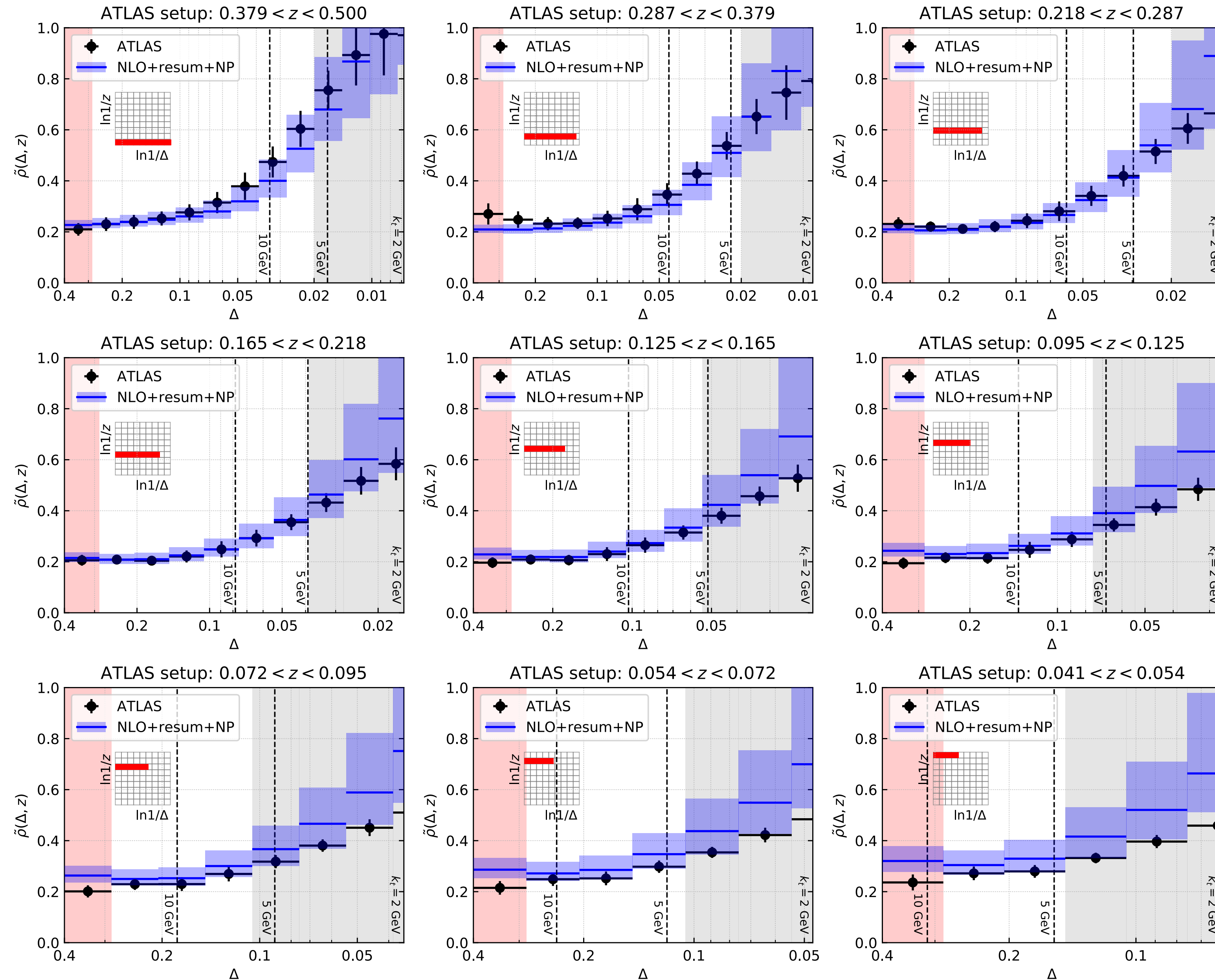
Same input used for ML can also be measured directly (validate / study QCD radⁿ pattern)



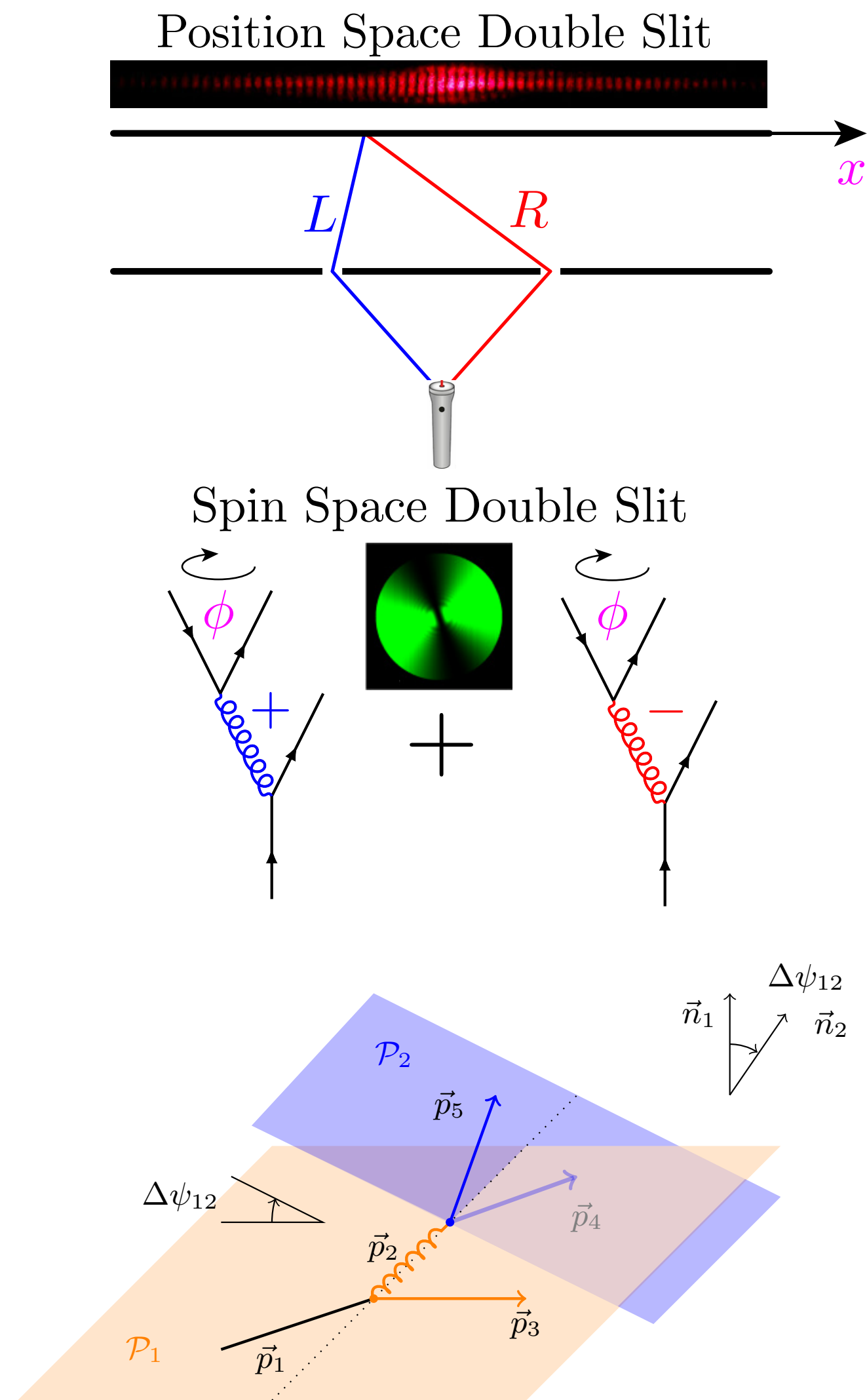
average over many jets:
Lund plane density

Lund plane: measured & calculated

- measurements by ATLAS & ALICE
- Good agreement between ATLAS & Lifson, GPS & Soyez, [arXiv:2007.06578](https://arxiv.org/abs/2007.06578)
- powerful tests also of Monte Carlos

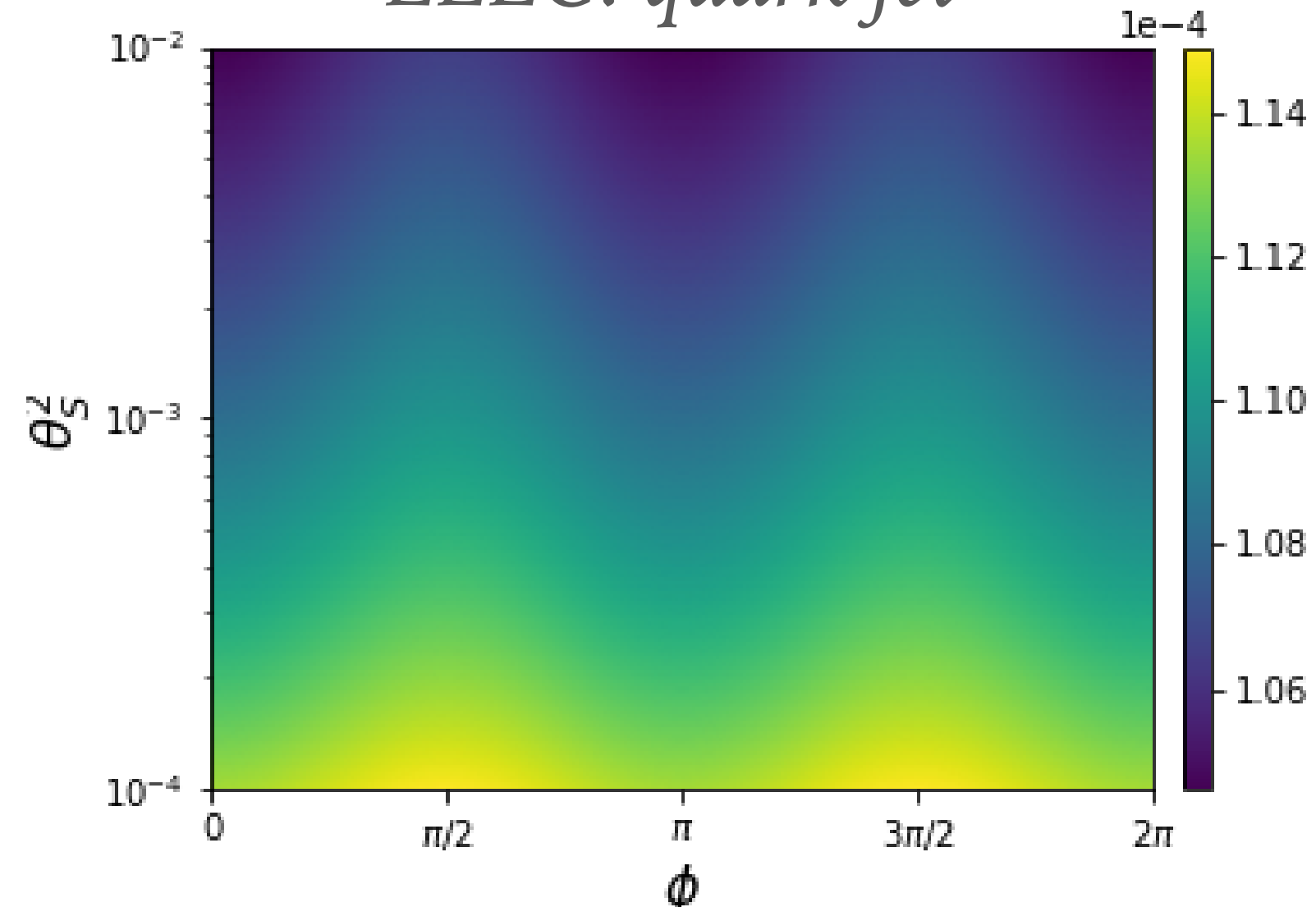


Collinear spin correlations within jets

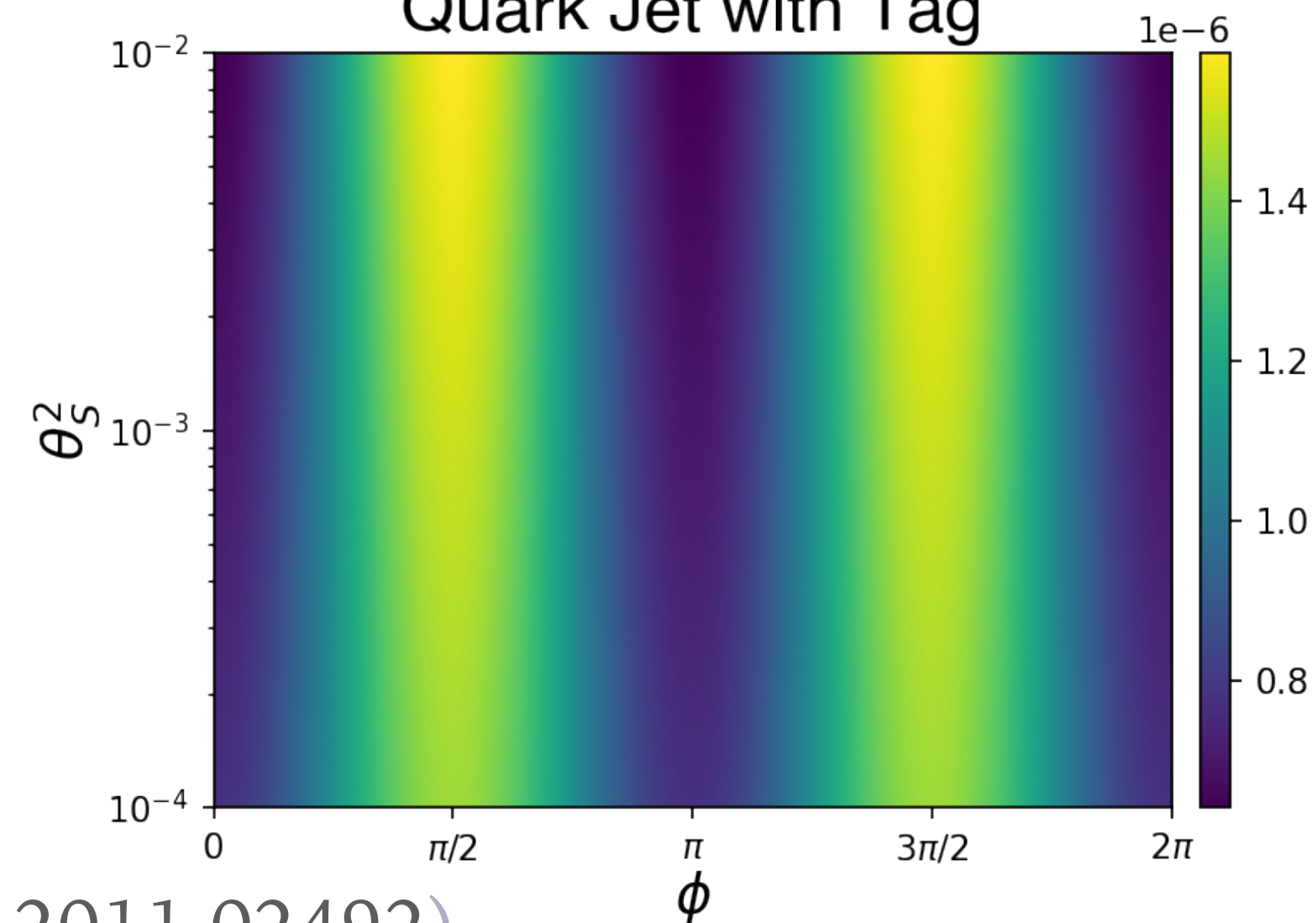


(Chen, Moult & Zhu, 2011.02492)

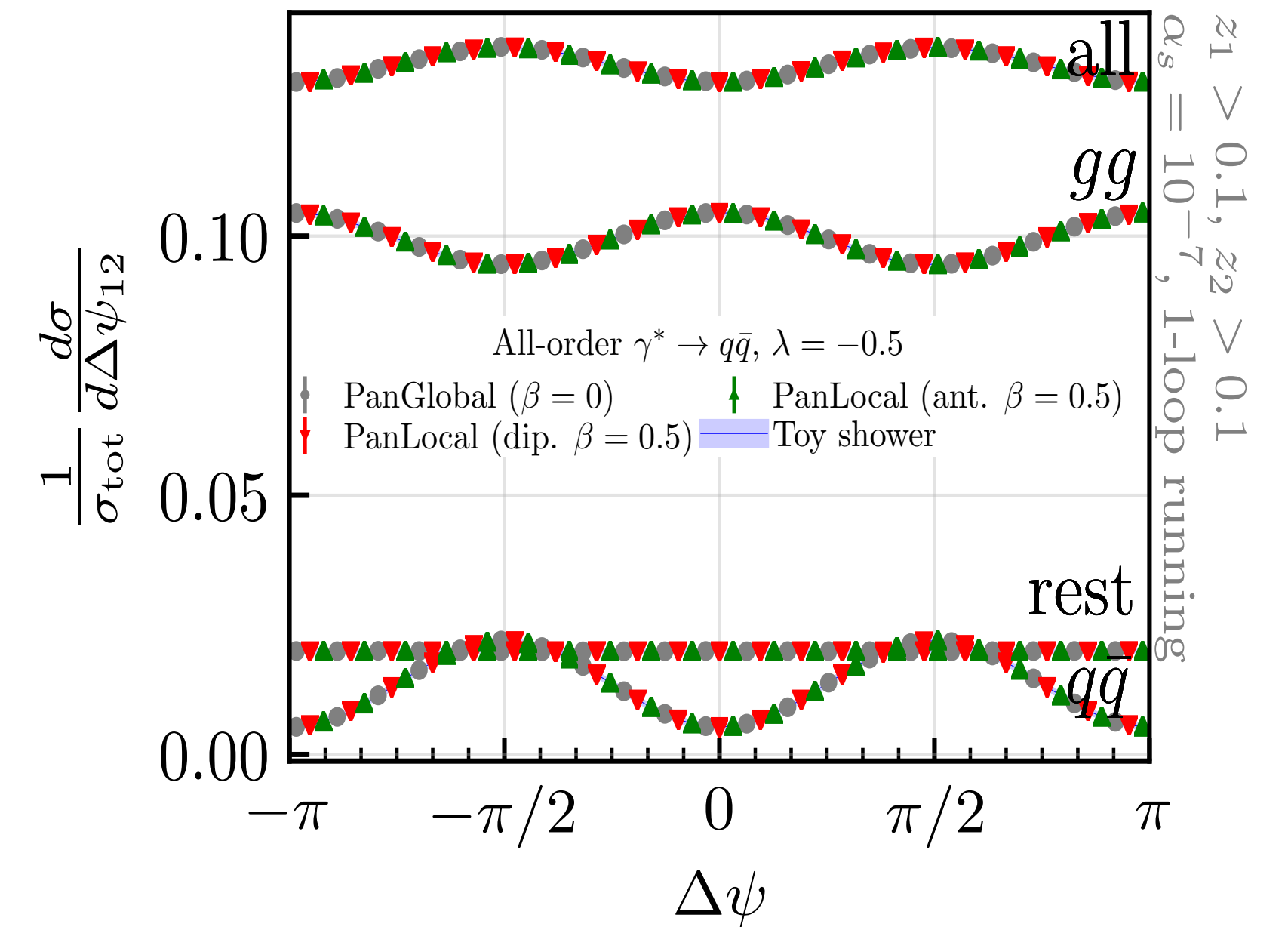
EEEC: quark jet



Quark Jet with Tag



Lund declustering $\Delta\psi_{12}$



magnitude of spin correlation effects

EEEC	-0.008
$\Delta\psi_{12}, z_1, z_2 > 0.1$	-0.025
$\Delta\psi_{12}, z_1 > 0.1, z_2 > 0.3$	-0.042

Karlberg, GPS, Scyboz &
Verheyen, 2103.16526

conclusions

Conclusions

- Jets are a crucial part of collider physics
Including broad programme to study new Higgs interactions and search for BSM physics
- Basic jet finding (1 quark = 1 jet) has simple, fast tools (anti- k_t , FastJet) that continue to work well 10-15 years since their inception
- Incredible how much information is hiding in jet substructure — every couple of years, people find that there is yet more info to be extracted
 - Lund declustering is one physical, powerful way of doing that (another is energy-flow polynomials)
- Will undoubtedly play major role in next 15 years of LHC, and at future $e^+e^-/pp/\mu\mu$ colliders
- The challenge is also on to make sure we can reliably predict the internal structure of jets and so make confident use of the associated information