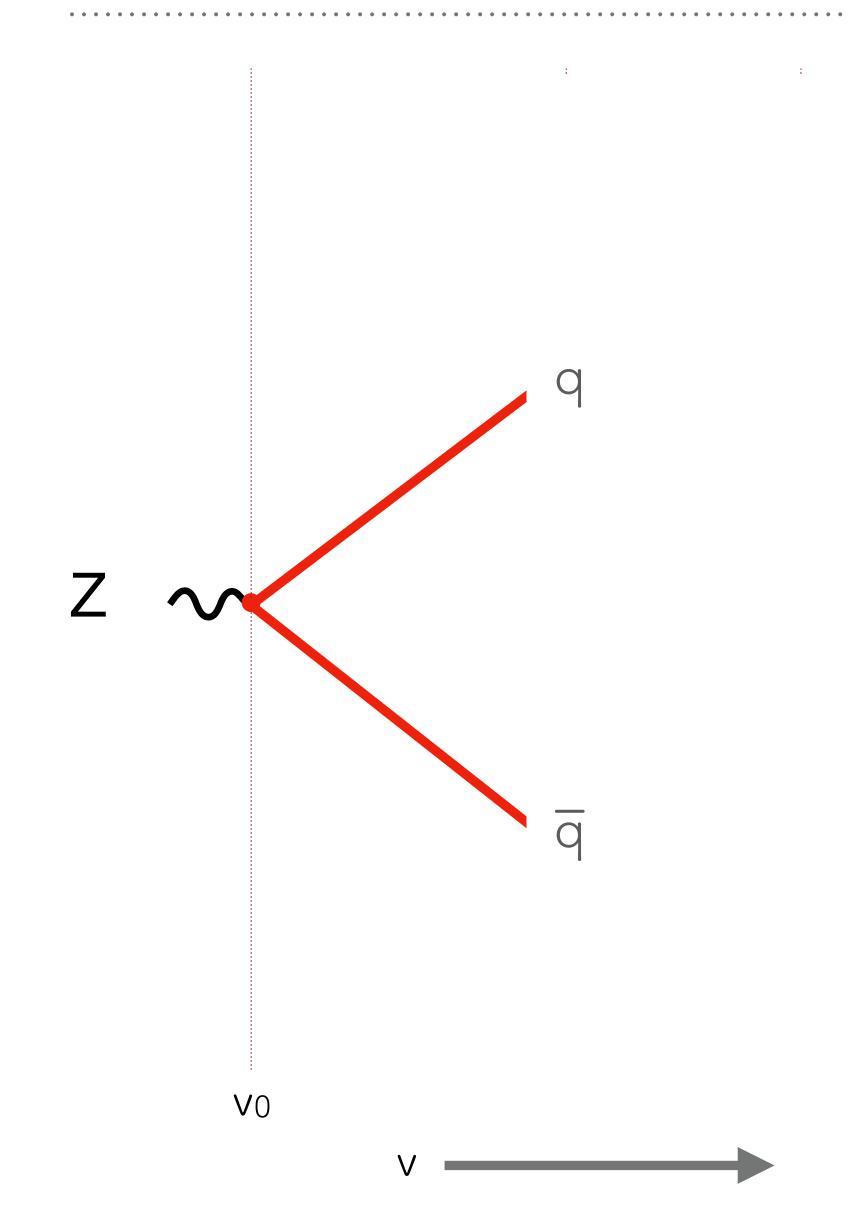


V

Start with  $q\bar{q}$  state.

Throw a random number to determine down to what scale state persists unchanged

$$\frac{dP_2(v)}{dv} = -f_{2\to 3}^{q\bar{q}}(v) P_2(v)$$



Start with q-qbar state.

Throw a random number to determine down to what scale state persists unchanged

$$\frac{dP_2(v)}{dv} = -f_{2\to 3}^{q\bar{q}}(v) P_2(v)$$

V0

Start with q-qbar state.

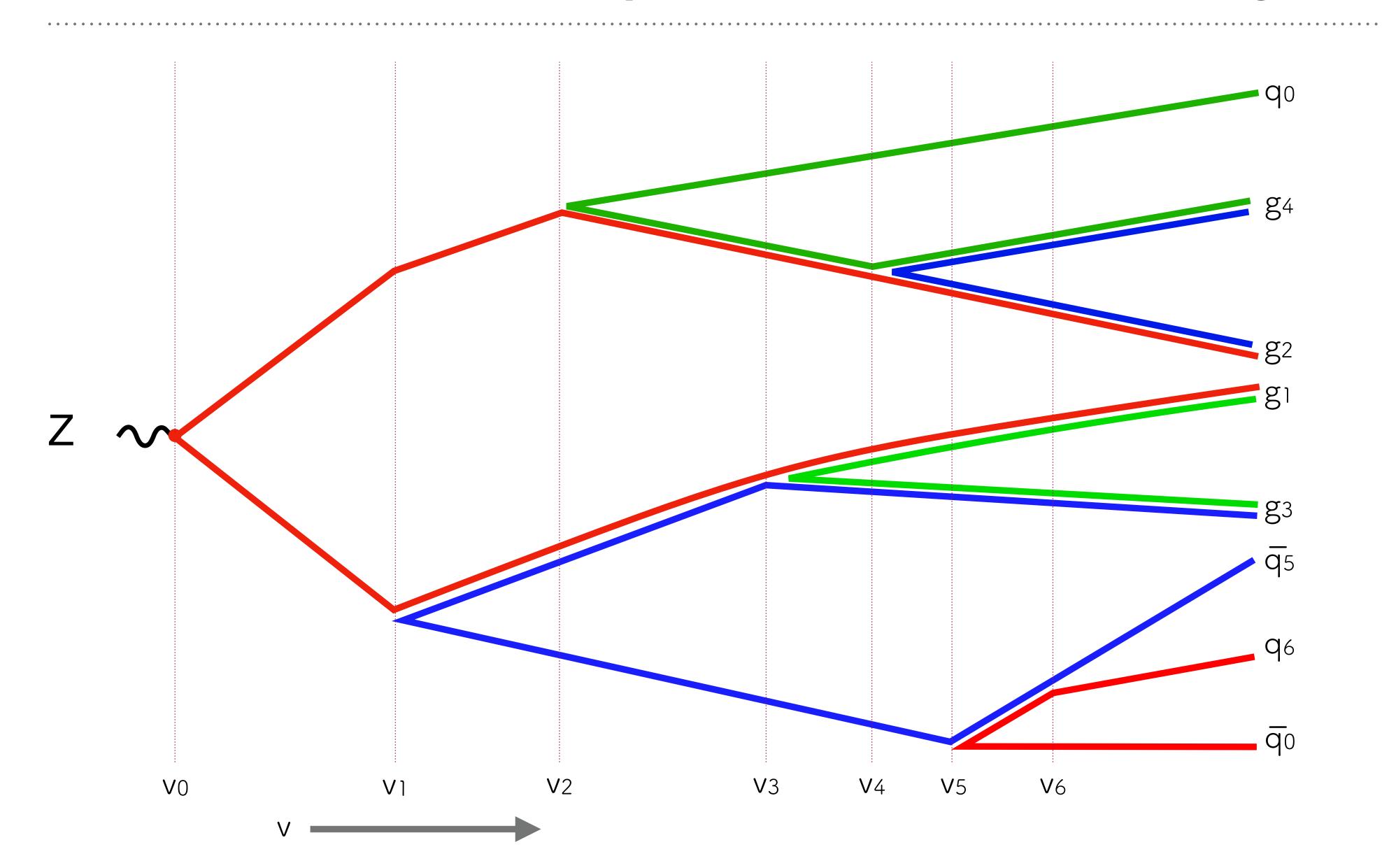
Throw a random number to determine down to what scale state persists unchanged

At some point, state splits  $(2\rightarrow 3, i.e. emits gluon)$ . Evolution equation changes

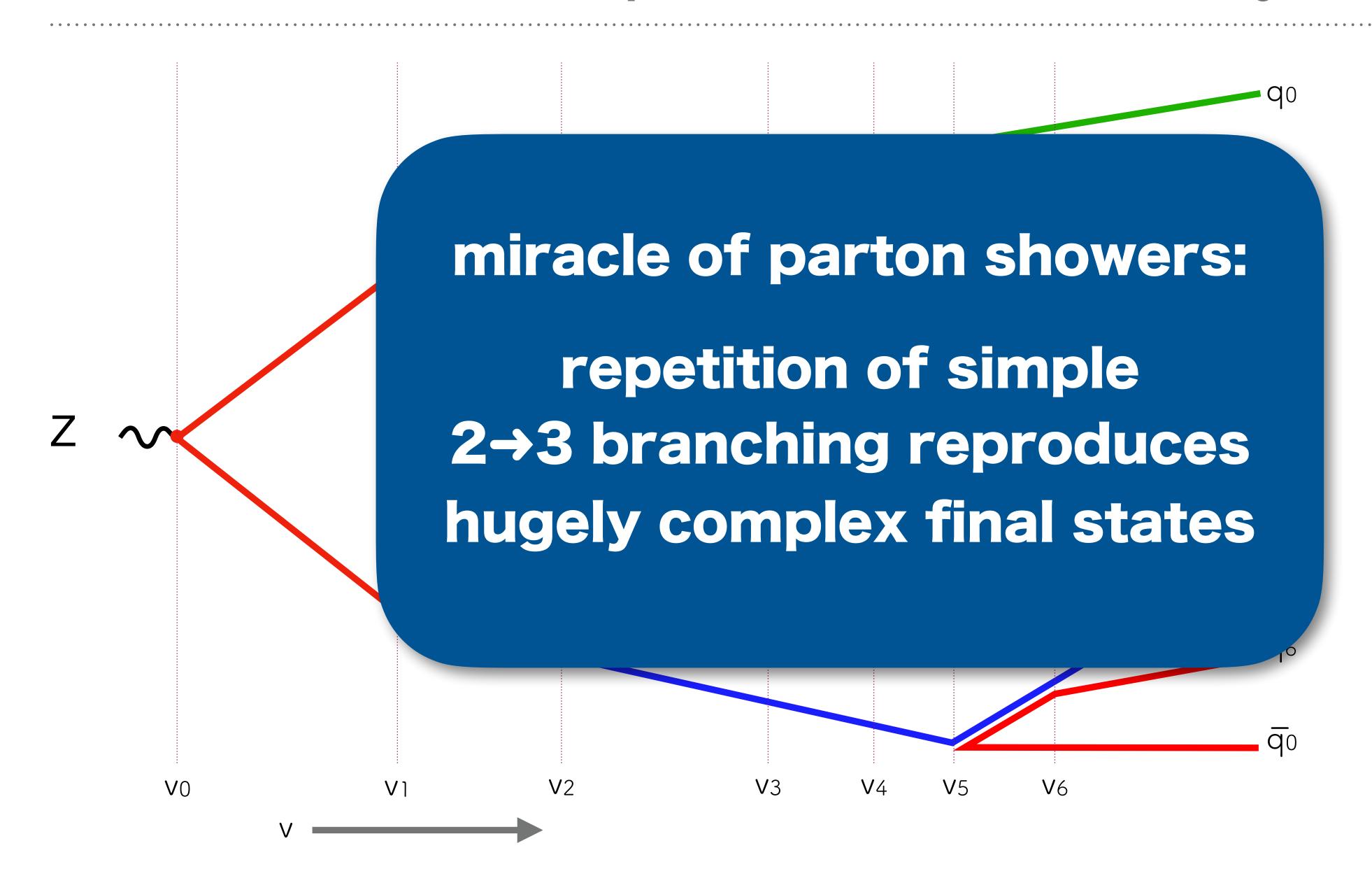
$$\frac{dP_3(v)}{dv} = -\left[f_{2\to 3}^{qg}(v) + f_{2\to 3}^{g\bar{q}}(v)\right] P_3(v)$$

gluon is part of two dipoles (qg),  $(g\bar{q})$ , each treated as independent

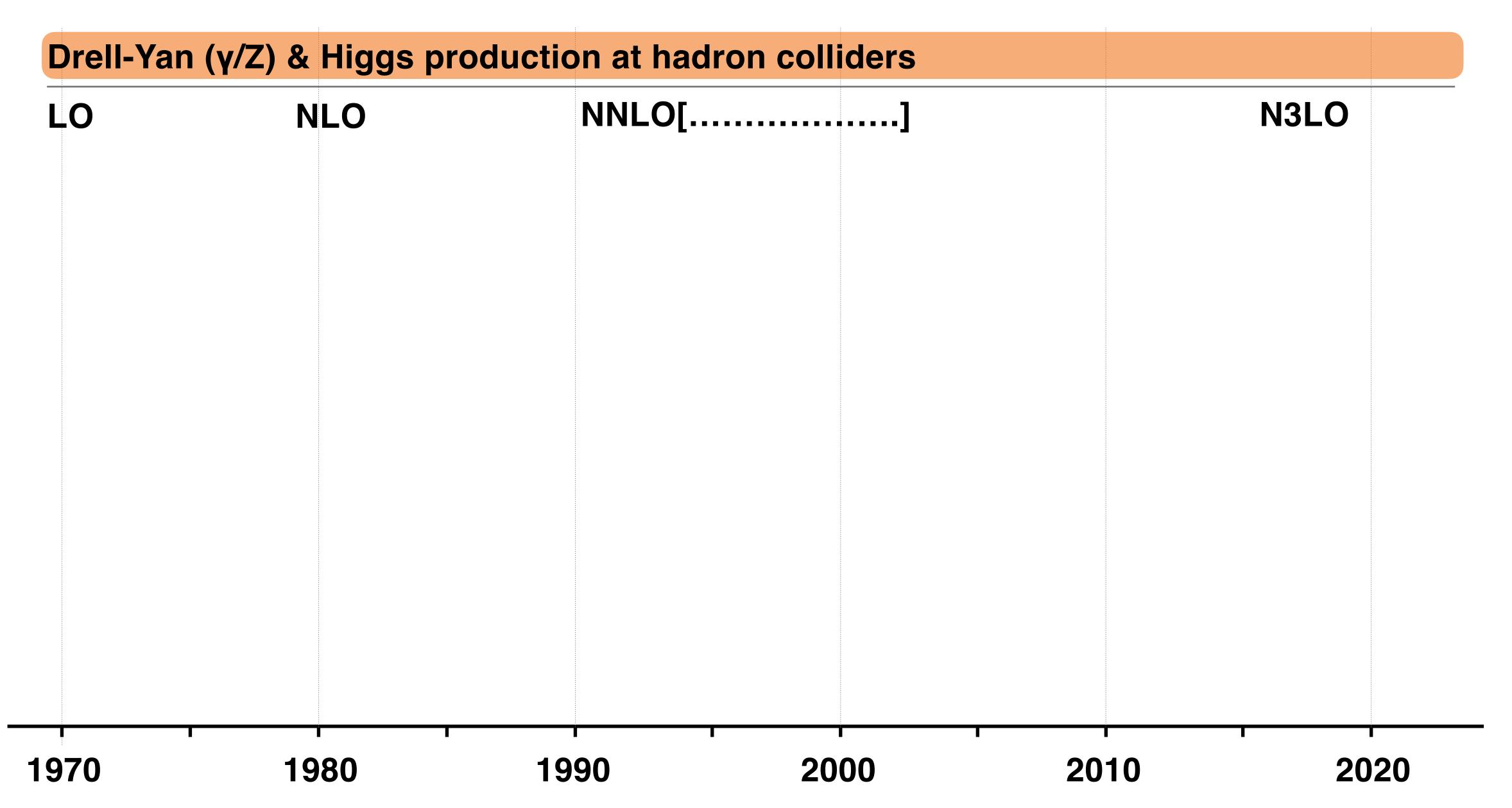
(many showers use a large N<sub>C</sub> limit)

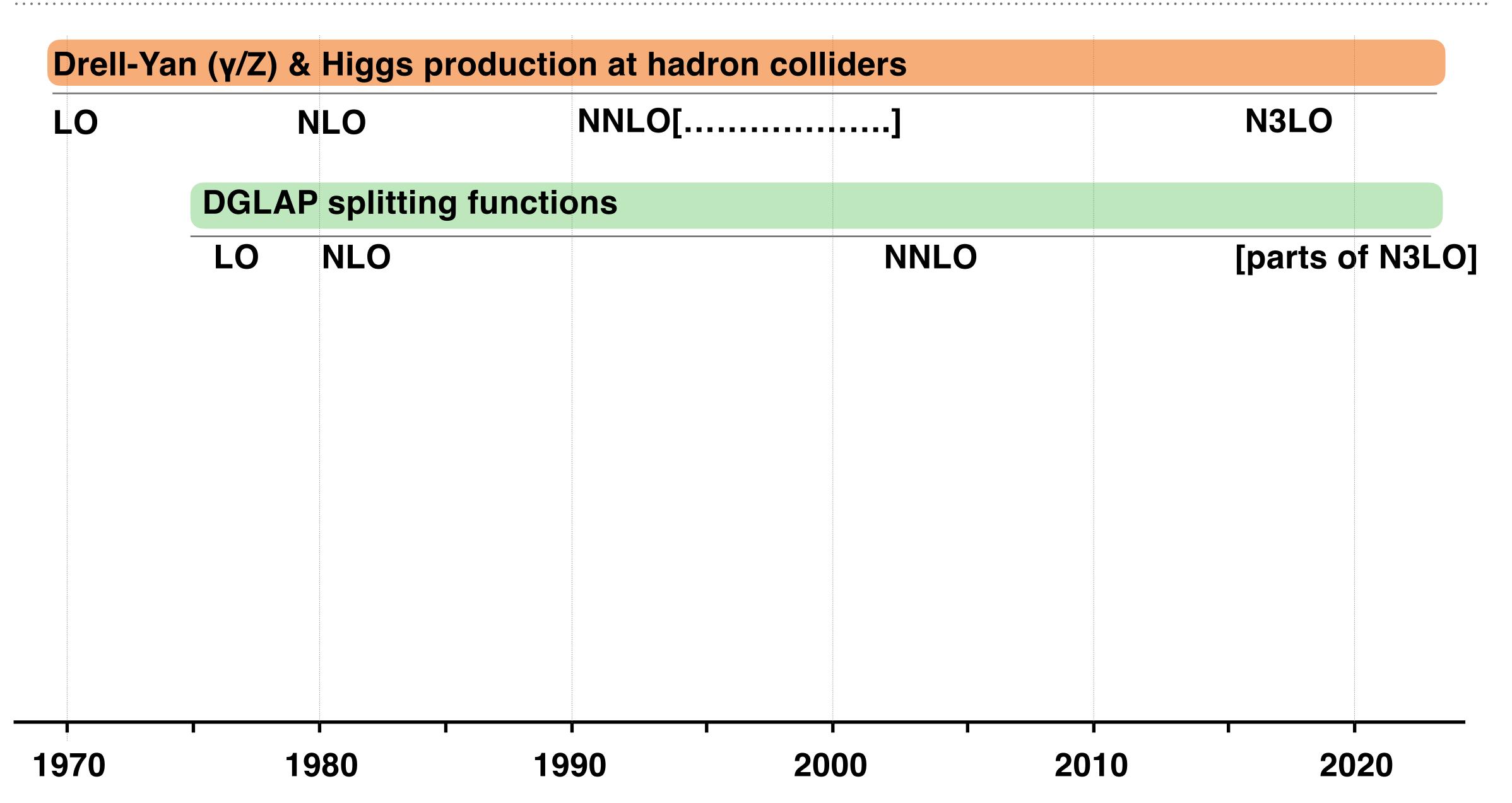


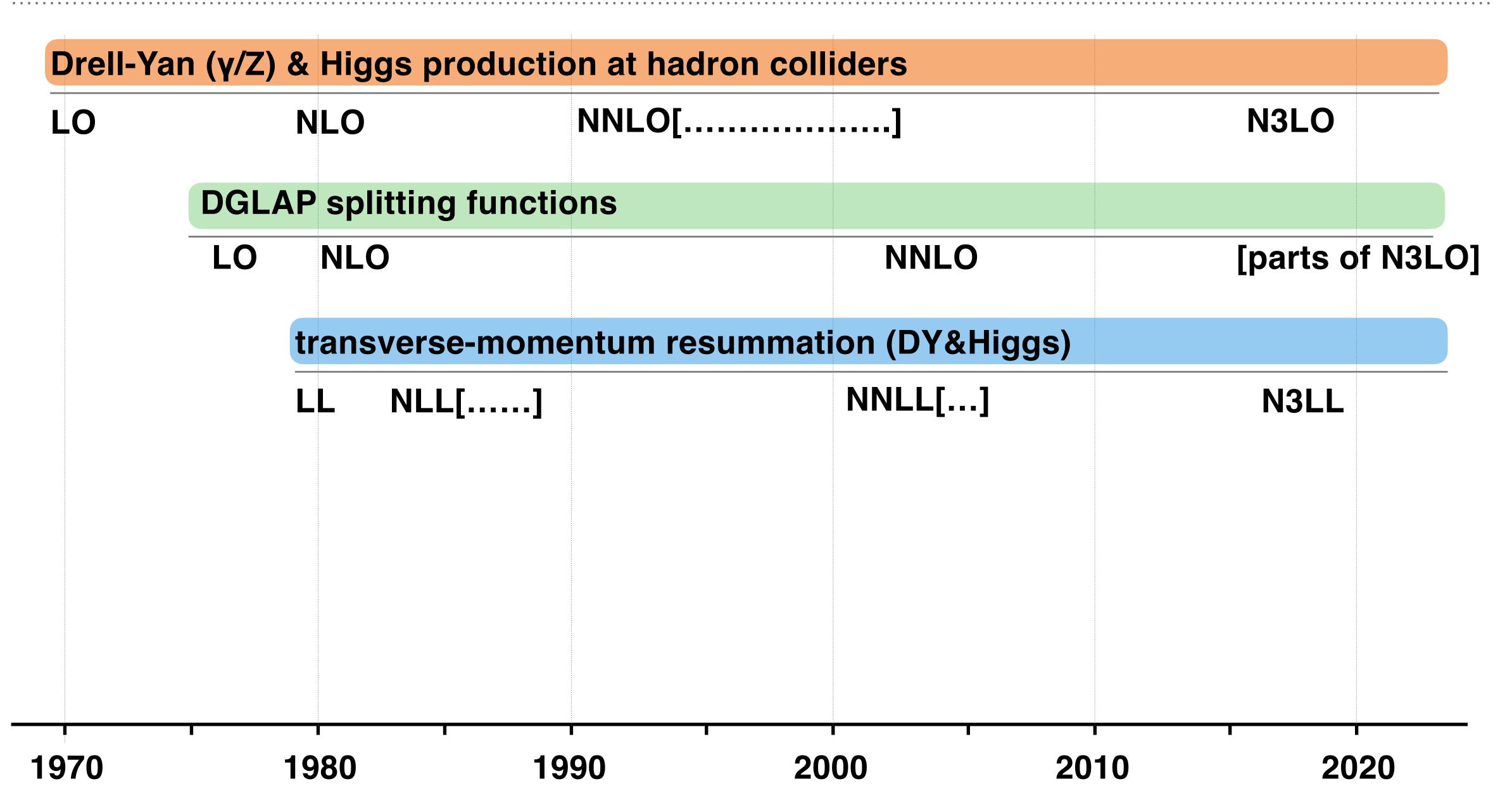
self-similar
evolution
continues until it
reaches a nonperturbative
scale

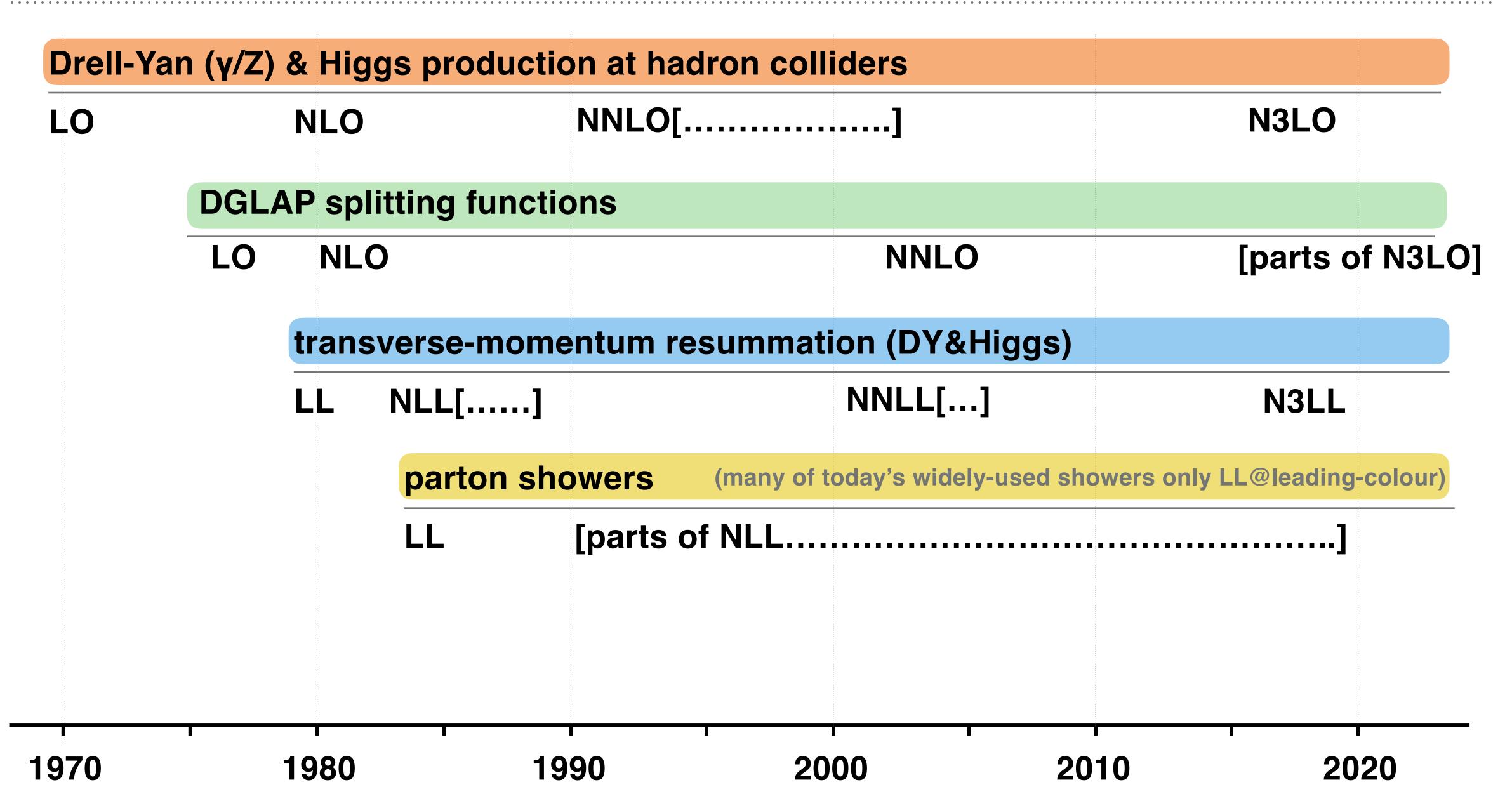


self-similar
evolution
continues until it
reaches a nonperturbative
scale



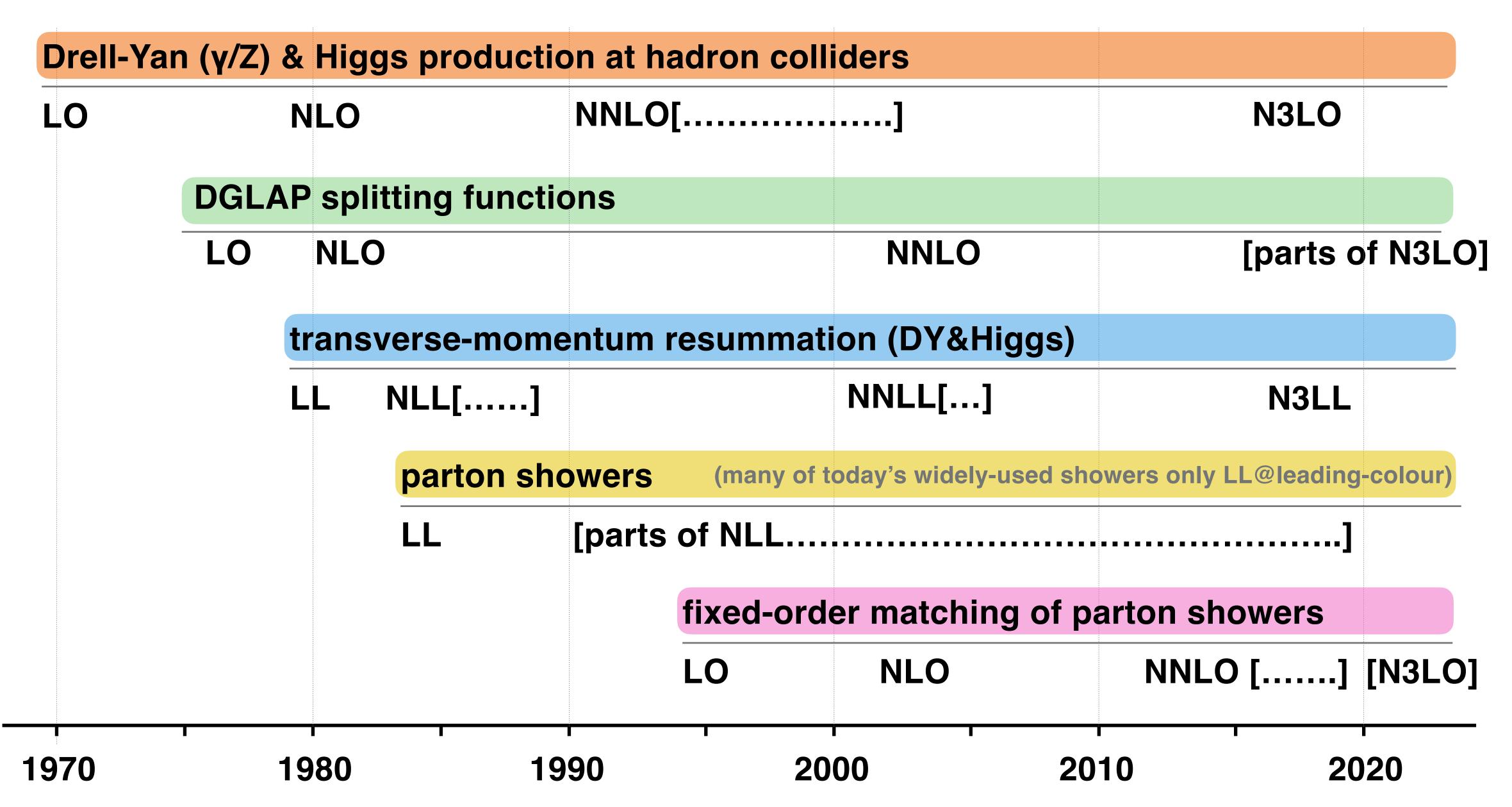


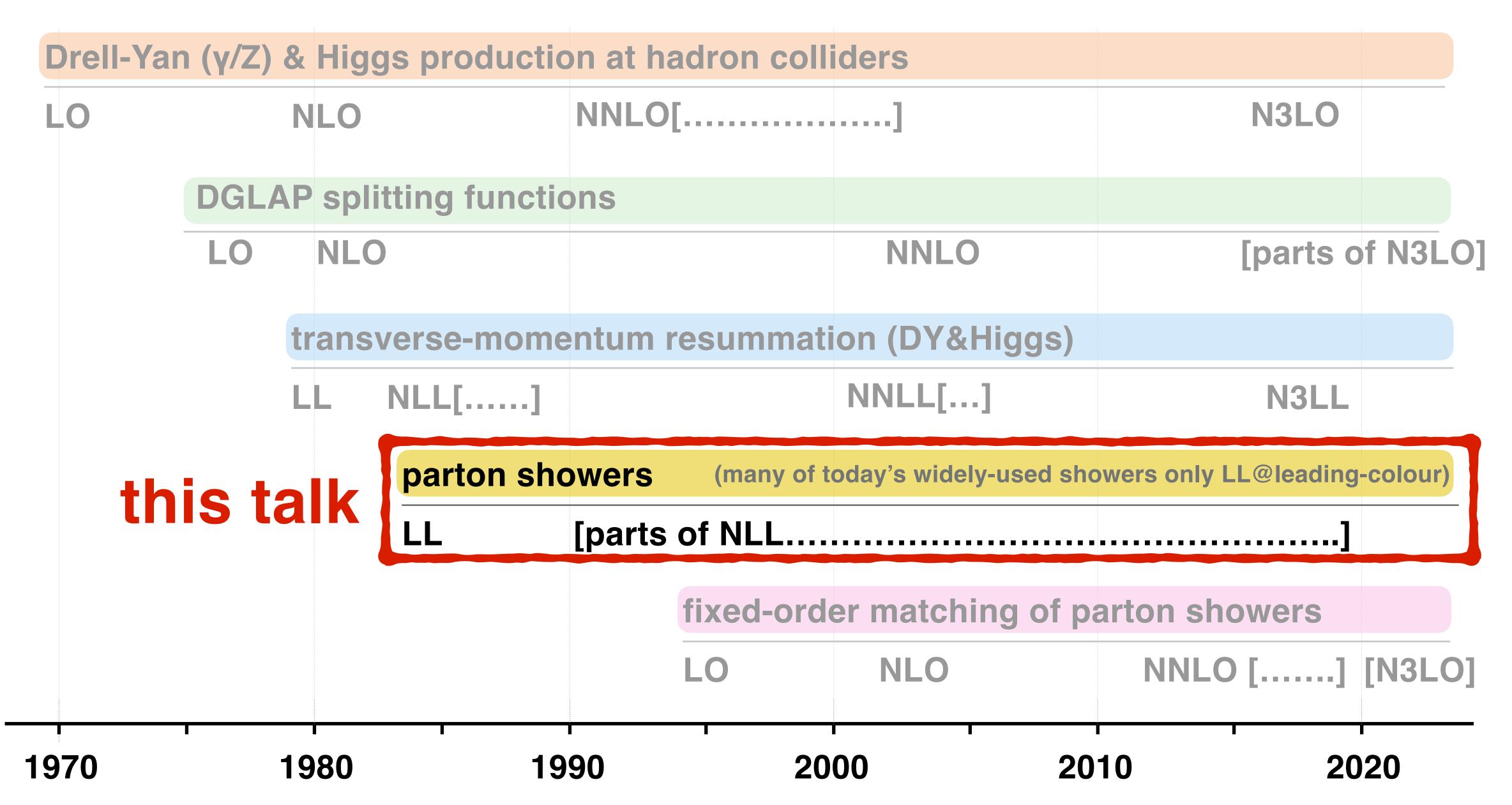




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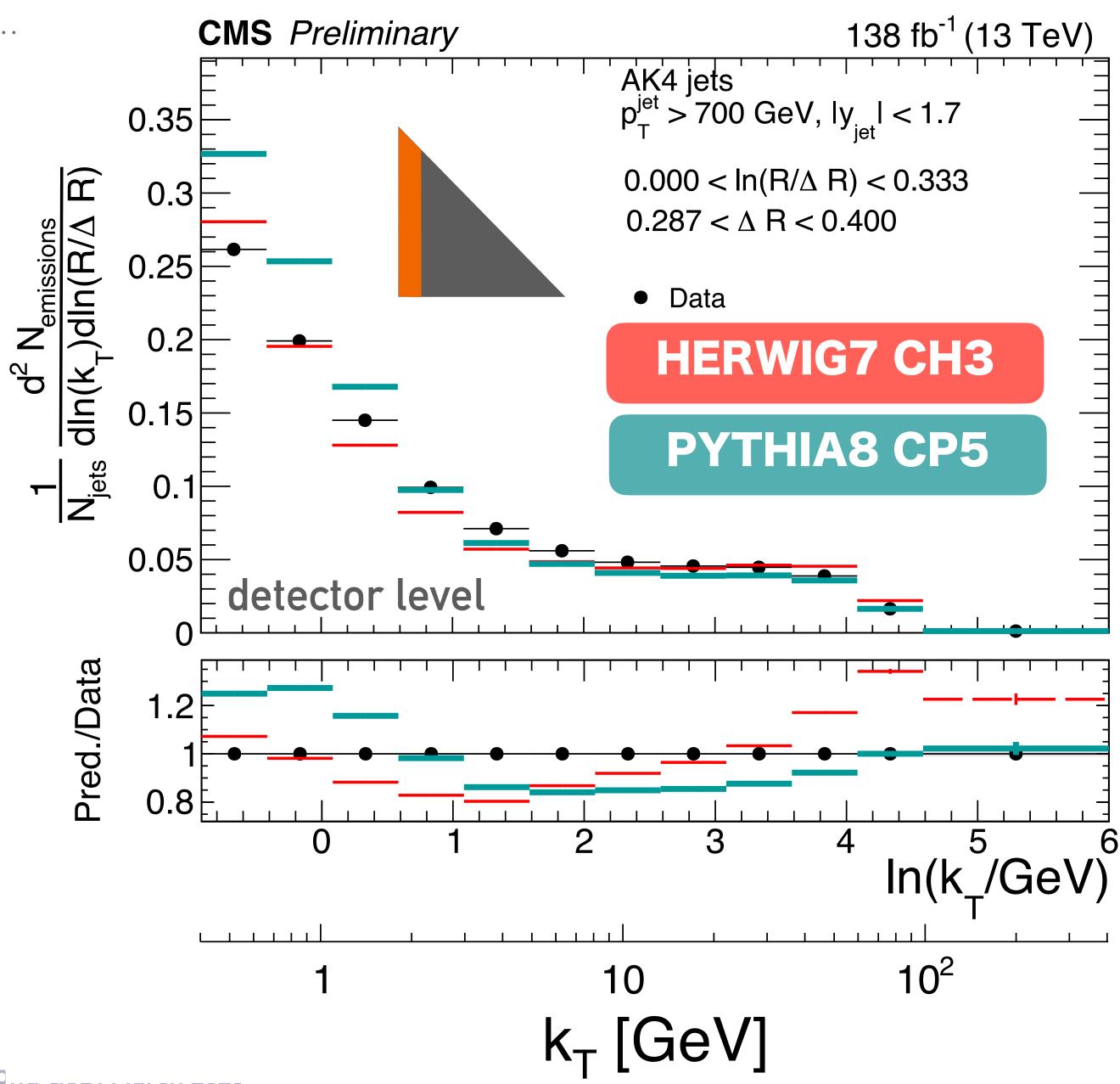




# Are showers good enough?

- > showers do an amazing job on many observables
- ➤ but various places see 10–30% discrepancies between showers and data
- > feeds into many analyses (e.g. via jet-energy scale)
- > as machine learning makes use of ever more information in jets & whole event, we want simulations to get it right





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# We want to design guaranteed NLL showers

#### A Matrix Element condition

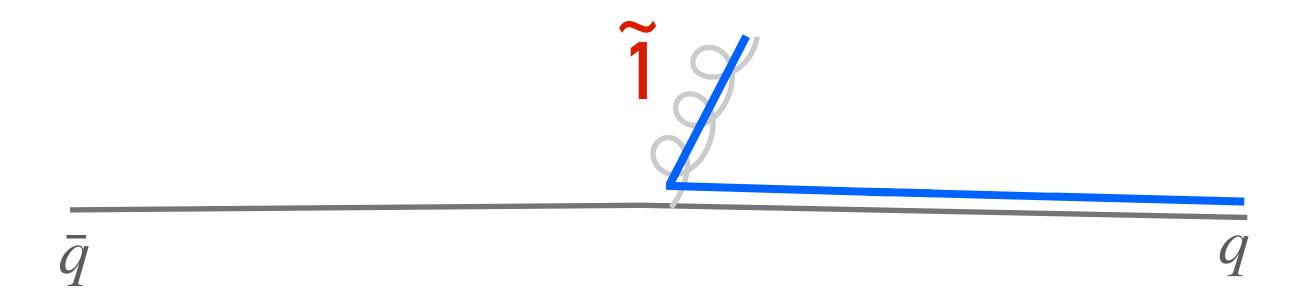
- ➤ correctly reproduce *n*-parton tree-level matrix element for arbitrary configurations, so long as all emissions well separated in the Lund diagram
- > supplement with unitarity, 2-loop running coupling & cusp anomalous dimension

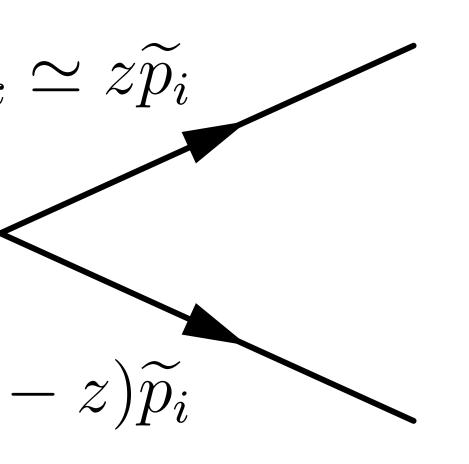
Resummation condition: reproduce NLL results for all standard resummations

- > global event shapes
- non-global observables
- ➤ fragmentation functions
- multiplicities



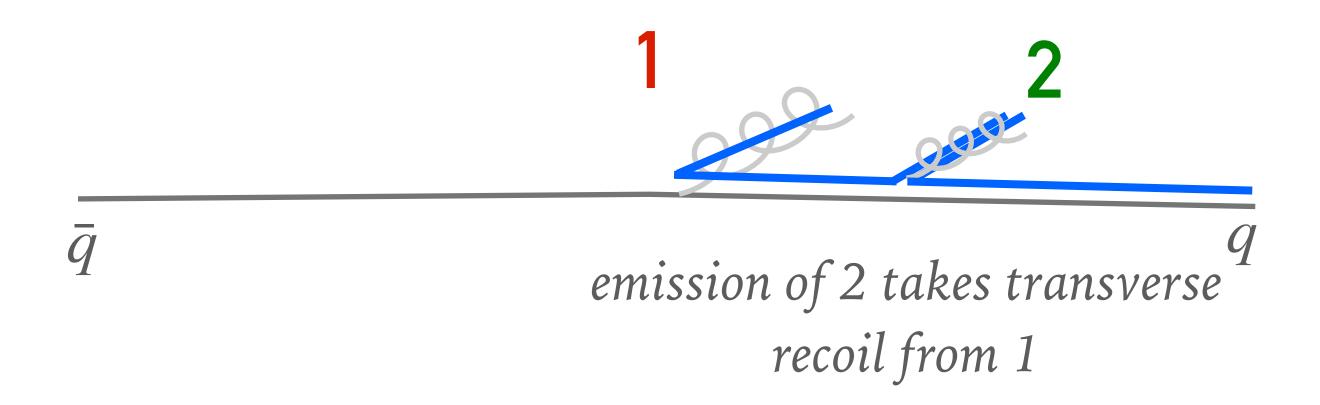
Dipole showers conserve momentum at each step. Traditional dipole-local recoil:





$$d\mathcal{P}_{\tilde{\imath}\to ik}^{\mathrm{FS}} = \frac{\alpha_s(k_{\perp}^2)}{2\pi} \frac{dk_{\perp}^2}{k_{\perp}^2} \frac{dz}{z} \frac{d\varphi}{2\pi} N_{ik}^{\mathrm{sym}} \left[ z P_{\tilde{\imath}\to ik}(z) \right]$$

Dipole showers conserve momentum at each step. Traditional dipole-local recoil:

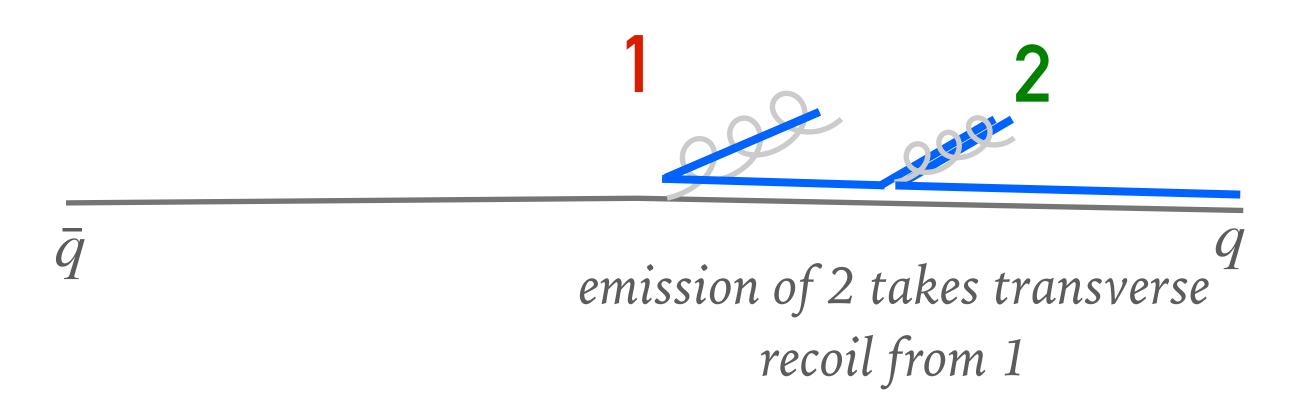


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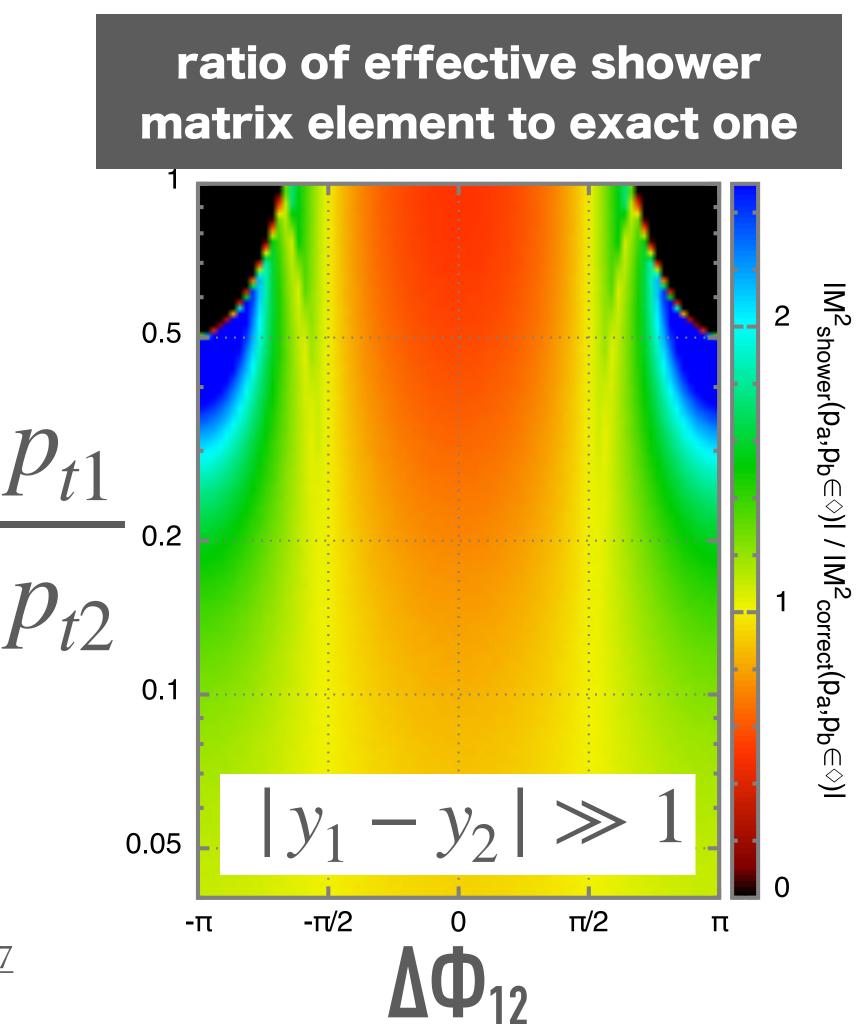
Dipole showers conserve momentum at each step. Traditional dipole-local recoil:



Shower initially generated matrix element for particle  $\tilde{1}$ , whose momentum differs (by  $\sim 50\%$ ) from final particle 1.

#### Matrix element is incorrect wrt final momentum 1.

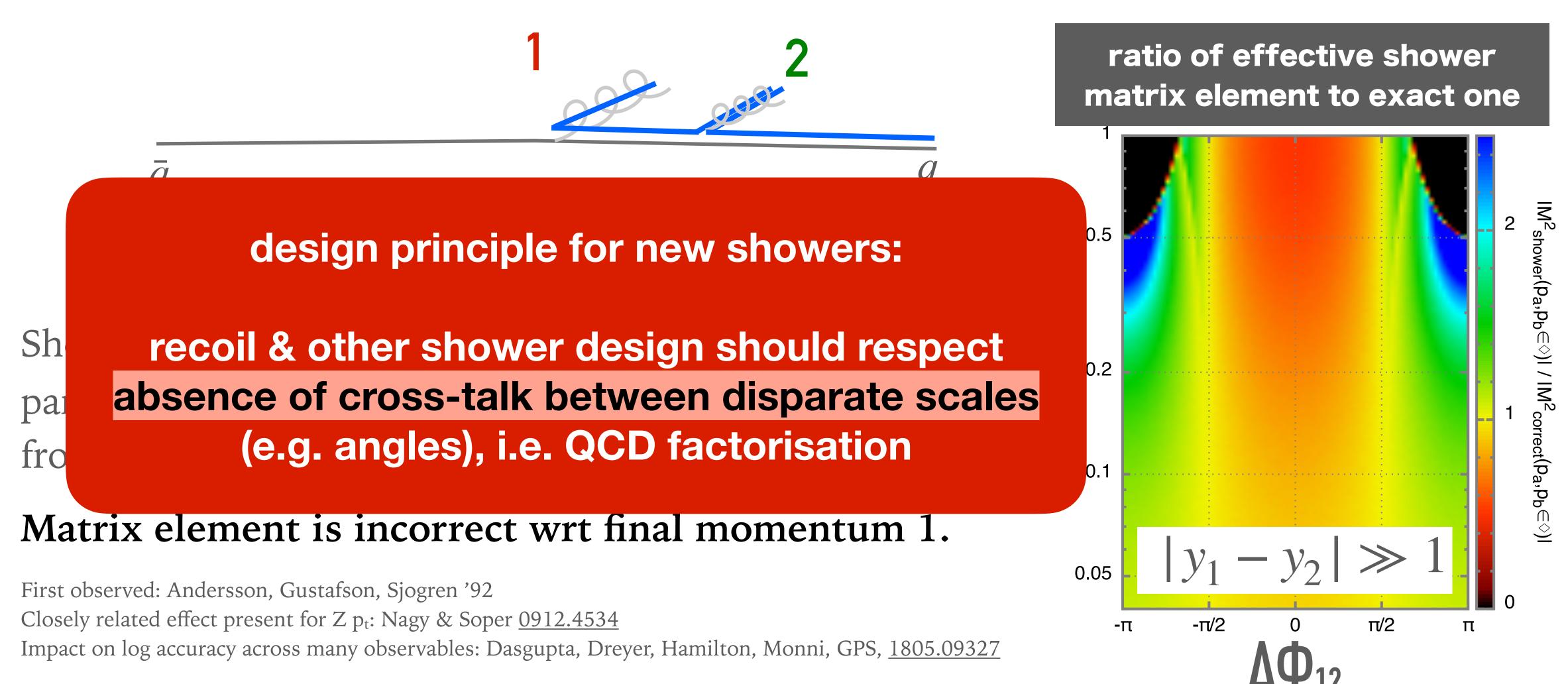
First observed: Andersson, Gustafson, Sjogren '92 Closely related effect present for Z p<sub>t</sub>: Nagy & Soper <u>0912.4534</u> Impact on log accuracy across many observables: Dasgupta, Dreyer, Hamilton, Monni, GPS, <u>1805.09327</u>



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Dipole showers conserve momentum at each step. Traditional dipole-local recoil:



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Melissa van Beekveld



Jack Helliwell



**Rok Medves** 



**Frederic Dreyer** 



**GPS** 



Ludo Scyboz

## UCL



**Keith Hamilton** 



Rob Verheyen

## Manchester



CERN



**Mrinal Dasgupta** 



**Gregory Soyez** 



**Pier Monni** 







Alba Soto Ontoso

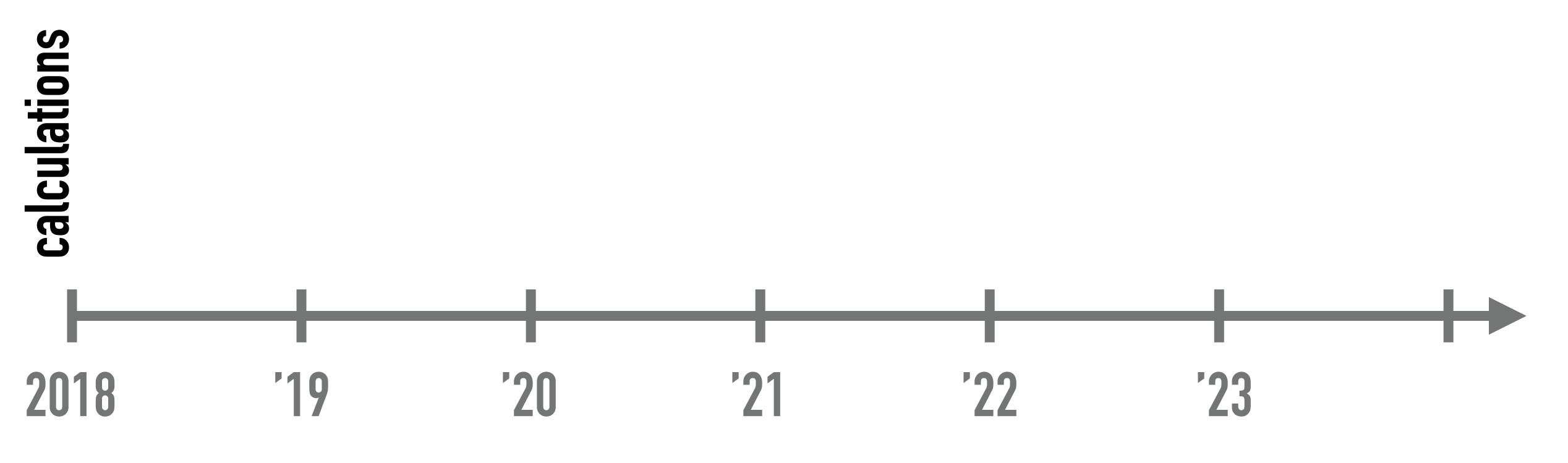


Silvia Ferrario Ravasio

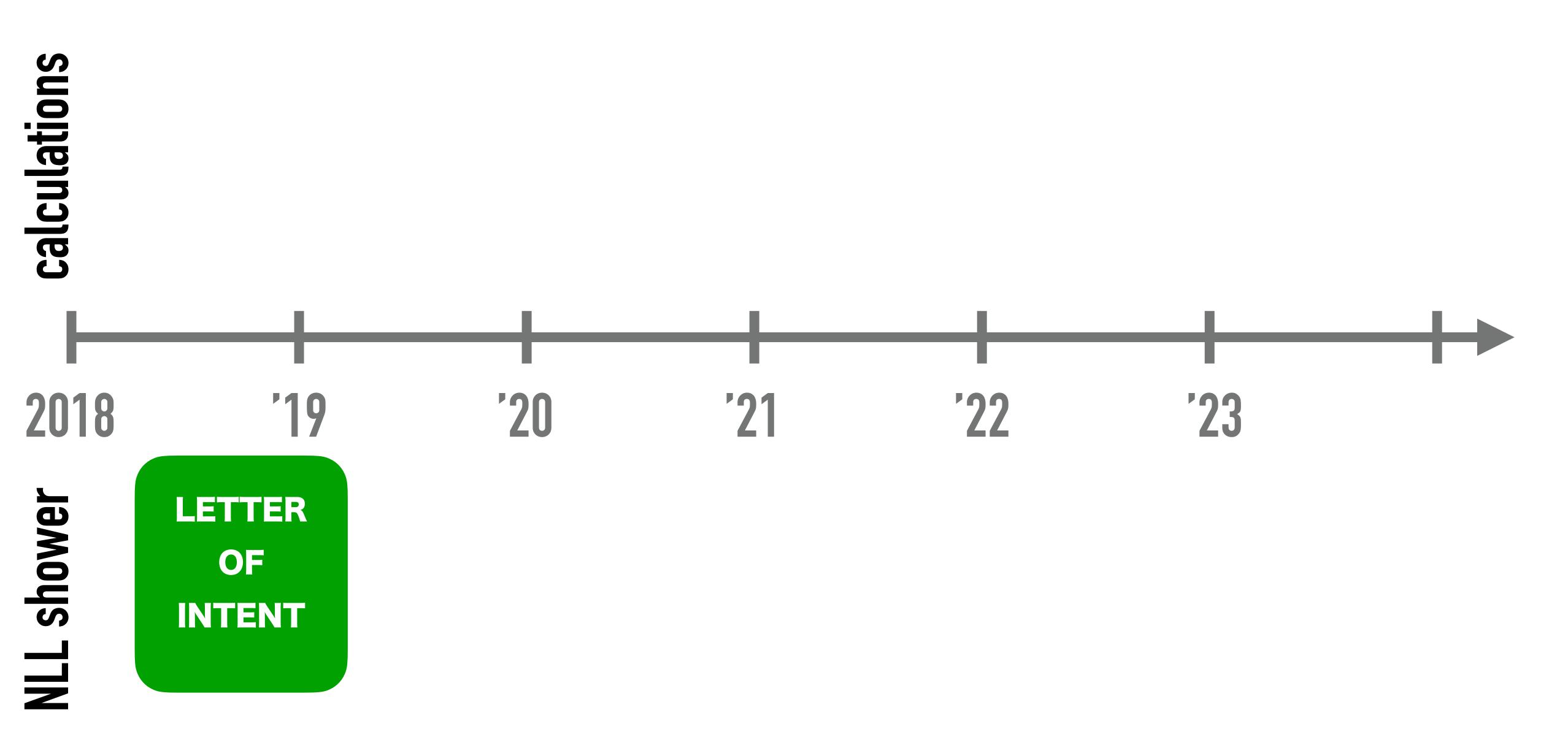
## PanScales

A project to bring logarithmic understanding and accuracy to parton showers



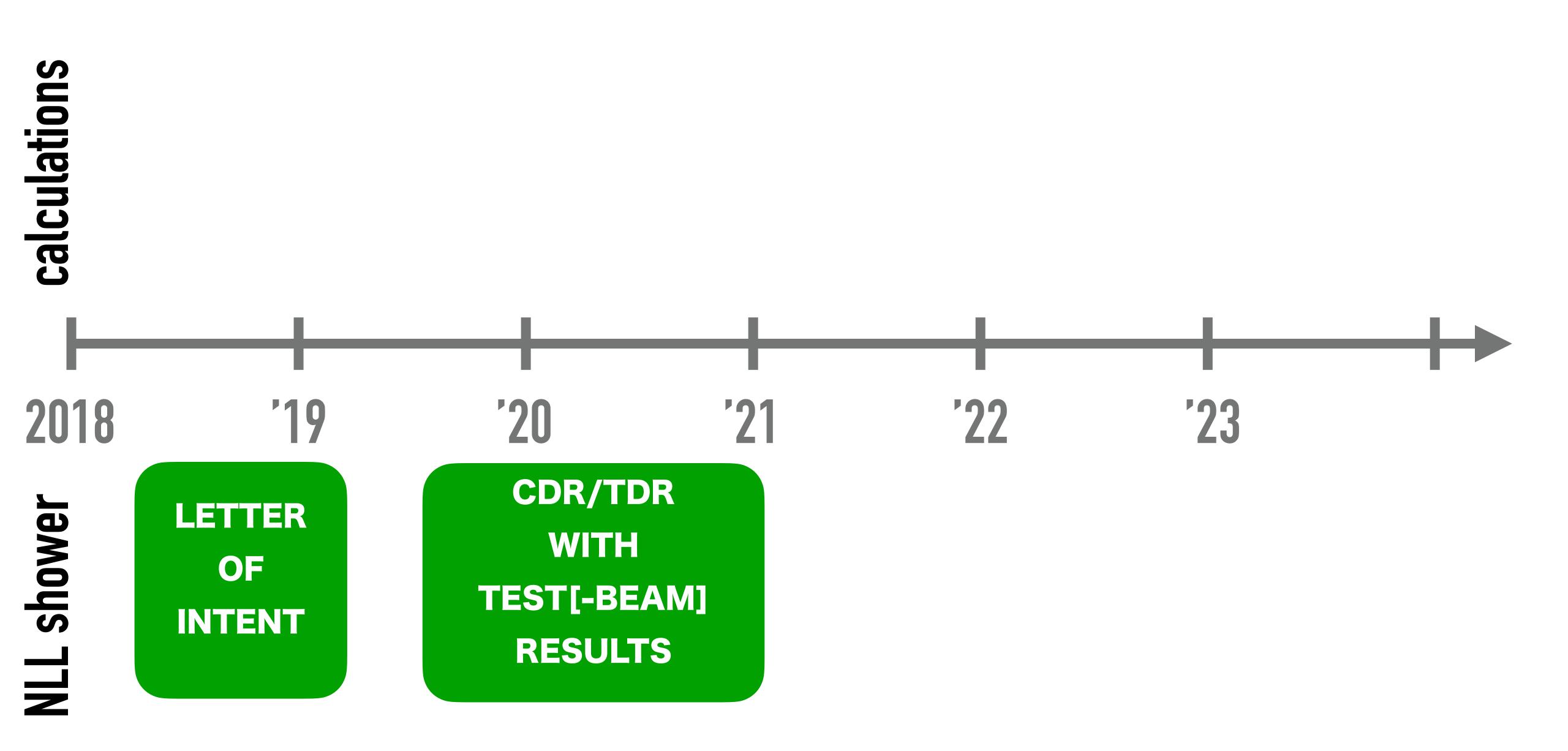


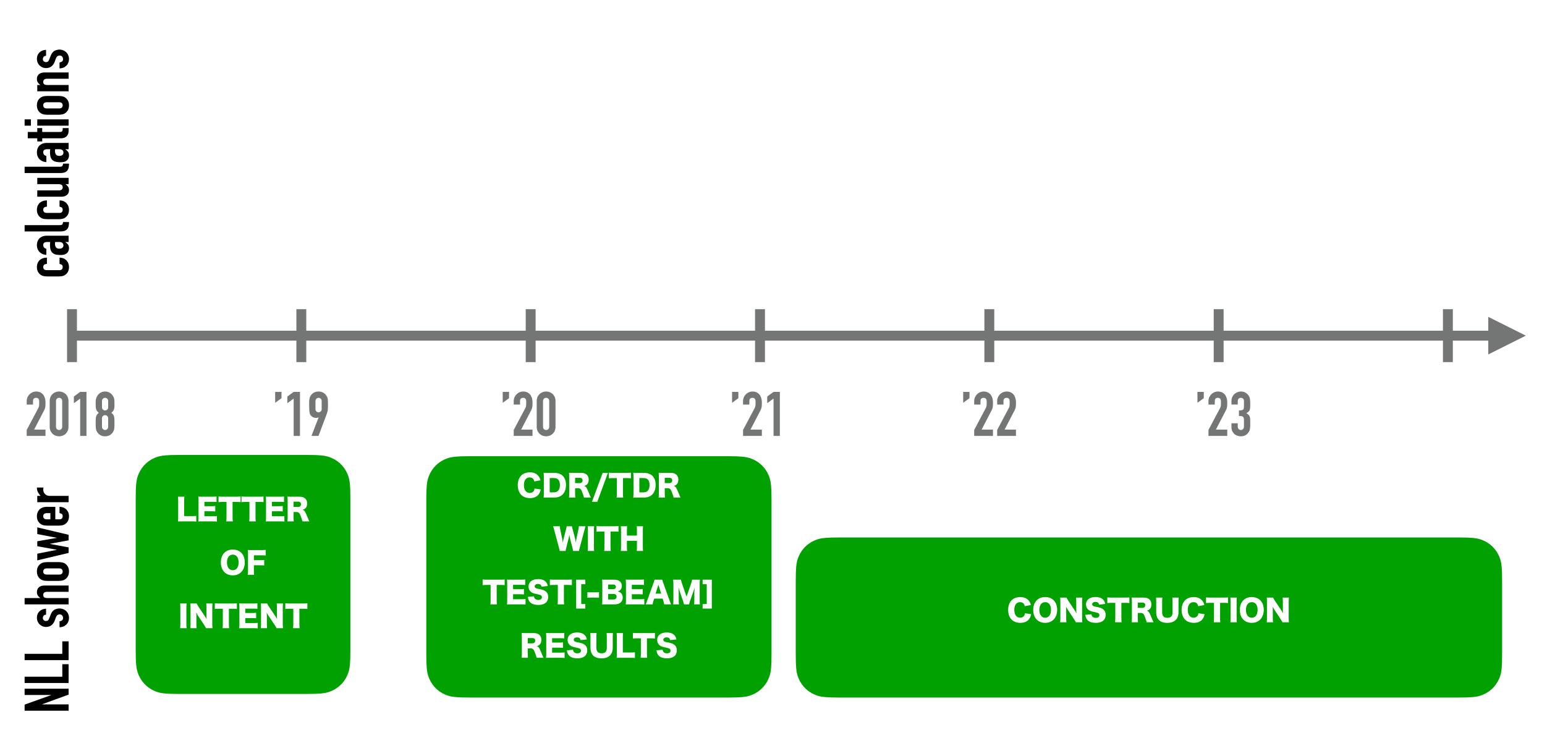
NLL shower



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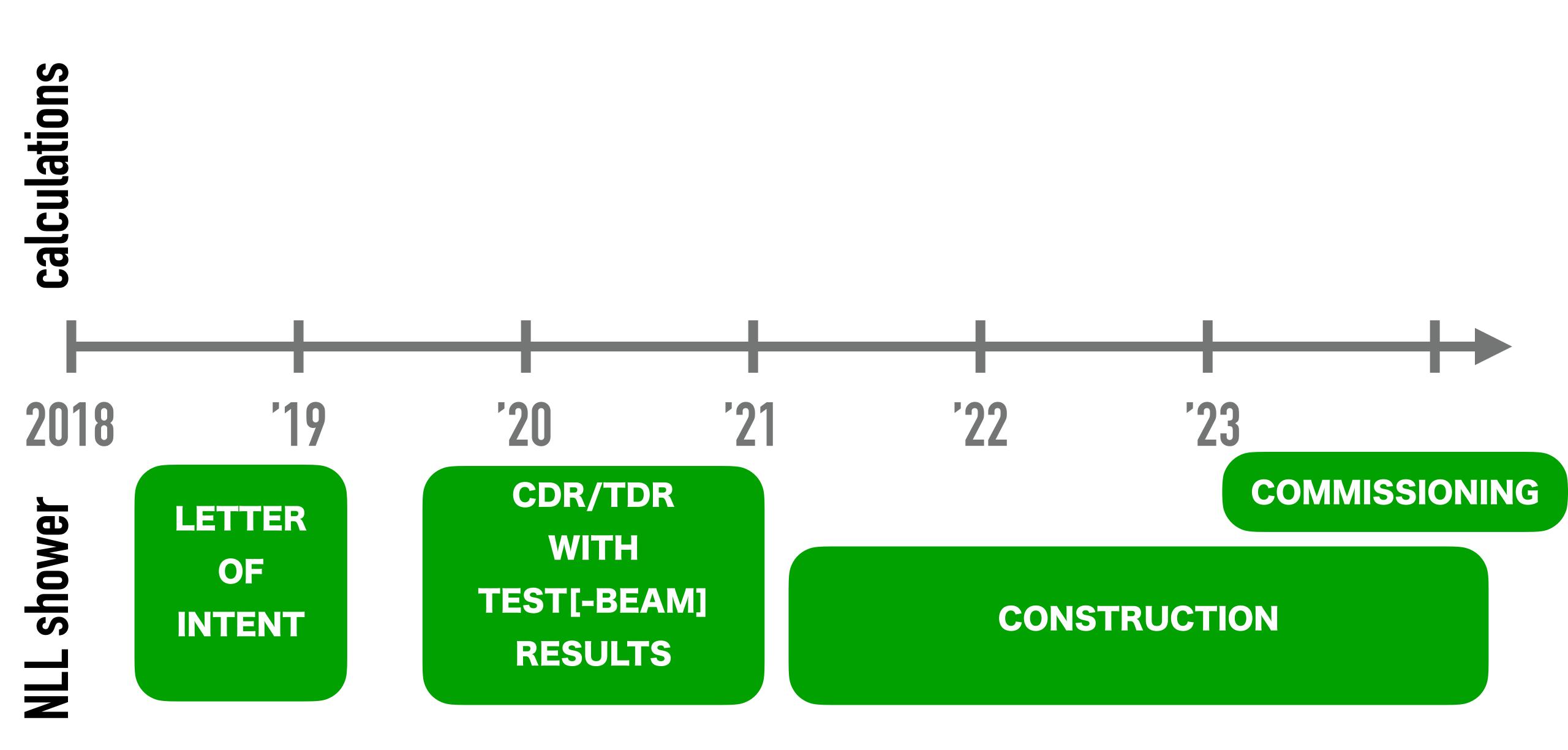
Moriond QCD, March 2023

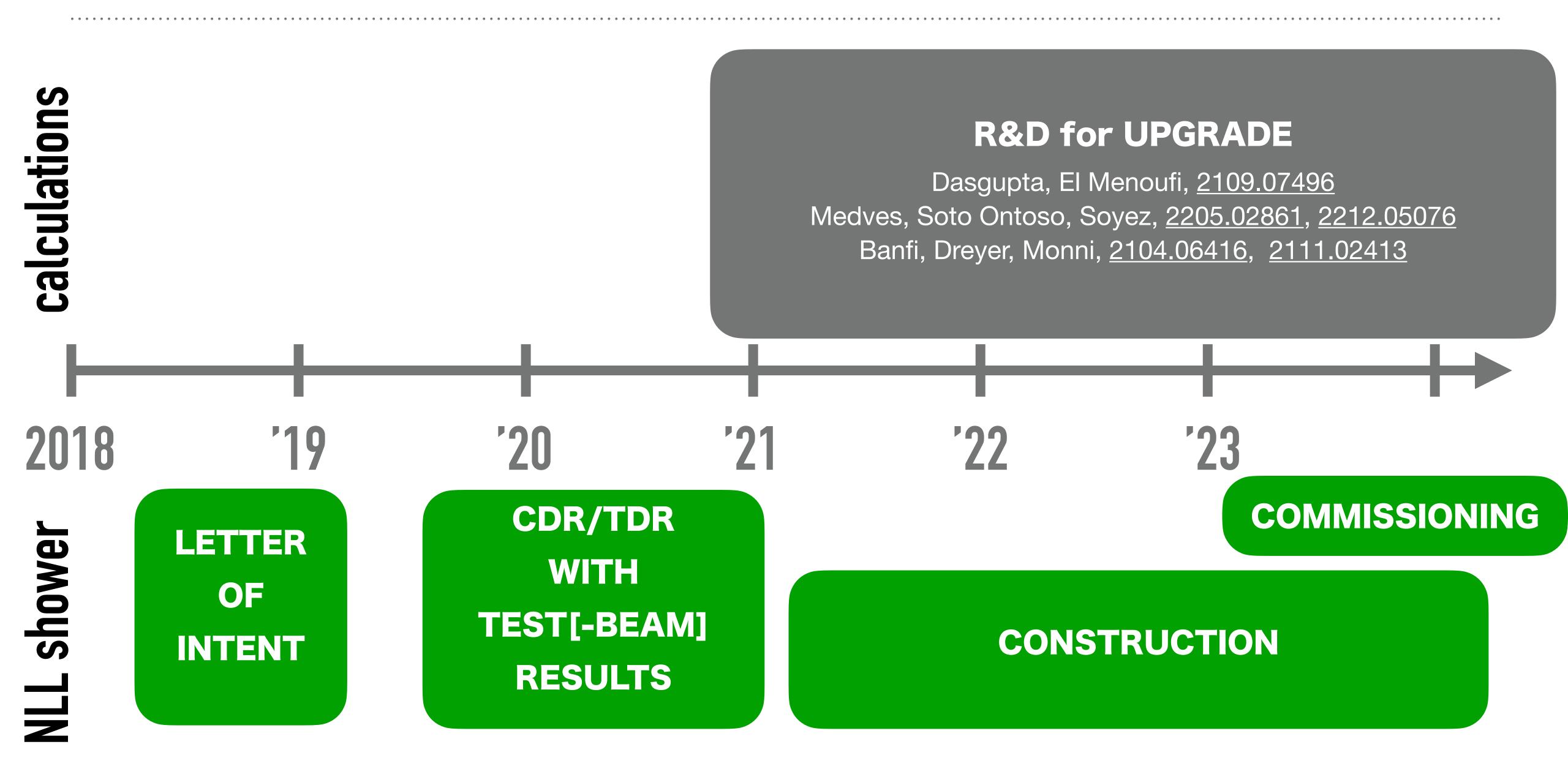




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#### PanLocal

 $k_t \sqrt{\theta}$  ordered

#### Recoil

⊥: local

+: local

-: local

Dipole partition event CoM

#### PanGlobal

 $k_t$  or  $k_t \sqrt{\theta}$  ordered

#### Recoil

1: global

+: local

-: local

Dipole partition event CoM

e+e-: Dasgupta, Dreyer, Hamilton, Monni, GPS & Soyez, 2002.11114; pp: van Beekveld, Ferrario Ravasio, GPS, Soto Ontoso, Soyez, Verheyen, 2205.02237; & pp tests, ibid + Hamilton: 2207.09467

#### Colour

nested ordered double soft (NODS)

Designed to
ensure LL are
full colour
(also gets many
NLL at full
colour)

Hamilton, Medves, GPS, Scyboz, Soyez, 2011.10054

## Spin

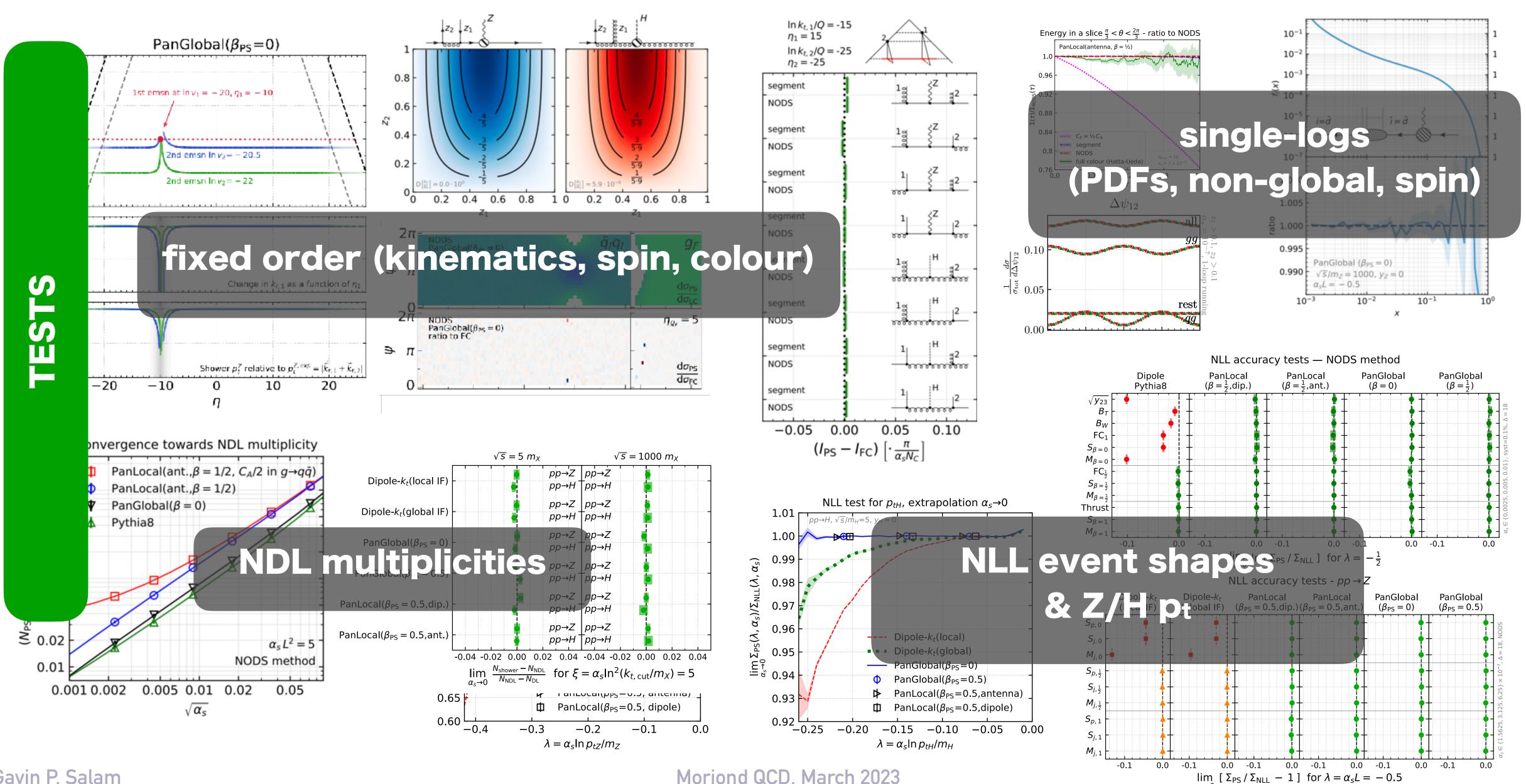
for correct
azimuthal
structure in
collinear and
soft—collinear

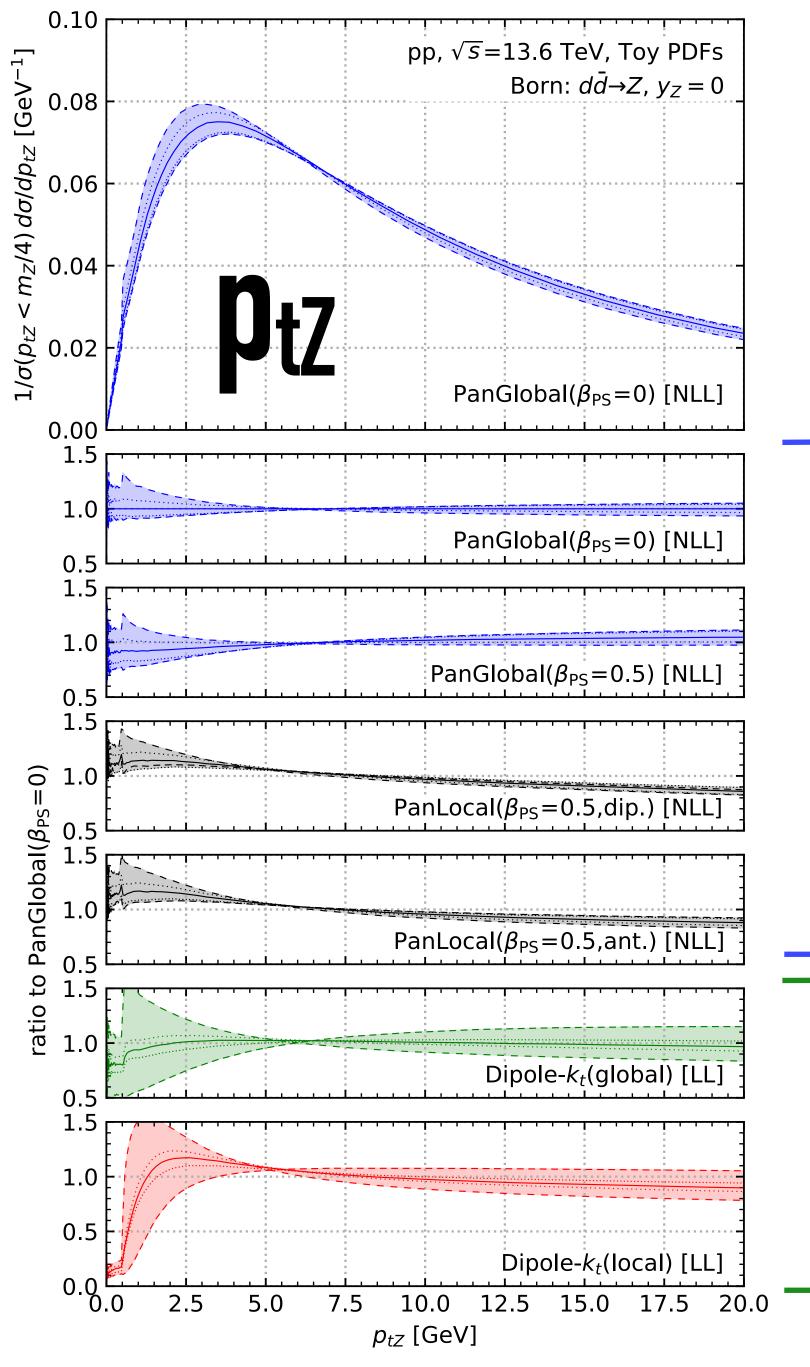
[Collins-Knowles extended to soft sector]

Karlberg, GPS, Scyboz, Verheyen, <u>2011.10054</u>; ibid + Hamilton, <u>2111.01161</u>

& pp extensions: van Beekveld et al, 2205.02237

# a selection of the logarithmic accuracy tests





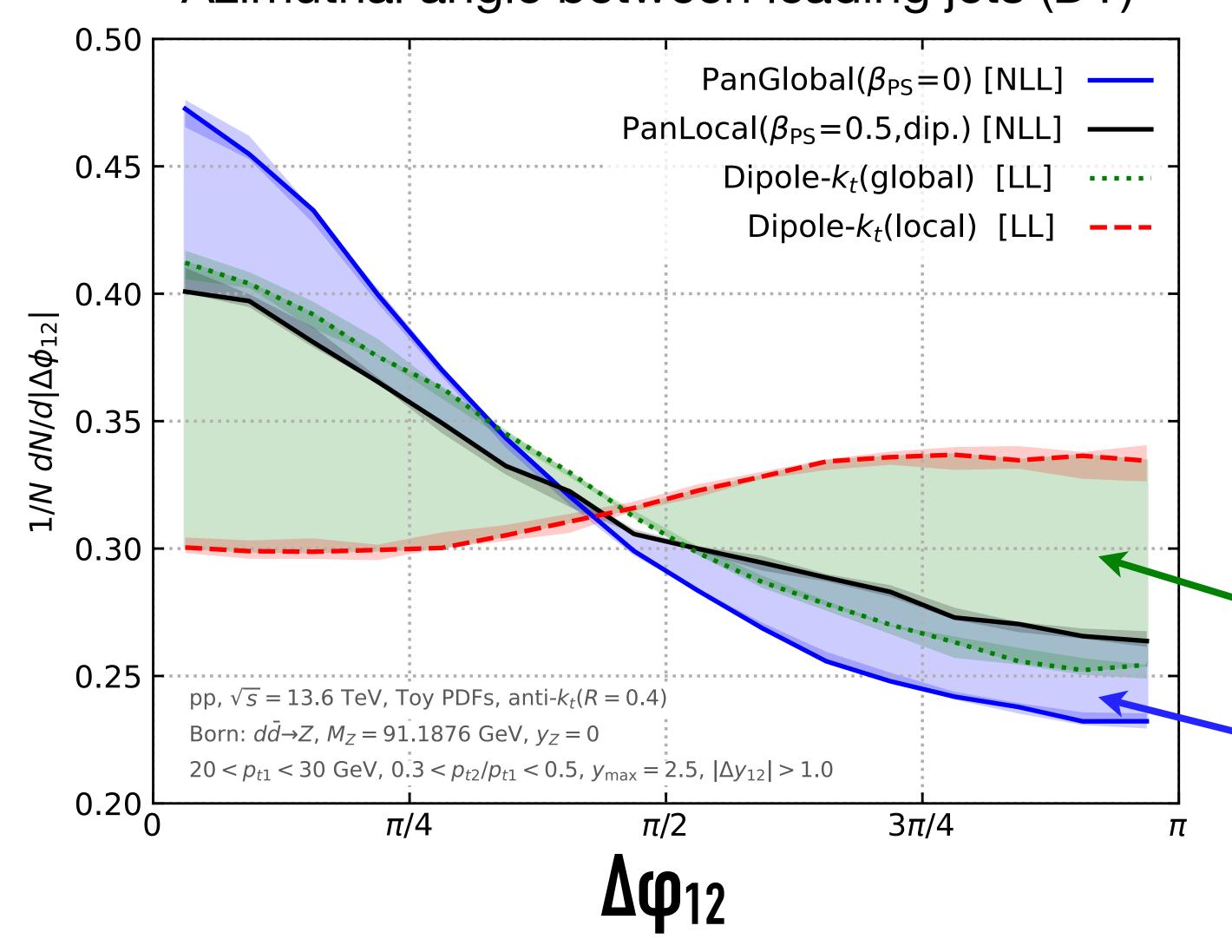
NLL showers

LL showers for inclusive quantities like ptz, advantage of NLL shower is partly in reduction of uncertainties

van Beekveld, Ferrario Ravasio, GPS, Soto Ontoso, Soyez, Verheyen, Hamilton: 2207.09467

# $m_{\ell\ell} = m_Z$

Azimuthal angle between leading jets (DY)



for more exclusive quantities, also see clear shape differences in going to NLL

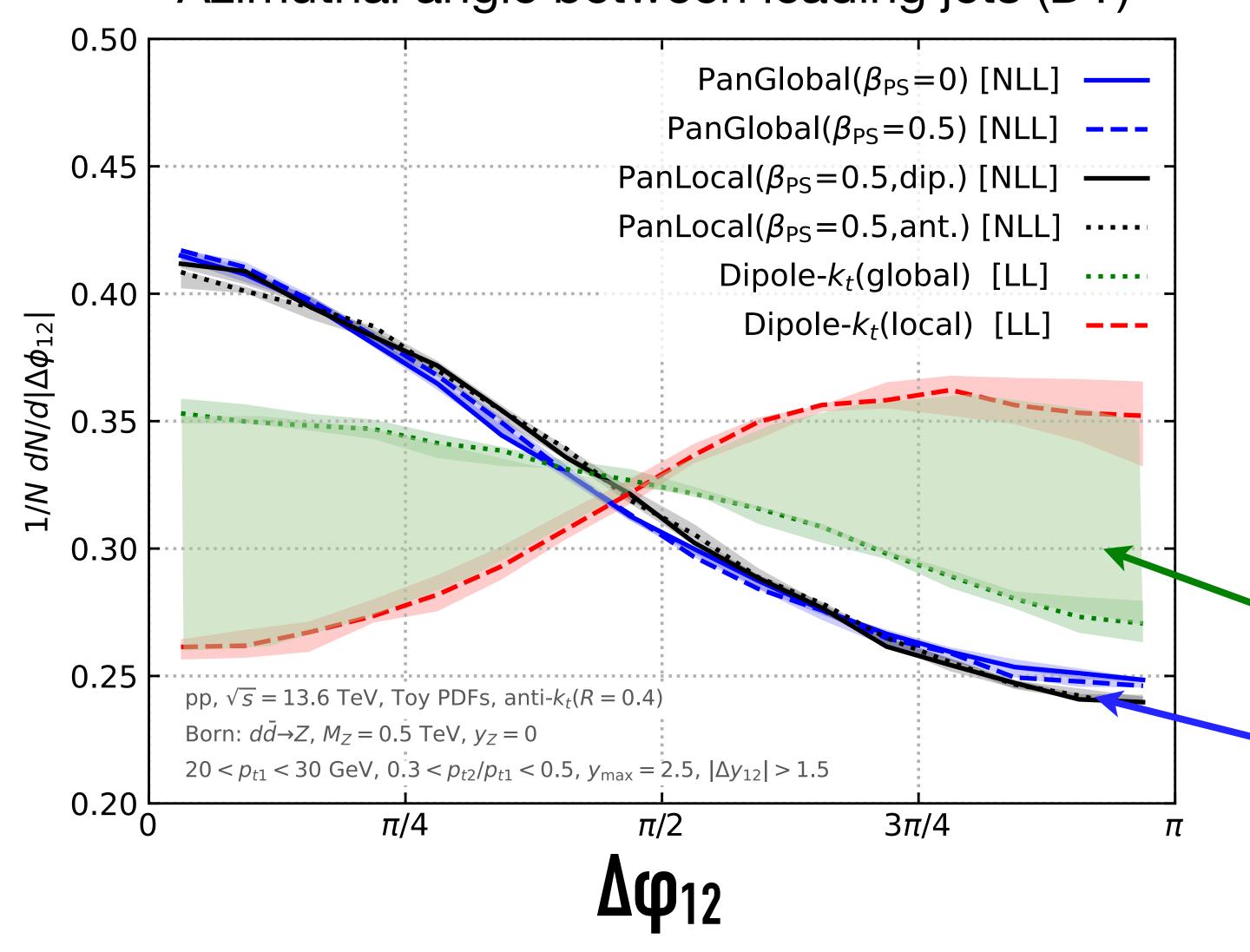
**LL** showers

**NLL showers** 

van Beekveld, Ferrario Ravasio, GPS, Soto Ontoso, Soyez, Verheyen, Hamilton: 2207.09467

# $m_{ee} = 500 \, \text{GeV}$

Azimuthal angle between leading jets (DY)



for more exclusive quantities, also see clear shape differences in going to NLL

especially at larger scales

**LL** showers

**NLL showers** 

van Beekveld, Ferrario Ravasio, GPS, Soto Ontoso, Soyez, Verheyen, Hamilton: 2207.09467

Hamilton, Karlberg, GPS, Scyboz, Verheyen, <u>2301.09645</u>

- ➤ Proof of concept explored for  $e^+e^- \rightarrow 2$  jets @ NLO
- ➤ some matching schemes supplement shower with pure  $\mathcal{O}(\alpha_s)$ , e.g. MC@NLO, KrKNLO, MAcNLOPS: these seem straightforward
- ➤ in other schemes, first emission is generated by an external program (POWHEG, MINNLO, Geneva, etc.): these need more care

cf also Corke, Sjostrand, 1003.2384  $\log k_t$ Lund plane double **HEG** contour counting

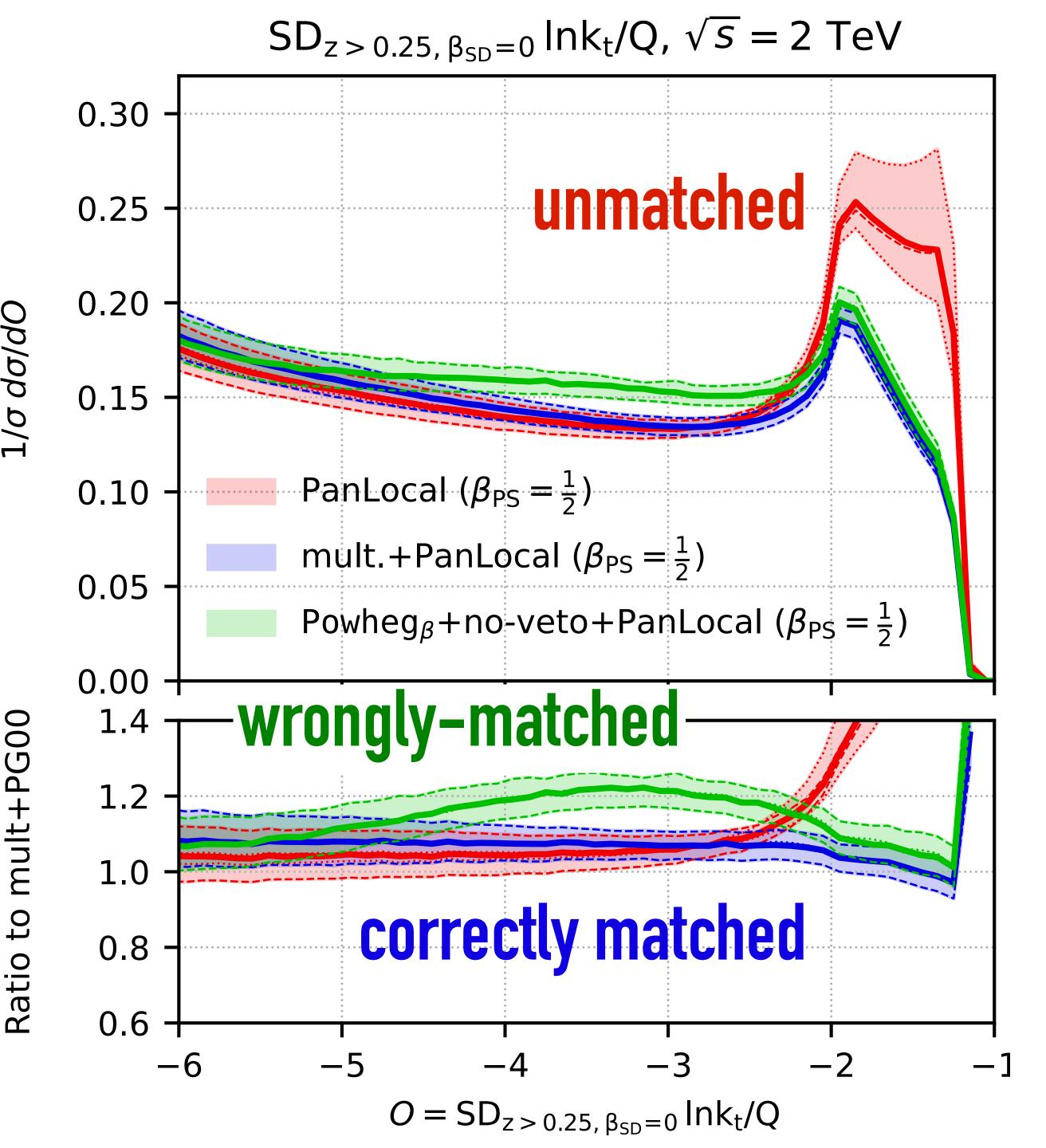
 $\log 1/\theta$ 

# Matching & log-accuracy

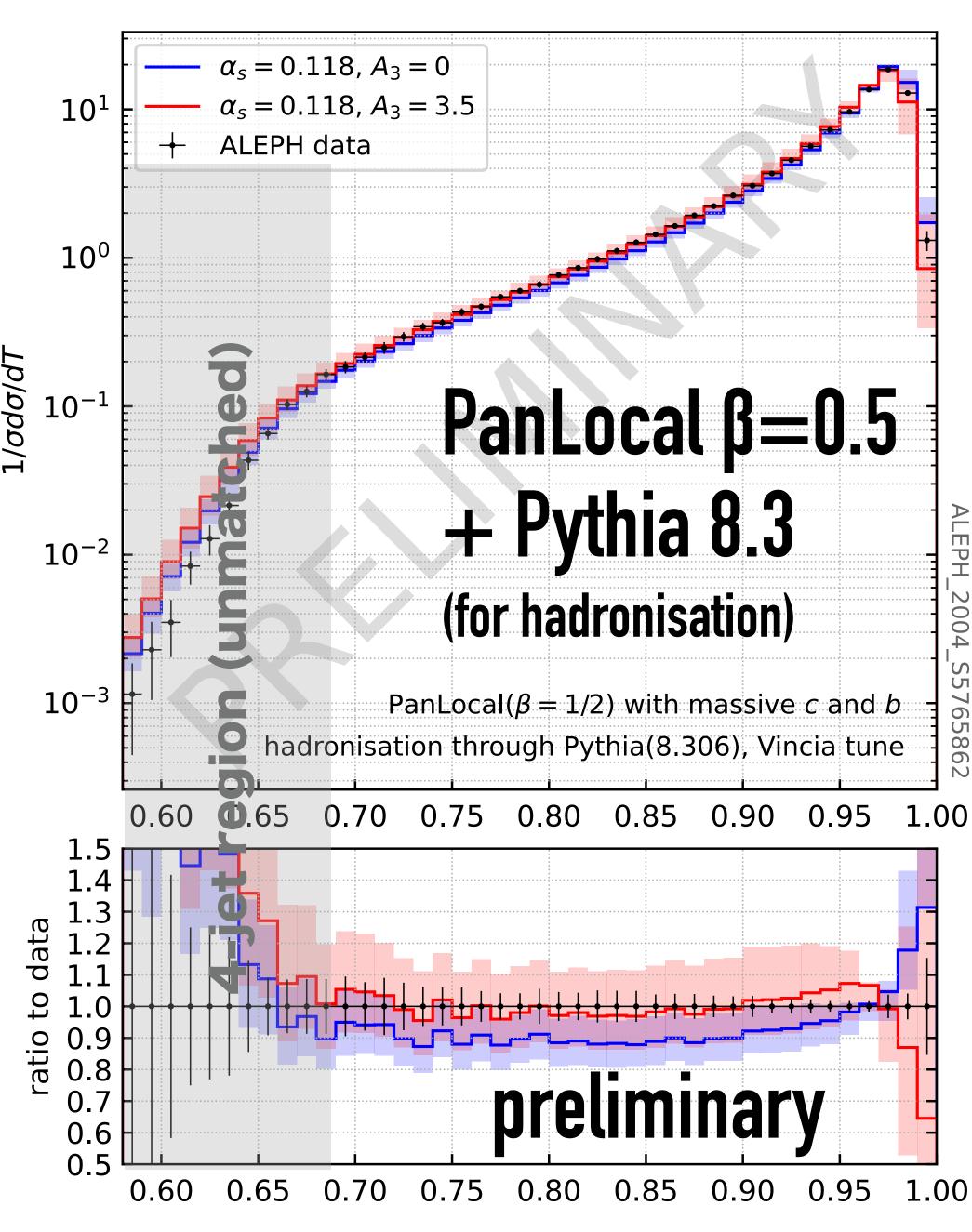
- ➤ Done correctly, matching augments accuracy of shower from NLL to NLL + NNDL (for event shapes)
- ➤ Done wrongly, it breaks exponentiation structure of shower (impact depends on observable)
- example with significant impact is
   SoftDrop transverse momentum
   (i.e. jet substructure)

$$\partial_L \Sigma_{\text{SD}}(L) = \bar{\alpha} c \, e^{\bar{\alpha} c L - \bar{\alpha} \Delta} - 2\bar{\alpha} L e^{-\bar{\alpha} L^2} (1 - e^{-\bar{\alpha} \Delta})$$

# spurious term from wrong matching



## e+e-thrust



# First comparisons to data

- $\triangleright$  we're starting with  $e^+e^-$  data
- ➤ aiming to understand nature of residual perturbative shower uncertainties
- and interplay with non-perturbative tuning
- plot includes preliminary treatment of heavy-quark masses

Medium term: making proper use of LEP data for tuning almost certainly requires NLO 3-jet accuracy.

## Conclusions

- > PanScales is first validated NLL shower (with spin & full-colour@LL/NLL)
  - ➤ benefits of LL → NLL include reduced uncertainties (and ability to reliably estimate uncertainties)
  - > multi-differential soft/collinear observables have enhanced sensitivity to NLL
  - > NLL is foundation for yet higher accuracies
- ➤ Matching is one of the next frontiers
  - ➤ first results with NLO  $e^+e^- \rightarrow 2$  jets
  - > for realistic applications we also need massive quarks and tuning
- > We're on the path towards public code
  - > exact timeline still fuzzy!

# backup

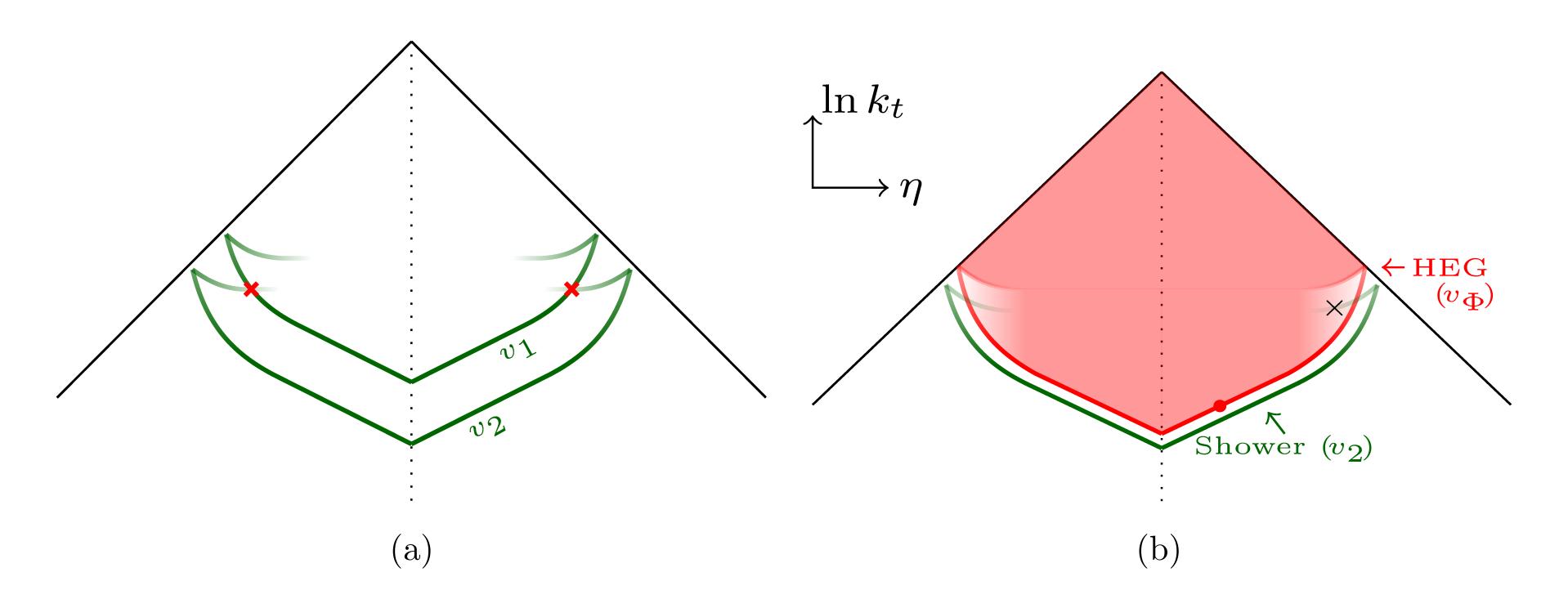
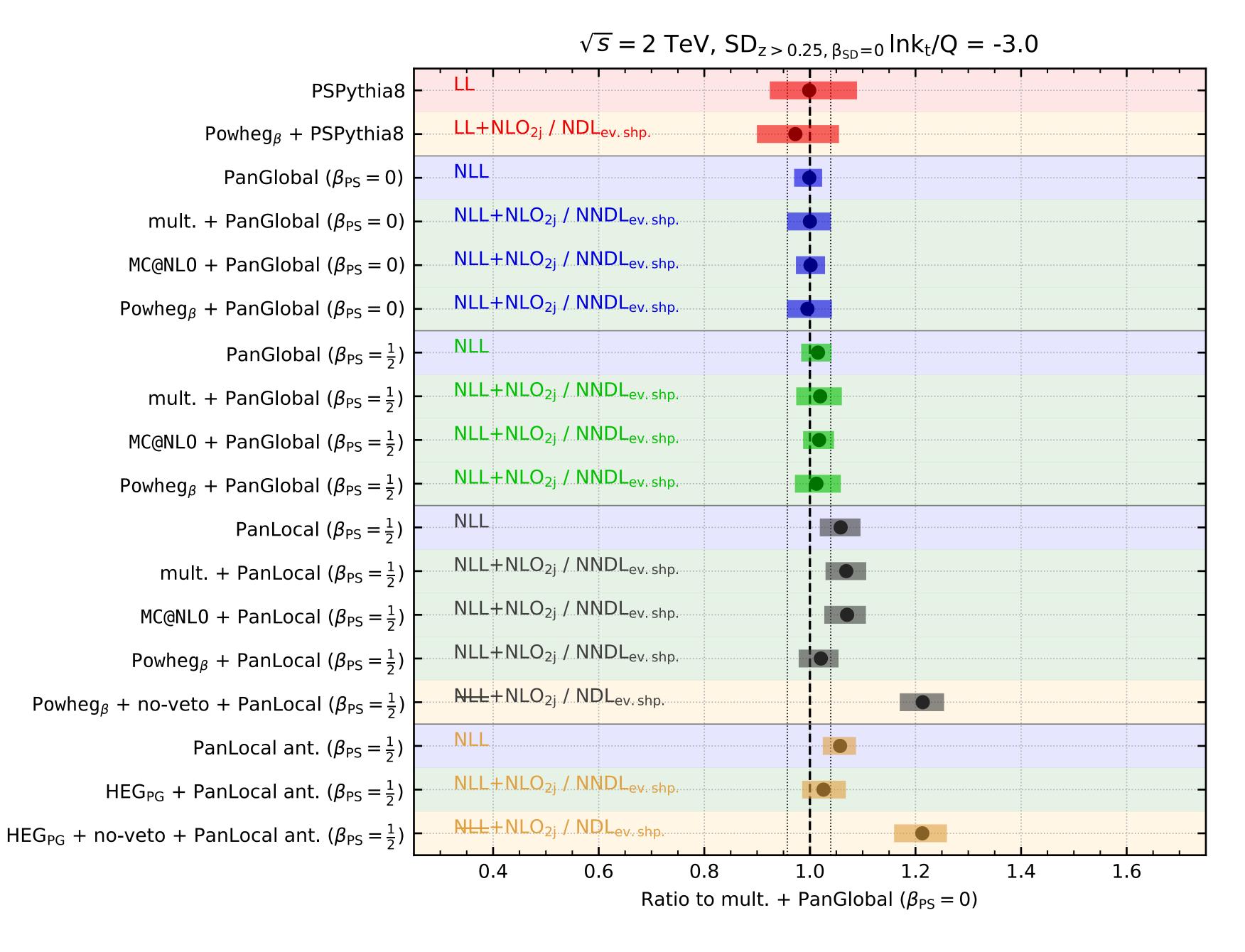


Figure 2: Schematic illustration of the issue associated with gluon asymmetrisation. (a) Contours on the Lund plane, in the PanLocal family of showers, highlighting the fact that a given physical point X in the Lund plane (highlighted with a red cross) can come from two different values of v. The shading of the green curves represents the variation in radiation intensity along the contour. (b) Density plot, at each point in the Lund plane, representing schematically the fraction of the emission intensity at that point that has been excluded once the HEG has reached a given v value ( $v_{\Phi}$ ) without emitting, and an illustration that as the shower continues there may still be phase-space points (such as that marked with a cross) where the Sudakov has only been partially accounted for. The implications are discussed in the text.



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# Matching — augment from NLL to NLL + NNDL?

Two ways of counting logarithms

$$\ln \Sigma = \alpha_s^n L^{n+1} + \underline{\alpha_s^n L^n} + \alpha_s^n L^{n-1} + \dots \text{ (relevant when } \alpha_s L \sim 1)$$
LL NLL NNLL

$$\Sigma = \underbrace{\alpha_s^n L^{2n}}_{\text{DL}} + \underbrace{\alpha_s^n L^{2n-1}}_{\text{NDL}} + \underbrace{\alpha_s^n L^{2n-2}}_{\text{NNDL}} + \dots \text{ (relevant when } \alpha_s L^2 \sim 1)$$

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